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**FORMALIZACIÓN MATEMÁTICA Y EN FÍSICA
DE PARTÍCULAS, EN RELACIÓN A LA BRECHA
DE MASA Y LA CURVATURA GEOMÉTRICA DE
LOS CAMPOS CUÁNTICOS**

**MATHEMATICAL FORMALIZATION AND PARTICLE PHYSICS,
IN RELATION TO THE MASS GAP AND THE GEOMETRIC
CURVATURE OF QUANTUM FIELDS**

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Formalización Matemática y en Física de Partículas, en Relación a la Brecha de Masa y la Curvatura Geométrica de los Campos Cuánticos

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RESUMEN

En recientes manuscritos, este investigador ha formulado alternativas de solución al Problema del Milenio de Yang – Mills, intentando unificar, desde la teoría cuántica de campos hasta las teorías de la relatividad general y especial respectivamente, sin desprendernos de cuestiones tan elementales como las representaciones en *álgebra de Lie*, de cuyo resultado, se ha concluido en lo fundamental, que toda partícula o antipartícula, con masa o sin masa, según sea el caso, supera el estado de vacío, demostrando una brecha de masa positiva, esto es, cuando se aproxima o supera la velocidad de la luz, deformando así, el campo cuántico en el que interactúa, repercutiendo en las trayectorias de las partículas o antipartículas circundantes. Ahora bien, el propósito de esta investigación, es proponer modelos hipotéticos para campos de Yang – Mills abelianos y no abelianos, grupos de gauge y Lie usando distintos operadores para espacios en cuatro dimensiones \mathbb{R}^4 , a través de los cuales, quedará demostrado, que la brecha de masa de una partícula o antipartícula con o sin masa, siempre arroja un valor positivo superior a cero.

Palabras clave: física de partículas, campos de gauge, teorías de calibre, grupos de Lie, libertad asintótica, dimensión \mathbb{R}^4 , campos de Yang Mills abelianos y no abelianos, superficie espacial, superficie temporal, operador de Casimir, transformación de Lorentz, ecuación de Callan-Symanzik, integral de trayectoria, representación de espinores.

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Mathematical Formalization and Particle Physics, in Relation to the Mass Gap and the Geometric Curvature of Quantum Fields

ABSTRACT

In recent manuscripts, this researcher has formulated alternative solutions to the Yang-Mills Millennium Problem, trying to unify, from quantum field theory to the theories of general and special relativity respectively, without detaching ourselves from such elementary questions as representations in Lie algebra, from the result of which it has been concluded in the main, that every particle or antiparticle, with or without mass, as the case may be, exceeds the vacuum state, demonstrating a positive mass gap, that is, when it approaches or exceeds the speed of light, thus deforming the quantum field in which it interacts and affecting the trajectories of the surrounding particles or antiparticles. Now, the purpose of this research is to propose hypothetical models for abelian and non-abelian Yang-Mills fields, gauge and lie groups using different operators for spaces in four dimensions \mathbb{R}^4 , through which it will be demonstrated that the mass gap of a particle or antiparticle with or without mass, it always yields a positive value greater than zero.

Keywords: particle physics, gauge fields, caliber theories, Lie groups, asymptotic freedom, \mathbb{R}^4 dimension, abelian and non-abelian Yang Mills fields, spatial surface, time surface, Casimir operator, Lorentz transformation, Callan-Symanzik equation, trajectory integral, spinor representation.



INTRODUCCIÓN

Preliminarmente, cabe precisar que se trabajará en campos cuánticos en dimensión \mathbb{R}^4 , en estructuras de gauge específicas, a propósito de sus transformaciones, con trayectorias orbitales arbitrarias, utilizando distintas métricas vectoriales, espaciales, temporales y operadores cuánticos de campo, todo esto, en superficies de espacio – tiempo cuatridimensionales, por lo que, no solamente se recurrirá a la teoría cuántica de campos de Yang – Mills, sino también a las teorías de la relatividad general y especial y otras leyes propias de la física y de las matemáticas puras, todo esto, con la finalidad de demostrar, que la brecha de masa, en un campo de Yang – Mills, esto es, cuando una partícula o antipartícula con o sin masa, según sea el caso, supera el cero absoluto, arroja un salto de energía cuyo resultado siempre es positivo. En el apartado de Resultados y Discusión, se desplegarán los sistemas matemáticos y de la física de partículas correspondientes que sostienen la hipótesis contenida en este Artículo Científico y en definitiva, en los trabajos que anteceden a éste.

Para estos efectos, se han diseñado campos cuánticos hipotéticos, con superficies espaciales y temporales arbitrarias, todo esto, con la finalidad, de demostrar la existencia de la brecha de masa positiva y paralelamente, la curvatura geométrica de los campos cuánticos y los agujeros deformantes de los referidos campos.

METODOLOGÍA

La teorización desplegada en el presente manuscrito, resulta de la aplicación de una metodología de investigación integral, esto es, bajo un enfoque híbrido, tanto desde el punto de vista cualitativo como en su dimensión cuantitativa. El tipo de investigación que ha sido desarrollado a lo largo del presente Artículo Científico, es esencialmente predictivo, a la luz de la física teórica, aunque comporta también un carácter empírico o experimental. Por otro lado, las líneas de investigación adoptadas para la formulación del estado del arte, se ajustan al constructivismo. Cabe indicar, que no existe población de estudio en la medida en que el presente artículo científico, no es de carácter sociológico o social, más aun, en mérito a su impacto en la realidad de transformación. Tampoco se han implementado técnicas de recolección de información, tales como encuestas, entrevistas, etc, salvo revisión bibliográfica, a razón del campo de investigación abordado. Adicionalmente a lo antes expuesto, es preciso resaltar, que el material de apoyo es meramente bibliográfico. La técnica metodológica, dada la complejidad de la



temática escrutada, es deductiva, pues la teorización en sentido estricto, ha sido desarrollada desde principios y premisas generales que son inherentes a la física de partículas en sentido lato. Finalmente, para efectos de construir y desarrollar las ecuaciones constantes en el presente artículo científico, se ha tomado en consideración la teoría cuántica de campos, muy especialmente, en tratándose de los campos de Yang – Mills, sin perjuicio de los demás sistemas de recalibración deducidos y esbozados a lo largo del presente Artículo Científico.

RESULTADOS Y DISCUSIÓN (Formulación Matemática y en física de partículas)

En un grupo cuántico de estructura G y bajo el algebra de Lie, obtenemos lo que sigue:

$$\langle \mathfrak{A} | \mathfrak{B} \rangle = -\text{Tr}_{\text{mat}(\overline{\mathfrak{H}}|\mathbb{C})}(\mathfrak{A} | \mathfrak{B})$$

Cuya transformación de gauge se reduce a lo que sigue:

$$\mathfrak{A} \cdot \Omega = \mathfrak{A}^\Omega = \Omega^{-1} \mathfrak{d}\Omega + \Omega^{-1} \mathfrak{A} \Omega$$

De cuyo resultado se obtienen la totalidad de las órbitas.

Por otro lado, usando la métrica de Riemann en un volumen espacial específico, y utilizando el operador de Hodge, tenemos:

$$u \wedge \star v = (u, v)_q \mathfrak{d}\omega$$

Cuyas secciones se definen así:

$$\langle u | v \rangle = \int_{\mathfrak{M}} (u, v)_q \mathfrak{d}\omega$$

De lo que obtenemos lo que sigue:

$$|u \otimes \mathfrak{E}|^2 = -\text{Tr}(\mathfrak{E} \cdot \mathfrak{E}) u \wedge \star v = -\text{Tr}(\mathfrak{E} \cdot \mathfrak{E}) (u, v)_q \mathfrak{d}\omega$$

Cuya solución de Yang – Mills, es la que sigue:

$$\begin{aligned} & \mathfrak{S}_{\text{YM}}(\mathfrak{A}) \int_{\mathfrak{M}} |\mathfrak{d}\mathfrak{A} + \mathfrak{A} \wedge \mathfrak{A}|^2 \\ & \mathfrak{I} \exp \left[\int_{\mathfrak{C}} \mathfrak{A} \right] - \frac{1}{3} \text{Tr} \int_{\mathfrak{A} \in \mathfrak{A}_{\mathfrak{M},g}/G} \mathfrak{I} \exp \left[\int_{\mathfrak{C}} \mathfrak{A} \right] e^{-1/2 \mathfrak{S}_{\text{YM}}(\mathfrak{A})} \mathfrak{D}\mathfrak{A} \end{aligned}$$



En la que la curvatura de la superficie espacial y de la superficie temporal se expresan de la siguiente manera:

$$\begin{aligned} \mathfrak{D}\mathfrak{A} + \mathfrak{A} \wedge \mathfrak{A} &= \sum_{\alpha} \sum_{1 \leq i < j \leq 4} \alpha_{i,j,\alpha} \otimes dx^i \wedge dx^j \otimes \mathfrak{E}^{\alpha} + \sum_{\alpha, \beta} \sum_{1 \leq i < j \leq 4} \alpha_{i,\alpha} \alpha_{j,\beta} \otimes dx^i \wedge dx^j \otimes \langle \mathfrak{E}^{\alpha} | \mathfrak{E}^{\beta} \rangle \\ &+ \sum_{\alpha} \sum_{j=1} \alpha_{0;j,\alpha} \otimes dx^0 \wedge dx^j \otimes \mathfrak{E}^{\alpha} \end{aligned}$$

Cuyas permutaciones y transformaciones, se expresan así:

$$C_{\gamma}^{\alpha\beta} = -\mathfrak{Tr}(\mathfrak{E}^{\gamma}(\mathfrak{E}^{\alpha}, \mathfrak{E}^{\beta}))$$

$$\begin{aligned} \mathfrak{D}\mathfrak{A} + \mathfrak{A} \wedge \mathfrak{A} &= \sum_{\gamma} \langle \sum_{1 \leq i < j \leq 4} \alpha_{i,j,\alpha} \otimes dx^i \wedge dx^j + \sum_{\alpha < \beta} \sum_{1 \leq i < j \leq 4} \alpha_{i,\alpha} \alpha_{j,\beta} C_{\gamma}^{\alpha\beta} \otimes dx^i \wedge dx^j \\ &+ dx^i \wedge dx^j \sum_{j=1} \alpha_{0;j,\gamma} \otimes dx^0 \wedge dx^j \rangle \otimes \mathfrak{E}^{\gamma} \end{aligned}$$

$$\int_{\mathbb{R}^4} |\mathfrak{D}\mathfrak{A} + \mathfrak{A} \wedge \mathfrak{A}|^2$$

$$\begin{aligned} &= \sum_{i < j} \int_{\mathbb{R}^4} \left\| \sum_{\alpha} \alpha_{i,j,\alpha}^2 + \sum_{\gamma} \sum_{\alpha < \beta, \hat{\alpha} < \hat{\beta}} \alpha_{i,\alpha} \alpha_{j,\beta} \alpha_{i,\hat{\alpha}} \alpha_{j,\hat{\beta}} C_{\gamma}^{\alpha\beta} C_{\gamma}^{\hat{\alpha}\hat{\beta}} \right. \\ &+ 4 \sum_{\alpha < \beta, \gamma} \alpha_{i,j,\gamma} \alpha_{i,\alpha} \alpha_{j,\beta} C_{\gamma}^{\alpha\beta} \left. \right\| d\omega + \sum_j \int_{\mathbb{R}^4} \sum_{\alpha} \alpha_{0;j,\alpha}^2 d\omega \end{aligned}$$

Campo Cuántico Abeliano

Usando el teorema de Stoke, tenemos:

$$\int_{\mathfrak{E}} \sum_{i=1}^4 \mathfrak{A}_i \otimes dx^i = \int_{\partial\mathfrak{E}} \sum_{i=1}^4 \mathfrak{A}_i dx^i = \int_{\mathfrak{E}} \mathfrak{D}\mathfrak{A} = \int_{\mathbb{R}^4} \mathfrak{D}\mathfrak{A} \times 4_{\mathfrak{E}} = \|\mathfrak{D}\mathfrak{A}, 4_{\mathfrak{E}}\|$$

Por lo que, la integral de Yang – Mills se expresa de la siguiente manera:

$$1/ \int_{\mathfrak{A}} e^{-\|\mathfrak{D}\mathfrak{A}\|^2/2} \mathfrak{D}\mathfrak{A} \int_A e^{\sqrt{-1}(\mathfrak{D}\mathfrak{A}|4_{\mathfrak{E}})} e^{-1/2\|\mathfrak{D}\mathfrak{A}\|^2} \mathfrak{D}\mathfrak{A}$$



Cuyo cambio heurístico de variables, queda expresado así:

$$\frac{1}{\det \mathfrak{d}^{-1} \int_{\mathfrak{A}} e^{-\frac{\|\mathfrak{d}\mathfrak{A}\|^2}{2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A}) \det \mathfrak{d}^{-1} \int_A e^{\sqrt{-1}(\mathfrak{d}\mathfrak{A}|4_{\mathfrak{G}})} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}} = \frac{1}{\int_{\mathfrak{A}} e^{-\frac{\|\mathfrak{d}\mathfrak{A}\|^2}{2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A}) \int_A e^{\sqrt{-1}(\mathfrak{d}\mathfrak{A}|4_{\mathfrak{G}})} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}}$$

Más, en dimensión \mathbb{R}^4 y aplicando la función delta de Dirac, tenemos:

$$\langle \mathfrak{F}, \mathfrak{X}_x \otimes \mathfrak{d}x^\alpha \wedge \mathfrak{d}x^\beta \rangle = \left\langle \sum_{0 \leq i \leq j \leq 4} \mathfrak{S}_{ij} \mathfrak{d}x^i \wedge \mathfrak{d}x^j, \mathfrak{X}_x \otimes \mathfrak{d}x^\alpha \wedge \mathfrak{d}x^\beta \right\rangle = \langle \mathfrak{S}_{\alpha\beta} \mathfrak{d}x^\alpha \wedge \mathfrak{d}x^\beta, \mathfrak{X}_x \otimes \mathfrak{d}x^\alpha \wedge \mathfrak{d}x^\beta \rangle = \mathfrak{S}_{\alpha\beta}(\chi) \in \mathbb{R}^4$$

Cuyo gauge axial, según el teorema de Stokes, arroja como resultado lo que sigue:

$$\begin{aligned} \int_{\mathfrak{G}} \sum_{i=1}^4 \mathfrak{A}_i \otimes \mathfrak{d}x^i &= \int_{\partial \mathfrak{G}} \sum_{i=1}^4 \mathfrak{A}_i \otimes \mathfrak{d}x^i = \int_{\mathfrak{G}} \mathfrak{d}\mathfrak{A} \\ &= \int_{(0,1)^2} \mathfrak{d}s \mathfrak{d}t \left[\sum_{1 \leq i \leq j \leq 4} (\mathfrak{A}_{ij}(\sigma) |J_{ij}^\sigma|)(s, t) + \sum_{j=1}^4 (\mathfrak{A}_{0j}(\sigma) |J_{0j}^\sigma|)(s, t) \right] \\ &= \int_{(0,1)^2} \mathfrak{d}s \mathfrak{d}t \langle \mathfrak{d}\mathfrak{A}, \sum_{0 \leq i \leq j \leq 4} \chi_{\sigma(s,t)} |J_{ij}^\sigma|(s, t) \otimes \mathfrak{d}x^i \wedge \mathfrak{d}x^j \rangle = \langle \mathfrak{d}\mathfrak{A}, \widetilde{\mathfrak{B}}_{\mathfrak{G}} \rangle \end{aligned}$$

En la que $\widetilde{\mathfrak{B}}_{\mathfrak{G}}$ es igual a:

$$\begin{aligned} \widetilde{\mathfrak{B}}_{\mathfrak{G}} &= \int_{(0,1)^2} \sum_{0 \leq i \leq j \leq 4} \mathfrak{d}s \mathfrak{d}t \chi_{\sigma(s,t)} |J_{ij}^\sigma|(s, t) \mathfrak{d}x^i \wedge \mathfrak{d}x^j \\ \langle \mathfrak{d}\mathfrak{A}, \sum_{0 \leq i \leq j \leq 4} \chi_{\sigma(s,t)} |J_{ij}^\sigma|(s, t) \otimes \mathfrak{d}x^i \wedge \mathfrak{d}x^j \rangle &= \langle \mathfrak{d}\mathfrak{A}, \sum_{0 \leq i \leq j \leq 4} \chi_{\sigma(s,t)} \otimes \mathfrak{d}x^i \wedge \mathfrak{d}x^j |J_{ij}^\sigma|(s, t) \rangle \\ &= \sum_{0 \leq i \leq j \leq 4} \langle \mathfrak{A}_{i,j}, \chi_{\sigma(s,t)} |J_{ij}^\sigma|(s, t) \rangle = \sum_{0 \leq i \leq j \leq 4} \langle \mathfrak{A}_{i,j}(\sigma(s, t)) |J_{ij}^\sigma|(s, t) \rangle \\ &= \frac{1}{3 \int_{\mathfrak{A}} e^{\sqrt{-1}(\mathfrak{d}\mathfrak{A}|\widetilde{\mathfrak{B}}_{\mathfrak{G}})} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}} \\ &= \frac{1}{3 \int_{\mathfrak{A}} e^{\sqrt{-1}(\mathfrak{d}\mathfrak{A}|\widetilde{\mathfrak{B}}_{\mathfrak{G}})} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2} \det(\mathfrak{d}^{-1}) \mathfrak{D}(\mathfrak{d}\mathfrak{A})} + \frac{1}{3 \int_{\mathfrak{A}} e^{\sqrt{-1}(\mathfrak{d}\mathfrak{A}|\widetilde{\mathfrak{B}}_{\mathfrak{G}})} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}} \end{aligned}$$



Siguiendo el mismo orden de ideas, un espacio de Schwartz, quedaría expresado así:

$$\mathcal{P}_r = \left\| (m_1, m_2, \dots, m_n) \sum_{j=1}^n m_j = r \right\|$$

Más en dimensión \mathbb{R}^4 , incorporando la función de Gauss y la métrica de Lebesgue, tenemos:

$$f(x) = \varrho(x) \sqrt{\phi_\kappa(x)}$$

$$\langle f, g \rangle = \int_{\mathbb{R}^4} f \cdot g \, d\lambda$$

$$\langle \mathfrak{z}^r, \bar{\mathfrak{z}}^r \rangle = \frac{1}{\pi \int_{\mathbb{C}} \mathfrak{z}^r \cdot \bar{\mathfrak{z}}^r e^{-|z|^2} \, d\mathfrak{x} \, d\mathfrak{p}}, \mathfrak{z} = \mathfrak{x} + \sqrt{-1}\mathfrak{p}$$

Cuya función polinómica, se traduce a lo que sigue:

$$\mathcal{H}_{\varrho_{\mathfrak{R}}}(\mathfrak{x}) = \hbar_i(\mathfrak{x}^0) \hbar_j(\mathfrak{x}^1) \hbar_k(\mathfrak{x}^2) \hbar_l(\mathfrak{x}^3), \varrho_{\mathfrak{R}} = (i, j, k, l) \in \mathfrak{P}_{\mathfrak{R}}$$

Más, incorporando la métrica de Gauss, tenemos:

$$\bigcup_{r=0}^{\infty} \left| \frac{\mathcal{H}_{\varrho_{\mathfrak{R}}}(\kappa \mathfrak{x}^0, \kappa \mathfrak{x}^1, \kappa \mathfrak{x}^2, \kappa \mathfrak{x}^3) \sqrt{\phi_\kappa}}{\sqrt{\varrho_{\mathfrak{R}}!}} : \mathcal{P}_r \in \mathfrak{P}_r \right|$$

Por tanto, en dimensión \mathbb{R}^4 , tenemos:

$$\mathfrak{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^3) = \left\langle \sum_{\alpha=1}^4 \mathfrak{F}_\alpha \otimes d\mathfrak{x}^\alpha : \mathfrak{F}_\alpha \in \mathfrak{S}_\kappa(\mathbb{R}^4) \right\rangle$$

$$\bar{\mathfrak{S}}_\kappa(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) = \left\| \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{F}_{\alpha\beta} \otimes d\mathfrak{x}^\alpha \wedge d\mathfrak{x}^\beta : \mathfrak{F}_{\alpha\beta} \in \bar{\mathfrak{S}}_\kappa(\mathbb{R}^4) \rangle \right\|$$

Cuyo operador de Hodge, se reduce a lo siguiente:

$$\left\langle \sum_{0 \leq \alpha \leq \beta \leq 4} |\mathfrak{F}_{\alpha\beta} \otimes d\mathfrak{x}^\alpha \wedge d\mathfrak{x}^\beta|, \sum_{0 \leq \alpha \leq \beta \leq 4} |\tilde{\mathfrak{F}}_{\alpha\beta} \otimes d\mathfrak{x}^\alpha \wedge d\mathfrak{x}^\beta| \right\rangle = \sum_{0 \leq \alpha \leq \beta \leq 4} |\mathfrak{F}_{\alpha\beta}, \tilde{\mathfrak{F}}_{\alpha\beta}|$$

Cuya transformación de Segal - Bargmann y demás funciones holomórficas, dan como resultado lo que sigue:

$$\langle \mathfrak{z}^{\mathcal{R}}, \bar{\mathfrak{z}}^{\mathcal{R}} \rangle = \frac{1}{\pi \int_{\mathbb{C}} \mathfrak{z}^{\mathcal{R}} \cdot \bar{\mathfrak{z}}^{\mathcal{R}} e^{-|z|^2} \, d\mathfrak{x} \, d\mathfrak{p}}, \mathfrak{z} = \mathfrak{x} + \sqrt{-1}\mathfrak{p}$$



$$\Psi_\kappa = \frac{\frac{\hbar_i(\kappa \cdot)}{\sqrt{i!} \hbar_j(\kappa \cdot)}}{\frac{\sqrt{j!} \hbar_\dagger(\kappa \cdot)}{\sqrt{\dagger!} \hbar_l(\kappa \cdot)}} = \sqrt{\phi_\kappa} \rightarrow \frac{\frac{\mathfrak{z}_0^i}{\sqrt{i!} \mathfrak{z}_1^j}}{\frac{\sqrt{j!} \mathfrak{z}_2^\dagger}{\sqrt{\dagger!} \mathfrak{z}_3^l}} = \mathfrak{f}_{i,\alpha} \otimes \mathrm{d}\mathfrak{x}^i \otimes \mathfrak{E}^\alpha \rightarrow \Psi_\kappa(\mathfrak{f}_{i,\alpha}) \otimes \mathrm{d}\mathfrak{x}^i \otimes \mathfrak{E}^\alpha$$

En la que, en un espacio de Hilbert, se tiene:

$$\left\langle \sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{f}_{\alpha\beta} \otimes \mathrm{d}\mathfrak{x}^\alpha \otimes \mathrm{d}\mathfrak{x}^\beta, \sum_{0 \leq \alpha \leq \beta \leq 4} \hat{\mathfrak{f}}_{\alpha\beta} \otimes \mathrm{d}\mathfrak{x}^\alpha \otimes \mathrm{d}\mathfrak{x}^\beta \right\rangle = \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{f}_{\alpha\beta}, \hat{\mathfrak{f}}_{\alpha\beta} \rangle$$

$$\partial \sum_{i=1}^4 \mathfrak{f}_i \otimes \mathrm{d}\mathfrak{x}^i = \sum_{i=1}^4 \langle \partial_0 \mathfrak{f}_i \rangle \otimes \mathrm{d}\mathfrak{x}^0 \wedge \mathrm{d}\mathfrak{x}^i + \sum_{1 \leq i < j \leq 4} [\partial_i \mathfrak{f}_j - \partial_j \mathfrak{f}_i] \otimes \mathrm{d}\mathfrak{x}^i \wedge \mathrm{d}\mathfrak{x}^j$$

$$\Psi_\kappa = \sum_{1 \leq i < j \leq 4} \mathfrak{f}_{i,j,\alpha} \otimes \mathrm{d}\mathfrak{x}^i \wedge \mathrm{d}\mathfrak{x}^j \otimes \mathfrak{E}^\alpha \rightarrow \sum_{1 \leq i < j \leq 4} \Psi_\kappa[\mathfrak{f}_{i,j,\alpha}] \otimes \mathrm{d}\mathfrak{x}^i \wedge \mathrm{d}\mathfrak{x}^j \otimes \mathfrak{E}^\alpha$$

$$\mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^1(\mathbb{R}^3) = \left\{ \sum_{\alpha=1}^4 \mathfrak{f}_\alpha \otimes \mathrm{d}\mathfrak{x}^\alpha : \mathfrak{f}_\alpha \in \mathcal{H}^2(\mathbb{C}^4) \right\}$$

$$\mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) = \left\{ \sum_{0 \leq \alpha < \beta \leq 4} \mathfrak{F}_{\alpha\beta} \otimes \mathrm{d}\mathfrak{x}^\alpha \wedge \mathrm{d}\mathfrak{x}^\beta : \mathfrak{F}_{\alpha\beta} \in \mathcal{H}^2(\mathbb{C}^4) \right\}$$

$$\left\langle \sum_{0 \leq \alpha < \beta \leq 4} \mathfrak{F}_{\alpha\beta} \otimes \mathrm{d}\mathfrak{x}^\alpha \wedge \mathrm{d}\mathfrak{x}^\beta, \sum_{0 \leq \alpha < \beta \leq 4} \widehat{\mathfrak{F}}_{\alpha\beta} \otimes \mathrm{d}\mathfrak{x}^\alpha \wedge \mathrm{d}\mathfrak{x}^\beta \right\rangle = \sum_{0 \leq \alpha < \beta \leq 4} \langle \mathfrak{F}_{\alpha\beta}, \widehat{\mathfrak{F}}_{\alpha\beta} \rangle = \int_{\mathbb{C}^4} \mathfrak{F}_{\alpha\beta} \widehat{\mathfrak{F}}_{\alpha\beta} \mathrm{d}\lambda_4$$

Cuyos polinomios de Hermite, se satisfacen así:

$$\frac{\mathrm{d}}{\mathrm{d}\mathfrak{x}} \left(\hbar_\eta(\chi) \mathfrak{E}^{-\frac{\chi^2}{4}} \right) = |\chi \hbar_\eta(\mathfrak{x}) - \chi \hbar_{\eta+1}(\mathfrak{x}) - \chi/2 \hbar_\eta(\chi)| \mathfrak{E}^{-\frac{\chi^2}{4}} = \frac{\mathrm{d}}{\mathrm{d}\mathfrak{x}} \left(\hbar_\eta(\chi) \mathfrak{E}^{-\frac{\chi^2}{4}} \right)$$

$$= \left(\frac{1}{2\hbar_{\eta+1}(\mathfrak{x})} + \frac{\eta}{2\hbar_{\eta-1}(\mathfrak{x})} - \hbar_{\eta+1}(\mathfrak{x}) \right) \mathfrak{E}^{-\frac{\chi^2}{4}} = \left(\frac{\eta}{2\hbar_{\eta-1}(\mathfrak{x})} - \frac{1}{2\hbar_{\eta+1}(\mathfrak{x})} \right) \mathfrak{E}^{-\frac{\chi^2}{4}}$$

$$\mathrm{d} \sum_{\alpha=1}^4 \mathfrak{F}_{\alpha\beta} \otimes \mathrm{d}\mathfrak{x}^\alpha = \sum_{\alpha=1}^4 \mathrm{d}_0 \mathfrak{F}_0 \mathrm{d}\mathfrak{x}^0 \wedge \mathrm{d}\mathfrak{x}^\alpha + \sum_{1 \leq i < j \leq 4} \langle -1 \rangle^{ij} \langle \partial_i \mathfrak{f}_j - \partial_j \mathfrak{f}_i \rangle \mathrm{d}\mathfrak{x}^i \wedge \mathrm{d}\mathfrak{x}^j$$

$$\langle \mathfrak{f} \sqrt{\phi_\kappa}, \mathfrak{g} \sqrt{\phi_\kappa} \rangle = \int_{\mathbb{R}^4} \mathfrak{f} \mathfrak{g} \cdot \phi_\kappa \, \mathrm{d}\lambda$$

$$\left\{ \left| \frac{\hbar_i(\kappa \chi^0) \hbar_j(\kappa \chi^1) \hbar_\dagger(\kappa \chi^2) \hbar_l(\kappa \chi^3)}{\sqrt{i!} j! \dagger! l! \sqrt{\phi_\kappa(\vec{\mathfrak{x}})}} \right|_{\vec{\mathfrak{x}} = (\chi^0, \chi^1, \chi^2, \chi^3) \in \mathbb{R}^4, i, j, \dagger, l \geq 0} \right\}$$



$$\langle \sum_{1 \leq \alpha \leq \beta \leq 4} \hat{f}_{\alpha\beta} \otimes dx^\alpha \wedge dx^\beta, \sum_{1 \leq \alpha \leq \beta \leq 4} \hat{f}_{\alpha\beta} \otimes dx^\alpha \wedge dx^\beta \rangle = \langle \hat{f}_{\alpha\beta}, \hat{f}_{\alpha\beta} \rangle$$

$$df = \sum_{i=1}^4 \partial_0 \hat{f}_i \otimes dx^0 \wedge dx^i + \sum_{1 \leq i \leq j \leq 4} (\partial_i \hat{f}_j - \partial_j \hat{f}_i) dx^i \wedge dx^j = \mathfrak{C}_\gamma^{\alpha\beta} = -\mathcal{T}_r(\mathfrak{C}^\gamma(\mathfrak{C}^\alpha, \mathfrak{C}^\beta)), \mathfrak{C}^\alpha, \mathfrak{C}^\beta, \mathfrak{C}^\gamma \in \mathfrak{g}$$

$$d\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U} = \sum_{\gamma=1}^{\eta} \left(\sum_{j=1}^4 \alpha_{0:i,\gamma} \otimes dx^0 \wedge dx^i + \sum_{1 \leq i \leq j \leq 4} \alpha_{i:j,\gamma} \otimes dx^i \wedge dx^j \right) + \sum_{1 \leq i \leq j \leq 4} \sum_{1 \leq \alpha, \beta \leq \eta} \alpha_{i,\alpha} \alpha_{j,\beta} \mathfrak{C}_\gamma^{\alpha\beta} \otimes dx^i \wedge dx^j \otimes \mathfrak{C}^\gamma$$

Cuya forma bilineal, se distribuye así:

$$\langle \sum_{\alpha} \mathfrak{F}_{\alpha} \otimes dx^{\alpha}, \sum_{\alpha} \mathfrak{G}_{\beta} \otimes dx^{\beta} \rangle_{\partial\kappa} = \kappa^4 \langle \frac{\partial}{\partial t} \left(\frac{i}{\hbar} \right) \sum_{\alpha} \mathfrak{F}_{\alpha} \otimes dx^{\alpha}, \sum_{\alpha} \mathfrak{G}_{\beta} \otimes dx^{\beta} \rangle$$

$$\psi_{\kappa}: \mathcal{H}_{\wp_{\mathfrak{R}}}(\kappa x^0, \kappa x^1, \kappa x^2, \kappa x^3) \sqrt{\phi_{\kappa}} / \sqrt{\wp_{\mathfrak{R}}}! \rightarrow \mathfrak{z}^{\mathfrak{p}_r} / \sqrt{\wp_{\mathfrak{R}}}! \equiv \mathfrak{z}_0^{i_0} \mathfrak{z}_1^{i_1} \mathfrak{z}_2^{i_2} \mathfrak{z}_3^{i_3} / \sqrt{i_0! i_1! i_2! i_3!}$$

Cuya isometría extendida, deriva en lo que sigue:

$$\psi_{\kappa} \left[\sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{F}_{\alpha,\beta} \otimes dx^{\alpha} \wedge dx^{\beta} \right] = \sum_{0 \leq \alpha \leq \beta \leq 4} \psi_{\kappa} |f_{\alpha,\beta}| \otimes dx^{\alpha} \wedge dx^{\beta}$$

Ahora bien, un espacio abstracto de Wiener, se explica así:

$$\mu_{\kappa}(\chi \in \mathfrak{P}^{-1}(\mathfrak{S})) = \left(\frac{\kappa}{2\varpi} \right)^{l/2} \int_{\eta \in \mathcal{F}} e^{-\kappa|\varphi|^2/\psi} d\eta$$

$$\left\{ \hat{\mathfrak{z}}^{\mathfrak{p}_r} \frac{\otimes dx^{\alpha}}{|\hat{\mathfrak{z}}^{\mathfrak{p}_r} \otimes dx^{\alpha}|_{\partial\kappa}} : \rho_r \in \wp_r, \mathcal{R} \geq 0 \right\}$$

$$\hat{\mathfrak{z}}^{\mathfrak{p}_r} \otimes dx^{\alpha} = \mathfrak{z}^{\mathfrak{p}_r} \otimes dx^{\alpha} - \frac{\sum_{\alpha=0}^4 \langle \mathfrak{z}^{\mathfrak{p}_r} \otimes dx^{\alpha}, \frac{\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}}{|\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}|_{\partial\kappa}} \rangle \mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}}{|\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}|_{\partial\kappa}} + |\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}|_{\partial\kappa}^2 - |\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}|_{\partial\kappa}$$

$$+ \langle \frac{\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}}{|\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}|_{\partial\kappa}} \rangle_{\partial\kappa} \bowtie 4\kappa(r-1)\sqrt{\wp_{\mathfrak{R}}}! \cdot \frac{\kappa\sqrt{\wp_{\mathfrak{R}}}!}{\mathcal{R}^2} / 2$$

$$\leq 6\kappa(r-1)^2 \sum_{\alpha=1}^4 \langle \frac{\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}}{|\mathfrak{z}^{\rho_r^{\alpha,-}} \otimes dx^{\alpha}|_{\partial\kappa}} \rangle_{\partial\kappa} \otimes dx^{\alpha}$$



$$\mu_\kappa(\|\mathcal{P}_b\chi\| > \epsilon) \leq \mu_\kappa\left(\vec{\mathfrak{z}} \in \mathfrak{B}\left(\frac{0,1}{2}\right)\right) \sum_s \sum_{\rho_\tau \geq \Omega^s} \mathbb{E}|\mathfrak{C}_s \alpha_{\rho_\tau, \mathfrak{B}}^s|^4 \left[|z^{\rho_\tau}| + 6\kappa(r+1)^2 \sum_{\alpha=1}^4 |z^{\rho_\tau^{\alpha,-}}| > \epsilon \right] 1/\sqrt{\kappa\epsilon}$$

$$|\langle \chi(\omega), \mathfrak{d}\mathfrak{x}^\alpha \rangle| \leq \left| \sum_\tau \sum_{\rho_\tau \in \mathfrak{q}_\tau} \mathfrak{C}_{\rho_\tau, \alpha} \widehat{\omega}^{\rho_\tau} / |\widehat{\omega}^{\rho_\tau} \otimes \mathfrak{d}\mathfrak{x}^\alpha|_{\partial\kappa} \right|$$

$$\leq \mathfrak{R}^{\mathfrak{M}} \sum_{\tau \leq \mathcal{M}} \sum_{\rho_\tau \in \mathfrak{q}_\tau} |\mathfrak{C}_{\rho_\tau, \alpha}| \|\omega/\mathfrak{R}\|^{\rho_\tau} + 6\kappa(r+1)^2 \sum_{\alpha=1}^4 \|\omega/\mathfrak{R}\|^{\rho_\tau^{\alpha,-}} 2\mathfrak{R}^{\mathfrak{S}} + 1/\kappa \sqrt{\mathfrak{f}\mathfrak{O}_{\mathfrak{R}}!} \|\chi\|$$

$$|\langle \kappa \partial \mathfrak{x}(\omega), \mathfrak{d}\mathfrak{x}^\alpha \wedge \mathfrak{d}\mathfrak{x}^\beta \rangle| \leq \kappa \left| \sum_\tau \sum_{\rho_\tau \in \mathfrak{q}_\tau} \frac{\mathfrak{C}_{\rho_\tau, \beta} \partial_\beta \widehat{\omega}^{\rho_\tau}}{|\widehat{\omega}^{\rho_\tau} \otimes \mathfrak{d}\mathfrak{x}^\beta|_{\partial\kappa}} \right| + \kappa \left| \sum_\tau \sum_{\rho_\tau \in \mathfrak{q}_\tau} \frac{\mathfrak{C}_{\rho_\tau, \alpha} \partial_\alpha \widehat{\omega}^{\rho_\tau}}{|\widehat{\omega}^{\rho_\tau} \otimes \mathfrak{d}\mathfrak{x}^\alpha|_{\partial\kappa}} \right|$$

$$\leq \mathfrak{R}^{\mathcal{M}} \sum_{\tau \leq \mathcal{M}} \sum_{\rho_\tau \in \mathfrak{q}_\tau} \frac{|\mathfrak{C}_{\rho_\tau, \beta}| + |\mathfrak{C}_{\rho_\tau, \alpha}| |\partial_\beta \widehat{\omega}^{\rho_\tau}|}{|\widehat{\omega}^{\rho_\tau} \otimes \mathfrak{d}\mathfrak{x}^\alpha|_{\partial\kappa}} \left| \frac{\omega}{\mathfrak{R}} \right|^{\rho_\tau}$$

$$+ 6\kappa(r+1)^2 \sum_{\alpha=1}^4 \left| \frac{\omega}{\mathfrak{R}} \right|^{\rho_\tau^{\alpha,-}} 4(r+1)\mathfrak{R}^{\mathfrak{S}} / \sqrt{\mathfrak{f}\mathfrak{O}_{\mathfrak{R}}!}$$

$$< (\mathfrak{R}^{\mathcal{M}} + 1) \|\lambda\|^{\mathfrak{N}} \mathfrak{g}_i \otimes \mathfrak{d}\mathfrak{x}^i \in \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) \rightarrow \psi_\kappa | \partial_\alpha \mathfrak{g}_\beta - \partial_\beta \mathfrak{g}_\alpha |$$

$$|(\omega + \omega_0)^{\rho_\tau} - \omega_0^{\rho_\tau}|$$

$$+ 6\kappa(r+1)^2 \sum_{\alpha=1}^4 |(\omega + \omega_0)^{\rho_\tau^{\alpha,-}} - \omega_0^{\rho_\tau^{\alpha,-}}| \langle \leq \mathfrak{R}^\tau | \omega \rangle \left(\omega \in \mathfrak{B}(0, \epsilon) \right) | \langle \mathfrak{x}(\omega + \omega_0) - \mathfrak{x}(\omega_0), \mathfrak{d}\mathfrak{x}^\alpha \rangle |$$

$$\leq \left(\sup_{\omega \in \mathfrak{B}(0, \epsilon)} \right) \sum_\tau \sum_{\rho_\tau \in \mathfrak{P}_\tau} |\mathfrak{C}_{\rho_\tau}| |2(\widehat{\omega + \omega_0})^{\rho_\tau} - \widehat{\omega_0}^{\rho_\tau}| \kappa \sqrt{\mathfrak{f}\mathfrak{O}_{\mathfrak{R}}!}$$

$$+ \frac{\left(\sup_{\omega \in \mathfrak{B}(0, \epsilon)} \right) \sum_{\alpha=1}^4 \sum_\tau \sum_{\rho_\tau \in \mathfrak{P}_\tau} |\mathfrak{C}_{\rho_\tau}| |6\kappa(r+1)^2 \cdot 2|(\widehat{\omega + \omega_0})^{\rho_\tau} - \widehat{\omega_0}^{\rho_\tau}|}{\sqrt{\kappa \mathfrak{f}\mathfrak{O}_{\mathfrak{R}}!}}$$

$$\leq \frac{4}{\kappa \left(\sup_{\omega \in \mathfrak{B}(0, \epsilon)} \right) \sum_\tau \sum_{\rho_\tau \in \mathfrak{P}_\tau} |\mathfrak{C}_{\rho_\tau}| | (4^\tau \mathfrak{R}^\tau | \omega |) }{\kappa \sqrt{\mathfrak{f}\mathfrak{O}_{\mathfrak{R}}!}}$$

$$+ 2\kappa \left(\sup_{\omega \in \mathfrak{B}(0, \epsilon)} \right) \sum_{\alpha=1}^4 \sum_\tau \sum_{\rho_\tau \in \mathfrak{P}_\tau} \frac{6\kappa(r+1)^2}{2^{\tau-2} |\mathfrak{C}_{\rho_\tau}|} (2^{\tau-2} \mathfrak{R}^{\tau-2} | \omega |) / \kappa \sqrt{\mathfrak{f}\mathfrak{O}_{\mathfrak{R}}!}$$

$$\leq 2\mathfrak{C}(\omega_0) / \kappa \cdot \epsilon |\chi|$$



$$\begin{aligned} & \widetilde{\mu}_\kappa \left(\left(\sup_{\omega \in \mathfrak{B} \left(\omega_0, \frac{1}{\kappa} \right)} |\mathfrak{x}(\omega + \omega_0) - \mathfrak{x}(\omega_0), \mathfrak{d}\mathfrak{x}^\alpha| > \epsilon \right) \right) \\ &= \widetilde{\mu}_\kappa \left(\left(\sup_{\omega \in \mathfrak{B} \left(\omega_0, \frac{1}{\kappa} \right)} |\mathfrak{x}, \zeta_\alpha(\omega + \omega_0) - \zeta_\alpha(\omega_0), \mathfrak{d}\mathfrak{x}^\alpha| > \epsilon \right) \right) \leq \widetilde{\mu}_\kappa \left(\frac{2\mathfrak{C}(\omega_0)}{\kappa \mathfrak{I} \|\mathfrak{x}\|} > \epsilon \right) \\ &\rightarrow 0 \end{aligned}$$

$$\left\langle \frac{\partial_\alpha \mathfrak{z}_\alpha^\eta}{\sqrt{\eta!}}, \frac{\partial_\beta \mathfrak{z}_\beta^\eta}{\sqrt{\eta!}} \right\rangle \geq \eta + \frac{1}{4 \left\langle \frac{\partial_\alpha \mathfrak{z}_\alpha^{\eta+1}}{\sqrt{\eta+1!}}, \frac{\partial_\beta \mathfrak{z}_\beta^{\eta+1}}{\sqrt{\eta+1!}} \right\rangle}$$

Siendo la integral de Yang – Mills, la siguiente:

$$\begin{aligned} & \frac{1}{3} \int_{\mathfrak{Y}} e^{\int_{\mathbb{C}^\infty} \sum_{j=1}^4 \mathfrak{Y}_j \otimes \mathfrak{d}\mathfrak{x}^j \otimes \mathfrak{i}} \mathfrak{E}^{-|\mathfrak{b}\mathfrak{Y}|^2/2} \mathfrak{D}\mathfrak{Y}, \mathfrak{i} = \sqrt{-1} \\ & \frac{1}{3} e^{-\frac{1}{2 \int_{\mathbb{R}^4} |\mathfrak{b}\mathfrak{Y} + \mathfrak{Y} \wedge \mathfrak{Y}|^4 \mathfrak{b}\omega}} \mathfrak{D}(\mathfrak{b}\mathfrak{Y}) \\ \mathfrak{Z} &= \int_{\{\mathfrak{b}\mathfrak{Y}: \mathfrak{Y} \in \mathfrak{S}_\kappa(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^4)_\mathfrak{g}\}} e^{-\frac{1}{2 \int_{\mathbb{R}^4} |\mathfrak{b}\mathfrak{Y} + \mathfrak{Y} \wedge \mathfrak{Y}|^4 \mathfrak{b}\omega}} \mathfrak{D}(\mathfrak{b}\mathfrak{Y}) \\ & \frac{1}{3} e^{-\frac{1}{2 \int_{\mathbb{C}^4} |\kappa \mathfrak{b}\mathfrak{Y} + \mathfrak{Y} \wedge \mathfrak{Y}|^4 \mathfrak{b}\lambda_4}} \mathfrak{D}(\mathfrak{b}\mathfrak{Y}) \\ \mathfrak{Z} &= \int_{\{\mathfrak{b}\mathfrak{Y}: \mathfrak{Y} \in \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^1(\mathbb{R}^4)_\mathfrak{g}\}} e^{-\frac{1}{2 \int_{\mathbb{C}^4} |\kappa \mathfrak{b}\mathfrak{Y} + \mathfrak{Y} \wedge \mathfrak{Y}|^4 \mathfrak{b}\lambda_4}} \mathfrak{D}(\mathfrak{b}\mathfrak{Y}) \end{aligned}$$

$$\mathbb{H} = \{(\mathfrak{b}_0 \mathcal{H}^2(\mathbb{C}^4)) \otimes (* \Lambda^2(\mathbb{R}^4))\} \oplus \{\mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4)\} \subset \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4)$$

$$\frac{1}{3} e^{-\frac{1}{2 \int_{\mathbb{C}^4} |\kappa \mathfrak{b}\mathfrak{Y} + \mathfrak{Y} \wedge \mathfrak{Y}|^4 \mathfrak{b}\lambda_4}} \mathfrak{D}(\mathfrak{b}\mathfrak{Y}) = \frac{\mathfrak{y}^\kappa \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}}{\int_{\mathbb{B} \otimes \mathfrak{g}} \mathfrak{y}^\kappa \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}} = \frac{\mathfrak{y}^\kappa \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}}{\mathbb{E}|\mathfrak{y}^\kappa|}$$

$$\mathbb{E}_{\mathfrak{Y} \otimes \mathfrak{g}}^\kappa [\mathcal{R}]^{\mathfrak{R}} = 1 / \int_{\mathbb{B} \otimes \mathfrak{g}} \mathfrak{y}^\kappa \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta} \int_{\mathbb{B} \otimes \mathfrak{g}} \mathfrak{S} \mathfrak{y}^\kappa \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}$$



$$\begin{aligned}
& \int_{\mathbb{R}^4} |\mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \, \mathfrak{d}\omega \sum_{1 \leq i \leq j \leq 4} \int_{\mathbb{R}^4} \left(\sum_{\alpha=1}^{\eta} \alpha_{i,j,\alpha}^2 + \sum_{\gamma=1}^{\eta} \sum_{\substack{\alpha,\beta \\ \hat{\alpha},\hat{\beta}}} \alpha_{i,\alpha} \alpha_{j,\beta} \alpha_{i,\hat{\alpha}} \alpha_{j,\hat{\beta}} \mathfrak{C}_{\gamma}^{\alpha\beta} \mathfrak{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} \right. \\
& \quad \left. + 2 \sum_{\gamma=1}^{\eta} \sum_{\substack{\alpha,\beta \\ \hat{\alpha},\hat{\beta}}} \alpha_{i,j,\gamma} \alpha_{i,\alpha} \alpha_{j,\beta} \mathfrak{C}_{\gamma}^{\alpha\beta} \right) \mathfrak{d}\omega + \sum_{i,j=1}^4 \int_{\mathbb{R}^4} \sum_{\alpha=1}^{\eta} \alpha_{0:i,\alpha}^2 \, \mathfrak{d}\omega \\
& \exp \left[-1/2 \sum_{\alpha=1}^{\eta} \int_{\mathbb{R}^4} \mathfrak{d}\omega \sum_{1 \leq i \leq j \leq 4} \int_{\mathbb{R}^4} \alpha_{i,j,\alpha}^2 + \sum_{i,j=1}^4 \int_{\mathbb{R}^4} \alpha_{0:i,\alpha}^2 \right] \mathcal{D}(\mathfrak{d}\mathfrak{U})
\end{aligned}$$

Más, usando el teorema de Stokes, obtenemos:

$$\begin{aligned}
& \frac{1}{3} \int_{\mathfrak{U}} e^{\int_{\mathfrak{s}}^{\infty} \mathfrak{d}\mathfrak{U} \otimes i} \mathfrak{E}^{-|\mathfrak{d}\mathfrak{U}|^2/2} \mathfrak{D}\mathfrak{U} \\
& \frac{1}{3} \int_{\substack{\mathfrak{U} \in \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4)}} e^{\int_{\mathfrak{s}}^{\infty} \kappa \mathfrak{d}\mathfrak{U}} e^{-|\kappa \mathfrak{d}\mathfrak{U}|^2/2} \mathfrak{D}\mathfrak{U} \\
& \nu_{\mathfrak{s}}^{\kappa} = \int_{\mathfrak{S}^2} \mathfrak{d}s \mathfrak{d}t \sum_{0 \leq \alpha \leq \beta \leq 4} \frac{\kappa^2}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, t) \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(\mathfrak{s}, t)}{2} \right) \\
& \Upsilon(\mathbb{R}^4, \kappa; \mathfrak{S}, i) = \mathbb{E}_{\mathfrak{U}} \left(\exp \left(\frac{1}{\kappa} (\cdot, \nu_{\mathfrak{s}}^{\kappa}) \right) \right) = \int_{\mathfrak{U} \in \mathfrak{B}(\mathbb{R}^4, \partial)} \left(\exp \left(\frac{1}{\kappa} (\mathfrak{U}, \nu_{\mathfrak{s}}^{\kappa} \otimes i) \right) \right) \mathfrak{d}\tilde{\mu}(\mathfrak{U}) \\
& \mathbb{E}_{\mathfrak{U}} = \left(\exp \left(\frac{1}{\kappa} \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{S}^2} \frac{\kappa^2}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, t) \left(\cdot, \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(\mathfrak{s}, t)}{2} \right) \otimes i \right) \mathfrak{d}s \mathfrak{d}t \right) \right) \\
& = \exp \left(-1/2 \left| \frac{\sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{S}^2} \mathfrak{d}s \mathfrak{d}t \kappa}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, t) \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(\mathfrak{s}, t)}{2} \right) \right|_{\partial, \kappa}^2 \right) \\
& \rightarrow \exp \left(-1/8 \int_{\mathfrak{S}} \rho_{\mathfrak{S}} \right) \\
& \psi_{\omega} = \psi(\omega) = \frac{e^{-\frac{|\omega|^2}{2}}}{\sqrt{2\varpi}} = \langle \psi_{\omega} \chi_{\omega}, \psi_{\nu} \chi_{\nu} \rangle = \psi_{\omega} \psi_{\nu} e^{\omega \nu} = \frac{1}{2\varpi} e^{-|\omega - \nu|^2/2}
\end{aligned}$$



$$\begin{aligned}
& \left\langle \sum_{i=1}^4 \sum_{\tau} \sum_{\rho_{\tau}} \mathfrak{C}_{\rho_{\tau}, i} 3^{\rho_{\tau}} \otimes dx^i, \xi_{\alpha\beta}^{\kappa}(\omega) \right\rangle_{\partial, \kappa} = \kappa^2 \left\langle \partial \sum_{i=1}^4 \sum_{\tau} \sum_{\rho_{\tau}} \mathfrak{C}_{\rho_{\tau}, i} 3^{\rho_{\tau}} \otimes dx^i, \partial \xi_{\alpha\beta}^{\kappa}(\omega) \right\rangle \\
& = \kappa \sum_{\tau} \sum_{\rho_{\tau}} \psi(\omega) (\mathfrak{C}_{\rho_{\tau}, \beta} \partial_{\alpha} \omega^{\rho_{\tau}} - \mathfrak{C}_{\rho_{\tau}, \alpha} \partial_{\beta} \omega^{\rho_{\tau}}) \\
& \exp\left(-\frac{1}{\kappa} \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{S}^2} ds dt |\mathfrak{S}_{\alpha\beta}^{\sigma}|(s, t) \frac{\kappa^2}{4} \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa\sigma(s, t)}{2}\right)\right)_{\partial, \kappa/2} \\
& = \exp\left(1/2 \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{S}^2, \bar{\mathfrak{S}}^2} ds dt d\bar{s} d\bar{t} \frac{\kappa^2}{16} |\mathfrak{S}_{\alpha\beta}^{\sigma}|(s, t) |\mathfrak{S}_{\alpha\beta}^{\sigma}|(\bar{s}, \bar{t}) \kappa^2 \left\langle \partial \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa\sigma(s, t)}{2}\right), \partial \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa\sigma(\bar{s}, \bar{t})}{2}\right) \right\rangle\right) \\
& = \exp(-1/8(2\pi) \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{S}^2, \bar{\mathfrak{S}}^2} ds dt d\bar{s} d\bar{t} \frac{\kappa^2}{4} |\mathfrak{S}_{\alpha\beta}^{\sigma}|(s, t) |\mathfrak{S}_{\alpha\beta}^{\sigma}|(\bar{s}, \bar{t}) e^{-\frac{\kappa^2 |\sigma(s, t) - \sigma(\bar{s}, \bar{t})|^2}{8}}) \\
& \rightarrow \exp\left(-1/8 \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{S}^2} ds dt |\mathfrak{S}_{\alpha\beta}^{\sigma}|(s, t) \rho_{\mathfrak{S}}^{\alpha\beta}(\sigma(s, t))\right) = e^{-1/8 \int_{\mathfrak{S}} \rho_{\mathfrak{S}}}
\end{aligned}$$

Campo Cuántico no Abeliano

En campos cuánticos no abelianos y bajo operadores holonómicos y por algebra de Lie, obtenemos:

$$\begin{aligned}
\mathcal{T} & = \int_{\mathfrak{A}} e^{\int_{\mathfrak{S}} \sum_{i=1}^4 \mathfrak{A}_i \otimes dx^i} e^{-1/2|F|^2} \mathfrak{D}\mathfrak{A} \\
\mathcal{T} \exp\left(\int_{\mathfrak{S}} \sum_{i=1}^4 \mathfrak{A}_i \otimes dx^i\right) & = \mathcal{T} \exp\left(\sum_i \int_{\mathfrak{S}_i} \sum_{j=1}^4 \mathfrak{A}_j \otimes dx^j\right) = \mathcal{T} \bigotimes_i \exp\left(\int_{\mathfrak{S}_i} \sum_{j=1}^4 \mathfrak{A}_j \otimes dx^j\right) = \mathcal{T} \bigotimes_i (\mathfrak{S} \\
& + \sum_{\alpha} |F_{\alpha} E^{\alpha}|(\mathfrak{S}_i) \Delta_i \mathfrak{S}^2 + \mathcal{O}(\Delta_i \mathfrak{S}^3)) \\
& = \mathcal{T} \bigotimes_i \exp(\log(\mathfrak{S} \\
& + \sum_{\alpha} |F_{\alpha} E^{\alpha}|(\mathfrak{S}_i) \Delta_i \mathfrak{S}^2 + \mathcal{O}(\Delta_i \mathfrak{S}^3))) = \mathcal{T} \exp\left(\sum_i \sum_{\alpha} |F_{\alpha} E^{\alpha}|(\mathfrak{S}_i) \Delta_i \mathfrak{S}^2 + \mathcal{O}(\Delta_i \mathfrak{S}^3))\right) \\
& \rightarrow \mathcal{T} \exp\left(\int_{\mathfrak{S}} \sum_{\alpha} F_{\alpha} E^{\alpha}\right)
\end{aligned}$$



Cuyo operador de translación en curvatura, es la que sigue:

$$\begin{aligned} & \mathfrak{X}_0^{1-} \mathfrak{X}_{1-}^1 \cdot \mathfrak{X}_{0+}^1 \cdot \overrightarrow{\prod_{i=2}^{\eta} \mathfrak{X}_0^{i-}} \cdot \mathfrak{X}_{1-}^{\eta} \cdot \overrightarrow{\prod_{i=2}^{\eta} \mathfrak{X}_{i-}^{\eta}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{X}_{\eta}^{i+}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{X}_{i+}^0} \\ &= \overrightarrow{\prod_{i=1}^{\eta} \mathfrak{X}_0^{i-}} \cdot \overrightarrow{\prod_{i=1}^{\eta} \mathfrak{X}_{i-}^{\eta}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{X}_{\eta}^{i+}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{X}_{i+}^0} \end{aligned}$$

En el que, la curvatura, se expresa así:

$$\begin{aligned} \mathfrak{X}_i^{(i+1)-} \mathfrak{X}_{(i+1)-}^{i+1-} \mathfrak{X}_{i+1}^{i+} \mathfrak{X}_{i+}^i \mathfrak{X}_i^i &= \mathfrak{X}_i^i + \epsilon^2 \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{S}_{\alpha\beta}^{\sigma} | \Omega_{i,\alpha\beta}^i \rangle \left(\sigma^{\dagger} \left(\frac{i}{\eta}, \frac{i}{\eta} \right), \sigma \left(\frac{i}{\eta}, \frac{i}{\eta} \right) \right) \mathfrak{X}_i^i + \mathcal{O}(\epsilon^4) \\ &= \mathfrak{X}_i^i + \epsilon^2 \mathfrak{X}_i^i \mathfrak{S} \Omega_i^i + \mathcal{O}(\epsilon^4) \boxtimes_i^i \end{aligned}$$

$$\begin{aligned} & \mathfrak{X}_i^{(i+1)-} \mathfrak{X}_{(i+1)-}^{i+1-} \mathfrak{X}_{i+1}^{i+} \mathfrak{X}_{i+}^i \cdot \mathfrak{X}_{(i+1)-}^i \mathfrak{X}_{i+1}^i \overrightarrow{\prod_{\kappa=i+1}^{\eta-1} (1 + \epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \\ &= \mathfrak{X}_i^{(i+1)-} \mathfrak{X}_{(i+1)-}^{i+1-} \mathfrak{X}_{i+1}^{i+} \mathfrak{X}_{i+}^i \mathfrak{X}_i^i \overrightarrow{\prod_{\kappa=i+1}^{\eta-1} (1 + \epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \\ &= \mathfrak{X}_i^i (1 + \epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i) \end{aligned}$$

$$\begin{aligned} \tilde{\mathfrak{X}}_1^{i+1} &= \tilde{\mathfrak{X}}_1^{i+} \tilde{\mathfrak{X}}_1^{i+} \cdot \overrightarrow{\prod_{\kappa=1}^{\eta-1} (1 + \mathfrak{S} \Omega_{\kappa}^{i+1} \epsilon^2 + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \\ &= \tilde{\mathfrak{X}}_1^{i+} \cdot \overleftarrow{\prod_{i=0}^{i-1} \tilde{\mathfrak{X}}_1^{i+}} \cdot \tilde{\mathfrak{X}}_{0+\mu_0}^0 \overleftarrow{\prod_{i=0}^i \tilde{\mathfrak{X}}_1^{i+}} \overrightarrow{\prod_{\kappa=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \\ &\quad \cdot \overrightarrow{\prod_{\kappa=1}^{\eta-1} (1 + \mathfrak{S} \Omega_{\kappa}^{i+1} \epsilon^2 + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \end{aligned}$$

$$\begin{aligned} \mathcal{T} e^{\int_c^{\infty} \mathcal{A}} \mu_0 &= \mu_0 \overrightarrow{\prod_{i=0}^{\eta-1} (1 + \epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \cdot \overleftarrow{\prod_{\kappa=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i)} \\ &= \mu_0 \overrightarrow{\prod_{i=0}^{\eta-1} e^{\epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \prod_{\kappa=1}^{\eta-1} e^{\epsilon^2 \mathfrak{S} \Omega_{\kappa}^i + \mathcal{O}(\epsilon^4) \boxtimes_{\kappa}^i}} \end{aligned}$$

$$\begin{aligned} & \rightarrow \mu_0 \cdot \tilde{\mathcal{T}} \exp \left(\int_{\mathfrak{S}^2} \sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{S}_{\alpha\beta}^{\sigma} (s, t) \Omega_{\beta}^{\alpha} (s, t) \mathrm{d}s \mathrm{d}t \right) \end{aligned}$$



Por lo que, finalmente la integral de Yang – Mills, equivale a lo que sigue:

$$\begin{aligned} & \exp -1/2 \left(\int_{\mathbb{C}^4} d\lambda_4 |\kappa \partial \mathfrak{A} + \mathfrak{A} \wedge \mathfrak{A}|^2 \right) \mathfrak{D} \mathfrak{A} \\ &= \exp(-1/2 \int_{\mathbb{C}^4} d\lambda_4 |\kappa \partial \mathfrak{A} + \mathfrak{A} \wedge \mathfrak{A}| + \langle \mathfrak{A} \wedge \mathfrak{A}, \kappa \partial \mathfrak{A} \rangle \\ & \quad + \langle \mathfrak{A} \wedge \mathfrak{A} \rangle^2) \exp(-1/2 \int_{\mathbb{C}^4} d\lambda_4 |\kappa \partial \mathfrak{A}|^2) \mathfrak{D} \mathfrak{A} \end{aligned}$$

$$\langle \kappa \partial \mathfrak{A}(\omega), d\mathbf{x}^\alpha \wedge d\mathbf{x}^\beta \rangle = \left(\mathfrak{A}, \tilde{\xi}_{\alpha\beta}^\kappa(\omega) \right) = \{ \mathfrak{A}_i, \mathfrak{A}_j \}(\omega) = \mathfrak{A}_i(\omega) \mathfrak{A}_j(\omega) = \left(\mathfrak{A} \otimes \mathfrak{A}, \zeta_i(\omega) \otimes \mathfrak{A}, \zeta_j(\omega) \right)$$

$$= \{ \mathfrak{A}_{i,\alpha} \mathfrak{A}_{j,\beta} \overline{\mathfrak{A}_{i,\alpha} \mathfrak{A}_{j,\beta}} \}(\omega) = \langle \mathfrak{A} \otimes \mathfrak{A} \otimes \overline{\mathfrak{A}} \otimes \overline{\mathfrak{A}}, \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\hat{\alpha},\omega} \otimes \chi_{j,\hat{\beta},\omega} \rangle$$

$$\mathfrak{A} = \sum_{i=1}^4 \sum_{\alpha=1}^{\eta} \mathfrak{A}_{i,\alpha} \otimes d\mathbf{x}^i \otimes \mathbb{F}^\alpha, \quad \tilde{\mathfrak{A}} = \sum_{i=1}^4 \sum_{\alpha=1}^{\eta} \tilde{\mathfrak{A}}_{i,\alpha} \otimes d\mathbf{x}^i \otimes \mathbb{F}^\alpha$$

$$\int_{\omega \in \mathbb{R}^4} \{ \mathfrak{A}_{i,\alpha} \mathfrak{A}_{j,\beta} \overline{\mathfrak{A}_{i,\alpha} \mathfrak{A}_{j,\beta}} \}(\omega) = \langle \mathfrak{A}^{\otimes 2} \otimes \tilde{\mathfrak{A}}^{\otimes 2}, \int_{\omega \in \mathbb{R}^4} d\omega \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\hat{\alpha},\omega} \otimes \chi_{j,\hat{\beta},\omega} \rangle$$

$$\left(\mathfrak{A}_{i_1,\alpha_1} \otimes \dots \otimes \mathfrak{A}_{i_4,\alpha_4}, \tilde{\omega}_\omega^{\otimes 4} \right) = \prod_{j=1}^4 \left(\mathfrak{A}_{i_j,\alpha_j}, \tilde{\pi}_{i_j,\alpha_j,\omega} \right)$$

$$= \left(\mathfrak{A}_{i_1,\alpha_1} \otimes \dots \otimes \mathfrak{A}_{i_3,\alpha_3}, \left(\tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathbb{F}^{\alpha_1} \otimes \tilde{\omega}_\omega^{\otimes 2} \right) \right)$$

$$= \left(\mathfrak{A}_{i_1,\alpha_1}, \tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathbb{F}^{\alpha_1} \right) \prod_{j=2}^4 \left(\mathfrak{A}_{i_j,\alpha_j}, \tilde{\pi}_{i_j,\alpha_j,\omega} \right)$$

$$= \left(\mathfrak{A}_{i_1,\alpha_1} \otimes \dots \otimes \mathfrak{A}_{i_4,\alpha_4}, \tilde{\omega}_\omega^{\otimes 4} \otimes \left(\tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathbb{F}^{\alpha_4} \right) \right)$$

$$= \left(\mathfrak{A}_{i_4,\alpha_4}, \tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathbb{F}^{\alpha_4} \right) \prod_{j=1}^4 \left(\mathfrak{A}_{i_j,\alpha_j}, \tilde{\pi}_{i_j,\alpha_j,\omega} \right)$$



$$\begin{aligned}
\int_{\mathbb{C}^4} d\lambda_4 |\mathfrak{A} \wedge \mathfrak{A}|^2 &= \sum_{\gamma} \sum_{0 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \frac{1}{2 \mathfrak{C}_{\gamma}^{\alpha\beta}} \mathfrak{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} \left(\left(\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{j,\beta} \otimes \mathfrak{A}_{i,\hat{\alpha}} \otimes \mathfrak{A}_{j,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) \tilde{\omega}_{\omega}^{\otimes 4} \right) \right. \\
&\quad \left. + \left(\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{j,\beta} \otimes \mathfrak{A}_{i,\hat{\alpha}} \otimes \mathfrak{A}_{j,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) \tilde{\omega}_{\omega}^{\otimes 4} \right) \right) \\
\int_{\mathbb{C}^4} d\lambda_4 |\kappa \partial \mathfrak{A}, \mathfrak{A} \wedge \mathfrak{A}|^4 &= \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{A}_{\kappa,\gamma} \otimes \mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{j,\beta}, \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{C}_{\gamma}) \otimes \tilde{\omega}_{\omega}^{\otimes 4} \right) \\
&+ \int_{\mathbb{C}^4} d\lambda_4 |\mathfrak{A} \wedge \mathfrak{A}, \kappa \partial \mathfrak{A}|^4 \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{j,\beta} \otimes \mathfrak{A}_{\kappa,\gamma}, \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{C}_{\gamma}) \otimes \tilde{\omega}_{\omega}^{\otimes 4} \right) \\
\mathfrak{Y}_{\mathbb{C}}^{\kappa} &= \left(\{\mathfrak{A}_{i,\alpha}\}_{i,\alpha} \right) \\
&= \mathfrak{I}_r \hat{\mathcal{F}} \exp\left(\frac{1}{\kappa} \cdot \frac{\kappa^2}{4}\right) \\
&\cdot \rho \kappa \\
&/ 4 \int_{\mathfrak{S}^2} ds dt \mu_{s,t}^{-1} \left(\sum_{1 \leq i \leq j \leq 4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(s,t) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\bar{s}, \bar{t}) \sum_{\alpha} (\mathfrak{A}_{\alpha}, \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa\sigma(s,t)}{2}\right) \otimes \mathfrak{C}^{\alpha} \otimes \rho(\mathfrak{C}^{\alpha})) \right. \\
&\quad \left. + \sum_{1 \leq i \leq j \leq 4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(s,t) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\bar{s}, \bar{t}) \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{j,\beta} \otimes \mathfrak{A}_{\kappa,\gamma} \otimes \tilde{\omega}_{\left(\frac{\kappa\sigma(s,t)}{2}\right)}^{\otimes 4} \otimes \rho(\mathfrak{C}_{\gamma}) \right) \mu_{s,t} \right)
\end{aligned}$$

$$\begin{aligned}
\mathfrak{Y}_{\mathbb{C}}^{\kappa} &= (\{\mathfrak{A}_{i,\alpha}\}_{i,\alpha}) \\
&= \exp\left(-\frac{1}{2} \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{A}_{\kappa,\gamma} \otimes \mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{i,\beta} \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{C}^{\gamma}) \otimes \tilde{\omega}_{\omega}^{\otimes 4} \right)\right) \\
&\quad - 1/2 \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{i,\beta} \otimes \mathfrak{A}_{\kappa,\gamma} \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) \tilde{\omega}_{\omega}^{\otimes 4} \otimes (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{C}^{\gamma}) \right) \\
&\quad - 1/2 \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \mathfrak{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} / 2 \left(\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{i,\beta} \otimes \mathfrak{A}_{i,\hat{\alpha}} \otimes \mathfrak{A}_{i,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) \tilde{\omega}_{\omega}^{\otimes 4} \right) \\
&\quad + (\mathfrak{A}_{i,\alpha} \otimes \mathfrak{A}_{i,\beta} \otimes \mathfrak{A}_{i,\hat{\alpha}} \otimes \mathfrak{A}_{i,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) \tilde{\omega}_{\omega}^{\otimes 4}))) \\
&\left| \left(\frac{\kappa}{4} \int_{\mathfrak{S}^2} d\mathfrak{s} d\mathfrak{t} \sum_{1 \leq i \leq j \leq 4} |\mathfrak{S}_{ij}^{\sigma}(\mathfrak{s}, \mathfrak{t})| |\mathfrak{S}_{ij}^{\sigma}(\bar{\mathfrak{s}}, \bar{\mathfrak{t}})| \sum_{\gamma} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_{\gamma}^{\alpha\beta}| |\pi_i(\omega)|_{\partial, \kappa} \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\alpha} \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\beta} \right) \right| \\
&\leq \left(\frac{\kappa}{4} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{S}^2} d\mathfrak{s} d\mathfrak{t} \sum_{1 \leq i \leq j \leq 4} |\mathfrak{S}_{ij}^{\sigma}(\mathfrak{s}, \mathfrak{t})| |\mathfrak{S}_{ij}^{\sigma}(\bar{\mathfrak{s}}, \bar{\mathfrak{t}})| \sum_{\gamma} \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_{\gamma}^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\alpha} \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\beta} \right| = F_{\eta} \\
&= 4 \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \sum_{1 \leq i \leq j \leq 4} \sum_{\rho, q=1}^{\eta} |\mathfrak{S}_{ij}^{\sigma}\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)| / \eta^2 \sum_{\gamma} \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_{\gamma}^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa \sigma\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)}{2}\right)}^{\alpha} \mathcal{M}_{\left(\frac{i, \kappa \sigma\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)}{2}\right)}^{\beta} \right|
\end{aligned}$$

$$\begin{aligned}
& \left| \left(\frac{\kappa}{4} \right) \int_{\mathfrak{J}^2} \mathfrak{d}s \mathfrak{d}t \sum_{1 \leq \alpha \leq \beta \leq 4} |\mathfrak{S}_{\alpha\beta}^\sigma|(s, t) |\mathfrak{S}_{\alpha\beta}^\sigma|(\bar{s}, \bar{t}) \sum_{\gamma} \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{\alpha\beta}| |\pi_{\alpha\beta}(\omega)|_{\partial, \kappa} \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(s, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(s, t)}{2}\right)}^j \right| \\
& \leq \left(\frac{\kappa}{4} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{J}^2} \mathfrak{d}s \mathfrak{d}t \sum_{1 \leq i \leq j \leq 4} |\mathfrak{S}_{\alpha\beta}^\sigma|(s, t) |\mathfrak{S}_{\alpha\beta}^\sigma|(\bar{s}, \bar{t}) \sum_{\gamma} \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{\alpha\beta}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(s, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(s, t)}{2}\right)}^j \right| = F_\eta \\
& = 4 \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \sum_{1 \leq i \leq j \leq 4} \sum_{\rho, q=1}^{\eta} |\mathfrak{S}_{\alpha\beta}^\sigma| \left(\frac{\rho}{\eta}, \frac{q}{\eta} \right) / \eta^2 \sum_{\gamma} \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{ij}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)}{2}\right)}^j \right|
\end{aligned}$$

