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FORMALIZACIÓN MATEMÁTICA Y EN FÍSICA DE PARTÍCULAS, EN RELACIÓN A LA BRECHA DE MASA Y LA CURVATURA GEOMÉTRICA DE LOS CAMPOS CUÁNTICOS

**MATHEMATICAL FORMALIZATION AND PARTICLE PHYSICS,
IN RELATION TO THE MASS GAP AND THE GEOMETRIC
CURVATURE OF QUANTUM FIELDS**

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Formalización Matemática y en Física de Partículas, en Relación a la Brecha de Masa y la Curvatura Geométrica de los Campos Cuánticos

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RESUMEN

En recientes manuscritos, este investigador ha formulado alternativas de solución al Problema del Milenio de Yang – Mills, intentando unificar, desde la teoría cuántica de campos hasta las teorías de la relatividad general y especial respectivamente, sin desprendernos de cuestiones tan elementales como las representaciones en *álgebra de Lie*, de cuyo resultado, se ha concluido en lo fundamental, que toda partícula o antipartícula, con masa o sin masa, según sea el caso, supera el estado de vacío, demostrando una brecha de masa positiva, esto es, cuando se aproxima o supera la velocidad de la luz, deformando así, el campo cuántico en el que interactúa, repercutiendo en las trayectorias de las partículas o antipartículas circundantes. Ahora bien, el propósito de esta investigación, es proponer modelos hipotéticos para campos de Yang – Mills abelianos y no abelianos, grupos de gauge y Lie usando distintos operadores para espacios en cuatro dimensiones $\mathbb{R}^4$, a través de los cuales, quedará demostrado, que la brecha de masa de una partícula o antipartícula con o sin masa, siempre arroja un valor positivo superior a cero.

Palabras clave: física de partículas, campos de gauge, teorías de calibre, grupos de Lie, libertad asintótica, dimensión $\mathbb{R}^4$, campos de Yang Mills abelianos y no abelianos, superficie espacial, superficie temporal, operador de Casimir, transformación de Lorentz, ecuación de Callan-Symanzik, integral de trayectoria, representación de espinores.

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Mathematical Formalization and Particle Physics, in Relation to the Mass Gap and the Geometric Curvature of Quantum Fields

ABSTRACT

In recent manuscripts, this researcher has formulated alternative solutions to the Yang-Mills Millennium Problem, trying to unify, from quantum field theory to the theories of general and special relativity respectively, without detaching ourselves from such elementary questions as representations in Lie algebra, from the result of which it has been concluded in the main, that every particle or antiparticle, with or without mass, as the case may be, exceeds the vacuum state, demonstrating a positive mass gap, that is, when it approaches or exceeds the speed of light, thus deforming the quantum field in which it interacts and affecting the trajectories of the surrounding particles or antiparticles. Now, the purpose of this research is to propose hypothetical models for abelian and non-abelian Yang-Mills fields, gauge and lie groups using different operators for spaces in four dimensions $\mathbb{R}^4$, through which it will be demonstrated that the mass gap of a particle or antiparticle with or without mass, it always yields a positive value greater than zero.

Keywords: particle physics, gauge fields, caliber theories, Lie groups, asymptotic freedom, $\mathbb{R}^4$ dimension, abelian and non-abelian Yang Mills fields, spatial surface, time surface, Casimir operator, Lorentz transformation, Callan-Symanzik equation, trajectory integral, spinor representation.



INTRODUCCIÓN

Preliminarmente, cabe precisar que se trabajará en campos cuánticos en dimensión $\mathbb{R}^4$, en estructuras de gauge específicas, a propósito de sus transformaciones, con trayectorias orbitales arbitrarias, utilizando distintas métricas vectoriales, espaciales, temporales y operadores cuánticos de campo, todo esto, en superficies de espacio – tiempo cuatridimensionales, por lo que, no solamente se recurrirá a la teoría cuántica de campos de Yang – Mills, sino también a las teorías de la relatividad general y especial y otras leyes propias de la física y de las matemáticas puras, todo esto, con la finalidad de demostrar, que la brecha de masa, en un campo de Yang – Mills, esto es, cuando una partícula o antipartícula con o sin masa, según sea el caso, supera el cero absoluto, arroja un salto de energía cuyo resultado siempre es positivo. En el apartado de Resultados y Discusión, se desplegarán los sistemas matemáticos y de la física de partículas correspondientes que sostienen la hipótesis contenida en este Artículo Científico y en definitiva, en los trabajos que anteceden a éste.

Para estos efectos, se han diseñado campos cuánticos hipotéticos, con superficies espaciales y temporales arbitrarias, todo esto, con la finalidad, de demostrar la existencia de la brecha de masa positiva y paralelamente, la curvatura geométrica de los campos cuánticos y los agujeros deformantes de los referidos campos.

METODOLOGÍA

La teorización desplegada en el presente manuscrito, resulta de la aplicación de una metodología de investigación integral, esto es, bajo un enfoque híbrido, tanto desde el punto de vista cualitativo como en su dimensión cuantitativa. El tipo de investigación que ha sido desarrollado a lo largo del presente Artículo Científico, es esencialmente predictivo, a la luz de la física teórica, aunque comporta también un carácter empírico o experimental. Por otro lado, las líneas de investigación adoptadas para la formulación del estado del arte, se ajustan al constructivismo. Cabe indicar, que no existe población de estudio en la medida en que el presente artículo científico, no es de carácter sociológico o social, más aun, en mérito a su impacto en la realidad de transformación. Tampoco se han implementado técnicas de recolección de información, tales como encuestas, entrevistas, etc, salvo revisión bibliográfica, a razón del campo de investigación abordado. Adicionalmente a lo antes expuesto, es preciso resaltar, que el material de apoyo es meramente bibliográfico. La técnica metodológica, dada la complejidad de la

temática escrutada, es deductiva, pues la teorización en sentido estricto, ha sido desarrollada desde principios y premisas generales que son inherentes a la física de partículas en sentido lato. Finalmente, para efectos de construir y desarrollar las ecuaciones constantes en el presente artículo científico, se ha tomado en consideración la teoría cuántica de campos, muy especialmente, en tratándose de los campos de Yang – Mills, sin perjuicio de los demás sistemas de recalibración deducidos y esbozados a lo largo del presente Artículo Científico.

RESULTADOS Y DISCUSIÓN (Formulación Matemática y en física de partículas)

En un grupo cuántico de estructura G y bajo el algebra de Lie, obtenemos lo que sigue:

$$\langle \mathfrak{U} | \mathfrak{B} \rangle = -\mathfrak{Tr}_{\text{mat}(\overline{n})}(\mathbb{C}) (\mathfrak{U} | \mathfrak{B})$$

Cuya transformación de gauge se reduce a lo que sigue:

$$\mathfrak{U} \cdot \Omega = \mathfrak{U}^\Omega = \Omega^{-1} \mathfrak{d} \Omega + \Omega^{-1} \mathfrak{U} \Omega$$

De cuyo resultado se obtienen la totalidad de las órbitas.

Por otro lado, usando la métrica de Riemann en un volumen espacial específico, y utilizando el operador de Hodge, tenemos:

$$u \wedge \star v = (u, v)_q \mathfrak{d} \omega$$

Cuyas secciones se definen así:

$$\langle u | v \rangle = \int_{\mathfrak{M}} (u, v)_q \mathfrak{d} \omega$$

De lo que obtenemos lo que sigue:

$$|u \otimes \mathfrak{E}|^2 = -\mathfrak{Tr}(\mathfrak{E} \cdot \mathfrak{E}) u \wedge \star v = -\mathfrak{Tr}(\mathfrak{E} \cdot \mathfrak{E}) (u, v)_q \mathfrak{d} \omega$$

Cuya solución de Yang – Mills, es la que sigue:

$$\begin{aligned} & \mathfrak{S}_{\mathfrak{YM}}(\mathfrak{U}) \int_{\mathfrak{M}} |\mathfrak{d} \mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^2 \\ & \mathfrak{I} \exp \left[\int_{\mathcal{C}} \mathfrak{U} \right] - \frac{1}{3} \mathfrak{Tr} \int_{\mathfrak{U} \in \mathfrak{U}_{\mathfrak{M},g}/G} \mathfrak{I} \exp \left[\int_{\mathcal{C}} \mathfrak{U} \right] e^{-1/2 \mathfrak{S}_{\mathfrak{YM}}(\mathfrak{U})} \mathfrak{D} \mathfrak{U} \end{aligned}$$