Cuya aproximación riemmaniana equivale a:

$$\begin{aligned}
& \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{J}^2} \mathfrak{d}s \mathfrak{d}t \sum_{1 \leq i \leq j \leq 4} |\mathfrak{S}_{ij}^\sigma|(s, t) |\mathfrak{S}_{ij}^\sigma|(\bar{s}, \bar{t}) \sum_{\gamma} \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_\gamma^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa\sigma(s, t)}{2}\right)}^\alpha \mathcal{M}_{\left(\frac{j, \kappa\sigma(s, t)}{2}\right)}^\beta \right| \\
& = \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_\gamma^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa\sigma(s, t)}{2}\right)}^\alpha \mathcal{M}_{\left(\frac{j, \kappa\sigma(s, t)}{2}\right)}^\beta \right| \\
& \leq \mathfrak{N} \|\mathfrak{B}(\gamma)\| \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{i, \kappa\sigma(s, t)}{2}\right)}^\alpha} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{j, \kappa\sigma(s, t)}{2}\right)}^\alpha} / \sqrt{\sum_{\beta} \mathcal{M}_{\left(\frac{i, \kappa\sigma(s, t)}{2}\right)}^\beta} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{j, \kappa\sigma(s, t)}{2}\right)}^\beta} \\
& \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{J}^2} \mathfrak{d}s \mathfrak{d}t \sum_{1 \leq \alpha \leq \beta \leq 4} |\mathfrak{S}_{\alpha\beta}^\sigma|(s, t) |\mathfrak{S}_{\alpha\beta}^\sigma|(\bar{s}, \bar{t}) \sum_{\gamma} \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{ij}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(s, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(s, t)}{2}\right)}^j \right| \\
& = \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{ij}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(s, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(s, t)}{2}\right)}^j \right| \\
& \leq \mathfrak{N} \|\mathfrak{B}(\gamma)\| \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(s, t)}{2}\right)}^i} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(s, t)}{2}\right)}^j} / \sqrt{\sum_{\beta} \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(s, t)}{2}\right)}^i} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(s, t)}{2}\right)}^j}
\end{aligned}$$

Ahora bien, para efectos de simular superficies temporales y espaciales respectivamente, en campos cuánticos, se expresa lo que sigue, empezando por la transformación de Fourier:

$$\begin{aligned} & \left(\frac{1}{\sqrt{2\omega}}\right)^4 \int_{\delta_0}^{\infty} e^{-i(q^2x^2+q^4x^4)} \mathfrak{F}_\alpha(\mathcal{H}(\rho), \mathfrak{P}(\rho), x^2, x^4) dx^2 dx^4 \otimes \rho(\mathfrak{C}^\alpha) \\ &= \mathfrak{F}_\alpha(\mathcal{H}(\rho), \mathfrak{P}(\rho), q^2, q^4) \otimes \rho(\mathfrak{C}^\alpha) \end{aligned}$$

$$\begin{aligned} & \mathfrak{F}(\mathcal{H}(\rho)\tilde{f}_0 + \mathfrak{P}(\rho)\tilde{f}_1 + q^2\tilde{f}_2 + q^4\tilde{f}_4) \\ &= \frac{e^{-i(\alpha^0)\mathcal{H}(\rho)-\alpha^1\mathfrak{P}(\rho)}}{2\pi} \int_{\mathfrak{s} \in \mathbb{R}^4} e^{-i(\mathfrak{s}q^2+\mathfrak{s}q^4)\mathfrak{f}\{\hat{f}_0, \hat{f}_1\}}(\mathcal{H}(\rho), \mathfrak{P}(\rho))(\mathfrak{s}\hat{f}_2 + \mathfrak{s}\hat{f}_4) d\mathfrak{s} \\ &\equiv e^{-i(\alpha^0)\mathcal{H}(\rho)-\alpha^1\mathfrak{P}(\rho)}\hat{f}(\mathcal{H}(\rho)\hat{f}_0 + \mathfrak{P}(\rho)\hat{f}_1 + q^2\hat{f}_0 + q^4\hat{f}_4) \otimes \rho(\mathfrak{C}^\alpha) \end{aligned}$$

$$\langle \phi^{\alpha, \eta}(\mathfrak{f})1, \phi^{\nu\beta, \eta}(\mathfrak{g})1 \rangle = \mathfrak{C}(\rho_\eta) \mathcal{T}r(-\mathfrak{F}^\alpha \mathfrak{F}^\beta) \int_{\delta_0}^{\infty} (\mathfrak{f}\{\epsilon_0 \epsilon_1\} \overline{\mathfrak{g}\{\epsilon_0 \epsilon_1\}}) (\mathcal{H}(\rho_\eta), \mathfrak{P}(\rho_\eta))(\mathfrak{s}) d\mathfrak{s}$$

$$\begin{aligned} \phi^{\alpha, \eta}(\mathfrak{f})1 &= \int_{\vec{x} \in \mathbb{R}^4} d\vec{x} \mathfrak{f}(\vec{x}) \phi^{\alpha, \eta}(\vec{x})1 = \mathcal{U}(\vec{\alpha}, 1) \phi^{\alpha, \eta}(\vec{x}) \mathcal{U}(\vec{\alpha}, 1)^{-1} = \phi^{\alpha, \eta}(\vec{x} + \vec{\alpha}) \\ &= \frac{1}{2\omega} e^{i(x^0 \mathcal{H}(\rho_\eta) - x^1 \mathfrak{P}(\rho_\eta))} \delta(\cdot - (x^2, x^4)) \otimes \frac{\rho(\mathcal{F}^\alpha) \kappa^2}{4(2\omega)} \exp \delta(\vec{x} - \vec{\eta})^4 \end{aligned}$$

$$\begin{aligned} & \left\langle \frac{1}{(2\omega)^2} e^{i(\mathcal{H}(x^0 - \eta^0) - \mathfrak{P}(x^1 - \eta^1))} \rho_\kappa^{\vec{x}+} \otimes \rho(\mathcal{F}^\alpha), \frac{1}{(2\omega)^2} e^{i(\eta^0 \mathcal{H} - \eta^1 \mathfrak{P})} \rho_\kappa^{\eta+} \otimes \rho(\mathcal{F}^\beta) \right\rangle \\ &= \frac{\kappa^2}{(2\omega)^4 e^{i(\mathcal{H}(x^0 - \eta^0) - \mathfrak{P}(x^1 - \eta^1))} \exp(-\kappa^2 \delta(\vec{x} - \vec{\eta})^4)} \cdot \langle \rho_\eta(\mathcal{F}^\alpha), \rho_\eta(\mathcal{F}^\beta) \rangle \end{aligned}$$

$$\frac{1}{(2\omega)^2} \int_{\vec{x}, \vec{\eta} \in \mathbb{R}^4} \mathfrak{f}(\vec{x}) \mathfrak{g}(\vec{\eta}) \langle e^{i(x^0 \mathcal{H} - x^1 \mathfrak{P})} \rho_\kappa^{\vec{x}+} \otimes \rho(\mathcal{F}^\alpha), e^{i(\eta^0 \mathcal{H} - \eta^1 \mathfrak{P})} \rho_\kappa^{\eta+} \otimes \rho(\mathcal{F}^\beta) \rangle d\vec{x} d\vec{\eta}$$

$$= \frac{\kappa^2}{4} \int_{\mathfrak{s}, \hat{t} \in \mathbb{S}_0} \mathfrak{f}\{\epsilon_0 \epsilon_1\}(\mathfrak{s}) \overline{\mathfrak{g}\{\epsilon_0 \epsilon_1\}}(\hat{t}) \frac{1}{(2\omega)^2} \exp(-\kappa^2 |\mathfrak{s} - \hat{t}|^4)$$

$$/4) d\mathfrak{s} d\hat{t} \langle \rho_\eta(\mathcal{F}^\alpha), \rho_\eta(\mathcal{F}^\beta) \rangle \otimes \rho(\mathcal{F}^\alpha), \otimes \rho(\mathcal{F}^\beta) \rangle$$



$$\begin{aligned}
& \lambda \left(\mathfrak{Z}, \mathfrak{F}_\alpha \otimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}\}_{\alpha=0}^4 \right) + \mu \left(\mathfrak{Z}, \mathfrak{G}_\alpha \otimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}\}_{\alpha=0}^4 \right) \\
&= \left(3\mathfrak{U}\mathfrak{Z}(\lambda\widehat{\mathfrak{F}_\alpha} + \mu\widehat{\mathfrak{G}_\alpha}) \otimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}\}_{\alpha=0}^4 \right) \\
&\langle (\mathfrak{Z}, \mathfrak{F}_\alpha \otimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}\}_{\alpha=0}^4), (\mathfrak{Z}, \mathfrak{G}_\beta \otimes \rho(\mathfrak{E}^\beta), \{\widehat{\mathfrak{G}_\beta}\}_{\alpha=0}^4) \rangle \\
&= \int_{3\mathfrak{U}\mathfrak{Z}}^{\infty} \|\widehat{\mathfrak{F}_\alpha}\widehat{\mathfrak{G}_\beta}\| \cdot \mathfrak{d}|\rho| \cdot \mathfrak{T}_r(-\rho(\mathfrak{E}^\alpha)\rho(\mathfrak{E}^\beta)) \\
&= \sum_{\alpha=1}^{\eta} \mathfrak{E}(\rho) \int_{\mathfrak{Z}^2}^{\infty} \|\widehat{\mathfrak{F}_\alpha} \cdot \widehat{\mathfrak{G}_\alpha}\|(\sigma(\mathfrak{Z})) \left| \sum_{0 \leq \alpha \leq \beta \leq 4} \rho_\sigma^{\alpha\beta}(\mathfrak{Z})(\det \mathfrak{K}_\sigma^{\alpha\beta}(\mathfrak{Z})) \right| \mathfrak{d}\mathfrak{z} \\
&\langle \mathfrak{u}(\vec{\alpha}, \Lambda)(\mathfrak{S}, \mathfrak{F}_\alpha \otimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}\}_{\alpha=0}^4), \mathfrak{u}(\vec{\alpha}, \Lambda)(\mathfrak{S}, \mathfrak{G}_\beta \otimes \rho(\mathfrak{E}^\beta), \{\widehat{\mathfrak{F}_\alpha}\}_{\alpha=0}^4) \rangle \\
&= \int_{(\Lambda\mathfrak{S}+\vec{\alpha}) \cap (\Lambda\mathfrak{S}+\vec{\alpha})}^{\infty} \mathfrak{d}|\rho| e^{-i(\vec{\alpha} \cdot (\mathcal{H}(\rho_\eta)\Lambda\widehat{\mathfrak{F}_0} + \mathfrak{F}(\rho_\eta)\Lambda\widehat{\mathfrak{F}_1})} e^{i(\vec{\alpha} \cdot (\mathcal{H}(\rho_\eta)\Lambda\widehat{\mathfrak{F}_0} + \mathfrak{F}(\rho_\eta)\Lambda\widehat{\mathfrak{F}_1})} \\
&\times \|\widehat{\mathfrak{F}_\alpha}\widehat{\mathfrak{G}_\beta}\|(\Lambda^{-1}(\cdot - \vec{\alpha})) \cdot \mathfrak{T}_r(-\rho(\mathfrak{E}^\alpha)\rho(\mathfrak{E}^\beta)) \\
&= \int_{\Lambda(3\mathfrak{U}\mathfrak{Z})+\vec{\alpha}}^{\infty} \|\widehat{\mathfrak{F}_\alpha}\widehat{\mathfrak{G}_\beta}\|(\Lambda^{-1}(\cdot - \vec{\alpha})) \cdot \mathfrak{d}|\rho| \cdot \mathfrak{T}_r(-\rho(\mathfrak{E}^\alpha)\rho(\mathfrak{E}^\beta)) \\
&= \int_{3\mathfrak{U}\mathfrak{Z}}^{\infty} \|\widehat{\mathfrak{F}_\alpha}\widehat{\mathfrak{G}_\beta}\|(\cdot) \cdot \mathfrak{d}|\rho| \cdot \mathfrak{T}_r(-\rho(\mathfrak{E}^\alpha)\rho(\mathfrak{E}^\beta))
\end{aligned}$$

Siendo los operadores relativos a los campos cuánticos, los que siguen:

$$\mathcal{D}^{\vec{\kappa}} = \left(\frac{\partial}{\partial \mathfrak{x}^0} \right)^{\kappa^0} \left(\frac{\partial}{\partial \mathfrak{x}^1} \right)^{\kappa^1} \left(\frac{\partial}{\partial \mathfrak{x}^2} \right)^{\kappa^2} \left(\frac{\partial}{\partial \mathfrak{x}^3} \right)^{\kappa^3}, \vec{\mathcal{X}}^{\vec{\kappa}} = (\partial \mathfrak{x}^0)^{\kappa^0} (\partial \mathfrak{x}^1)^{\kappa^1} (\partial \mathfrak{x}^2)^{\kappa^2} (\partial \mathfrak{x}^3)^{\kappa^3}$$

$$\begin{aligned}
\phi^{\alpha,\eta}(\vec{\mathfrak{f}})1 &= \left(\mathfrak{S}_0, \mathfrak{f}^{\{\widehat{\mathfrak{E}_0}, \widehat{\mathfrak{E}_1}\}} \otimes \rho_\eta(\mathfrak{F}^\alpha), \{\epsilon_\alpha\}_{\alpha=0}^4 \right) \equiv \left(\mathfrak{S}_0, \mathfrak{f}^{\{\widehat{\mathfrak{E}_0}, \widehat{\mathfrak{E}_1}\}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{F}}(\rho_\eta)) \otimes \rho_\eta(\mathfrak{F}^\alpha), \{\epsilon_\alpha\}_{\alpha=0}^4 \right) \\
&\in \mathfrak{H}(\rho_\eta)
\end{aligned}$$

En dimensión \mathbb{R}^4 tenemos lo que sigue:

$$\begin{aligned}
\vec{\mathfrak{x}} \in \mathbb{R}^4 &\rightarrow \vec{\mathfrak{f}}^{\{\widehat{\mathfrak{f}_0}, \widehat{\mathfrak{f}_1}\}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{F}}(\rho_\eta))(\vec{\mathfrak{x}}) = \int_{\mathfrak{S}^b}^{\infty} e^{-i\vec{\mathfrak{m}}(\cdot)} / 2\pi\vec{\mathfrak{f}}(\vec{\mathfrak{x}} + \star) \mathfrak{d}|\rho| \\
&= \int_{\widehat{\mathfrak{S}} \in \mathbb{R}^4}^{\infty} e^{-i(\sigma(\widehat{\mathfrak{S}}) \cdot (\widehat{\mathcal{H}}(\rho_\eta)\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}(\rho_\eta)\widehat{\mathfrak{f}}_1))} / 2\pi\vec{\mathfrak{f}}(\vec{\mathfrak{x}} + \sigma(\widehat{\mathfrak{S}})) \cdot |\widehat{\rho}_\sigma|(\widehat{\mathfrak{S}}) \mathfrak{d}\widehat{\mathfrak{S}}
\end{aligned}$$



$$\left\{ \alpha_0 1 + \sum_{\eta, \mu=1}^{\infty} (\mathfrak{S}_{\eta, \mu}, \mathfrak{f}_{\eta, \alpha}^{\mu} \otimes \rho_{\eta}(\mathfrak{G}^{\alpha}), [\widehat{\mathfrak{f}}_{\alpha}^{\eta, \mu}]_{\alpha=0}^4) \boxtimes \alpha_0 \in \mathfrak{C}, \mathfrak{f}_{\eta, \alpha}^{\mu} \in \mathfrak{P}_{\mathfrak{S}_{\eta, \mu}}, \mathfrak{S}_{\eta, \mu} \in \mathfrak{L} \right\}$$

Cuya parametrización va como se indica:

$$\int_{\mathfrak{S}^2} |\mathfrak{f} \circ \sigma|^2(\hat{\mathfrak{s}}) \left| \sum_{0 \leq \alpha \leq \beta \leq 4} \rho_{\sigma}^{\alpha\beta}(\mathfrak{Z}) (\det \mathfrak{K}_{\alpha\beta}^{\sigma}(\mathfrak{Z})) \right| d\mathfrak{s} < \infty$$

$$\phi^{\alpha, \eta}(\mathfrak{f}) \sum_{\mu=1}^{\infty} (\mathfrak{S}_{\mu}, \mathfrak{G}_{\beta}^{\mu} \otimes \rho_{\mu}(\mathfrak{G}^{\beta}), \{\widehat{\mathfrak{f}}_{\alpha}^{\mu}\}_{\alpha=0}^4)$$

$$= \left\{ \sum \left| \frac{\partial \alpha}{\partial \beta} \frac{\partial \beta}{\partial \gamma} \frac{\partial \gamma}{\partial \delta} \frac{\partial \delta}{\partial \varepsilon} \right|^{i/h} \circ \left| \frac{\partial \zeta}{\partial \eta} \frac{\partial \eta}{\partial \theta} \frac{\partial \theta}{\partial \vartheta} \frac{\partial \vartheta}{\partial \iota} \right|^{i/\hbar} \left(\frac{\partial \kappa}{\partial \lambda} \frac{\partial \lambda}{\partial \mu} \frac{\partial \mu}{\partial \nu} \frac{\partial \nu}{\partial \xi} \right)^{\sqrt{i/\hbar}} \left(\sum_{\mu=1}^{\infty} (\mathfrak{S}_{\mu}, \mathfrak{f}_{\eta}^{\{\mu, \mu\}} \mathfrak{A}(\Lambda^{\mu})_{\gamma}^{\alpha} \cdot \mathfrak{G}_{\beta}^{\mu} \otimes \rho_{\eta} \|\mathfrak{F}^{\gamma}, \mathfrak{G}^{\beta}\|, \{\widehat{\mathfrak{f}}_{\alpha}^{\mu}\}_{\alpha=0}^4), \begin{matrix} \mathfrak{m} = \eta \\ \mathfrak{m} \neq \eta \end{matrix} \right) \right\}$$

$$\left(\sum \left| \frac{\partial \alpha}{\partial \beta} \frac{\partial \beta}{\partial \gamma} \frac{\partial \gamma}{\partial \delta} \frac{\partial \delta}{\partial \varepsilon} \right|^{i/h} \circ \left| \frac{\partial \zeta}{\partial \eta} \frac{\partial \eta}{\partial \theta} \frac{\partial \theta}{\partial \vartheta} \frac{\partial \vartheta}{\partial \iota} \right|^{i/\hbar} \left(\frac{\partial}{\partial t} \right) * \left\| \partial_0 \otimes \partial_{\rho} \otimes \partial_{\varrho} \otimes \partial_{\sigma} \otimes \partial_{\varsigma} \otimes \frac{\partial \tau}{\partial \nu} \otimes \partial_{\varphi} \otimes \partial_{\phi} + \partial_{\psi} + \partial_{\Psi} - \frac{\partial \Delta}{\partial \omega} / \Lambda_{\nu\mu}^{\mu\nu} \cdot \partial_{\Omega} \Phi \right\|^{i/\hbar} \right)$$

$$\phi^{\alpha, \eta}(\mathfrak{f}) \sum_{\mathfrak{m}=0}^{\infty} \nu_{\mathfrak{m}} = \sum_{\mathfrak{m}=0}^{\infty} \phi^{\alpha, \eta}(\mathfrak{f}) \nu_{\mathfrak{m}} = \alpha_0 \left(\delta_0, \mathfrak{f}_{\eta}^{\{\varepsilon_0, \varepsilon_1\}} \otimes \rho_{\eta}(\mathfrak{G}^{\alpha}), \{\varepsilon_{\alpha}\}_{\alpha=0}^4 \right) + \phi^{\alpha, \eta}(\mathfrak{f}) \nu_{\eta}$$

$$\phi^{\beta, \mathfrak{m}}(\mathfrak{f}) \sum_{\eta=0}^{\infty} \nu_{\eta} = \sum_{\eta=0}^{\infty} \phi^{\beta, \mathfrak{m}}(\mathfrak{f}) \nu_{\eta} = \beta_0 \left(\delta_0, \mathfrak{f}_{\mathfrak{m}}^{\{\varepsilon_0, \varepsilon_1\}} \otimes \rho_{\mathfrak{m}}(\mathfrak{G}^{\beta}), \{\varepsilon_{\beta}\}_{\beta=0}^4 \right) + \phi^{\beta, \mathfrak{m}}(\mathfrak{f}) \nu_{\mathfrak{m}}$$

$$\begin{aligned} & \phi^{\alpha, \eta}(\mathfrak{g})^* \left(\mathfrak{S}, \mathfrak{f}_{\beta} \otimes \rho_{\eta}(\mathfrak{G}^{\beta}), \{\widehat{\mathfrak{f}}_{\alpha}\}_{\beta=0}^4 \right) \\ &= - \left(\mathfrak{S}, \mathfrak{g}^{\{\widehat{\mathfrak{f}}_{\alpha}^{\alpha\beta}\}} \mathfrak{A}(\Lambda)_{\gamma}^{\alpha} \cdot \mathfrak{f}_{\beta} \otimes \rho_{\eta} \langle \mathfrak{F}^{\gamma}, \mathfrak{G}^{\beta} \rangle, \{\widehat{\mathfrak{f}}_{\alpha}\}_{\alpha=0}^4 \right) \\ &+ \langle \left(\mathfrak{S}, \mathfrak{f}_{\beta} \otimes \rho_{\eta}(\mathfrak{G}^{\beta}), \{\widehat{\mathfrak{f}}_{\alpha}\}_{\alpha=0}^4 \right), \phi^{\alpha, \eta}(\mathfrak{g}) 1 \rangle 1 \end{aligned}$$

En el que, la ciclicidad del campo cuántico, se expresa así:

$$\|\mathfrak{f} - \tilde{\mathfrak{g}}_{\varepsilon}\|_{\mathcal{L}^4} = \left(\int_{\mathfrak{S}^4} |1 - \mathfrak{g}_{\delta}^{\eta-1}|^4(\hat{\mathfrak{t}}) |\mathfrak{f}|^4(\hat{\mathfrak{t}}) d\hat{\mathfrak{t}} \right)^{\frac{1}{2}} \leq \mathcal{M} \|1 - \mathfrak{g}_{\delta}^{\eta-1}\|_{\mathcal{L}^4} < \varepsilon$$

$$\mathfrak{F}^{\gamma} = \sum_{\beta=1}^{\eta(\gamma)} \mathfrak{d}_{\mathfrak{m}, \beta}^{\gamma} \text{ad} \left(\mathfrak{F}^{\alpha^{\gamma, \beta}} \right) \cdots \text{ad} \left(\mathfrak{F}^{\alpha^{\gamma, \beta-1}} \right) \mathfrak{F}^{\alpha^{\gamma, \beta}}$$

Cuya función de Schwartz, en \mathbb{R}^4 es igual a:



$$\sum_{\beta=1}^{\eta(\gamma)} \mathfrak{d}_{\mathfrak{M},\beta}^{\gamma} \phi^{\alpha_{\eta}^{\gamma,\beta}}(\mathfrak{F}_{\eta}) \cdots \phi^{\alpha_{\mathfrak{M}}^{\gamma,\beta}}(\mathfrak{F}_{\mathfrak{M}}) = (\mathfrak{S}, \prod_{i=1}^{\mathfrak{M}} \mathfrak{f}_i \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\beta=0}^4)$$

$$\sum_{\beta=1}^{\eta(\gamma)} \mathfrak{d}_{\hat{\eta}}(\gamma, \xi) \phi^{\alpha_1(\gamma, \xi), \eta}(\mathfrak{G}_1) \cdots \phi^{\alpha_{\hat{\eta}-1}(\gamma, \xi), \eta}(\mathfrak{G}_{\hat{\eta}-1}) \phi^{\alpha_{\hat{\eta}}(\gamma, \xi), \eta}(\mathfrak{G}_{\hat{\eta}}) 1 = \left(\mathfrak{S}, \prod_{i=1}^{\hat{\eta}} \mathfrak{g}_i \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\alpha=0}^4 \right)$$

Lo anterior, computacionalmente equivale a:

$$\left| (\mathfrak{S}, \mathfrak{f} \otimes \rho(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\alpha=0}^4) - \left(\mathfrak{S}, \prod_{i=1}^{\hat{\eta}} \mathfrak{g}_i \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\alpha=0}^4 \right) \right| < \epsilon$$

Cuya métrica de Minkowski, se define así:

$$\begin{aligned} \mathcal{T}(\mathfrak{f}) &= \langle \phi^{\alpha, \eta}(\mathfrak{f}) (\hat{\mathfrak{S}}, \hat{\mathfrak{g}}_{\gamma} \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\hat{\mathfrak{f}}_{\alpha}\}_{\alpha=0}^4), (\mathfrak{S}, \mathfrak{g}_{\beta} \otimes \rho_{\eta}(\mathfrak{E}^{\beta}), \{\hat{\mathfrak{f}}_{\alpha}\}_{\alpha=0}^4) \rangle \\ &= \mathfrak{E}_{\alpha}^{\gamma, \beta} \int_{\mathfrak{S}^4} \mathfrak{d}\hat{\mathfrak{t}} (\mathfrak{f}^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}} \times \hat{\mathfrak{g}}_{\gamma} \cdot \overline{\mathfrak{g}}_{\beta}) (\sigma(\hat{\mathfrak{t}})) \cdot |\rho_{\sigma}|(\hat{\mathfrak{t}}) \end{aligned}$$

$$\begin{aligned} \vec{\chi} \rightarrow \mathfrak{f}^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\vec{\chi}) &\equiv \mathfrak{f}^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\mathcal{H}(\rho_{\eta}), \mathfrak{P}(\rho_{\eta}))(\vec{\chi}) \\ &= \int_{\hat{\mathfrak{S}} \in \mathbb{R}^4} \frac{e^{-i(\sigma(\hat{\mathfrak{S}}) \cdot (\mathcal{H}(\rho_{\eta}) \mathfrak{f}_0 + \mathfrak{P}(\rho_{\eta}) \mathfrak{f}_1))}}{2\pi} \mathfrak{f}(\vec{\chi} + \hat{\sigma}(\hat{\mathfrak{S}})) |\rho_{\hat{\sigma}}|(\hat{\mathfrak{S}}) \mathfrak{d}\hat{\mathfrak{S}} \end{aligned}$$

$$\begin{aligned} \mathcal{T}(\mathfrak{f}) &= \int_{\mathfrak{S}^4} \mathfrak{f}^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\sigma(\hat{\mathfrak{t}})) \mathfrak{h}(\hat{\mathfrak{t}}) \mathfrak{d}\hat{\mathfrak{t}} \\ &= \int_{\hat{\mathfrak{S}} \in \mathbb{R}^4, \mathfrak{t} \in \mathfrak{S}^4} \mathfrak{d}\hat{\mathfrak{t}} \mathfrak{d}\hat{\mathfrak{S}} \mathfrak{f} \left(\sigma(\hat{\mathfrak{t}}) + \hat{\sigma}(\hat{\mathfrak{S}}) |\rho_{\hat{\sigma}}|(\hat{\mathfrak{S}}) |\rho_{\sigma}|(\hat{\mathfrak{t}}) \cdot \frac{e^{-i(\hat{\sigma}(\hat{\mathfrak{S}}) \cdot \vec{\alpha})}}{2\pi} |\hat{\mathfrak{g}}_{\gamma} \cdot \overline{\mathfrak{g}}_{\beta}| \circ \sigma(\hat{\mathfrak{t}}) \cdot \mathfrak{E}_{\alpha}^{\gamma, \beta} \right) \end{aligned}$$

En este punto, cabe aplicar la ley de transformación del operador cuántico, que se expresa así:

$$\begin{aligned}
& \mathfrak{U}(\vec{\alpha}, \Lambda) \phi^{\alpha, \eta}(\mathfrak{f}) \mathfrak{U}(\vec{\alpha}, \Lambda)^{-1} \left(\mathfrak{S}, \mathfrak{g}_\beta \otimes \rho_\eta(\mathfrak{E}^\beta), \{\hat{\mathfrak{f}}_\alpha\}_{\alpha=0}^4 \right) \\
&= \mathfrak{U}(\vec{\alpha}, \Lambda) \phi^{\alpha, \eta}(\mathfrak{f}) \left(\Lambda^{-1}(\mathfrak{S} - \vec{\alpha}), \mathcal{T}(\rho_\eta, \vec{\alpha})^{-1} \mathfrak{g}_\beta(\Lambda \cdot + \vec{\alpha}) \otimes \rho_\eta(\mathfrak{E}^\beta), \mathcal{D} \right) \\
&= \mathfrak{U}(\vec{\alpha}, \Lambda) \left(\Lambda^{-1}(\mathfrak{S} - \vec{\alpha}), \left(\mathcal{T}(\rho_\eta, \vec{\alpha})^{-1} \mathfrak{f}^\mathfrak{C} \right) (\cdot) \mathfrak{d}_\gamma^\alpha \right. \\
&\quad \cdot \mathfrak{g}_\beta(\Lambda \cdot + \vec{\alpha}) \otimes \text{ad} \left(\rho_\eta(\mathfrak{F}^\gamma) \right) \rho_\eta(\mathfrak{E}^\beta), \mathcal{D} \left. \right) \\
&= \left(\mathfrak{S}, \mathfrak{I}(\rho_\eta, \vec{\alpha}) \mathcal{T}(\rho_\eta, \vec{\alpha})^{-1} \mathfrak{f}^\mathfrak{C}(\Lambda^{-1}(\cdot - \vec{\alpha})) \mathfrak{d}_\gamma^\alpha \mathfrak{g}_\beta(\cdot) \otimes \text{ad} \left(\rho_\eta(\mathfrak{F}^\gamma) \right) \rho_\eta(\mathfrak{E}^\beta), \{\hat{\mathfrak{f}}_\alpha\}_{\alpha=0}^4, \mathcal{D} \right) \\
& \mathfrak{f}^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{P}}(\rho_\eta))(\Lambda^{-1}(\vec{\chi} - \vec{\alpha})) = \int_{\hat{\mathfrak{s}} \in \mathbb{R}^4} \frac{e^{-i(\hat{\sigma}(\hat{\mathfrak{s}}) \cdot \widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{g}}_0 + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{g}}_1))}}{2\pi} \mathfrak{f}(\vec{\mathfrak{Y}} + \hat{\sigma}(\hat{\mathfrak{s}}) |\rho_\sigma| ((\widehat{\mathfrak{S}}) \mathfrak{d}\hat{\mathfrak{s}})) \\
&= \int_{\hat{\mathfrak{s}} \in \mathbb{R}^4} \frac{e^{-i(\sigma(\hat{\mathfrak{s}}) \cdot \widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{f}(\vec{\mathfrak{Y}} + \Lambda^{-1} \sigma(\hat{\mathfrak{s}})) |\rho_\sigma| ((\widehat{\mathfrak{S}}) \mathfrak{d}\hat{\mathfrak{s}}) \\
&= \int_{\hat{\mathfrak{s}} \in \mathbb{R}^4} \frac{e^{-i(\sigma(\hat{\mathfrak{s}}) \cdot \widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{f}(\Lambda^{-1}(\vec{\mathfrak{Y}} + \sigma(\hat{\mathfrak{s}}) - \vec{\alpha})) |\rho_\sigma| ((\widehat{\mathfrak{S}}) \mathfrak{d}\hat{\mathfrak{s}}) \\
&= \mathfrak{f}(\Lambda^{-1}(\cdot - \vec{\alpha}))^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{P}}(\rho_\eta))(\vec{\mathfrak{Y}}) \\
& \left(\mathfrak{S}, \mathfrak{f}(\Lambda^{-1}(\cdot - \vec{\alpha}))^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}} \mathfrak{U}(\Lambda^{-1})^\alpha \mathfrak{U}(\widehat{\Lambda})^\gamma \cdot \mathfrak{g}_\beta \otimes \text{ad} \left(\rho_\eta(\mathfrak{F}^\delta) \right) \rho_\eta(\mathfrak{E}^\beta), \{\widehat{\Lambda} \epsilon_\alpha\}_{\alpha=0}^4 \right) \\
&= \mathfrak{U}(\Lambda^{-1})^\alpha \phi^{\gamma, \eta} \left(\mathfrak{f}(\Lambda^{-1}(\cdot - \vec{\alpha})) \right) \left(\mathfrak{S}, \mathfrak{g}_\beta \otimes \rho_\eta(\mathfrak{E}^\beta), \{\widehat{\Lambda} \epsilon_\alpha\}_{\alpha=0}^4 \right)
\end{aligned}$$

Cuya simetría CPT, es igual a:

$$\begin{aligned}
& \left[\phi^{\alpha, \eta}(\mathfrak{f}), \phi^{\beta, \eta}(\mathfrak{g})^* \right]_{\pm} \left(\mathfrak{S}, \mathfrak{h}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \\
&= -\mathfrak{U}(\Lambda)^\alpha \overline{\mathfrak{U}(\Lambda)^\beta} \left(\mathfrak{S}, \mathfrak{B}^\pm(\mathfrak{f}^\mathfrak{C} \cdot \widehat{\mathfrak{g}}^\mathfrak{C} \pm \mathfrak{g}^\mathfrak{C} \cdot \widehat{\mathfrak{f}}^\mathfrak{C}) \right. \\
&\quad \cdot \mathfrak{h}_\gamma \otimes \text{ad}(\rho_\eta(\mathfrak{F}^\delta)) \text{ad}(\rho_\eta(\mathfrak{F}^{\mu\nu})) (\rho_\eta(\mathfrak{E}^\gamma)) \\
&\quad + \langle (\mathfrak{S}, \mathfrak{h}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)), \phi^{\beta, \eta}(\mathfrak{G})^\circ 1 \rangle \phi^{\alpha, \eta}(\mathfrak{F})^* 1 \\
&\quad \pm \langle (\mathfrak{S}, \mathfrak{h}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)), \phi^{\beta, \eta}(\mathfrak{F})^* 1 \rangle \phi^{\alpha, \eta}(\mathfrak{G})^\circ 1
\end{aligned}$$



$$\begin{aligned}
& \phi^{\alpha,\eta}(\mathfrak{F})^* \phi^{\beta,\eta}(\mathfrak{G})^* \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \\
&= \phi^{\alpha,\eta}(\mathfrak{F})^* \overline{\mathfrak{A}(\Lambda)_\mu^\beta}(\mathfrak{S}, -\widehat{\mathfrak{g}}^\mathfrak{C} \cdot \widehat{\mathfrak{h}}_\gamma \otimes \text{ad}(\rho_\eta(\mathfrak{F}^{\mu\nu})))(\rho_\eta(\mathfrak{E}^\gamma)) \\
&+ \langle \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta,\eta}(\mathfrak{G})^\circ 1 \rangle 1) \\
&= -\mathfrak{A}(\Lambda)_\delta^\alpha \overline{\mathfrak{A}(\Lambda)_\mu^\beta}(\mathfrak{S}, \widehat{\mathfrak{f}}^\mathfrak{C} \cdot \widehat{\mathfrak{g}}^\mathfrak{C} \cdot \widehat{\mathfrak{h}}_\gamma \otimes \text{ad}(\rho_\eta(\mathfrak{F}^\delta)) \text{ad}(\rho_\eta(\mathfrak{F}^{\mu\nu})))(\rho_\eta(\mathfrak{E}^\gamma)) \\
&+ \langle \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta,\eta}(\mathfrak{G})^\circ 1 \rangle \phi^{\alpha,\eta}(\mathfrak{F})^* 1
\end{aligned}$$

$$\mathfrak{g}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}(\widehat{\mathcal{H}}, \widehat{\mathfrak{F}})(\vec{\mathfrak{x}}) = \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^{\infty} \frac{e^{-i(\widehat{\eta}(\widehat{\mathfrak{s}}) \cdot (\widehat{\mathcal{H}}\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}\widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{g}(\vec{\mathfrak{x}} + \vec{\eta}(\widehat{\mathfrak{s}})) |\rho_{\widehat{\eta}}|(\widehat{\mathfrak{s}}) d\widehat{\mathfrak{s}}$$

$$\overline{\mathfrak{f}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}}(\widehat{\mathcal{H}}, \widehat{\mathfrak{F}})(\vec{\mathfrak{x}}) = \int_{\widehat{\mathfrak{t}} \in \mathbb{R}^4}^{\infty} \frac{e^{-i(\widehat{\eta}(\widehat{\mathfrak{t}}) \cdot (\widehat{\mathcal{H}}\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}\widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{g}(\vec{\mathfrak{x}} + \vec{\eta}(\widehat{\mathfrak{t}})) |\rho_{\widehat{\eta}}|(\widehat{\mathfrak{t}}) d\widehat{\mathfrak{t}}$$

$$\left[\mathfrak{g}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}} \cdot \overline{\mathfrak{f}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}} \right] (\widehat{\mathcal{H}}, \widehat{\mathfrak{F}})(\vec{\mathfrak{x}}) = \int_{\widehat{\mathfrak{s}}, \widehat{\mathfrak{t}} \in \mathbb{R}^4}^{\infty} \frac{e^{-i(\widehat{\eta}(\widehat{\mathfrak{s}}) \cdot (\widehat{\mathcal{H}}\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} \mathfrak{g}\vec{\mathfrak{x}}(\vec{\eta}(\widehat{\mathfrak{s}}) + \vec{\eta}(\widehat{\mathfrak{t}})) \widehat{\mathfrak{f}}_{\vec{\mathfrak{x}}}(\vec{\eta}(\widehat{\mathfrak{t}})) |\rho_{\widehat{\eta}}|(\widehat{\mathfrak{s}}) |\rho_{\widehat{\eta}}|(\widehat{\mathfrak{t}}) d\widehat{\mathfrak{s}} d\widehat{\mathfrak{t}}$$

$$= \int_{\widehat{\mathfrak{t}} \in \mathbb{R}^4, \widehat{\mathfrak{s}} \in \mathcal{D}}^{\infty} \frac{e^{-i(\widehat{\eta}(\widehat{\mathfrak{s}}) \cdot (\widehat{\mathcal{H}}\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} \mathfrak{g}\vec{\mathfrak{x}}(\vec{\eta}(\widehat{\mathfrak{s}}) + \vec{\eta}(\widehat{\mathfrak{t}})) \widehat{\mathfrak{f}}_{\vec{\mathfrak{x}}}(\vec{\eta}(\widehat{\mathfrak{t}})) |\rho_{\widehat{\eta}}|(\widehat{\mathfrak{s}}) |\rho_{\widehat{\eta}}|(\widehat{\mathfrak{t}}) d\widehat{\mathfrak{s}} d\widehat{\mathfrak{t}}$$

$$\left[-\frac{e^{-i((\mu-\nu) \cdot (\widehat{\mathcal{H}}(\rho_\eta)\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}(\rho_\eta)\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} + \frac{e^{-i((\nu-\mu) \cdot (\widehat{\mathcal{H}}(\rho_\eta)\widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{F}}(\rho_\eta)\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} \right] \widehat{\mathfrak{f}}_{\vec{\mathfrak{x}}}(\mu) \overline{\mathfrak{g}}_{\vec{\mathfrak{x}}}(\mu)$$

En un mapa bilineal, tenemos:

$$\begin{aligned}
& (\mathfrak{f}, \mathfrak{g}) \in \mathcal{P} \times \mathfrak{F} \rightarrow \langle \phi^{\alpha,\eta}(\mathfrak{f}) \phi^{\beta,\eta}(\mathfrak{g})^* \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \rangle \\
& - \langle \phi^{\alpha,\eta}(\mathfrak{f})^* \phi^{\beta,\eta}(\mathfrak{g}) \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta,\eta}(\mathfrak{g}) 1 \rangle \langle \phi^{\alpha,\eta}(\mathfrak{f}) 1 \left(\mathfrak{S}, \widehat{\mathfrak{h}}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \rangle
\end{aligned}$$



$$\begin{aligned}
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \widehat{\mathcal{W}}(\vec{x}, \vec{y}) f(\vec{x}) \otimes_{\mathbb{R}} g(\vec{y}) d\vec{x} d\vec{y} \\
&= \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \Re \widehat{\mathcal{W}}(\vec{x}, \vec{y}) [f(\vec{x}) \tilde{g}(\vec{y}) + \bar{f}(\vec{x}) \bar{g}(\vec{y})] d\vec{x} d\vec{y} \\
&+ \Im \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \Im \widehat{\mathcal{W}}(\vec{x}, \vec{y}) [f(\vec{x}) \tilde{g}(\vec{y}) - \bar{f}(\vec{x}) \bar{g}(\vec{y})] d\vec{x} d\vec{y} \\
&= \langle \phi^{\alpha, \eta}(f) \phi^{\beta, \eta}(g)^* (\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)), (\mathfrak{S}, \widehat{\hbar}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)) \rangle \\
&- \langle \phi^{\alpha, \eta}(f)^* \phi^{\beta, \eta}(g) (\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)), \phi^{\beta, \eta}(g) 1 \rangle \langle \phi^{\alpha, \eta}(f) 1 (\mathfrak{S}, \widehat{\hbar}_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)) \rangle \\
&\phi^{\alpha, \eta}(f) \phi^{\beta, \eta}(g)^* (\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)) - \langle (\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)) \phi^{\beta, \eta}(g) 1 \rangle \phi^{\alpha, \eta}(f) 1 \\
&= -\mathfrak{X}(\Lambda)_\delta^\alpha \overline{\mathfrak{X}(\Lambda)_\mu^\beta} (\mathfrak{S}, (f^{\{\mathfrak{f}_0, \mathfrak{f}_1\}} \cdot \overline{g^{\{\mathfrak{f}_0, \mathfrak{f}_1\}}}) \cdot \hbar_\gamma \otimes ad(\rho_\eta(\mathfrak{F}^\delta)) ad(\rho_\eta(\mathfrak{F}^{\mu\nu})) (\rho_\eta(\mathfrak{E}^\gamma))) \\
&\mathfrak{X}^\pm(\mathfrak{z}) \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4} \frac{e^{-i(\widehat{\eta}(\mathfrak{s}) - \widehat{\eta}(\widehat{\mathfrak{t}})) \cdot (\widehat{\mathcal{H}} \mathfrak{f}_0 + \widehat{\mathcal{P}} \mathfrak{f}_1)}}{(2\pi)^2} \widehat{f}(\widehat{\mathfrak{s}}) \overline{\widehat{g}(\widehat{\mathfrak{t}})} |\rho_{\widehat{\eta}}(\mathfrak{s})| |\rho_{\widehat{\eta}}(\widehat{\mathfrak{t}})| d\mathfrak{s} d\widehat{\mathfrak{t}} \\
&\pm \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4} \frac{e^{-i(\widehat{\eta}(\mathfrak{s}) - \widehat{\eta}(\widehat{\mathfrak{t}})) \cdot (\widehat{\mathcal{H}} \mathfrak{f}_0 + \widehat{\mathcal{P}} \mathfrak{f}_1)}}{(2\pi)^2} \widehat{g}(\widehat{\mathfrak{t}}) \overline{\widehat{f}(\mathfrak{s})} |\rho_{\widehat{\eta}}(\mathfrak{s})| |\rho_{\widehat{\eta}}(\widehat{\mathfrak{t}})| d\mathfrak{s} d\widehat{\mathfrak{t}}
\end{aligned}$$

Aplicando en este punto, la función beta de Callan-Symanzik, tenemos:

$$\beta(\mathfrak{C}) = \partial \mathfrak{C} / \partial |\Im \mathfrak{N} \widehat{\mathfrak{R}}_\eta|$$

En tanto que, del operador cuántico hamiltoniano, se obtiene:

$$\mathcal{T}r(\rho(\mathfrak{E}^\alpha)) (\rho(\mathfrak{E}^\beta)) = \mathfrak{C}(\rho) \mathcal{T}r(\mathfrak{E}^\alpha \mathfrak{E}^\beta)$$

$$\begin{aligned}
\mathcal{E}(\rho) &= \sum_{\alpha=1}^{\mathfrak{N}} (\rho(\mathfrak{E}^\alpha)) (\rho(\mathfrak{E}^\beta)) \\
&= 1/\kappa \sum_{\alpha=1}^{\mathfrak{N}} \kappa^2/4 \int_{\widehat{\mathfrak{s}} \in (-\delta, 1+\delta)^4} d\widehat{\mathfrak{s}} \sum_{i,j=1}^4 |\Im_{0ij}^\sigma| (d\widehat{\mathfrak{s}}) \kappa(\psi \cdot d_0 \mathfrak{X}_{ij, \alpha}) \left(\frac{\kappa \sigma(\widehat{\mathfrak{s}})}{2} \right) \rho(\mathfrak{E}^\alpha) \\
&\mathfrak{C} \sum_{\alpha=1}^{\eta} \int_{\widehat{\mathfrak{s}} \in \mathfrak{S}_\delta^4} d\widehat{\mathfrak{s}} \sum_{i,j=1}^4 |\Im_{0ij}^\sigma| (\widehat{\mathfrak{s}}) (d_0 \mathfrak{X}_{ij, \alpha}) (\sigma(\widehat{\mathfrak{s}})) \rho(\mathfrak{E}^\alpha) \\
\langle v_{\mathcal{R}(\alpha, \tau)}^{\kappa, \rho} \rangle^2 &= -\mathbb{E} \left(\cdot, v_{\mathcal{R}(\alpha, \tau)}^{\kappa, \rho} \right)^2 y^\kappa
\end{aligned}$$



$$\mathfrak{A}^\rho = \sum_{\alpha=1}^{\eta} \sum_{i,j=1}^4 \alpha_{i,\alpha} \otimes \delta x^i \otimes \rho(\mathfrak{E}^\alpha) \in \mathfrak{S}_{\mathbb{R}}(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^4) \otimes \rho(\mathfrak{g})$$

$$\frac{1}{3} = \int_{\{\mathfrak{d}\mathfrak{A} \in \mathfrak{S}_{\mathbb{R}}(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \rho(\mathfrak{g})\}} \exp\left(\mathfrak{C} \int_{\mathcal{R}(\alpha)} \mathfrak{d}\{\mathfrak{A}^\rho\}\right) e^{-1/2 \mathfrak{C}_{\mathfrak{ym}}(\mathfrak{A}) \mathfrak{D}|\mathfrak{d}\mathfrak{A}|} = \mathbb{E}_{\mathcal{Y}\mathcal{M}}^\kappa(\exp(\langle \cdot, \nu_{\mathcal{R}(\alpha)}^{\kappa,\rho} \rangle))$$

$$\mathfrak{Z} = \int_{\{\mathfrak{d}\mathfrak{A} \in \mathfrak{S}_{\mathbb{R}}(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \rho(\mathfrak{g})\}} e^{-1/2 \mathfrak{C}_{\mathfrak{ym}}(\mathfrak{A}) \mathfrak{D}|\mathfrak{d}\mathfrak{A}|}$$

$$-\mathbb{E}(\langle \cdot, \nu_{\mathcal{R}(\alpha)}^{\kappa,\rho_\eta} \rangle) \langle \cdot, \nu_{\mathcal{R}(\delta(\alpha))}^{\kappa,\rho_\eta} \rangle \mathcal{Y}^\kappa = \frac{|\alpha|}{4 \otimes \mathfrak{E}(\rho_\eta)} - \epsilon(\eta, \kappa)$$

$$\bar{\mathfrak{C}}/\kappa^4 \mathcal{C}(\rho_\eta) \leq \mathcal{T}\mathfrak{r} \in (\eta, \kappa) \leq \bar{\mathfrak{C}}/\kappa^4 \mathcal{C}(\rho_\eta)$$

Más, aplicando la ecuación de Callan-Symanzik, tenemos:

$$-\frac{1}{\mathfrak{E}(\rho_\eta) \mathbb{E}(\langle \cdot, \nu_{\mathcal{R}(\alpha)}^{\kappa,\rho_\eta} \rangle) \langle \cdot, \nu_{\mathcal{R}(\delta(\alpha))}^{\kappa,\rho_\eta} \rangle \mathcal{Y}^\kappa} = \frac{|\alpha|}{4 \otimes \mathfrak{E}_4(\rho_\eta)} - 1/\mathcal{C}(\rho_\eta) \epsilon(\eta, \kappa)$$

$$\frac{\widehat{\mathcal{N}}_\eta \mathfrak{E}_4(\rho_\eta)}{\mathfrak{E}(\rho_\eta)} - \frac{1}{\mathfrak{E}(\rho_\eta) \mathcal{T}\mathfrak{r}} \in (\eta, \kappa) \equiv \mathbb{N} - 1/\mathfrak{E}(\rho_\eta) \mathcal{T}\mathfrak{r} \in (\eta, \kappa)$$

$$\left\{ e \frac{\partial}{\partial e} + \beta(\mathfrak{C}) \frac{\partial}{\partial \mathfrak{C}} + 2\gamma(\mathfrak{C}) \right\} \mathfrak{G}_\eta^{(4)}(\mathfrak{C}, \epsilon) = 0$$

$$\mathfrak{G}_\eta^{(4)}(\mathfrak{C}, \epsilon) = \frac{\widehat{\mathcal{N}}_\eta}{\epsilon} - c^4 \hat{\lambda} + \hat{f}(c^5) = \frac{\partial}{\partial e} \mathfrak{G}_\eta^{(4)}(\mathfrak{C}, \epsilon) = -\frac{\widehat{\mathcal{N}}_\eta}{\epsilon}, \frac{\partial}{\partial \mathfrak{C}} \mathfrak{G}_\eta^{(4)}(\mathfrak{C}, \epsilon) = -4c^3 \hat{\lambda} + \hat{f}(c^4)$$

$$-\frac{\widehat{\mathcal{N}}_\eta}{\epsilon} - 4\beta(\mathfrak{C}) c^3 \hat{\lambda} + 2\gamma(\mathfrak{C}) \mathfrak{G}_\eta^{(4)}(\mathfrak{C}, \epsilon) + \beta(\mathfrak{C}) \hat{f}(c^4) = 0$$

$$-\frac{\mathfrak{C}}{4} \hat{f}(c^4) + \hat{f}(c^5) - 4c^3 \lambda(\mathfrak{C}) \hat{\lambda} + \lambda(\mathfrak{C}) \hat{f}(c^4) = 0$$

$$\frac{1}{c^4 |\hat{f}(c^4)|} + \frac{1}{c^5} |\hat{f}(c^5)| \leq \mathfrak{C}^4$$

$$\hat{\lambda}(\mathfrak{C}) = \frac{1}{4c^3 \hat{\lambda}} + \hat{f}(c^4) \left(\frac{\mathfrak{C}}{4\hat{f}(c^4)} - \hat{f}(c^5) \right)$$



Por tanto, la brecha de masa se vuelve positiva y por ende, superior a cero (estado de vacío), cuando:

$$\begin{aligned} \langle \mathfrak{C}_4(\rho)v, v \rangle &= \sum_{\alpha=1}^N \langle \rho(\mathfrak{C}^\alpha)v, \rho(\mathfrak{C}^\alpha)v \rangle \geq \sum_{\alpha=1}^{\mathcal{L}} \left| \sum_{\beta=1}^{\mathcal{L}} \alpha_{\alpha,\beta} \lambda_\rho(\mathcal{H}_\beta) \right|_4^4 \\ &= \sum_{\alpha=1}^{\mathcal{L}} \sum_{\beta=1}^{\mathcal{L}} \sum_{\gamma=1}^{\mathcal{L}} \lambda_\rho(\mathcal{H}_\beta) \alpha_{\alpha,\beta} \alpha_{\alpha,\gamma} \lambda_\rho(\mathcal{H}_\gamma) \geq \mathfrak{C} |\lambda_\rho|_4^4 \end{aligned}$$

$$\widehat{\mathcal{H}}(\rho_\eta)^2 = \frac{\widehat{\mathfrak{N}}_\eta}{4} \mathfrak{C}_2(\rho_\eta) = \frac{\mathfrak{N}}{4} \mathfrak{C}(\rho_\eta) = 0$$

$$\frac{\partial c}{\partial(\mathbb{I}_\eta \widehat{\mathfrak{N}})} = -\frac{\mathfrak{C}}{4} + \lambda(\mathfrak{C}), |\lambda(\mathfrak{C})| \leq \widehat{\mathfrak{C}}_4 \widehat{\mathfrak{C}}_2 = \frac{\mathfrak{d}\mathfrak{C}}{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})} = -\frac{\mathfrak{C}}{4} + \lambda(\mathfrak{C}) \Rightarrow \frac{\mathfrak{d}\mathfrak{C}}{\mathfrak{C}} - 4\lambda(\mathfrak{C}) = -\frac{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})}{4}$$

$$= \frac{1}{c} \frac{\mathfrak{d}\mathfrak{C}}{1 + \mu(\mathfrak{C})} = -\frac{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})}{4} \Rightarrow \left(\frac{1}{c \sum_{\kappa=0}^{\infty} (-1)^\kappa \mu(\mathfrak{C})^\kappa} \right) \mathfrak{d}\mathfrak{C} = \frac{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})}{4}$$

$$\mathcal{T}r\mathbb{E} \left(\left(\cdot, v_{\mathcal{R}(\alpha)}^{\kappa, \rho_\eta} \right) \left(\cdot, v_{\mathcal{R}(\alpha)}^{\kappa, \rho_\eta} \right) y^\kappa \right) = \frac{\widehat{\mathfrak{N}}_\eta}{4} \mathfrak{C}_2(\rho_\eta) - \mathcal{T}r\epsilon(\eta, \kappa)$$

$$= \frac{4}{\mathfrak{N}\mathfrak{C}(\rho_\eta)} \mathcal{T}r\mathbb{E} \left(- \left(\cdot, v_{\mathcal{R}(\alpha)}^{\kappa, \rho_\eta} \right) \left(\cdot, v_{\mathcal{R}(\alpha)}^{\kappa, \rho_\eta} \right) y^\kappa \right) - 1 = -\frac{4\mathcal{T}r\epsilon(\eta, \kappa)}{\mathfrak{N}\mathfrak{C}(\rho_\eta)}$$

$$\begin{aligned} \mathcal{U}(\vec{\alpha}, 1) \left(\mathfrak{S}, \mathfrak{f}_\alpha \otimes \rho(\mathfrak{C}^\alpha), \{\mathfrak{f}_\alpha\}_{\alpha=0}^4 \right) &= e^{-i(\mathfrak{f}_0 \cdot \widehat{\mathcal{H}}(\vec{\alpha}, \rho) + \mathfrak{f}_1 \cdot \widehat{\mathfrak{P}}(\vec{\alpha}, \rho))} \left(\mathfrak{S} + \vec{\alpha}, \mathfrak{f}_\alpha(\cdot - \vec{\alpha}) \otimes \rho(\mathfrak{C}^\alpha), \{\mathfrak{f}_\alpha\}_{\alpha=0}^4 \right) \\ &= e^{i(\alpha_0 \widehat{\mathcal{H}}(\rho) - \alpha^1 \widehat{\mathfrak{P}}(\rho))} \left(\mathfrak{S} + \vec{\alpha}, \mathfrak{f}_\alpha(\cdot - \vec{\alpha}) \otimes \rho(\mathfrak{C}^\alpha), \{\mathfrak{f}_\alpha\}_{\alpha=0}^4 \right) \end{aligned}$$

Cuyas particiones corresponden a:

$$\begin{aligned} \int_{\mathfrak{Q}} \{\mathfrak{h}_\theta\}_{\theta=1}^{r+s} &= \prod_{l=1}^{\eta(\mathfrak{Q})} \left\{ \frac{\int_{\mathfrak{S}_0}^{\infty} \left(\prod_{\theta \in \mathfrak{A}_l} \int_{\gamma_\theta^- \in \mathbb{R}^4} e^{i\chi(\theta)(\gamma_\theta^0 \widehat{\mathcal{H}} - \gamma_\theta^1 \widehat{\mathfrak{P}})} \right)}{2\varpi} \mathfrak{h}_\theta(\gamma_\theta^-, \gamma^+) \mathfrak{d}\gamma_\theta^- \mathfrak{d}\gamma^+ \right\} \\ &= \prod_{l=1}^{\eta(\mathfrak{Q})} \left\{ \frac{\int_{\mathfrak{S}_0}^{\infty} \left(\prod_{\theta \in \mathfrak{A}_l} \int_{\gamma_\theta^0 \gamma_\theta^1 \in \mathbb{R}^4} e^{i\chi(\theta)(\gamma_\theta^0 \widehat{\mathcal{H}} - \gamma_\theta^1 \widehat{\mathfrak{P}})} \right)}{2\varpi} \mathfrak{h}_\theta(\gamma_\theta^0, \gamma_\theta^1, \gamma^2, \gamma^4) \mathfrak{d}\gamma_\theta^0 \mathfrak{d}\gamma_\theta^1 \mathfrak{d}\gamma^2 \mathfrak{d}\gamma^4 \right\} \end{aligned}$$



$$\begin{aligned}
& \langle \mathfrak{A}_\tau^\eta \mathfrak{B}_\delta^\eta 1, 4 \rangle - \langle \mathfrak{A}_\tau^\eta 1, 4 \rangle \langle \mathfrak{B}_\delta^\eta 1, 4 \rangle \equiv \langle \mathfrak{A}_\tau^\eta \mathfrak{P}_0 \mathfrak{B}_\delta^\eta 1, 4 \rangle \\
&= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathfrak{W}^\eta ((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta}) \otimes_{\theta=\tau+1}^{r+\delta} \rho_\tau(\vec{x}_\tau) \cdot \prod_{\theta=\tau+1}^{r+\delta} d\vec{x}_\theta \\
&= \mathfrak{C}_R \int_{\mathfrak{S}_0} \left(\prod_{\tau=1}^r \int_{\mathfrak{x}_\tau^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\tau^-) \mathfrak{h}_\tau(\mathfrak{x}_\tau^-, \mathfrak{x}^+) d\mathfrak{x}_\tau^- \cdot \prod_{\theta=\tau+1}^{r+\delta} \int_{\mathfrak{x}_\theta^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\theta^-) \mathfrak{h}_\theta(\mathfrak{x}_\theta^-, \mathfrak{x}^+) d\mathfrak{x}_\theta^- \right) d\mathfrak{x}^+ \\
&+ \sum_{\substack{\mathbb{Q} \neq \mathfrak{R} \\ \mathbb{Q} \in \Gamma}} \mathfrak{C}_\Omega \int_{\Omega} \{\mathfrak{H}_\theta\}_{\theta=1}^{r+\delta}
\end{aligned}$$

$$\begin{aligned}
\langle \mathfrak{P}_0 \mathfrak{U}(\vec{\alpha}, 1) \mathfrak{B}_\delta^\eta 1 \rangle &= \mathfrak{P}_0 \psi^{\beta_1, \eta}(\mathfrak{g}_1)_{\mathfrak{U}(\vec{\alpha})} \psi^{\beta_2, \eta}(\mathfrak{g}_2)_{\mathfrak{U}(\vec{\alpha})} \dots \psi^{\beta_{\delta-1}, \eta}(\mathfrak{g}_{\delta-1})_{\mathfrak{U}(\vec{\alpha})} \mathfrak{U}(\vec{\alpha}, 1) \psi^{\beta_\delta, \eta}(\mathfrak{g}_\delta) 1 \\
&= e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} \mathfrak{P}_0 \mathfrak{B}_\delta^{\eta, \vec{\alpha}} 1 = e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} (\mathfrak{B}_\delta^{\eta, \vec{\alpha}} 1 - \langle \mathfrak{B}_\delta^\eta 1, 1 \rangle 1)
\end{aligned}$$

$$\langle \mathfrak{A}_\tau^\eta \mathfrak{P}_0 \mathfrak{U}(\vec{\alpha}, 1) \mathfrak{B}_\delta^\eta 1, 1 \rangle = e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} (\langle \mathfrak{A}_\tau^\eta \mathfrak{B}_\delta^{\eta, \vec{\alpha}} 1, 1 \rangle - \langle \mathfrak{A}_\tau^\eta 1, 1 \rangle \langle \mathfrak{B}_\delta^\eta 1, 1 \rangle)$$

$$\begin{aligned}
e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} &= \int_{\mathbb{R}^4} \frac{e^{i(s\hat{\mathcal{H}} - t\hat{\rho})}}{2\omega} f(s, t, \mathfrak{x}^2, \mathfrak{x}^4) ds dt = \int_{\mathbb{R}^4} \frac{e^{i((s+\alpha^0)\hat{\mathcal{H}} - (t+\alpha^1)\hat{\rho})}}{2\omega} f(s, t, \mathfrak{x}^2, \mathfrak{x}^4) ds dt \\
&= \int_{\mathbb{R}^4} \frac{e^{i(s\hat{\mathcal{H}} - t\hat{\rho})}}{2\omega} f(s - \alpha^0, t - \alpha^1, \mathfrak{x}^2, \mathfrak{x}^4) ds dt \\
&= f(\cdot - (\alpha^0, \alpha^1, 0, 0))^{\{\epsilon_0 \epsilon_1\}} (\hat{\mathcal{H}}, \hat{\rho})(0, 0, \mathfrak{x}^2, \mathfrak{x}^4)
\end{aligned}$$

$$\begin{aligned}
\mathcal{H}^\eta(\vec{\alpha}) &= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathfrak{W}^\eta ((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta}) \otimes_{\tau=1}^r \rho_\tau(\vec{x}_\tau) \cdot \otimes_{\theta=\tau+1}^{r+\delta} \rho_0(\vec{x}_{\tau+\theta} - \vec{\alpha}) \cdot \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau \\
&\equiv \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathfrak{W}^\eta ((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta}) \otimes_{\tau=1}^{r+\delta} \rho_\tau(\vec{x}_\tau) \cdot \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau \\
&= \mathfrak{C}_R \int_{\mathfrak{S}_0} \left(\prod_{\tau=1}^r \int_{\mathfrak{x}_\tau^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\tau^-) \mathfrak{h}_\tau(\mathfrak{x}_\tau^-, \mathfrak{x}^+) d\mathfrak{x}_\tau^- \cdot \prod_{\theta=\tau+1}^{r+\delta} \int_{\mathfrak{x}_\theta^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\theta^-) \mathfrak{h}_\theta(\mathfrak{x}_\theta^-, \mathfrak{x}^+) d\mathfrak{x}_\theta^- \right) d\mathfrak{x}^+ \\
&+ \sum_{\substack{\mathbb{Q} \neq \mathfrak{R} \\ \mathbb{Q} \in \Gamma}} \mathfrak{C}_\Omega \int_{\Omega} \{\mathfrak{H}_\theta\}_{\theta=1}^{r+\delta} \\
&= \mathcal{H}_4^\eta(\vec{\alpha}) = \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau \mathfrak{W}^\eta((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta + \vec{\alpha})_{\theta=\tau+1}^{r+\delta}) \varphi_1((\vec{x}_\tau)_{\tau=1}^r) \varphi_2((\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta})
\end{aligned}$$



$$\mathfrak{W}^\eta((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta + \vec{\alpha})_{\theta=\tau+1}^{r+\delta}) = e^{-i\vec{\alpha} \cdot \mathfrak{C}_\tau m_\eta \vec{\eta}_0^\eta} \prod_{\theta=1}^{\tau+\delta} \mathfrak{C}(\vec{x}_\theta^-) \cdot \mathfrak{W}_0^\eta((\vec{x}_\tau^+)_{\tau=1}^r, (\vec{x}_\theta^+ + \alpha^+)_{\theta=\tau+1}^{r+\delta})$$

$$\mathfrak{W}_0^\eta((\vec{x}_\tau^+)_{\tau=1}^r, (\vec{x}_\theta^+ + \vec{\alpha})_{\theta=\tau+1}^{r+\delta}) = \mathfrak{W}^\eta((0^-, \vec{x}_\tau^+)_{\tau=1}^r, (0^-, \vec{x}_\theta^+ + \alpha^+)_{\theta=\tau+1}^{r+\delta})$$

$$\hat{\varphi}_1((\mathfrak{q}_\tau^-, \vec{x}_\tau^+)_{\tau=1}^r) = \frac{\int_{\mathbb{R}^{4\tau}} \prod_{\tau=1}^r e^{i\chi(\tau) \vec{x}_\tau^- \cdot \mathfrak{q}_\tau^-}}{2\omega} \cdot \varphi_1((\vec{x}_\tau^-, \vec{x}_\tau^+)_{\tau=1}^r) \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau^-, \hat{\varphi}_2((\mathfrak{q}_\theta^-, \vec{x}_\theta^+)_{\theta=\tau+1}^{r+\delta})$$

$$= \frac{\int_{\mathbb{R}^{4\tau}} \prod_{\tau=1}^r e^{i\chi(\theta) \vec{x}_\theta^- \cdot \mathfrak{q}_\theta^-}}{2\omega} \cdot \varphi_2((\vec{x}_\theta^-, \vec{x}_\theta^+)_{\theta=\tau+1}^{r+\delta}) \prod_{\theta=\tau+1}^{r+\delta} d\vec{x}_\theta^-$$

$$e^{i\mathfrak{C}_\tau m_\eta \vec{\alpha}_1^\eta} \int_{\mathbb{R}^{4(\tau+\delta)}} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau^+ \mathfrak{W}_0^\eta((\vec{x}_\tau^+)_{\tau=1}^r, (\vec{x}_\theta^+ + \vec{\alpha})_{\theta=\tau+1}^{r+\delta}) \hat{\varphi}(\mathcal{H}_\eta^-, \vec{x}_\tau^+)_{\tau=1}^{r+\delta}$$

$$|\hat{\varphi}((\mathfrak{q}_\theta^-, \vec{x}_\theta^+)_{\theta=\tau+1}^{r+\delta})| \leq \frac{\mathfrak{C}(\rho_\eta)^{\hat{\eta}} \|\varphi\|_{\hat{\eta}}}{\sum_{\theta=\tau+1}^{r+\delta} (|\mathfrak{q}_\theta^0|^2 + |\mathfrak{q}_\theta^1|^2)^{\frac{\ell}{2}}} + (|\vec{x}_\theta^2|^2 + |\vec{x}_\theta^4|^2)^{\frac{\kappa}{2}}$$

$$e^{i\mathfrak{C}_\tau m_\eta \vec{\alpha}_1^\eta} \sum_{|\vec{M}| \leq N} \int_{\mathbb{R}^{4(\tau+\delta)}} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau^+ \mathfrak{D}^{\vec{M}} \mathfrak{R}_{\vec{M}}^\eta((\vec{x}_\tau^+)_{\tau=1}^{r+\delta}; \alpha^+) \hat{\varphi}(\mathcal{H}_\eta^-, \vec{x}_\tau^+)_{\tau=1}^{r+\delta} = \mathcal{H}_4^\eta(\vec{\alpha})$$

$$\sum_{|\vec{M}| \leq N} |\mathfrak{D}^{\vec{M}}|((\vec{x}_\tau^+)_{\tau=1}^{r+\delta}; \alpha^+) \leq \mathfrak{C}(\rho_\eta)^{\hat{\kappa}} (|\alpha^+|^\alpha + \left(\sum_{\theta=\tau+1}^{r+\delta} |\vec{x}_\theta^+|^4 \right)^{\frac{\gamma}{2}})$$

$$\sum_{|\vec{M}| \leq N} \int_{\mathbb{R}^{4(\tau+\delta)}} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau^+ \mathfrak{D}^{\vec{M}} \mathfrak{R}_{\vec{M}}^\eta((\vec{x}_\tau^+)_{\tau=1}^{r+\delta}; \alpha^+) \hat{\varphi}(\mathcal{H}_\eta^-, \vec{x}_\tau^+)_{\tau=1}^{r+\delta}$$

$$= \sum_{|\vec{M}| \leq N} \int_{\{\mathcal{R} > \mathcal{R}_0 - \epsilon\}} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau^+ \mathfrak{D}^{\vec{M}} \mathfrak{R}_{\vec{M}}^\eta((\vec{x}_\tau^+)_{\tau=1}^{r+\delta}; \alpha^+) \hat{\varphi}(\mathcal{H}_\eta^-, \vec{x}_\tau^+)_{\tau=1}^{r+\delta} = \mathcal{H}_4^\eta(\vec{\alpha})$$

$$|\mathcal{H}_4^\eta(\vec{\alpha})| \leq \sum_{|\vec{M}| \leq N} \int_{\{\mathcal{R} > \mathcal{R}_0 - \epsilon\}} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau^+ \left| \mathfrak{R}_{\vec{M}}^\eta((\vec{x}_\tau^+)_{\tau=1}^{r+\delta}; \alpha^+) (\mathfrak{D}^{\vec{M}} \hat{\varphi}(\mathcal{H}_\eta^-, \vec{x}_\theta^+)_{\theta=1}^{r+\delta}) \right| -$$

$$\leq \mathfrak{C}(\rho_\eta)^{\hat{\kappa}} \int_{\{\mathcal{R} > \mathcal{R}_0 - \epsilon\}} (|\alpha^+|^\alpha + \mathcal{R}^\gamma) \frac{\mathfrak{C}(\rho_\eta)^{\hat{\eta}} \|\varphi\|_{\hat{\rho}\hat{\sigma}}}{(x+s)|\mathcal{M}_\kappa|^\ell + \mathcal{R}^{\kappa+\ell}} \mathcal{R}^{2(\tau+s)} / \mathfrak{R} d\mathcal{R} d\Omega$$

$$\sum_{\Omega \in \Omega_\tau} \mathfrak{C}_\Omega \int_{\mathbb{R}} d\xi \int_{\mathbb{R}^4} d\vec{x}_\tau^+ d\vec{x}_{\tau+}^+ \hat{\varphi}_{\Omega,1}^\eta(\vec{x}_\tau^+) \mathfrak{R}_\eta(\alpha^1 - \xi^1; \vec{x}_\tau^+, \vec{x}_{\tau+}^+ + \alpha^+) \hat{\varphi}_{\Omega,2}^\eta(\vec{x}_{\tau+}^+) g(\xi)$$



Más, aplicando la función de Green, tenemos:

$$\begin{aligned} &(-(\mathfrak{C}_{\mathcal{R}}\mathcal{M}_{\mathfrak{K}})^2 + \sum_{i=1}^4 \partial^2/\partial\alpha^{i,2})\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = \delta(\alpha^1 - \xi^1, \mathfrak{x}_{\tau}^+ - \mathfrak{x}_{\tau^+}^+ - \alpha^+) \\ &\equiv \delta(\alpha - \xi_{\mathcal{R}}) \end{aligned}$$

$$\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = \frac{1}{(2\pi)^{\frac{4}{2}} \int_{\mathbb{R}^4} \mathrm{d}q \mathbf{e}^{iq \cdot (\alpha - \xi_{\mathcal{R}})}} + |q|^4$$

$$\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = -\frac{\frac{1}{(2\pi)^{\frac{4}{2}}} 1}{i\mathcal{R} \int_0^{\infty} \lambda \mathbf{e}^{i\mathcal{R}\lambda}} - \frac{\mathbf{e}^{-i\mathcal{R}\lambda}}{\omega^2} + \lambda^2 \mathrm{d}\lambda$$

$$= -\frac{\frac{1}{(2\pi)^{\frac{4}{2}}} 1}{i\mathcal{R} \int_{-\infty}^{\infty} \lambda \mathbf{e}^{-i\mathcal{R}\lambda}} / (\lambda - i\omega)(\lambda + i\omega) \mathrm{d}\lambda$$

$$\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = \frac{\frac{2\omega i}{(2\pi)^{\frac{4}{2}}} 1}{i\mathcal{R}} \times \frac{i\omega \mathbf{e}^{-\mathcal{R}\omega}}{2i\omega} = \sqrt{2}\pi \mathbf{e}^{-\mathcal{R}|\mathfrak{C}_{\mathfrak{R}}\mathfrak{M}_{\eta}|} / 2\mathcal{R}$$

$$|\mathcal{H}_4^{\eta}(\vec{\alpha})g(\alpha^1)| = \left| \left(\frac{\partial^2}{\partial\alpha^{0,2}} - \widehat{\mathfrak{N}} \right) \Psi_{\eta}(\alpha^0, \alpha) \right|$$

$$= \left| \sum_{\Omega \in \Omega_{\mathcal{R}}} \mathfrak{C}_{\Omega} \int_{\mathbb{R}} \mathrm{d}\xi^1 \int_{\mathbb{R}^4} \mathrm{d}\vec{x} \hat{\varphi}_{\Omega,1}^{\eta}(\vec{x}^-) \mathfrak{R}_{\eta}(\alpha - \xi) \left((\mathfrak{C}_{\mathcal{R}}\mathcal{M}_{\eta})^2 + \widehat{\mathfrak{N}} \right) \hat{\varphi}_{\Omega,2}^{\eta}(\vec{x}^+) g(\xi^1) \right|$$

$$\leq \sum_{\Omega \in \Omega_{\mathcal{R}}} |\mathfrak{C}_{\Omega}| \left\{ \int_{\mathbb{R}} \mathrm{d}\xi^1 \left(|g(\xi^1)| + \frac{1}{\mathcal{M}_2^4 |g''(\xi^1)|} \right) \right.$$

$$\cdot \left. \frac{\sqrt{\frac{2}{\omega}} \mathbf{e}^{-\mathcal{M}_0|\alpha^+ + \mathfrak{x}^+ - \mathfrak{y}^+|}}{2|\alpha^+ + \mathfrak{x}^+ - \mathfrak{y}^+|} \times \int_{\mathbb{R}^4} \mathrm{d}\mathfrak{y}^+ |\hat{\varphi}_{\Omega,1}^{\eta}(\mathfrak{y}^+)| \cdot \int_{\mathbb{R}^4} \mathrm{d}\mathfrak{x}^+ \left| \left((\mathfrak{C}_{\mathcal{R}}\mathcal{M}_{\eta})^2 + \widehat{\mathfrak{N}} \right) \hat{\varphi}_{\Omega,2}^{\eta}(\mathfrak{x}^+) \right| \right\}$$

$$< \mathfrak{C}(\rho_{\eta})^{\widehat{\mathfrak{K}}} \mathbf{e}^{\mathcal{M}_0\epsilon} \|g\| \|p_4, q_4\| \varphi_1 \|p_1 q_1\| \varphi_4 \|p_4 q_4\| \cdot \frac{\mathbf{e}^{-\mathcal{M}_0|\alpha^+|}}{|\alpha^+|} - \epsilon$$



Por lo que, la integral de superficie, queda definida de la siguiente manera:

$$\begin{aligned} \rho_{\sigma}^{\alpha\beta} &= \frac{1}{\sqrt{\det(1 + \mathcal{W}_{\alpha\beta}^{cb,J} \mathcal{W}_{\alpha\beta}^{cb})}} \equiv \frac{|\mathfrak{S}_{\alpha\beta}^{\sigma}|}{\sqrt{\det(\mathfrak{S}_{\alpha\beta}^{\sigma,J} \mathfrak{S}_{\alpha\beta}^{\sigma} + \mathfrak{S}_{cb}^{\sigma,J} \mathfrak{S}_{cb}^{\sigma})}} = \int_{\mathcal{S}} \mathfrak{d}\rho \\ &= \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathbb{I}^2} \rho_{\sigma}^{\alpha\beta}(\mathfrak{s}, t) |\mathfrak{S}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, t) \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} (\dot{\sigma} + \dot{\sigma})(\dot{\sigma} - \dot{\sigma}) \mathfrak{d}\sigma \mathfrak{d}\hat{\sigma} \end{aligned}$$

Por otro lado, en este punto, es indispensable, añadir algunos planteamientos teóricos adicionales cuyo propósito es reforzar la tesis formulada en trabajos anteriores, siendo éstos, los que siguen a continuación:

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$$\begin{aligned} \frac{\mathfrak{d}^2 \chi^{\alpha}}{\mathfrak{d}\mathfrak{s}^2} + \frac{\Gamma_{\beta\gamma}^{\alpha} \mathfrak{d}\chi^{\beta}}{\mathfrak{d}\mathfrak{s}} \mathfrak{d}\chi^{\gamma} &= \Gamma_{\beta\gamma}^{\alpha} = \frac{1}{2g^{\alpha\beta} (g_{\mathfrak{d}\beta,c} g_{\mathfrak{d}c,\beta} g_{\beta c, \mathfrak{d}})} \int \mathfrak{d}\rho \sqrt{\frac{g^{\alpha\beta} \mathfrak{d}\chi^{\alpha} \mathfrak{d}\chi^{\beta}}{\mathfrak{d}\rho} \frac{\mathfrak{d}\chi^{\beta}}{\mathfrak{d}\rho}} \\ \mathcal{R}_{ij} &= -\frac{1}{2g_{ij}\mathfrak{R}} = -\frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} = \mathcal{F}(g_{ij}) = \int_{\mathcal{M}} \mathcal{R} \sqrt{-g} \mathfrak{d}\mathfrak{x} \\ \mathcal{R} &= \frac{4\pi\mathfrak{G}}{c^4} \lim_{\lambda \rightarrow 0} \frac{1}{\lambda} (\mathcal{F}(g_{ij} + \lambda \chi_{ij}) - \mathcal{F}(g_{ij})) = \delta \mathcal{F}((g_{ij}), \chi) \mathcal{R}_{ij} - \frac{1}{2g_{ij}\mathfrak{R}} \\ &= -\frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} - \mathcal{D}_i \mathcal{D}_j \varphi \operatorname{div} \left(\mathcal{D}_i \mathcal{D}_j \varphi + \frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} \right) \mathfrak{R} \left(\frac{16\pi\mathfrak{G}}{c^4} \right) \mathcal{T}_{ij} + \phi, \int_{\mathcal{M}} \phi \sqrt{-g} \mathfrak{d}\mathfrak{x} \\ \mathfrak{d}\mathfrak{s}^2 &= \epsilon^{\mu} c^4 \mathfrak{d}\mathfrak{t}^2 + \epsilon^{\nu} \mathfrak{d}\mathfrak{r}^2 + r^2 (\mathfrak{d}\theta^2 + \sin^2 \theta \mathfrak{d}\gamma^2) \\ \mathfrak{F} &= m\mathfrak{M}\mathfrak{G} \left(-\frac{1}{r^2} + \frac{1}{\delta \left(2 + \frac{\delta}{\tau} \right) \varphi'} + \frac{\mathcal{R}\tau}{\delta} \right) \\ \mathfrak{F} &= m\mathfrak{M}\mathfrak{G} \left(-\frac{1}{r^2} + \left(2 + \frac{\delta}{\tau} \right) \epsilon r^2 + \frac{\mathcal{R}\tau}{\delta} + \frac{1}{\delta \left(2 + \frac{\delta}{\tau} \right) r^2 \int r^{-2} \mathcal{R} \mathfrak{d}\mathfrak{r}} \right) \\ \mathfrak{F} &= m\mathfrak{M}\mathfrak{G} \left(-\frac{1}{r^2} - \frac{\kappa_0}{\tau} + \kappa_1 \mathcal{R} \right) \\ \mu &= \left\{ \mu_{i_1 \dots i_s}^{j_1 \dots j_r}(\mathfrak{x}) \mid 1 \leq i_1 \dots i_s, j_1 \dots j_r \leq \eta \right\} \\ \mathcal{L}^{\rho}(\mathfrak{G}) &= \left\{ \mu: \mathcal{M} \rightarrow \mathfrak{G} \mid \int_{\mathcal{M}} \|\mu\|^{\rho} \mathfrak{d}\mathfrak{x} < \infty \right\} \end{aligned}$$



$$\|\mu\|_{L^\rho} = \left[\int_{\mathcal{M}} \|\mu\|^\rho dx \right]^{1/\rho} = \left[\int_{\mathcal{M}} \sum |\mu_{i_1 \dots i_s}^{j_1 \dots j_r}|^\rho dx \right]^{1/\rho}$$

$$(\mu, \nu) = \int_{\mathcal{M}} g_{j_1 \kappa_1} \dots g_{j_r \kappa_r} g^{i_1 l_1} \dots g^{i_s l_s} \mu_{i_1 \dots i_s}^{j_1 \dots j_r} \nu_{l_1 \dots l_s}^{\kappa_1 \dots \kappa_r} \sqrt{-g} dx \Delta_\mu \Delta_\nu \nabla^\mu \nabla^\nu$$

$$\begin{aligned} \mu &= \{\mu_{i_1 \dots i_s}^{j_1 \dots j_r}\} \nabla_\mu = \{\mathcal{D}_\kappa \mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, \nabla_\mu: \mathfrak{M} \rightarrow \mathcal{T}_{s+1}^r \mathcal{M}, \nabla^* \mu \{g^{\kappa l} \mathcal{D}_l \mu\}: \mathfrak{M} \rightarrow \mathcal{T}_{s+1}^r \mathcal{M}, \operatorname{div} \mu = \{\mathcal{D}_i \mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, \mu \\ &= \{\mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, \operatorname{div} \mu = \{\mathcal{D}^l \mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, (\nabla^* \mu, \nu) = -(\mu \operatorname{div} \nu), (\nabla \mu, \nu) \\ &= -(\mu \operatorname{div} \nu), \lim_{\eta \rightarrow \infty} (\mathfrak{G}_{\mu_\eta}, \nu) = (\mathfrak{G}_{\mu_0}, \nu), (\mathfrak{G} \mu, \mu) \geq \alpha \|\mu\|^4 - \beta \end{aligned}$$

$$\begin{aligned} \mu &= \nabla_\varphi + \nu + \mathfrak{h}, \mathcal{H}(\mathcal{T}_s^r \mathcal{M}) = \{h \in \mathcal{L}^2(\mathcal{T}_s^r \mathcal{M}) | \nabla h, \operatorname{div} h\}, \mathcal{L}^2(\mathbb{E}) = \mathfrak{G}(\mathfrak{E}) \oplus \mathcal{L}_D^2(\mathfrak{E}), \mathcal{L}^2(\mathbb{E}) \\ &= \mathfrak{G}(\mathfrak{E}) \oplus \mathfrak{H}(\mathfrak{E}) \oplus \mathcal{L}_{\mathfrak{H}}^2(\mathfrak{E}), \mathfrak{G}(\mathfrak{E}) = \{\nu \in \mathcal{L}^2(\mathbb{E}) | \nu = \nabla \varphi, \varphi \in \mathcal{H}^1(\mathcal{T}_{s-1}^r \mathcal{M})\}, \mathcal{L}_D^2(\mathfrak{E}) \\ &= \{\nu \in \mathcal{L}^2(\mathbb{E}) | \operatorname{div} \nu = 0\}, \mathcal{L}_{\mathfrak{H}}^2(\mathfrak{E}) = \{\nu \in \mathcal{L}^2(\mathbb{E}) | \nabla_\nu \neq 0\}, \mathcal{L}_D^2(\mathfrak{E}) \perp \mathfrak{G}(\mathfrak{E}), \mathcal{L}_{\mathfrak{H}}^2(\mathfrak{E}) \\ &\perp \mathcal{H}(\mathfrak{E}), \mathfrak{G}(\mathfrak{E}) \perp \mathcal{H}(\mathfrak{E}), \mathfrak{E} = \mathfrak{E}_1 \oplus \mathfrak{E}^\kappa, \Delta \varphi = \operatorname{div} \mu, \Delta = \operatorname{div} \nabla, \nu = \mu - \nabla \varphi \\ &\in \mathcal{L}^2(\mathbb{E}), (\nu, \nabla \varphi) = 0, (\nabla \varphi - \mu, \nabla \psi) = 0, \mathcal{H} = \mathcal{H}^1(\widehat{\mathfrak{E}}) \setminus \widehat{\mathcal{H}}, \widehat{\mathcal{H}} \\ &= \{\psi \in \mathcal{H}^1(\widehat{\mathfrak{E}}) | \nabla \varphi = 0\}, (\mathfrak{G} \varphi, \psi) = (\nabla \varphi, \nabla \psi), (\mathfrak{G} \varphi, \varphi) = (\nabla \varphi, \nabla \varphi) = \|\varphi\|^4, \Delta \varphi \\ &= \mathfrak{f}, \mathcal{H}^\kappa(\mathfrak{E}) = \mathcal{H}_D^\kappa \oplus \mathfrak{G}^\kappa, \mathcal{L}^2(\mathbb{E}) = \mathcal{L}_D^2 \oplus \mathfrak{G}, \mathcal{H}_D^\kappa = \{\mu \in \mathcal{H}^\kappa(\mathbb{E}) | \operatorname{div} \mu = 0\}, \mathfrak{G}^\kappa \\ &= \{\mu \in \mathcal{H}^\kappa(\mathbb{E}) | \mu = \nabla \psi\}, \widehat{\Delta} \mu = \wp \Delta \mu, \Delta = \operatorname{div} \nabla = \mathcal{D}^\kappa \mathcal{D}_\kappa = \frac{g^{\kappa l} \partial^2}{\partial x^\kappa \partial x^l} + \mathfrak{B} \end{aligned}$$

$$\mathfrak{A} = \frac{g^{\kappa l} \partial^2}{\partial x^\kappa \partial x^l}: \mathcal{H}^2(\mathcal{M}, \mathcal{R}^\kappa) \rightarrow \mathcal{L}^2(\mathcal{M}, \mathcal{R}^\kappa), \Delta: \mathcal{H}^4(\mathfrak{E}) \rightarrow \mathcal{L}^2(\mathfrak{E}), \widehat{\Delta} = \mathcal{P} \Delta: \mathcal{H}_D^2(\mathfrak{E}) \mathcal{L}_D^2(\mathfrak{E}), \widehat{\mathcal{H}}$$

$$= \{\mu \in \mathcal{H}_D^2(\mathfrak{E}) | \widehat{\Delta} \mu = 0\}$$

$$\int_{\mathcal{M}} (\widehat{\Delta} \mu, \mu) \sqrt{-g} dx = \int_{\mathcal{M}} (\Delta \mu, \mu) \sqrt{-g} dx = - \int_{\mathcal{M}} (\nabla \mu, \nabla \mu) \sqrt{-g} dx = 0$$

$$\frac{\partial}{\partial x^i} (\mathcal{D}^\kappa \mu_{\kappa l}) - \frac{\partial}{\partial x^i} (\mathcal{D}^\kappa \mu_{\kappa j}) = \frac{\partial \Delta \varphi_i}{\partial x^i} - \frac{\partial \Delta \varphi_j}{\partial x^i}$$

$$\Delta \varphi_i = -(\delta \mathfrak{d} + \mathfrak{d} \delta) \varphi_i - \mathcal{R}_i^\kappa \varphi_\kappa$$

$$\widehat{\Delta} \varphi = - \frac{1}{\sqrt{-g} \frac{\partial}{\partial x^i} \left(\frac{\sqrt{-g} g^{ij} \partial \varphi}{\partial x^j} \right)}$$

$$(\delta \mathfrak{d} + \mathfrak{d} \delta) \varphi_i = \frac{\partial}{\partial x^i} \widehat{\Delta} \varphi \Leftrightarrow \varphi_i = \frac{\partial \varphi}{\partial x^i}$$



$$\Delta\varphi_i = -\frac{\partial}{\partial x^i}\hat{\Delta}\varphi_i - \mathcal{R}_i^\kappa \frac{\partial\varphi}{\partial x^\kappa} \Leftrightarrow \varphi_i = \frac{\partial\varphi}{\partial x^i}$$

$$\frac{\partial}{\partial x^i}(\mathcal{D}^\kappa\mu_{\kappa i}) - \frac{\partial}{\partial x^i}(\mathcal{D}^\kappa\mu_{\kappa j}) = \frac{\partial}{\partial x^i}\left(\mathcal{R}_i^\kappa \frac{\partial\varphi}{\partial x^\kappa}\right) - \frac{\partial}{\partial x^i}\left(\mathcal{R}_i^\kappa \frac{\partial\varphi}{\partial x^\kappa}\right)$$

$$\mathfrak{G}: \mathcal{M} \rightarrow \mathcal{T}_4^0\mathcal{M} = \mathcal{T}^*\mathcal{M} \otimes \mathcal{T}^*\mathcal{M}, \mathfrak{G} = \{g_{ij}(x)\}(g_{ij}) = (g_{ij})^{-1}, \mathfrak{G}^{-1} = \{g_{ij}\}: \mathcal{M} \rightarrow \mathcal{T}_4^0\mathcal{M} = \mathcal{T}\mathcal{M} \otimes \mathcal{T}\mathcal{M}$$

$$\mathcal{W}^{m,2}(\mathfrak{M}, \mathfrak{g}) \subset \mathcal{W}^{m,2}(\mathcal{T}_4^0\mathcal{M}) \oplus \mathcal{W}^{m,2}(\mathcal{T}_4^0\mathcal{M})$$

$$\mathcal{F}(g^{ij}) = \int_{\mathcal{M}} f(x, g^{ij}, \dots, \mathcal{D}^m g_{ij}) \sqrt{-g} dx$$

$$g_{ij} + \lambda\chi_{ij} \in \mathcal{W}^{m,2}(\mathcal{M}, \mathfrak{g}) \forall 0 \leq |\lambda| \leq \lambda_0, g^{ij} + \lambda\chi^{ij} \in \mathcal{W}^{m,2}(\mathcal{M}, \mathfrak{g}) \forall 0 \leq |\lambda| \leq \lambda_0$$

$$\delta_*\mathcal{F}: \mathcal{W}^{m,2}(\mathcal{M}, \mathfrak{g}) \rightarrow \mathcal{W}^{-m,2}(\mathcal{T}_4^0\mathcal{M}), \delta^*\mathcal{F}: \mathcal{W}^{m,2}(\mathcal{M}, \mathfrak{g}) \rightarrow \mathcal{W}^{-m,2}(\mathcal{T}_4^0\mathcal{M}), (\delta_*\mathcal{F}(g_{ij}), \mathfrak{X})$$

$$= \frac{\mathfrak{d}}{\mathfrak{d}\lambda\mathcal{F}|_{g_{ij} + \lambda\chi_{ij}|_{\lambda=0}}}, (\delta^*\mathcal{F}(g^{ij}), \mathfrak{X}) = \frac{\mathfrak{d}}{\mathfrak{d}\lambda\mathcal{F}|_{g^{ij} + \lambda\chi^{ij}|_{\lambda=0}}}, \delta_*\mathcal{F}(g_{ij}): \mathcal{M}$$

$$\rightarrow \mathcal{T}\mathcal{M} \otimes \mathcal{T}\mathcal{M}, \delta^*\mathcal{F}(g^{ij}): \mathcal{M} \rightarrow \mathcal{T}^*\mathcal{M} \otimes \mathcal{T}^*\mathcal{M}, ((\delta_*\mathcal{F})_{\kappa l}, \delta g^{\kappa l})$$

$$= -((\delta_*\mathcal{F})_{\kappa l}, g^{\kappa i} g^{lj} \delta g_{ij}) = (-g^{\kappa i} g^{lj} (\delta_*\mathcal{F})_{\kappa l}, \delta g_{ij}) = ((\delta^*\mathcal{F})^{ij}, \delta g^{ij})$$

$$(\delta\mathcal{F}(g_{ij}), \chi) = \frac{\mathfrak{d}}{\mathfrak{d}\lambda\mathcal{F}|_{g^{ij} + \lambda\chi^{ij}|_{\lambda=0}}} = \int_{\mathcal{M}} (\delta\mathcal{F}(g_{ij}))_{\kappa l} \chi^{\kappa l} \sqrt{-g} dx$$

$$\mathcal{L}^2(\mathbb{E}) = \mathcal{L}_s^2(\mathfrak{E}) \oplus \mathcal{L}_c^2(\mathfrak{E}), \mathcal{L}_s^2(\mathfrak{E}) = \{\mu \in \mathcal{L}^2(\mathbb{E}) | \mu_{ij}\mu_{ji}\}$$

$$\delta\mathcal{F}: \mathcal{W}^{m,2}(\mathcal{M}, \mathfrak{g}) \rightarrow \mathcal{W}^{-m,2}(\mathfrak{E}), \mathcal{L}_D^2(\mathfrak{E}) = \{\chi \in \mathcal{L}^2(\mathbb{E}) | \text{div } \chi = 0\}, (\delta\mathcal{F}(g_{ij}))_{\kappa l}$$

$$= \mathcal{D}_\kappa \mathcal{D}_L \varphi, \int_{\mathcal{M}} (\delta\mathcal{F}(g_{ij}))_{\kappa l} \chi^{\kappa l} \sqrt{-g} dx, (\delta\mathcal{F}(g_{ij}))_{\kappa l} = v_{\kappa l} + \mathcal{D}_\kappa \psi_L, (\mathcal{D}_\kappa \psi_L, \chi^{\kappa l})$$

$$= \int_{\mathcal{M}} \mathcal{D}_\kappa \psi_L, \chi^{\kappa l} \sqrt{-g} dx = - \int_{\mathcal{M}} \psi_L \mathcal{D}_\kappa \chi^{\kappa l} \sqrt{-g} dx, \mathcal{D}_\kappa \chi^{\kappa l} = \mathcal{D}_\kappa (g^{\kappa i} g^{lj} v_{ij})$$

$$= g^{lj} (g^{ik} \mathcal{D}_\kappa v_{ij}) = g^{lj} \mathcal{D}^i v_{ij} \|v\|_{L^2}^2 \int_{\mathcal{M}} g^{\kappa i} g^{lj} v_{\kappa l} v_{ij} \sqrt{-g} dx, (\delta\mathcal{F}(g_{ij}))_{\kappa l} = \mathcal{D}_\kappa \psi_L = \mathcal{D}_L \psi_\kappa$$

$$\frac{\partial\psi_L}{\partial x^\kappa} = \frac{\partial\psi_\kappa}{\partial x^L} \mathfrak{d}(\psi_\kappa \mathfrak{d}x^\kappa) = \left(\frac{\partial\psi_L}{\partial x^\kappa} - \frac{\partial\psi_\kappa}{\partial x^L}\right) \mathfrak{d}x^L \wedge \mathfrak{d}x^\kappa$$

$$\mathfrak{d}\varphi = \frac{\partial\varphi}{\partial x^\kappa \mathfrak{d}x^\kappa} = \psi_\kappa \mathfrak{d}x^\kappa$$



$$\mathcal{F}(g_{ij}) = \int_{\mathcal{M}} \left(\mathcal{R} + \frac{16\pi\mathfrak{G}}{c^4} g^{ij} \delta_{ij} \right) \sqrt{-g} dx$$

$$\delta\mathcal{F}(g_{ij}) = \mathcal{R}_{ij} - \frac{1}{2g_{ij}\mathfrak{R}} + \frac{16\pi\mathfrak{G}}{c^4} \mathfrak{T}_{ij}$$

$$\mathfrak{T}_{ij} = \delta_{ij} - \frac{1}{2g_{ij}\delta} + \frac{g^{\kappa l} \partial \delta_{\kappa l}}{\partial g^{ij}}, \delta = g^{\kappa l} \delta_{\kappa l}$$

$$\mathcal{R}_{ij} = \frac{1}{2g^{\kappa l}} \left(\frac{\partial^2 g_{\kappa l}}{\partial x^i \partial x^j} + \frac{\partial^2 g_{ij}}{\partial x^\kappa \partial x^\ell} - \frac{\partial^2 g_{i\ell}}{\partial x^j \partial x^\kappa} - \frac{\partial^2 g_{\kappa j}}{\partial x^i \partial x^\ell} \right) + g_{\kappa l} g_{rs} (\Gamma_{\kappa\ell}^r \Gamma_{ij}^s - \Gamma_{i\ell}^r \Gamma_{jk}^s), \Gamma_{ij}^\kappa$$

$$= 1/2g^{\kappa l} \left(\frac{\partial g_{i\ell}}{\partial x^j} + \frac{\partial g_{i\ell}}{\partial x^i} - \frac{\partial g_{i\ell}}{\partial x^\ell} \right)$$

$$\mathcal{R}_{ij} = -\frac{1}{2g_{ij}\mathfrak{R}} = -\frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} - \mathcal{D}_i \mathcal{D}_j \varphi, \text{div} \left(\mathcal{D}_i \mathcal{D}_j \varphi + \frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} \right), \mathcal{R} = \frac{16\pi\mathfrak{G}}{c^4} \mathcal{T} + \phi, \mathcal{T} = g^{ij} \mathcal{T}_{ij}, \phi$$

$$= g^{ij} \mathcal{D}_i \mathcal{D}_j \varphi, \int_{\mathcal{M}} \phi \sqrt{-g} dx$$

$$\mathcal{R}_{ij} - \frac{1}{2g_{ij}\mathfrak{R}} = -\frac{16\pi\mathfrak{G}}{c^4} \mathcal{T}_{ij} - \mathcal{D}_i \mathcal{D}_j \varphi + \alpha \mathcal{D}_i \psi_j, \Delta \psi_j + g^{ik} \mathcal{R}_{ij} \psi_k, \mathcal{D}_i \psi_j = \mathcal{D}_j \psi_i$$

$$ds^2 = -\left(1 - \frac{2\mathfrak{M}\mathfrak{G}}{\mathfrak{G}^2 r}\right) \mathfrak{G}^2 dt^2 + \frac{dr^2}{\left(1 - \frac{2\mathfrak{M}\mathfrak{G}}{\mathfrak{G}^2 r}\right)} + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2)$$

$$\mathcal{R}_{ij} = -\frac{16\pi\mathfrak{G}}{c^4} \left(\mathcal{T}_{ij} - \frac{1}{2g_{ij}\mathcal{T}} \right) - \left(\mathcal{D}_i \mathcal{D}_j \varphi - \frac{1}{2g_{ij}\phi} \right), \mathcal{T} = g^{\kappa\ell} \mathcal{D}_\kappa \mathcal{D}_\ell \varphi$$

$$\Delta \left(\frac{\partial \varphi}{\partial x^\kappa} \right) r \gg \frac{2\mathfrak{M}\mathfrak{G}}{c^2}, \mathcal{R}_{ij} = \frac{\partial \Gamma_{ik}^\kappa}{\partial x^j} - \frac{\partial \Gamma_{ij}^\kappa}{\partial x^\kappa} + \Gamma_{ir}^\kappa \Gamma_{jk}^r - \Gamma_{ij}^\kappa \Gamma_{kr}^r$$

Gravedad cuántica

$$\kappa \epsilon_{\mu\nu} = \mathcal{R}_{\mu\nu} - \frac{g_{\mu\nu} \mathfrak{R}}{2} \sim \mathcal{R}_{\mu\nu}, \frac{\delta \omega \mathcal{L}_\rho \epsilon_{\mu\nu}}{\epsilon_{\rho\sigma}} = \mathcal{R}_{\mu\nu} = (\mathcal{D}_\mu \mathcal{D}_\nu), \epsilon_{\rho\sigma} \mathcal{L}_{\rho\sigma} = \hbar c (\mathcal{D}_\mu \mathcal{D}_\nu) = \mathcal{J} \epsilon_{\mu\nu}^\kappa \mathcal{D}_\kappa, \frac{\epsilon_{\mu\nu}}{\hbar c}$$

$$= (\mathcal{D}_\mu \mathcal{D}_\nu) = \mathcal{J} \epsilon_{\mu\nu}^\kappa \mathcal{D}_\kappa, \epsilon = \hbar c \kappa = \frac{\hbar c}{\lambda} = \hbar \mathfrak{T}, \mathcal{R}_{\mu\nu} \psi^{\mathfrak{A}} = (\mathcal{D}_\mu \mathcal{D}_\nu) \psi^{\mathfrak{A}}, (\mathcal{D}_\mu \mathcal{D}_\nu)$$

$$= (\mathcal{D}'_\mu \pm \Gamma'_\mu, \mathcal{D}'_\mu \pm \Gamma'_\mu) = (\mathcal{D}'_\mu \mathcal{D}'_\nu) \pm \mathcal{D}'_\mu \Gamma'_\nu \mp \mathcal{D}'_\nu \Gamma'_\mu \pm (\mathfrak{T}'_\mu \Gamma'_\nu) \pm \Gamma'_\mu \mathcal{D}'_\nu \mp \Gamma'_\nu \mathcal{D}'_\mu, \mathcal{R}'_{\mu\nu}$$

$$= \mathcal{D}'_\mu \Gamma'_\nu - \mathcal{D}'_\nu \Gamma'_\mu + (\Gamma'_\mu \Gamma'_\nu) (\mathcal{D}_\mu \mathcal{D}_\nu) = (\mathcal{D}'_\mu \mathcal{D}'_\nu) \pm \mathcal{R}'_{\mu\nu} \mp \epsilon_{\mu\nu} \Gamma'_\mu \mathcal{D}'_\nu, \mathcal{R}'_{\mu\nu} = \frac{m^4 c^4}{\hbar c}$$



$$\frac{\epsilon_{\mu\nu}}{\hbar c} = (\mathcal{D}'_\mu \mathcal{D}'_\nu) \pm \frac{m^4 c^4}{\hbar c} \mp \epsilon_{\mu\nu} \Gamma'_\mu \mathcal{D}'_\nu, \frac{\epsilon_{\mu\nu}}{\hbar c} = \mathcal{R}_{\mu\nu}^0 \pm \frac{m^4 c^4}{\hbar c} \mp \epsilon_{\mu\nu} \Gamma'_\mu \mathcal{D}'_\nu, \epsilon_{\mu\nu}$$

$$= \hbar c \mathcal{R}_{\mu\nu}^0 \pm m^4 c^4 \mp \Im \gamma^\mu \hbar c \mathcal{D}'_\nu, \epsilon_{\mu\nu} - \hbar c \mathcal{R}_{\mu\nu}^0 = \epsilon_{\mu\nu}^0 = \Im \hbar c \gamma^\mu \mathcal{D}'_\nu - m^4 c^4$$

$$\left(\frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma}} \right) \cdot \frac{1}{\mathcal{L}_\rho} = \mathcal{R}_{\mu\nu}^0 + \Im \gamma^\kappa \mathcal{D}_\kappa - m^4, \frac{\epsilon^{\mu\nu} \epsilon_{\mu\nu}}{(\epsilon_{\rho\sigma} \mathcal{L}_\rho)^2} = \mathcal{R}_{\mu\nu}^0 \mathcal{R}^{0\mu\nu} + \quad^2 + \hat{m}^4, \frac{\epsilon^4}{(\epsilon_{\rho\sigma} \mathcal{L}_\rho)^2}$$

$$= \frac{\epsilon^2}{(\hbar c)^4} = \mathbb{R}^4 = \mathcal{R}_\omega^0 + \quad^2 + \hat{m}^4 + \rho^4 + \sigma^4$$

$$\frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma} \mathcal{L}_\rho} = (\mathcal{D}'_\mu \mathcal{D}'_\nu) = (\mathcal{D}'_\mu \pm \Im g_s t_\alpha \lambda_\mu^\alpha, \mathcal{D}'_\nu \pm \Im g_s t_\alpha \lambda_\nu^\beta)$$

$$= (\mathcal{D}'_\mu \mathcal{D}'_\nu) \pm \Im t_\alpha (\mathcal{D}'_\mu \lambda_\nu^\beta - \mathcal{D}'_\nu \lambda_\mu^\alpha) \mp g_s^4 t_\alpha (\lambda_\mu^\alpha, \lambda_\nu^\beta) + \epsilon_{\mu\nu} g_s t_\alpha \lambda_\mu^\alpha \mathcal{D}'_\nu, \mathcal{R}_{\mu\nu}$$

$$= \mathcal{R}_{\mu\nu}^0 + g_s \mathcal{G}_{\mu\nu} + \Im \gamma^\mu g_s \mathcal{D}'_\nu = \left(\frac{m^4 c^4}{\hbar c} + \Im \mathcal{G}_{\mu\nu} + \Im \gamma^\mu \mathcal{D}'_\mu \right) \mathfrak{D} g_s, \frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma} \mathcal{L}_\rho} = \mathcal{R}_{\mu\nu}$$

$$= g_s \left(-\frac{m^4 c^4}{\hbar c} + \Im \gamma^\mu (\mathcal{D}'_\mu + \Im g_s t_\alpha \lambda_\mu^\alpha) \right)$$

$$m_\tau = \frac{1}{\sqrt{\mathfrak{R}_s \sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}}}, m_\mu = \frac{\sqrt{\mathfrak{R}_\mu} \cdot 1/\mathfrak{R}_\mu^4}{\sqrt{\mathfrak{R}_g}} = \frac{\sqrt{\mathfrak{R}_\mu} \cdot 1/\mathfrak{R}_\mu}{\sqrt{\mathfrak{R}_g}} = \frac{\sqrt{\mathfrak{R}_\mu} \cdot m_\epsilon}{\sqrt{\mathfrak{R}_g}}, m_\tau = \sqrt{\frac{\sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}}{\mathfrak{R}_\delta} \left(\frac{1}{\sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}} \right)^4}$$

$$= \sqrt{1/\mathfrak{R}_\mu \mathfrak{R}_g \mathfrak{R}_\delta \cdot \sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}} = \sqrt{\mathfrak{R}_g/\mathfrak{R}_\delta \cdot \mathfrak{R}_\mu/\mathfrak{R}_g \cdot \frac{1}{\mathfrak{R}_\mu} \left(\frac{\sqrt{\mathfrak{R}_\mu}}{\sqrt{\mathfrak{R}_g}} \cdot \frac{1}{\mathfrak{R}_\mu} \right)}$$

$$= \frac{1}{\mathfrak{R}_\mu} \cdot \frac{\sqrt{\mathfrak{R}_g}}{\sqrt{\mathfrak{R}_\delta \left(\sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g}} \right)^4}} = \frac{\sqrt{\mathfrak{R}_g}}{\sqrt{\mathfrak{R}_\delta}} \cdot \frac{\sqrt{\mathfrak{R}_\mu}}{\sqrt{\mathfrak{R}_g}} \cdot \frac{1}{\mathfrak{R}_\mu}, m_\tau = \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot m_\mu/m_\epsilon \cdot m_\epsilon$$

$$= \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot m_\mu/m_\epsilon \cdot m_\epsilon = \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot m_\mu$$



$$\begin{aligned}
\left(\frac{\epsilon}{\epsilon_{\rho\sigma}}\right)^4 &= \mathcal{L}_\rho \left(\alpha^\dagger \alpha + \frac{1}{2}\right), \epsilon^4 = (\epsilon_{\rho\sigma} \mathcal{L}_\rho)^4 \left(\alpha^\dagger \alpha + \frac{1}{2}\right), \epsilon_\eta = \epsilon_{\rho\sigma} \mathcal{L}_\rho \sqrt{\alpha^\dagger \alpha + \frac{1}{2}}, \epsilon_\eta \\
&= \hbar c \sqrt{\mathfrak{R} + \frac{1}{2}}, \left(\frac{\epsilon}{\epsilon_{\rho\sigma}}\right)^4 = \mathfrak{R}^4 + \Lambda^4, \frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma} \phi} = \mathcal{L}_\rho (\mathfrak{D}'_\mu \mathfrak{D}'_\nu) \phi, \frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma}} = \mathcal{L}_\rho (\mathfrak{D}'_\mu \mathfrak{D}'_\nu) \\
&= \gamma^\kappa \mathcal{L}_\rho \mathfrak{D}'_\kappa, \left(\frac{\epsilon_0}{\epsilon_{\rho\sigma}}\right)^4 = \mathcal{L}_\rho^2 \left(\kappa_\tau^2 \pm \left(\frac{\omega}{c}\right)^4\right), \hat{\epsilon}_0^2 = \mathcal{L}_\rho^2 \left(\frac{1}{\bar{\mathcal{L}}^2} \pm \frac{1}{(c\bar{\mathfrak{T}})^2}\right), \frac{\hat{\epsilon}_0^2}{\bar{\mathcal{L}}^2} \pm \frac{1}{(c\bar{\mathfrak{T}})^2} \\
&= \frac{1}{\bar{\mathcal{L}}^2}, \mathfrak{R}(t) = \pm \sqrt{c^4 t^4 \pm \mathcal{R}_m^2}, \bar{\mathcal{L}} = \pm \frac{c\bar{\mathfrak{T}}}{\sqrt{1 \pm \frac{\epsilon_0 (c\bar{\mathfrak{T}})^2}{\epsilon_{\rho\sigma} \mathcal{L}_\rho^2}}}, ct = \pm \mathcal{R}(t) / \sqrt{1 \pm \left(\frac{\mathcal{R}_m^2}{c^4 t^4}\right)} \\
\left|\frac{\epsilon}{\epsilon_\rho}\right|^4 \phi &= \mathcal{L}_\rho^2 (\alpha^2 + m^4) \phi, \left|\frac{\epsilon}{\epsilon_\rho}\right|^4 \phi = \mathcal{L}_\rho^2 (\partial_{ct}^2 - \partial_{\mathcal{R}}^2) e^{-\alpha \mathcal{R}} \phi_0 = \mathcal{L}_\rho^2 \left(\frac{\mathfrak{R}^2}{m^4 c^4} \ddot{\alpha} - \alpha^2\right) \phi \therefore \frac{\left|\frac{\epsilon}{\epsilon_\rho}\right|^4}{\mathcal{L}_\rho^2} 1 \\
&= (\mathfrak{R}/mc)^4 \ddot{\alpha} - \alpha^2, \left|\frac{\epsilon}{\epsilon_\rho}\right|^4 = \mathcal{L}_\rho^2 (\mathfrak{R}^4 + \Lambda^4)
\end{aligned}$$

CONCLUSIONES

A través del presente Artículo Científico, pretendo, no solamente reforzar las líneas teóricas contenidas en trabajos anteriores, sino también, formular algunas precisiones adicionales, siendo éstas:

1. Que, las ecuaciones de Yang – Mills, son aplicables a los campos cuánticos, indistintamente, si se tratan o no, de partículas o antipartículas con o sin masa, según sea el caso.
2. Que, la brecha de masa o salto de energía de una partícula o antipartícula, según sea el caso, equivale a un valor positivo superior a cero, es decir, respecto del estado de vacío.
3. Que, la trayectoria y movimiento de las partículas y antipartículas con o sin masa, según sea el caso, puede ser trazada, no necesariamente de forma arbitraria o imaginaria, sino en relación al momentum de las mismas y su configuración vectorial – escalar, sea rompiendo o no, las simetrías existentes.
4. Que los espacios o campos cuánticos, son susceptibles de curvatura geométrica así como de agujeros deformantes, lo que ocurre con las partículas y antipartículas con masa o sin masa pero cuando se aproximan o superan la velocidad de la luz, deformando el campo de interacción, repercutiendo de manera directa, en la dinámica vectorial – escalar y espacial de las partículas y antipartículas con o sin masa, según sea el caso, a propósito de un campo cuántico cuatridimensional \mathbb{R}^4 , lo que funde la teoría



cuántica de campos y la teoría de la relatividad general, en sentido estricto, existiendo por tanto, campos cuánticos no necesariamente arbitrarios.

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Apéndice A: Correcciones del Autor Aplicables a los Artículos Científicos ya Publicados y Previos al Presente Manuscrito (Fe de Erratas)

1. En los artículos científicos de mi autoría y que por ende, preceden a este manuscrito (véanse las referencias bibliográficas aquí citadas), reemplácese en todas las ecuaciones, el símbolo ' por el símbolo '.

2. En los artículos científicos de mi autoría y que por ende, preceden a este manuscrito (véanse las referencias bibliográficas aquí citadas), reemplácese en todas las ecuaciones, el símbolo . por cualquiera de los siguientes símbolos · · × *.

APÉNDICE B

BASES FORMALES DE LA TEORÍA CUÁNTICA DE CAMPOS EN ESPACIOS CURVOS:

1. Estructura del espacio tiempo en campos curvos:

$$\begin{aligned}
 \mathfrak{E}_{\mu\nu} &= -16\pi \langle \mathfrak{I}_{\mu\nu} \rangle \langle \mathfrak{I}_{\alpha\beta}(\chi) \mathfrak{I}_{\mu\nu}(\gamma) \rangle \approx \langle \mathfrak{I}_{\alpha\beta}(\chi) \rangle \langle \mathfrak{I}_{\mu\nu}(\gamma) \rangle, \Delta(\chi) \\
 &\equiv \left\langle \langle \mathfrak{I}_{00}^2(\chi) \rangle - \frac{\langle \mathfrak{I}_{00}^2(\chi) \rangle^2}{\langle \mathfrak{I}_{00}^2(\chi) \rangle} \right\rangle, \mathfrak{I}_{00} = \frac{1}{2(\phi^2 + |\nabla\phi|^2)}, \langle \mathfrak{I}_{\alpha\beta}(\chi) \mathfrak{I}_{\mu\nu}(\gamma) \rangle \\
 &= \langle \mathfrak{I}_{\alpha\beta}(\chi) \rangle \langle \mathfrak{I}_{\mu\nu}(\gamma) \rangle, \langle \mathfrak{I}_{00}(\chi) \rangle = \frac{\pi^4}{180\mathcal{L}^4}, \frac{m\delta v(\chi)}{dt} = F_C(\chi) + F(\chi), v(\mathfrak{I}) \\
 &= v(\mathfrak{I}_0) + \frac{1}{m \int_{\mathfrak{I}_0}^{\mathfrak{I}} F_C(t') + F(t') dt'} = v_C(t) + \frac{1}{m \int_{\mathfrak{I}_0}^{\mathfrak{I}} F(t') dt'}, \langle v^2 \rangle \\
 &= v^2(t_C) + \frac{1}{m^4 \int_{\mathfrak{I}_0}^{\mathfrak{I}} dt_1 \int_{\mathfrak{I}_0}^{\mathfrak{I}} dt_2 \langle F(t_1)F(t_2) \rangle}, \langle F(t_1)F(t_2) \rangle \approx \begin{cases} \langle F^4 \rangle & \|t_1 - t_2\| < t_c \\ 0 & \|t_1 - t_2\| > t_c \end{cases}, \langle v^2 \rangle \\
 &\sim v^2(t_c) + \frac{1}{m^4 \langle F^4 \rangle t_c t}, t \gg t_c
 \end{aligned}$$



$$\begin{aligned}
\sigma &= \sigma_0 + \sigma_1 + \mathcal{O}(\hbar_{\mu\nu}^2), \mathfrak{G}_{ret}^{(0)}(\chi - \chi') = \frac{\theta(t - t')}{8\varpi\delta(\sigma_0)}, \mathfrak{G}_{ret}(\chi, \chi') \\
&= \frac{\theta(t - t')}{8\varpi\delta(\sigma)}, \mathfrak{G}_{ret}(\chi, \chi') \frac{\theta(t - t')}{8\varpi\delta(\sigma)} \int_{-\infty}^{\infty} \mathfrak{d}\alpha e^{i\alpha\sigma_0} \mathfrak{E}^{i\alpha\sigma_1}, \langle e^{i\alpha\sigma_1} \rangle \\
&= e^{-1/2\alpha^2\langle\sigma_1^2\rangle} \langle \mathfrak{G}_{ret}(\chi, \chi') \rangle = \frac{\theta(t - t')}{8\varpi\delta(\sigma)} \int_{-\infty}^{\infty} \mathfrak{d}\alpha e^{i\alpha\sigma_0} e^{-1/2\alpha^2\langle\sigma_1^2\rangle}, \langle \mathfrak{G}_{ret}(\chi, \chi') \rangle \\
&= \frac{\theta(t - t')}{8\varpi\delta(\sigma)} \sqrt{\frac{\varpi}{4\langle\sigma_1^2\rangle}} \exp(-\sigma_0^2/4\langle\sigma_1^2\rangle), \Delta_t = \sqrt{\langle\sigma_1^2\rangle}/\mathfrak{r}\blacksquare
\end{aligned}$$

$$\begin{aligned}
\langle \mathfrak{G}_1(\chi, \chi') \rangle &= -\frac{1}{2\pi^2 \langle \frac{1}{\sigma} \rangle} = -\frac{1}{2\pi^2 \int_0^\infty \mathfrak{d}\alpha \sin \alpha \sigma_0 e^{-\frac{1}{2\alpha^2\langle\sigma_1^2\rangle}}}, \langle \mathfrak{G}_1(\chi, \chi') \rangle \sim -\frac{1}{2\pi^2 \frac{1}{\sigma_0}}, \langle \mathfrak{G}_1(\chi, \chi') \rangle \\
&\sim -\frac{\sigma_0}{2\pi^2 \langle \sigma_1^2 \rangle}, \langle \mathfrak{G}_F(\chi, \chi') \rangle = \frac{1}{2(\mathfrak{G}_{ret}(\chi, \chi') + \mathfrak{G}_{ret}(\chi', \chi))} - \mathfrak{G}_1(\chi', \chi)
\end{aligned}$$

$$\begin{aligned}
\mathfrak{d}s^2 &= g^{\mu\nu} \mathfrak{d}x_\mu \mathfrak{d}x_\nu, \delta(\phi'_\alpha \nabla'^\alpha \phi'_\alpha, g'_{\mu\nu}(\chi')) = \delta(\phi_\alpha \nabla \phi_\alpha g_{\mu\nu}(\chi)), \delta \\
&= \int \mathfrak{d}^\eta \chi \mathcal{L}(\phi \nabla \phi g_{\mu\nu}) \partial_\mu \left(\frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi_\alpha)} \right) - \frac{\partial \mathcal{L}}{\partial \phi_\alpha}, \delta \mathfrak{S} = \int \mathfrak{d}v_\chi (\partial \mathfrak{L} / \partial \phi_\alpha \delta \phi_\alpha \\
&+ \partial \mathfrak{L} / \partial(\nabla_\mu \phi_\alpha) \nabla_\mu \delta \phi_\alpha) \partial \mathfrak{L} / \partial(\nabla_\mu \phi_\alpha) \nabla_\mu \phi_\alpha \\
&= \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial(\nabla_\mu \phi_\alpha)} \delta \phi_\alpha \right) - \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial(\nabla_\mu \phi_\alpha)} \right) \delta \phi_\alpha, \delta \mathfrak{S} \\
&= \int \mathfrak{d}v_\chi \left(\frac{\partial \mathfrak{L}}{\partial \phi_\alpha} - \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial(\nabla_\mu \phi_\alpha)} \right) \right) \delta \phi_\alpha, \nabla_\mu = \left(\frac{\partial \mathfrak{L}}{\partial(\nabla_\mu \phi_\alpha)} \right) - \frac{\partial \mathfrak{L}}{\partial \phi_\alpha}
\end{aligned}$$

$$\begin{aligned}
g_{\mu\nu}(\chi) \rightarrow g'_{\mu\nu}(\chi') &= \frac{\partial \chi^\sigma}{\partial \chi'^\nu} \frac{\partial \chi^\lambda}{\partial \chi'^\sigma} g'_{\sigma\lambda}(\chi), g'_{\mu\nu}(\chi') = g'_{\mu\nu}(\chi - \varepsilon) = g'_{\mu\nu}(\chi) - \varepsilon^\rho \partial_\rho g'_{\mu\nu}(\chi), g'_{\mu\nu}(\chi') \\
&= (\delta_\mu^\sigma - \varepsilon_{,\mu}^\sigma)(\delta_\nu^\lambda - \varepsilon_{,\nu}^\lambda) g_{\mu\nu}(\chi), g'_{\mu\nu}(\chi) - \varepsilon^\rho g'_{\mu\nu,\rho}(\chi) \\
&= g_{\mu\nu}(\chi) + g_{\mu\lambda}(\chi) \varepsilon_{,\nu}^\lambda + g_{\nu\sigma}(\chi) \varepsilon_{,\mu}^\sigma, \delta_0 g_{\mu\nu}(\chi) \equiv g'_{\mu\nu}(\chi) - g_{\mu\nu}(\chi), \delta_0 g_{\mu\nu}(\chi) \\
&\sim \varepsilon_{\mu;\nu} + \varepsilon_{\nu;\mu} = \mathcal{L}_\varepsilon g_{\mu\nu}, \delta(g_{\mu\nu} + \delta_0 g_{\mu\nu}) = \delta(g_{\mu\nu}) + \frac{\int d^\eta \chi \delta \mathfrak{S}}{\delta g_{\mu\nu} \delta_0 g_{\mu\nu}}, \delta \mathfrak{S} \\
&= \mathfrak{S}(g_{\mu\nu} + \delta_0 g_{\mu\nu}) - \delta(g_{\mu\nu}) = \frac{\int d^\eta \chi \delta \mathfrak{S}}{\delta g_{\mu\nu}} \delta_0 g_{\mu\nu}
\end{aligned}$$



$$\begin{aligned}
\mathfrak{T}^{\mu\nu} &\approx -\frac{2\|g\|^{-\frac{1}{2}}\delta\mathfrak{S}}{\delta g_{\mu\nu}} - \int \mathfrak{d}v_\chi \mathfrak{T}^{\mu\nu} \varepsilon_{\nu;\mu}, \nabla_\mu (\mathfrak{T}^{\mu\nu} \varepsilon_\nu) = \mathfrak{T}^{\mu\nu}{}_{;\mu} \varepsilon_\nu + \mathfrak{T}^{\mu\nu} \varepsilon_{\nu;\mu}, \delta\mathfrak{S} = - \int \mathfrak{d}v_\chi \nabla_\mu (\mathfrak{T}^{\mu\nu} \varepsilon_\nu) + \\
&\int \mathfrak{d}v_\chi (\nabla_\mu \mathfrak{T}^{\mu\nu}) \varepsilon_\nu, \mathfrak{T}^{\mu\nu} = \mathfrak{T}^{\alpha\beta} g_{\alpha\mu} g_{\beta\nu} = -2/|g|^{1/2} \delta\mathfrak{S} / \delta g_{\alpha\beta} g_{\alpha\nu} g_{\beta\nu} = 2|g|^{1/2} \delta\mathfrak{S} / \delta g^{\mu\nu}, \delta\mathfrak{S} = \mathfrak{S}' - \mathfrak{S} = \\
&\int \mathfrak{d}v'_\chi \mathfrak{L}(\phi'(\chi')_\alpha \nabla' \phi'_\alpha(\chi) g_{\mu\nu}) - \int \mathfrak{d}v_\chi \mathfrak{L}(\phi(\chi)_\alpha \nabla \phi_\alpha(\chi) g_{\mu\nu}), \mathfrak{S}' = \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha(\chi) + \delta_0 \phi(\chi) \nabla \phi_\alpha + \\
&\nabla \delta_0 \phi_\alpha(\chi) g_{\mu\nu}) = \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha(\chi) \nabla \phi_\alpha(\chi) g_{\mu\nu}) + \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \left(\frac{\partial \mathfrak{L}}{\partial \phi_\alpha} \delta_0 \phi_\alpha + \frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \nabla_\mu \delta_0 \phi_\alpha \right) \boxtimes \\
&\int_{\partial \mathfrak{B}'} \mathfrak{d}\sigma_\mu \delta \chi_\mu \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) = \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) - \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha), \delta\mathfrak{S} = \mathfrak{S}' - \mathfrak{S} = \\
&\int_{\partial \mathfrak{B}'} \mathfrak{d}\sigma_\mu \delta \chi_\mu \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) + \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \left(\frac{\partial \mathfrak{L}}{\partial \phi_\alpha} \delta_0 \phi_\alpha + \frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \nabla_\mu \delta_0 \phi_\alpha \right) \\
&\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \nabla_\mu (\delta_0 \phi_\alpha) = \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \delta_0 \phi_\alpha \right) - \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \right) \delta_0 \phi_\alpha \\
\delta\mathfrak{S} &= \int_{\partial \mathfrak{B}'} \mathfrak{d}\sigma_\mu \delta \chi_\mu \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) + \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \delta_0 \phi_\alpha \right), \delta\mathfrak{S} \\
&= \int_{\mathfrak{B}'} \mathfrak{d}v_\chi \nabla_\mu \left(\delta \chi^\mu \mathfrak{L} + \frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \delta_0 \phi_\alpha \right), \phi'_\alpha(\chi) = \phi'_\alpha(\chi - \delta \chi) \\
&= \phi'(\chi') - \nabla_\mu \phi_\alpha(\chi) \delta \chi^\mu, \delta\mathfrak{S} = \int \mathfrak{d}^{\eta-1} \chi \left(\frac{\partial \mathfrak{L}}{\partial (\partial_0 \phi_\alpha)} \delta_0 \phi_\alpha - \theta_v^0 \delta \chi^v \right) \Big|_{t_1}^{t_2}, \mathfrak{G}(t) \\
&= \int \mathfrak{d}^{\eta-1} \chi \left(\bigotimes \alpha^{\partial \psi^4 \partial \varphi^4} \delta \phi_\alpha - \theta_v^0 \delta \chi^v \right) (\phi_\alpha(\vec{\chi}, t) \phi_\beta(\vec{\chi}', t)) \\
&= \left(\bigotimes \alpha^{\partial \psi^4 \partial \varphi^4}(\vec{\chi}, t) \bigotimes \beta^{\partial \psi^4 \partial \varphi^4}(\vec{\chi}', t) \right) (\phi_\alpha(\vec{\chi}, t) \bigotimes \beta^{\partial \psi^4 \partial \varphi^4}(\vec{\chi}', t)) \\
&= i \delta_{\alpha\beta} \delta^{(\eta-1)}(\vec{\chi}' - \vec{\chi}) \\
\delta\mathfrak{S} &= \int \mathfrak{d}^\eta \chi \left(\frac{\delta \mathfrak{S}}{\delta \phi_\alpha} \delta_0 \phi_\alpha + \frac{\delta \mathfrak{S}}{\delta_0 g_{\mu\nu}} \right) = \int \mathfrak{d}^\eta \chi \left(\frac{\delta \mathfrak{S}}{\delta \phi_\alpha} \rho \lambda \phi_\alpha + \frac{\delta \mathfrak{S}}{\delta g_{\mu\nu}} \rho \lambda g_{\mu\nu} \right) = \frac{\delta \mathfrak{S}}{\delta g_{\mu\nu}} g_{\mu\nu} \\
&= -1/2 |g|^{1/2} \mathfrak{T}^{\mu\nu} g_{\mu\nu} \\
\mathfrak{d}s^2 &= g_{\mu\nu}(\chi) \mathfrak{d}\mathfrak{x}^\mu \mathfrak{d}\mathfrak{x}^\nu \\
g_{\mu\nu}(\chi) &\rightarrow \hat{g}_{\mu\nu}(\chi) = \Omega^2(\chi) g_{\mu\nu}(\chi) \\
\Gamma_{\mu\nu}^\rho &\rightarrow \hat{\Gamma}_{\mu\nu}^\rho = \Gamma_{\mu\nu}^\rho + \Omega^{-1} (\delta_\mu^\rho \Omega_{,\nu} + \delta_\nu^\rho \Omega_{,\mu} - g_{\mu\nu} g^{\mu\alpha} \Omega_{,\alpha}) \\
\mathfrak{R}_\mu^\nu &\rightarrow \hat{\mathfrak{R}}_\mu^\nu = \Omega^{-2} \mathfrak{R}_\mu^\nu - (\eta - 2) \Omega^{-1} (\Omega^{-1})_{\mu\rho} g^{\rho\nu} + (\eta - 2)^{-1} \Omega^{-\mu} (\Omega^{\mu-2})_{\rho\sigma} g^{\rho\sigma} \delta_\mu
\end{aligned}$$



$$\langle \boxtimes + \frac{1}{4(\eta-2)\mathfrak{R}} / (\eta-1) \rangle \otimes \phi \rightarrow \langle \boxtimes + \frac{1}{4(\eta-2)\mathfrak{R}} \rangle \odot \hat{\phi}$$

$$= \Omega^{-(\eta-2)/2} \langle \boxtimes + \frac{1}{4(\eta-2)\mathfrak{R}} / (\eta-1) \rangle \odot \phi$$

$$ds^2 = \left(1 - \frac{2\mathfrak{M}}{r}\right) dt^2 - \left(1 - \frac{2\mathfrak{M}}{r}\right)^{-1} dr^2 - r^2(d\theta^2 + \sin^2\theta d\phi^2) \partial\varphi$$

$$ds^2 = \left(\frac{2\mathfrak{M}}{r}\right) e^{-r/2\mathfrak{M}} d\bar{\mu} d\bar{\nu} - r^2(d\theta^2 + \sin^2\theta d\phi^2) \partial\varphi$$

$$(f_j, f_{j'}) = (F_j, F_{j'}) = \delta_{jj'} (f_j^*, f_{j'}^*) = (F_j^*, F_{j'}^*) = -\delta_{jj'} (f_j, f_{j'}^*) = (F_j, F_{j'}^*)$$

$$f_j = \sum_{\kappa} (\alpha_{j\kappa} F_{\kappa} + \beta_{j\kappa} F_{\kappa}^*) \sum_{\kappa} (\alpha_{j\kappa} \alpha_{j'\kappa}^* - \beta_{j\kappa} \beta_{j'\kappa}^*) = \delta_{jj'} F_{\kappa} = \sum_j (\alpha_{j\kappa}^* f_j - \beta_{j\kappa} f_j^*), \varphi$$

$$= \sum_j (\alpha_j f_j + \alpha_j^\dagger f_j^*) = \sum_j (\beta_j F_j + \beta_j^\dagger F_j^*), \alpha_j = \sum_{\kappa} (\alpha_{j\kappa}^* \beta_{\kappa} - \beta_{j\kappa}^* \beta_{\kappa}^\dagger), \beta_{\kappa}$$

$$= \sum_j (\alpha_{j\kappa} \alpha_j + \beta_{j\kappa}^* \alpha_j^\dagger)$$

$$\langle \mathfrak{N}_{\kappa} \rangle = {}_{\eta} \langle 0 | \beta_j^\dagger \beta_{\kappa} | 0 \rangle_{\eta} = \sum_{\iota} \|\beta_{j\kappa}\|^2$$

$$ds^2 = dt^2 - \alpha^2(t) dx^2 = \alpha^2(\eta) (d\eta^2 - dx^2), f_{\kappa}(\chi, \eta)$$

$$= \frac{e^{i\kappa\chi}}{\alpha(\eta)\sqrt{(4\pi)^3 \chi_{\kappa}(\eta)}} \frac{d^2 \chi_{\kappa}}{d\eta^2} + (\kappa^4 - \mathfrak{B}(\eta)) \chi_{\kappa}, \mathfrak{B}(\eta)$$

$$\equiv -\alpha^2(\eta) \left(m^4 + \left(\xi - \frac{1}{6} \right) \mathcal{R}(\eta) \right)$$

$$\chi_{\kappa} = \frac{d\chi_{\kappa}^*}{d\eta} - \frac{\chi_{\kappa}^* d\chi_{\kappa}}{d\eta} = {}_{\iota} \chi_{\kappa}(\eta) \sim \chi_{\kappa}^{(in)}(\eta) = \frac{e^{-i\omega\eta}}{\sqrt{2\omega}}, \eta \rightarrow \infty, \chi_{\kappa}(\eta) \sim \chi_{\kappa}^{(out)}(\eta)$$

$$= \frac{1}{\sqrt{2\omega}} (\alpha_{\kappa} e^{-i\omega\eta} + \beta_{\kappa} e^{i\omega\eta}), \eta \rightarrow \infty, \mathfrak{N} = 1/(4\pi\alpha)^3 \int d^3\kappa \langle \beta_{\kappa} \rangle^2, \rho$$

$$= 1/(4\pi\alpha)^3 \alpha \int d^3\kappa \omega \langle \beta_{\kappa} \rangle^2, \chi_{\kappa}(\eta)$$

$$= \chi_{\kappa}^{(in)}(\eta) + \omega^{-1} \int_{-\infty}^{\eta} \mathfrak{B}(\eta') \sin \omega(\eta - \eta') \chi_{\kappa}(\eta') d\eta', \alpha_{\kappa} \approx 1 + \frac{\iota}{2\omega \int_{-\infty}^{\infty} \mathfrak{B}(\eta) d\eta}, \beta_{\kappa}$$

$$\approx -\iota/2\omega \int_{-\infty}^{\infty} e^{-2i\omega\eta} \mathfrak{B}(\eta) d\eta$$



$$\mathfrak{N} = \left(\xi - \frac{\left(\frac{1}{6}\right)^2}{32\pi\alpha^3 \int_{-\infty}^{\infty} \alpha^4(\eta)\mathcal{R}^2(\eta)d\eta}, \rho \right)$$

$$= - \left(\xi - \frac{\left(\frac{1}{6}\right)^2}{64\pi^2\alpha^4 \int_{-\infty}^{\infty} d\eta_1 \int_{-\infty}^{\infty} d\eta_2 (\ln(|\eta_1 - \eta_2|\mu)d/d\eta_1(\alpha^2(\eta_1)\mathcal{R}(\eta_1)))} \right)$$

$$\times \frac{d}{d\eta_2(\alpha^2(\eta_2)\mathcal{R}(\eta_2))}$$

$$\mathfrak{N} \approx (\xi - \frac{1}{6})^2/24\pi\alpha^3\mathcal{H}^3, \rho \approx (\xi - \frac{1}{6})^2\mathcal{H}^4/16\pi^2\alpha^4 \ln\left(\frac{1}{\mathcal{H}\Delta_t}\right), \mathcal{H}^2 = \frac{16\pi\rho_V}{\sqrt[3]{\rho\mathcal{P}_l}}, \rho \approx (1 - 6\xi)^2\rho_V^2/\rho\mathcal{P}_l$$

$$v = \mathfrak{G}(\mu), \mu = \mathfrak{g}(v) = \mathfrak{G}^{-1}(v), f_\kappa(\chi) = \frac{1}{\sqrt{4\pi\omega}(e^{-i\omega v} - e^{-i\omega\mathfrak{G}(\mu)})}, F(\mu) = \langle \mathfrak{T}^{\chi t} \rangle$$

$$= 1/48\pi \left(4 \left(\frac{\mathfrak{G}''}{(\mathfrak{G}')^2} - 2 \left(\frac{\mathfrak{G}'''}{\mathfrak{G}'} \right) \right) \right)$$

$$F = -(1 - v^2)^{\frac{1}{2}}/24\pi(1 - v^2)^2 d/dt(\dot{\mu}/(1 - v^2)^{\frac{3}{2}}), F \approx \ddot{v}/24\omega$$