En la que la curvatura de la superficie espacial y de la superficie temporal se expresan de la siguiente manera:

$$\begin{aligned} \mathfrak{D}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U} &= \sum_{\alpha} \sum_{1 \leq i < j \leq 4} \alpha_{i,j,\alpha} \otimes \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j \otimes \mathfrak{E}^{\alpha} + \sum_{\alpha, \beta} \sum_{1 \leq i < j \leq 4} \alpha_{i,\alpha} \alpha_{j,\beta} \otimes \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j \otimes \langle \mathfrak{E}^{\alpha} | \mathfrak{E}^{\beta} \rangle \\ &+ \sum_{\alpha} \sum_{j=1} \alpha_{0;j,\alpha} \otimes \mathfrak{d}\mathfrak{x}^0 \wedge \mathfrak{d}\mathfrak{x}^j \otimes \mathfrak{E}^{\alpha} \end{aligned}$$

Cuyas permutaciones y transformaciones, se expresan así:

$$\begin{aligned} \mathcal{C}_{\gamma}^{\alpha\beta} &= -\mathfrak{I}\mathfrak{r}(\mathfrak{E}^{\gamma}(\mathfrak{E}^{\alpha}, \mathfrak{E}^{\beta})) \\ \mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U} &= \sum_{\gamma} \langle \sum_{1 \leq i < j \leq 4} \alpha_{i,j,\alpha} \otimes \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j + \sum_{\alpha < \beta} \sum_{1 \leq i < j \leq 4} \alpha_{i,\alpha} \alpha_{j,\beta} \mathcal{C}_{\gamma}^{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j \\ &+ \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j \sum_{j=1} \alpha_{0;j,\gamma} \otimes \mathfrak{d}\mathfrak{x}^0 \wedge \mathfrak{d}\mathfrak{x}^j \rangle \otimes \mathfrak{E}^{\gamma} \end{aligned}$$

$$\int_{\mathbb{R}^4}^{\infty} |\mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^2$$

$$\begin{aligned} &= \sum_{i < j} \int_{\mathbb{R}^4}^{\infty} \left\| \sum_{\alpha} \alpha_{i,j,\alpha}^2 + \sum_{\gamma} \sum_{\alpha < \beta, \hat{\alpha} < \hat{\beta}} \alpha_{i,\alpha} \alpha_{j,\beta} \alpha_{i,\hat{\alpha}} \alpha_{j,\hat{\beta}} \mathcal{C}_{\gamma}^{\alpha\beta} \mathcal{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} \right. \\ &\quad \left. + 4 \sum_{\alpha < \beta, \gamma} \alpha_{i,j,\gamma} \alpha_{i,\alpha} \alpha_{j,\beta} \mathcal{C}_{\gamma}^{\alpha\beta} \right\| \mathfrak{d}\omega + \sum_j \int_{\mathbb{R}^4}^{\infty} \sum_{\alpha} \alpha_{0;j,\alpha}^2 \mathfrak{d}\omega \end{aligned}$$

Campo Cuántico Abeliano

Usando el teorema de Stoke, tenemos:

$$\int_{\mathfrak{C}}^{\infty} \sum_{i=1}^4 \mathfrak{U}_i \otimes \mathfrak{d}\mathfrak{x}^i = \int_{\partial\mathfrak{C}}^{\infty} \sum_{i=1}^4 \mathfrak{U}_i \mathfrak{d}\mathfrak{x}^i = \int_{\mathfrak{C}}^{\infty} \mathfrak{d}\mathfrak{U} = \int_{\mathbb{R}^4}^{\infty} \mathfrak{d}\mathfrak{U} \times 4_{\mathfrak{C}} = \|\mathfrak{d}\mathfrak{U}, 4_{\mathfrak{C}}\|$$

Por lo que, la integral de Yang – Mills se expresa de la siguiente manera:

$$1/\int_{\mathfrak{U}}^{\infty} e^{-\|\mathfrak{d}\mathfrak{U}\|^2/2} \mathfrak{D}\mathfrak{U} \int_A^{\infty} e^{\sqrt{-1}(\mathfrak{d}\mathfrak{U}|4_{\mathfrak{C}})} e^{-1/2\|\mathfrak{d}\mathfrak{U}\|^2} \mathfrak{d}\mathfrak{U}$$