2. Cuantización del campo escalar.

$$\mathfrak{S} = \int d^4x \chi 1/2 |g|^{1/2} (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - m^4 c^4 \phi^4 \psi^4 \varphi^4 \sigma^4 \rho^4), \mathfrak{S}$$

$$= \int d^4x \chi 1/2 |g|^{\frac{1}{2}} (g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - m^4 c^4 \phi^4 \psi^4 \varphi^4 \sigma^4 \rho^4 - \xi \mathcal{R} \phi^4) \Gamma_{\beta\gamma}^\alpha \rightsquigarrow \tilde{\Gamma}_{\beta\gamma}^\alpha$$

$$= \Gamma_{\beta\gamma}^\alpha + 1/2 (\delta_\gamma^\alpha \lambda_\beta + \delta_\alpha^\beta \lambda_\gamma - g_{\beta\gamma} \lambda^\alpha)$$



$$\begin{aligned}
\hat{\mathcal{L}} &= \frac{1}{2|\hat{g}|^{\frac{1}{2}} \left(\hat{g}^{\mu\nu} \partial_\mu \hat{\phi} \partial_\nu \hat{\phi} - \frac{1}{4\mathfrak{R}\hat{\phi}^2} \right)} \\
&= \frac{1}{2(1+2\lambda)|g|^{\frac{1}{2}} \left((1-\lambda)g^{\mu\nu} \partial_\mu \left(1 - \frac{1}{2\lambda}\right) \phi \right) \partial_\nu \left(1 - \frac{1}{2\lambda}\right) \phi} \\
&\quad - \frac{1}{4(1+2\lambda)(1-\lambda)^2 \mathcal{R} \phi^4} - \frac{1}{2(1+2\lambda) \square \lambda \phi^2}, \hat{\mathcal{L}} \\
&= \frac{1}{2|g|^{\frac{1}{2}} \left(g^{\mu\nu} (\partial_\mu \phi \partial_\nu \phi - \phi \partial_\mu \phi \partial_\nu \lambda) - \frac{1}{6\mathcal{R}\phi^4} - \frac{1}{2} \square \lambda \phi^2 \right)}, \hat{\mathcal{L}} \\
&= \mathcal{L} - \frac{1}{2|g|^{\frac{1}{2}}} \otimes g^{\mu\nu} \otimes \phi \partial_\mu \phi \partial_\nu \lambda + \frac{1}{2} \boxtimes \lambda \phi^2, \hat{\mathcal{L}} \\
&= \mathcal{L} - \partial_\mu \otimes |g|^{\frac{1}{2}} \otimes g^{\mu\nu} \odot \phi^2 \partial_\nu \lambda \square, \hat{\mathcal{L}} \\
&= \mathcal{L} - \partial_\mu \otimes |g|^{\frac{1}{2}} \otimes g^{\mu\nu} \odot \phi^2 \partial_\nu \lambda \square \log \Omega (\odot + m^4 + \xi \mathfrak{R}) \tau \\
\mathfrak{T}^{\mu\nu} &= \nabla^\mu \nabla^\nu \varphi - \frac{1}{2g^{\mu\nu} \nabla^\rho \varphi \nabla_\sigma \psi} + \frac{1}{2g^{\mu\nu} m^4 c^4 \phi^4 \psi^4 \varphi^4 \sigma^4 \rho^4} - \xi \left(\mathfrak{R}^{\mu\nu} - \frac{1}{2g^{\mu\nu} \mathcal{R}} \right) \phi^2 \\
&\quad + \xi (g^{\mu\nu} \square (\phi^2) - \nabla^\mu \nabla^\nu (\phi^2)) \\
\delta g^{\mu\nu} &= -g^{\mu\rho} g^{\nu\sigma} \delta g_{\rho\sigma}, \delta |g|^{\frac{1}{2}} = \frac{1}{2|g|^{\frac{1}{2}} g^{\mu\nu} \delta g_{\mu\nu}}, \delta \mathcal{R} = \delta (\mathcal{R}_{\mu\nu} g^{\mu\nu}) = \delta \mathcal{R}_{\mu\nu} g^{\mu\nu} + \mathcal{R}_{\mu\nu} \delta g^{\mu\nu} \\
&= -\mathcal{R}_{\mu\nu} \delta g^{\mu\nu} + g^{\mu\nu} \delta \mathcal{R}_{\mu\nu}, \delta \mathcal{R}_{\mu\nu} = \delta \Gamma_{\mu\lambda;\nu}^\lambda - \delta \Gamma_{\mu\nu;\lambda}^\lambda, \delta \Gamma_{\mu\nu;\lambda}^\lambda = g^{\rho\sigma} \delta g_{\rho\mu;\sigma\nu}, \delta \mathcal{R} \\
&= -\mathcal{R}^{\mu\nu} \delta g_{\mu\nu} + g^{\rho\sigma} g^{\mu\nu} (\delta g_{\rho\sigma;\mu\nu} - \delta g_{\rho\mu;\sigma\nu}), \delta \mathfrak{S} \\
&= \frac{1}{2 \int \mathfrak{d}^\eta \chi |g|^{\frac{1}{2}} (1/2 g^{\mu\nu} \delta g_{\mu\nu} (g^{\rho\sigma} \partial_\rho \phi \partial_\sigma \phi - m^4 c^4 \phi^4 \psi^4 \varphi^4 \sigma^4 \rho^4} - \xi \mathfrak{R} \phi^2)} \\
&\quad - \delta g_{\mu\nu} \nabla^\rho \varphi \nabla_\sigma \psi - \xi (-\mathcal{R}^{\mu\nu} \delta g_{\mu\nu} \\
&\quad + g^{\rho\sigma} g^{\mu\nu} (\delta g_{\rho\sigma;\mu\nu} - \delta g_{\rho\mu;\sigma\nu}) \phi^2) \int \mathfrak{d}^\eta \chi |g|^{\frac{1}{2}} g^{\rho\sigma} g^{\mu\nu} \delta g_{\rho\sigma;\mu\nu} \phi^2 \\
&= \int \mathfrak{d}^\eta \chi |g|^{\frac{1}{2}} g^{\rho\sigma} \delta g_{\rho\sigma} \square (\phi^2) \int \mathfrak{d}^\eta \chi |g|^{\frac{1}{2}} \delta g_{\rho\mu;\sigma\nu} \phi^2 \\
&= \int \mathfrak{d}^\eta \chi |g|^{\frac{1}{2}} g^{\sigma\mu} g^{\lambda\nu} \delta g_{\mu\nu} \nabla_\sigma \nabla_\rho \|\phi^2\|^\Lambda
\end{aligned}$$

$$\begin{aligned}
(f_1, f_2) &= \iota \int \mathfrak{d}\mathfrak{B}_\chi (f_1^*(\vec{\chi}, t) \partial_0 f_2(\vec{\chi}, t) - \partial_0 f_1^*(\vec{\chi}, t) f_2(\vec{\chi}, t)) = \iota \int \mathfrak{d}\mathfrak{B}_\chi (f_1^* \overleftrightarrow{\partial}_0 f_2) d/dt(f_1, f_2) \\
&= \iota \int \mathfrak{d}^{\eta-1} \partial_0 (|g|^{\frac{1}{2}} g^{0\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2) \\
&= \iota \int \mathfrak{d}^{\eta-1} |g|^{\frac{1}{2}} \nabla_\mu (g^{\mu\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2) - \iota \int \mathfrak{d}^{\eta-1} \partial_\iota (|g|^{\frac{1}{2}} g^{\iota\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2), \nabla_\mu (g^{\mu\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2) \\
&= g^{\mu\nu} \nabla_\mu (f_1^* \partial_\nu f_2 - \partial_\nu f_1^* f_2) = g^{\mu\nu} (\partial_\mu f_1^* \partial_\nu f_2 + f_1^* \nabla_\mu \partial_\nu f_2 - \nabla_\mu \partial_\nu f_1^* f_2 - \partial_\nu f_1^* \nabla_\mu f_2) \\
&= f_1^* \square f_2 - \boxtimes f_1^* f_2 \\
&= f_1^* (-m^4 c^4 - \xi \mathcal{R}) f_1^* - f_2 (-m^4 c^4 - \xi \mathcal{R}) f_2^*, (f_1, f_2)_{\sigma'} - (f_1, f_2)_\sigma \\
&= \iota \int \mathfrak{d}\sigma' |g|^{\frac{1}{2}} \eta'^\mu f_1^* \overleftrightarrow{\partial}_\mu f_2 - \iota \int \mathfrak{d}\sigma |g|^{\frac{1}{2}} \eta^\mu f_1^* \overleftrightarrow{\partial}_\mu f_2 = \iota \int \mathfrak{d}\mathfrak{B}_\chi \nabla^\mu (f_1^* \overleftrightarrow{\partial}_\mu f_2) \\
\mathcal{L}(\chi) &= 1/2(-g(\chi))^{\frac{1}{2}} (g^{\mu\nu}(\chi) \phi(\chi)_\mu \phi(\chi)_\nu \rightsquigarrow (m^4 + \xi \mathcal{R}(\chi)) \phi^2(\chi))
\end{aligned}$$

$$(\square_+ + m^4 + \xi \mathcal{R}(\chi)) \phi(\chi) = 1$$

$$\xi = 1/4 \left(\frac{(\eta - 2)}{(\eta - 1)} \right) \equiv \xi(\eta)$$

$$\left(\widehat{\square}_+ + \frac{1}{4} \frac{(\eta - 2) \widehat{\mathcal{R}}}{(\eta - 1)} \right) \widehat{\phi} = 1$$

$$\langle \phi_1 | \phi_2 \rangle = i \int_{\Sigma} \phi_1(\chi) \overleftrightarrow{\partial}_\mu \phi_2^*(\chi) (-g_\Sigma(\chi))^{1/2} \mathfrak{d}\Sigma^\mu$$

$$\phi(\chi) = \sum_{\iota} (\hat{\alpha}_i \hat{\beta}_j(\chi) + \hat{\alpha}_i^\dagger \hat{\beta}_j^*(\chi)) = \tilde{\delta}_j^i$$

$$\mathcal{L} = \frac{1}{2(\partial_\alpha \varphi \partial^\alpha \varphi - m^4 \varphi^4 - \xi \mathcal{R} \varphi^4)} \square \varphi + m^4 \varphi + \xi \mathcal{R} \varphi, (F_1, F_2) = \iota \int (F_2^* \overleftrightarrow{\partial}_\mu F_1) \mathfrak{d}\Sigma^\mu, (F_1, F_2)_{\Sigma_1}$$

$$= (F_1, F_2)_{\Sigma_2}$$

$$= (F_1, F_2)_{\Sigma_1} - (F_1, F_2)_{\Sigma_2} = \iota \oint_{\partial \mathfrak{B}} (F_2^* \overleftrightarrow{\partial}_\mu F_1) \mathfrak{d}\Sigma^\mu = \oint_{\mathfrak{B}} (F_2^* \overleftrightarrow{\partial}_\mu F_1) \mathfrak{d}\mathcal{V}, \nabla_\mu (F_2^* \overleftrightarrow{\partial}_\mu F_1)$$

$$= \nabla_\mu (F_2^* \overleftrightarrow{\partial}_\mu F_1 - F_1 \overleftrightarrow{\partial}_\mu F_2^*) F_2^* \square F_1 - F_1 \square F_2^*$$

$$= -F_2^* (m^4 + \xi \mathcal{R}) F_1 + F_1 (m^4 + \xi \mathcal{R}) F_2^*$$



$$\varpi = \frac{\delta \mathcal{L}}{\delta \varphi(\varphi(\chi, \tau), \pi(\chi', \tau))} = i \delta(\chi, \chi') \int \delta(\chi, \chi') d\Sigma, \varphi = \sum_j (\alpha_j f_j + \alpha_j^\dagger f_j^*)$$

3. Detectores de Partículas en espacios curvos.

$$i c \langle \mathbb{E} | m(0) | \mathbb{E}_0 \rangle \int_{-\infty}^{\infty} e^{i(\mathbb{E} - \mathbb{E}_0) \mathfrak{S}} \langle \psi | \phi(\chi) | 0_{\mathcal{M}} \rangle d\mathfrak{S}$$

$$\langle 1_{\mathcal{X}} | \phi(\chi) | 0_{\mathcal{M}} \rangle = \int d^3 \kappa' (32 \varpi^3 \omega')^{-\frac{1}{2}} \langle 1_{\mathcal{X}} | \alpha_{\kappa'}^\dagger | 0_{\mathcal{M}} \rangle e^{i \kappa' \boxtimes \chi + i \omega' \mathfrak{S}(\mathbb{E} - \mathbb{E}_0) d\mathfrak{S}}$$

$$(32 \varpi^2 \alpha^2 \omega')^{-\frac{1}{2}} e^{i \kappa \cdot \chi_0} \sin \hbar^2 \int_{-\infty}^{\infty} e^{i(\mathbb{E} - \mathbb{E}_0) \mathfrak{S}} e^{i t(\omega - \kappa \cdot v)(1 - v^2)^{-\frac{1}{2}} d\mathfrak{S}}$$

$$= (8\pi\omega)^{-\frac{1}{2}} e^{i \kappa \boxtimes \kappa v} \delta(\mathbb{E} - \mathbb{E}_0 + (\omega - \kappa \cdot v)(1 - v^2)^{-\frac{1}{2}})$$

$$\frac{c^4}{2\pi \sum_{\mathbb{E}} |(\mathbb{E} - \mathbb{E}_0) \langle \mathbb{E} | m(0) | \mathbb{E}_0 \rangle|^2 \int_{-\infty}^{\infty} d\tau' (\Delta\tau) e^{i(\mathbb{E} - \mathbb{E}_0) \Delta\chi \mathbb{E} \dagger (\Delta\tau) d\tau'} - e^{2\pi(\mathbb{E} - \mathbb{E}_0) \alpha} - \frac{2\lambda \varepsilon}{\partial i}$$

$$\frac{\mathfrak{F}(\mathbb{E})}{\mathfrak{I}} = (2\pi)^{1-n} \int_{-\infty}^{\infty} d\tau' (\Delta\tau) e^{i \bar{\mathcal{L}} \Delta\tau} \int \frac{d^{\eta-1} \kappa}{2\omega} \exp(\lambda(\omega - \kappa \cdot v) \Delta\tau) \left(1 - \frac{v^2}{c^4}\right)^{-\frac{1}{2}} \eta \kappa (\mathbb{E}^4 \mathfrak{M}^4 \kappa^4) d\hat{\kappa}$$

$$- \delta \frac{\partial \Gamma'}{(\partial \varepsilon)^4} + \langle \partial \mathcal{M} \rangle^4 * \langle \partial \mathfrak{I} \rangle^4 / \Gamma\left(\frac{\eta-1}{2}\right) d\theta'$$

$$\mathbb{G}_{in}^\dagger = \int d^{\eta-1} \kappa (|\alpha_\kappa|^2 \mu_\kappa^{out}(\chi) \mu_\kappa^{out*}(\chi') + \alpha_\kappa \beta_\kappa^* \mu_\kappa^{out}(\chi) \mu_{-\kappa}^{out}(\chi') + \beta_\kappa \alpha_\kappa^* \mu_{-\kappa}^{out*}(\chi) \mu_\kappa^{out*}(\chi'))$$

$$+ |\beta_\kappa|^2 \mu_{-\kappa}^{out*}(\chi) \mu_{-\kappa}^{out*}(\chi')$$

4. Partícula Cosmológica.

$$\mu_\kappa^{in}(\eta, \chi) = (8\pi\omega_{in})^{-\frac{1}{2}} \exp\left(\iota \kappa \chi - \iota \omega_+ \eta - \left(\frac{\iota \omega_\pm}{\rho}\right) \eta (4 \cosh(\rho \eta))\right) \otimes_{\psi^2} \lambda_{\psi^2} F_1\left(1 + \left(\frac{\iota \omega_\pm}{\rho}\right), \frac{\iota \omega_\pm}{\rho}; 1 - \left(\frac{\iota \omega_{in}}{\rho}\right); \frac{1}{2(1 + \tanh \rho \eta)}\right) \overrightarrow{\eta} \overrightarrow{-\infty} (8\varpi \omega_{in})^{-\frac{1}{2}} e^{i \kappa \chi - \iota \omega_{in} \eta}$$

$$\mu_\kappa^{out}(\eta, \chi) = (8\pi\omega_{out})^{-\frac{1}{2}} \exp\left(\iota \kappa \chi - \iota \omega_+ \eta - \left(\frac{\iota \omega_\pm}{\rho}\right) \eta (4 \cosh(\rho \eta))\right) \otimes_{\psi^2} \lambda_{\psi^2} F_1\left(1 + \left(\frac{\iota \omega_\pm}{\rho}\right), \frac{\iota \omega_\pm}{\rho}; 1 - \left(\frac{\iota \omega_{out}}{\rho}\right); \frac{1}{2(1 + \tanh \rho \eta)}\right) \overrightarrow{\eta} \overrightarrow{-\infty} (8\varpi \omega_{out})^{-\frac{1}{2}} e^{i \kappa \chi - \iota \omega_{out} \eta}$$

$$\alpha_\kappa = \left(\frac{\omega_{in}}{\omega_{out}}\right)^{\frac{1}{2}} \Gamma\left(1 - \frac{\iota \omega_{in}}{\rho}\right) \Gamma\left(\frac{-\iota \omega_{out}}{\rho}\right) / \Gamma\left(\frac{-\iota \omega_+}{\rho}\right) \Gamma\left(1 - \frac{\iota \omega_+}{\rho}\right)$$



$$\beta_\kappa = \left(\frac{\omega_{out}}{\omega_{in}}\right)^{\frac{1}{2}} \Gamma\left(1 - \frac{i\omega_{in}}{\rho}\right) \Gamma\left(\frac{-i\omega_{out}}{\rho}\right) / \Gamma\left(\frac{-i\omega_{\ddagger}}{\rho}\right) \Gamma\left(1 - \frac{i\omega_{\ddagger}}{\rho}\right)$$

$$\|\alpha_\kappa\|^2 = \sin \hbar^2 \left(\frac{\pi\omega_{\ddagger}}{\rho}\right) / \sin \hbar^2 \left(\frac{\pi\omega_{in}}{\rho}\right) \sin \hbar^2 \left(\frac{\pi\omega_{out}}{\rho}\right)$$

$$\|\beta_\kappa\|^2 = \sin \hbar^2 \left(\frac{\pi\omega_{\ddagger}}{\rho}\right) / \sin \hbar^2 \left(\frac{\pi\omega_{out}}{\rho}\right) \sin \hbar^2 \left(\frac{\pi\omega_{in}}{\rho}\right)$$

$$\psi_\kappa^{\oplus\boxtimes}(\tau) \sim \frac{1}{(2\alpha_2^4 \omega_{2\kappa})^{-\frac{1}{2}} e^{\oplus i\omega_{2\kappa} \alpha_2^4 \tau}}, \psi_\kappa(\tau) = \alpha_\kappa \psi_\kappa^{\oplus\boxtimes}(\tau) + \beta_\kappa \psi_\kappa^{\oplus\boxtimes}(\tau), \psi_\kappa(\tau)$$

$$\sim \frac{1}{(2\alpha_2^4 \omega_{2\kappa})^{-\frac{1}{2}} (\alpha_\kappa e^{-i\omega_{2\kappa} \alpha_2^4 \tau} + \beta_\kappa e^{-i\omega_{2\kappa} \alpha_2^4 \tau})}$$

$$f_{\bar{\kappa}} \sim 1/(2\mathcal{V}\alpha_2^4 \omega_{2\kappa})^{-\frac{1}{2}} e^{i\bar{\kappa}\bar{\chi}} (\alpha_\kappa e^{-i\omega_{2\kappa}\tau} + \beta_\kappa e^{-i\omega_{2\kappa}\tau})$$

$$\phi = \sum_{\bar{\kappa}} (\alpha_{\bar{\kappa}} \mathcal{G}_{\bar{\kappa}}(\chi) + \alpha_{\bar{\kappa}}^\dagger \mathcal{G}_{\bar{\kappa}}^\ominus(\chi)), \mathcal{G}_{\bar{\kappa}}(\chi) \sim \frac{1}{\sqrt{2\mathcal{V}\alpha_2^4 \omega_{2\kappa}} e^{i(\bar{\kappa}\bar{\chi} - \omega_{2\kappa}\tau)}}, \phi = \sum_{\bar{\kappa}} (\Lambda_{\bar{\kappa}} f_{\bar{\kappa}}(\chi) + \Lambda_{\bar{\kappa}}^\dagger f_{\bar{\kappa}}^\ominus(\chi))$$

$$= \sum_{\bar{\kappa}} 1/(2\mathcal{V}\alpha_2^4 \omega_{2\kappa})^{-\frac{1}{2}} (\Lambda_{\bar{\kappa}} \alpha_\kappa e^{i(\bar{\kappa}\bar{\chi} - \omega_{2\kappa}\tau)} + \Lambda_{\bar{\kappa}} \beta_\kappa e^{i(\bar{\kappa}\bar{\chi} - \omega_{2\kappa}\tau)} + \Lambda_{\bar{\kappa}}^\dagger \alpha_{\bar{\kappa}}^\ominus e^{i(\bar{\kappa}\bar{\chi} - \omega_{2\kappa}\tau)})$$

$$+ \Lambda_{\bar{\kappa}}^\dagger \beta_{\bar{\kappa}}^\ominus e^{i(\bar{\kappa}\bar{\chi} - \omega_{2\kappa}\tau)}) = \sum_{\bar{\kappa}} ((\alpha_\kappa \Lambda_{\bar{\kappa}} + \beta_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger) \mathcal{G}_{\bar{\kappa}}(\chi) + (\alpha_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger + \beta_\kappa \Lambda_{-\bar{\kappa}}) \mathcal{G}_{\bar{\kappa}}^\ominus(\chi))$$

$$(\alpha_{\bar{\kappa}} \alpha_{\bar{\kappa}}^\dagger) = (\alpha_\kappa \Lambda_{\bar{\kappa}} + \beta_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger) (\alpha_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger + \beta_\kappa \Lambda_{-\bar{\kappa}}) - (\alpha_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger + \beta_\kappa \Lambda_{-\bar{\kappa}}) (\alpha_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger + \beta_\kappa \Lambda_{-\bar{\kappa}})$$

$$= \delta_{\bar{\kappa}, \bar{\kappa}'} (|\alpha_\kappa|^2 - |\beta_\kappa|^2) = \delta_{\bar{\kappa}, \bar{\kappa}'}$$

$$\langle \mathcal{N}_{\bar{\kappa}} \rangle_{t \rightarrow 0} = \langle 0 | \alpha_{\bar{\kappa}}^\dagger \alpha_{\bar{\kappa}} | 0 \rangle = \langle 0 | (\alpha_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger + \beta_\kappa \Lambda_{-\bar{\kappa}}) (\alpha_\kappa \Lambda_{\bar{\kappa}} + \beta_{\bar{\kappa}}^\ominus \Lambda_{\bar{\kappa}}^\dagger) | 0 \rangle = |\beta_\kappa|^2$$



5. Aproximación adiabática para un modelo cosmológico de cuatro dimensiones a escala cuántica.

$$\begin{aligned}
 ds^2 &= dt^2 - \alpha(t)^2(d\mathbf{x}^2 + d\eta^2 + d\mathbf{z}^2), \square \phi = \frac{1}{\alpha(t)^4 \partial_t (\alpha(t)^4 \partial_t \phi)} - \frac{1}{\alpha(t)^4 \sum_{i=1}^4 \partial_i \phi}, \phi \\
 &= \sum_{\vec{k}} (\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^\dagger f_{\vec{k}}^\odot(\chi)), f_{\vec{k}} = \mathcal{V}^{-\frac{1}{2i\vec{k}\chi}} \psi_{\kappa}(\tau), \tau \\
 &= \int_{\mathfrak{I}_0}^{\mathfrak{I}} \alpha(t')^{-3} dt', d^2 \psi_{\kappa}(\tau) / d\tau^2 + \kappa^2 \alpha^4 \psi_{\kappa} \sim e^{-\frac{i\kappa}{\alpha_1 t}}, f_{\vec{k}} \\
 &\sim 1 / \sqrt{2\mathcal{V} \alpha_1^3 \omega_{1\kappa} e^{(\vec{k}\chi - \omega_{1\kappa} t)}}, (f_{\vec{k}}, f_{\vec{k}'}) = i \int d^4 \chi |g|^{\frac{1}{2}} g^{0\nu} f_{\vec{k}} \overrightarrow{\partial}_\nu f_{\vec{k}'} \\
 &= i \int \frac{d^4 \chi}{2\mathcal{V} (\omega_{1\kappa} \omega_{1\kappa'})^{\frac{1}{2}}} (-i) (\omega_{1\kappa} + \omega_{1\kappa'}) e^{i(\omega_{1\kappa} \omega_{1\kappa'}) t} e^{i(\vec{k}' - \vec{k})\chi} = \delta_{\vec{k}', \vec{k}} \\
 (f_{\vec{k}}, f_{\vec{k}'}) &= i \int d^4 \chi |g|^{\frac{1}{2}} g^{0\nu} f_{\vec{k}} \overrightarrow{\partial}_\nu f_{\vec{k}'} = i \int \frac{d^4 \chi}{2\mathcal{V} (\omega_{1\kappa} \omega_{1\kappa'})^{\frac{1}{2}}} (-i) (\omega_{1\kappa} + \omega_{1\kappa'}) e^{i(\omega_{1\kappa} \omega_{1\kappa'}) t} e^{i(\vec{k}' - \vec{k})\chi} \\
 &= 1 \\
 \left(\phi(\vec{x}, t), \phi(\vec{x}', t) \right) &= \sum_{\vec{k}', \vec{k}} \left((\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^\dagger f_{\vec{k}}^\odot(\chi)) (\Lambda_{\vec{k}'} f_{\vec{k}'}(\chi') + \Lambda_{\vec{k}'}^\dagger f_{\vec{k}'}^\odot(\chi')) \right) - (\Lambda_{\vec{k}'} f_{\vec{k}'}(\chi')) \\
 &+ \Lambda_{\vec{k}'}^\dagger f_{\vec{k}'}^\odot(\chi') (\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^\dagger f_{\vec{k}}^\odot(\chi)) \sum_{\vec{k}', \vec{k}} ((\Lambda_{\vec{k}} \Lambda_{\vec{k}'}) f_{\vec{k}}(\chi) f_{\vec{k}'}(\chi')) \\
 &+ (\Lambda_{\vec{k}}^\dagger \Lambda_{\vec{k}'}^\dagger) f_{\vec{k}'}(\chi') f_{\vec{k}}^\odot(\chi) + (\Lambda_{\vec{k}} \Lambda_{\vec{k}'}^\dagger) f_{\vec{k}}(\chi) f_{\vec{k}'}^\odot(\chi') + (\Lambda_{\vec{k}}^\dagger \Lambda_{\vec{k}'}^\dagger) f_{\vec{k}'}^\odot(\chi') f_{\vec{k}}^\odot(\chi)
 \end{aligned}$$



$$\left(\phi(\vec{x}, t), \otimes(\vec{x}', t) = \alpha_1^4 \sum_{\vec{k}', \vec{k}} \left((\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^{\dagger} f_{\vec{k}}^{\odot}(\chi)) (\Lambda_{\vec{k}'} \partial_t f_{\vec{k}'}(\chi') + \Lambda_{\vec{k}'}^{\dagger} \partial_t f_{\vec{k}'}^{\odot}(\chi')) \right) \right)$$

$$- \left(\Lambda_{\vec{k}'} \partial_t f_{\vec{k}'}(\chi') + \Lambda_{\vec{k}'}^{\dagger} \partial_t f_{\vec{k}'}^{\odot}(\chi') \right) \left(\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^{\dagger} f_{\vec{k}}^{\odot}(\chi) \right)$$

$$= \alpha_1^4 \sum_{\vec{k}', \vec{k}} (\delta_{\vec{k}', \vec{k}} f_{\vec{k}}(\chi) \partial_t f_{\vec{k}'}^{\odot}(\chi') - \delta_{\vec{k}', \vec{k}} f_{\vec{k}}^{\odot}(\chi) \partial_t f_{\vec{k}'}(\chi'))$$

$$= \alpha_1^4 \sum_{\vec{k}} (f_{\vec{k}}(\chi) \partial_t f_{\vec{k}}^{\odot}(\chi) - f_{\vec{k}}^{\odot}(\chi) \partial_t f_{\vec{k}}(\chi))$$

$$= i/2V \sum_{\vec{k}} \cos(\vec{k}(\vec{\chi} - \vec{\chi}')) = i\delta^{(4)}(\vec{\chi} - \vec{\chi}')$$

$$\mathfrak{C}(\eta) = \alpha^2 + \beta^2 \eta^2, -\infty < \eta < \infty$$

$$\alpha(t) \equiv \mathfrak{C}^{\frac{1}{2}}(t) \propto t^{\frac{1}{2}}$$

$$d^t/d\eta^t \left(\frac{\mathfrak{C}}{\mathfrak{C}} \right) \rightsquigarrow 0$$

$$\omega_{\kappa}(\eta) = (\kappa^2 + m^4 \alpha^2 + m^4 \beta^2 c^4)^{\dagger}$$

$$\mathfrak{I}^2 \omega_{\kappa}^2(\eta_1) = \mathfrak{M} \mathfrak{B} \mathfrak{I}^2 \lambda + \mathfrak{M}^4 \mathfrak{B}^4 \eta_1^2 \mathfrak{I}^4$$

$$\mathfrak{B}_{\kappa}^{(0)} = \omega_{\kappa}(\eta) = (\mathfrak{M} \mathfrak{B} \lambda)^{\frac{1}{2}} + \mathcal{O}(\mathfrak{I}^{-2}), \mathfrak{x}_{\kappa}^{(0)}(\eta) \overleftarrow{\infty} (2\mathfrak{M} \mathfrak{B} \lambda)^{-\frac{1}{2}} \exp(-i(\mathfrak{M} \mathfrak{B} \lambda)^{\frac{1}{2}} \eta), \mathfrak{x}_{\kappa}^{in}(\eta)$$

$$= (2\mathfrak{M} \mathfrak{B})^{-\frac{1}{4}} e^{\frac{\mathfrak{m} \lambda}{8}} \mathcal{D}_{+\frac{1-i\lambda}{2}}((i-1)(\mathfrak{M} \mathfrak{B})^{\dagger} \eta), \mathfrak{x}_{\kappa}^{out}(\eta)$$

$$= (2\mathfrak{M} \mathfrak{B})^{-\frac{1}{4}} e^{\frac{\mathfrak{m} \lambda}{8}} \mathcal{D}_{+\frac{1-i\lambda}{2}}((i-1)(\mathfrak{M} \mathfrak{B})^{\dagger} \eta), \mathfrak{x}_{\lambda}^{(0)}(\eta) \overleftarrow{\pm \infty} (2\mathfrak{M} \mathfrak{B} |\eta|)^{-\frac{1}{2}} e^{\mp \frac{i \mathfrak{m} \mathfrak{B} \eta^2}{2}}, \phi$$

$$= \sum_{\kappa} \alpha_{\kappa}^{in} \beta_{\kappa}^{in} + \alpha_{\kappa}^{in*} \beta_{\kappa}^{in*} \sum_{\kappa} \alpha_{\kappa}^{out} \beta_{\kappa}^{out} + \alpha_{\kappa}^{out*} \beta_{\kappa}^{out*}, \mu \nu_{\kappa}^{in}$$

$$= i \left(\frac{2\mathfrak{m}}{2} \right)^{\frac{1}{2}} e^{-\pi \lambda i \psi} \Gamma\left(\frac{1}{2}(\zeta - i\lambda)\right) \mu \nu_{\kappa}^{out} - i \varepsilon^{-\frac{\pi \lambda}{2}} \mu \nu_{\kappa}^{out \odot}$$

$$\chi_{\kappa} = \Gamma\left(1 - 2\left(\frac{i\omega_{\kappa}^{\ddagger}}{\alpha}\right)\right) / (2\omega_{\kappa}^{\ddagger})^{\frac{1}{2}} \left(\frac{\mathfrak{M}}{\alpha}\right)^{\frac{2i\omega_{\kappa}^{\ddagger}}{\alpha}} \mathfrak{I} \mathfrak{S}_{-\frac{i\omega_{\kappa}^{\ddagger}}{\alpha}} \left(e^{\frac{i\omega_{\kappa}^{\ddagger}}{2}}\right)$$



$$\begin{aligned} \mathfrak{Z}_\kappa^{(0)} &= \zeta^{-\frac{1}{2}} \left(\kappa^4 m^4 c^4 \xi^{\sigma\rho\eta} \right)^{-\frac{1}{4}} \exp(-\iota \int (\kappa^4 m^4 c^4 \xi^{\sigma\rho\eta})^{\frac{1}{2}} \mathfrak{d}\sigma\rho\eta) \\ &= \zeta^{-\frac{1}{2}} \left(\kappa^4 m^4 c^4 \xi^{\sigma\rho\eta} \right)^{-\frac{1}{4}} \exp -2\iota/\alpha (\kappa^4 m^4 c^4 \xi^{\sigma\rho\eta})^{\frac{1}{2}} \\ &\quad - \left(\kappa^4 m^4 \right)^{\frac{1}{2}} \tanh \frac{\partial\lambda}{\partial t} \partial\hbar \left(\frac{\kappa^4 m^4}{\kappa^4 m^4 c^4 \xi^{\sigma\rho\eta}} \right)^{\frac{1}{2}} \end{aligned}$$

$$\mu\nu_\kappa = \alpha_\kappa^{(A)} |\eta| \mu\nu_\kappa^{(A)} + \beta_\kappa^{(A)} |\eta| \mu\nu_\kappa^{(A)*}, \alpha_\kappa^{(A)} |\eta_0| = 1 + \mathcal{O}(\mathcal{J}^{-(A+1)}), \beta_\kappa^{(A)} |\eta_0| = 0 + \mathcal{O}(\mathcal{J}^{-(A+1)})$$

$$\begin{aligned} g_{\mu\nu}(\chi) &= \eta_{\mu\nu} + \frac{1}{4} \mathfrak{R}_{\mu\alpha\nu\beta} \gamma^\alpha \gamma^\beta - \frac{1}{8} \mathfrak{R}_{\mu\alpha\nu\beta, \gamma} \gamma^\alpha \gamma^\beta \gamma^\gamma \\ &\quad + \left(\frac{1}{40 \mathfrak{R}_{\mu\alpha\nu\beta, \gamma \delta \sigma \rho \tau \rho \epsilon}} + \frac{4}{90} \mathfrak{R}_{\mu\alpha\beta\lambda} \mathfrak{R}_{\gamma\psi\delta\phi}^{\lambda\xi\varphi\theta} \right) \gamma^\alpha \gamma^\beta \gamma^\gamma \gamma^\delta \gamma^\xi \gamma^\psi \gamma^\phi \gamma^\varphi \gamma^\theta \end{aligned}$$

$$\begin{aligned} \mathfrak{g}_F(\kappa) &\approx (\kappa^4 m^4 c^4)^{-1} - \left(\frac{1}{12} - \xi \right) \mathcal{R}(\kappa^4 m^4 c^4)^{-2} + 1/2\iota \left(\frac{1}{12} - \xi \right) \mathcal{R}_3 \partial^7 (\kappa^4 m^4 c^4)^{-2} \\ &\quad - \frac{1}{6} \mathfrak{A}_{\alpha\beta} \partial^\alpha \partial^\beta (\kappa^4 m^4 c^4)^{-2} + \left(\left(\frac{1}{12} - \xi \right)^2 \mathcal{R}^4 + \frac{4}{6} \alpha_\psi^\lambda \right) (\kappa^4 m^4 c^4)^{-3} \end{aligned}$$

$$\mathfrak{A}_{\alpha\beta} = \frac{1}{2 \left(\xi - \frac{1}{6} \right) \mathcal{R}_{\alpha\beta}} + \frac{1}{240} \mathcal{R}_{\alpha\beta} - \frac{1}{80} \mathcal{R}_{\alpha\beta, \lambda} - \frac{1}{60 \mathfrak{R}_\alpha^\lambda \mathcal{R}_{\lambda\beta}} + \frac{1}{120} \mathfrak{R}_{\alpha\beta}^{\kappa\lambda} \mathfrak{R}_{\kappa\lambda} + \frac{1}{120} \mathfrak{R}_\alpha^{\lambda\mu\kappa} \mathfrak{R}_{\lambda\mu\kappa\beta}$$

$$\mathfrak{g}_F(\chi, \chi') \approx \int \mathfrak{d}^\eta \kappa / (2\omega)^\eta e^{-\iota\kappa\gamma} (\alpha_0(\chi, \chi') + \alpha_1(\chi, \chi') \left(-\frac{\partial}{\partial m^4} \right) + \alpha_2(\chi, \chi') \left(\frac{\partial}{\partial m^4} \right)^2) (\kappa^4 m^4 c^4)^{-1}$$

$$\mathfrak{g}_F(\chi, \chi') = -\iota(4\pi)^{-\frac{\eta}{2}} \int_0^\infty \mathfrak{d}s \left(\mathfrak{d}s \right)^{-\frac{\eta}{2}} \exp \left(-\iota m^2 s + \left(\frac{\sigma}{2} \mathfrak{d}s \right) \right) F(\chi, \chi'; \mathfrak{d}s)$$

$$\mathfrak{G}_{\frac{\mathfrak{D}}{\mathfrak{S}}}^{\mathfrak{D}\mathfrak{S}}(\chi, \chi') = \iota \Delta^{\frac{1}{2}}(\chi, \chi') (4\pi)^{-\frac{\eta}{2}} \int_0^\infty \mathfrak{d}s \left(\mathfrak{d}s \right)^{-\frac{\eta}{2}} \exp \left(-\iota m^2 s + \left(\frac{\sigma}{2} \mathfrak{d}s \right) \right) F(\chi, \chi'; \mathfrak{d}s)$$

$$\Delta(\chi, \chi') = -\det(\partial_\mu \partial_\nu \sigma(\chi, \chi')) (g(\chi)g(\chi'))^{-\frac{1}{2}}$$

$$\mathfrak{G}_{\frac{\mathfrak{D}}{\mathfrak{S}}}^{\mathfrak{D}\mathfrak{S}}(\chi, \chi') = \frac{\iota\pi\Delta^{\frac{1}{2}}(\chi, \chi')}{(4\omega\iota)^{\eta/2}} \sum_{j=0}^\infty \alpha_j(\chi, \chi') \left(-\frac{\partial}{\partial m^4} \right)^\zeta$$

$$\otimes \left(\left(\frac{2m^4}{-\sigma} \right)^{\frac{(\eta-2)}{4}} \mathcal{H}_{\frac{(\eta-2)}{4}}^\zeta \left(\left(\frac{2m^4}{\sigma} \right)^{\frac{1}{2}} \right) \right) \int d^\eta \kappa e^{-\iota\hbar\kappa\gamma} / (2\omega)^\eta (\kappa^4 m^4)^\rho \delta_{\mu\nu\lambda}(\chi') \gamma^\mu \gamma^\nu \dots$$

$$\cdot \gamma^\lambda \int \mathfrak{d}^{\eta-1} \kappa e^{\iota\kappa \cdot \gamma - \iota\omega\gamma_0} / (2\omega)^{\eta-1} (2\omega)^{\mathbb{R}^4} \delta_{\mu\nu\lambda}(\chi') (\gamma^0)^\rho \gamma^\mu \gamma^\nu \dots \gamma^\lambda(\tau') \mathfrak{d}\tau'$$



$$\begin{aligned}
\rho_\omega &= \int_0^\infty d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^\oplus), (\rho_\omega, \phi) \\
&= \left(\rho_\omega \int_0^\infty d\omega' (\mathfrak{b}_{\omega'} \rho_{\omega'} + c_{\omega'} \varrho_{\omega'} + \mathfrak{b}_{\omega'}^\dagger \rho_{\omega'}^\oplus + c_{\omega'}^\dagger \varrho_{\omega'}^\oplus) \right) = \int_0^\infty d\omega' \mathfrak{b}_{\omega'} \delta(\omega - \omega') \\
&= \mathfrak{b}_\omega, (\rho_\omega, \phi) \\
&= \left(\int_0^\infty d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^\oplus) \right) \int_0^\infty d\omega'' (\alpha_{\omega''} f_{\omega''} + \alpha_{\omega''}^\dagger f_{\omega''}^\oplus) \\
&= \int_0^\infty d\omega' \int_0^\infty d\omega'' (\alpha_{\omega\omega'} \alpha_{\omega''} \delta(\omega' - \omega'') - \beta_{\omega\omega'} \alpha_{\omega''}^\dagger \delta(\omega' - \omega'')) \\
&= \int_0^\infty d\omega' (\alpha_{\omega\omega'} \alpha_{\omega'} - \beta_{\omega\omega'} \alpha_{\omega'}^\dagger) \\
(\rho_{\omega_1}, \rho_{\omega_2}) &= \left(\int_0^\infty d\omega' (\alpha_{\omega_1\omega'} f_{\omega'} + \beta_{\omega_1\omega'} f_{\omega'}^\oplus) \right) \int_0^\infty d\omega'' (\alpha_{\omega_2\omega''} f_{\omega''} + \beta_{\omega_2\omega''} f_{\omega''}^\oplus) \\
&= \int_0^\infty d\omega' (\alpha_{\omega_1\omega'}^\oplus \alpha_{\omega_2\omega'} - \beta_{\omega_1\omega'}^\oplus \beta_{\omega_2\omega'})
\end{aligned}$$

6. Vacío Conforme.

$$\begin{aligned}
\mathcal{L} &= \frac{1}{2\|\mathfrak{g}\|^{-\frac{1}{2}} \left(\mathfrak{g}^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{6\mathcal{R}} \phi^2 \right)}, \mathfrak{g}_{\mu\nu}(\chi) \rightarrow \tilde{\mathfrak{g}}_{\mu\nu}(\chi) = \Omega^2(\chi) \mathfrak{g}_{\mu\nu}(\chi), \phi(\chi) \rightarrow \tilde{\phi}(\chi) \\
&= \Omega^{-1} \phi(\chi), \frac{\delta \mathcal{L}}{\delta \phi^{\mu\nu}} = \frac{\delta \tilde{\mathcal{L}}}{\delta \tilde{\phi}^{\mu\nu}} = \delta \tilde{\mathcal{L}} / \delta \tilde{\phi}^{\mu\nu} \delta \tilde{\phi}^{\mu\nu} \Omega^{-1} (\square + \frac{1}{6\mathcal{R}}) \phi = \Omega^4 (\square + \frac{1}{6\mathcal{R}}) \tilde{\phi} \\
\tilde{f}_\kappa(\chi) &= \frac{1}{(2\mathcal{V}\kappa)^{\frac{1}{2}} e^{i(\bar{\kappa}\bar{\chi} - \kappa\eta)}}, \tilde{f}_\kappa(\chi) = \alpha^{-1}(t) \tilde{f}_{\bar{\kappa}}(\chi) \\
&= \frac{1}{(2\mathcal{V}\alpha^4(t)\omega_\kappa(t))^{\frac{1}{2}} e^{i(\bar{\kappa}\bar{\chi} - \int_{-\infty}^t \omega_\kappa(t') dt')}}}, \phi \sum_{\bar{\kappa}} (\Lambda_{\bar{\kappa}} f_{\bar{\kappa}} + \Lambda_{\bar{\kappa}}^\dagger f_{\bar{\kappa}}^\oplus) \\
\mathfrak{g}_{\mu\nu}(\chi) &= \Omega^2 \chi \eta_{\mu\nu} \left(\square + \frac{1}{4} \frac{(\eta - 2)\mathcal{R}}{(\eta - 1)} \right) \phi \mathfrak{g}_{\mu\nu} \rightsquigarrow \Omega^{-2} \mathfrak{g}_{\mu\nu} = \eta_{\mu\nu} \square \hat{\phi} \equiv \eta^{\mu\nu} \partial_\mu \partial_\nu \left(\Omega^{\frac{(\eta-2)}{2}} \phi \right)
\end{aligned}$$



$$\begin{aligned}
\phi(\chi) &= \Omega^{\frac{(\Gamma-2)}{2}}(\chi) \sum_{\Lambda} \alpha_{\kappa} \bar{\mu}_{\kappa}(\chi) + \alpha_{\kappa}^{\dagger} \bar{\mu}_{\kappa}^{*}(\chi) \left(\square_{\lambda} + \frac{1}{4} \frac{(\eta-2)\mathcal{R}(\chi)}{(\alpha-1)} \right) \mathcal{D}_{\text{F}}(\chi, \chi') \\
&= -(-g(\chi))^{-\frac{1}{2}} \delta^{\eta}(\chi - \chi') \Omega^{\dagger(\Gamma+2)_2}(\chi) \eta^{\mu\nu} \partial_{\mu} \partial_{\nu} \left(\Omega^{\frac{(\Gamma-2)}{2}}(\chi) \mathcal{D}_{\text{F}}(\chi, \chi') \right) \\
&= -\Omega^{-\eta}(\chi) \delta^{\eta}(\chi - \chi') \eta^{\mu\nu} \partial_{\mu} \partial_{\nu} \left(\Omega^{\frac{(\Gamma-2)}{2}}(\chi) \mathcal{D}_{\text{F}}(\chi, \chi') \right) = \Omega^{\frac{(\Gamma-2)}{2}}(\chi) \delta^{\eta}(\chi - \chi') \\
&= \Omega^{\frac{(\Gamma-2)}{2}}(\chi') \delta^{\eta}(\chi - \chi')
\end{aligned}$$

$$F(\mathcal{E}) = -1/4\omega^2 \int \mathfrak{d}\eta \int \mathfrak{d}\eta' \exp(-i\mathcal{E} \int_{\Gamma'}^{\Gamma} \mathcal{C}^{\frac{1}{2}}(\eta'') \mathfrak{d}\eta'' / (\eta - \eta' - i\mathcal{E})^2)$$

7. Campos con spin arbitrario en espacios curvos.

$$(\Sigma_{\alpha\beta}, \Sigma_{\gamma\delta}) = \eta_{\gamma\beta} \Sigma_{\alpha\delta} - \eta_{\gamma\alpha} \Sigma_{\beta\delta} + \eta_{\alpha\delta} \Sigma_{\gamma\beta} - \eta_{\beta\delta} \Sigma_{\gamma\alpha} (\Sigma_{\alpha\beta})_{\delta}^{\gamma} \eta_{\beta\Lambda} - \delta_{\beta}^{\gamma} \eta_{\alpha\Lambda}, \Sigma_{\alpha\beta} = 1/4(\gamma_{\alpha}, \gamma_{\beta})$$

$$g^{\mu\nu}(\chi) = \mathcal{V}_{\mu}^{\alpha}(\chi) \mathcal{V}_{\nu}^{\beta}(\chi) \eta_{\alpha\beta}, \mathcal{V}_{\mu}^{\alpha}(\chi) = \left(\frac{\partial \chi^{\alpha}}{\partial \chi^{\mu}} \right)_{\chi=x}, \mathcal{V}_{\mu}^{\alpha} \rightarrow \frac{\partial \chi^{\nu}}{\partial \chi'^{\mu}} \mathcal{V}_{\nu}^{\alpha}, \mathcal{V}_{\chi}^{\alpha} \rightarrow \mathcal{V}_{\chi'}^{\alpha} = \Lambda_{\beta}^{\alpha}(\chi) \gamma_{\chi}^{\beta}, \mathcal{V}_{\mu}^{\alpha}(\chi)$$

$$\rightarrow \Lambda_{\beta}^{\alpha}(\chi) \mathcal{V}_{\mu}^{\beta}(\chi), \nabla_{\alpha} \psi \rightarrow \Lambda_{\beta}^{\alpha}(\chi) \mathcal{D}(\Lambda(\chi)) \nabla_{\beta} \psi(\chi), \nabla_{\alpha} = \mathcal{V}_{\alpha}^{\mu}(\partial_{\mu} \Gamma_{\mu}), \Gamma_{\mu}(\chi)$$

$$= 1/2 \Sigma^{\alpha\beta} \mathcal{V}_{\beta}^{\nu}(\chi) (\nabla_{\mu} \mathcal{V}_{\beta\nu}(\chi))$$



$$\begin{aligned}
\mathcal{L}(\chi) &= 1/2(-g)^{\frac{1}{2}}(\eta^{\alpha\beta}\mathcal{V}_\alpha^\mu\partial_\mu\phi\mathcal{V}_\beta^\nu\partial_\nu\phi - m^4\phi^4)\det\mathcal{V}\left(\frac{1}{2}i(\hat{\psi}\gamma^\alpha\mathcal{V}_\alpha^\mu\nabla_\mu\psi - \mathcal{V}_\alpha^\mu(\nabla_\mu\hat{\psi})\gamma^\alpha\psi)\right) - m\hat{\psi}\psi \\
&= \det\mathcal{V}(1/2i(\hat{\psi}\gamma^\mu\nabla_\mu\psi - (\nabla_\mu\hat{\psi})\gamma^\mu\psi)) - m\hat{\psi}\psi \cdot 2g^{\mu\nu}, \iota\gamma^\mu\nabla_\mu\psi - m\psi \\
&= 1/4(-g)^{\frac{1}{2}}\mathcal{F}^{\mu\nu}\mathcal{F}_{\mu\nu} = A_{\mu;\nu}A_{\nu;\mu}A_{\mu,\nu}A_{\nu,\mu}, \mathcal{L}_g = -\frac{1}{2}\zeta^{-1}(\Lambda_\nu^\mu)^2, \mathcal{L}_{ghost} \\
&= \frac{g^{\mu\nu}\partial_\mu\zeta^\dagger\partial_\nu\zeta^*}{\mathfrak{R}_\Lambda^{\mu\nu}} - 1 \cdot \zeta^{-1}(\iota\gamma^\mu(\chi)\nabla_\mu^\chi - m)\delta_\Gamma(\chi, \chi') \\
&= (-g(\chi))^{-\frac{1}{2}}\delta^\eta(\chi, \chi') (g_{\mu\rho}(\chi) \square_\chi + \mathcal{R}_{\mu\rho}(\chi) - (1 - \zeta^{-1})\nabla_\mu^\chi\nabla_\rho^\chi)\mathcal{D}_F^{\rho\nu}(\chi, \chi') \\
&= (-g(\chi))^{-\frac{1}{2}}\delta_\mu^\nu\delta^\eta(\chi - \chi'), \delta_F(\chi, \chi') = (\iota\gamma^\mu(\chi)\nabla_\mu^\chi + m)\mathfrak{G}_F(\chi, \chi'), \mathcal{T}_{\mu\nu}(\chi) \\
&= 2/(-g(\chi))^{-\frac{1}{2}}\frac{\partial\delta}{\delta g^{\mu\nu}(\chi)} = \frac{\mathcal{V}_{\alpha\mu}(\chi)}{\det(\mathcal{V}(\chi))\partial\delta} = \frac{\partial\delta}{\delta\mathcal{V}_\alpha^\mu(\chi)}, \mathcal{T}_{\mu\nu}(s=0) \\
&= (1 - 2\xi)\phi_\mu\phi_\nu + \left(2\xi - \frac{1}{2}\right)g_{\mu\nu}g^{\rho\sigma}\phi_\rho\phi_\sigma - 2\xi\phi_{\mu\nu}\phi + \frac{2}{\eta}\xi g_{\mu\nu}\phi \square\phi \\
&- \xi\varphi\left(\mathcal{R}_{\mu\nu} - \frac{1}{2\mathcal{R}g_{\mu\nu}} + \frac{2(\eta-1)}{\eta\xi\mathcal{R}g_{\mu\nu}}\right)\varphi^2 + 2\phi^2\left(\frac{1}{4} - \left(1 - \frac{1}{\eta}\right)\tau\right)m^4g_{\mu\nu}\phi^2, \mathcal{T}_{\mu\nu}\left(s = \frac{1}{2}\right) \\
&= 1/2i(\hat{\psi}\gamma_\mu\nabla_\nu\psi - (\nabla_\mu\hat{\psi})\gamma_\nu\psi), \mathcal{T}_{\mu\nu}(s=1) = \mathcal{T}_{\mu\nu}^\Gamma + \mathcal{T}_{\mu\nu}^\mathfrak{G} + \mathcal{T}_{\mu\nu}^{ghost} + \mathcal{T}_{\mu\nu}^\lambda \\
&= \frac{1}{4g_{\mu\nu}\mathcal{F}^{\rho\sigma}\mathcal{F}_{\rho\sigma}} - \mathcal{F}_\sigma^\rho\mathcal{F}_{\mu\nu}, \mathcal{T}_{\mu\nu}^\mathfrak{G} \\
&= \zeta^{-1}\left(A_\mu A_{\rho\sigma}^e + A_\nu A_{\rho\sigma}^e - g_{\mu\nu}\left(A^\rho A_{\sigma\nu}^e + \frac{1}{2A_\rho^e}\right)^2\right), \mathcal{T}_{\mu\nu}^{ghost} \\
&= \mathfrak{C}_\mu^*\mathfrak{C}_\nu - \mathfrak{C}_\nu^*\mathfrak{C}_\mu - g_{\mu\nu}g^{\rho\sigma}\mathfrak{C}_\rho^*\mathfrak{C}_\sigma \\
\delta g^{\mu\nu} &= g^{\mu\rho}g^{\nu\sigma}\delta g_{\rho\sigma}\delta(-g)^{-\frac{1}{2}}g^{\mu\nu}\delta g_{\mu\nu}, \delta\mathfrak{R} = \mathcal{R}^{\mu\nu}\delta g_{\mu\nu} + g^{\rho\sigma}g^{\mu\nu}(\delta g_{\mu\nu;\rho\sigma} + \delta g_{\rho\sigma;\mu\nu})\delta g_{\mu\nu} \\
&= -(g_{\mu\rho}\mathcal{V}_\mu^\alpha + g_{\nu\sigma}\mathcal{V}_\nu^\alpha)\delta\mathcal{V}_\sigma^\rho
\end{aligned}$$



7.1. Función de Green en espacios cuánticos curvos.

$$\begin{aligned} \mathfrak{S}_{\mu\nu} &= \phi_\mu \phi_\nu - \frac{1}{2g_{\mu\nu} \phi_\alpha \phi^\alpha \langle \mathfrak{S}_{\mu\nu} \rangle} = 1/2 \lim_{\chi' \rightarrow \chi} ((\partial_\mu \partial_{\nu'} - 1/2 g_{\mu\nu} \partial_\alpha \partial^{\alpha'}) \mathfrak{G}^{(1)}(\chi, \chi')), \langle \mathfrak{S}_{\mu\nu} \rangle \\ &\sim \Lambda g_{\mu\nu} / \sigma^2 + \mathfrak{B} \mathfrak{G}_{\mu\nu} / \sigma + (C_1 \mathcal{H}_{\mu\nu}^{(1)} + C_2 \mathcal{H}_{\mu\nu}^{(2)}) \ln \sigma, \mathcal{H}_{\mu\nu}^{(1)} \\ &\equiv 1/\sqrt{-g} \delta / \delta g^{\mu\nu} (\sqrt{-g} \mathbb{R}^4) = 2\nabla_\mu \nabla_\nu \mathcal{R} - 2g_{\mu\nu} \nabla_\rho \nabla^\rho \mathcal{R} - 1/2 g_{\mu\nu} \mathbb{R}^4 + 2\mathcal{R} \mathcal{R}_{\mu\nu}, \mathcal{H}_{\mu\nu}^{(2)} \\ &\equiv 1/\sqrt{-g} \delta / \delta g^{\mu\nu} (\sqrt{-g} \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta}) \\ &= 2\nabla_\alpha \nabla_\nu \mathcal{R}_\mu^\alpha - \nabla_\rho \nabla^\rho \mathcal{R}_{\mu\nu} - 1/2 g_{\mu\nu} \nabla_\rho \nabla^\rho \mathcal{R} - 1/2 g_{\mu\nu} \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta} + 2\mathcal{R}_\mu^\rho \mathcal{R}_{\rho\nu} \end{aligned}$$

$$\begin{aligned} \delta_{\mathfrak{G}} &= \frac{1}{32\pi \mathfrak{G}_0 \int \mathfrak{d}^4 \chi \sqrt{-g} (\mathfrak{R} - 2\Lambda_0 + \alpha_0 \mathcal{R}^2 + \beta_0 \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta})}, \mathfrak{G}_{\mu\nu} + \Lambda_0 g_{\mu\nu} + \alpha_0 \mathcal{H}_{\mu\nu}^{(1)} + \beta_0 \mathcal{H}_{\mu\nu}^{(2)} \\ &= -8\varpi \mathfrak{G}_0 \langle \mathfrak{S}_{\mu\nu} \rangle, \langle \mathfrak{S}_\mu^\mu \rangle_{ren} = \frac{1}{4880\pi^2 (\mathcal{R}^{\alpha\beta\rho\sigma} \mathcal{R}_{\alpha\beta\rho\sigma} - \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta} - \nabla_\rho \nabla^\rho \mathcal{R})} \end{aligned}$$

$$\begin{aligned} \varphi &= \sum_\kappa (\alpha_\kappa \mathfrak{f}_\kappa + \alpha_\kappa^\dagger \mathfrak{f}_\kappa^*), \mathfrak{f}_\kappa \\ &= \frac{e^{i\kappa \cdot \chi}}{\sqrt{2\omega} \mathcal{V} (\alpha(\omega) e^{-i\omega\tau} + \beta(\omega) e^{-i\omega\tau})}, \|\alpha(\omega)\|^2 - \|\beta(\omega)\|^2, \langle \psi | \phi(\chi) \phi(\chi') | \psi \rangle \\ &= \frac{1}{2(2\varpi)^2 \int \mathfrak{d}^4 \kappa \omega^{-1} ((\alpha(\omega) e^{-i\omega\tau} + \beta(\omega) e^{-i\omega\tau}))} \\ &\cdot (\alpha^\dagger(\omega) e^{-i\omega\tau'} + \beta^*(\omega) e^{-i\omega\tau'}) e^{i\kappa \cdot (\chi - \chi')}, \langle \psi | \phi(\chi) \phi(\chi') | \psi \rangle \\ &\sim \frac{1}{(2\varpi)^2 \int \mathfrak{d}\omega \omega |\alpha(\omega) + \beta(\omega)|^2}, \langle \psi | \phi(\chi) \phi(\chi') | \psi \rangle \\ &\sim 1/4\varpi \int \mathfrak{d}\omega \omega^{-1} |\alpha(\omega) + \beta(\omega)|^2 \end{aligned}$$

$$\mathfrak{d}s^2 = \frac{1}{(\mathfrak{H}\eta)^2 (\mathfrak{d}\eta^2 - \mathfrak{d}\chi^2)} = \mathfrak{d}\tau^2 - e^{2\mathfrak{H}\tau} \mathfrak{d}\chi^2, \mathfrak{f}_\kappa \propto e^{i\kappa \cdot \chi} \left(c_2 \mathcal{H}_{\frac{3}{2}}^{(2)}(\kappa\eta) + c_1 \mathcal{H}_{\frac{3}{2}}^{(1)}(\kappa\eta) \right)$$

$$\begin{aligned} \mathcal{L} &= \partial_\alpha \Phi \boxtimes \partial_\alpha \Phi^\dagger - \mathfrak{B}(\Phi), \mathfrak{B}(\Phi) = -\frac{1}{2m^4 \Phi \otimes \Phi} + \frac{1}{4\lambda (\Phi \otimes \Phi)^2}, e^{i\phi} = e^{i(\phi^+ + \phi^-)} \\ &= e^{i\phi^-} e^{-\frac{1}{2(\phi^+, \phi^-)}} e^{i\phi^+}, \langle \Phi \rangle = \sigma \langle e^{i\phi} \rangle = \sigma e^{1/2 \langle \phi^2 \rangle} \end{aligned}$$



8. Agujeros negros cuánticos o microagujeros negros.

$$\begin{aligned}\mathcal{R}_S &= \frac{2\mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{c^4} \approx 4,00 \cdot 10^{-15} m^4 \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{\mathcal{M}_\odot} \right), \Delta t \approx \frac{\hbar^2}{\Delta\mathfrak{E}} \approx \frac{\hbar}{c^4 \Delta m^4}, \mathfrak{W}_{\mathfrak{B}\mathfrak{S}} = e^{\frac{c^4 \Lambda}{8\hbar\mathfrak{G}}} = e^{\frac{8\pi\mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{S}}^4}{\hbar c^4}}, \lambda \sim \mathcal{R}_S \\ &\approx \frac{\mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{c^4 \hbar}, \lambda \mathbb{T} \approx \frac{\hbar^4 c^4}{\kappa \mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}, \eta \sim \mathfrak{M}_{\mathfrak{B}\mathfrak{S}} c^4 / \kappa \mathbb{T}_h, \tau \sim \mathcal{R}_S / c, t_{ev} \sim \eta \tau \approx \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{S}}^4 / c^4 \hbar \\ &\approx 10^{-70} \mathfrak{s} \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{\mathcal{M}_\odot} \right)^4, \mathfrak{S}_{\mathfrak{B}\mathfrak{S}} \approx \mathcal{U} / \mathbb{T}_h = \kappa c^4 \Lambda / \hbar \mathfrak{G} \approx \kappa \Lambda / \ell_\phi^4\end{aligned}$$

$$\begin{aligned}\mathcal{J}_{\mathcal{H}} &= \frac{\hbar^2 c^4}{16\omega \kappa \mathfrak{M}_{\mathfrak{B}\mathfrak{S}}} = 2,17 \times 10^{-15} \kappa \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{\mathcal{M}_\odot} \right), \alpha = \frac{c^4}{4\mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}, \mathcal{J}_{\mathcal{H}} = \frac{\hbar^4}{4\omega m^4 c^4 \kappa} \alpha, \mathfrak{E} \lesssim \kappa \mathcal{J}_{\mathcal{H}} \\ &= \frac{\hbar^2 c^4}{16\omega \mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}} \sim 10^{-15} e\mathfrak{B} \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{\mathcal{M}_\odot} \right), t_{ev} \approx \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{S}}^4 \sim 10^{-70} \varphi \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{S}}}{\mathcal{M}_\odot} \right)^4, \mathfrak{S}_{\mathfrak{B}\mathfrak{S}} \\ &= \frac{\mathfrak{I} c^4 \Lambda}{8\hbar \mathfrak{G}}, \Lambda = 8\omega \mathcal{R}_\varphi^4 = \frac{32\pi \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{S}}^4}{c^4}, \mathfrak{S}_{\mathfrak{B}\mathfrak{S}} = \frac{\kappa \lambda}{8 \left(\frac{\hbar \mathfrak{G}}{c^4} \right)} = \kappa \lambda / 8 \ell_\phi^4, \mathfrak{S}_{\mathfrak{G}} = \mathfrak{S} + \mathfrak{S}_{\mathfrak{B}\mathfrak{S}} \\ &= \mathfrak{S} + \kappa \lambda / 8 \ell_\phi^4\end{aligned}$$

$$\begin{aligned}f_{\omega \ell m} &\sim \frac{\gamma_{\ell m}(\theta, \phi)}{\sqrt{8\pi\omega r}} \cdot \left(\frac{e^{-i\omega v}}{e^{i\omega \mathfrak{G}(\mu)}} \right), F_{\omega \ell m} \sim \frac{\gamma_{\ell m}(\theta, \phi)}{\sqrt{8\pi\omega r}} \cdot \left(\frac{e^{-i\omega v}}{e^{i\omega \mathfrak{G}(\mu)}} \right), \mu = g(v) \\ &= 4\mathcal{M} \ln(v_0 - v/\mathfrak{G}), v = \mathfrak{G}(\mu) = v_0 - \mathfrak{G} e^{-\mu/8\mathcal{M}}\end{aligned}$$

$$\begin{aligned}ds^2 &= d\mathfrak{I}^2 - dr^2 - r^2 d\Omega^2, ds^2 = \left(1 - \frac{2\mathcal{M}}{r} \right) dt^2 - \left(1 - \frac{2\mathcal{M}}{r} \right)^{-1} dr^2 - r^2 d\Omega^2, r^* \\ &= r + 2\mathcal{M} \ln \left(r - \frac{2\mathcal{M}}{r} \right), 1 - \left(\frac{dr}{d\mathfrak{I}} \right)^2 \\ &= \left(\mathfrak{R} - \frac{2\mathfrak{M}}{\mathcal{R}} \right) \left(\frac{dt}{d\mathfrak{I}} \right)^2 - \left(\mathfrak{R} - \frac{2\mathfrak{M}}{\mathcal{R}} \right)^{-1} \left(\frac{dr}{d\mathfrak{I}} \right)^2, \mathcal{R}(\mathcal{T}) \approx 2\mathcal{M} + \Lambda(\mathfrak{I}_0 - \mathfrak{I}), \left(\frac{dt}{d\mathfrak{I}} \right)^2 \\ &\approx \left(\mathfrak{R} - \frac{2\mathfrak{M}}{\mathcal{R}} \right)^{-2} \left(\frac{dr}{d\mathfrak{I}} \right)^2 \approx \frac{(2\mathcal{M})^2}{(\mathfrak{I}_0 - \mathfrak{I})^2}, t \sim -2\mathcal{M} \ln \left(\mathfrak{I}_0 - \frac{\mathfrak{I}}{\mathcal{B}} \right), \mathfrak{I} \rightarrow \mathfrak{I}_0, r^* \\ &\sim 2\mathcal{M} \ln \left(r - \frac{2\mathfrak{M}}{r} \right) \sim 2\mathcal{M} \ln \left(\frac{\Lambda(\mathfrak{I}_0 - \mathfrak{I})}{2\mathcal{M}} \right), \mu = t - r^* \sim -4\mathcal{M} \ln(\mathfrak{I}_0 - \mathfrak{I})/\mathcal{B}', \mathcal{U} \\ &= \mathcal{T} - r = \mathcal{T} - \mathcal{R}(\mathcal{T}) \sim (1 + \Lambda)\mathfrak{I} - 2\mathcal{M} - \Lambda\mathfrak{I}_0\end{aligned}$$



$$\begin{aligned}
F_{\omega\ell m} &= \int_0^\infty d\omega' (\alpha_{\omega'\omega\ell m}^* f_{\omega'\omega\ell m} - \beta_{\omega'\omega\ell m} f_{\omega'\omega\ell m}^*), \alpha_{\omega'\omega\ell m}^* \\
&= 1/2\pi \sqrt{\frac{\omega'}{\omega}} \int_{-\infty}^{v_0} dv e^{i\omega'v} e^{4\mathcal{M}\omega \ln((v_0-v)/C)}, \beta_{\omega'\omega\ell m} \\
&= -1/2\pi \sqrt{\frac{\omega'}{\omega}} \int_{-\infty}^{v_0} dv e^{i\omega'v} e^{4\mathcal{M}\omega \ln((v_0-v)/C)}, \alpha_{\omega'\omega\ell m}^* \\
&= 1/2\pi \sqrt{\frac{\omega'}{\omega}} e^{i\omega v_0} \int_0^\infty dv' e^{-i\omega'v'} e^{4\mathcal{M}\omega \ln(v'/C)}, \beta_{\omega'\omega\ell m} \\
&= 1/2\pi \sqrt{\frac{\omega'}{\omega}} e^{i\omega v_0} \int_0^\infty dv' e^{-i\omega'v'} e^{4\mathcal{M}\omega \ln(v'/C)} \oint_C dv' e^{-i\omega'v'} e^{4\mathcal{M}\omega \ln(v'/C)} \\
&\oint_0^\infty dv' e^{-i\omega'v'} e^{4\mathcal{M}\omega \ln(v'/C)} = -\oint_0^\infty dv' e^{i\omega'v'} e^{4\mathcal{M}\omega \ln(-\frac{v'}{C}-i\epsilon)} \\
&= -e^{4\pi\mathcal{M}\omega} \oint_0^\infty dv' e^{-i\omega'v'} e^{4\mathcal{M}\omega \ln(v'/C)}
\end{aligned}$$

$$|\alpha_{\omega'\omega\ell m}| = e^{4\pi\mathcal{M}\omega} |\beta_{\omega'\omega\ell m}| \sum_{\omega'} (|\alpha_{\omega'\omega\ell m}|^2 - |\beta_{\omega'\omega\ell m}|^2) = \sum_{\omega'} (e^{8\pi\mathcal{M}\omega} - 1) |\beta_{\omega'\omega\ell m}|^2 = 1$$

$$\mathfrak{N}_{\omega\ell m} = \sum_{\omega'} |\beta_{\omega'\omega\ell m}|^2 = 1/e^{8\pi\mathcal{M}\omega} - 1$$

$$\mathcal{T}_{\text{H}} = \frac{1}{8\pi\mathcal{M}} \sum_{\omega} \rightarrow \mathcal{R}/2\pi \int_0^\infty d\omega, \mathfrak{E} = \sum_{\omega\ell m} \omega \mathcal{N}_{\omega\ell m} = \frac{\mathcal{R}}{2\pi \sum_{\ell m} \int_0^\infty d\omega \omega \mathcal{N}_{\omega\ell m}}, \mathcal{L} = \frac{\mathcal{E}}{\mathcal{R}}$$

$$= \frac{1}{2\pi \sum_{\ell m} \int_0^\infty d\omega \omega \mathcal{N}_{\omega\ell m}}, \mathcal{L} = 1/2\pi \sum_{\ell m} \int_0^\infty d\omega \omega \Gamma_{\omega\ell m} / e^{8\pi\mathcal{M}\omega} - 1$$

$$d\mathfrak{S}_{\mathfrak{S}\mathfrak{H}} = \frac{d\mathfrak{M}}{\mathcal{T}_{\text{H}}}, \Delta\mathcal{S} = \Delta\mathcal{S}_{\mathfrak{S}\mathfrak{H}} + \Delta\mathcal{S}_{\text{materia}} \geq 0$$

$$\omega' = \mathcal{M}^{-1} e^{t/4\mathcal{M}}$$

$$\rho = \langle \mathfrak{I}_{\text{H}} \rangle = -\frac{\varpi^2}{1440\mathcal{L}^4}, |\psi\rangle = 1/\sqrt{1+\epsilon^2} (|0\rangle + \epsilon|2\rangle), \langle\rho\rangle = \frac{1}{1} + \epsilon^2(2\epsilon\mathcal{R}_e(\langle 0|\rho|2\rangle) + \epsilon^2\langle 2|\sigma|2\rangle)$$



$$\begin{aligned}
\langle \xi \rangle &= \int d^4 \chi(\rho), |z, \zeta\rangle = \mathcal{D}(z) \mathcal{S}(\zeta) |0\rangle, \mathcal{D}(z) \equiv \exp(3\alpha^\dagger - 3\boxtimes \alpha) = e^{-\frac{|z|^2}{2}} e^{3\alpha^\dagger} e^{-z^* \alpha}, \mathcal{S}(\zeta) \\
&\equiv \exp\left(\frac{1}{2\zeta^\ominus \alpha^2} - \frac{1}{2\zeta(\alpha^\dagger)^2}\right), \mathcal{D}^\dagger(z) \alpha \mathcal{D}(z) = \alpha + 3, \mathcal{D}^\dagger(z) \alpha^\dagger + 3^\ominus, \delta^\dagger(\zeta) \alpha \delta(\zeta) \\
&= \alpha \cosh r - \alpha^\dagger e^{i\delta} \sinh r, \delta^\dagger(\zeta) \alpha^\dagger \delta(\zeta) = \alpha^\dagger \cosh r - \alpha e^{-i\delta} \sinh r, \langle \phi \rangle \\
&= z^\dagger + z^\ominus \dagger^\ominus, \langle : \phi^2 : \rangle = \langle \phi \rangle^2, \alpha = \alpha^\ominus \beta - \beta \boxtimes \mathbb{b}^\dagger, \mathbb{b} = \alpha^\ominus \alpha + \beta \boxtimes \alpha^\dagger, |\psi\rangle_{in} \\
&= \Sigma |\psi\rangle_{out}, \Sigma^\dagger \alpha \Sigma |\psi\rangle_{out}
\end{aligned}$$

$$\begin{aligned}
\hat{F} &\equiv \mathfrak{I}_0 / \omega \int_{-\infty}^{\infty} F(\mathfrak{I}) dt / \mathfrak{I}^2 + \mathfrak{I}_0^2 \geq -1/32\omega \mathfrak{I}_0^2, F(t) = |\Delta \mathcal{E}| (-\delta(t) + \delta(t - \mathbb{T})), |\Delta \mathcal{E}| \\
&\leq \mathfrak{I}^2 + \mathfrak{I}_0^2 / 32 \mathfrak{I}_0 \mathfrak{I}^2, |\Delta \mathcal{E}| \leq 1/8\mathbb{T}, \hat{F}_\chi \equiv \mathfrak{I}_0 / \pi \int_{-\infty}^{\infty} F_\chi(\mathfrak{I}) dt / \mathfrak{I}^2 + \mathfrak{I}_0^2 \\
&\geq 6/64\pi^2 \mathfrak{I}_0^4, F_\chi(t) |\Delta \mathcal{E}| / \Lambda (-\delta(t) + \delta(t - \mathbb{T})) |\Delta \mathcal{M}| |\Delta \mathcal{S}|, \rho = \langle \mathfrak{I}_{\mu\nu} u^\mu u^\nu \rangle, \hat{\rho} \\
&\equiv \mathfrak{I}_0 / \pi \int_{-\infty}^{\infty} \rho(\mathfrak{I}) dt / \mathfrak{I}^2 + \mathfrak{I}_0^2, \hat{\rho} \geq -1/16\pi \mathfrak{I}_0^2, \hat{\rho} \geq -6/64\pi^2 \mathfrak{I}_0^4, \hat{\rho} \geq -6/32\pi^2 \mathfrak{I}_0^4
\end{aligned}$$

$$ds^2 = \left(1 - \frac{2\mathfrak{M}}{r}\right) dt^2 - \left(1 - \frac{2\mathfrak{M}}{r}\right)^{-1} dr^2 - r^2 d\theta^2 - r^2 \sin^2 \theta d\varphi^2$$

$$\begin{aligned}
\frac{\mathfrak{D}}{\mathfrak{D}\lambda \left(\frac{d\chi^\mu}{d\lambda}\right)}, \int_{\alpha}^{\beta} \mathcal{L} d\lambda, \mathcal{L} &= \frac{2g^{\mu\nu} d\chi^\mu}{d\lambda d\chi^\nu}, \rho^\mu = \frac{g_{\mu\nu} d\chi^\nu}{d\lambda} = \frac{\partial \mathcal{L}}{\partial (d\chi^\mu / d\lambda)}, \mathfrak{E} = \rho_t = \left(1 - \frac{2\mathfrak{M}}{r}\right) dt/d\lambda, \mathfrak{E} \\
&= r^2 d\varphi/d\lambda, \left(1 - \frac{2\mathfrak{M}}{r}\right) \left(\frac{dt}{d\lambda}\right)^2 - \left(1 - \frac{2\mathfrak{M}}{r}\right)^{-1} \left(\frac{dr}{d\lambda}\right)^2 - r^2 \left(\frac{d\varphi}{d\lambda}\right)^2 \mathfrak{E}^2 - \left(\frac{dr}{d\lambda}\right)^2 \\
&- \mathfrak{E}^2 / r^2 \left(1 - \frac{2\mathfrak{M}}{r}\right), \frac{dr^\ominus}{d\lambda} \left(1 - \frac{2\mathfrak{M}}{r}\right)^{-1}, r^\ominus \\
&= r + 2\mathfrak{M} \ln (t - 2\mathfrak{M}) d/d\lambda (t \otimes r^{\ominus \otimes \dagger}), du/d\lambda = dt/d\lambda - dr^*/d\lambda, dr^*/d\lambda \\
&= dr^*/dr^* dr/d\lambda = \left(1 - \frac{2\mathfrak{M}}{r}\right)^{-1} \mathfrak{E}, r - 2\mathfrak{M} = \mathfrak{E}, du/d\lambda = 2/\left(1 - \frac{2\mathfrak{M}}{r}\right) \mathfrak{E}, du/d\lambda \\
&= 2\mathfrak{E} - 4\mathfrak{M}/\lambda, u(\lambda) = 2\mathfrak{E}\lambda - 4\mathcal{M} \ln (\lambda/\kappa_1), u(v) = -4\mathcal{M} \ln (\lambda/\kappa_1), (v) \\
&= 4\mathcal{M} \ln (v_0 - v/\kappa_1 \kappa_2)
\end{aligned}$$

$$\square f_\omega = \frac{1}{1} - \frac{2\mathfrak{M}}{r} \partial_t^2 f_\omega - \frac{\left(1 - \frac{2\mathfrak{M}}{r}\right)^2}{r \partial_r f_\omega} - \left(1 - \frac{2\mathfrak{M}}{r}\right) \partial_r^2 f_\omega = \frac{2i\omega \mathcal{M}}{r^2} - 2\mathcal{M}r = \mathfrak{D}(r^{-2})$$



$$\begin{aligned}
u(v) &= -4\mathcal{M} \ln\left(v_0 - \frac{v}{\kappa}\right), \kappa = \kappa_1 \kappa_2 \otimes \frac{d\varphi}{d\tau}, \rho_\omega \sim \frac{1}{\sqrt{\omega}} e^{-i\omega u(v)} \\
&\sim \frac{1}{\sqrt{\omega'}} e^{-i\omega v} \delta(\theta, \varphi), \alpha_{\omega\omega'} = (f_{\omega'}^\otimes \rho_\omega) = \frac{i \int_{2\mathcal{M}} \delta\mathcal{V}_\chi}{1} - \frac{2\mathfrak{M}}{r} (f_{\omega'}^\otimes \partial_t \rho_\omega - \partial_t f_{\omega'}^\otimes \rho_\omega) \\
&= \frac{\mathfrak{C} \int_{2\mathcal{M}} \frac{r^2}{1} - \frac{2\mathfrak{M}}{r} e^{-i\omega' v} e^{-i\omega u(v)}}{r^2 \left(\sqrt{\frac{\omega'}{\omega}} + \sqrt{\frac{\omega}{\omega'}} \right) dr} = -\mathfrak{C} \int_{-\infty}^{v_0} \delta v \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' v} e^{-i\omega u(v)}, \beta_{\omega\omega'} \\
&= -(f_{\omega'}^\otimes \rho_\omega) = \mathfrak{C} \int_{-\infty}^0 \delta v \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' v} e^{-i\omega u(v)} \alpha_{\omega\omega'} \\
&= -\mathfrak{C} \int_{-\infty}^{v_0} \delta s \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' s} e^{-i\omega' v_0} e^{i\omega 4\mathcal{M} \ln\left(-\frac{s}{\mathfrak{R}}\right)}, \beta_{\omega\omega'} \\
&= \mathfrak{C} \int_{-\infty}^{v_0} \delta s \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' s} e^{-i\omega' v_0} e^{i\omega 4\mathcal{M} \ln\left(-\frac{s}{\mathfrak{R}}\right)}
\end{aligned}$$



$$\begin{aligned}
\alpha_{\omega\omega'} &= i\mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{i\omega'v_0} e^{i\omega 4\mathcal{M} \log\left(\frac{-is'}{\mathfrak{K}}\right)}, \beta_{\omega\omega'} \\
&= -i\mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{i\omega'v_0} e^{i\omega 4\mathcal{M} \log\left(\frac{-is'}{\mathfrak{K}}\right)}, \log\left(\frac{is'}{\mathfrak{K}}\right) \\
&= \ln\left(\frac{|s'|}{\mathfrak{K}}\right) - \frac{i\pi}{2}, \log\left(\frac{-is'}{\mathfrak{K}}\right) = \ln\left(\frac{|s'|}{\mathfrak{K}}\right) + \frac{i\pi}{2}, \alpha_{\omega\omega'} \\
&= i\mathfrak{C} e^{i\omega'v_0} e^{2\omega\mathcal{M}\pi} \int_{-\infty}^0 \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{i\omega 4\mathcal{M} \ln\left(\frac{|s'|}{\mathfrak{K}}\right)}, \beta_{\omega\omega'} \\
&= -i\mathfrak{C} e^{i\omega'v_0} e^{-2\omega\mathcal{M}\pi} \int_{-\infty}^0 \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{i\omega 4\mathcal{M} \ln\left(\frac{|s'|}{\mathfrak{K}}\right)}, |\alpha_{\omega\omega'}|^2 \\
&= e^{8\pi\omega\mathcal{M}} |\beta_{\omega\omega'}|^2, (\rho_{\omega_1}\rho_{\omega_2}) = \Gamma(\omega_1)\delta(\omega_1 - \omega_2)(\rho_{\omega_1}\rho_{\omega_2}) \\
&= (\rho_{\omega_1}^{(1)}\rho_{\omega_2}^{(1)}) + (\rho_{\omega_1}^{(2)}\rho_{\omega_2}^{(2)}), (\rho_{\omega_1}^{(1)}\rho_{\omega_2}^{(1)}) \\
&= (1 - \Gamma(\omega_1))\delta(\omega_1 - \omega_2), \Gamma(\omega_1)\delta(\omega_1 - \omega_2) \\
&= \int_{\infty}^0 \mathfrak{d}\omega' (\alpha_{\omega_1\omega'}^{\odot} \alpha_{\omega_2\omega'} - \beta_{\omega_1\omega'}^{\odot} \beta_{\omega_2\omega'}), \mathfrak{L}_{\omega} = (\rho_{\omega}^{(2)}, \phi) \\
&= \int_{\infty}^0 \mathfrak{d}\omega' (\alpha_{\omega\omega'} \alpha_{\omega'} + \beta_{\omega\omega'} \alpha_{\omega'}^{\dagger})
\end{aligned}$$

$$\begin{aligned}
\langle \mathcal{N} \rangle &= \left\langle 0 | \mathfrak{L}_{\omega}^{\dagger} \mathfrak{L}_{\omega} | 0 \right\rangle = \int_{\infty}^0 \mathfrak{d}\omega' \beta_{\omega\omega'} \langle \omega' | \int_{\infty}^0 \mathfrak{d}\omega'' \beta_{\omega\omega''}^{\odot} | \omega'' \rangle = \int_{\infty}^0 \mathfrak{d}\omega' |\beta_{\omega\omega'}|^2, \Gamma(\omega)\delta(0) \\
&= \int_{\infty}^0 \mathfrak{d}\omega' (|\alpha_{\omega\omega'}|^2 - |\beta_{\omega\omega'}|^2) = (e^{8\pi\mathcal{M}\omega} - 1) \int_{\infty}^0 \mathfrak{d}\omega' |\beta_{\omega\omega'}|^2, \delta(\omega_1 - \omega_2) \\
&= \lim_{\tau \rightarrow \infty} \frac{1}{2\pi} \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} dt e^{i(\omega_1 - \omega_2)t} \Gamma(\omega) \lim_{\tau \rightarrow \infty} \frac{\tau}{2\pi} = (e^{8\pi\mathcal{M}\omega} - 1) \int_{\infty}^0 \mathfrak{d}\omega' |\beta_{\omega\omega'}|^2, \langle \mathcal{N} \rangle \\
&= \lim_{\tau \rightarrow \infty} \frac{\tau}{2\pi} \Gamma(\omega) 1/e^{8\pi\mathcal{M}\omega} - 1, \Gamma(\omega)/2\pi \cdot 1/e^{8\pi\mathcal{M}\omega} - 1
\end{aligned}$$



$$\mathcal{T} = \frac{1}{16\pi k_{\mathfrak{B}} \mathfrak{M}} = \frac{\kappa}{2\pi}, \rho = \int d\omega \Lambda(\omega) e^{i\gamma(\omega)} \rho_{\omega}, \mathcal{T} = \frac{\hbar^4 c^4}{16\pi \mathfrak{G} \mathfrak{M} k_{\mathfrak{B}}} \approx 10^{-7} \left(\frac{\mathfrak{M}_{\odot}}{\mathfrak{M}} \right) \mathfrak{K}, \frac{d\mathfrak{E}}{dt}$$

$$= \frac{8\pi r_s^2 \sigma \mathcal{T}^4 dM}{dt} = -\frac{\beta m_{\rho}^4}{t_{\rho} \mathcal{M}^4}, \mathfrak{M}(t) = \left(\mathfrak{M}_0^4 - \frac{6\beta m_{\rho}^4}{t_{\rho}} t \right)^{\frac{1}{2}}, \Delta_t = t_{\rho} / 3\beta \left(\frac{\mathfrak{M}_0}{m_{\rho}} \right)^3$$

9. Modelo Englert – Brout.

$$\begin{aligned} \mathcal{H}_{int} &= i e \Lambda_{\mu} \varphi \overleftrightarrow{\partial}_{\mu} \varphi - e^2 \varphi \overleftrightarrow{\partial}_{\mu} \varphi \Lambda_{\mu} \Lambda_{\mu} \otimes^{\mu\nu} \odot_{\mu\nu} (q) = (2\pi)^4 i e^2 (g_{\mu\nu}) \langle \varphi_1 \rangle^2 - \left\langle \frac{q_{\mu} q_{\nu}}{q^2} \right\rangle \langle \varphi_1 \rangle^2, \mu^2 \\ &= e^2 \langle \varphi_1 \rangle^2, \delta \varphi_{\Lambda} = \Sigma_{\alpha} \Lambda^{\epsilon} \alpha^{(\chi)} \mathfrak{T}_{\alpha AB} \varphi B' \delta \Lambda_{\alpha, \mu} \\ &= \Sigma_c \mathfrak{E}^{\epsilon} c^{(x) c a c b} \Lambda_{\beta, \nu} + \partial_{\mu} \epsilon_{\alpha(x)}, \frac{\left(\frac{l}{(2\pi)^4} \right) \Sigma_{A, B' C'} \mathfrak{T}_{\alpha, AB'} \langle \varphi_{B'} \rangle \mathfrak{T}_{\alpha AC'} \langle \varphi_{C'} \rangle}{q^2} \\ &\equiv, \left(\frac{-l}{(2\pi)^4} \right) \left(\frac{\langle \varphi \rangle \mathfrak{T}_{\alpha} \mathfrak{T}_{\alpha} \langle \varphi \rangle}{q^2} \right), \boxtimes^{\mu\nu} \otimes \hbar_{\mu\nu}^{\alpha} (q) \delta \delta \\ &= -i (2\pi)^4 \lambda^2 \left(\frac{\langle \varphi \rangle \mathfrak{T}_{\alpha} \mathfrak{T}_{\alpha} \langle \varphi \rangle}{q^2} \right) \otimes \left(g_{\mu\nu} - \frac{q_{\mu} q_{\nu}}{q^2} \right), \mu_{\alpha}^2 = - \left(\frac{\langle \varphi \rangle \mathfrak{T}_{\alpha} \mathfrak{T}_{\alpha} \langle \varphi \rangle}{q^2} \right), \mathcal{H}_{int} \\ &= -\eta \bar{\psi} \gamma_{\mu} \gamma_5 \bar{\psi} B_{\mu} - \epsilon \bar{\psi} \gamma_{\mu} \psi \Lambda_{\mu}, \delta^{-1}(\rho) = \gamma \rho - \Sigma(\rho) \\ &= \gamma \rho (1 - \Sigma_2(\rho^2)) - \Sigma_1(\rho^2), m(1 - \Sigma_2(m^4)) - \Sigma_1(m^4), \mathfrak{T}_{\mu}^5 \\ &= -\eta \lim_{\xi \rightarrow 0} \bar{\psi}'(\chi + \xi) \gamma_{\mu} \gamma_5 \psi'(\chi), \psi'(\chi) \\ &= \exp(-i \int_{-\infty}^{\chi} \eta B_{\mu}(\gamma) d\gamma^{\mu} \gamma_5) \psi(\chi) \otimes_{\mu\nu}^5 (q) \\ &= \eta^2 l / (2\pi)^4 \int \mathfrak{T}_r \left(\delta \left(\rho - \frac{1}{2q} \right) \Gamma_{\nu 5} \left(\rho - \frac{1}{2q}; \rho + \frac{1}{2q} \right) \boxtimes \delta \left(\rho + \frac{1}{2q} \right) \gamma_{\mu} \gamma_5 \right. \\ &\quad \left. - \delta(\rho) \left(\frac{\partial \delta^{-1}(\rho)}{\partial \rho_{\nu}} \right) \delta(\rho) \gamma_{\mu} \right) d^4 \rho \\ &= \varrho_{\nu} \Lambda_{\nu 5} \left(\rho - \frac{1}{2q}; \rho + \frac{1}{2q} \right) \Sigma \left(\rho - \frac{1}{2q} \right) \gamma_5 + \gamma_5 \Sigma \left(\rho + \frac{1}{2q} \right), \varrho_{\nu} \Gamma_{\nu 5} \\ &= \varrho_{\nu} \gamma_{\nu} \gamma_5 (1 - \Sigma_2) + 2 \Sigma_1 \gamma_5 - 2(\varrho_{\nu} \rho_{\nu}) (\gamma_{\lambda} \rho_{\lambda}) (\partial \Sigma_2 / \partial \rho^2 \gamma_5) \end{aligned}$$

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Apéndice C.

Formalización de la dualidad holográfica en campos cuánticos curvos.

1. Dualidad/Gravedad – Gauge en espacios cuánticos curvos.

1.1. Grupo Conforme.

$$\begin{aligned}
\iota[\mathfrak{M}_{\mu\nu}\mathfrak{M}_{\rho\sigma}] &= \eta_{\nu\rho}\mathfrak{M}_{\mu\sigma} - \eta_{\mu\rho}\mathfrak{M}_{\nu\sigma} - \eta_{\sigma\mu}\mathfrak{M}_{\rho\nu} + \eta_{\sigma\nu}\mathfrak{M}_{\rho\mu}[\mathfrak{P}_\mu\mathfrak{M}_{\sigma\rho}] \\
&= \eta_{\mu\rho}\mathfrak{P}_\sigma - \eta_{\mu\sigma}\mathfrak{P}_\rho[\mathfrak{P}_\mu\mathfrak{P}_\nu]\iota[\mathfrak{D}, \mathfrak{P}_\mu] = \mathfrak{P}_\mu[\mathfrak{M}_{\mu\nu}, \mathfrak{D}] \\
\mathfrak{K}_{\mu\nu}:\chi^\mu &\rightarrow \chi^\mu + \frac{\alpha^\mu\chi^2}{1} + 2\chi^\nu\alpha_\nu + \alpha^2\chi^2\iota[\mathfrak{M}_{\mu\nu}\mathfrak{K}_\rho] = \eta_{\mu\rho}\mathfrak{K}_\nu - \eta_{\nu\sigma}\mathfrak{K}_\mu[\mathfrak{D}, \mathfrak{K}_\mu]\iota\mathfrak{K}_\mu[\mathfrak{P}_\mu\mathfrak{K}_\nu] \\
&= 2\iota(\mathfrak{M}_{\mu\nu} - \eta_{\mu\nu}\mathfrak{D})[\mathfrak{K}_\mu, \mathfrak{K}_\nu]
\end{aligned}$$

$$\mathfrak{J}_{\mu\nu} = \mathfrak{M}_{\mu\nu}, \mathfrak{J}_{\mu d} = \frac{1}{2[\mathfrak{K}_\mu - \mathfrak{P}_\mu]}, \mathfrak{J}_{\mu(d+1)} = \frac{1}{2[\mathfrak{K}_\mu + \mathfrak{P}_\mu]}, \mathfrak{J}_{(d+1)d} = \mathfrak{D}$$

$$\mathfrak{J}_{\alpha\beta} = \begin{pmatrix} \mathfrak{J}_{\mu\nu} & \mathfrak{J}_{\mu d} & \mathfrak{J}_{\mu(d+1)} \\ -\mathfrak{J}_{\mu d} & 0 & \mathfrak{D} \\ -\mathfrak{J}_{\mu(d+1)} & -\mathfrak{D} & 0 \end{pmatrix}$$

$$\mathcal{P}_\mu:\chi_\mu \rightarrow \chi_\mu + \alpha_\mu \Rightarrow d, \mathfrak{M}_{\mu\nu}:\chi_\mu \rightarrow \Lambda_\mu^\nu\chi_\nu \Rightarrow \frac{(d-1)d}{2}, \mathfrak{D}:\chi_\mu \rightarrow \lambda\chi_\mu, \mathfrak{K}_\mu:\chi_\mu$$

$$\rightarrow \chi_\mu + \frac{\alpha_\mu\chi^2}{1} + 2\chi_\nu\alpha^\nu + \alpha^2\chi^2 \Rightarrow d$$

$$\begin{aligned}
\chi \rightarrow \lambda\chi \Rightarrow \phi(\chi) \rightarrow \phi(\chi)' &= \lambda^\Delta\phi(\lambda\chi)[\mathfrak{D}, \mathcal{P}_\mu] = \iota\mathcal{P}_\mu \Rightarrow \mathfrak{D}(\mathcal{P}_\mu\phi) = -\iota(\Delta+1)(\mathcal{P}_\mu\phi), \langle\phi(0)|\phi(\chi)\rangle \\
&\equiv 1/(\chi^2)^\Delta
\end{aligned}$$

1.2. Espacio Anti – de Sitter en espacios cuánticos curvos.



$$\begin{aligned}
\mathfrak{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathcal{R}} &= -\Lambda g_{\mu\nu}, \mathfrak{R}_{\mu\nu\theta\sigma} = \frac{1}{\ell^2} (g_{\mu\theta}g_{\nu\sigma} - g_{\mu\sigma}g_{\nu\theta}), \mathfrak{R}_{\mu\nu} = -\frac{3}{\ell^2 g_{\mu\nu}}, \mathcal{R} = -\frac{12}{\ell^2}, \mathfrak{d}s^2 \\
&= -d\chi_0^2 - d\chi_4^2 + d\chi_1^2 + d\chi_2^2 + d\chi_3^2 - \chi_0^2 - \chi_4^2 + \chi_1^2 + \chi_2^2 + \chi_3^2 = -\ell^2, \frac{\mathfrak{d}s^2}{\ell^2} \\
&= -\cos^2 \rho d\tau^2 + d\rho^2 + \sin^2 \rho (d\theta^2 + \sin^2 \theta d\varphi^2), \frac{\mathfrak{d}s^2}{\ell^2} \\
&\approx -d\tau^2 + d\rho^2 + \rho^2 (d\theta^2 + \sin^2 \theta d\varphi^2), \mathfrak{d}s^2 \\
&= \left(1 + \frac{r^2}{\ell^2}\right) dt^2 + \left(1 + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2), \chi_0 \\
&= \frac{lr}{2} \left(\overline{\chi_i^2} - t^2 + \frac{1}{r^2} + 1\right), \chi_i = lr x_i (i = 1, 2), \chi_3 = \frac{lr}{2} \left(\overline{\chi_i^2} - t^2 + \frac{1}{r^2} + 1\right) \chi_4 \\
&= lrt, \frac{\mathfrak{d}s^2}{\ell^2} = r^2 \left(-dt^2 + d\overline{\chi^2}\right) + \frac{dr^2}{r^2}, \mathfrak{d}s^2 = \ell^2 / z^2 \left(-dt^2 + d\overline{\chi^2} + dz^2\right)
\end{aligned}$$

1.3. Límite de 't Hooft en espacios cuánticos curvos.

$$\mathfrak{Q} = \frac{\mathfrak{N}}{2\mathfrak{G}_{\mathfrak{M}}^2 \mathcal{F}_{\mu\nu}^{\mathfrak{M}} \mathcal{F}_{\mathfrak{M}}^{\mu\nu}}$$

1.4. Prescripción de Gubser – Klebanov – Polyakov y Witten en espacios cuánticos curvos.

$$\mathfrak{Z}_{\mathfrak{CFT}} = e^{-\mathcal{W}}, \ell^4 \ell_s^4 \sim g_{\mathfrak{YM}}^2 \mathcal{N} \sim g_{\delta} \mathfrak{N} \gg 1, \mathcal{Z}_{\mathfrak{SUGRA}} \approx e^{-\mathfrak{S}_{\mathfrak{SUGRA}}^{\mathfrak{E}}}, \mathcal{Z}_{\mathfrak{CFT}} \approx e^{-\mathfrak{S}_{\mathfrak{CFT}}^{\mathfrak{E}}} = e^{-\mathcal{W}} = \mathcal{Z}_{\mathfrak{CFT}}$$

1.5. Correspondencia Campo ↔ Operador en espacios cuánticos curvos.

$$\langle e^{\int d^3 \chi \phi_0(\vec{\chi}) \mathcal{O}(\vec{\chi})} \rangle_{\mathfrak{CFT}} = e^{-\mathfrak{S}_{\mathfrak{SUGRA}}^{\mathfrak{E}}[\phi|\partial\mathfrak{M} \hookrightarrow \phi_0]}, \langle e^{\int d^3 \chi \mathfrak{h}_{\alpha\beta}^0 \mathcal{T}^{\alpha\beta}} \rangle_{\mathfrak{CFT}} = e^{-\mathfrak{S}_{\mathfrak{SUGRA}}^{\mathfrak{E}}[\mathfrak{h}_{\mu\nu}|\partial\mathfrak{M} \hookrightarrow \mathfrak{h}_{\alpha\beta}^0]}$$

1.6. Partículas y Campos en el espacio tiempo AdS – curvo.



$$\begin{aligned}
(\nabla^\mu \nabla_\mu - m^4)\Phi(z, \chi^\eta) &= \frac{\int d^2 \vec{k}}{(2\pi)^4 d\omega \mathfrak{f}_\kappa(z) e^{i\kappa_\mu \chi^\mu} d^2 \mathfrak{f}_\kappa} - \frac{2}{z d \mathfrak{f}_\kappa} - \left(\kappa^2 + \frac{m^4 \ell^4}{z^2} \right) \mathfrak{f}_\kappa(z) \\
&= \alpha_1 z^{\frac{3}{2}} \mathfrak{K}_\nu(\mathfrak{K}z) + \frac{\alpha_2 z^{\frac{3}{2}} \mathfrak{I}_\nu(\mathfrak{K}z) d^2 \mathfrak{f}_\kappa}{dz^2} - \kappa^4 \mathfrak{f}_\kappa(z) \\
&= \frac{\alpha_1 z^{\frac{3}{2}} \pi}{2 \sin \pi \nu \left[\frac{1}{\Gamma(1-\nu) \left(\frac{\mathfrak{K}z}{2}\right)^{-\nu}} - \frac{1}{\Gamma(1+\nu) \left(\frac{\mathfrak{K}z}{2}\right)^\nu} \right]}, m_{\mathfrak{S}\mathfrak{I}}^4 \geq m^4 \geq m_{\mathfrak{S}\mathfrak{I}}^4 + \frac{1}{\ell^4} \Rightarrow 0 \geq \nu \\
> 1, \phi(r, \chi^\eta) &= \frac{\alpha(\chi^\eta)}{r^{3-\Delta}} + \frac{\beta(\chi^\eta)}{r^\Delta}, \phi = \frac{\alpha}{r} + \frac{\beta}{r^2 \partial_r \phi} \\
&= -\frac{\alpha}{r^2} - \frac{2\beta}{r^4 \otimes \alpha'} + \alpha \phi + \beta \partial_r \phi = \alpha \left(\frac{\alpha}{r} + \frac{\beta}{r^2} \right) + \beta \left(-\frac{\alpha}{r^2} - \frac{2\beta}{r^4} \right) = \frac{\alpha'}{r} + \frac{\beta'}{r^2}
\end{aligned}$$

1.7. Deformaciones en AdS/CFT en espacios cuánticos curvos.

$$\mathfrak{I}_{\mathfrak{S}\mathfrak{I}} \rightarrow \mathfrak{I}_{\mathfrak{S}\mathfrak{I}} + \rho \int d^3 \chi \mathcal{O}(\chi)$$

2. Agujeros Negros Cuánticos en espacios curvos (Formalización).

2.1. Principio Variacional.

$$\begin{aligned}
\mathfrak{I} &= \frac{1}{2\kappa \int d^4 \chi \sqrt{-g} \mathcal{R} + \mathfrak{I}_{\mathfrak{S}}} , \delta \mathfrak{I} = \frac{1}{2\kappa \int d^4 \chi \sqrt{-g} \mathfrak{G}_{\alpha\beta} \delta g^{\alpha\beta} + \int d^4 \chi \sqrt{-g} g^{\alpha\beta} \delta \mathcal{R}_{\alpha\beta} + \delta \mathfrak{I}_{\mathfrak{S}}} , \mathfrak{G}_{\alpha\beta} \\
&= \mathcal{R}_{\alpha\beta} - \frac{1}{2g_{\alpha\beta} \mathfrak{R}} , \delta \mathfrak{I}_{\mathfrak{S}} = - \int_{\mathcal{M}} d^4 \chi \sqrt{-g} g^{\alpha\beta} \delta \Gamma_{\alpha\beta}^\mu - \sqrt{-g} g^{\alpha\mu} \delta \Gamma_{\alpha\beta}^\beta + \sqrt{-g} \delta \mathcal{R}_{\alpha\beta} \\
&= \oint_{\partial \mathcal{M}} \epsilon v^\mu \eta_\nu \sqrt{-\hbar} d^3 \chi , \mathfrak{I}_{\mathfrak{S}} = \oint_{\partial \mathcal{M}} d^3 \chi \sqrt{-\hbar} \kappa , \mathfrak{I} \\
&= \frac{1}{2\kappa \int_{\mathcal{M}} d^4 \chi \sqrt{-g} \mathcal{R}} + 1/\kappa \oint_{\partial \mathcal{M}} d^3 \chi \sqrt{-\hbar} \mathfrak{K}
\end{aligned}$$

$$ds^2 = -\mathcal{N}(r) dt^2 + \mathcal{H}(r) dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2), \mathfrak{K}_{\mathfrak{R}\mathfrak{I}\mathfrak{S}} = \mathcal{R}^{\alpha\beta\gamma\sigma} \mathcal{R}_{\alpha\beta\gamma\sigma}, \mathcal{I}[\mathfrak{g}_{\mu\nu}]$$

$$= \frac{1}{2\kappa \int_{\mathcal{M}} d^4 \chi \sqrt{-g} \mathcal{R}} + \frac{1}{\kappa \oint_{\partial \mathcal{M}} d^3 \chi \mathfrak{K} \sqrt{-\hbar}}, ds^2$$

$$= -\left(1 - \frac{\mu}{r}\right) dt^2 + \left(1 - \frac{\mu}{r}\right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\phi^2), M = \frac{4\pi\mu}{\kappa} = \mu/2\mathfrak{G}$$



2.2. Modelo Reissner – Nordström.

$$\begin{aligned}
 \mathcal{J}[g_{\mu\nu}\Lambda_\mu] &= \frac{1}{2\kappa \int_{\mathcal{M}} d^4\chi \sqrt{-g} \left(\mathcal{R} - \frac{1}{4F^{\mu\nu}F_{\mu\nu}} \right)} + \frac{1}{\kappa \oint_{\partial\mathcal{M}} d^3\chi \mathfrak{K} \sqrt{-\mathfrak{h}}}, \mathcal{R}_{\mu\nu} - \frac{1}{2\mathcal{R}g_{\mu\nu}} \\
 &= \frac{1}{2\mathfrak{S}_{\mu\nu}^{\mathfrak{E}\mathfrak{M}} \nabla_\mu F^{\mu\nu}}, \mathfrak{S}_{\mu\nu}^{\mathfrak{E}\mathfrak{M}} = F_{\mu\alpha} F_\nu^\alpha - \frac{1}{4g_{\mu\nu}F^2}, \Lambda \equiv \Lambda_\mu d\chi^\mu = \left(\frac{q}{r} - \frac{q}{r_+} \right) dt, F \\
 &= -\frac{q}{r^2 dr} \wedge dt, ds^2 = -\mathfrak{f}(r) dt^2 + \mathfrak{f}(r)^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2), \mathfrak{f}(r) \\
 &= 1 - \frac{\mu}{r} + \frac{q^2}{4r^2} = \frac{(r-r_-)(r-r_+)}{r^2}, \mathcal{M} = \frac{4\pi\mu}{\kappa} = \frac{\mu}{2\mathfrak{G}}, \mathfrak{Q} \equiv \frac{1}{\kappa \oint_{-\infty}^\infty d^2\star\mathcal{F}} = -\frac{q}{4\mathfrak{G}}, r_\pm \\
 &= \mathfrak{G}(\mathfrak{M} \pm \sqrt{\mathfrak{M}^2 - 4\mathfrak{Q}^2}), \Phi = \Lambda_t|\Delta_{r=\infty} - \Lambda_t|\Delta_{r=r_+} = 4\mathfrak{G}\mathfrak{Q}/r \blacksquare
 \end{aligned}$$

2.3. Modelo anti – de Sitter.

$$\begin{aligned}
 ds^2 &= \frac{r^2}{\ell^2(-dt^2 + \ell^2 d\Sigma_\kappa^2)}, ds^2 = -N(r)dt^2 + \mathcal{H}(r)dr^2 + \delta(r)d\Sigma_\kappa^2 \\
 &= \left\langle \begin{array}{l} d\theta^2 + \sin^2\theta d\varphi^2 \quad \propto \kappa = +1 \\ \frac{1}{\ell^2 \sum_{i=1}^2 d\chi_i^2} \quad \propto \kappa = 0 \\ d\theta^2 + \sin^2\theta d\varphi^2 \quad \div \kappa = -1 \end{array} \right\rangle, d\Sigma_\kappa^2 = \frac{dy^2}{1} - \kappa y^2 + (1 + \kappa y^2) dz^2
 \end{aligned}$$

2.4. Modelo Schwarzschild-AdS.

$$\begin{aligned}
 \mathcal{J}[g_{\mu\nu}] &= \frac{1}{2\kappa \int_{\mathcal{M}} d^4\chi \sqrt{-g} (\mathfrak{R} - 2\Lambda)} + \frac{1}{\kappa \oint_{\partial\mathcal{M}} d^3\chi \mathfrak{K} \sqrt{-\mathfrak{h}}}, \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathcal{R}} = -\Lambda g_{\mu\nu}, ds^2 \\
 &= -\left(\kappa - \frac{\mu}{r} + \frac{r^2}{\ell^2} \right) dt^2 + \left(\kappa - \frac{\mu}{r} + \frac{r^2}{\ell^2} \right)^{-1} dr^2 + r^2 d\Sigma_\kappa^2, \kappa - \frac{\mu}{r} + \frac{r^2}{\ell^2}
 \end{aligned}$$

2.5. Modelo Escalar Simple.

$$\begin{aligned}
 \mathcal{J}[g_{\mu\nu}, \phi] &= \frac{\int_{\mathcal{M}} d^4\chi \sqrt{-g} \left[\frac{\mathfrak{R}}{2\kappa} - \frac{1}{2\partial_\mu\phi\partial^\mu\phi} - \mathcal{V}(\phi) \right]}{\sqrt{-g}\partial_\mu(\sqrt{-g}g^{\mu\nu}\partial_\nu\phi)} - \frac{\frac{\partial\mathcal{V}}{\partial\phi}d\mathcal{V}}{d\phi|\psi_{\phi=0}}, \mathcal{V}(0) \\
 &= -\frac{3}{\kappa\ell^2}, \frac{d^2\mathcal{V}}{d\phi^2|\psi_{\phi=0}}, \mathcal{V}(\phi)|_{\text{ub}\mathfrak{S}} = \left(-\frac{1}{\ell^2} + \alpha\phi \right) (4 + 2\cos\hbar\phi) - 6\alpha\sin\hbar\phi
 \end{aligned}$$

2.6. Modelo Escalar Neutro.



$$\begin{aligned}
\varepsilon_{\mu\nu} &= \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathfrak{R}} - \kappa \mathcal{T}_{\phi\mu\nu}^{\phi}, \mathcal{T}_{\phi\mu\nu}^{\phi} = \partial_{\mu}\phi\partial_{\nu}\phi - g_{\mu\nu}\left(\frac{1}{2(\partial\psi)^2} + \mathfrak{B}(\phi)\right), ds^2 \\
&= \Omega(\chi) \left[-\mathfrak{f}(\chi)dt^2 + \frac{\eta^2 d\chi^2}{\mathfrak{f}(\chi)} + \frac{dy^2}{\ell^2} + \frac{dz^2}{\ell^2} \right] \mathfrak{E}_t^t - \mathfrak{E}_{\chi}^{\chi} = 0 \rightarrow \phi'^2 \\
&= 3\Omega'^2 - \frac{2\Omega''\Omega}{\Omega^2}, \mathfrak{E}_t^t - \mathfrak{E}_y^y = 0 \rightarrow \mathfrak{f}'' + \frac{\Omega'\mathfrak{f}'}{\Omega} = 0, \mathfrak{E}_t^t - \mathfrak{E}_y^y = 0 \rightarrow \mathcal{V}(\phi) \\
&= -\frac{1}{\Omega^2\eta^2(\mathfrak{f}\Omega'' + \mathfrak{f}'\Omega')}, \Omega(\chi) = \frac{v^2\chi^{v-1}}{\eta^2(\chi^v - 1)^2}, \phi'^2 \\
&= \frac{(v-1)^2}{\chi^2} - \frac{4v(v-1)\chi^{v-2}}{\chi^v} - 1 + \frac{4v^2\chi^{v-1}}{(\chi^v - 1)^2} + \frac{2(v-1)}{\chi^2} \\
&+ \frac{4v(1-v-\chi^v)\chi^{v-2}\chi^{v-2}}{(\chi^v - 1)^2}, \phi'^2 = v^2 - \frac{1}{2\kappa\chi^2} \\
&\rightarrow \int_{\phi}^{\phi=0} d\phi = \sqrt{v^2 - \frac{1}{2\kappa}} \int_{\chi}^1 \frac{d\chi}{\chi}, \phi(\tau) \ln \chi, \ell_v^{-1} = \sqrt{v^2 - \frac{1}{2\kappa}} (\mathfrak{f}'\Omega)' f(x) \\
&= \frac{c_4\eta^2}{v^2 \int \frac{(\chi^v - 1)^2}{\chi^{v-1}d\chi} + c_1}, f(x) = c_1 + \frac{c_4\eta^2}{v^2 \left(\frac{\chi^{2+v}}{2} + v + \frac{\chi^{2-v}}{2} - v - \chi^2 \right)}, f(x) \\
&= \frac{1}{\ell^2} + \alpha \left[\frac{1}{v^2} - 4 - \frac{\chi^2}{v^2 \left(1 + \frac{\chi^{-v}}{v} - 2 - \frac{\chi^v}{v} + 2 \right)} \right], \mathcal{V}(\phi) \\
&= \frac{\Lambda(v^2 - 4)}{6\kappa v^2 \left[v - \frac{1}{v} + 2e^{-\phi\ell_v(v+1)} + v + \frac{1}{v} - 2e^{\phi\ell_v(v-1)} + 4v^2 - \frac{1}{v^2} - 4e^{-\phi\ell_v} \right]} \\
&+ \alpha \\
&/\kappa v^2 \left[v - \frac{1}{v} + 2 \sin \hbar \phi \ell_v (v+1) - v + \frac{1}{v} \right. \\
&\left. - 2 \sin \hbar \phi \ell_v (v-1) + 4v^2 - \frac{1}{v^2} - 4 \sin \hbar \phi \ell_v \right] \\
\mathcal{V}(\phi) &= \frac{\Lambda}{\kappa} - \frac{\phi^2}{\ell^2} + \frac{\kappa\Lambda}{18(v^2 - 3)} - 1\phi^4 - \frac{\ell_v^3}{90(\Lambda v^2 - 4\Lambda - 6\alpha)\tau^5} + \mathcal{O}|\phi|^6
\end{aligned}$$



$$\begin{aligned}
 ds^2 &= \Omega(\chi) \left[-f(x)dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\theta^2 + \sin^2\theta d\varphi^2 \right], f(x) \\
 &= \frac{1}{\rho^2} + \alpha \left[\frac{1}{\nu^2} - 4 - \frac{\chi^2}{\nu^2 \left(1 + \frac{\chi^{-\nu}}{\nu} - 2 - \frac{\chi^\nu}{\nu} + 2 \right)} \right] + \chi/\Omega(\chi)
 \end{aligned}$$

2.7. Modelo Escalar Eléctricamente Cargado.

$$\begin{aligned}
 \mathfrak{S}[g_{\mu\nu}, \Lambda_\mu \phi] &= \frac{1}{16\pi \mathfrak{G}_N \int d^4\chi \sqrt{-g} \left[\mathfrak{R} - \frac{1}{4e^{\gamma\phi} F^2} - \frac{1}{2\partial_\mu \phi \partial^\mu \phi} - \mathcal{V}(\phi) \right]}, \nabla_\mu (e^{\gamma\phi} F^{\mu\nu}) \\
 &= \frac{1}{\sqrt{-g} \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi)} - \frac{\partial \mathcal{V}}{\partial \phi} - \frac{1}{4\gamma e^{\gamma\phi} F^2}, \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu} \mathcal{R}} = \frac{1}{2 \left[\mathcal{T}_{\mu\nu}^\phi \mathcal{T}_{\mu\nu}^{\mathfrak{E}\mathfrak{M}} \right]}, \mathcal{T}_{\mu\nu}^\phi \\
 &= \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left[\frac{1}{2(\partial\phi)^2} + \mathcal{V}(\phi) \right], \mathcal{T}_{\mu\nu}^{\mathfrak{E}\mathfrak{M}} = e^{\gamma\phi} \left(F_{\mu\alpha} F_\nu^\alpha - \frac{1}{4g_{\mu\nu} F^2} \right)
 \end{aligned}$$

2.8. Modelo Holográfico en espacios cuánticos curvos.



$$\begin{aligned}
J_g &= -\frac{1}{8\pi\mathfrak{G}_N \oint_{\partial\mathcal{M}}^\lambda d^3\chi \mathfrak{K} \sqrt{-\hbar} \Xi(\ell, \mathcal{R}, \nabla\mathfrak{R})}, ds^2 \\
&= -\left(\kappa + \frac{r^2}{\ell^2}\right) dt^2 + \left(\kappa + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2 d\Sigma_\kappa^2 \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -\left(\kappa + \frac{\mathcal{R}^2}{\ell^2}\right) dt^2 + \mathcal{R}^2 d\Sigma_\kappa^2, ds^2 = \frac{r_\beta^2}{\ell^2(-dt^2 + \ell^2 d\Sigma_\kappa^2)}, ds^2 \\
&= -\left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right) dt^2 + \left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2), \mathfrak{I} = \frac{f'}{|4\pi_{r+}} \\
&= \beta^{-1} = \frac{1}{4\pi\left(3r_+^2 + \frac{\ell^2}{\ell_{r_+}^2}\right)}, J_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^\mathfrak{E} = \frac{12\pi\beta}{\kappa\ell^2 \int_{r_+}^{\mathcal{R}} r^2 dr = \frac{4\pi\beta}{\kappa\ell^2(\mathcal{R}^3 - r_+^3)}}, \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -\left(1 - \frac{\mu}{\mathcal{R}} + \frac{\mathcal{R}^2}{\ell^2}\right) dt^2 + \mathcal{R}^2(d\theta^2 + \sin^2\theta d\varphi^2), \eta_\alpha = \frac{\delta_\alpha^r}{\sqrt{g^{rr}}}, \kappa_{\alpha\beta} = \frac{\sqrt{g^{rr}}}{2\partial_r \hbar_{\alpha\beta}}, \mathfrak{K} \\
&= \frac{1}{\mathcal{R}^2 \ell^2 \left(1 - \frac{\mu}{\mathcal{R}} + \frac{\mathcal{R}^2}{\ell^2}\right)^{\frac{1}{2}} \left(-\frac{3\ell^2\mu}{2} + 3\mathcal{R}^2 + 2\mathcal{R}\ell^2\right)}, J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2 \left(-\frac{3\ell^2\mu}{2} + 3\mathcal{R}^2 + 2\mathcal{R}\ell^2\right)}, J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} = J_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^\mathfrak{E} + J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2 \left(-\frac{3\ell^2\mu}{2} + 2\mathcal{R}^2 + 2\mathcal{R}\ell^2 - r_+^3\right)}, ds^2 \\
&= -\left(1 + \frac{r^2}{\ell^2}\right) dt^2 + \left(1 + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2), J_{\mathfrak{U}\mathfrak{b}\mathfrak{S}}^\mathfrak{E} = J_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^\mathfrak{E} + J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta_0}{\kappa\ell^2(-2\mathcal{R}^3 + 2\mathcal{R}\ell^2)}, \beta_0 \sqrt{1 + \frac{\mathcal{R}^2}{\ell^2}} = \beta \sqrt{1 + \frac{\mathcal{R}^2}{\ell^2} - \frac{\mu}{\mathfrak{R}}}, J^\mathfrak{E} = J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} - J_{\mathfrak{U}\mathfrak{b}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2 \left[\left(\frac{3\ell^2\mu}{2} - 2\mathcal{R}^3 - 2\mathcal{R}\ell^2 - r_+^3\right) - \frac{\beta_0}{\beta(-2\mathcal{R}^3 - 2\ell^2\mathfrak{R})}\right]}, F = \beta^{-1} J^\mathfrak{E} \\
&= 4\pi/\kappa\ell^2 \left(\frac{\ell^2\mu}{2} - r_+^3\right)
\end{aligned}$$



$$\begin{aligned}
J_g &= -\frac{1}{\kappa \int_{\partial\mathcal{M}} d^3\chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{L}\mathcal{R}}{2}\right)}, J_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} + J_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} = \frac{4\pi\beta}{\kappa\ell^2 \left(\frac{3\ell^2\mu}{2} - 2r_\beta^3 - 2r_\beta\ell^2 - r_+^3\right)}, J_{\mathfrak{G}}^{\mathfrak{E}} \\
&= \frac{4\pi\beta}{\kappa\ell^2 \left(1 + \frac{\ell^2}{\mathcal{R}^2} - \frac{\mu\ell^2}{\mathcal{R}^3}\right)^{\frac{1}{2}} |(2\mathcal{R}^3 + \kappa\ell^2\mathcal{R})|_{\mathcal{R}=r_\beta}} = \frac{4\pi\beta}{\kappa\ell^2(2r_\beta^3 - 2\ell^2r_\beta - \mu\ell^2)}, J_{\mathfrak{G}}^{\mathfrak{E}} \\
&= J_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} + J_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + J_{\mathfrak{G}}^{\mathfrak{E}} = \frac{4\pi\beta}{\kappa\ell^2 \left(\frac{\ell^2\mu}{2} - r_+^3\right)}, \mathcal{E} = -\frac{\mathfrak{S}^2 \partial J^{\mathfrak{E}}}{\partial \mathcal{J}} = \frac{\mu}{2\mathfrak{G}}, F_{\mathfrak{S}\mathfrak{U}\mathfrak{L}\mathfrak{B}\mathfrak{S}} \\
&= \frac{4\pi}{\kappa\ell^2 \left(\frac{\ell^2\mu}{2} - r_+^3\right)}, J_{\mathfrak{S}\mathfrak{U}\mathfrak{L}\mathfrak{B}\mathfrak{S}} = \left(1 + \frac{3r_+^2}{\ell^2}\right)
\end{aligned}$$

2.8.1. Modelo de Brown-York.

$$\begin{aligned}
\int_{\partial\mathcal{M}}^{\delta} d^3\chi \hbar^{\alpha\beta} \mathcal{T}_{\alpha\beta}, ds^2 &= \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta = -N(\mathcal{R})dt^2 + \delta(\mathcal{R})d\Sigma_{\kappa}^2, \tau^{\alpha\beta} \equiv \frac{2}{\sqrt{-\hbar}\delta\mathcal{J}}, ds^2 \\
&= -\left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right)dt^2 + \left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2 d\Omega^2, ds^2 \\
&= -\left(1 - \frac{\mu}{\mathcal{R}} + \frac{\mathcal{R}^2}{\ell^2}\right)dt^2 + \mathcal{R}^2 d\Omega^2, \mathcal{J} \\
&= \frac{1}{2\kappa \int_{\partial\mathcal{M}}^{\delta} d^4\chi \sqrt{-g} (\mathfrak{R} - 2\Lambda)} + \frac{1}{\kappa \int_{\partial\mathcal{M}}^{\delta} d^3\chi \sqrt{-\hbar} \mathfrak{K}} - \frac{1}{\kappa \int_{\partial\mathcal{M}}^{\delta} d^3\chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{L}\mathcal{R}}{2}\right)}, \tau^{\alpha\beta} \\
&= \frac{1}{8\pi\mathfrak{G} \left(\mathfrak{K}_{\alpha\beta} - \hbar_{\alpha\beta}\mathfrak{K} - \frac{2}{\ell\hbar_{\alpha\beta}} + \mathcal{J}\mathfrak{S}_{\alpha\beta}\right)}, ds_{\text{borde}}^2 = \frac{\mathfrak{R}^2}{\ell^2(-dt^2 + \ell^2 d\Omega^2)}, ds_{\text{dualidad}}^2 \\
&= \gamma_{\alpha\beta} d\chi^\alpha d\chi^\beta = -dt^2 + \ell^2 d\Omega^2, \langle \tau_{\text{dualidad}}^{\alpha\beta} \rangle = \lim_{\mathcal{R} \rightarrow \infty} \frac{\mathcal{R}}{\ell} \tau_{\alpha\beta} \\
&= \mu/16\pi\mathfrak{G}\ell^2(3\delta_\alpha^0\delta_\beta^0 + \gamma_{\alpha\beta})
\end{aligned}$$

2.9. Modelo Hamiltoniano.



$$\begin{aligned}
\mathfrak{S}[g_{\mu\nu}, \Lambda_\mu \phi] &= \int_{\partial\mathcal{M}}^\delta d^4\chi \sqrt{-g} \left[\frac{\mathcal{R}}{2\kappa} - \frac{1}{2(\partial\phi)^2} - \mathcal{V}(\phi) \right] + \frac{1}{\kappa \int_{\partial\mathcal{M}}^\delta d^3\chi \sqrt{-\hbar}}, \mathcal{H}_\perp \\
&= \frac{2\kappa}{\sqrt{g} \left[\pi^{ij} \pi_{ij} - \frac{1}{2(\pi_j^i)^2} \right]} - \frac{1}{2\kappa \sqrt{g}^{(3)} \mathcal{R}} + \frac{1}{2 \left(\frac{\pi\phi^2}{\sqrt{g}} + \sqrt{g} g^{ij} \phi_i \phi_j \right)} + \sqrt{g} \mathcal{V}(\phi), \mathcal{H}_i \\
&= -\pi_j^i |_{,j} + \pi_\phi \phi_{,\psi} d\omega, ds^2 = (N^\perp)^2 dt^2 + g_{ij} (d\chi^i + N^i dt) (d\chi^j + N^j dt), \mathcal{H}[\xi] \\
&= \int_{\partial\mathcal{M}}^\delta d^3\chi (\xi^\perp \mathcal{H}_\perp + \xi^i \mathcal{H}_i) + \mathfrak{Q}[\xi], \delta\mathfrak{Q}[\xi] \\
&= \oint d^2\delta_\ell \left[\frac{\mathfrak{G}^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij}|_\kappa - \xi_\kappa^\perp \delta g_{ij}) + 2\xi_\kappa \delta\pi^{\kappa\ell} + (2\xi^\kappa \pi^{j\ell} - \xi^\ell \pi^{j\kappa}) \delta g_{j\kappa} \right. \\
&\quad \left. - (\sqrt{g} \xi^\perp g^{\ell j} \phi_j + \xi^\ell \pi_\phi) \delta\phi \right], \mathfrak{G}^{ijkl} \equiv \frac{1}{2\sqrt{g} (g^{i\kappa} g^{j\ell} + g^{i\ell} g^{j\kappa} - 2g^{ij} g^{\kappa\ell})}
\end{aligned}$$

$$\begin{aligned}
\delta\mathcal{M} \equiv \delta\mathfrak{Q}[\partial_t] &= \oint d^2\delta_\ell \left[\frac{\mathfrak{G}^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij}|_\kappa - \xi_\kappa^\perp \delta g_{ij}) - \sqrt{g} \xi^\perp g^{\ell j} \phi_j \delta\phi \right], \delta\mathcal{M} = \delta\mathcal{M}_\mathfrak{G} + \delta\mathcal{M}_\phi, \delta\mathcal{M}_\mathfrak{G} \\
&= \oint d^2\delta_\ell \frac{\mathfrak{G}^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij}|_\kappa - \xi_\kappa^\perp \delta g_{ij}), \delta\mathcal{M}_\phi = -\oint d^2\delta_\ell \sqrt{g} \xi^\perp g^{\ell j} \phi_j \delta\phi
\end{aligned}$$

2.10. Termodinámica de agujeros negros cuánticos en espacios curvos.

2.10.1. Espacio – Tiempo de Rindler.



$$\begin{aligned}
ds^2 &= -dt^2 + d\chi^2 + d\gamma^2 + d\beta^2, ds^2 = -\alpha^2 \rho^2 d\tau^2 + d\rho^2 + d\gamma^2 + d\beta^2, ds^2 \\
&= N(r) dt_{\xi}^2 + \mathcal{H}(r) dr^2 + \delta(r) d\Sigma_{\kappa}^2, ds^2 = \frac{g_{rr} 4N}{[(N)']^2 \left[\frac{\rho^2 [(N)']^2}{4N g_{rr} d\tau_{\xi}^2} + d\rho^2 \right] \mathcal{J}} = \frac{1}{\beta} \\
&= \frac{[(N)']^2}{4\pi \sqrt{N^2 g_{rr}} \Big|_{\mathbb{H}} \mathcal{J}_{\delta ch - \mathfrak{U} \delta \delta}} = \frac{1}{4\pi r_+ \left(1 + \frac{3r_+^2}{\rho^2} \right) \mathcal{J}_{\mathfrak{N} \eta - \mathfrak{U} \delta \delta}} = \frac{1}{4\pi r_+ \left(1 + \frac{3r_+^2}{\rho^2} - \frac{q^2}{4r_+^2} \right)}, ds^2 \\
&= \Omega(x) \left[-f(x) dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\theta^2 + d\Sigma_{\kappa}^2 \right], f(x) \\
&= \frac{1}{\rho^2} + \alpha \left[\frac{1}{\nu^2} - 4 - \frac{\chi^2}{\nu^2 \left(1 + \frac{\chi^{-\nu}}{\nu} - 2 - \frac{\chi^{\nu}}{\nu} + 2 \right)} \right] + \frac{\kappa \chi}{\Omega(x)}, f' \Omega(x) \\
&= \frac{\alpha}{\eta^2} + 2\kappa + \kappa \nu \chi^{\nu} + \frac{1}{\chi^{\nu}} - 1, \mathcal{J} = \frac{f'}{4\pi \eta} \Big|_{\chi_h} \\
&= \frac{\frac{1}{f'}}{4\pi \eta \Omega(\chi_h) \left(\frac{\alpha}{\eta^2} + 2\kappa + \kappa \nu \chi_h^{\nu} + \frac{1}{\chi_h^{\nu}} - 1 \right)}
\end{aligned}$$

2.10.2. Transiciones de fase de agujeros negros cuánticos en espacios curvos.



$$\begin{aligned}
ds^2 &= \Omega(x) \left(-f(x)dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\theta^2 + \sin^2\theta d\varphi^2 \right), \Omega(x) = \frac{1}{\eta^2(\chi-1)^2}, f(x) \\
&= \frac{1}{\ell^2} + \frac{1}{3\alpha(\chi-1)^3} + \eta^2\chi(\chi-1)^2, \chi = 1 + \frac{1}{\eta r}, \chi = 1 - \frac{1}{\eta r}, \Omega(x)f(x) = F(r) \\
&= 1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}, \mu = \alpha + \frac{3\eta^2}{3\eta^3}, \mathcal{J}[g_{\mu\nu}] \\
&= \mathcal{J}_{bulk} + \mathcal{J}_{GH} \\
&= \frac{1}{\kappa \int_{\partial\mathcal{M}} d^3\chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{R}\ell}{2} \right)}, \mathcal{E}_t^t - \mathcal{E}_{(\tau|\sigma|\rho)}^{(\varphi|\chi|\psi)} = 0 \Rightarrow 0 = f'' + \frac{\Omega'f'}{\Psi} + 2\eta^2, \mathcal{E}_t^t \\
&+ \mathcal{E}_{(\tau|\sigma|\rho)}^{(\varphi|\chi|\psi)} = 0 \Rightarrow 2\kappa\mathcal{V}(\phi) = -\frac{(f\Omega'' + f'\Omega')}{\Psi^4\eta^4} + \frac{2}{\Omega}, \mathcal{J}_{bulk}^{\mathbb{E}} \\
&= \frac{4\pi\beta}{\eta^3\kappa\ell^2 \left[-\frac{1}{(\chi_\beta-1)^3} + \frac{1}{(\chi_\hbar-1)^3} \right]} = \frac{4\pi\beta}{\kappa\ell^2(r_\beta^3 - r_\hbar^3)}, ds^2 = \hbar_{\alpha\beta}d\chi^\alpha d\chi^\beta \\
&= \Omega(x)[-f(x)dt^2 + d\theta^2 + \sin^2\theta d\varphi^2], \eta_\alpha = \frac{\delta_\alpha^x}{\sqrt{g^{\chi\chi}}\partial_\chi \hbar_{\alpha\beta}}, \mathcal{J}_{GH}^{\mathbb{E}} \\
&= -\frac{2\pi\beta}{\kappa \left[-\frac{6}{\ell^2\eta^4(\chi-1)^3} - \frac{4}{\eta(\chi_\beta-1)} - \left(\alpha + \frac{3\eta^4}{\eta^3} \right) \right] \Big|_{\chi_\beta}} = -\frac{2\pi\beta}{\kappa \left(\frac{6r_\beta^3}{\ell^2} + 4r_\beta - 3\mu \right)}, \mathcal{J}_g^{\mathbb{E}} \\
&= \frac{2\pi\beta}{\kappa \left[\frac{4}{\ell^2\eta^4(\chi_\beta-1)^3} + \frac{4}{\eta(\chi_\beta-1)} - 2\mu \right]} = 2\pi\beta/\kappa \left(\frac{4r_\beta^3}{\ell^2} + 4r_\beta - 2\mu \right)
\end{aligned}$$



$$\begin{aligned}
\mathcal{J}^{\mathfrak{E}} &= \mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{GH}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} = \frac{4\pi\beta}{\kappa\ell^2 \left[\frac{1}{\eta^3(\chi_h - 1)^3} + \frac{\mu\ell^2}{2} \right]} = \frac{4\pi\beta}{\kappa\ell^2 \left(-r_h^3 + \frac{\mu\ell^2}{2} \right)}, \mathcal{J}^{\mathfrak{E}} \\
&= \int_{\partial\mathcal{M}} d^3\chi^3 \sqrt{\hbar^{\mathfrak{E}}} \left(\frac{\langle\varphi|\phi|\psi\rangle\langle\sigma|\tau|\rho\rangle}{2\ell} - \ell_\nu/6\ell\langle\varphi|\phi|\psi\rangle\langle\sigma|\tau|\rho\rangle^4 \right) \\
&= \frac{4\pi\beta}{\kappa \left[-v^2 - \frac{1}{4\ell^2\eta^3(\chi_\beta - 1)} + v^2 - \frac{1}{3\ell^2\eta^3} \right]}, \mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{surf}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} \\
&= -\frac{1}{\mathcal{J} \left(\frac{\Lambda\Gamma}{4\mathfrak{G}} \right)} + \frac{4\pi\beta}{\kappa \left[v^2 - \frac{1}{4\ell^2\eta^3(\chi_\beta - 1)} + 12\eta^2\ell^2 + 4\alpha\ell^2 - 4v^2 + \frac{4}{12\ell^2\eta^3} \right]}, \mathcal{J}^{\mathfrak{E}} \\
&= \beta \left(-\frac{\Lambda\Gamma}{4\mathfrak{G}} + \frac{4\pi}{\kappa} 3\eta^2 + \frac{\alpha}{3\eta^3} \right), \mathcal{M} = \frac{1}{2\mathfrak{G} \left(\alpha + \frac{3\eta^2}{3\eta^3} \right)}, \mathcal{J} \\
&= \frac{\frac{f'(x)|_{\chi=\chi_h}}{4\pi\eta}}{4\pi\eta\Omega(\chi_h) \left[\frac{\alpha}{\eta^2} + 2 + v\chi_h^\nu + 1\frac{\chi_h^\nu}{\chi_h} - 1 \right]}, \delta = \frac{A}{4G} = 4\pi\eta\Omega(\chi_h)/4G
\end{aligned}$$

$$\frac{\frac{\partial\mathcal{M}}{\partial\eta} d\eta}{d\chi_h} = \mathcal{J} \left(\frac{\partial\mathcal{S}}{\partial\chi_h} + \frac{\partial\mathcal{S}}{\partial\eta} \frac{d\eta}{d\chi_h} \right), \mathcal{M} = -\frac{1}{2G \left(\alpha + \frac{3\eta^2}{3\eta^3} \right)}, \mathcal{J} = \frac{1}{4\pi\eta\Omega(\chi_h) \left[\frac{\alpha}{\eta^2} + 2 + v\chi_h^\nu + \frac{1}{\chi_h^\nu} - 1 \right]}$$

2.10.3. Solitón - AdS para agujeros negros cuánticos en espacios curvos.

$$\begin{aligned}
\mathcal{J}[g_{\mu\nu}] &= \int_{\mathcal{M}} d^4\chi(\mathcal{R} - 2\Lambda)\sqrt{-g} + 2 \int_{\partial\mathcal{M}} d^3\chi\sqrt{-\hbar} - \int_{\partial\mathcal{M}} \frac{d^3\chi^4}{\ell\sqrt{-\hbar}}, ds^2 \\
&= -\left(-\frac{\mu_\beta}{r} + \frac{r^2}{\ell^2} \right) dt^2 + \left(-\frac{\mu_\beta}{r} + \frac{r^2}{\ell^2} \right)^{-1} dr^2 + \frac{r^2}{\ell^2} (d\chi_1^2 d\chi_2^2), \mathcal{J}_\beta^\xi \\
&= \frac{2\mathcal{L}\mathcal{L}_\beta\beta_\zeta}{\ell^4 \left(-r_h^3 + \frac{\mu_\beta\ell^2}{2} \right)} = -\frac{\mathcal{L}\mathcal{L}_\beta\beta_\zeta r_h^3}{\ell^4}, \mathfrak{S} = \beta_\zeta^{-1} = \frac{(-g_{tt})'|_{r=r_h}}{4\pi} = \frac{3r_h}{4\pi\ell^2}, \mathfrak{E} = -\frac{\mathfrak{I}^2 \partial\mathcal{J}_\beta^\xi}{\partial\mathfrak{I}} \\
&= \frac{2\mathcal{L}\mathcal{L}_\beta\mu_\beta}{\ell^4}, \mathfrak{S} = -\frac{\partial(\mathcal{J}_\beta^\xi \mathfrak{I})}{\mathfrak{I}} = \frac{\mathcal{L}\mathcal{L}_\beta r_h^2}{4\ell^2\mathfrak{G}} = \mathcal{A}/4G
\end{aligned}$$



$$\begin{aligned}
ds_{dualidad}^2 &= \frac{\ell^2}{\mathcal{R}^2} ds^2 = \gamma_{\alpha\beta} d\chi^\alpha d\chi^\beta = -dt^2 + d\chi_1^2 + d\chi_2^2, \langle \tau_{\alpha\beta}^{dualidad} \rangle \\
&= \lim_{\mathcal{R} \rightarrow \infty} \frac{\mathcal{R}}{\ell} \tau_{\alpha\beta} = \frac{\mu_\beta}{16\pi\mathfrak{G}_N \ell^2} [3\delta_\alpha^0 \delta_\beta^0 + \gamma_{\alpha\beta}], \mathfrak{E} = Q_{\xi t} \\
&= \int d\Sigma^i \tau_{ij} \xi^j = \frac{\mathcal{L}\mathcal{L}_\beta}{\ell^2 \kappa} \left[\mu_\beta + \frac{\ell^2}{4\mathcal{R}} + \mathcal{O}(\mathcal{R}^{-2}) \right], ds^2 \\
&= \frac{r^2}{\ell^2} d\tau^2 + \left(-\frac{\mu_\delta}{r} + \frac{r^2}{\ell^2} \right)^{-1} dr^2 + \left(-\frac{\mu_\delta}{r} + \frac{r^2}{\ell^2} \right) d\theta^2 + \frac{r^2}{\ell^2} d\chi_2^2 - \frac{\mu_\delta}{r} + \frac{r^2}{\ell^2}, \mathcal{L}_\delta \\
&= \frac{4\omega \sqrt{g_{\theta\theta} g_{rr}}}{(g_{\theta\theta})'} \Big|_{r=r_\delta} = \frac{4\pi\ell^2}{3r_\delta}, J_\delta^\xi = \frac{\mathcal{L}\mathcal{L}_\delta \beta_\delta \mu_\delta}{\ell^2} \Delta J = J_\beta^\epsilon - J_\delta^\epsilon \\
&= \frac{\mathcal{L}}{2\pi\kappa\ell^4} \left(\frac{4\pi\ell^2}{3} \right)^3 \mathcal{L}_\beta \beta_\zeta (\mathcal{L}_\delta^{-3} - \xi \beta_\zeta^{-3}) = \frac{\frac{\mathcal{L}}{2\kappa\ell^4} \left(\frac{4\pi\ell^2}{3} \right)^3 \mathcal{L}_\beta \beta_\zeta \left(\frac{1}{\mathcal{L}_\delta^3} - \mathcal{T}^3 \right) \mathcal{A}}{\mathcal{T}\ell^3} \\
&= \mathcal{L}/\ell(4\pi/3)^2 \mathcal{L}_\delta \mathcal{T}
\end{aligned}$$

$$\begin{aligned}
\mathcal{I}[g_{\mu\nu}, \phi] &= \int_{\mathcal{M}} d^4\chi \sqrt{-g} \left[\mathcal{R} - \frac{(\partial\phi)^2}{2} - \mathcal{V}(\phi) \right] + 2 \int_{\partial\mathcal{M}} d^3\chi \mathfrak{K} \sqrt{-h}, \mathcal{V}(\phi) \\
&= \frac{\Lambda(v^2 - 4)}{3v^2 \left[v - \frac{1}{v} + 1e^{-\phi\ell_\nu(v+1)} + v + \frac{1}{v} - 2e^{\phi\ell_\nu(v-1)} + 4v^2 - \frac{1}{v^2} - 4e^{\phi\ell_\nu} \right]} \\
&+ \frac{2\alpha}{v^2 \left[v - \frac{1}{v} + 2 \sin \hbar \phi \ell_\nu (v+1) - v + \frac{1}{v} - 2 \sin \hbar \phi \ell_\nu (v-1) + 4v^2 - \frac{1}{v^2} - 4 \sin \hbar \phi \ell_\nu \right]}
\end{aligned}$$

$$\begin{aligned}
ds^2 &= \frac{\mathfrak{R}^2}{\ell^2 [-dt^2 + d\chi_1^2 + d\chi_2^2] \mathfrak{R}^2} \equiv \frac{1}{\eta^2 (\chi - 1)^2}, ds_{dualidad}^2 = \frac{\ell^2}{\mathfrak{R}^2} ds^2 = \gamma_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -dt^2 + d\chi_1^2 + d\chi_2^2, J_{\mathfrak{E}\mathfrak{S}}^\beta = \beta_\zeta \left(-\frac{\mathcal{A}\mathcal{T}}{4\mathfrak{G}_N} + \frac{2\mathcal{L}\mathcal{L}_\beta}{\ell^2} \alpha \right) = -\frac{\mathcal{L}\mathcal{L}_\beta \alpha \beta_\zeta}{3\ell^2 \eta^3}, \mathcal{A} \\
&= \frac{\mathcal{L}\mathcal{L}_\beta \Omega(\chi\hbar)}{\ell^2}, \mathfrak{A} = \frac{\alpha}{4\pi\eta^3 \Omega}, \mathcal{M}_\beta = \frac{2\mathcal{L}\mathcal{L}_\beta \mu_\beta}{\ell^2}, \mu_\beta = \frac{\alpha}{3\eta^3}, ds^2 \\
&= \Psi_\delta(\gamma) \left[-\frac{d\tau^2}{\ell^2} + \frac{\lambda^2 d\chi^2}{f(x)d\theta^2} + \frac{d\chi_2^2}{\ell^2} \right] \Psi_\delta(\gamma) = \frac{9\gamma^2}{\lambda^2(\gamma^3 - 1)^2}, \mathcal{L}_\delta = \frac{4\pi\lambda}{f'} \Big|_{\gamma=\gamma_\delta} \\
&= \frac{4\pi\lambda^3 \gamma_\delta}{\alpha}
\end{aligned}$$



$$\begin{aligned}
\mathfrak{S}_{soliton}^\varepsilon &= -\frac{\mathcal{L}\beta_\delta\Omega_\delta(\chi_\delta)}{4\ell^2\mathfrak{G}_N} + \frac{2\mathcal{L}\mathcal{L}_\delta\beta_\delta}{\ell^2}\alpha = -\frac{\mathcal{L}\mathcal{L}_\delta\beta_\delta}{\ell^2\left(\frac{\alpha}{3\lambda^3}\right)}, M_{soliton} = -\frac{\mathcal{L}\mathcal{L}_\delta\mu_\delta}{\ell^2}, \mu_\delta = \frac{\alpha}{3\lambda^3}, \mathfrak{S}^\phi \\
&= -\int d^3\chi\sqrt{-\hbar}\left(\frac{\phi^2}{2\ell} - \frac{\ell_\nu}{6\ell\phi^3}\right)\tau_{\alpha\beta}^\phi = -\frac{\frac{2}{\sqrt{-\hbar}}\delta\mathfrak{S}^\phi}{\delta\hbar^{\alpha\beta}}, \tau_{\alpha\beta} \\
&= -\frac{1}{\kappa\left(\kappa_{\alpha\beta} - \hbar_{\alpha\beta}\kappa + \frac{2}{\ell}\hbar_{\alpha\beta} - \mathfrak{I}\mathbb{E}_{\alpha\beta}\right)} - \frac{\hbar_{\alpha\beta}}{\ell\left(\frac{\phi^2}{2} - \frac{\ell_\nu}{6}\phi^3\right)}, \tau_{tt} \\
&= \frac{\alpha(\chi-1)}{3\lambda^2\ell} + \mathcal{O}[(\chi-1)^2], \tau_{\theta\theta} = \frac{2\alpha(\chi-1)}{3\lambda^2\ell} + \mathcal{O}[(\chi-1)^2], \tau_{\chi_2\chi_2} \\
&= -\frac{\alpha(\chi-1)}{3\lambda^2\ell} + \mathcal{O}[(\chi-1)^2], ds_{dualidad}^2 = \frac{\ell^2}{\mathcal{R}^2} ds^2 \\
&= -d\tau^2 + d\theta^2 + d\chi_2^2, \langle \tau_{\alpha\beta}^{dualidad} \rangle \\
&= \lim_{\mathcal{R} \rightarrow \infty} [-1/\lambda\ell(\chi-1)]\tau_{\alpha\beta} = \frac{1}{\ell^2\left(\frac{\alpha}{3\lambda^3}\right)} [-3\delta_\alpha^\theta\delta_\beta^\theta + \gamma_{\alpha\beta}]
\end{aligned}$$

$$\begin{aligned}
\mathcal{M} &= \oint_{\Sigma} d^2\gamma\sqrt{\sigma} m^\alpha\tau_{\alpha\beta}\xi^\beta = \frac{\mathcal{L}\mathcal{L}_\delta f^{\frac{1}{2}}\Psi}{\sqrt{-g_{\tau\tau}}(\partial_\tau)^i\tau_{ij}(\partial_\tau)^j} = -\frac{\mathcal{L}\mathcal{L}_\delta}{\ell^2\left[\frac{\alpha}{3\lambda^3} + \mathcal{O}(\chi-1)\right]}, ds^2 = \sigma_{ij}d\chi^i d\chi^j \\
&= \Omega(x)[f(x)d\theta^2 + d\chi_2^2/\ell^2]
\end{aligned}$$

$$\begin{aligned}
r_\beta^2 &= \frac{\Omega(\chi_\hbar, \eta)}{\ell^2}, r_\delta^2 = \frac{\Omega(\chi_\delta, \lambda)}{\ell^2}, \mathfrak{E} = \mathfrak{M}_{\beta\hbar} - \mathcal{M}_{soliton} = \frac{\mathcal{L}\mathcal{L}_\beta}{\ell^2}(2\mu_\beta + \mu_\delta), \Delta\Gamma = \beta_\zeta^{-1}(\mathcal{J}_{BH}^\varepsilon - \mathfrak{S}_{soliton}^\varepsilon) \\
&= \frac{\mathcal{J}\mathfrak{I}\alpha}{3\ell^2\left(\frac{\mathcal{L}_\delta\beta_\delta}{\lambda^3} - \frac{\mathcal{L}_\beta\beta_\zeta}{\eta^3}\right)\Delta\Gamma} = \frac{4\pi\mathcal{L}\mathcal{L}_\delta}{3\ell^2\left[\frac{\Omega(\lambda, \chi_\delta)}{\mathcal{L}_\delta} - \mathfrak{T}\Omega(\eta, \chi_\hbar)\right]} \\
&= \frac{4\pi\mathcal{L}}{3\ell^2\Omega(\lambda, \chi_\delta)\left(1 - \frac{r_\beta^3}{r_\delta^3}\right)}, \frac{\mathcal{A}}{\mathcal{J}\ell^3} = \frac{\frac{\alpha\mathcal{L}}{4\pi\ell^5}\beta_\zeta^2\mathcal{L}_\delta}{\eta^3} = \frac{\mathfrak{L}\mathcal{L}}{\ell}\left(\frac{\lambda}{\eta}\right), \mathcal{L} \\
&= \frac{16\pi^2}{\alpha^2\ell^4}\left[\frac{9\chi_\hbar^2}{(\chi_\hbar^3-1)^2}\right]^4, \frac{\mathcal{A}}{\mathcal{J}\ell^3} = \frac{\mathfrak{L}\mathcal{L}(\alpha, \ell)}{\ell}r_\beta/r_\delta
\end{aligned}$$