Cuyo cambio heurístico de variables, queda expresado así:

$$\frac{1}{\det \mathfrak{d}^{-1} \int_{\mathfrak{A}}^{\infty} e^{-\frac{\|\mathfrak{d}\mathfrak{A}\|^2}{2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A}) \det \mathfrak{d}^{-1} \int_A^{\infty} e^{\sqrt{-1} \langle \mathfrak{d}\mathfrak{A} | 4_{\mathfrak{G}} \rangle} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}$$

$$= \frac{1}{\int_{\mathfrak{A}}^{\infty} e^{-\frac{\|\mathfrak{d}\mathfrak{A}\|^2}{2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A}) \int_A^{\infty} e^{\sqrt{-1} \langle \mathfrak{d}\mathfrak{A} | 4_{\mathfrak{G}} \rangle} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}$$

Más, en dimensión $\mathbb{R}^4$ y aplicando la función delta de Dirac, tenemos:

$$\langle \mathfrak{F}, \mathfrak{X}_x \otimes \mathfrak{d}x^{\alpha} \wedge \mathfrak{d}x^{\beta} \rangle = \left\langle \sum_{0 \leq i \leq j \leq 4} \mathfrak{S}_{ij} \mathfrak{d}x^i \wedge \mathfrak{d}x^j, \mathfrak{X}_x \otimes \mathfrak{d}x^{\alpha} \wedge \mathfrak{d}x^{\beta} \right\rangle = \langle \mathfrak{S}_{\alpha\beta} \mathfrak{d}x^{\alpha} \wedge \mathfrak{d}x^{\beta}, \mathfrak{X}_x \otimes \mathfrak{d}x^{\alpha} \wedge \mathfrak{d}x^{\beta} \rangle = \mathfrak{S}_{\alpha\beta}(\chi)$$

$$\in \mathbb{R}^4$$

Cuyo gauge axial, según el teorema de Stokes, arroja como resultado lo que sigue:

$$\int_{\mathfrak{G}}^{\infty} \sum_{i=1}^4 \mathfrak{A}_i \otimes \mathfrak{d}x^i = \int_{\partial \mathfrak{G}}^{\infty} \sum_{i=1}^4 \mathfrak{A}_i \otimes \mathfrak{d}x^i = \int_{\mathfrak{G}}^{\infty} \mathfrak{d}\mathfrak{A}$$

$$= \int_{(0,1)^2}^{\infty} \mathfrak{d}s \mathfrak{d}t \left[\sum_{1 \leq i \leq j \leq 4} (\mathfrak{A}_{ij}(\sigma) |J_{ij}^{\sigma}|)(s, t) + \sum_{j=1}^4 (\mathfrak{A}_{0j}(\sigma) |J_{0j}^{\sigma}|)(s, t) \right]$$

$$= \int_{(0,1)^2}^{\infty} \mathfrak{d}s \mathfrak{d}t \langle \mathfrak{d}\mathfrak{A}, \sum_{0 \leq i \leq j \leq 4} \chi_{\sigma(s,t)} |J_{ij}^{\sigma}|(s, t) \otimes \mathfrak{d}x^i \wedge \mathfrak{d}x^j \rangle = \langle \mathfrak{d}\mathfrak{A}, \widetilde{\mathfrak{B}}_{\mathfrak{G}} \rangle$$

En la que $\widetilde{\mathfrak{B}}_{\mathfrak{G}}$ es igual a:

$$\widetilde{\mathfrak{B}}_{\mathfrak{G}} = \int_{(0,1)^2}^{\infty} \sum_{0 \leq i \leq j \leq 4} \mathfrak{d}s \mathfrak{d}t \chi_{\sigma(s,t)} |J_{ij}^{\sigma}|(s, t) \mathfrak{d}x^i \wedge \mathfrak{d}x^j$$

$$\langle \mathfrak{d}\mathfrak{A}, \sum_{0 \leq i \leq j \leq 4} \chi_{\sigma(s,t)} |J_{ij}^{\sigma}|(s, t) \otimes \mathfrak{d}x^i \wedge \mathfrak{d}x^j \rangle = \langle \mathfrak{d}\mathfrak{A}, \sum_{0 \leq i \leq j \leq 4} \chi_{\sigma(s,t)} \otimes \mathfrak{d}x^i \wedge \mathfrak{d}x^j |J_{ij}^{\sigma}|(s, t)$$

$$= \sum_{0 \leq i \leq j \leq 4} \langle \mathfrak{A}_{i,j}, \chi_{\sigma(s,t)} \rangle |J_{ij}^{\sigma}|(s, t) = \sum_{0 \leq i \leq j \leq 4} \langle \mathfrak{A}_{i,j}(\sigma(s,t)) \rangle |J_{ij}^{\sigma}|(s, t)$$