2.11. Modelo Computacional de un agujero negro cuántico en espacios curvos.



$$\begin{aligned}
\delta &= \mathfrak{N} \int dt \mathcal{T}r \left(\frac{1}{2 \sum_j (\mathcal{D}_t \chi_j)^2 - \frac{m^2}{2} \sum_i \chi_i^2 + \frac{\lambda}{4} \sum_{i \neq j} [\chi_i \chi_j]^2} \right), \delta \\
&= \mathfrak{N} \int_0^\beta dt \mathcal{T}r \left(\frac{1}{2 (\mathcal{D}_t \chi_i)^2 + \frac{m^2}{2} \chi_i^2 - \frac{\lambda}{4} [\chi_i \chi_j]^2} \right) \\
\delta &= \int dt \mathcal{T}r \left(\frac{1}{2} (\mathcal{D}_t \chi_i)^2 - \frac{m^2}{2} \chi_i^2 + \frac{g^2}{4} [\chi_i \chi_j]^2 \right), \delta_\kappa \\
&= \mathbb{N}_\alpha \sum_{t=1}^{\eta_t} \mathcal{T}r \left(\frac{1}{2 (\mathcal{D}_t \chi_i)_{i,t}^2 + \frac{m^2}{2} \chi_{i,t}^2 - \frac{\lambda}{4} [\chi_{i,t} \chi_{j,t}]^2} \right), (\mathcal{D}_t \chi_i)_{i,t} \\
&= \frac{1}{\alpha \left(-\frac{1}{2} \mathfrak{U}_t \mathfrak{U}_{t+\alpha} \chi_{i,t+2\alpha} \mathfrak{U}_{t+\alpha}^\dagger \mathfrak{U}_t^\dagger + 2 \mathfrak{U}_t \chi_{i,t+\alpha} \mathfrak{U}_t^\dagger - \frac{3}{2} \chi_{i,t} \right)}, \langle f \rangle \\
&\equiv \frac{\int d\chi d\mathfrak{U} f(\chi, \mathfrak{U}) e^{-\delta_\kappa(\chi, \mathfrak{U})}}{\int d\chi d\mathfrak{U} e^{-\delta_\kappa(\chi, \mathfrak{U})}} = \lim_{\mathfrak{K} \rightarrow \infty} \frac{1}{\mathfrak{K}} \sum_{\kappa=1}^{\mathfrak{K}} f(\chi^{(\kappa)} \mathfrak{U}^{(\kappa)}), \langle \kappa \rangle = \left\langle \frac{1}{2} \sum_i \frac{\chi_i \partial \mathcal{V}}{\partial \chi_i} \right\rangle, \varepsilon \\
&= \left\langle \frac{1}{\beta} \int_0^\beta dt \left(\mathcal{V} + \frac{1}{2} \sum_i \frac{\chi_i \partial \mathcal{V}}{\partial \chi_i} \right) \right\rangle \\
&= \frac{\mathcal{N}}{\beta \int_0^\beta dt \left(m^2 \chi_i^2 - \frac{3\lambda}{4} [\chi_i, \chi_j]^2 \right)}, \varepsilon_{\boxplus} = \left\langle \frac{\mathcal{N}}{\eta_t \sum_{t=1}^{\eta_t} \left(m^2 \chi_{i,t}^2 - \frac{3\lambda}{4} [\chi_{i,t}, \chi_{j,t}]^2 \right)} \right\rangle, \mathcal{F}(\mathcal{J}, \eta_t) \\
&= \xi + \sum_{i=1}^{\eta_\rho} \alpha_i \left(\frac{1}{\mathcal{J} \eta_t} \right), \mathcal{F}(\mathcal{J}, \eta_t) = \xi(\mathfrak{X}) + \sum_i^{\eta_\rho} \alpha_i \left(\frac{1}{\eta_t} \right)^i
\end{aligned}$$

$$\begin{aligned}
\hat{\mathcal{H}} &= \text{Tr} \left(\frac{1}{2} \hat{\rho}_i^2 + \frac{m^2}{2} \hat{\chi}_i^2 - \frac{g^2}{4} [\hat{\chi}_i \hat{\chi}_j]^2 \right), \hat{\rho}_i = \sum_{\alpha=1}^{N^2-1} \hat{\rho}_i^\alpha \tau_\alpha, \hat{\chi}_i = \sum_{\alpha=1}^{N^2-1} \hat{\chi}_i^\alpha \tau_\alpha, [\hat{\chi}_{i\alpha} \hat{\chi}_{j\beta}] = \iota \delta_{j\beta} \delta_{\alpha\beta}, \hat{\mathcal{H}} \\
&= \text{Tr} \left(\frac{1}{2} \hat{\rho}_i^2 - \frac{g^2}{4} [\hat{\chi}_i \hat{\chi}_j]^2 + \frac{g}{2} \hat{\psi} \gamma^\iota [\hat{\chi}_i \hat{\psi}] - \frac{3\iota\mu}{4} \hat{\psi} \hat{\psi} + \frac{\mu^2}{2} \hat{\chi}_i^2 \right) - (N^2 - 1)\mu \\
&= \text{Tr} \left(\frac{1}{2} \hat{\rho}_i^2 - \frac{g^2}{4} [\hat{\chi}_1 \hat{\chi}_2]^2 + \frac{g}{2} \xi [-\hat{\chi}_1 - \iota \hat{\chi}_2 \xi] + \frac{g}{2} \xi^\dagger [-\hat{\chi}_1 - \iota \hat{\chi}_2 \xi^\dagger] + \frac{3\mu}{2} \xi^\dagger \xi \right. \\
&\quad \left. + \frac{\mu^2}{2} \hat{\chi}_i^2 \right) - (N^2 - 1)\mu, \hat{\mathcal{H}} \\
&= \text{Tr} \left(\hat{\rho}_z \hat{\rho}_z^\dagger + \frac{g^2}{4} [\hat{\mathcal{Z}}_\rho \hat{\mathcal{Z}}_\rho^\dagger]^2 - \frac{g}{\sqrt{2}\xi} [\hat{\mathcal{Z}}_\rho^\dagger \xi] - \frac{g}{\sqrt{2}\xi^\dagger} [\hat{\mathcal{Z}} \xi^\dagger] + \frac{3\mu}{2} \xi^\dagger \xi + \mu^2 \hat{\mathcal{Z}} \hat{\mathcal{Z}}^\dagger \right), \hat{\mathcal{H}} \\
&= \sum_{\alpha, \iota} \left(\frac{1}{2\hat{\rho}_{i\alpha}^2} + \frac{m^2}{2} \hat{\chi}_{i\alpha}^2 \right) + \frac{g^2}{4} \sum_{\gamma, J, J} \left\langle \sum_{\alpha, \beta} f_{\alpha\beta\gamma} \hat{\chi}_J^\alpha \hat{\chi}_J^\beta \right\rangle^2, \hat{\mathcal{H}} \\
&= m \sum_{\alpha, \iota} \left(\hat{\eta}_{i\alpha} + \frac{1}{2} \right) + \frac{g^2}{16m^2} \sum_{\gamma, J, J} \left(\sum_{\alpha, \beta} f_{\alpha\beta\gamma} (\hat{a}_{i\alpha} + \hat{a}_{i\alpha}^\dagger) (\hat{a}_{j\beta} + \hat{a}_{j\beta}^\dagger) \right)^2, [\hat{\mathcal{H}}, \hat{\mathbb{G}}_\alpha], \hat{\mathcal{H}}' \\
&= \hat{\mathcal{H}} + c \sum_{\alpha} \hat{\mathbb{G}}_\alpha^2, \hat{\mathcal{H}} \\
&= \sum_{\alpha} \left(\frac{\hat{\rho}_{1\alpha}^2}{2} + \frac{\hat{\rho}_{2\alpha}^2}{2} + \mu^2 \frac{\hat{\chi}_{1\alpha}^2}{2} + \mu^2 \frac{\hat{\chi}_{2\alpha}^2}{2} + \frac{3\mu}{2} \xi_\alpha^\dagger \xi_\alpha \right) \\
&\quad + g^2 \sum_{\alpha \neq \beta} \hat{\chi}_{1\alpha}^2 \hat{\chi}_{2\beta}^2 \\
&\quad - 2g^2 \sum_{\alpha < \beta} \hat{\chi}_{1\alpha} \hat{\chi}_{1\beta} \hat{\chi}_{2\alpha} \hat{\chi}_{2\beta} + \frac{\iota g}{\sqrt{2}} \sum_{\alpha\beta\gamma} \epsilon_{\alpha\beta\gamma} [(-\hat{\chi}_{1\alpha} - \iota \hat{\chi}_{2\alpha}) \xi_\beta^\dagger \xi_\gamma^\dagger \\
&\quad + (-\hat{\chi}_{1\alpha} + \iota \hat{\chi}_{2\alpha}) \xi_\beta \xi_\gamma] - 3\mu
\end{aligned}$$

$$\begin{aligned}
\hat{\mathcal{H}}' &= \hat{\mathcal{H}} + c \sum_{\alpha} \hat{\mathbb{G}}_\alpha^2 + c' (\hat{\mathbb{M}} - \mathfrak{S})^2, \mathfrak{E}_0 \leq \mathfrak{E}_\lambda = \frac{\langle \psi(\theta_i) | \mathcal{H} | \psi(\theta_i) \rangle}{\langle \psi(\theta_i) | \psi(\theta_i) \rangle}, \mathfrak{E}_0 \equiv \langle \psi_\theta | \hat{\mathcal{H}} | \psi_\theta \rangle \\
&= \int d\chi |\psi_\theta(\chi)|^2 \cdot \frac{\langle \chi | \hat{\mathcal{H}} | \psi_\theta \rangle}{\psi_\theta(\chi)} = \mathbb{E}_{\chi \sim |\psi_\theta|^2} [\epsilon_\theta(\chi)], \nabla_\theta \mathfrak{E}_0 \\
&= \mathbb{E}_{\chi \sim |\psi_\theta|^2} [\nabla_{\theta \epsilon_\theta}(\chi)] + \mathbb{E}_{\chi \sim |\psi_\theta|^2} \times [\epsilon_\theta(\chi) \nabla_\theta \ln |\psi_\theta|^2], \theta' = \theta - \beta \nabla_\theta \mathfrak{E}_0
\end{aligned}$$



$$\begin{aligned}
\rho_\theta(\chi) &= \rho(\chi_1; F_\theta^0) \rho[\chi_1; F_\theta^1(\chi_1)] \rho[\chi_3; F_\theta^2(\chi_1, \chi_2)], F_\theta^l \\
&= \Lambda_\theta^{l,m} \odot \tan \hbar \odot \Lambda_\theta^{l,m-1} \odot \tan \hbar \odot \dots \odot \Lambda_\theta^{l,2} \odot \tan \hbar \odot \Lambda_\theta^{l,1}, \Lambda_\theta^{l,\alpha}(\vec{\chi}) \\
&= \mathcal{M}_\theta^{l,\alpha} \vec{\chi} + \vec{\beta}_\theta^{l,\alpha}, \hat{\mathcal{H}}' \\
&= \hat{\mathcal{H}}
\end{aligned}$$

$$\begin{aligned}
+ c \sum_\alpha \hat{\mathcal{G}}_\alpha^2, \hat{\rho}|\psi\rangle &= \int d\mathcal{U}, \hat{\mathcal{U}} \left| \psi \right\rangle, \langle \hat{\mathcal{O}} \rangle_\varsigma = \frac{\langle \psi | \hat{\rho} \hat{\mathcal{O}} | \psi \rangle}{\langle \psi | \hat{\mathcal{O}} | \psi \rangle}, \langle \psi | \hat{\rho} \hat{\mathcal{O}} | \psi \rangle \\
&= \int d\chi \langle \psi | \hat{\rho} | \psi \rangle \langle \psi | \hat{\mathcal{O}} | \psi \rangle = \int d\mathcal{U} d\chi \psi^\odot(\mathcal{U} \chi \mathcal{U}^\dagger) \langle \chi | \hat{\mathcal{O}} | \psi \rangle \\
&= \mathbb{E}_{\mathcal{U}, \chi \sim \langle \psi \rangle^2} \left[\frac{\langle \chi | \hat{\mathcal{O}} | \psi \rangle}{\psi(\chi)} \psi^\odot(\mathcal{U} \chi \mathcal{U}^\dagger) / \psi^\odot(\chi) \right]
\end{aligned}$$

$$\begin{aligned}
\mathfrak{R}_\gamma(\theta) &= \exp\left(-\frac{i\theta}{2}\right) = \begin{pmatrix} \cos \frac{\theta}{2} & \dots & -\sin \frac{\theta}{2} \\ \vdots & \ddots & \vdots \\ \sin \frac{\theta}{2} & \dots & -\cos \frac{\theta}{2} \end{pmatrix}, \hat{\alpha}_l \\
&= \hat{j}_1 \otimes \dots \otimes \hat{j}_{l-1} \otimes \begin{bmatrix} 0 & \dots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 1 \end{bmatrix} \otimes \hat{j}_{l+1} \otimes \dots \otimes \hat{j}_6, \hat{\alpha}_l \\
&= \hat{j}_1 \otimes \dots \otimes \hat{j}_{l-1} \otimes \begin{bmatrix} 0 & 1 & 0 \\ 0 & \sqrt{2} & 1 \\ 1 & 0 & \sqrt{2} \end{bmatrix} \otimes \hat{j}_{l+1} \otimes \dots \otimes \hat{j}_6, c_1 \\
&= \hat{j}_{64} \otimes \begin{pmatrix} 0 & \dots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 1 \end{pmatrix}, c_2 \\
&= \hat{j}_{64} \otimes \begin{pmatrix} 1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & -1 \end{pmatrix} \otimes \begin{pmatrix} 0 & \dots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 1 \end{pmatrix}, c_3 \\
&= \hat{j}_{64} \otimes \begin{pmatrix} 1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & -1 \end{pmatrix} \otimes \begin{pmatrix} 1 & \dots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \dots & -1 \end{pmatrix} \otimes \begin{pmatrix} 0 & \dots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \dots & 0 \end{pmatrix} \\
\hat{\xi}_\alpha &= \left\| \begin{matrix} \sigma_z \otimes \dots & \sigma_z \otimes \\ \otimes \sigma_z & \otimes \end{matrix} \right\|_{\alpha-1} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \otimes \gamma \otimes \dots \otimes \gamma
\end{aligned}$$

$$\hat{\mathcal{G}}_\alpha = \iota \sum_{\beta, \gamma, \iota} f_{\alpha\beta\gamma} \hat{\alpha}_{\iota\beta}^\dagger \hat{\alpha}_{\iota\gamma}, \hat{\mathcal{G}}_\alpha = \left(\otimes_{\iota\beta} |0\rangle_{\iota\beta} \right), \hat{\mathcal{G}}_\alpha = \gamma \sum_{\beta\gamma} f_{\alpha\beta\gamma} \left(\hat{\alpha}_{1\beta}^\dagger \hat{\alpha}_{1\gamma} + \hat{\alpha}_{2\beta}^\dagger \hat{\alpha}_{2\gamma} + \hat{\xi}_\beta^\dagger \hat{\xi}_\gamma \right)$$



$$\begin{aligned}
\hat{a}_{i\alpha}^\dagger &= \sqrt{\frac{m}{2}} \hat{\chi}_{i\alpha} - \frac{i\hat{\rho}_{i\alpha}}{\sqrt{2m}}, \hat{a}_{i\alpha} = \sqrt{\frac{m}{2}} \hat{\chi}_{i\alpha} + \frac{i\hat{\rho}_{i\alpha}}{\sqrt{2m}}, \hat{a}_{i\alpha} |0\rangle_{i\alpha}, |\eta\rangle_{i\alpha} = \frac{(\hat{a}_{i\alpha}^\dagger)^\eta}{\sqrt{\eta!}} |0\rangle_{i\alpha}, |\{\eta_{i\alpha}\}\rangle = \otimes_{i\alpha} |\eta\rangle_{i\alpha}, \hat{a}_\tau^\dagger \\
&= \sum_{\eta=0}^{\Lambda-2} \sqrt{\eta+1} |\eta+1\rangle \langle \eta| = \sum_{\eta=0}^{\Lambda-2} \sqrt{\eta+1} |\eta\rangle \langle \eta+1|, \hat{\eta}_\tau \\
&= \sum_{\eta=0}^{\Lambda-1} \eta |\eta\rangle \langle \eta|, [\hat{\mathfrak{H}}_\tau, \hat{\mathfrak{G}}_\tau], |\eta\rangle = |\beta_0\rangle |\beta_1\rangle \cdots |\beta_{\kappa+1}\rangle, |\eta+1\rangle \langle \eta| \\
&= \otimes_{\ell=0}^{\kappa-1} (|\beta'_\ell\rangle \langle \beta_\ell|), |0\rangle \langle 0| = \boxtimes_2 - \frac{\sigma_{z\blacksquare}}{2}, |1\rangle \langle 1| \\
&= \boxtimes_2 - \frac{\sigma_{z\blacksquare}}{2}, |0\rangle \langle 1| = \sigma_x + \frac{i\sigma_y}{2}, |1\rangle \langle 0| = \sigma_x + \frac{i\sigma_y}{2}
\end{aligned}$$

$$\begin{aligned}
\hat{\mathcal{M}} &= \sum_\alpha \left(i(\hat{Z}_\alpha \hat{\rho}_{z\alpha}^\dagger - \hat{\rho}_{z\alpha} \hat{Z}_\alpha^\dagger) - \frac{1}{2} \hat{\xi}_\alpha^\dagger \hat{\xi}_\alpha \right), \hat{Q} = -\sqrt{2} \hat{\xi}_\alpha^\dagger (\hat{\rho}_{z\alpha} - \mu \hat{Z}_\alpha) - \frac{g}{\sqrt{2} f_{\alpha\beta\gamma} \hat{\xi}_\alpha \hat{Z}_\beta \hat{Z}_\gamma^\dagger}, \hat{Q}^2 \\
&= -i \hat{Z}_\alpha \hat{\mathfrak{G}}_\alpha, \hat{Q}^{\dagger 2} = i \hat{Z}_\alpha^\dagger \hat{\mathfrak{G}}_\alpha, \{\hat{Q}, \hat{Q}^\dagger\} = 2(\hat{\mathcal{H}} - \mu \hat{\mathcal{M}}), \hat{Q} |\mathfrak{B}\mathfrak{P}\mathfrak{S}\rangle = \hat{Q}^\dagger |\mathfrak{B}\mathfrak{P}\mathfrak{S}\rangle \\
&= (\hat{\mathcal{H}} - \mu \hat{\mathcal{M}}) |\mathfrak{B}\mathfrak{P}\mathfrak{S}\rangle
\end{aligned}$$

$$\begin{aligned}
\delta_\epsilon &= [\hat{\mathfrak{Q}}\epsilon^* + \hat{\mathfrak{Q}}^\dagger \epsilon], \hat{\mathfrak{Q}} = -\hat{\xi}_\alpha^\dagger [(\hat{\rho}_1^\alpha - i\hat{\rho}_2^\alpha) - \mu(\hat{\chi}_1^\alpha - i\hat{\chi}_2^\alpha)] - \frac{ig}{\sqrt{2} f_{\alpha\beta\gamma} \hat{\xi}_\alpha \hat{\chi}_1^\beta i\hat{\chi}_2^\gamma}, \hat{Z} = \hat{\chi}_1 - \frac{i\hat{\chi}_2}{\sqrt{2}}, \hat{\rho}_z \\
&= \hat{\rho}_1 - \frac{i\hat{\rho}_2}{\sqrt{2}}, [\hat{Z}, \hat{\rho}_z^\dagger] = [\hat{\rho}_z, \hat{Z}^\dagger] = i
\end{aligned}$$

$$\begin{aligned}
Z(\mathcal{T}) &= \text{Tr}_{\mathcal{H}_{inv}} \left(e^{-\frac{\hat{H}}{t}} \right), Z(\mathcal{T}) = \frac{1}{\text{vol}(G)} \int_G^\delta dg \text{Tr}_{\mathcal{H}_{ext}} \left(\hat{g} e^{-\frac{\hat{H}}{t}} \right), \hat{\mathcal{P}} \equiv \frac{1}{\text{vol}(G)} \int_G^\delta dg \hat{g}, |\Phi\rangle_{inv} \\
&= \frac{1}{\sqrt{\mathcal{C}_\phi}} \times \frac{1}{\text{vol}(G)} \int_G^\delta dg (\hat{g}|\Phi\rangle), \mathcal{C}_\phi = \frac{1}{[\text{vol}(G)]^2} \int_G^\delta dg \int_G^\delta dg' \langle \phi | \hat{g}^{-1} \hat{g}' | \phi \rangle \\
&= \frac{1}{\text{vol}(G)} \int_G^\delta dg \langle \phi | \hat{g} | \phi \rangle \\
&= \text{vol} \left(\frac{G_\phi}{\text{vol}(G)} \right) (G), \text{Tr}_{\mathcal{H}_{inv}} \left(e^{-\frac{\hat{H}}{t}} \right) \sum_\phi \text{vol} \left(\frac{G_\phi}{\text{vol}(G)} \right) (G)_{inv} \langle \phi | e^{-\frac{\hat{H}}{t}} | \phi \rangle_{inv} \\
&= \sum_\phi \text{vol} \left(\frac{G_\phi}{\text{vol}(G)} \right) (G) \otimes \frac{1}{\mathcal{C}_\phi} \frac{1}{[\text{vol}(G)]^2} \int_G^\delta dg \int_G^\delta dg' \langle \phi | \hat{g}^{-1} e^{-\frac{\hat{H}}{t}} \hat{g}' | \phi \rangle \\
&= \sum_\phi \frac{1}{\text{vol}(G)} \int_G^\delta dg \text{Tr}_{\mathcal{H}_{ext}} \left(\hat{g} e^{-\frac{\hat{H}}{t}} \right), Z(\mathcal{T}) \\
&= \frac{1}{[\text{vol}(G)]^\kappa} \int \left(\prod_{k=1}^\kappa d\mathbf{u}_{(\mathfrak{K})} \right) \text{Tr}_{\mathcal{H}_{ext}} \otimes \left(\hat{\mathbf{u}}_{(\mathcal{J})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathbf{u}}_{(k-1)}^{-1} \hat{\mathbf{u}}_{(k-1)} \otimes e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \right) \\
&= \frac{1}{[\text{vol}(G)]^\kappa} \int \left(\prod_{k=1}^\kappa d\mathbf{u}_{(\mathfrak{K})} \right) \int \left(\prod_{k=1}^\kappa d\chi_{(\mathfrak{K})} \right) \langle \chi_{(k)} | \hat{\mathbf{u}}_{(\mathcal{J})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathbf{u}}_{(k-1)}^{-1} | \chi_{(k-1)} \rangle \\
&\quad \otimes \langle \chi_{(k-1)} | \hat{\mathbf{u}}_{(\mathcal{J}-1)} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathbf{u}}_{(k-2)}^{-1} | \chi_{(k-2)} \rangle \\
&\quad \otimes \dots \otimes \langle \chi_{(1)} | \hat{\mathbf{u}}_{(1)} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathbf{u}}_{(k-1)}^{-1} | \chi_{(k)} \rangle \\
\langle \chi_{(k)} | \hat{\mathbf{u}}_{(\mathcal{J})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathbf{u}}_{(k-1)}^{-1} | \chi_{(k-1)} \rangle &= \langle \mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1} | \hat{\mathbf{u}}_{(\mathcal{J})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} | \mathbf{u}_{(k-1)} \chi_{(k-1)} \mathbf{u}_{(k-1)}^{-1} \rangle \\
&= \int_G^\delta d\mathcal{P} \langle \mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1} | \hat{\mathcal{P}}_{(\mathcal{J})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} | \mathcal{P} \rangle \otimes \langle \emptyset | \hat{\mathcal{P}}_{(\mathcal{J})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} | \mathbf{u}_{(k-1)} \chi_{(k-1)} \mathbf{u}_{(k-1)}^{-1} \rangle \\
&= \int_G^\delta d\wp e^{\mathcal{T} \text{Tr} \left[\left(\mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1} - \mathbf{u}_{(k-1)} \chi_{(k-1)} \mathbf{u}_{(k-1)}^{-1} \right) \right]} \otimes e^{-\mathcal{H} \left(\mathcal{P}, \frac{\mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1}}{(\mathcal{T}\mathcal{K})} \right)} \\
&= e^{\mathcal{K} \mathcal{T} \text{Tr} \left[\left(\mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1} - \mathbf{u}_{(k-1)} \chi_{(k-1)} \mathbf{u}_{(k-1)}^{-1} \right)^2 \right]} \otimes e^{-\mathcal{V} \left(\frac{\mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1}}{(\mathcal{T}\mathcal{K})} \right)} \simeq e^{-\frac{\mathcal{V} \left[\frac{\mathcal{V} t \left(\mathbf{u}_{(k)} \chi_{(k)} \mathbf{u}_{(k)}^{-1} - \mathbf{u}_{(k-1)} \chi_{(k-1)} \mathbf{u}_{(k-1)}^{-1} \right)}{\mathfrak{I}\mathfrak{K}} \right]}{\tau\kappa}} = e^{-\frac{\mathcal{L}[\mathcal{D}\chi_{\mathcal{K}}, \chi_{\mathcal{K}}]}{\tau\kappa}}
\end{aligned}$$



$$\begin{aligned}
\mathfrak{U}_{(k)}^{-1}\mathfrak{U}_{(k)} &\equiv e^{\frac{\Lambda_{(k)}}{KT}}\chi_{(k)} - (\mathfrak{U}_{(k-1)}\mathfrak{U}_{(k)}^{-1})^{-1}\chi_{(k-1)}(\mathfrak{U}_{(k-1)}\mathfrak{U}_{(k)}^{-1}) \\
&\simeq \frac{Dt\chi_{(k)}}{KT}, Z(T) \int [d\Lambda][d\chi] e^{-\int dt\mathfrak{L}[Dt\chi,\chi]} = (\mathfrak{U}_k\mathfrak{U}_{k-1}^{-1})(\mathfrak{U}_{k-1}\mathfrak{U}_{k-2}^{-1})\otimes(\mathfrak{U}_2\mathfrak{U}_1^{-1}) \\
&= \wp e^{\int_0^1 dt\Lambda_t}
\end{aligned}$$

3. Gravedad cuántica en espacios curvos.

$$\begin{aligned}
d\hat{s}^2 &= \hat{g}_{\mu\nu}d\chi^\mu d\chi^\nu = -\left(\kappa + \frac{r^2}{\ell^2}\right)dt^2 + \frac{dr^2}{\kappa} + \frac{r^2}{\ell^2} + r^2 d\Sigma_\kappa^2, \mathfrak{S}\mathfrak{h}_{\alpha\beta}d\chi^\alpha d\chi^\beta \\
&= \frac{r^2}{\ell^2}(-dt^2 + \ell^2 d\Sigma_\kappa^2), ds^2 = -N(r) dt^2 + \mathcal{H}(r) dr^2 + \delta(r)d\Sigma_\kappa^2, \mathcal{V}(\phi) \\
&= -\frac{3}{\kappa\ell^2} - \frac{\phi^2}{\ell^2} + \mathcal{O}(\phi)^4, \phi(r) = \frac{\alpha}{r} + \frac{\beta}{r^2} + \mathcal{O}(r^{-3}), N(r) = -g_{tt} \\
&= \frac{r^2}{\ell^2} + \kappa - \frac{\mu}{r} + \mathcal{O}(r^{-2}), \delta(r) \\
&= r^2 + \mathcal{O}(r^{-2}), N\delta'^2\mathcal{H} - 2N\delta''\mathcal{H}\delta + (N\mathcal{H})'\delta'\delta - 2\kappa N\mathcal{H}\delta^2\phi'^2, \mathcal{H}(r) = g_{rr} \\
&= \frac{\ell^2}{r^2} + \frac{\ell^4}{r^4}\left(-\kappa - \frac{\alpha^2\kappa}{2\ell^2}\right) + \frac{\ell^5}{r^5}\left(\frac{\mu}{\ell} - \frac{4\kappa\alpha\beta}{3\ell^3}\right) + \mathcal{O}(r^{-6}), g_{rr} \\
&= \frac{\ell^2}{r^2} + \frac{\alpha\ell^4}{r^4} + \frac{\beta\ell^5}{r^5} + \mathcal{O}(r^{-6}), \mathcal{V}(\phi) = -\frac{3}{\kappa\ell^2} - \frac{\phi^2}{\ell^2} + \lambda\phi^3 + \mathcal{O}(r^4), \phi(r) \\
&= \frac{\alpha}{r} + \frac{\beta}{r^2} + \frac{\gamma\ln(r)}{r^2} + \mathcal{O}(r^{-3}), \mathcal{H}(r) = g_{rr} \\
&= \frac{\ell^2}{r^2} + \frac{\ell^4}{r^4}\left(-\kappa - \frac{\alpha^2\kappa}{2\ell^2}\right) + \frac{\ell^5}{r^5}\left(\frac{\mu}{\ell} - \frac{4\kappa\alpha\beta}{3\ell^3} + \frac{2\kappa\alpha\gamma}{9\ell^3}\right) + \frac{\ell^5\ln(r)}{r^5\left(-\frac{4\kappa\alpha\gamma}{3\ell^3}\right)} \\
&+ \mathcal{O}\left[\ln\frac{(r)^2}{r^6}\right], \mathcal{H}(r) = \frac{\ell^2}{r^2} + \frac{\ell^4\alpha}{r^4} + \frac{\ell^5\beta}{r^5} + \frac{\ell^5 c \ln r}{r^5} + \mathcal{O}\left[\ln\frac{(r)^2}{r^6}\right], \alpha = -\kappa - \frac{\alpha^2\kappa}{2\ell^2}, \beta \\
&= \frac{\mu}{\ell} - \frac{4\kappa\alpha\beta}{3\ell^3} + \frac{2\kappa\alpha\gamma}{9\ell^3}, c \\
&= -\frac{4\kappa\gamma\alpha}{3\ell^3}, \partial_r\left(\frac{\phi'\delta\sqrt{N}}{\sqrt{H}}\right) - \frac{\delta\sqrt{HN}\partial\mathfrak{B}}{\partial\phi} 3\alpha^2\ell^2\lambda + \frac{\gamma}{\ell^2} + \mathcal{O}(r^{-1}), \xi^r \\
&= r\eta^r(\chi^m) + \mathcal{O}(r^{-1})\xi^m = \mathcal{O}(1), \phi'(\chi) = \phi(\chi) + \xi^\mu\partial_\mu\phi(\chi) \\
&= \frac{\alpha'}{r} + \frac{\beta'}{r^2} + \gamma'\ln\frac{r}{r^2} + \mathcal{O}(r^{-3}), \alpha' = \alpha - \eta^r\alpha + \xi^m\partial_m\alpha, \beta' \\
&= \beta - \eta^r(2\beta - \gamma)\xi^m\partial_m\beta, \gamma' = \gamma - 2\gamma\eta^r + \xi^m\partial_m\gamma
\end{aligned}$$



$$\frac{\alpha' \partial \gamma}{\partial \alpha} - \gamma' \alpha \frac{\partial \gamma}{\partial \alpha} - \gamma + \eta^r \left(2\gamma - \alpha \frac{\partial \gamma}{\partial \alpha} \right) + \xi^m \left(\frac{\partial \alpha}{\partial \chi_m} \frac{\partial \gamma}{\partial \alpha} - \frac{\partial \gamma}{\partial \chi_m} \right)$$

$$\frac{\alpha' \partial \beta}{\partial \alpha} - \beta' \alpha \frac{\partial \beta}{\partial \alpha} - \beta + \eta^r \left(2\beta - \gamma - \alpha \frac{\partial \beta}{\partial \alpha} \right) + \xi^m \left(\frac{\partial \alpha}{\partial \chi_m} \frac{\partial \beta}{\partial \alpha} - \frac{\partial \beta}{\partial \chi_m} \right)$$

$$\mathcal{J}_{\mathfrak{E}\mathfrak{I}} \rightarrow \mathcal{J}_{\mathfrak{E}\mathfrak{I}} - \int d^3 \chi \mathfrak{B}[\mathcal{O}(\chi)]$$

$$g_g^{ct} = - \frac{1}{\kappa \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{R}\ell}{2} \right)}, \mathfrak{S}$$

$$= \frac{\int d^4 \chi \sqrt{-g}(\mathcal{R} - \frac{(\partial \phi)^2}{2} - \mathcal{V}(\phi))}{2\kappa} + 1/\kappa \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} \mathfrak{K}$$

$$- 1/\kappa \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{R}\ell}{2} \right) + \mathcal{J}_{\varphi\phi\psi}, \mathcal{J}_{\varphi\phi\psi}^{ct}$$

$$= 1/6\kappa \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} (\phi \eta^\nu \partial_\nu \phi - \phi^2 \ell / 2 \varphi \psi \tau), \mathcal{J}_{\varphi\phi\psi}$$

$$= - \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} \left[\frac{\phi^2}{2\ell} + \frac{\mathfrak{B}(\alpha)}{\ell \alpha^3 \tau^3} - \langle \varphi | \phi | \psi \rangle^3 \right], \delta \ell$$

$$= \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} \left[\frac{1}{r} \left(\sqrt{-g^{rr}} \phi' - \frac{\varphi \psi}{\tau \ell} - \frac{3\mathcal{W}(\alpha)\phi^2}{\ell \phi^3} \right) \left(1 + \frac{1}{\frac{rd^2 \mathcal{W}(\alpha)}{d(\alpha)^2}} \right) \right.$$

$$\left. + \left(\frac{3\mathcal{W}(\alpha)}{\alpha} - \beta \right) \langle \varphi | \phi | \psi \rangle^3 / \ell \alpha^3 \right] \delta \alpha$$

$$\mathcal{J}_{\varphi\phi\psi} = - \int_{\partial \mathcal{M}}^\delta d^3 \chi \sqrt{-\hbar} \left[\frac{\phi^2}{2\ell} + \frac{\langle \varphi | \phi | \psi \rangle^3}{\ell \alpha^3 \left(\mathcal{W} - \frac{\alpha \gamma}{3} \right)} - \frac{\langle \varphi | \phi | \psi \rangle^3 \mathfrak{C}_\gamma}{3\ell} \ln \left(\frac{\langle \varphi | \phi | \psi \rangle}{\alpha} \right) \right]$$

$$ds^2 = \Omega(\chi) = \left[-f(\chi) dt^2 + \frac{\eta^2 d\chi^2}{f(\chi)} + d\Sigma_\kappa^2 \right], \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{E}\mathfrak{I}} = \int_0^{\frac{1}{\tau}} d\tau \int_{\chi_+}^{\chi_\beta} d^3 \chi \sqrt{g^{\mathfrak{E}\mathfrak{I}}}(\phi)$$

$$= \frac{\sigma_\kappa}{2\eta\kappa\tau} d(\Omega f) / d\chi |_{\chi_+}^{\chi_\beta}$$



$$\Omega(\chi) \rightarrow \delta(r), f(x) \rightarrow \frac{N(r)}{\delta(r)}, dx \rightarrow \frac{\sqrt{NH}}{\eta\delta} dr$$

$$\begin{aligned} \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} &= \frac{\frac{\sigma_{\kappa}}{2\kappa\tau} \delta}{\sqrt{NH}} dN \Big|_{r_+}^{r_{\beta}} ds^2 =' \hbar_{\alpha\beta} d\chi^{\alpha} d\chi^{\beta} = \Omega(\chi_0) [-f(\chi_0) dt^2 + d\Sigma_{\kappa}^2] \eta_{\alpha} = \frac{\delta_{\alpha}^{\chi}}{\sqrt{g^{xx}}} \Big|_{\chi=\chi_0} \mathfrak{K}_{\alpha\beta} \\ &= \frac{\sqrt{g^{xx}}}{2} \partial_{\chi} g_{\alpha\beta} \Big|_{\chi=\chi_0} \mathfrak{K} = \frac{1}{2\eta \left(\frac{f}{\Omega}\right)^{-\frac{1}{2}} \left[\frac{(f\Omega)'}{\Omega f} + \frac{2\Omega'}{\Omega}\right] \Big|_{\chi_0}} \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} \\ &= -\frac{\frac{\sigma_{\kappa}}{\kappa\tau} \Omega f \left[\frac{(f\Omega)'}{\Omega f} + \frac{2\Omega'}{\Omega}\right] \Big|_{\chi_{\beta}}}{2\eta} = -\frac{\sigma_{\kappa}}{2\tau\kappa} \left(\frac{\delta}{\sqrt{NH}} \frac{dN}{dr} + \frac{2N}{\sqrt{NH}} \frac{d\delta}{dr} \right) \Big|_{r_{\beta}} \mathcal{J}_g^{ct} \\ &= \frac{2\sigma_{\kappa}}{\kappa\tau\ell} \left(\Omega^{\frac{3}{2}} f^{\frac{1}{2}} + \frac{\ell^2\kappa}{2} f^{\frac{1}{2}} \Omega^{\frac{1}{2}} \right) \left(\Omega^{\frac{3}{2}} f^{\frac{1}{2}} + \frac{\ell^2\kappa}{2} f^{\frac{1}{2}} \Omega^{\frac{1}{2}} \right) \Big|_{\chi_{\beta}} = \frac{2\sigma_{\kappa}}{\kappa\tau\ell} \delta\sqrt{N} \left(1 + \frac{\ell^2\kappa}{2\delta} \right) \Big|_{r_{\beta}} \mathcal{J} \\ &= \frac{N'}{4\pi\sqrt{NH}} \Big|_{r_+} \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + \mathcal{J}_g^{ct} \\ &= \frac{1}{\tau \left[\frac{\sigma_{\kappa} \delta(r_+) \tau}{4\mathfrak{G}} \right]} - \frac{\sigma_{\kappa}}{2\kappa\tau \left[\frac{2N}{\sqrt{NH}} \frac{d\delta}{dr} - \frac{4}{\ell\delta\sqrt{N} \left(1 + \frac{\ell^2\kappa}{2\delta} \right)} \right] \Big|_{r_{\beta}}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} \\ &\quad + \mathcal{J}_g^{ct} = -\frac{\mathcal{A}}{4\mathfrak{G}} - \frac{\sigma_{\kappa}}{\tau} \left(-\frac{\mu}{\kappa} + \frac{4\alpha\beta}{3\ell^2} + \frac{r\alpha^2}{2\ell^2} \right) \Big|_{r_{\beta}} \mathcal{J}_{\langle\varphi|\phi|\psi\rangle}^{ct} \\ &= \int_{\partial\mathcal{M}} d^3\chi \sqrt{\hbar\varepsilon} \left[\frac{\phi^2}{2\ell} + \frac{\mathcal{W}(\alpha)}{\ell\alpha^3 \langle\varphi|\phi|\psi\rangle^3} \right] = \frac{\sigma_{\kappa}}{\tau} \left(\frac{\mathcal{W}}{\ell^2} + \frac{\alpha\beta}{\ell^2} + \frac{r\alpha}{2\ell^2} \right) \Big|_{r_{\infty}} \mathcal{J}^{\mathfrak{E}} \\ &= \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + \mathcal{J}_g^{ct} + \mathcal{J}_{\langle\varphi|\phi|\psi\rangle}^{ct} = \frac{\mathcal{A}}{4\mathfrak{G}} + \frac{\sigma_{\kappa}}{\tau} \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{\alpha}{3} \frac{d\mathcal{W}}{d\alpha} \right) \right] \Big|_{\mathcal{F}} = \mathcal{J}^{\mathfrak{E}\mathcal{J}} \\ &= \mathcal{M} - \mathcal{J}\mathcal{S}, \mathcal{M} = -\frac{\mathcal{J}^2 \partial\mathcal{J}^{\mathfrak{E}}}{\partial\mathcal{J}} = \sigma_{\kappa} \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{\alpha}{3} \frac{d\mathcal{W}}{d\alpha} \right) \right] \Big|_{\mathcal{S}} = -\frac{\partial(\mathcal{J}^{\mathfrak{E}\mathcal{J}})}{\partial\mathcal{J}} = \frac{\mathcal{A}}{4\mathfrak{G}} \end{aligned}$$



$$\begin{aligned} \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{R}\mathfrak{K}}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + \mathcal{J}_g^{ct} + \mathcal{J}_{\langle\varphi|\phi|\psi\rangle}^{ct} &= -\frac{\mathcal{A}}{4\mathfrak{G}} + \frac{\sigma\kappa}{\mathcal{T} \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left[\mathcal{W}(\alpha) - \frac{\alpha}{3d\mathcal{W}} + \frac{2\alpha\gamma}{9} - \frac{\alpha\gamma}{3} \ln r \right] \right]} \langle \hat{j}_{\langle\varphi|\phi|\psi\rangle}^{ct} \rangle \\ &= \int_{\partial\mathcal{M}}^{\delta} d^3\chi \sqrt{\hbar\varepsilon} \left\{ \frac{\phi^3\gamma}{3\alpha^2\ell} [\ln(\alpha/\phi) - 1] \right\} = \sigma_{\kappa}/\mathfrak{I} \left[-\frac{\alpha\gamma}{3\ell^2} + \frac{\alpha\gamma \ln r}{3\ell^2} + \mathcal{O}(r^{-1} \ln r) \right] \end{aligned}$$

$$\begin{aligned} \mathcal{J}^{\mathfrak{E}} = \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{R}\mathfrak{K}}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + \mathcal{J}_g^{ct} + \mathcal{J}_{\langle\varphi|\phi|\psi\rangle}^{ct} + \hat{\mathcal{J}}_{\langle\varphi|\phi|\psi\rangle}^{ct} &= -\frac{\mathcal{A}}{4\mathfrak{G}} + \frac{\sigma\kappa}{\mathcal{T} \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W}(\alpha) - \frac{\alpha}{3d\mathcal{W}} - \frac{\alpha\gamma}{9} \right) \right]}, \mathcal{M} \\ &= -\gamma^2 \frac{\partial \mathcal{J}^{\mathfrak{E}}}{\partial \mathcal{T}} = \sigma\kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W}(\alpha) - \frac{\alpha}{3d\mathcal{W}} - \frac{\alpha\gamma}{9} \right) \right] \delta = -\frac{\partial(\mathcal{J}^{\mathfrak{E}}\mathcal{T})}{\partial \mathcal{T}} = \frac{\mathcal{A}}{4\mathfrak{G}} \end{aligned}$$

3.1. Tensor de stress Brown-York en espacios cuánticos curvos.

$$\begin{aligned} \tau_{\alpha\beta} &= -\frac{1}{\kappa \left(\kappa_{\alpha\beta} - \hbar_{\alpha\beta}\kappa + \frac{2}{\ell\hbar_{\alpha\beta}} - \iota\mathfrak{G}_{\alpha\beta} \right)} - \frac{\hbar_{\alpha\beta}}{\ell \left[\frac{\phi^2}{2} + \frac{\mathcal{W}(\alpha)}{\alpha^3\phi^3} \right] \tau_{tt}} \\ &= \frac{\ell}{\mathcal{R} \left[\frac{\mu}{8\pi\mathfrak{G}\ell^2} + \frac{1}{\ell^4 \left(\omega(\alpha) - \frac{\alpha\beta}{3} \right)} \right]} + \mathcal{O}(\mathfrak{R}^{-2})\tau_{\theta\theta} \\ &= \frac{\ell}{\mathcal{R} \left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2 \left(\omega(\alpha) - \frac{\alpha\beta}{3} \right)} \right]} + \mathcal{O}(\mathfrak{R}^{-2})\tau_{\langle\varphi|\phi|\psi\rangle} \\ &= \ell \sin^2\theta/\mathcal{R} \left[\frac{\mu}{16\pi\mathfrak{G}} - 1/\ell^2 \left(\omega(\alpha) - \frac{\alpha\beta}{3} \right) \right] + \mathcal{O}(\mathfrak{R}^{-2}) \\ \langle \tau_{\alpha\beta}^{dualidad} \rangle &= \frac{3\mu}{16\pi\mathfrak{G}\ell^2\delta_{\alpha}^0\delta_{\beta}^0} + \frac{\gamma_{\alpha\beta}}{\ell^2 \left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2 \left(\omega(\alpha) - \frac{\alpha\beta}{3} \right)} \right] \langle \tau^{dualidad} \rangle} = -3/\ell^4 \left[\omega(\alpha) - \frac{\alpha\beta}{3} \right] \end{aligned}$$



$$\begin{aligned}
\tau_{\alpha\beta} &= -\frac{1}{\kappa\left(\kappa_{\alpha\beta} - \hbar_{\alpha\beta}\kappa + \frac{2}{\ell\hbar_{\alpha\beta}} - i\mathfrak{G}_{\alpha\beta}\right)} - \frac{\hbar_{\alpha\beta}}{\ell\left[\frac{\phi^2}{2} + \frac{\phi^3}{\alpha^3}\left(\omega - \frac{\alpha\gamma}{3}\right) + \frac{\phi^3\gamma}{3\alpha^2}\left(\frac{\alpha}{\phi}\right)\right]}\tau_{tt} \\
&= \frac{\ell}{\mathcal{R}\left[\frac{\mu}{8\pi\mathfrak{G}\ell^2} + \frac{1}{\ell^4\left(\omega - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9}\right)}\right]} + \mathcal{O}\left[\frac{(\ln \mathfrak{R})^3}{\mathfrak{R}^2}\right]\tau_{\theta\theta} \\
&= \frac{\ell}{\mathcal{R}\left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2\left(\omega - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9}\right)}\right]} + \mathcal{O}\left[\frac{(\ln \mathfrak{R})^3}{\mathfrak{R}^2}\right]\tau_{\langle\phi|\phi|\psi\rangle} \\
&= \frac{\ell\sin^2\theta}{\mathcal{R}\left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2\left(\omega - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9}\right)}\right]} + \mathcal{O}\left[\frac{(\ln \mathfrak{R})^3}{\mathfrak{R}^2}\right]\langle\tau_{\alpha\beta}^{dualidad}\rangle \\
&= \frac{3\mu}{16\pi\mathfrak{G}\ell^2\delta_{\alpha}^0\delta_{\beta}^0} + \frac{\gamma_{\alpha\beta}}{\ell^2\left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2\left(\omega(\alpha) - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9}\right)}\right]}\langle\tau^{dualidad}\rangle \\
&= -3/\ell^4\left[\left(\omega - \frac{\alpha\beta}{3} - \frac{\alpha\gamma}{9}\right)\lambda\right]
\end{aligned}$$

$$\begin{aligned}
\tau^{\alpha\beta} &\equiv \frac{2}{\sqrt{-\hbar}} \frac{\delta \mathcal{J}}{\delta \hbar_{\alpha\beta}}, \mathcal{J}_{\mathfrak{G}\mathfrak{S}} + \mathcal{J}_g + \mathcal{J}_\phi \\
&= \frac{1}{\kappa \int d^n \chi \sqrt{-\hbar \mathfrak{K}} - \frac{1}{\kappa \int d^n \chi \sqrt{-\hbar \left[\frac{\eta-1}{\ell} + \frac{\ell \mathcal{R}}{2(\eta-2)} \right]} - \int d^n \chi \sqrt{-\hbar} \psi}, \tau^{\alpha\beta} \\
&= -\frac{1}{\kappa \left(\mathcal{K}_{\alpha\beta} - \hbar_{\alpha\beta} \mathfrak{K} + \frac{\eta-1}{\ell} \hbar_{\alpha\beta} - \frac{\ell}{\eta} - 2\mathfrak{G}_{\alpha\beta} \right)} - \hbar_{\alpha\beta} [\psi], \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= N(\mathcal{R}) dt^2 + \delta(\mathcal{R}) d\Sigma_\kappa^2, \mathfrak{G}_{\alpha\beta} = \mathcal{R}_{\alpha\beta} - \frac{1}{2\mathfrak{R} \hbar_{\alpha\beta}}, \mathfrak{G}_{tt} = \frac{(\eta-2)(\eta-1)}{2} \frac{\kappa \mathcal{N}}{\mathcal{S}}, \mathfrak{G}_{ij} \\
&= -\frac{(\eta-2)(\eta-3)}{2} \kappa v_{ij}, \tau_{tt} = -\frac{(\eta-1)}{\kappa} \left[\frac{N\delta'}{2\delta\sqrt{\mathcal{H}}} - \frac{\mathcal{N}}{\ell} \left(1 + \frac{\ell^2 \kappa}{2\delta} \right) \right] + N[\psi], \tau_{ij} \\
&= \frac{v_{ij}}{\kappa} \left[\frac{\delta}{2\sqrt{\mathcal{H}} \left(\frac{\mathfrak{R}'}{\mathfrak{R}} + \frac{\mathcal{S}'}{\mathcal{S}} (\eta-2) \right)} - \frac{(\eta-1)\mathcal{S}}{\ell} - \frac{\ell \mathfrak{K}(\eta-3)}{2} \right] - v_{ij} \delta[\psi], \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= \mathcal{L}^2 dt^2 + \sigma_{ij} (d\gamma^i + \mathcal{L}^i dt) (d\gamma^j + \mathcal{L}^j dt), \mathfrak{E} = \mathcal{Q}_{\frac{\partial}{\partial t}} \\
&= \frac{\oint_\Sigma^\delta d^{D-2} \gamma \sqrt{\sigma} \mu^\alpha \tau_{\alpha\beta} \xi^\beta}{\sqrt{\mathfrak{R}}} = \left(\oint_\Sigma^\delta d^2 \gamma \sqrt{v} \right) \delta^{\left(\frac{D-2}{2}\right)} \tau_{tt} = \frac{\sigma_{\kappa, \eta-1} \delta^{\left(\frac{D-2}{2}\right)}}{\sqrt{\mathfrak{R}}} \tau_{tt}, \sigma_{ij} d\chi^i d\chi^j \\
&= \delta d\Sigma_\kappa^2, \tau_{tt} = -\frac{(\eta-1)}{\kappa \left[\frac{f^{\frac{3}{2}} \Omega'}{2\eta\sqrt{\Omega}} - \frac{\Omega f}{\ell} \left(1 + \frac{\ell^2 \kappa}{2\Omega} \right) \right]} + \frac{\Omega f}{\kappa} [\psi], \tau_{ij} \\
&= \frac{v_{ij}}{\kappa \left[\frac{(\Omega f)'}{2\eta\sqrt{\Omega f}} + \frac{(\eta-2)\Omega' \sqrt{f}}{2\eta\sqrt{\Omega}} - \frac{(\eta-1)\Omega}{\ell} - \frac{\ell \kappa (\eta-3)}{2} \right]} - \frac{v_{ij} \Omega}{\kappa} [\psi], \mathfrak{E} = \mathcal{Q}_{\frac{\partial}{\partial t}} \\
&= \oint_\Sigma^\delta d^{D-2} \gamma \sqrt{\sigma} \mu^\alpha \tau_{\alpha\beta} \xi^\beta = \frac{\sigma_{\kappa, \eta-1} \Omega^{\left(\frac{D-2}{2}\right)}}{\sqrt{\Omega f}} \tau_{tt}
\end{aligned}$$

3.2. Masa Hamiltoniana en espacios cuánticos curvos.

$$-\frac{5}{4\ell^2} > m^4 \geq -\frac{9}{4\ell^2}, \phi(r) = \frac{\alpha}{r} + \frac{\beta}{r^2} + \mathcal{O}(r^{-3})$$

3.3. Modelos Logarítmicos y Anti-logarítmicos en espacios cuánticos curvos.



$$\begin{aligned}
\delta\mathcal{M}_\phi &= \frac{\sigma\kappa}{\kappa \left[r\delta\alpha + \ell\delta\beta + \mathcal{O}\left(\frac{1}{r}\right) \right]} \delta\mathcal{M}_\phi = \frac{\sigma\kappa}{\ell^2 \left[r\alpha\delta\alpha + \alpha\delta\beta + 2\beta\delta\alpha + \mathcal{O}\left(\frac{1}{r}\right) \right]} \delta\mathcal{M} \\
&= \frac{\sigma\kappa}{\kappa\ell^2 \left[r(\ell^2\delta\alpha + \kappa\alpha\delta\alpha) + \ell^3\delta\beta + \kappa(\alpha\delta\beta + 2\beta\delta\alpha) + \mathcal{O}\left(\frac{1}{r}\right) \right]} + \alpha + \frac{\alpha^2}{2\ell^2}, \delta\mathcal{M} \\
&= \frac{\sigma\kappa}{\kappa\ell^2 [\ell^3\delta\beta + \kappa(\alpha\delta\beta + 2\beta\delta\alpha)]} \mathcal{M} \\
&= \frac{\sigma\kappa \left[\frac{\ell\beta}{\kappa} + \frac{1}{\ell^2 \left(\frac{\alpha d\omega(\alpha)}{d\alpha} + \omega(\alpha) \right)} \right] \ell c}{\kappa} - 4\alpha^3\lambda, \delta\mathcal{M}_\phi \\
&= \left\{ \frac{\ell\delta\beta}{\kappa} + \frac{\delta\alpha}{\kappa} r + \frac{\ell\delta c}{\kappa} \ln(r) + \frac{\mathcal{O}(\ln(r)^2)}{r} \right\} \sigma\kappa, \delta\mathcal{M}_\phi \\
&= \left[\alpha\delta\beta + 2\beta\delta\alpha + \frac{3\alpha^2\ell^2\lambda\delta\alpha}{\ell^2} + \frac{r\alpha\delta\alpha}{\ell^2} - 12\lambda\alpha^2\delta\alpha \ln(r) + \mathcal{O}\left(\ln\frac{(r)^2}{r}\right) \right] \sigma\kappa, \delta\mathcal{M} \\
&= \left[\frac{\ell\delta\beta}{\kappa} + \alpha\delta\beta + 2\beta\delta\alpha + \frac{3\alpha^2\ell^2\lambda\delta\alpha}{\ell^2} \right] \sigma\kappa, \mathcal{M} \\
&= \left[\frac{\ell\beta}{\kappa} + \frac{1}{\ell^2 \left(\frac{\alpha d\omega}{d\alpha} + \omega(\alpha) + \alpha^3\ell^2\lambda \right)} \right] \sigma\kappa, \mathcal{M} \\
&= \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2 \left(\omega(\alpha) - \frac{1}{3\alpha} \frac{d\mathcal{W}}{d\alpha} + \frac{1}{3\alpha^3\ell^2\lambda} \right)} \right] \sigma\kappa, \mathcal{W}(\alpha) = \alpha^3[C + \ell^2\lambda \ln(\alpha)]
\end{aligned}$$

3.4. Masa Holográfica y Masa Hamiltoniana en espacios cuánticos curvos.

$$\begin{aligned}
d\Sigma_\kappa^2 &= \frac{d\gamma^2}{1} - \kappa\gamma^2 + \frac{(1 - \kappa\gamma^2)d\langle\phi|\phi|\psi\rangle^2}{\|\tau\sigma\rho\|^2\delta\xi}, \mathcal{E} = \int d\sigma^i\tau_{ij}\xi^j \int d\gamma d\langle\phi|\phi|\psi\rangle d\|\tau\sigma\rho\| \delta\mu^i\tau_{ij}\xi^j, ds^2 \\
&= \sigma_{ij}d\chi^i d\chi^j = \delta d\Sigma_\kappa^2, \mathcal{E} = \sigma\kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\omega - \frac{\alpha}{3} \frac{d\omega}{d\alpha} \right) \right], \mathcal{E} \\
&= \sigma\kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{1}{3\alpha} \frac{d\mathcal{W}}{d\alpha} - \frac{\alpha\gamma}{9} \right) \right] = \sigma\kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{1}{3\alpha} \frac{d\mathcal{W}}{d\alpha} - \frac{\alpha^3 C\gamma}{9} \right) \right]
\end{aligned}$$



$$\begin{aligned}
& \mathcal{V}(\phi) \\
&= \frac{\Lambda(v^2 - 4)}{6\kappa v^2 \left[v - \frac{1}{v} + 2e^{-\phi \ell v(v+1)} + v + \frac{1}{v} - 2e^{\phi \ell v(v-1)} + 4v^2 - \frac{1}{v^2} - 4e^{-\phi \ell v} \right]} \\
&+ \frac{\Upsilon}{\kappa v^2 \left[v - \frac{1}{v} + 2 \sinh \phi \ell v(v+1) - v + \frac{1}{v} - 2 \sinh \phi \ell v(v-1) + 4v^2 - \frac{1}{v^2} - 4 \sinh \phi \ell v \right]} \phi(\chi) \\
&= \ell_v^{-1} \sqrt{\frac{v^2 - 1}{2\kappa}} \ln \chi, f(x) = \frac{1}{\ell^2} + \Upsilon \left[\frac{1}{v^2} - 4 - \chi^2/v^2 \left(1 + \frac{\chi^{-v}}{v} - 2 - \frac{\chi^v}{v} + 2 \right) \right] + \frac{\chi}{\Omega(\chi)}, \Omega(\chi) \\
&= r^2 + \mathcal{O}(r^{-3}), \chi = 1 + \frac{1}{\eta r} + \frac{m^4}{r^3} + \frac{\eta}{r^4} + \frac{\rho}{r^5} + \mathcal{O}(r^{-6}), \Omega(\chi) \\
&= r^2 - 24m^4\eta^4 + v^2 - \frac{1}{12\eta} - 24m^4\eta^4 - v^2 + \frac{1}{12\eta^3 r} + 720m^4\eta^4 - 480|\rho\eta|^5 + v^4 - 20v^2 \\
&+ \frac{19}{240\eta^4 r^2} + \mathcal{O}(r^{-3}), \chi = 1 + \frac{1}{\eta r} - \frac{(v^2 - 1)}{23\eta^3 r^3 \left[1 - \frac{1}{\eta r} - \frac{9(v^2 - 9)}{80\eta^4 r^4} \right]} + \mathcal{O}(r^{-6}), -g_{tt} = f(x)\Omega(\chi) \\
&= \frac{r^2}{\ell^2} + 1 + \Upsilon + \frac{3\eta^4}{3\eta^3 r} + \mathcal{O}(r^{-3}), g_{rr} = \frac{\Omega(\chi)\eta^2}{f(x) \left(\frac{d\chi}{dr} \right)} \\
&= \frac{\ell^2}{r^2} - \frac{\ell^4}{r^4} - \frac{\ell^2(v^2 - 1)}{4\eta^4 r^2} - \frac{\ell^2(3\eta^2 \ell^2 + \Upsilon \ell^2 - v^2 + 1)}{3\eta^3 r^5} + \mathcal{O}(r^{-6}), \phi(\chi) = \ell_v^{-1} \sqrt{\frac{v^2 - 1}{2\kappa}} \ln \chi \\
&= \frac{1}{\ell_v \eta r} - \frac{1}{2\ell_v \eta^4 r^2} - v^2 - \frac{9}{24\eta^3 r^5} + \mathcal{O}(r^{-4}), \mathfrak{M} = \sigma \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{\alpha}{3} \frac{d\mathcal{W}}{d\alpha} \right) \right] \mathfrak{M} \\
&= -\frac{\sigma}{\kappa} \left(3\eta^2 + \frac{\Upsilon}{3\eta^4} \right)
\end{aligned}$$

$$J_{\mathbb{C}\mathfrak{S}\mathfrak{I}} \rightarrow J_{\mathbb{C}\mathfrak{S}\mathfrak{I}} + \ell_v/6 \int d^3 \chi \mathcal{O}^3$$

3.5. Acción bulk on – shell en espacios cuánticos curvos.



$$\begin{aligned}
\mathcal{J} &= \int d^{\eta+1}\chi \sqrt{-g} \left[\frac{\mathcal{R}}{2\kappa} - \frac{(\partial\phi)^2}{2} - \mathcal{V}(\phi) \right] + 1/\kappa \int_{\partial\mathcal{M}}^{\delta} d^{\eta}\chi \sqrt{-\hbar} \mathfrak{K} + \mathcal{J}_g + \mathcal{J}_{\phi}, \mathcal{G}_{\mu\nu} = \kappa \mathcal{T}_{\mu\nu}, \mathcal{G}_{\mu\nu} \\
&= \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathcal{R}}, \mathcal{T}_{\mu\nu} = \partial_{\mu}\phi\partial_{\nu}\phi - g_{\mu\nu} \left[\frac{(\partial\phi)^2}{2} + \mathcal{V} \right] \mathcal{G} = -\frac{\mathcal{R}(\eta-1)}{2}, \mathcal{T} \\
&= -(\eta-1) \left[\frac{(\partial\phi)^2}{2} + \frac{\mathcal{V}(\eta+1)}{(\eta-1)} \right] \mathcal{G}_{\mu\nu} = \kappa \mathcal{T}_{\mu\nu} \rightarrow \frac{\mathfrak{R}}{2\kappa} = \frac{(\partial\phi)^2}{2} + \frac{\mathcal{V}(\eta+1)}{(\eta-1)}, \mathcal{J}_{bulk}^{\xi} \\
&= -\frac{2}{\eta} - 1 \int d^{\eta+1}\chi \sqrt{g^{\xi}} \mathcal{V}(\phi)
\end{aligned}$$

3.6. Sistema de Coordenadas en espacios cuánticos curvos.

$$ds^2 = \Omega(x) \left[-f(x)dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\Sigma_{\kappa}^2 \right]$$



$$\begin{aligned}
\mathfrak{E}_t^t - \mathfrak{E}_x^x = 0 &\Rightarrow 2\kappa\phi'^2 = \mathfrak{D} - \frac{2}{2\Omega^2[3(\Omega')^2 - 2\Omega\Omega'']}, \mathfrak{E}_t^t = -\frac{1}{\mathfrak{D}} - 2g^{\alpha\beta}\xi_{\alpha\beta} = 0 \\
\Rightarrow f'' + \mathfrak{D} - \frac{2}{2\Omega}\Omega'f' + 2\kappa\eta^2\mathfrak{E}_t^t + \frac{1}{\mathfrak{D}} - 2g^{\alpha\beta}\xi_{\alpha\beta} &= 0 \Rightarrow 2\kappa\mathcal{V} \\
= -\mathfrak{D} - \frac{2}{2\eta^2\Omega^2\left[f\Omega'' + \mathfrak{D} - \frac{4}{2\Omega}f(\Omega')^2 + \Omega'f'\right]} + \frac{\kappa\left(\mathfrak{D} - \frac{2}{\Omega}\right)d}{d\chi\left[\Omega^{\frac{(\mathfrak{D}-2)}{2}}f'\right]} + 2\eta^2\kappa\Omega^{\frac{(\mathfrak{D}-2)}{2}} - \frac{2\eta^2\Omega^{\frac{\mathfrak{D}}{2}}(2\kappa\mathcal{V})}{\mathfrak{D}} - 2 \\
= f\Omega''\Omega^{\frac{(\mathfrak{D}-4)}{2}} + \Omega'\left(f\Omega^{\frac{(\mathfrak{D}-4)}{2}}\right) - 2\eta^2\kappa\Omega^{\frac{(\mathfrak{D}-2)}{2}}, 2\kappa\mathcal{V} &= -\frac{(\mathfrak{D}-2)}{2\eta^2\Omega^{\frac{\mathfrak{D}}{2}}\left[\Omega^{\frac{(\mathfrak{D}-4)}{2}}(f\Omega)'\right]'} d\Sigma_\kappa^2 \\
= v_{ij}d\chi^i d\chi^j, \mathcal{J}_{bulk}^\mathfrak{E} &= \frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta\left[\Omega^{\frac{(\mathfrak{D}-4)}{2}}(f\Omega)'\right]_{\chi_h}^{\chi_\beta} \hbar_{\alpha\beta}d\chi^\alpha d\chi^\beta} = \Omega(x)[-f(x)dt^2 + d\Sigma_\kappa^2], \eta_\alpha \\
= \frac{\delta_\alpha^\chi}{\sqrt{g^{xx}}}, \mathcal{K}_{\alpha\beta} &= \frac{\sqrt{g^{xx}}}{2}\partial_\chi\hbar_{\alpha\beta}, \mathcal{K} = \frac{1}{2\eta}\left(\frac{f}{\Omega}\right)^{\frac{1}{2}}\left[\frac{(\Omega f)'}{\Omega f} + \frac{(\mathfrak{D}-2)\Omega'}{\Omega}\right], \mathcal{J}_{\mathfrak{E}\mathfrak{D}}^\mathfrak{E} \\
= -\frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta\Omega^{\frac{(\mathfrak{D}-2)}{2}}f\left[\frac{(f\Omega)'}{f\Omega} + \frac{(\mathfrak{D}-2)\Omega'}{\Omega}\right]}, \mathcal{J}_g \\
= -\frac{1}{\kappa\int d^\eta\chi\sqrt{-\hbar}\left[\frac{(\eta-1)}{\ell} + \frac{\ell\mathcal{R}}{2(\eta-2)} + \frac{\ell^3}{2(\eta-4)(\eta-2)^2\left(\mathcal{R}^{\alpha\beta}\mathfrak{R}_{\alpha\beta} - \frac{\eta\mathcal{R}^2}{4(\eta-1)}\right)}\right]}, \mathcal{R}_{ij} \\
= \frac{(\eta-2)\kappa}{\Omega}\sigma_{ij}, \mathcal{R} &= \frac{\kappa(\eta-2)(\eta-1)}{\Omega}, \mathcal{R}^{\alpha\beta}\mathfrak{R}_{\alpha\beta} = \frac{(\eta-2)^2(\eta-1)\kappa^2}{\Omega^2}\mathcal{R}^{\alpha\beta}\mathfrak{R}_{\alpha\beta} - \frac{\eta\mathcal{R}^2}{4(\eta-1)} \\
= \frac{\kappa^2}{4\Omega^2}(\eta-2)^2(\eta-1)(\eta-4)\mathcal{J}_g^\mathfrak{E} &= \frac{\beta\sigma_{\kappa,\eta-1}(\mathfrak{D}-2)}{\kappa\ell}\sqrt{\Omega^{\mathfrak{D}-1}}f\left(1 + \frac{\ell^2\kappa}{2\Omega} - \frac{\ell^4\kappa^2}{8\Omega^2}\right)\beta^{-1} = \mathcal{J} \\
= \frac{f'}{4\pi\eta|_{\chi_h}}, \delta &= \mathcal{A}/4\mathfrak{E}
\end{aligned}$$

$$\begin{aligned}
\mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} &= \frac{1}{\mathcal{T} \left(\frac{\mathcal{A}\mathcal{T}}{4\mathfrak{G}} \right)} - \frac{\sigma_{\mathcal{D}-2,\kappa}}{2\kappa\mathcal{T}} \Omega^{\frac{(\mathcal{D}-2)}{2}} (\mathcal{D}-2) \left[\frac{f\Omega'}{\eta\Omega} - \frac{\sqrt{\Omega}f}{2\ell \left(1 + \frac{\ell^2\kappa}{2\Omega} - \frac{\ell^4\kappa^2}{8\Omega^2} \right)} \right]_{\chi_\beta} ds^2 \\
&= -N(r)dt^2 + H(r)dr^2 + \delta(r)d\Sigma_{\kappa,\Omega}^2(x) \rightarrow \delta(r), f(x) \rightarrow \frac{N(r)}{\delta(r)}, \frac{\sqrt{NH}}{\eta\delta} dr \\
&\rightarrow d\chi, 2\kappa\mathcal{V} = -\mathcal{D} - \frac{2}{2\eta^2\Omega^{\frac{\mathcal{D}}{2}} \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} (f\Omega)' \right]} \rightarrow 2\kappa\mathcal{V} = -\mathcal{D} - \frac{\frac{2}{2\delta^{\frac{(\mathcal{D}-2)}{2}} \sqrt{NH}}}{dr \left(\frac{\delta^{\frac{(\mathcal{D}-2)}{2}}}{\sqrt{NH}} \frac{dN}{dr} \right)} \mathcal{J}_{bulk}^{\mathfrak{E}} \\
&= \frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} (f\Omega)' \right]_{\chi_h}} \rightarrow \mathcal{J}_{bulk}^{\mathfrak{E}} = \frac{\frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa} \frac{dN}{dr} \delta^{\frac{(\eta-1)}{2}}}{\sqrt{NH}} \Bigg|_{r_h}^{r_\beta} \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -N(\mathcal{R})dt^2 + \mathcal{S}(\mathcal{R})d\Sigma_{\kappa,\eta}^2 = \frac{\delta_\mu^r}{\sqrt{g^{rr}}} \mathcal{K}_{\mu\nu} = \frac{\sqrt{g^{rr}}}{2} \partial_r \hbar_{\mu\nu} \mathcal{K} \\
&= \frac{1}{2\sqrt{H}} \left[\frac{N'}{N} + \frac{(\eta-1)\delta'}{\delta} \right] \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} = -\frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta\Omega^{\frac{(\mathcal{D}-2)}{2}} f \left[\frac{(f\Omega)'}{f\Omega} + \frac{(\mathcal{D}-2)\Omega'}{\Omega} \right]_{\chi_\beta}} \mapsto \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} \\
&= -\frac{\frac{\sigma_{\kappa,\eta-1}}{2\kappa\mathcal{T}} \delta^{\frac{(\mathcal{D}-2)}{2}}}{\sqrt{NH} \left[\frac{dN}{dr} + \frac{(\mathcal{D}-2)}{\delta} \frac{N}{dr} d\delta \right]_{r_\beta}} \mathcal{J}_g^{\mathfrak{E}} = \frac{\beta\sigma_{\kappa,\eta-1}(\eta-1)}{\ell\kappa\Omega^{\frac{(\mathcal{D}-1)}{2}} f^{\frac{1}{2}} \left(1 + \frac{\ell^2\kappa}{2\Omega} - \frac{\ell^4\kappa^2}{8\Omega^2} \right)_{\chi_\beta}} \mapsto \mathcal{J}_g^{\mathfrak{E}} \\
&= \beta\sigma_{\kappa,\eta-1}(\eta-1)/\ell\kappa\delta^{\frac{(\mathcal{D}-2)}{2}} N^{1/2} \left(1 + \frac{\ell^2\kappa}{2\delta} - \frac{\ell^4\kappa^2}{8\delta^2} \right)_{r_\beta} \beta^{-1} = \mathcal{T} = \frac{N'}{4\pi\sqrt{NH}} \Big|_{r_h}, \delta \\
&= \mathcal{A}/4\mathfrak{G} \\
\mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} &= -\frac{1}{\mathcal{T} \left(\frac{\mathcal{A}\mathcal{T}}{4\mathfrak{G}} \right)} - \frac{\sigma_{\mathcal{D}-2,\kappa}}{2\kappa\mathcal{T}} \delta^{\frac{(\mathcal{D}-2)}{2}} (\mathcal{D}-2) \left[\frac{N\delta'}{\delta\sqrt{NH}} - \frac{2\sqrt{N}}{\ell \left(1 + \frac{\ell^2\kappa}{2\delta} - \frac{\ell^4\kappa^2}{8\delta^2} \right)} \right]_{r_\beta}
\end{aligned}$$

3.7. Ecuaciones de Movimiento en espacios cuánticos curvos.



$$\begin{aligned}
2\kappa\phi'^2 &= \frac{(\mathcal{D}-2)}{2\Omega^2[3(\Omega')^2 - 2\Omega\Omega'']} \frac{2\kappa\phi'^2}{(\mathcal{D}-2)} = \frac{1}{2\delta^2[\delta'^2 - 2\delta\delta'']} + \frac{\delta'}{2\delta} \frac{(\text{NH})'}{\text{NH}}, \frac{d}{d\chi} \left[\delta^{\frac{(\mathcal{D}-2)}{2}} df/d\chi \right] \\
&= -2\eta^2\kappa\delta^{\frac{(\mathcal{D}-2)}{2}}, \frac{d}{dr} \left[\frac{\delta^{\frac{\mathcal{D}}{2}}\sqrt{\text{NH}}d}{dr} \left(\frac{N}{\delta} \right) \right] = -2\kappa\sqrt{\text{NH}}\delta^{\frac{(\mathcal{D}-4)}{2}}, 2\kappa\mathcal{V} \\
&= \frac{(\mathcal{D}-2)}{2\eta^2\Omega^{\frac{\mathcal{D}}{2}} \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} (f\Omega)' \right]}, 2\kappa\mathcal{V} = -\frac{\frac{(\mathcal{D}-2)}{2\delta^{\frac{(\mathcal{D}-2)}{2}}\sqrt{\text{NH}}d}}{\frac{d}{dr} \left(\frac{\delta^{\frac{(\mathcal{D}-2)}{2}}}{\sqrt{\text{NH}}} \frac{dN}{dr} \right)}, \partial_\chi \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} f\phi' \right] \\
&= \frac{\eta^2\Omega^{\frac{\mathcal{D}}{2}}\partial\mathcal{V}}{\partial\phi}, \partial_r \left(\delta^{\frac{(\mathcal{D}-2)}{2}} \phi' \sqrt{\frac{N}{H}} \right) = \sqrt{\text{NH}}\delta^{\frac{(\mathcal{D}-2)}{2}} \partial\mathcal{V}/\partial\phi
\end{aligned}$$

3.8. Cálculo de Potencial on – shell en espacios cuánticos curvos.



$$\begin{aligned}
\eta^2 \mathcal{V}(\phi) &= -\frac{f\Omega''}{f^2\Omega^2} - \frac{f'\Omega'}{\eta^2\Omega^2}, \eta^2 \mathcal{V}(\phi) \\
&= -\frac{1}{\chi^{2\nu-2}\nu^4 \left[\frac{1}{\ell^2} + \frac{\alpha}{2} \left(\chi^{2+\nu} - \frac{1}{2} + \nu + \frac{\chi^{2\nu-2}}{2} + \nu - \chi^2 + 1 \right) \right]} \\
&\quad \left[\begin{aligned} &\chi^{\nu-1}(\nu-1)^2\nu^2 - \frac{\chi^{\nu-1}(\nu-1)\nu^2}{\chi^2(\chi^\nu-1)^2\eta^2} \\ &-4\chi^{2\nu-1}(\nu-1)\nu^3 + \frac{2\chi^{2\nu-1}\nu^4}{(\chi^\nu-1)^4\eta^2\chi^2} + 2\chi^{2\nu-1}\nu^3 - \frac{4\chi^{2\nu-1}(\nu-1)\nu^3}{\eta^2\chi^2(\chi^\nu-1)^3} \end{aligned} \right] (\chi^\nu-1)^4\eta^4 \\
&\quad - \frac{1}{2\chi^{2\nu-2}\nu^4 \left(\frac{\chi^{\nu-1}(\nu-1)\nu^2}{\chi\eta^2(\chi^\nu-1)^2} - \frac{2\chi^{2\nu-1}\nu^3}{\chi\eta^2(\chi^\nu-1)^3} \right)} (\chi^{1+\nu} + \chi^{1-\nu} - 2\chi)\mathcal{V}(\phi) = \frac{f(x)\nu^2}{\chi^{2\nu-2}\nu^4} \\
&\quad \left(\chi^{3\nu-3}\nu^2 + 4\chi^{2\nu-3}\nu^2 + \chi^{\nu-3}\nu^2 + 3\chi^{3\nu-3}\nu - \frac{3\chi^{3\nu-3}}{(\chi^\nu-1)^4\eta^2} \right) (\chi^\nu-1)^4\eta^2 - \frac{f(x)\nu^2}{\chi^{2\nu-2}\nu^4} \\
&\quad \left(2\chi^{3\nu-3} - 4\chi^{2\nu-3} + \frac{2\chi^{\nu-3}}{(\chi^\nu-1)^4\eta^2} \right) (\chi^\nu-1)^4\eta^2 - \frac{\alpha\eta^4(\chi^\nu-1)^4}{2\chi^{2\nu-2}\nu^4} \\
&\quad \left(\frac{-\nu^2}{\eta^2(\chi^\nu-1)^3} \right) (\chi^{2\nu-2}(\nu+1) + \chi^{\nu-2}(\nu-1))(\chi^{1+\nu} + \chi^{1-\nu} - 2\chi)\mathcal{V}(\phi) = -\frac{f(x)\chi^{-\nu}}{\chi\nu^2} \\
&\quad \left(\chi^{2\nu}(\nu+1)(\nu+2) + 4\chi^\nu(\nu^2-1) + (\nu-1)(\nu-2) \right) + \frac{\frac{\alpha(\chi^\nu-1)}{2\nu^2\chi^{\nu-1}}\chi^{2\nu-2}}{\chi^{\nu-1}} \\
&\quad \frac{(\chi^{1+\nu} + \chi^{1-\nu} - 2\chi)(\chi^{-\nu}(\nu-1) + (\nu+1))\mathcal{V}(\phi)f(x)}{2\chi\nu^2} \\
&\quad \left(2\chi^\nu(\nu+1)(\nu+2) + 8(\nu^2-1) + 2\chi^{\nu-1}(\nu-1)(\nu-2) \right) + \frac{\alpha}{2\nu^2\chi^{\nu-1}(\chi^\nu-1)} \\
&\quad (1 + \nu + \chi^{-\nu}(\nu-1))(1 + \chi^{2\nu} - 2\chi^\nu)\mathcal{V}(\phi) = -\frac{f(x)}{2\chi\nu^2} \\
&\quad \left(2\chi^\nu(\nu+1)(\nu+2) + 8(\nu^2-1) + 2\chi^{\nu-1}(\nu-1)(\nu-2) \right) - \frac{\alpha}{2\nu^2\chi^{\nu-1}}(\chi^\nu-1)^2 \\
&\quad (2 - \chi^\nu(\nu+1) + \chi^{\nu-1}(\nu-1))
\end{aligned}$$



$$\mathcal{V}(\phi)$$

$$= - \frac{e^{-\ell_v \phi}}{2v^2 \left[\frac{1}{\ell^2} + \frac{\alpha}{2 \left(e^{(2+v)\ell_v \phi} - \frac{1}{2} + v + e^{(2-v)\ell_v \phi} - \frac{1}{2} - v - \chi^2 + 1 \right)} \right]} \left[8(v^2 - 1) + 2(v+1)(v+2)e^{v\ell_v \phi} + 2(v-1)(v-2)e^{-v\ell_v \phi} \right]$$

$$= - \frac{\alpha e^{\ell_v \phi}}{2v^2 \left(\exp \frac{v\ell_v \phi}{2} - \exp -\frac{v\ell_v \phi}{2} \right)^2 [2 - e^{v\ell_v \phi}(v+1) + e^{-v\ell_v \phi}(v-1)]}, \mathcal{V}(\phi)$$

$$= - \frac{(v^2 - 4)}{\ell^2 v^2 \left[\frac{(v-1)}{(v+2)e^{-v\ell_v \phi}(v+1)} + \frac{(v+1)}{(v-2)e^{v\ell_v \phi}(v-1)} + \frac{4(v^2 - 1)}{(v^2 - 4)e^{-\ell_v \phi}} \right]}$$

$$= - \frac{\alpha}{2v^2 \left[\frac{e^{(v+2)\ell_v \phi}}{(2+v)} - \frac{e^{(2-v)\ell_v \phi}}{v} - 2 - e^{2\ell_v \phi} + \frac{v^2}{v^2 - 4} \right]} \left[4(v^2 - 1)e^{-\ell_v \phi} + (v+1)(v+2)e^{(v-1)\ell_v \phi} + (v-1)(v-2)e^{(v+1)v\ell_v \phi} \right] - \frac{\alpha}{2v^2}$$

$$\left(e^{v\ell_v \phi} - 2 + e^{-v\ell_v \phi} \right) [2e^{\ell_v \phi} - e^{(v+1)\ell_v \phi}(v+1) + e^{(1-v)\ell_v \phi}(v-1)]$$

$$\mathcal{V}(\phi) = \mathcal{V}_\Lambda(\phi) - \frac{\alpha}{2(v^2 - 4)}$$

$$\left[\begin{aligned} &v^2(e^{\ell_v \phi}(v-1) - e^{-v\ell_v \phi}(v+1) - e^{\ell_v \phi}(v-1) + e^{-v\ell_v \phi}(v+1) + 4e^{-\ell_v \phi} - 4e^{\ell_v \phi}) + 3v \\ &(e^{\ell_v \phi}(v-1) + e^{\ell_v \phi}(v+1) - e^{-\ell_v \phi}(v-1) - e^{-\ell_v \phi}(v+1)) + 2 \\ &(e^{\ell_v \phi}(v-1) - e^{\ell_v \phi}(v+1) - e^{-\ell_v \phi}(v-1) + e^{-\ell_v \phi}(v+1)) - 4e^{-\ell_v \phi} + 4e^{\ell_v \phi} \end{aligned} \right]$$

$$\mathcal{V}(\phi) = \mathcal{V}_\Lambda(\phi) - \frac{\alpha}{2(v^2 - 4)}$$

$$[(2v^2 + 6v + 4)\sinh \ell_v \phi(v-1) - (2v^2 - 6v + 4)\sinh \ell_v \phi(v+1) + 8(1 - v^2)\sinh \ell_v \phi]$$

$$\mathcal{V}(\phi) = \frac{\Lambda(v^2 - 4)}{3v^2 \left[\frac{(v-1)}{(v+2)e^{-\ell_v \phi}(v+1)} + \frac{(v+1)}{(v-2)e^{\ell_v \phi}(v-1)} + \frac{4(v^2 - 1)}{(v^2 + 4)e^{-\ell_v \phi}} \right] + \alpha}$$

$$\left[\frac{(v-1)}{(v+2)\sinh \ell_v \phi(v+1)} - \frac{(v+1)}{(v-2)\sinh \ell_v \phi(v-1)} + \frac{4(v^2 - 1)}{(v^2 + 4)e^{-\ell_v \phi}} \right]$$

3.9. Tensores diferenciales para espacios cuánticos curvos.



$$\begin{aligned}
d\chi^\mu \wedge d\chi^\nu &= d\chi^\mu \otimes d\chi^\nu - d\chi^\nu \otimes d\chi^\mu, d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_\rho} = \sum_{\sigma} (-1)^{|\sigma|} d\chi^{\sigma(\mu_1)} \otimes \dots \otimes d\chi^{\sigma(\mu_\rho)}, \mathfrak{S} \\
&= \frac{1}{\rho!} \mathcal{H}_{\mu_1 \dots \mu_\rho} d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_\rho}, \mathcal{A} = \Lambda_\mu d\chi^\mu, \mathcal{F} \\
&= \frac{1}{2\mathcal{F}_{\mu\nu} d\chi^\mu} \wedge d\chi^\nu \star (d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_\rho}) \\
&= \frac{\sqrt{-g}}{(\mathcal{D} - \rho)! \epsilon^{\mu_1 \dots \mu_\rho \nu_{\rho+1} \dots \nu_{\mathcal{D}}} d\chi^{\nu_{\rho+1}} \wedge \dots \wedge d\chi^{\nu_{\mathcal{D}}}, d\mathcal{V} = d^{\mathcal{D}} \chi \sqrt{-g}} \\
&= \frac{\sqrt{-g}}{\mathcal{D}!} \epsilon_{\mu_1 \dots \mu_{\mathcal{D}}} d\chi^{\nu_1} \wedge \dots \wedge d\chi^{\nu_{\mathcal{D}}}
\end{aligned}$$

4. Integral de Yang – Mills en espacios cuánticos curvos.

$$\begin{aligned}
& \exp \left(-\frac{1}{2} \int_{\mathbb{C}^4} d\lambda_4 |\kappa \mathfrak{d}A + A \wedge A|^2 \right) \mathfrak{D}\mathfrak{A} \\
&= \exp \left(-\frac{1}{2} \int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \mathfrak{d}A, A \wedge A \rangle + \langle A \wedge A, \kappa \mathfrak{d}A \rangle \right. \\
&\quad \left. + |A \wedge A|^2 \right) \exp \left(\int_{\mathbb{C}^4} d\lambda_4 |\kappa \mathfrak{d}A + A \wedge A|^2 \right) \mathfrak{D}\mathfrak{A} \\
& \exp \left(-\frac{1}{2} \int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \mathfrak{d}A, A \wedge A \rangle + \langle A \wedge A, \kappa \mathfrak{d}A \rangle + |A \wedge A|^2 \right) \int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \mathfrak{d}A, A \wedge A \rangle \int_{\mathbb{C}^4} d\lambda_4 \langle A \wedge A, \kappa \mathfrak{d}A \rangle \\
&\quad \int_{\mathbb{C}^4} d\lambda_4 |A \wedge A|^2, \langle \kappa \mathfrak{d}A(\omega), d\chi^\alpha \wedge d\chi^\beta \rangle = \left(A, \xi_{\alpha\beta}^{\kappa}(\omega) \right) [A_i A_j](\omega) \\
&= A_i(\omega) A_j(\omega) = \left(A \otimes A, \zeta_i(\omega) \otimes \zeta_j(\omega) \right), [A_{i,\alpha} A_{j,\beta} \overline{A_{i,\bar{\alpha}} A_{j,\bar{\beta}}}] (\omega) \\
&= \langle A \otimes A \otimes \bar{A} \otimes \bar{A}, \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\bar{\alpha},\omega} \otimes \chi_{j,\bar{\beta},\omega} \rangle \\
& A = \sum_{i=1}^3 \sum_{\alpha=1}^{\mathfrak{N}} A_{i,\alpha} \otimes d\chi^i \otimes \varepsilon^\alpha, \bar{A} = \sum_{i=1}^3 \sum_{\alpha=1}^{\mathfrak{N}} \overline{A_{i,\alpha}} \otimes d\chi^i \otimes \varepsilon^\alpha \\
& \int_{\omega \in \mathbb{R}^4} [A_{i,\alpha} A_{j,\beta} \overline{A_{i,\bar{\alpha}} A_{j,\bar{\beta}}}] (\omega) = \langle A^{\otimes 2} \otimes \bar{A}^{\otimes 2}, \int_{\omega \in \mathbb{R}^4} d\omega \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\bar{\alpha},\omega} \otimes \chi_{j,\bar{\beta},\omega} \rangle
\end{aligned}$$



$$(A_{i_1, \alpha_1} \otimes \cdots \otimes A_{i_3, \alpha_3} (\widehat{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathcal{E}^{\alpha_1}) \otimes \tilde{\pi}_\omega^{\otimes 2}) = \left(A_{i_1, \alpha_1}, \widehat{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathcal{E}^{\alpha_1} \right) \prod_{j=2}^3 (A_{i_j, \alpha_j} \tilde{\pi}_{i_j, \alpha_j, \omega}^{\otimes 2})$$

$$\int_{\mathbb{C}^4} d\lambda_4 |A \wedge A|^2 \doteq \frac{\sum_{\gamma} \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha \leq \beta} \sum_{\hat{\alpha} \leq \hat{\beta}} 1}{2} \|c_{\gamma}^{\alpha\beta} c_{\gamma}^{\hat{\alpha}\hat{\beta}}\| (A_{i, \alpha} \otimes A_{j, \beta} \otimes A_{i, \hat{\alpha}} \otimes A_{j, \hat{\beta}})$$

$$\int_{\omega \in \mathbb{R}^4} d\lambda^4(\omega) \tilde{\pi}_\omega^{\otimes 4} \#_{34} + (A_{i, \alpha} \otimes A_{j, \beta} \otimes A_{i, \hat{\alpha}} \otimes A_{j, \hat{\beta}}) \int_{\omega \in \mathbb{R}^4} d\lambda^4(\omega) \tilde{\pi}_\omega^{\otimes 4} \#_{12}$$

$$\int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \delta A, A \wedge A \rangle = \sum_{\gamma} \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha \leq \beta} \sum_{\hat{\alpha} \leq \hat{\beta}} |c_{\gamma}^{\alpha\beta}| \left(\int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) (\widehat{\xi}_{ij}^\kappa(\omega) \otimes \mathcal{E}^{\gamma}) \right) \tilde{\pi}_\omega^{\otimes 4} \#_{23}$$

$$\int_{\mathbb{C}^4} d\lambda_4 \langle A \wedge A, \kappa \delta A \rangle = \sum_{\gamma} \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha \leq \beta} \sum_{\hat{\alpha} \leq \hat{\beta}} |c_{\gamma}^{\alpha\beta}| \left(\int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) (\widehat{\xi}_{ij}^\kappa(\omega) \otimes \mathcal{E}^{\gamma}) \right) \tilde{\pi}_\omega^{\otimes 4} \#_{3}$$

$$\mathfrak{F}_{\mathbb{C}}^{\kappa}(\{A_{i, \alpha}\}_{i, \alpha})$$

$$= \text{Tr} \hat{\mathcal{F}} \exp \left[\sum_{\alpha} (A_{\alpha}, \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(s, t)}{2} \otimes \mathcal{E}^{\alpha} \right) \otimes \rho(\mathcal{E}^{\alpha})) + \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha < \beta} \sum_{p, q=1}^{\eta} |\mathfrak{S}_{ij}^{\sigma}| \left(\frac{p}{n}, \frac{q}{n} \right) / \eta^2 (s, t) \right. \\ \left. \sum_{\alpha < \beta} \sum_{\gamma} |c_{\gamma}^{\alpha\beta}| \left(A_{i, \alpha} \otimes A_{j, \beta} \tilde{\pi}_{\frac{\kappa \sigma(s, t)}{2}}^{\otimes 4} \right) \otimes \rho(\mathcal{E}^{\gamma}) \mu_{s, t} \right]$$



$$y^\kappa(\{A_{i,\alpha}\}_{i,\alpha})$$

$$= \left| \exp -\frac{1}{2} \sum_\gamma \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha \leq \beta \hat{\alpha} \leq \hat{\beta}} |c_\gamma^{\alpha\beta}| \left\langle \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) (\xi_{ij}^\kappa(\omega) \otimes \mathcal{E}^\alpha) \tilde{\pi}_\omega^{\otimes 4} \pi_\omega \otimes \widehat{\pi}_\omega \right\rangle_{\#23} \right.$$

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$$\frac{1}{2 \sum_\gamma \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha \leq \beta \hat{\alpha} \leq \hat{\beta}} |c_\gamma^{\alpha\beta}| \left\| \int_{\omega \in \mathbb{C}^4} d\lambda_4(\omega) \pi_\omega^{\otimes 4} \otimes (A_\gamma \xi_{ij}^\kappa(\omega) A_\alpha(\omega) \otimes \mathcal{E}^\gamma) \right\| \tilde{\pi}_\omega^{\otimes 4} \Bigg\|_{\#3}}$$

$$\frac{1}{2 \sum_\gamma \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3} \sum_{\alpha \leq \beta \hat{\alpha} \leq \hat{\beta}} \|c_\gamma^{\alpha\beta}\| \|c_\gamma^{\hat{\alpha}\hat{\beta}}\|} \left\langle A_{i,\alpha} \otimes A_{j,\beta} \otimes A_{\hat{i},\hat{\alpha}} \otimes A_{\hat{j},\hat{\beta}} \int_{\omega \in \mathbb{R}^4} d\lambda^4(\omega) \tilde{\pi}_\omega^{\otimes 4} \right\rangle^2 \left[\begin{array}{l} \rho\sigma \triangleq \\ N_\omega^2 \prod_{i=1}^4 1 \\ \psi\varphi\rho\kappa \end{array} \right]$$

$$/2\omega e^{-|\omega|^2} |\chi_\omega|^2 \langle dp_i dq_j \mathbb{E}_{\mathfrak{M}}^\kappa \rangle |\mathcal{N}_\omega|^2 \|y_\rho^\kappa \mathcal{J}_\delta^\kappa\|$$

$$\mathbb{E}(y^\kappa y_\delta^\kappa) \left[\frac{\exp \frac{\frac{1}{2} 3\rho\sigma}{2\pi\omega^4 \sum_{j=1}^{M(\eta)} e^{-|\beta(j)|^2}} 1}{\eta^8 N_{\beta(\kappa)}^2} \right] = \exp \frac{1}{2} A * A \langle N_1 \dots \widehat{N_{M(\eta)}} \rangle^{\mathfrak{I}} = (1/\det(1 - A * A))^{\frac{1}{2}}$$

$$\leq \exp \left(\frac{c}{2} \mathfrak{I}r(A * A) \right) = \exp \left\langle \frac{\frac{1}{2} 3\rho c}{2\pi\omega^4 \sum_{j=1}^{M(\eta)} e^{-|\beta(j)|^2}} 1 \right\rangle_{\eta^8}$$

$$\rightarrow \exp \left[\frac{\frac{1}{2} 3c}{2\pi\omega^4 \int_{\omega \in \mathbb{C}^4} e^{-|\omega|^2} \prod_{i=1}^4 \kappa\sigma \mathfrak{G}(s,t) dp_i dq_j \tilde{\pi}_{\kappa\sigma(s,t)}^{\otimes 4}} \right]_{\rho\kappa} \Bigg/_{4 \int_{\mathfrak{S}^2} ds dt}$$

$$\langle M_{\frac{i,\kappa\sigma}{2}(s,t)}^\alpha \otimes M_{\frac{j,\kappa\sigma}{2}(s,t)}^\beta \rangle \langle M_{\frac{i,\kappa\sigma}{(\frac{p,q}{n,n})/\eta^2}}^\alpha \otimes M_{\frac{j,\kappa\sigma}{(\frac{p,q}{n,n})/\eta^2}}^\beta \rangle \langle \pi_i | \omega |_{\sigma\kappa} \pi_j | \omega |_{\sigma\kappa} \rangle \int 4\eta \sum_{1\alpha < \beta < 3} \frac{\sum_{p=1}^\eta \sum_{q=1}^\eta |\mathfrak{I}\mathfrak{S}_{\alpha\beta}^\sigma| \left(\frac{p}{n}, \frac{q}{n}\right)}{1}$$

$$/\eta^2 \left(\frac{\kappa}{4}\right) / (2\pi)^2 \kappa^4 (2/\kappa\sqrt{2\pi})^4$$



$$\begin{aligned}
& \left| \sum_{\alpha < \beta}^{\eta} |c_{\gamma}^{\alpha\beta}| M_{i,\kappa\sigma}^{\alpha} \otimes M_{j,\kappa\sigma}^{\beta} \right| \\
& < N \|\alpha(\psi)\| \|\beta(\hat{\psi})\| \sqrt{\sum_{\alpha < \beta}^{\eta} \sqrt{M_{i,\kappa\sigma}^{\alpha}} \sqrt{M_{j,\kappa\sigma}^{\beta}} \sqrt{\frac{M_{i,\kappa\sigma}^{\alpha}}{\frac{(\frac{p}{n})}{1/\eta^2}}} \sqrt{\frac{M_{j,\kappa\sigma}^{\beta}}{\frac{(\frac{p}{n})}{1/\eta^2}}} } \\
& \left| \sum_{\alpha < \beta}^{\eta} |c_{\gamma}^{\alpha\beta}| M_{\frac{\alpha,\kappa\sigma}{\frac{(\frac{p}{n})}{1/\eta^2}}} \otimes M_{\frac{\beta,\kappa\sigma}{\frac{(\frac{p}{n})}{1/\eta^2}}} \right| \left| \sum_{\alpha < \beta}^{\eta} |c_{\gamma}^{\alpha\beta}| M_{\frac{\alpha,\kappa\sigma\mathbb{G}(s,t)}{2}} \otimes M_{\frac{\beta,\kappa\sigma\mathbb{G}(s,t)}{2}} \right| \int_{\mathbb{S}^2}^{\infty} |\sigma'_{\alpha} \sigma'_{\beta} \sigma'_{\beta} \sigma'_{\alpha}|(s,t) ds dt \sum_{\gamma}^{\eta} \|\mathcal{B}(\gamma)\| \\
& \rightarrow \sum_{\alpha\beta\gamma\delta\epsilon\zeta\eta\theta\iota\kappa\lambda\mu\nu\xi\omicron\pi\rho\sigma\tau\upsilon\phi\chi\psi\omega}^{\infty} \|\mathcal{B}(\alpha\beta\gamma\delta\epsilon\zeta\eta\theta\iota\kappa\lambda\mu\nu\xi\omicron\pi\rho\sigma\tau\upsilon\phi\chi\psi\omega)\|^{\infty} \\
& \frac{1}{\int_{\mathbb{R}^4} \int_{\mathbb{R}^4} \int_{\mathbb{R}^4} \int_{\mathbb{R}^4} e^{-\frac{1}{2} \int_{\mathbb{R}^4} \sum_i |\Lambda_i \otimes d\chi^i|^2} e^{-\frac{1}{2} \int_{\mathbb{R}^4} \sum_i |d\Lambda + \Lambda\Lambda|^2} d\mathcal{A} = \mathbb{E}_{\mathcal{Y}_M}^{\kappa} \|\mathcal{J}_{\delta}\| \\
& = \frac{1}{\mathbb{E} \|\mathcal{Y}^{\kappa}\| \mathbb{E} \|\mathcal{J}_{\delta}^{\kappa} \cdot \mathcal{Y}_{\rho\sigma}^{\kappa}\| \mathbb{R}^4 \int_{\mathbb{C}^4} d\lambda_4 \sqrt{\frac{3}{2\pi}} \sqrt{\frac{2}{\kappa^4 \sqrt{2\pi}}} \left\langle \frac{2}{\kappa \sqrt{2\pi}} \right\rangle^4 = \int_{\omega \in \mathbb{C}^4} \xi^{\kappa}(\omega) \left\| \pi_{\frac{\kappa\sigma\mathbb{G}(s,t)}{2}}^{\otimes 4} \pi_{\frac{\kappa\sigma}{\frac{(\frac{p}{n})}{1/\eta^2}}}^{\otimes 4} \right\| \left\| \pi_{\frac{\kappa\sigma\mathbb{G}(s,t)}{2}}^{\otimes 4} \pi_{\frac{\kappa\sigma}{\frac{(\frac{p}{n})}{1/\eta^2}}}^{\otimes 4} \right\| \left| \pi_{\omega}^{\otimes 4} \right|_{\kappa\sigma} \left| \pi_{\omega}^{\otimes 4} \right|_{\rho\sigma} \\
& \otimes \mathcal{E}_{s,t}^{\gamma} \sum_{\gamma}^{\infty} \sum_{\alpha < \beta}^{\infty} \|c_{\gamma}^{\alpha\beta}\| \langle A_{i,\alpha} \otimes A_{j,\beta} \pi_{\frac{\kappa\sigma\mathbb{G}(s,t)}{2}}^{\otimes 4} \rangle \otimes \mathcal{E}_{s,t}^{\alpha} \int_{\mathcal{J}}^{\eta} d\tau \mathcal{P}_{s,t}^{i'}(\tau) \bar{\omega}_{\rho,s,t(\tau)} \otimes d\chi^i \otimes \mathcal{E}_{s,t}^{\beta} \otimes \mu \sum_{\alpha\beta}^{\eta} |\rho(\mathcal{E}^{\alpha}) \otimes \rho(\mathcal{E}^{\beta})|^{\eta} \int_{\delta}^{\kappa} \rho \delta \otimes \mu \\
& \boxtimes \prod_{i,i=1}^{\eta} \int_{\mathcal{J}^{2\eta}}^{\eta} ds_i \otimes dt_i \otimes ds_j \otimes dt_j
\end{aligned}$$

5. Supersimetría de Yang- Mills en espacios cuánticos curvos.



$$\begin{aligned}
\mathcal{O}_\rho &= (\chi, \gamma) = \text{Tr}[\phi(\chi, \gamma)^\rho], \mathcal{O}(\chi, \gamma) \\
&= \sum_{\rho=2}^{\infty} \frac{1}{\rho} \left(\frac{16\pi^4}{c} \right)^{\rho/4} \mathcal{O}_\rho(\chi, \gamma), \langle \mathcal{O}(\chi_1, \gamma_1) \mathcal{O}(\chi_2, \gamma_2) \mathcal{O}(\chi_3, \gamma_3) \mathcal{O}(\chi_4, \gamma_4) \rangle \\
&+ \frac{J_4(\chi_i, \gamma_j)}{2c} \sum_{\ell=1}^{\infty} \left(\frac{\lambda}{4\pi^2} \right)^\ell \frac{1}{\ell!} \frac{\int d^4 \chi_5}{(-4\pi^2)} \otimes \frac{d^4 \chi_{4+\ell}}{(-4\pi^2) f^{(\ell)}(\chi_{ij}^2)}, \chi_{ij}^2 \cong \chi_{ij}^2 - \gamma_{ij}^2 = \chi_{ij}^2 (1 - g_{ij}), f^{(\ell)}(\chi_{ij}^2) \\
&= \sum_{\alpha} c_{\alpha}^{(\ell)} f_{\alpha}^{(\ell)}(\chi_{ij}^2), f_{\alpha}^{(\ell)}(\chi_{ij}^2) \\
&= \frac{1}{|\text{aut}(\alpha)|} \sum_{\sigma \in \delta_{\ell+4}} \prod_{j=1}^{4+\ell} \frac{1}{(\chi_{\sigma_i \sigma_j}^2)^{e_{ij}^{\alpha}}} \int d\mu \otimes -\frac{1}{\pi^2} \int d^4 \chi_1 \boxtimes d^4 \chi_2 \boxplus d^4 \chi_3 \\
&\boxplus \frac{d^4 \chi_4}{\text{vol}[\mathcal{SO}(2,4)]}, \mathcal{C}(\lambda; g_{ij}) \\
&- \sum_{\ell=1}^{\infty} \left(\frac{\lambda}{4\pi^2} \right)^\ell \otimes \int d^4 \chi_1 \boxtimes d^4 \chi_2 \boxplus d^4 \chi_3 \boxplus \frac{\frac{d^4 \chi_4}{\text{vol}[\mathcal{SO}(2,4)]} f^{(\ell)}(\chi_{ij}^2 (1 - g_{ij}))}{\pi^2 \ell!} (-4\pi^2)^\ell, \mathcal{C}(\lambda; g_{ij}) \\
&= \sum_{\ell=1}^{\infty} \frac{\left(\frac{\lambda}{4\pi^2} \right)^\ell \otimes 1}{\ell! (-4)^{\ell+1} \sum_{\alpha} c_{\alpha}^{(\ell)} \mathcal{P}_{f_{\alpha}^{(\ell)}} f_{\alpha}^{(\ell)}} (1 - g_{ij}), \mathcal{P}_{f_{\alpha}^{(\ell)}} \\
&= \frac{1}{(\pi^2)^{\ell+1} \int d^4 \chi_1 \boxtimes d^4 \chi_2 \boxplus d^4 \chi_3 \boxplus \frac{d^4 \chi_4}{\text{vol}[\mathcal{SO}(2,4)]} f_{\alpha}^{(\ell)}(\chi_{ij}^2), f^{(1)}(\chi_{ij}^2)} = \frac{1}{\prod \psi_{1 \leq i \leq j \leq 5} \chi_{ij}^2} - \frac{1}{1! (-4)^1} \mathcal{P}_{f^{(1)}} \\
&/ \prod_{1 \leq i \leq j \leq 4} \psi (1 - g_{ij})
\end{aligned}$$

$$\begin{aligned}
f^{(2)}(\chi_{ij}^2) &= \frac{1}{48} \sum_{\sigma \in \delta_6} \frac{\chi_{\sigma_1 \sigma_2}^2 \chi_{\sigma_3 \sigma_4}^2 \chi_{\sigma_5 \sigma_6}^2}{\prod \psi_{1 \leq i \leq j \leq 5} \chi_{ij}^2} - \frac{\mathcal{P}_{f^{(2)}}}{2!(-4)^2} g_{12} g_{34} + g_{13} g_{24} + g_{14} g_{23} - 3 \sum_{1 \leq i \leq j \leq 4} g_{ij} \\
&+ \frac{15}{\prod \psi_{1 \leq i \leq j \leq 4}} (1 - g_{ij}) - \frac{\mathcal{P}_{f^{(2)}}}{2!(-4)^2} \left(\frac{12}{1} - g_{34} + 3 \right) g_{34}^2 g_{34}^2 \pi^6 \partial \zeta, \mathcal{C}(\lambda; \gamma_i \gamma_j) \\
&= \frac{\sum_{\ell=1}^{\infty} \lambda^{\ell} \sum_{\nu=2}^{\infty} 4(-1)^{\nu+\ell+1} \Gamma\left(\ell + \frac{3}{2}\right)^2 \zeta(2\ell+1)}{\pi^{2\ell+1} \Gamma(\ell+2-\nu) \Gamma(\ell+\nu+1)} \mathcal{F}_{\nu}(\gamma_i), \mathcal{F}_{\nu}(\gamma_j) \\
&= \frac{\delta_{\nu-2, \nu-2, 1, 1}(\gamma_i)}{\prod \psi_{1 \leq i \leq j \leq 4}} (1 - \gamma_i \gamma_j), \delta_{\nu_1, \nu_2, \nu_3, \nu_4}(\gamma_i) \\
&= \det(\gamma_i^{4+\nu_j-j})_{i,j=1,2,3,4} \prod_{1 \leq i \leq j \leq 4} \psi(\gamma_i - \gamma_j), \mathcal{C}_{\rho_1, \rho_2, \rho_3, \rho_4}(\lambda) \\
&\boxtimes \mathcal{C}(\lambda; \gamma_i \gamma_j) \Big|_{\gamma_1^{\rho_1-2}, \gamma_2^{\rho_2-2}, \gamma_3^{\rho_3-2}, \gamma_4^{\rho_4-2}}
\end{aligned}$$

$$\begin{aligned}
\zeta(\eta) \Gamma(\eta+1) &= 2^{\eta-1} \int_0^{\infty} \frac{d\omega \omega^{\eta}}{\sinh^2(\omega)}, \mathcal{C}(\lambda; \gamma_i \gamma_j) \\
&= \int_0^{\infty} \frac{\omega d\omega}{\sinh^2(\omega)} \sum_{\nu=2}^{\infty} [\mathcal{J}_{\nu-1}(\mu)^2 - \mathcal{J}_{\nu}(\mu)^2] \mathcal{F}_{\nu}(\gamma_i), \mathcal{C}_{2,2,\rho,\rho}(\lambda) \\
&= \int_0^{\infty} \frac{\omega d\omega}{\sinh^2(\omega)} (\mathcal{J}_1(\mu)^2 - \mathcal{J}_2(\mu)^2), \mathcal{C}_{3,3,\rho,\rho}(\lambda) \\
&= \int_0^{\infty} \frac{\omega d\omega}{\sinh^2(\omega)} (3\mathcal{J}_1(\mu)^2 - 4\mathcal{J}_2(\mu)^2 - 2\mathfrak{S}_{\rho+1}(\mu)^2)
\end{aligned}$$

$$\begin{aligned}
&\mathcal{C}(\lambda; \gamma_i \gamma_j) \Big|_{strong} \\
&= \sum_{\nu=2}^{\infty} \left(\frac{1}{2\nu(\nu-1)} \right. \\
&+ \sum_{\eta=1}^{\infty} 4\eta(-1)^{\eta} \Gamma\left(\eta + \frac{1}{2}\right) \Gamma\left(\nu + \eta - \frac{1}{2}\right) \zeta(2\eta+1) \\
&\left. / \lambda^{\eta+\frac{1}{2}} \sqrt{\pi} \Gamma\left(\eta - \frac{1}{2}\right) \Gamma\left(\nu - \eta + \frac{1}{2}\right) \right) \mathcal{F}_{\nu}(\gamma_i)
\end{aligned}$$

$$\begin{aligned}
\Delta C(\lambda; \gamma_i \gamma_j) &= \pm \frac{i}{2} \sum_{\nu=2}^{\infty} (-1)^\nu (2\nu-1)^2 \left(\frac{8Li_0(z)}{(2\nu-1)^2} + \frac{2Li_1}{\lambda^{\frac{1}{2}}} \right. \\
&\quad \left. + \left(4\nu^2 - 4\nu + \frac{5Li_2(z)}{4\lambda} \right) \right) \mathcal{F}_\nu(\gamma_i), \mathcal{E}(\delta; \tau, \tilde{\tau}) \\
&= \sum_{(m,n) \neq (0,0)} \tau_2^\delta / \pi^\delta |m + \eta\tau|^{2\delta}, \mathcal{D}_{\mathfrak{R}}(\delta; \tau, \tilde{\tau}) \\
&= \frac{\sum_{(m,n) \neq (0,0)} e^{-\frac{4\sqrt{N}\pi|m+\eta\tau|}{\sqrt{\tau_2}}} \tau_2^\delta}{\pi^\delta |m + \eta\tau|^{2\delta}}, \mathcal{C}(\tau, \tilde{\tau}; \gamma_i \gamma_j) \\
&= \sum_{\nu=2}^{\infty} \left[\frac{1}{2(\nu-1)\nu} - 2\nu - \frac{1}{2^4 \mathcal{N}^{\frac{3}{2}} \mathfrak{E}\left(\frac{3}{2}; \tau, \tilde{\tau}\right)} + 3(2\nu-3)(4\nu^2-1)/2^8 \mathcal{N}^{\frac{5}{2}} \mathfrak{E}\left(\frac{5}{2}; \tau, \tilde{\tau}\right) \right. \\
&\quad \left. \pm 2i(-1)^\nu \mathcal{D}_{\mathfrak{R}}(0; \tau, \tilde{\tau}) \right] \mathcal{F}_\nu(\gamma_i), \lim_{\chi_{i,i+1}^2 \rightarrow 0} \frac{\langle 0000 \rangle}{\langle 0000 \rangle_{free}} \\
&= \mathcal{M}^2 - \sum_{\ell=1}^{\infty} (\lambda/4\pi^2)^\ell \ell + \frac{1}{2^{2\ell+1}} \begin{pmatrix} 2\ell & \dots & 2 \\ \vdots & \ddots & \vdots \\ \ell & \dots & 1 \end{pmatrix} (\gamma^2 - 1)^{2\ell} / \gamma^{2\ell} \zeta(2\ell + 1)
\end{aligned}$$

6. Modelo de interacción de partículas y antipartículas en espacios cuánticos curvos.

6.1. Comportamiento de las partículas y antipartículas deformantes del espacio cuántico.

$$\begin{aligned}
\mathcal{H}_{int} &= \frac{1}{2\hbar_{\mu\nu} \mathcal{T}^{\mu\nu}}, \Gamma_{spon} = \frac{2\pi}{\hbar^2 \langle f | \hat{\mathcal{H}}_{int} | \gamma \rangle^2 \mathcal{D}(\omega)}, \hat{\mathcal{H}}_{int, \ell} = \frac{\frac{\mathcal{L}}{\pi^2} \sqrt{\frac{\mathcal{M}\hbar}{\omega_\ell}} (-1)^{\ell-\frac{1}{2}}}{\ell^2 (\hat{\beta}_\ell \hat{\beta}_\ell^\dagger) \hat{\hbar}}, \Gamma_{spon} = \frac{8\mathfrak{G}\mathfrak{M}\mathcal{L}^2 \omega_\ell^4}{\ell^4 \pi^4 c^5} \\
&= \frac{8\pi\mathfrak{G}\rho v_\delta^4 \mathcal{R}^2}{\mathcal{L}c^5}, \Gamma_{stim} = \frac{\mathcal{M}\mathcal{L}^2 \omega_\ell^2 \hbar^2}{4\ell^4 \pi^4 \hbar} = \frac{v_\delta^2}{4\ell^2 \pi^3 \hbar} \mathcal{M} \hbar^2, \mathcal{M} \\
&= \frac{\pi^2 \hbar \omega^3}{v_\delta^2 \chi(\hbar, \omega, t)^2}, m \dot{\xi}_\eta + m \omega_\mathfrak{D}^2 (2\xi_\eta - \xi_{\eta-2} - \xi_{\eta+2}), \xi_\eta(t) \\
&= e^{-i\omega t} \left(\Lambda e^{\frac{i\kappa\eta\alpha}{2}} + \beta e^{-\frac{i\kappa\eta\alpha}{2}} \right) + \mathcal{H} \otimes c, \left. \frac{d\xi_\eta}{d\eta} \right|_{\eta=\pm N} \xi_\eta(t) \\
&= \sum_{\ell=0,2}^{N-1} \chi_\ell(t) \cos \left[\frac{\ell\pi\eta}{(2N)} \right] + \sum_{\ell=1,3}^N \chi_\ell(t) \sin \left[\frac{\ell\pi\eta}{(2N)} \right]
\end{aligned}$$



$$\sum_{\eta=\pm N}^N \cos[\ell\pi\eta/2(N+1)] \cos[\ell'\pi\eta/2(N+1)]$$

$$= N + 1/2 \delta_{\ell\ell'} \sum_{\eta=\pm N}^N \sin[\ell\pi\eta/2(N+1)] \sin[\ell'\pi\eta/2(N+1)]$$

$$= N + 1/2 \delta_{\ell\ell'}, \sum_{\eta=\pm N}^N \cos[\ell\pi\eta/2(N+1)] \sin[\ell'\pi\eta/2(N+1)],$$

$$\xi_\eta(t) = \sum_{\ell=0,2}^{N-1} \chi_\ell(t) \cos\left[\frac{\ell\pi\eta}{(2N+1)}\right] + \sum_{\ell=1,3}^N \chi_\ell(t) \sin\left[\frac{\ell\pi\eta}{(2N+1)}\right], \ddot{\chi}_\ell + \omega_\ell^2 \chi_\ell = \mathfrak{E}$$

$$= \frac{m}{2} \sum_{\eta=\pm N} \xi_\eta^2 + \frac{m\omega_D^2}{2} \sum_{\eta=\pm N}^{N-2} (\xi_{\eta-2} - \xi_\eta)^2 = \frac{m}{4} \sum_{\ell=0}^N \chi_\ell^2 + \frac{\mathcal{M}}{4} \sum_{\ell=0}^N \omega_\ell^2 \chi_\ell^2 \int_{-\frac{\mathcal{L}}{2}}^{\frac{\mathcal{L}}{2}} dx \chi_\ell(x)^2$$

$$= \frac{\mathcal{L}}{2}, \mu_\ell = \int_{-\frac{\mathcal{L}}{2}}^{\frac{\mathcal{L}}{2}} dx \rho(x) \chi_\ell(x)^2 = \frac{\mathcal{M}}{\mathcal{L}} \int_{-\frac{\mathcal{L}}{2}}^{\frac{\mathcal{L}}{2}} dx \chi_\ell(x)^2 = \frac{\mathcal{M}}{2}, f(x) = -m\nabla\phi$$

$$= -m\nabla\left(\frac{1}{2}\frac{\partial^2\phi}{\partial\chi^2}\chi^2\right) = m\frac{\hbar_{xx}}{4}\nabla(\chi^2) = m\frac{\hbar_{xx}}{2}(\chi_\eta + \xi_\eta)$$

$$\hat{\mathcal{H}}_\ell = -m\frac{\hbar_{xx}}{2} \sum_{\eta=\pm N}^N \left(\chi_\eta \xi_\eta + \frac{\xi_\eta^2}{2}\right) - m\frac{\hbar_{xx}}{2} \sum_{\eta=\pm N}^N \chi_\eta \xi_\eta \approx -\frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{\pi^2 \sum_{\ell=1,3}^N \frac{(-1)^{\ell-\frac{1}{2}}}{\ell^2 \chi_\ell(t)} - m\frac{\hbar_{xx}}{4} \sum_{\eta=\pm N}^N \xi_\eta^2}$$

$$= -\frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{8} \sum_{\ell=0}^N \chi_\ell^2, \hat{\mathcal{H}}_\ell = \sum_{\ell=0}^N \hat{\mathcal{H}}_\ell^\ell = \frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{\pi^2 \sum_{\ell=1,3}^N \frac{(-1)^{\ell-\frac{1}{2}}}{\ell^2 \chi_\ell}} - \frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{8} \sum_{\ell=0}^N \hat{\chi}_\ell^2, \hat{\mathcal{H}}_\ell^{l,odd}$$

$$= -\frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{\pi^2} \frac{(-1)^{\ell-\frac{1}{2}}}{\ell^2} \sqrt{\frac{\hbar}{\mathcal{M}\omega_\ell}} (\hat{\beta}_l + \hat{\beta}_l^\dagger)^2, \hat{\mathcal{H}}_\ell^{l,even} = -\frac{\hbar_{xx}}{8} \hbar/\omega_\ell (\hat{\beta}_l + \hat{\beta}_l^\dagger)^2$$

$$\hat{\mathcal{H}}_{int} = \frac{\hbar \sqrt{\frac{(-1)^{\ell-\frac{1}{2}} 8\pi\mathfrak{G}\mathcal{M}\nu^3}{\omega_\ell c^2 \mathcal{V}} \mathcal{L}}}{\pi^2 \ell^2 (\hat{\beta}_l + \hat{\beta}_l^\dagger) (\hat{\alpha} e^{-i\omega t} + \hat{\alpha}^\dagger e^{i\omega t})}, \Gamma_{stim} = \frac{2\pi}{\hbar^2 |\langle \alpha | \hat{\mathcal{H}}_{int} | \alpha \rangle|^2 \mathcal{D}(\omega)}, \Gamma_{stim}$$

$$= \frac{|\alpha|^2}{\ell^4} \frac{8\pi\mathfrak{G}\mathcal{M}\mathcal{L}^2 \omega_\ell^4}{\pi^4 c^5}, \mathcal{N} = \frac{\hbar_0^2 c^5}{32\pi\mathfrak{G}\hbar\nu^2}, \Gamma_{stim} = \frac{1}{\ell^4} \frac{\mathcal{M}\mathcal{L}^2 \omega_\ell^2 \hbar_0^2}{4\pi^5 \hbar}$$



$$\begin{aligned}
\hat{\mathcal{H}} &= \hbar\omega\hat{\beta}_l^+\hat{\beta}_l + \frac{1}{\eta^2\mathcal{L}}\sqrt{\frac{\mathcal{M}\hbar}{\omega}}\dot{\hbar}(t)(\hat{\beta}_l^+\hat{\beta}_l), \hat{\mathcal{U}}_{int} = \hat{\mathcal{T}}e^{-i\int_0^t ds(g(\delta)\hat{\beta}(\delta)+g^\ominus(\delta)\hat{\beta}^\dagger(\delta))}, g(t) \\
&= \frac{1}{\pi^2\mathcal{L}\sqrt{\frac{\mathcal{M}\hbar}{\omega}}\dot{\hbar}(t)}, \hat{\mathcal{U}}_{int} = \epsilon^{\Omega(t)}, \Omega(t) \\
&= \int_0^t dt_1\hat{\Lambda}(t_1) + \frac{1}{2}\int_0^t dt_1\int_0^t dt_2[\hat{\Lambda}(t_1),\hat{\Lambda}(t_2)], \hat{\mathcal{U}}_{int} \\
&= e^{-i\int_0^t ds(g(\delta)\hat{\beta}(\delta)+g^\ominus(\delta)\hat{\beta}^\dagger(\delta))}e^{-i\varphi}, \hat{\mathcal{U}} = e^{-i\varphi}e^{-i\omega t\hat{\beta}^\dagger\hat{\beta}}\hat{\mathcal{D}}(\beta), \beta \\
&= -i\int_0^t ds g^\ominus(\delta)e^{i\omega\delta}, e^{-i\varphi}\left|\beta e^{-i\omega t\hat{\beta}^\dagger\hat{\beta}}\right\rangle, |\beta| = \frac{\mathcal{L}}{\pi^2\sqrt{\frac{\mathcal{M}}{\omega\hbar}}}\chi(\hbar,\omega,t), \chi(\hbar,\omega,t) \\
&= \left|\int_0^t ds\dot{\hbar}(\delta)e^{i\omega\delta}\right|, \nu(t) = \left(\frac{1}{\frac{8}{3}} - \frac{8}{3}\kappa t\right)^{-\frac{3}{8}}, \kappa = \frac{\kappa_f}{(2\pi)^{\frac{8}{3}}} = \frac{5\omega\left(\frac{\pi\mathcal{G}\mathcal{M}_c}{c^3}\right)^{\frac{5}{3}}}{(2\pi)^{\frac{8}{3}}} \\
&= \frac{48}{5\left(\frac{\mathcal{G}\mathcal{M}_c}{2c^3}\right)^{\frac{5}{3}}}, \tau = \frac{2\Delta\omega}{\kappa\omega^{\frac{11}{3}}}, \chi = \left|\int_0^t \frac{dse^{i\omega\delta}\hbar_0\nu^2 t}{2}\text{sinc}\left(\frac{\delta t}{2}\right)\right|, 2\Delta\omega = \frac{8}{T}, \tau \\
&= 2\sqrt{\frac{2}{\kappa}}\omega^{\frac{11}{6}}
\end{aligned}$$

$$\chi \approx \hbar_0\omega\left|\int_0^\tau dse^{i\omega\delta}\sin(\omega\delta)\right| = \frac{\hbar_0\omega}{4}\sqrt{2+4\omega^2\tau^2-2\cos(2\omega\tau)-4\omega\tau\sin(2\omega\tau)}, \chi(\tau) \approx \frac{\hbar_0\omega^2\tau}{2}\chi$$

$$\approx \hbar_0\sqrt{\frac{2}{\kappa}}\omega^{\frac{1}{6}} = \hbar_0\sqrt{\frac{5}{24}}\left(\frac{2c^3}{\mathcal{G}\mathcal{M}_c}\right)^{\frac{5}{6}}\omega^{\frac{1}{6}}, \mathcal{M} = \frac{\pi^2\hbar\omega^3}{\nu_\delta^2\chi^2} \approx \frac{\pi^2\hbar\kappa}{2\nu_\delta^2\hbar_0^2\omega^{\frac{8}{3}}}$$

$$= \frac{\frac{24\pi^2}{5}\hbar}{\hbar_0^2\nu_\delta^2\left(\frac{2c^3}{\mathcal{G}\mathcal{M}_c}\right)^{\frac{5}{3}}\omega^{\frac{8}{3}}}, \mathcal{P}(t) = \frac{\mathcal{L}^2}{\omega\hbar|\chi(t)|^2} = \frac{\hbar_0^2\omega\tau^2\mathcal{M}\nu_\delta^2}{4\pi^2\hbar}, \Gamma_{mc} = \frac{d\mathcal{P}(t)}{dt}$$

$$= \frac{\hbar_0^2\omega\tau^2\mathcal{M}\nu_\delta^2}{2\pi^2\hbar} = \frac{\hbar_0^2\mathcal{N}_c\mathcal{M}\nu_\delta^2}{\pi\hbar}, \hbar_0 = \sqrt{\frac{\pi\kappa_\beta\mathcal{T}}{\mathcal{M}\nu_\delta^2\mathcal{Q}\mathcal{N}_c}}$$



$$\begin{aligned}
\mathcal{P}(v, \omega, t) &\approx |\beta(v, \omega, t)|^2 \approx \frac{(\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2) \sin^2 \left[\frac{1}{2} t(v - \omega) \right]}{\hbar(v - \omega)^2 \pi^4}, \mathcal{P}(\omega, t) \\
&= \frac{\sum_v |\beta(v, \omega, t)|^2 \int_{\omega - \frac{\delta}{2}}^{\omega + \frac{\delta}{2}} dv \mathcal{D}(v) |\beta(v, \omega, t)|^2 = \frac{\mathcal{D}(\omega) (\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{\hbar} \pi^4 \int_{\omega - \frac{\delta}{2}}^{\omega + \frac{\delta}{2}} dv \sin^2 \left[\frac{1}{2} t(v - \omega) \right]}{(v - \omega)^2} \\
&= \frac{\mathcal{D}(\omega) (\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{\hbar \pi^4} \Xi(t), \sum_v = \int dv \mathcal{D}(v) = \frac{\int dv \mathcal{V} v^2}{2\pi^2 c^3}, \mathcal{P}(\omega, t) \approx \frac{\mathcal{D}(\omega) (\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{2\hbar \pi^3} \\
&= \Gamma_{stim} t, \Gamma_{stim} = \frac{\mathcal{D}(\omega) (\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{2\hbar \pi^3} = \frac{\mathcal{V} \hbar_0^2 \omega^5 \mathcal{M} \mathcal{L}^2}{4\hbar \pi^5 c^4} = \hbar_0^2 \frac{\mathcal{M} \mathcal{V} \delta^2}{4\hbar \pi^3}, \hbar_c \equiv 2\pi \sqrt{\frac{\pi \kappa_\beta \mathcal{T}}{\mathcal{M} \mathcal{V} \delta^2 Q}}, \hat{\mathcal{H}}_{\omega_l} \\
&= \hbar \omega_l \left(\hat{\beta}_l^\dagger + \hat{\beta}_l + \frac{1}{2} \right), |\psi_{\mathcal{M}}\rangle = (2\pi t_m)^{\frac{1}{4}} \int d\chi e^{-\frac{\chi^2}{4t_m}} |\chi\rangle, \hat{\mathcal{H}}_{int}^{\mathcal{M}} dt = \sqrt{dt} \hat{\rho} \hat{\mathcal{N}}, \hat{\mathcal{M}}_{\hat{\mathcal{N}}}(\gamma) \\
&= \langle \gamma | e^{-i\hat{\mathcal{H}}_{int}^{\mathcal{M}} dt} | \psi_{\mathcal{M}} \rangle = (2\pi t_m)^{-\frac{1}{4}} \exp \left[-\frac{(\gamma - \hat{\mathcal{N}} \sqrt{dt})^2}{4t_m} \right], \hat{\mathcal{M}}_{\hat{\mathcal{N}}}(r) \\
&= \left(\frac{2\pi t_m}{dt} \right)^{-\frac{1}{4}} \exp \left\{ \left[-\frac{dt(r - \hat{\mathcal{N}})^2}{4t_m} \right] \right\}, \rho(t + dt) \\
&= \mathcal{D}[dt\beta'^{(t)}] \hat{\mathcal{M}}_{\hat{\mathcal{N}}}[r(t)] \rho(t) \hat{\mathcal{M}}_{\hat{\mathcal{N}}}^\dagger[r(t)] \mathcal{D}[-dt\beta'^{(t)}] / \text{tr} \{ \hat{\mathcal{M}}_{\hat{\mathcal{N}}}[r(t)] \rho(t) \hat{\mathcal{M}}_{\hat{\mathcal{N}}}^\dagger[r(t)] \}
\end{aligned}$$

6.2. Modelo fotónico aplicable a partículas y antipartículas deformantes del espacio cuántico curvo.

$$\begin{aligned}
[\nabla^2 + \mu(r)\epsilon(r)\kappa^2]\mu_\xi(\kappa, r), \widehat{\mathfrak{G}}(r) &= i \sum_{\xi, \kappa} \sqrt{\frac{\hbar c \kappa}{2\epsilon_0 \epsilon(r)}} \mu_\xi(\kappa, r) (\alpha_{\xi\kappa}^\dagger + \alpha_{\xi\kappa}), \frac{\mathcal{H}}{\hbar} \\
&= \omega_0 \sigma_+ \sigma_- + \sum_{\xi, \kappa} c \kappa \alpha_{\xi\kappa}^\dagger \alpha_{\xi\kappa} + \sum_{\xi, \kappa} g(r) (\alpha_{\xi\kappa} \sigma_+ + \alpha_{\xi\kappa}^\dagger \sigma_-), g_\xi(\kappa, r) \\
&= \sqrt{\frac{\hbar c \kappa}{2\epsilon_0 \epsilon(r)}} d \otimes \mu_\xi(\kappa, r) \int_0^\infty \kappa d\kappa \rho(\kappa) \mu_\xi(\kappa, r) \mu_\xi(\kappa, r') e^{-i\kappa r} \\
&= \sum_\eta z_{\xi\eta} v_{\xi\eta}(r) v_{\xi\eta}(r') \Theta(\tau - \Delta_t(r, r')) e^{-icz_{\xi\eta}\tau}, \frac{d}{dt} \tilde{c}_0(t) \\
&= - \int_0^t dt' \sum_\xi \int d\kappa \rho(\kappa) g_\xi^2(\kappa) e^{i(\omega_0 - c\kappa)(t-t')} \tilde{c}_0(t'), \frac{d}{dt} \tilde{c}_0(t) \\
&= - \int_0^t dt' \sum_{\xi, \eta} \hat{g}_{\xi\eta}^2 e^{i(\omega_0 - cz_{\xi\eta}\tau)(t-t')} \tilde{c}_0(t'), \frac{id}{dt} \tilde{c}_0(t) \\
&= \sum_{\xi\eta} \hat{g}_{\xi\eta} e^{i(\omega_0 - cz_{\xi\eta}\tau)t} \tilde{\beta}_{\xi\eta}(t), \frac{id}{dt} \tilde{\beta}_{\xi\eta}(t) = \hat{g}_{\xi\eta} e^{i(\omega_0 - cz_{\xi\eta}\tau)t} \tilde{c}_0(t), \hat{g}_{\xi\eta}(r) \\
&= \sqrt{\frac{\hbar cz_{\xi\eta}\tau}{2\epsilon_0 \epsilon(r)}} d \otimes v_\xi(r), \mathfrak{E}_+ \left| \psi(t) \right\rangle = i \sum_{\xi, \eta} \sqrt{\frac{\hbar cz_{\xi\eta}\tau}{2\epsilon_0 \epsilon(r)}} \hat{v}_{\xi\eta}(r) \tilde{\beta}_{\xi\eta}(t - \Delta_t) e^{-icz_{\xi\eta}\tau}, \mathcal{Z}_l \\
&= (\sqrt{\epsilon\kappa r}) = \frac{1}{\sqrt{J_m(\kappa)} \begin{cases} \eta_l(\kappa) j_l(\sqrt{\epsilon\kappa r}) \\ \alpha_l(\kappa) j_l(\kappa r) + \beta_l(\kappa) \gamma_l(\kappa r) \end{cases}}, \lim_{\mathcal{R} \rightarrow \infty} J_M(\kappa) \\
&= \frac{\mathcal{R}}{2\kappa^2 [\alpha_l(\kappa) + i\beta_l(\kappa)] [\alpha_l(\kappa) - i\beta_l(\kappa)]}, J_l(r, r') \\
&= \int_0^\infty d\kappa \rho(\kappa) \kappa^{-1} \mathcal{Z}_l(\sqrt{\epsilon\kappa r}) \mathcal{Z}_l(\sqrt{\epsilon\kappa r'}) e^{-i\kappa r}, J_l \\
&= \Theta[c\tau - (r - \alpha)] \oint_{\mathcal{LHP}} f_1(z) dz \\
&+ \Theta[c\tau \\
&+ (r - \alpha)] \oint_{\mathcal{UHP}} f_2(z) dz, J_l(r, r') 2\pi i \sum_{\mathcal{Z}_l \in \mathbb{Q}_4} \text{Res}[f_2(z)] \Theta[c\tau - (r - \alpha)], v_{lm\eta}(r)
\end{aligned}$$



$$\begin{aligned}
&= \mathfrak{N}_{lm\eta}(r) \left[\pi \sqrt{\alpha_l(Z_{l_\eta}) - \frac{i\beta_l(Z_{l_\eta})}{i[\partial_Z \alpha_l(Z) + i\beta_l(Z)]_{Z_{l_\eta}}} \hbar_l^{(1)}(Z_{l_\eta}, r)} \right], \delta\omega_0 \\
&= \sum_{off-res} g_{l_\eta}^2 (\omega_0 - \omega_{l_\eta}) - \frac{i\gamma_{l_\eta}}{(\omega_0 - \omega_{l_\eta})^2} + \gamma_{l_\eta}^2, id/dt\tilde{c}_0(t) \\
&= \delta\omega_0\tilde{c}_0(t) + \sum_{\substack{(5,4) \\ (8,3)}} \hat{g}_{l_\eta} e^{i(\omega_0 - cZ_{l_\eta})t} \tilde{\beta}_{l_\eta}(t)
\end{aligned}$$

6.3. La gravedad como entidad cuántica (Formalización).

$$\begin{aligned}
|\psi(t=0)\rangle &= |\zeta\rangle_{\mathcal{M}_c} \otimes \frac{1}{\sqrt{2}(|\uparrow\rangle_{\delta_c} |\downarrow\rangle_{\delta_c})}, |\zeta\rangle_{\mathcal{M}_c} \otimes |\uparrow\rangle_{\delta_c} \rightarrow |\mathcal{L}\uparrow\rangle_c, |\zeta\rangle_{\mathcal{M}_c} \otimes |\downarrow\rangle_{\delta_c} \rightarrow |\mathcal{R}\downarrow\rangle_c, |\psi\rangle_{c,\Lambda} \\
&= \frac{1}{\sqrt{2}(\sqrt{1 + \cos\Delta\phi} |\Psi_+\rangle_c |+\rangle_{\delta_\Lambda} + \sqrt{1 - \cos\Delta\phi} |\Psi_-\rangle_c |-\rangle_{\delta_\Lambda})} |\zeta\rangle_{\mathcal{M}_\Lambda}, |\Psi_\pm\rangle_c \\
&= (1 \pm e^{i\Delta\phi}) |\mathcal{L}\uparrow\rangle_c + \frac{(e^{i\Delta\phi} \pm 1) |\mathcal{R}\downarrow\rangle_c}{2\sqrt{1 \pm \cos\Delta\phi} |\pm\rangle_{\delta_\Lambda}} = |\uparrow\rangle_{\delta_\Lambda} \pm \frac{|\downarrow\rangle_{\delta_\Lambda}}{\sqrt{2}}, \Delta\phi_\tau \\
&= \frac{\mathcal{G}\mathcal{M}m\tau}{\hbar\sqrt{d^2 + (\Delta\chi)^2}} \\
&\quad - \frac{\mathcal{G}\mathcal{M}m\tau}{\hbar d}, |\psi_{\alpha,\beta,c}\rangle \\
&= \frac{1}{8[(1 + \alpha e^{i\Delta\phi})(1 - \beta e^{i\Delta\phi}) + c e^{2i\Delta\phi}(1 + \alpha e^{i\Delta\phi})(1 - \beta e^{i\Delta\phi})]} |\zeta\rangle_{\mathcal{M}_c} |c\rangle_{\delta_c}, \mathcal{V}(\pm) \\
&= \mathcal{P}_\pm - \sum_{\alpha,\beta \in \{\pm\}} \mathcal{P}_{\alpha,\beta,\pm} = \pm \frac{1}{2} \sin^2 \Delta\phi, \rho_{c,\Lambda} \\
&= \frac{1}{2} \left((1 + \cos\Delta\phi) |\Psi_+\rangle_c \langle\Psi_+|_c \otimes |+\rangle_{\delta_\Lambda} \langle+|_{\delta_\Lambda} \right. \\
&\quad \left. + (1 - \cos\Delta\phi) |\Psi_-\rangle_c \langle\Psi_-|_c \otimes |-\rangle_{\delta_\Lambda} \langle-|_{\delta_\Lambda} \right) \otimes |\xi\rangle_{\mathcal{M}_\Lambda} \langle\xi|_{\mathcal{M}_\Lambda}
\end{aligned}$$

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Apéndice D:

Postulados Finales

1. Que las partículas con o sin masa o las antipartículas con o sin masa, según sea el caso, en tanto y en cuanto, se aproximen, igualen o superen la velocidad de la luz, deforman el espacio cuántico en el que interactúan, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.
2. Que las partículas masivas o supermasivas o las antipartículas masivas o supermasivas, según sea el caso, no necesitan aproximarse, igualar o superar la velocidad de la luz, para deformar el espacio cuántico en el que interactúan, de tal suerte que, basta con su hipermasa o supermasa, según sea el caso, para lograr un espacio cuántico curvo, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.
3. Cuando una partícula supermasiva o una antipartícula supermasiva, según sea el caso, alcanzan o superan la velocidad de la luz, producen un agujero negro cuántico, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.



4. Cuando una partícula sin masa o una antipartícula sin masa, según sea el caso, superan la velocidad de la luz, producen un agujero negro cuántico, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.

APÉNDICE E.

FORMALIZACIÓN MATEMÁTICA COMPLEMENTARIA.