$$\frac{1}{3 \int_{\mathfrak{A}}^{\infty} e^{\sqrt{-1} \langle \mathfrak{d}\mathfrak{A} | \widetilde{\mathfrak{B}}_{\mathfrak{G}} \rangle} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}$$

$$- \frac{1}{3 \int_{\mathfrak{A}}^{\infty} e^{\sqrt{-1} \langle \mathfrak{d}\mathfrak{A} | \widetilde{\mathfrak{B}}_{\mathfrak{G}} \rangle} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2}} \det(\mathfrak{d}^{-1}) \mathfrak{D}(\mathfrak{d}\mathfrak{A}) + \frac{1}{3 \int_{\mathfrak{A}}^{\infty} e^{\sqrt{-1} \langle \mathfrak{d}\mathfrak{A} | \widetilde{\mathfrak{B}}_{\mathfrak{G}} \rangle} e^{-\frac{1}{2|\mathfrak{d}\mathfrak{A}|^2}} \mathfrak{D}(\mathfrak{d}\mathfrak{A})}$$



Siguiendo el mismo orden de ideas, un espacio de Schwartz, quedaría expresado así:

$$\mathcal{P}\mathfrak{r} = \left\| (m_1 m_2, \dots, m_n) \sum_{j=1}^n m_j = \mathfrak{r} \right\|$$

Más en dimensión $\mathbb{R}^4$, incorporando la función de Gauss y la métrica de Lebesgue, tenemos:

$$\mathfrak{f}(\mathfrak{x}) = \wp(\mathfrak{x}) \sqrt{\phi_{\kappa}}(\mathfrak{x})$$

$$\langle \mathfrak{f}, \mathfrak{g} \rangle = \int_{\mathbb{R}^4} \mathfrak{f} \cdot \mathfrak{g} d\lambda$$

$$\langle \mathfrak{z}^{\mathfrak{r}}, \hat{\mathfrak{z}}^{\mathfrak{r}} \rangle = \frac{1}{\pi \int_{\mathbb{C}} \mathfrak{z}^{\mathfrak{r}} \cdot \overline{\hat{\mathfrak{z}}^{\mathfrak{r}}} e^{-|z^2|} d\mathfrak{x} d\mathfrak{p}}, \mathfrak{z} = \mathfrak{x} + \sqrt{-1}\rho$$

Cuya función polinómica, se traduce a lo que sigue:

$$\mathcal{H}_{\wp_{\mathfrak{R}}}(\mathfrak{x}) = \hbar_i(\mathfrak{x}^0) \hbar_j(\mathfrak{x}^1) \hbar_k(\mathfrak{x}^2) \hbar_l(\mathfrak{x}^3), \wp_{\mathfrak{R}} = (i, j, k, l) \in \mathfrak{P}_{\mathfrak{R}}$$

Más, incorporando la métrica de Gauss, tenemos:

$$\bigcup_{\mathfrak{r}=0}^{\infty} \left| \frac{\mathcal{H}_{\wp_{\mathfrak{R}}}(\kappa \mathfrak{x}^0, \kappa \mathfrak{x}^1, \kappa \mathfrak{x}^2, \kappa \mathfrak{x}^3) \sqrt{\phi_{\kappa}}}{\sqrt{\wp_{\mathfrak{R}}!}} : \mathcal{P}_{\mathfrak{r}} \in \mathfrak{P}_{\mathfrak{r}} \right|$$

Por tanto, en dimensión $\mathbb{R}^4$, tenemos:

$$\mathfrak{S}_{\kappa}(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^3) = \langle \sum_{\alpha=1}^4 \mathfrak{F}_{\alpha} \otimes d\mathfrak{x}^{\alpha} : \mathfrak{F}_{\alpha} \in \mathfrak{S}_{\kappa}(\mathbb{R}^4) \rangle$$

$$\overline{\mathfrak{S}}_{\kappa}(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) = \left\| \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{F}_{\alpha\beta} \otimes d\mathfrak{x}^{\alpha} \wedge d\mathfrak{x}^{\beta} : \mathfrak{F}_{\alpha\beta} \in \overline{\mathfrak{S}}_{\kappa}(\mathbb{R}^4) \rangle \right\|$$