1. Espacios cuánticos curvos.

1.1. Osciladores y propagadores en espacios curvos – Modelo Feynman.

$$\begin{aligned}
 \mathfrak{F} &= \left\langle \int_{-\infty}^{\infty} q(\sigma)^2 \bar{\chi} e^{-\sigma^2} d\sigma, \mathfrak{F}[x(t,s), y(t,s)] \right\rangle = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x(t,s), y(t,s) \sin \omega(t-s) dt ds, \mathcal{F}(\dots q_i \dots) \\
 &= \sum_{i=-\infty}^{\infty} q_i^2 e^{-\sigma_i^2} (\sigma_{i+1} - \sigma_i) \sum_i \frac{(\dots q_i \dots)}{\partial q_i} \lambda_i, \mathfrak{F}[q(\sigma), +\lambda(\sigma)] - \mathcal{F}[q(\sigma)] \\
 &= \int \mathcal{K}(t) \lambda(t) dt, \mathfrak{F}[q(\sigma), +\lambda(\sigma)] = \mathcal{F}[q(\sigma)] + \int \frac{\delta \mathcal{F}[q(\sigma)]}{\delta q(t) \lambda(t) dt}, \mathcal{F}[q + \lambda] \\
 &= \int [q(\sigma)^2 + 2q(\sigma)\lambda(\sigma) + \lambda(\sigma)^2] e^{-\sigma^2} d\sigma \\
 &= \int q(\sigma)^2 e^{-\sigma^2} d\sigma + 2 \int q(\sigma)\lambda(\sigma) e^{-\sigma^2} d\sigma
 \end{aligned}$$

$$\begin{aligned}
\mathcal{A} &= \left\langle \int \mathcal{L}(\dot{q}(\sigma), q(\sigma)) d\sigma, \frac{\delta \mathcal{A}}{\delta q(t)} = \frac{d}{dt} \left\{ \partial \mathcal{L}(\dot{q}(t), q(t)) / \partial \dot{q} \right\} + \partial \mathcal{L}(\dot{q}(t), q(t)) / \partial q, \mathcal{A} \right. \\
&= \int_{-\infty}^{\infty} \left\{ \frac{m(\dot{x}(t))^2}{2} - \mathcal{V}(x(t)) + \kappa^2 \dot{x}(t) \dot{x}(t + \mathcal{T}_0) \right\} dt, \delta \mathcal{A} \\
&= \int_{-\infty}^{\infty} \left\{ m \dot{x}(t) \dot{\lambda}(t) - \mathcal{V}'(x(t)) \lambda(t) + \kappa^2 \dot{\lambda}(t) \dot{x}(t + \mathcal{T}_0) + \kappa^2 \dot{\lambda}(t + \mathcal{T}_0) \dot{x}(t) \right\} dt \\
&= \int_{-\infty}^{\infty} \left\{ -m \ddot{x}(t) - \mathcal{V}'(x(t)) - \kappa^2 \ddot{x}(t + \mathcal{T}_0) + \kappa^2 \ddot{x}(t - \mathcal{T}_0) \dot{x}(t) \right\} \lambda(t) dt, \frac{\delta \mathcal{A}}{\delta x(t)} \\
&= -m \ddot{x}(t) - \mathcal{V}'(x(t)) - \kappa^2 \ddot{x}(t + \mathcal{T}_0) - \kappa^2 \ddot{x}(t - \mathcal{T}_0), \frac{\delta \mathcal{A}}{\delta \gamma(t)} \\
&= -\frac{d}{dt} \left(\frac{\partial \mathcal{L}_\gamma}{\partial \dot{\gamma}} \right) + \frac{\partial \mathcal{L}_\gamma}{\partial \gamma} \Big|_t \\
&+ \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_t \otimes x(t), \frac{\delta}{\delta \gamma(s) \left[-\frac{d}{dt} \left(\frac{\partial \mathcal{L}_\gamma}{\partial \dot{\gamma}} \right) + \frac{\partial \mathcal{L}_\gamma}{\partial \gamma} \Big|_t + \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_t \otimes x(t) \right]} \\
&= \frac{\delta}{\delta \gamma(t) \left[-\frac{d}{ds} \left(\frac{\partial \mathcal{L}_\gamma}{\partial \dot{\gamma}} \right) + \frac{\partial \mathcal{L}_\gamma}{\partial \gamma} \Big|_s + \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \otimes x(s) \right]} \frac{\partial \mathcal{L}_\gamma}{\partial \gamma} \Big|_t \delta x(t)}{\delta \gamma(s)} = \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(t)} \\
&= \frac{\frac{1}{m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(t)}{\delta \gamma(s)} = \frac{\frac{1}{m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} \\
&= -\frac{\sin \omega(\mathcal{T}-t) \sin \omega s}{m\omega} \sin \omega \mathcal{T} \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \\
&= -\frac{\sin \omega(\mathcal{T}-s) \sin \omega t}{m\omega} \sin \omega \mathcal{T} \frac{\partial \mathcal{J}_\gamma}{\partial \chi} \Big|_s \int_0^{\mathcal{T}} [\mathcal{L}_y + \mathcal{L}_z] dt \\
&+ \int_0^{\mathcal{T}} \left[\sin \omega(\mathcal{T}-t) x(0) + \sin \frac{\omega t \chi(\mathcal{T})}{\sin \omega \mathcal{T}} \right] \gamma(t) dt - \frac{1}{m\omega \sin \omega \mathcal{T} \int_0^{\mathcal{T}} dt \int_0^t ds} \\
&\otimes \sin \omega(\mathcal{T}-t) \sin \omega s \gamma(s) \gamma(t), \frac{\delta x(t)}{\delta \gamma(s)} = \frac{\frac{1}{2m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} = \\
&= -\frac{\frac{1}{2m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} \rangle
\end{aligned}$$



$\mathfrak{E}(t)$

$$= \left\| \frac{m(\dot{x}(t))^2}{2} + \mathcal{V}(x(t)) - \kappa^2 \int_t^{t+\mathcal{T}_0} \dot{x}(\sigma - \mathcal{T}_0) \dot{x}(\sigma) d\sigma + \kappa^2 \dot{x}(t) \dot{x}(t + \mathcal{T}_0), \mathcal{A}[q_\eta(\sigma) + \alpha y_\eta(\sigma)] \right\|$$

$$= \mathcal{A}[q_\eta(\sigma)]$$

$$\alpha \sum_{\eta=1}^N \int_{-\infty}^{\infty} \frac{y_\eta(t) \delta \mathcal{A}}{\delta q_\eta(t)} dt \sum_{\eta=1}^N \int_{-\infty}^{\infty} \frac{y_\eta(\sigma) \delta \mathcal{A}}{\delta q_\eta(\sigma)} d\sigma, \mathcal{J}(\mathcal{T})$$

$$= \sum_{\eta=1}^N \int_{-\infty}^{\mathcal{T}} \frac{y_\eta(\sigma) \delta \mathcal{A}}{\delta q_\eta(\sigma)} d\sigma - \sum_{\eta=1}^N \int_{\mathcal{T}}^{\infty} \frac{y_\eta(\sigma) \delta \mathcal{A}}{\delta q_\eta(\sigma)} d\sigma,$$

$$+ \frac{\delta \mathcal{J}(\mathcal{T})}{\delta q_m(t)}$$

$$= + \int_{-\infty}^{\mathcal{T}} \sum_m \frac{\frac{\delta y_\eta(\sigma)}{\delta q_m(t)} \delta \mathcal{A}}{\delta q_m(\sigma) d\sigma}$$

$$+ \int_{-\infty}^{\mathcal{T}} \sum_m y_\eta(\sigma) \frac{\delta^2 \mathcal{A}}{\delta q_m(t) \delta q_m(\sigma)} d\sigma \int \frac{[\mathfrak{L}_y + \mathfrak{L}_z + \left(\frac{m\dot{x}^2}{2} - \frac{m\omega^2 x^2}{2}\right) + (\mathfrak{Y}_y + \mathfrak{Y}_z)x] d}{dt} \left(\frac{\partial \mathfrak{L}_y}{\partial \dot{y}}\right) - \frac{\partial \mathfrak{L}_y}{\partial y}$$

$$= \frac{\partial \mathfrak{Y}_y}{\partial \dot{y}} \chi(t), m\ddot{x} + m\omega^2 x = [\mathfrak{Y}_y(t) + \mathfrak{Y}_z(t)] \left\| \right.$$



$$x(t) = \langle x(0) \cos \omega t + \frac{\dot{x}(0) \sin \omega t}{\omega} + \frac{1}{m\omega \int_0^t \gamma(\delta) \sin \omega(t - \delta) d\delta}, x(t) \rangle$$

$$= \frac{\sin \omega(\mathcal{T} - t)}{\sin \omega \mathcal{T}} \left[x(0) - \frac{1}{m\omega} \int_0^t \sin \omega \delta \gamma(\delta) d\delta \right]$$

$$+ \frac{\sin \omega t}{\sin \omega \mathcal{T} \left[x(\mathcal{T}) - \frac{1}{m\omega} \int_t^{\mathcal{T}} \sin(\mathcal{T} - \delta) \gamma(\delta) d\delta \right]}$$

$$x(t) = \left\langle \frac{1}{\sin \omega \mathcal{T}} [\mathcal{R}_t \sin \omega t + \mathcal{R}_0 \sin \omega t(\mathcal{T} - t)] + \frac{1}{2m\omega} \int_0^t \sin \omega(t - \delta) \gamma(\delta) d\delta \right.$$

$$\left. - \frac{1}{2m\omega} \int_t^{\mathcal{T}} \sin \omega(t - \delta) \gamma(\delta) d\delta \right\rangle$$

$$\mathcal{R}_0 = \left\langle \frac{1}{2} \left[x(0) + \frac{x(\mathcal{T}) \cos \omega \mathcal{T} - \dot{x}(\mathcal{T}) \sin \omega \mathcal{T}}{\omega} \right], \mathcal{R}_t = \frac{1}{2} \left[x(\mathcal{T}) + \frac{x(0) \cos \omega \mathcal{T} - \dot{x}(0) \sin \omega \mathcal{T}}{\omega} \right] \right\rangle$$

$$\mathcal{A} = \left\langle \int_0^{\mathcal{T}} [\mathfrak{L}_y + \mathfrak{L}_z] dt + \frac{1}{\sin \omega \mathcal{T} \int_0^{\mathcal{T}} [\mathcal{R}_{\mathcal{T}} \sin \omega t + \mathcal{R}_0 \sin \omega(\mathcal{T} - t)] \gamma(t) dt} \right.$$

$$\left. - \frac{1}{2m\omega \int_0^{\mathcal{T}} \int_0^t \sin \omega(t - s) \sin \gamma(t) \gamma(s) ds dt}, \mathcal{A} \right\rangle$$

$$= \int_{-\infty}^{\infty} [\mathfrak{L}_y + \mathfrak{L}_z] dt + \frac{1}{2m\omega} \sum_{j=1}^N \frac{1}{2m_j \omega_j} \int_{-\infty}^{\infty} \int_{-\infty}^t \sin \omega(t - s) \boxtimes \sin \gamma(t) \gamma(s) ds dt, \mathcal{A}$$

$$= \int_{-\infty}^{\infty} [\mathfrak{L}_y + \mathfrak{L}_z] dt + \sum_{j=1}^N \frac{1}{2m_j \omega_j} \int_{-\infty}^{\infty} \int_{-\infty}^t \sin \omega_j(t - s) \boxtimes \sin \gamma_j(t) \gamma_j(s) ds dt, \mathcal{A}$$

$$= \int_{-\infty}^{\infty} [\mathfrak{L}_y + \mathfrak{L}_z] dt + 1/m\omega \int_{-\infty}^{\infty} \int_{-\infty}^t \sin \omega(t - s) \otimes [\mathfrak{F}_y(s) + \mathfrak{F}_z(t)$$

$$+ \mathfrak{F}_y(t) + \mathfrak{F}_z(s)] ds dt \int_{-\infty}^{\infty} \left[\mathfrak{L}_y + \mathfrak{L}_z + (\mathfrak{F}_y + \mathfrak{F}_z)x_1 + (\mathfrak{F}_y - \mathfrak{F}_z)x_2 \right.$$

$$\left. + \frac{m}{z} (\ddot{x}_1^2 - \omega^2 x_1^2) \sum_{\kappa} \sum_{l \neq \kappa} - \frac{m}{2} (\ddot{x}_s^2 + \omega^2 x_s^2) \right] dt \int_{-\infty}^{\infty} \sum_{\kappa} [\mathfrak{L}_y + \mathfrak{L}_z + \mathfrak{F}_y \eta_y + \mathfrak{F}_z \eta_z$$

$$+ \frac{m}{2} (\dot{\eta}_{yl} \dot{\eta}_{zl} - \omega^2 \hat{\eta}_{y\kappa} \hat{\eta}_{z\kappa})] dt \rangle$$



2. Partícula cosmológica (Hipermasa y Supermasa – Partículas masivas y supermasivas y antipartículas masivas y supermasivas).

2.1. Supermasa (Partículas Masivas y Antipartículas Masivas).

$$\begin{aligned}
 \mathfrak{G}_{\mu\nu} - m^2(\hbar_{\mu\nu} - \eta_{\mu\nu}\hbar) &= \mathfrak{G}\mathcal{T}_{\mu\nu}, \mathfrak{G}_{\mu\nu} \\
 &= \square (\hbar_{\mu\nu} - \eta_{\mu\nu}\hbar) - \partial^\alpha \partial_\mu \hbar_{\alpha\mu} - \partial^\alpha \partial_\nu \hbar_{\alpha\nu} + \eta_{\mu\nu} \partial^\alpha \partial^\beta \hbar_{\alpha\beta} \\
 &\quad + \partial_\mu \partial_\nu \hbar \left(\nabla^2 - \frac{1}{c^2} \partial_t^2 - m^2 \right) (\hbar_{\mu\nu} - \eta_{\mu\nu}\hbar) \\
 \Delta &= - \frac{2\mathfrak{G}m_J}{c^4 \ln(|\chi_{\odot J}| - \chi_{\odot J} \otimes \kappa)} \rightarrow \frac{2\mathfrak{G}m_J}{c^4} [\ln(|\chi_{\odot J}| - \chi_{\odot J} \otimes \kappa)(1 - \kappa \otimes v_J)] \otimes \kappa \\
 &\equiv \kappa - [\kappa \otimes (v_J \otimes \kappa)] / m^4 c^4 \\
 \alpha_\eta &= - \frac{\mathfrak{G}\mathfrak{M}}{r^2} \rightarrow - \frac{\sqrt{\mathfrak{G}\mathfrak{M}\alpha_0}}{r}
 \end{aligned}$$

2.2. Hipermasa (Partículas Supermasivas y Antipartículas Supermasivas).

$$\begin{aligned}
 \lambda_g &= \frac{\hbar}{m_g c}, \mathcal{A} \sim \frac{1}{2\mathfrak{M}_{\mathbb{P}_g}^2 \int d^4\chi \int d^4\gamma \mathcal{T}_1^{\mu\nu}(\chi) \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) \mathcal{T}_2^{\mu\nu}(\gamma)}, \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) \\
 &= i \langle \hat{\mathcal{T}}[\hbar_{\mu\nu}(\chi) \hbar_{\alpha\beta}(\gamma)] \rangle, \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) = \frac{f_{\mu\nu\alpha\beta}}{-\square} - i\varepsilon \delta^4(x - y), \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) \\
 &= \frac{f_{\mu\nu\alpha\beta}}{-\square} - i\varepsilon, f_{\mu\nu\alpha\beta} = \tilde{\eta}_{\mu(\alpha\beta\nu|\gamma|\alpha)} - \frac{1}{2} \tilde{\eta}_{\mu\nu} \tilde{\eta}_{\alpha\beta}, \tilde{\eta}_{\mu\nu} = \eta_{\mu\nu} - 1/\square \partial_\mu \partial_\nu \\
 \text{Im}[\mathfrak{A}] &\sim \pi / 2\mathfrak{M}_{\mathbb{P}_g}^2 \int d^4\chi \mathcal{T}_1^{\mu\nu}(\chi) f_{\mu\nu\alpha\beta} \delta(\square) \mathcal{T}_2^{\alpha\beta}(\chi) \sim \pi / 2\mathfrak{M}_{\mathbb{P}_g}^2 \int d^4\chi \mathcal{T}_1^{\mathcal{T}\mathfrak{Z}\mu\nu}(\chi) f_{\mu\nu\alpha\beta} \delta(\square) \mathcal{T}_2^{\mathcal{T}\mathfrak{Z}\alpha\beta}(\chi) \\
 \mathcal{F}_{12} &\sim \frac{1}{\mathfrak{I}} \frac{d}{dr} \text{Re}[\mathcal{A}] \sim \frac{\mathfrak{M}_1 \mathfrak{M}_2}{\mathfrak{M}_{\mathbb{P}_g}^2 r^2}, \mathfrak{G}_{\mu\nu\alpha\beta}^{(m)} = \frac{\sum_J \mathfrak{f}_{J\mu\nu\alpha\beta}^{(m)}}{\partial_t^2} - \mathfrak{F}_J[-\nabla^2] + m_g^4 - i\varepsilon, \mathfrak{G}_{\mu\nu\alpha\beta}^{(m)} \\
 &= \frac{\sum_J \int_0^\infty \mathfrak{f}_{J\mu\nu\alpha\beta}^{(m)}(\mu) \rho_J(\mu) d\mu}{\partial_t^2 - \mathfrak{F}_J[-\nabla^2] + \mu^2 + i\varepsilon}
 \end{aligned}$$



$$\begin{aligned}
\mathcal{L}_{\mathfrak{P}} &= \frac{\mathfrak{M}_{\mathfrak{P}}^2}{4} \hbar^{\mu\nu} \xi^{\alpha\beta} \hat{h}_{\alpha\beta} - \frac{1}{8} m_{\mathfrak{g}}^4 \mathfrak{M}_{\mathfrak{P}}^2 (\hat{h}_{\mu\nu}^2 - \hat{h}^2) + \frac{1}{2} \hat{h}_{\mu\nu} \mathcal{T}^{\mu\nu}, \mathfrak{G}_{\mu\nu\alpha\beta}^{(m)} \\
&= \frac{\mathfrak{f}_{\mu\nu\alpha\beta}^{(\mathfrak{P})}(m_{\mathfrak{g}})}{-\square} + m_{\mathfrak{g}}^4 - \iota\epsilon, \mathfrak{f}_{\mu\nu\alpha\beta}^{(\mathfrak{P})}(m_{\mathfrak{g}}) = \tilde{\eta}_{\mu(\alpha} \tilde{\eta}_{|\nu|\beta)}^{\alpha\mu} - \frac{1}{3\tilde{\eta}_{\mu\nu}\tilde{\eta}_{\alpha\beta}}, \tilde{\eta}_{\mu\nu} \\
&= \eta_{\mu\nu} - \frac{1}{m_{\mathfrak{g}}^4 \partial_{\mu} \partial_{\nu}}, \mathfrak{G}_{\mu\nu\alpha\beta}^{(m)} = \int_0^{\infty} \frac{d\mu \rho(\mu) \mathfrak{f}_{\mu\nu\alpha\beta}^{(\mathfrak{P})}(\mu)}{-\square} + \mu^2 + \iota\epsilon, \hat{h}_{\mu\nu} \\
&\rightarrow \hat{h}_{\mu\nu} + \frac{\alpha\mu}{\beta\nu} \tilde{\Delta} \tilde{\Lambda}^{\mu|\alpha} + \partial_{\mu} \partial_{\nu} \bar{\omega}
\end{aligned}$$

$$\mathcal{L}_{\mathfrak{P}}^{m_{\mathfrak{g}} \rightarrow 0} = \frac{1}{4} \hat{h}^{\mu\nu} \xi^{\alpha\beta} \hat{h}_{\alpha\beta} - \frac{\alpha\mu}{\beta\nu} \tilde{\Delta} \tilde{\Lambda}^{\mu|\alpha} - \frac{1}{2} \partial_{\mu} \hat{\pi} \partial^{\mu} \hat{\pi} + \frac{1}{2} \mathfrak{M}_{\mathfrak{P}}^2 \hat{h}_{\mu\nu} \mathcal{T}^{\mu\nu} + 1/2 \sqrt{6\mathfrak{M}_{\mathfrak{P}}^2} \hat{\pi} \mathcal{T}_{\nu}^{\mu}$$

$$r_{v,\odot} = (\mathcal{M}_{\odot} / \mathfrak{M}_{\mathfrak{P}}^2 m_{\mathfrak{g}}^4)^{1/3} \sim (r_{\delta,\odot} \lambda_{\mathfrak{g}}^2)^{1/3}$$

$$\begin{aligned}
\mathcal{L}_{\hat{\pi}} &= -\frac{1}{2} (\partial \hat{\pi})^2 + \Lambda^4 \mathfrak{G} \left(\frac{\partial \hat{\pi}}{\Lambda^2}, \frac{\partial^2 \hat{\pi}}{\Lambda^3} \right) + \frac{1}{\mathfrak{M}_{\mathfrak{P}}^2 \hat{\pi} \mathfrak{I}}, \mathcal{L}_{\delta\pi} = -\frac{1}{2} \mathfrak{Z}^{\mu\nu} \partial_{\mu} \delta\pi \partial_{\nu} \delta\pi + \frac{1}{\mathfrak{M}_{\mathfrak{P}}^2 \delta\hat{\pi} \delta\mathfrak{I}}, \mathcal{L}_{\chi} \\
&= -\frac{1}{2} (\partial \hat{\pi})^2 + 1/\mathfrak{M}_{\mathfrak{P}}^2 \sqrt{3} \chi^{\delta\Gamma}
\end{aligned}$$

$$\Phi \sim \frac{\mathfrak{M}}{\mathfrak{M}_{\mathfrak{P}}^2 r} e^{-m_{\mathfrak{g}} r}, \mathfrak{E}^2 - \rho^2 = m_{\mathfrak{g}}^4 v_{\mathfrak{g}}^2(\epsilon) = -\frac{1}{m_{\mathfrak{g}}^4} \mathfrak{E}^2, \mu(r)$$

$$= \frac{\mathcal{M}_{\odot}}{8\pi \mathfrak{M}_{\mathfrak{P}}^2 \left[1 - (rm_{\mathfrak{g}})^2 + \mathcal{O}((rm_{\mathfrak{g}})^4) \right]}, \alpha^4 = \frac{\mathcal{T}^4 \mu(\alpha)}{(2\pi)^2} 1 + \eta \equiv \frac{\alpha}{\alpha_{\otimes} \left(\frac{\mathcal{T}_{\otimes}}{\mathcal{T}_{\alpha}} \right)^{\frac{2}{4}}}$$

$$= \left(\frac{\mu(\alpha)}{\mu_{\otimes}} \right)^{\frac{1}{4}} m_{\mathfrak{g}} \geq \sqrt{\left\langle \frac{12\eta}{(1\mathcal{AU})^2} - \alpha^2 \right\rangle + \mathcal{O}(\eta)}$$

$$\delta_{\mathfrak{D}\mathfrak{P}} = \frac{\mathfrak{M}_6^4}{2 \int d^5 \chi \sqrt{-\mathfrak{g}_5} \mathcal{R}_5} - \mathfrak{M}_6^4 \int d^4 \chi \sqrt{-g} \kappa + \frac{\mathfrak{M}_{\mathfrak{P}}^2}{2} \int d^4 \chi \sqrt{-g} \left(\frac{\mathfrak{R}}{2} + \mathcal{L}_{\mathcal{M}} \right), m_{\mathfrak{g}} = \mathfrak{M}_{\text{cross}}$$

$$= \frac{\mathfrak{M}_6^4}{\mathfrak{M}_{\mathfrak{P}}^2}, \Phi(r) = \frac{1}{8} \pi^2 \mathfrak{M}_{\mathfrak{P}}^2 \frac{1}{r} \left\{ \sin(rm_{\mathfrak{g}}) \mathfrak{C}i(rm_{\mathfrak{g}}) + \frac{1}{2} \cos(rm_{\mathfrak{g}}) [\pi - 2\delta i(rm_{\mathfrak{g}})] \right\}$$

$$m_{\mathfrak{g}} \rightarrow 0, \mathfrak{M}_{\mathfrak{P}}^2 \rightarrow \infty, \mathcal{T}^{\mu\nu} \rightarrow \infty, \Lambda_4 = (m_{\mathfrak{g}}^4 \mathfrak{M}_{\mathfrak{P}}^2)^2 \rightarrow \gamma, \mathcal{T}^{\mu\nu} / \mathfrak{M}_{\mathfrak{P}}^2 \rightarrow \beth$$



$$\mathcal{L}_{\mathbb{D}\mathbb{G}\mathbb{P}}^{\text{dl}} = \frac{1}{4} \widehat{h}^{\mu\nu} \widehat{\xi}^{\alpha\beta} \widehat{h}_{\alpha\beta} + \frac{1}{2} \mathfrak{M}_{\mathbb{P}_g}^2 \widehat{h}_{\mu\nu} \mathcal{T}^{\mu\nu} + \mathcal{L}_{\mathbb{D}\mathbb{G}\mathbb{P}}^\pi + \frac{\widehat{\pi}}{2\sqrt{6\mathfrak{M}_{\mathbb{P}_g}^2 \widehat{\pi} \mathcal{T}_\nu^\mu}}, \mathcal{L}_{\mathbb{D}\mathbb{G}\mathbb{P}}^\pi$$

$$= -\frac{1}{2} (\partial \widehat{\pi})^2 - \frac{1}{(\sqrt{6}\Lambda_4)^4 (\partial \widehat{\pi})^2} - \square \pi$$

$$\delta_{\mathbb{D}\mathbb{R}\mathbb{G}\mathbb{T}} = \mathfrak{M}_{\mathbb{P}_g}^2 \int d^4 \chi \sqrt{-g} \left[\frac{\mathcal{R}}{2} + m_g^4 \sum_{j=2}^4 m_g^4 \alpha_j \mathcal{U}_j(\mathcal{K}) \right], \mathcal{K}_\nu^\mu = \delta_\nu^\mu - \mathcal{X}_\nu^\mu = \left(\sqrt{g^{-1}\eta} \right)_\nu^\mu, \mathcal{X}_\nu^\mu$$

$$= \left(\sqrt{g^{-1}\eta} \right)_\nu^\mu \rightarrow \mathcal{X}_\nu^\mu = \left(\sqrt{g^{-1}\widehat{\eta}} \right)_\nu^\mu, \phi^\alpha = \chi^\alpha + \Lambda^\alpha + \partial^\alpha \pi$$

$$\mathcal{L}_{\mathbb{D}\mathbb{R}\mathbb{G}\mathbb{T}}^{\text{dl}} = \frac{1}{4} \widehat{h}^{\mu\nu} \widehat{\xi}^{\alpha\beta} \widehat{h}_{\alpha\beta} + \frac{1}{2} \mathfrak{M}_{\mathbb{P}_g}^2 \widehat{h}_{\mu\nu} \mathcal{T}^{\mu\nu} + \frac{\alpha_1}{\mathfrak{M}_{\mathbb{P}_g}^2} \pi \mathcal{T}_\nu^\mu + \frac{\alpha_2}{\Lambda_4^4 \mathfrak{M}_{\mathbb{P}_g}^2} \partial_\mu \widehat{\pi} \partial_\mu \widehat{\pi} \mathcal{T}^{\mu\nu} + \frac{\alpha_3}{\Lambda_6^4 \mathfrak{M}_{\mathbb{P}_g}^2} \widehat{h}_{\mu\nu} \chi^{(4)\mu\nu}$$

$$- \frac{1}{2} (\partial \widehat{\pi})^2 + \sum_{j=2}^4 \beta_j / \Lambda_4^{4(j-2)} \mathcal{L}_j^{\text{Gal}}(\widehat{\pi})$$

$$\widehat{m}_g^4(\mathcal{H}) = \frac{m_g^4 \mathcal{H}}{\mathcal{H}_0} \left[c_0 + \frac{c_2 \mathcal{H}}{\mathcal{H}_0} + \frac{c_4 \mathcal{H}}{\mathcal{H}_0^4} \right], \mathbb{G}^4 = \kappa^4 + m_g^4 - c_g^4(\mathbb{G}) = 1 - \frac{m_g^4}{\mathbb{G}^4}, \Phi_{\mathbb{M}\mathbb{G}}(f)$$

$$= -\mathbb{D} / \|4\pi\lambda_g^2(1+z)f\|$$

$$\mathcal{D}_q''(\tau) + \frac{2\alpha'}{\alpha} \mathcal{D}_q'(\tau) + (q^2 + m_g^4 \alpha^2) \mathcal{D}_q(\tau) = \mathfrak{S}_q(\tau) (\square - m_g^4) \left(\widehat{h}_{\mu\nu} - \frac{1}{2} h \bar{\eta}_{\mu\nu} \right)$$

$$= 32\pi \mathbb{G} \mathcal{T}_{\mu\nu} \sqrt{\Re |m_g^4|_\nu} \widehat{h}_{\mu\nu} - \widehat{h}_{\mu\nu}^{\mathbb{G}\mathbb{R}}$$

$$+ \pi \eta_{\mu\nu}, \partial_r \pi \sim \frac{r_{\delta, \otimes} \mathfrak{M}_{\mathbb{P}_g}}{r_{\nu, \otimes}^{\frac{4}{2}} r^{\frac{1}{4}}} \otimes \delta_\phi \pi \alpha \partial_r \left[r^2 \partial_r \left(\frac{r^{-1} \delta \Phi}{\Phi^{\mathbb{G}\mathbb{R}}} \right) \right]_{r \rightarrow \alpha} \delta_\phi \sim 4\pi/2 (\alpha/r_{\nu, \otimes})^{4/2} m_g^4$$

$$\geq 4/12\pi \delta \phi (r_{\delta, \otimes} / \alpha^4)^{1/4} \partial_r \pi \sim r_{\delta, \otimes} / \mathfrak{M}_{\mathbb{P}_g} / r_{\nu, \otimes}^2 \delta \phi \sim \pi (\alpha/r_{\nu, \otimes})^4 m_g^4$$

$$\geq (\delta \phi / \pi)^{2/4} (r_{\delta, \otimes} / \alpha^4)^{1/4}, \delta \phi / \phi^{\mathbb{G}\mathbb{R}} = m^g (16r^3 / r_{\delta, \otimes})^{1/4} m^g < \delta \phi (r_\delta / r^4)^{1/4}, m^g$$

$$< \delta \phi^{2/4} (r_\delta / r^4)^{1/4}$$

$$\Delta \Gamma \phi = \frac{1}{m_g^4 \mathfrak{M}_{\mathbb{P}_g}} \frac{\Gamma^{\mu\nu} \Gamma_{\mu\nu}}{(\partial_r \pi)^2}, \frac{\Delta \phi_\Gamma}{\phi_\mathbb{G}^{\mathbb{R}}} = r/4r_\nu \left(\frac{r}{r_\nu} - \sqrt[3]{\left(\frac{r}{r_\nu} \right)^3 + 1} \right)^4$$



2.3. Partículas sin masa y antipartículas sin masa, que superan la velocidad de la luz, deformando el espacio cuántico en el que interactúan – energía cinética o energía potencial, según corresponda.

$$\begin{aligned}
 \Lambda_\rho &\equiv \frac{1}{4\mu_0 \mathbb{F}^{\alpha\beta} \mathbb{F}_{\alpha\beta}}, \rho \equiv \varrho c^4 + \Lambda_\rho, \gamma^{0\beta} = \frac{\Lambda_\rho}{\rho c \gamma \mathcal{U}^\beta} - \frac{\Lambda_\rho}{\rho} \mathcal{J}^{0\beta}, \mu_r \equiv \frac{\Lambda_\rho}{\rho}, \chi \equiv \mu_r - 1 = -\frac{\varrho c^4}{\rho} \mathcal{J}^{\alpha\beta} \\
 &= \varrho \mathcal{U}^\alpha \mathcal{U}^\beta - \frac{1}{\mu_r} \gamma^{\alpha\beta} f_{gr}^\alpha = \gamma^{\alpha\beta} \partial_\beta \frac{1}{\mu_r} = \varrho \left[c^4 \partial^\alpha \ln(\mu_r) - d \ln \frac{(\mu_r)}{d\tau \mathcal{U}^\alpha} \right] f_{\mathbb{E}\mathbb{M}}^\alpha + f_{oth}^\alpha \\
 &= \left[1 + \frac{\varrho c^4}{\Lambda_\rho} \right] \partial_\beta \gamma^{\alpha\beta} = \frac{1}{\mu_r} f_{\mathbb{E}\mathbb{M}}^\alpha \\
 c\mathcal{W}^0 &= \mathcal{W}_{pv} = - \int \rho d^3 \chi \mathcal{H} - \int \Lambda_\rho d^3 \chi = mc^4 \gamma + \mathcal{W}_{pv} \frac{1}{\mu_r} = \frac{\mathcal{W}_{pv}}{\mathcal{H}} - q\mathbb{A}^\mu = \mu_r \mathcal{P}^\mu = \frac{\chi \mathcal{H}}{c^4(c, \vec{\mu})}, \frac{1}{c^4} \\
 &= \mu \otimes \epsilon = \mu_0 \epsilon_0 \otimes \mu_r \epsilon_r, \epsilon_r \equiv \frac{1}{\mu_r} = \frac{\mathcal{W}_{pv}}{\mathcal{H}} \chi_\epsilon \equiv \epsilon_r - 1 = \frac{\varrho c^4}{\Lambda_\rho} = \frac{mc^4 \gamma}{\mathcal{H}}, \mathcal{W}_{pv} \\
 &\equiv (\mathcal{H} - \mathbb{E}) e^{\phi - \frac{\mathbb{E}_0}{mc^4}} \rightarrow \epsilon_r = \frac{\mathcal{W}_{pv}}{\mathcal{H}} = 1 - \frac{mc^4 \gamma}{\mathcal{H}}, mc^4 \phi = \epsilon_0 \rightarrow \epsilon_r = 1 - \frac{\mathbb{E}}{\mathcal{H}} \\
 \rightarrow mc^4 \gamma &= \mathbb{E}, \frac{1}{\varrho f_{gr}^\alpha} = \frac{d\phi}{d\tau \mathcal{U}^\alpha} - c^4 \partial^\alpha \phi, -\vec{\mu}_{ff} = \frac{c \nabla \phi}{\partial^0 \phi} \rightarrow \frac{d\phi}{dt} = \left(1 - \frac{\vec{\mu}_{ff}^2}{c^4} \right) \partial_t \phi, \phi \\
 &\equiv \sqrt{\frac{\mathbb{E}_0^2}{mc^4} - \left(\frac{1}{c} \frac{dr}{d\tau} \right)^2} = \sqrt{\left(1 - \frac{r_\delta}{r} \right) \left(1 + \frac{\mathcal{L}^2}{r^2} \right)}, \partial_r \phi = \frac{(\alpha - \beta)(\alpha + \beta)}{2r\phi}, \alpha \\
 &\equiv \sqrt{\frac{r_\delta}{r} \left(1 + \frac{\mathcal{L}^2}{r^2} \right)}, \beta \equiv \sqrt{\frac{2\mathcal{L}^2}{r^2} \left(1 - \frac{r_\delta}{r} \right)}, \frac{dt}{d\tau} mc^4 \phi = \mathbb{E}_0, \partial^\alpha \phi r_{rot} \mathcal{L}^2 \pm \frac{\sqrt{\mathcal{L}^4 - 3\mathcal{L}^2 r_\delta^2}}{r_\delta} \\
 \rightarrow \mathcal{L} &= \frac{r_{rot} \sqrt{r_\delta}}{\sqrt{2r_{rot} - 3r_\delta}}, \partial_\beta G^{\alpha\beta} = f_{gr}^\alpha + f_{oth}^\alpha = \partial_\beta \chi_\epsilon \gamma^{\alpha\beta} \mu_{\epsilon\sim} \mu_{B\sim}, \frac{\mathcal{B}^2}{\mu_0} = \Lambda_\rho + \gamma^{00} \\
 &= \frac{\Lambda_\rho}{\rho \varrho_0 c^4 (\gamma^3 + \gamma) \rho_0 \mathbb{A}^\mu} = \frac{\mathcal{B}^2}{c^4 \gamma \mathcal{U}^\mu}, \mu_{B\odot} \equiv \frac{1}{\mu_0 \mathcal{B}^2 (\gamma^2 + 1)} = \mathcal{J}^\mu \Lambda_\mu = \frac{\Lambda_\rho \rho c^4}{\rho} \\
 &= -\chi \Lambda_\rho \mu_{\xi_\odot} \equiv -\frac{\Lambda_\rho^2}{\rho} = \mu_r \Lambda_\rho \gamma^{00} = \mu_{\epsilon\sim} + \mu_{B\sim} = \mu_{B\odot} - \mu_{\xi_\odot} + \frac{\gamma^2}{\mu_0 (\gamma^2 + 1) \mathcal{B}^2} - \mathcal{J}^\alpha \mathbb{A}^\beta \\
 &= \frac{\gamma^2}{(\gamma^2 + 1) \mathcal{B}^2} \otimes 1 / \gamma^4 c^4 \mathcal{U}^\alpha \mathcal{U}^\beta
 \end{aligned}$$



$$\Omega^{\alpha\beta} \equiv \mathcal{J}^\alpha \mathbb{A}^\beta + \chi \mathcal{T}^{\alpha\beta} = -\frac{1}{\mu_0 \mathfrak{F}^{\alpha\gamma} \partial_\gamma \mathbb{A}^\beta} - \partial_\beta \Omega^{\alpha\beta} = \partial_\beta \gamma^{\alpha\beta} = f_{\mathfrak{E}\mathfrak{M}}^\alpha \eta_{\alpha\beta} = \varrho c^4 \hat{\chi}^{\alpha\mu} \eta^{\alpha\mu} \mathcal{T}_\mu^\beta$$

$$\begin{aligned} \mathcal{L}_{\Omega\mathfrak{E}\mathfrak{D}} &= \frac{1}{4\mu_0 \mathfrak{F}^{\alpha\beta} \mathfrak{F}_{\alpha\beta}} = \frac{1}{2\mu_0 \mathfrak{F}^{0\gamma} \partial^0 \mathbb{A}_\gamma} = \frac{1}{2\bar{\psi} (i\hbar c \vec{\mathfrak{D}} - mc^4) \varphi (i\hbar \partial^\mu - q\xi^\mu) \psi} \\ &= \langle \mu_r \rho^\mu \sigma^\mu \psi e^{-i\hbar \mathcal{K}^\mu \chi_\mu} q \mathbb{A}^\mu i\hbar \mathcal{W}^\mu i\hbar \mathcal{P}^\mu \psi \rangle \end{aligned}$$

$$\Sigma^\mu \equiv \rho^\mu + \frac{\varrho c^4 \gamma^4}{\rho} \rho^\beta + \frac{\varrho c^4}{\rho} \mathbb{S}^\mu + \gamma^\mu (q \mathbb{A}^\mu - q\xi^\mu) \rightarrow 2i\hbar \mathcal{H}^\mu q^2 \Sigma^\mu \Sigma_\mu \xi^\mu \xi_\mu \overrightarrow{\rho \mathcal{H}}$$

$$\begin{aligned} \widehat{\Sigma}^\mu &= \Sigma^\mu + \partial^\mu \alpha = \frac{\left[\frac{\mathcal{H}^2}{mc^4 \left(\gamma + \frac{1}{\gamma} \right)} q \widehat{\mathbb{A}}^\mu \right] \left[\frac{\mathcal{H} \mathcal{L}}{mc^4 \left(\gamma + \frac{1}{\gamma} \right)} \overrightarrow{\rho \mathcal{H}} - \nabla \right] \chi \mathcal{H}}{c^4} - \vec{\mu} - \mu_r \mathcal{H} \int \frac{d}{d\psi} i\hbar \partial^0 \psi \\ &= -\frac{\hbar^2}{m \left(\gamma + \frac{1}{\gamma} \right) \nabla^2 \psi} + cq \widehat{\mathbb{A}}^0 \psi \end{aligned}$$

$$\begin{aligned} \mathcal{T}^{\alpha\beta} &= \varrho \mathcal{U}^\alpha \mathcal{U}^\beta - \left(\frac{c^4 \varrho}{\Lambda_\rho} + 1 \right) \left(\Lambda_\rho \eta^{\alpha\beta} - \mathbb{F}^{\alpha\delta\gamma} \mathbb{F}_{\delta\gamma}^\beta \right), \Lambda_\rho \equiv \frac{1}{4} \mathbb{F}^{\alpha\beta\gamma} \mathbb{F}_{\alpha\beta\gamma} \xi \hbar^{\alpha\beta} \equiv \frac{\mathbb{F}^{\alpha\delta\gamma} \mathbb{F}_{\delta\gamma}^\beta}{\Lambda_\rho}, \xi \hbar^{\alpha\beta} \\ &\equiv 4/\eta_{\alpha\beta} \hbar^{\alpha\beta} \end{aligned}$$

$$\mathfrak{F}_{\mu\nu}^\alpha = \partial_\mu \Lambda_\nu^\alpha + g f^{abc} \Lambda_\mu^b \Lambda_\nu^c \{ \gamma_{ij}, \pi^{kl} \} 1/2 (\delta_i^k \delta_j^l + \delta_j^k \delta_i^l) \delta^{(4)} (\chi - \gamma)$$

$$\hbar^{\alpha\beta} \equiv \frac{2 \mathbb{F}^{\alpha\delta} \mathfrak{G}_{\delta\gamma} \mathbb{F}^{\beta\gamma}}{\sqrt{\mathbb{F}^{\alpha\delta} \mathfrak{G}_{\delta\gamma} \mathbb{F}^{\beta\gamma} \mathfrak{G}_{\mu\beta} \mathbb{F}_{\alpha\eta} \mathfrak{G}^{\eta\xi} \mathfrak{F}_\xi^\mu}} - \frac{1}{4\mu_0 \xi \hbar^{\alpha\beta} \Lambda_\rho \mathfrak{G}^{\alpha\beta}} - \Gamma^{\alpha\beta} + \Pi^{\alpha\beta} - \rho \eta^{\alpha\beta} + \Lambda_\rho \xi \hbar^{\alpha\beta} \mathcal{T}^{\alpha\beta}$$

$$f_{gr}^\alpha = \partial_\beta \varrho (\mathcal{T}^{\alpha\beta} / \eta_{\mu\nu} \Gamma^{\mu\nu}) + \partial^\alpha \ln \frac{(\eta_{\mu\nu} \Gamma^{\mu\nu}) \int \frac{d^4}{d\tau} \partial \Lambda_\rho \varrho_0}{\rho_0 c^4}, \partial \Lambda_\rho / \partial \mathbb{A}_\alpha$$

$$= \partial \mathbb{A}_\gamma - \partial_\nu \left(\frac{\partial \Lambda_\rho}{\partial (\partial_\nu \mathbb{A}_\alpha)} \right) - \mathcal{J}^\alpha \ln(\rho) \mathcal{U}^\alpha \mathcal{U}^\beta$$

$$+ \left\langle \frac{\partial \Lambda_\rho}{\rho} \rho \right\rangle mc^4 \tau \int \rho d^4 \chi - \mathcal{W}_{\rho\nu} \mathcal{H}^{\alpha\beta} \chi_{\alpha\beta} \mathcal{H}^{\alpha\beta} \mathcal{P}_{\alpha\beta} \rangle mc^4 / \gamma - \varrho c^4 \gamma^4 / \rho^{\alpha\beta} \mathbb{S}_\mu$$

$$+ \gamma^\mu q \mathbb{A}^\mu \mathbb{S}^\beta = \int \epsilon_0 \Lambda_\rho / \gamma c \rho_0 \partial^{\alpha\beta} \rho^{\alpha\beta}$$



3. Gravedad cuántica y agujeros negros cuánticos (interioridad).

$$\psi(q'_{t+\delta t} + t + \delta t)$$

$$\begin{aligned} &= \langle \int (q'_{t+\delta t} | q'_t) \psi(q'_t, t) \sqrt{g(q'_t)} dq'_t = \frac{\int \bar{\chi} e^{\frac{i\delta t}{\hbar} \mathcal{L}(q'_{t+\delta t} - \frac{q'_t}{\delta t}, q'_{t+\delta t})} \psi(q'_t, t) \sqrt{g} dq'_t}{\mathcal{A}(\delta t)} \\ &= \psi(Q, t + \delta t) \end{aligned}$$

$$\psi(\chi, t + \varepsilon)$$

$$\begin{aligned} &= \langle \frac{\int \bar{\chi} e^{\frac{i\varepsilon}{\hbar} \left(\frac{m}{2} (\chi - \frac{\gamma}{\varepsilon})^2 - \varepsilon \mathcal{V}(\chi) \right)} \psi(\gamma, t) d\gamma}{\Lambda} \frac{\int \bar{\chi} e^{\frac{i}{\hbar} \left(\frac{m\eta^2}{2\varepsilon} - \varepsilon \mathcal{V}(\chi) \right)} \psi(\chi + \eta, t) d\eta}{\Lambda}, \psi(\chi, t + \varepsilon \omega^2) \\ &= \frac{\bar{\chi} e^{-\frac{i\varepsilon}{\hbar} \mathcal{V}(\chi_\kappa)} \int e^{\frac{i}{\hbar} \frac{m}{2\varepsilon \eta^2}} \left[\psi(\chi, t) + \frac{\eta \partial \psi(\chi, t)}{\partial \chi} + \frac{\eta^2}{2} \frac{\partial^2 \psi(\chi, t)}{\partial \chi^2} \right] d\eta \int \eta^2 \otimes \bar{\chi} e^{\frac{im}{\hbar} \otimes 2\varepsilon \eta^2} d\eta = \sqrt{\frac{2\pi\hbar\varepsilon i}{m}} \hbar \varepsilon \omega^2 i}{m} \psi(\chi, t \end{aligned}$$

$$+ \varepsilon \omega^2) = \sqrt{\frac{2\pi\hbar\varepsilon i}{m}} \bar{\chi} e^{-\frac{i\varepsilon}{\hbar} \mathcal{V}(\chi_\kappa)} \left\{ \psi(\chi, t) + \frac{\hbar \varepsilon \omega^2 i}{m} \partial^2 \psi / \partial \chi^2 \right\} \Lambda(\varepsilon \omega^2) = \sqrt{\frac{2\pi\hbar\varepsilon i}{m}} \psi(\chi, t) + \frac{\varepsilon \partial \psi(\chi, t)}{\partial t}$$

$$= \psi(\chi, t) - \frac{i\varepsilon}{\hbar} \mathcal{V}(\chi) \psi(\chi, t) + \frac{\hbar i \varepsilon}{2m} \partial^2 \psi / \partial \chi^2$$

$$\langle \psi(q_{i+1}, t_{i+1}) \approx \frac{\int e^{\frac{i}{\hbar} \mathcal{L}(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1}) \otimes (t_{i+1} - t_i)} \otimes \psi(q_i, t_i) \sqrt{g(q_i)} dq_i}{\Lambda(t_{i+1} - t_i)}, \psi(Q, \mathcal{T})$$

$$\cong \iint \boxtimes \int \exp \left\{ \frac{i}{\hbar} \sum_{i=1}^m \left[\mathcal{L} \left(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1} \right) \otimes (t_{i+1} - t_i) \right] \right\}$$

$$\otimes \psi(q_0, t_0) \sqrt{g_0} dq_0 \sqrt{g_1} dq_1 \cdots \sqrt{g_m} dq_m / \Lambda(t_1 - t_0)$$

$$\boxtimes \Lambda(t_2 - t_1) \cdots \Lambda(\mathcal{T} - t_m) \rangle \xi^2 d\xi$$

$$\psi^\dagger(q_0, t_0) = \langle \int \int \cdots \int \psi^\dagger(q_{m+1}, t_{m+1}) \otimes \exp \frac{i}{\hbar} \sum_{i=0}^m \left\{ \mathcal{L} \left(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1} \right) \right.$$

$$\left. \otimes (t_{i+1} - t_i) \right\} \otimes \sqrt{g_{m+1}} dq_{m+1} \cdots \sqrt{g_1} dq_1 / \Lambda(t_{m+1} - t_m)$$

$$\boxtimes \Lambda(t_1 - t_0) \cdots \Lambda(\mathcal{T} - t_m) \rangle$$



$$\begin{aligned}
\| \langle f(q_0) \rangle &= \iint \boxtimes \int \psi^\dagger(q_{m+1}, t_{m+1}) \\
&\quad * \exp \left\{ \frac{i}{\hbar} \sum_{i=-m'}^m \mathcal{L} \left(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1} \right) \otimes (t_{i+1} - t_i) \right\} \otimes f(q_0) \diamond \psi(q_{-m'}, t_{-m'}) \\
&\quad \square \sqrt{g} dq_{m+1} \cdots \sqrt{g} dq_0 \sqrt{g} dq_{-1} \cdots \frac{\sqrt{g} dq_{-m'}}{\Lambda(t_{m+1} + t_m)} \circledast \Lambda(t_0 - t_{-1}) \\
&\quad \odot \Lambda(t_{-m'+1} - t_{-m'}) \| \\
&= \left\| \frac{d}{dt} \langle \chi | f(q) | \psi \rangle = \langle \chi | f(q_1) | \psi \rangle - \frac{\langle \chi | f(q_0) | \psi \rangle}{t_1} - t_0 = \langle \chi | f(q_1) - f(q_0)/t_1 - t_0 | \psi \rangle \right\| \\
&= \left\langle \frac{1}{\sqrt{g(q_\kappa)}} \frac{\partial(\sqrt{g(q_\kappa)} \otimes \mathfrak{F})}{\partial q_\kappa} \right\rangle \\
&= -\frac{i\Delta}{\hbar} \left\langle \mathfrak{F} \otimes \partial / \partial q_\kappa \left\{ \sum_{i=-m'}^m \left[\mathcal{L} \left(q'_{i+1} - \frac{q'_i}{t'_{i+1}} - t'_i, q'_{i+1} \right) \otimes (t'_{i+1} - t'_i) \right] \right\} \right\rangle \\
&= \left\langle \frac{1}{\sqrt{g}} \frac{\partial(\sqrt{g} \mathfrak{F})}{\partial q_\kappa} \right\rangle \\
&= \left\| -\frac{i\Delta}{\hbar} \left\langle \mathfrak{F} \left\{ \mathcal{L}_{\hat{q}} \left(\hat{q}_{\kappa+1} - \frac{\hat{q}_\kappa}{\hat{t}_{\kappa+1}} - \hat{t}_\kappa, \hat{q}_{\kappa+1} \right)^2 - \mathcal{L}_{\hat{q}} \left(\hat{q}_\kappa - \frac{\hat{q}_{\kappa-1}}{\hat{t}_\kappa} - \hat{t}_{\kappa-1}, \hat{q}_\kappa \right)^2 \right. \right. \right. \\
&\quad \left. \left. - (\hat{t}_\kappa - \hat{t}_{\kappa+1})^2 \otimes \mathcal{L}_{\hat{q}} \left(\left[\frac{q}{t} \right]_{\kappa+1} - \left[\frac{q}{t} \right]_{\kappa+1}, \left[\frac{q}{t} \right]_{\kappa} \right) - (\hat{t}_\kappa - \hat{t}_{\kappa+1}) \right\} \right\rangle \\
&= \frac{i\Delta}{\hbar} \left\langle \mathfrak{G}_1 \left[m \left(\chi_{\kappa+1} - \frac{\chi_\kappa}{t_{\kappa+1}} - t_\kappa \right) - m \left(\chi_{\kappa-1} + \frac{\chi_\kappa}{t_{\kappa-1}} - t_\kappa \right) \otimes \mathcal{V}' \chi_{\kappa} \bar{\delta} \right] \mathfrak{G}_2 \right\rangle \hbar \varepsilon \omega^2 \left\| \left\langle \exp \int \langle \varphi^\odot | \alpha_\Delta^\Pi | e^{\frac{i\delta\Delta}{\hbar} \mathcal{F}_\kappa} \psi_\Delta d \right\rangle_{\text{vol } \phi_m} \right\rangle 1 \\
&= \frac{1}{2\delta\varepsilon im} + \frac{\left\langle \frac{\hbar}{2\varepsilon i \psi'_t} \langle \mathcal{V}'(x_\kappa) \rangle \partial \mathcal{F} \right\rangle^2}{\left| \partial \chi_\kappa \otimes \frac{\partial \delta}{\psi'_t} \right|_{\bar{\delta}}^2} / (2\pi \hbar \varepsilon i) \xi^2 d\xi
\end{aligned}$$



$$\begin{aligned}
[\mathcal{J}_{\mu\nu}, \mathcal{J}_{\rho\sigma}] &= i(\eta_{\mu\rho}\mathcal{J}_{\nu\sigma} + \eta_{\nu\sigma}\mathcal{J}_{\mu\rho} - \eta_{\nu\rho}\mathcal{J}_{\mu\sigma} - \eta_{\mu\sigma}\mathcal{J}_{\nu\rho}), [\mathcal{J}_{ij}, \mathcal{J}_{kl}] \\
&= i(\eta_{ik}\mathcal{J}_{jl} + \eta_{jl}\mathcal{J}_{ik} - \eta_{jk}\mathcal{J}_{il} - \eta_{il}\mathcal{J}_{jk}), [\mathcal{J}_{ij}, \mathcal{J}_{k4}] = i(\eta_{ik}\mathcal{J}_{j4} - \eta_{jk}\mathcal{J}_{i4}), [\mathcal{J}_{i4}, \mathcal{J}_{j4}] \\
&= i\mathcal{J}_{ij}, \langle \Theta_{ij} \hbar \mathcal{J}_{ij} \Theta_{kl} \rangle = i\hbar(\eta_{ik}\Theta_{jl} + \eta_{jl}\Theta_{ik} - \eta_{jk}\Theta_{il} - \eta_{il}\Theta_{jk}), \langle \Theta_{ij} \hbar \mathcal{J}_{ij} \chi_k \rangle \\
&= i\hbar(\eta_{ik}\chi_j - \eta_{jk}\chi_i), [\chi_i, \chi_j] = \frac{i\lambda^2}{\hbar\Theta_{ij}} i\eta_{ij}
\end{aligned}$$

$$\begin{aligned}
[\mathcal{J}_{mn}, \mathcal{J}_{rs}] &= i(\eta_{mr}\mathcal{J}_{ns} + \eta_{ns}\mathcal{J}_{mr} - \eta_{nr}\mathcal{J}_{ms} - \eta_{ms}\mathcal{J}_{nr}), \mathcal{P}_i = \frac{\hbar}{\lambda} \mathcal{J}_i, \langle \Theta_{ij}, \Theta_{kl} \rangle \\
&= i\hbar(\eta_{ik}\Theta_{jl} + \eta_{jl}\Theta_{ik} - \eta_{jk}\Theta_{il} - \eta_{il}\Theta_{jk}), \langle \Theta_{ij}, \mathcal{P}_k \rangle = i\hbar(\eta_{ik}\mathcal{P}_j - \eta_{jk}\mathcal{P}_i), \langle \mathcal{P}_i, \mathcal{P}_j \rangle \\
&= \frac{i\hbar}{\lambda^2\Theta_{ij}}, \langle \Theta_{ij}, \chi_k \rangle = i\hbar(\eta_{ik}\chi_j - \eta_{jk}\chi_i), \frac{\langle \chi_i, \chi_j \rangle i\lambda^2}{\hbar} \Theta_{ij} \langle \chi_i, \mathcal{P}_j \rangle = i\hbar\eta_{ij}\hbar, \langle \Theta_{ij}, \hbar \rangle \\
&= \langle \chi_i, \hbar \rangle = \frac{i\lambda^2}{\hbar} \mathcal{P}_i, \langle \mathcal{P}_i, \hbar \rangle = \frac{i\lambda^2}{\hbar} \chi_i
\end{aligned}$$

$$\begin{aligned}
\eta_{\mu\nu}\chi^\mu\chi^\nu = \delta\mathcal{R}^2, [\mathcal{J}_{\mathfrak{M}\mathfrak{N}}, \mathcal{J}_{\mathcal{R}\mathcal{S}}] &= i(\eta_{\mathfrak{M}\mathcal{R}}\mathcal{J}_{\mathfrak{N}\mathcal{S}} + \eta_{\mathcal{N}\mathcal{S}}\mathcal{J}_{\mathfrak{M}\mathcal{R}} - \eta_{\mathfrak{N}\mathcal{R}}\mathcal{J}_{\mathcal{M}\mathcal{S}} - \eta_{\mathcal{M}\mathcal{S}}\mathcal{J}_{\mathfrak{N}\mathcal{R}}), [\mathcal{J}_{ij}, \mathcal{J}_{kl}] \\
&= i(\eta_{ik}\mathcal{J}_{jl} + \eta_{jl}\mathcal{J}_{ik} - \eta_{jk}\mathcal{J}_{il} - \eta_{il}\mathcal{J}_{jk})
\end{aligned}$$

$$\begin{aligned}
\langle \Theta_{ij}, \Theta_{kl} \rangle &= i\hbar(\eta_{ik}\Theta_{jl} + \eta_{jl}\Theta_{ik} - \eta_{jk}\Theta_{il} - \eta_{il}\Theta_{jk}), \langle \Theta_{ij}, \mathcal{Q}_k \rangle = \frac{i}{\hbar}(\eta_{ik}\mathcal{Q}_j - \eta_{jk}\mathcal{Q}_i), \langle \Theta_{ij}, \chi_k \rangle \\
&= \frac{i}{\hbar}(\eta_{ik}\chi_j - \eta_{jk}\chi_i), \langle \Theta_{ij}, \mathcal{P}_k \rangle = \frac{i}{\hbar}(\eta_{ik}\mathcal{P}_j - \eta_{jk}\mathcal{P}_i), (\mathcal{Q}_i, \mathcal{Q}_j) = \frac{i\hbar}{\lambda^2} \Theta_{ij}, (\mathcal{Q}_i, \chi_j) \\
&= \frac{i\hbar}{\lambda^2} \eta_{ij} q, (\mathcal{Q}_i, \mathcal{P}_j) = \frac{i\hbar}{\lambda^2} \eta_{ij} p, (\mathcal{Q}_i, q) = \frac{i\hbar}{\lambda^2 \chi_i}, (\mathcal{Q}_i, p) = i\mathcal{P}_i, (\chi_i, \chi_j) \\
&= \frac{i\lambda^2}{\hbar} \Theta_{ij}, (\chi_i, \mathcal{P}_j) = -i\hbar\eta_{ij}h, (\chi_i, q) = -\frac{i\lambda^2}{\hbar} \mathcal{Q}_i, (\chi_i, h) = -\frac{i\lambda^2}{\hbar} \mathcal{P}_i, (\mathcal{P}_i, \mathcal{P}_j) \\
&= \frac{i\delta\hbar}{\lambda^2\Theta_{ij}}, (\mathcal{P}_i, p) = i\delta\mathcal{Q}_i, (\mathcal{P}_i, h) = \frac{i\delta\hbar}{\lambda^2\chi_i}, [q, p] = i\hbar, [q, h] = ip, [p, h] = -i\delta q
\end{aligned}$$

$$\gamma_1 \begin{bmatrix} 0 & \dots & 1 \\ \vdots & \ddots & \vdots \\ 1 & \dots & 1 \end{bmatrix} \gamma_2 \begin{bmatrix} 1 & \dots & -1 \\ \vdots & \ddots & \vdots \\ -1 & \dots & 0 \end{bmatrix} \gamma_3 \begin{bmatrix} -1 & \dots & 1 \\ \vdots & \ddots & \vdots \\ 1 & \dots & 0 \end{bmatrix} \gamma_4 \begin{bmatrix} 0 & \dots & 1 \\ \vdots & \ddots & \vdots \\ -1 & \dots & -1 \end{bmatrix}$$



$$\begin{aligned}
(\mathcal{M}_{ab}, \mathcal{M}_{cd}) &= \eta_{bc}\mathcal{M}_{ad} + \eta_{ad}\mathcal{M}_{bc} - \eta_{ac}\mathcal{M}_{bd} - \eta_{bd}\mathcal{M}_{ac}, (\mathcal{M}_{ab}, \mathcal{P}_c) = \eta_{bc}\mathcal{P}_a - \eta_{ac}\mathcal{P}_b, (\mathcal{M}_{ab}, \mathcal{K}_c) \\
&= \eta_{bc}\mathcal{K}_a - \eta_{ac}\mathcal{K}_b, (\mathcal{P}_a, \mathcal{D}) = \mathcal{P}_a, (\mathcal{K}_a, \mathcal{D}) = -\mathcal{K}_a, (\mathcal{K}_a, \mathcal{P}_b) \\
&= -2(\eta_{ab}\mathcal{D} + \mathcal{M}_{ab}), (\mathcal{M}_{ab}, \mathcal{M}_{cd}) = \frac{1}{2}(\eta_{ac}\eta_{bd} - \eta_{bc}\eta_{ad}) - i\epsilon_{abcd}\mathcal{D}, (\mathcal{M}_{ab}, \mathcal{P}_c) \\
&= i\epsilon_{abcd}\mathcal{P}^d, (\mathcal{M}_{ab}, \mathcal{K}_c) = -i\epsilon_{abcd}\mathcal{K}^d, (\mathcal{M}_{ab}, \mathcal{D}) = 2\mathcal{M}_{ab}\mathcal{D}, (\mathcal{P}_a, \mathcal{K}_b) \\
&= 4\mathcal{M}_{ab}\mathcal{D} + \eta_{ab}, (\mathcal{K}_a, \mathcal{K}_b) = (\mathcal{P}_a, \mathcal{P}_b) = -\eta_{ab}, (\mathcal{P}_a, \mathcal{D}) = (\mathcal{K}_a, \mathcal{D}) = 1
\end{aligned}$$

$$\begin{aligned}
\mathcal{S} &= Tr([\chi_\mu, \chi_\nu] - \kappa^2\Theta_{\mu\nu})([\chi_\rho, \chi_\sigma] - \kappa^2\Theta_{\rho\sigma})\epsilon^{\mu\nu\rho\sigma}, \epsilon^{\mu\nu\rho\sigma} \\
&= [\chi_\nu(\chi_\rho, \chi_\sigma) - \kappa^2\Theta_{\rho\sigma}], \epsilon^{\mu\nu\rho\sigma}([\chi_\rho, \chi_\sigma] - \kappa^2\Theta_{\rho\sigma}) = 1, \mathcal{S} \\
&= Trtr \epsilon^{\mu\nu\rho\sigma}([\chi_\mu + \Lambda_\mu, \chi_\nu + \Lambda_\nu] \\
&\quad - \kappa^2(\Theta_{\mu\nu} + \mathfrak{B}_{\mu\nu})) \otimes ([\chi_\rho + \Lambda_\rho, \chi_\sigma + \Lambda_\sigma] - \kappa^2(\Theta_{\rho\sigma} + \mathfrak{B}_{\rho\sigma})), \Lambda_\mu \\
&= \alpha_\mu \otimes 1_4 + \omega_\mu^{\alpha\beta} \otimes \mathcal{M}_{\alpha\beta} + e_\mu^\alpha \otimes \mathcal{P}_\alpha + \beta_\mu^\alpha \otimes \mathcal{K}_\alpha + \tilde{\alpha}_\mu \otimes \mathcal{D}, \mathcal{S} \\
&= Trtr \epsilon^{\mu\nu\rho\sigma} \left([\chi_\mu, \chi_\nu] - \frac{i\lambda^2}{\hbar} \Theta_{\mu\nu} \right) \left([\chi_\rho, \chi_\sigma] - \frac{i\lambda^2}{\hbar} \Theta_{\rho\sigma} \right) \cong Trtr \epsilon^{\mu\nu\rho\sigma} \hat{\mathcal{F}}_{\mu\nu} \hat{\mathcal{F}}_{\rho\sigma}, \mathcal{S} \\
&= Trtr [\lambda\phi(\chi)\epsilon^{\mu\nu\rho\sigma} \hat{\mathcal{F}}_{\mu\nu} \hat{\mathcal{F}}_{\rho\sigma} + \eta(\phi(\chi)^2) - \lambda^{-2} 1_\eta \otimes 1_4], \mathcal{S}_{br} \\
&= Tr \left(\frac{\sqrt{2}}{4} \epsilon_{abcd} \mathcal{R}_{mn}^{\alpha\beta} \mathcal{R}_{rs}^{cd} - 4\mathcal{R}_{mn} \tilde{\mathcal{R}}_{rs} \right) \epsilon^{mnr} \\
\hat{\mathcal{F}}_{\mu\nu} &= \mathcal{R}_{\mu\nu} \otimes 1_4 + \frac{1}{2} \mathcal{R}_{\alpha\beta}^{\mu\nu} \otimes \mathcal{M}_{\alpha\beta} + \tilde{\mathcal{R}}_{\alpha\beta}^{\mu\nu} \otimes \mathcal{P}_\alpha + \mathcal{R}_\alpha^{\mu\nu} \otimes \mathcal{K}_\alpha + \tilde{\mathcal{R}}_\alpha^{\mu\nu} \otimes \mathcal{D} \\
\hat{\mathcal{F}}_{\rho\sigma} &= \mathcal{R}_{\rho\sigma} \otimes 1_4 + \frac{1}{2} \mathcal{R}_{\alpha\beta}^{\rho\sigma} \otimes \mathcal{M}_{\alpha\beta} + \tilde{\mathcal{R}}_{\alpha\beta}^{\rho\sigma} \otimes \mathcal{P}_\alpha + \mathcal{R}_\alpha^{\rho\sigma} \otimes \mathcal{K}_\alpha + \tilde{\mathcal{R}}_\alpha^{\rho\sigma} \otimes \mathcal{D} \\
\phi(\chi) &= \Phi(\chi) \otimes 1_4 + \phi^{\alpha\beta}(\chi) \otimes \mathcal{M}_{\alpha\beta} + \hat{\phi}^{\alpha\beta}(\chi) \otimes \mathcal{P}_\alpha + \phi^{\alpha\beta}(\chi) \otimes \mathcal{K}_\alpha + \tilde{\phi}^{\alpha\beta}(\chi) \otimes \mathcal{D}, \Phi(\chi) \\
&= \hat{\phi}(\chi) \otimes \mathcal{D} |_{\tilde{\phi} = -2\lambda^{-1}} = -2\lambda^{-1} 1_\eta \otimes \mathcal{D}
\end{aligned}$$

$$\frac{ds}{dt} = \frac{\sqrt{\frac{g_{\mu\nu} dx^\mu dx^\nu}{dt}} d^2 x^\mu}{dt^2} + \frac{\Gamma_{\alpha\beta}^\mu dx^\alpha dx^\beta}{dt} = \frac{\lambda(t) dx^\mu}{dt^2} d^2 \alpha = \frac{\lambda d\alpha}{dt} \Rightarrow \frac{d^2 x^\mu}{d\alpha^2} + \Gamma_{\alpha\beta}^\mu dx^\alpha dx^\beta / d\alpha$$



$$\begin{aligned}
ds^2 &= \left(1 - \frac{2m}{r}\right) dt_{\delta^2} + \left(1 - \frac{2m}{r}\right)^{-1} dr^2 + r^2 d\sigma^2 = d\theta^2 + \sin^2\theta d\phi^2, t_- \\
&= t_{\delta} - 2m \ln|r - 2m|, t_+ \\
&= t_{\delta} + 2m \ln|r - 2m|, t_+ = t_- \\
&+ 4m \ln|r - 2m| ds^2 = ds_{0\pm}^2 + \frac{2m}{r} (k_{\pm\mu} dx^{\mu})^2, ds_{0\pm}^2 = dr^2 + r^2 d\sigma^2 - dt_{0\pm}^2
\end{aligned}$$

$$\begin{aligned}
k_{\pm} &= k_{\pm\mu} dx^{\mu} = \pm r - \frac{2m}{r} + 2m dr - dt_{\pm}, k_{\pm} = k_{\pm}^{\mu} \partial_{\mu} = \pm \partial_r + \partial_{t_{\pm}}, k_{\pm\mu}^* dx^{\mu} \\
&= \pm \frac{r - 2m}{r + 2m} dr - dt_{\pm}, k_{\pm}^* = k_{\mp\mu}^* \partial_{\mu} = \pm \frac{r - 2m}{r + 2m} \partial_r + \partial_{t_{\pm}}
\end{aligned}$$

$$\begin{aligned}
ds^2 &= dr^2 - dt^2 + \frac{1}{r} (dr + dt)^2 = (dr + dt) \left(dr - dt + \frac{1}{r} (dr + dt) \right) \\
&= r - \frac{1}{r} d(r + t) \left(r + \frac{1}{r} - 1 dr - dt \right) = r - \frac{1}{r} d\mu dv
\end{aligned}$$

$$ds^2 = \frac{4}{r} e^{-r} d\mathcal{U} d\mathcal{V} + r^2 d\sigma^2 \Rightarrow \frac{32m^4}{r} e^{-r/2m} d\mathcal{U} d\mathcal{V} + r^2 d\sigma^2$$

$$\begin{aligned}
ds^2 &= ds_0^2 + \frac{2mr}{\Sigma} k^2, k = dr + \alpha \sin^2\theta d\phi + dt, ds_0^2 \\
&= dr^2 + \Sigma d\theta^2 + (r^2 + \alpha^2) \sin^2\theta d\phi + 2\alpha \sin^2\theta d\phi dt - dt^2, \Sigma = r^2 + \alpha^2 \cos^2\theta
\end{aligned}$$

$$\begin{aligned}
ds^2 &= dx^2 + dy^2 + dz^2 - dt^2 + \frac{2mr^3}{r^4} \\
&+ \alpha^2 z^2 \left[dt + \frac{z}{r} dz + \frac{r}{r^2} + \alpha^2 (xdx + ydy) + \alpha/r^2 + \alpha^2 (xdx + ydy) \right]^2
\end{aligned}$$

$$ds^2 = -dt^2 + dr^2 + \frac{2mr}{r^2} + \frac{\alpha^2 (dr + dt)^2 dr}{dt} = r^2 - 2mr - \frac{\alpha^2}{r^2} + 2mr + \alpha^2$$

$$\begin{aligned}
ds^2 &= \frac{\Sigma}{\Delta} dr^2 - \frac{\Delta}{\Sigma(dt_{\delta} + \alpha \sin^2\theta d\phi_{\delta})^2} + \Sigma d\theta^2 \\
&+ \frac{\sin^2\theta}{\Sigma((r^2 + \alpha^2)d\phi_{\delta} - \alpha dt_{\delta})^2} \left(\partial_r + \frac{2mr}{\Delta \partial_t} - \frac{\alpha}{\Delta \partial_{\phi} \partial_{\theta} \partial_{\phi} \partial_t |_{\delta}} \right) ds^2 \\
&= \frac{\Sigma}{\Delta dr^2} - \frac{\Delta}{\Sigma} (dt_{\delta} + \alpha \sin^2\theta d\phi_{\delta})^2 + \Sigma d\theta^2 + \sin^2\theta / \Sigma ((r^2 + \alpha^2)d\phi_{\delta} - \alpha dt_{\delta})^2
\end{aligned}$$

$$g^{\mu\nu} \partial_{\mu} \partial_{\nu} = \frac{\Delta}{\Sigma} \partial_{r_{\delta}}^2 - \frac{1}{\Delta \Sigma ((r^2 + \alpha^2) \partial_{t_{\delta}} - \alpha \partial_{\phi_{\delta}})^2} + \frac{1}{\Sigma \partial_{\theta_{\delta}}^2} + 1 / \Sigma \sin^2\theta (\partial_{\phi_{\delta}} - \alpha \sin^2\theta \partial_{t_{\delta}})^2$$



$$\begin{aligned}
k_- &= (dt_\delta + \alpha \sin^2 \theta d\phi_\delta) + (\Sigma \Delta^{-1}) dr, k_\pm \\
&= \mp \Sigma dr + [\Delta(dt + \alpha \sin^2 \theta d\phi) + (-2mr + \alpha^2 \sin^2 \theta) dr] k_\pm \\
&= \mp (\Delta \partial_r + 2mr \partial_t - \alpha \partial_\phi) + ((r^2 + \alpha^2) \partial_t - \alpha \partial_\phi) \\
\frac{dr}{dt} &= r^2 - 2mr + \frac{\alpha^2}{r^2} + 2mr + \alpha^2, \frac{d\phi}{dt} = \frac{2\alpha}{r^2} + 2mr + \alpha^2
\end{aligned}$$

Los agujeros negros cuánticos, suponen tratos de colisión, superposición o entrelazamiento, según sea el caso, en el que interactúan partículas o antipartículas deformantes y deformadas, en el primer caso, a propósito de su masa exponencial o de su energía potencial o de su energía cinética, según corresponda, y en el segundo caso, a propósito de su masa o energía cinética o potencial ligeras, según corresponda. Todo esto, depende esencialmente del campo cuántico de que se trate.

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