Cuyo operador de Hodge, se reduce a lo siguiente:

$$\langle \sum_{0 \leq \alpha \leq \beta \leq 4} |\mathfrak{F}_{\alpha\beta} \otimes d\mathfrak{x}^{\alpha} \wedge d\mathfrak{x}^{\beta}|, \sum_{0 \leq \alpha \leq \beta \leq 4} |\widehat{\mathfrak{F}}_{\alpha\beta} \otimes d\mathfrak{x}^{\alpha} \wedge d\mathfrak{x}^{\beta}| \rangle = \sum_{0 \leq \alpha \leq \beta \leq 4} |\mathfrak{F}_{\alpha\beta}, \widehat{\mathfrak{F}}_{\alpha\beta}|$$

Cuya transformación de Segal - Bargmann y demás funciones holomórficas, dan como resultado lo que sigue:

$$\langle \mathfrak{z}^{\mathcal{R}}, \mathfrak{z}^{\mathcal{R}'} \rangle = \frac{1}{\varpi \int_{\mathbb{C}} \mathfrak{z}^{\mathcal{R}} \cdot \overline{\mathfrak{z}^{\mathcal{R}'}} e^{-|\mathfrak{z}^2|} d\mathfrak{x} d\mathfrak{p}}, \mathfrak{z} = \chi + \sqrt{-1}\rho$$

$$\Psi_{\kappa} = \frac{\frac{\hbar_{\mathfrak{i}}(\kappa \cdot)}{\sqrt{\mathfrak{i}!} \hbar_{\mathfrak{j}}(\kappa \cdot)}}{\frac{\sqrt{\mathfrak{j}!} \hbar_{\mathfrak{k}}(\kappa \cdot)}{\sqrt{\mathfrak{k}!} \hbar_{\mathfrak{l}}(\kappa \cdot)}} = \sqrt{\phi_{\kappa}} \rightarrow \frac{\frac{\mathfrak{z}_0^{\mathfrak{i}}}{\sqrt{\mathfrak{i}!} \mathfrak{z}_1^{\mathfrak{j}}}}{\frac{\sqrt{\mathfrak{j}!} \mathfrak{z}_2^{\mathfrak{k}}}{\sqrt{\mathfrak{k}!} \mathfrak{z}_3^{\mathfrak{l}}}} = \mathfrak{f}_{\mathfrak{i},\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} \otimes \mathfrak{E}^{\alpha} \rightarrow \Psi_{\kappa}(\mathfrak{f}_{\mathfrak{i},\alpha}) \otimes \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} \otimes \mathfrak{E}^{\alpha}$$

En la que, en un espacio de Hilbert, se tiene:

$$\langle \sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{f}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\beta}, \sum_{0 \leq \alpha \leq \beta \leq 4} \hat{\mathfrak{f}}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\beta} \rangle = \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{f}_{\alpha\beta}, \hat{\mathfrak{f}}_{\alpha\beta} \rangle$$

$$\partial \sum_{\mathfrak{i}=1}^4 \mathfrak{f}_{\mathfrak{i}} \otimes \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} = \sum_{\mathfrak{i}=1}^4 \langle \partial_0 \mathfrak{f}_{\mathfrak{i}} \rangle \otimes \mathfrak{d}\mathfrak{x}^0 \wedge \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} + \sum_{1 \leq \mathfrak{i} \leq \mathfrak{j} \leq 4} [\partial_{\mathfrak{i}} \mathfrak{f}_{\mathfrak{j}} - \partial_{\mathfrak{j}} \mathfrak{f}_{\mathfrak{i}}] \otimes \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} \wedge \mathfrak{d}\mathfrak{x}^{\mathfrak{j}}$$

$$\Psi_{\kappa} = \sum_{1 \leq \mathfrak{i} \leq \mathfrak{j} \leq 4} \mathfrak{f}_{\mathfrak{i},\mathfrak{j},\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} \wedge \mathfrak{d}\mathfrak{x}^{\mathfrak{j}} \otimes \mathfrak{E}^{\alpha} \rightarrow \sum_{1 \leq \mathfrak{i} \leq \mathfrak{j} \leq 4} \Psi_{\kappa}[\mathfrak{f}_{\mathfrak{i},\mathfrak{j},\alpha}] \otimes \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} \wedge \mathfrak{d}\mathfrak{x}^{\mathfrak{j}} \otimes \mathfrak{E}^{\alpha}$$

$$\mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^1(\mathbb{R}^3) = \left\{ \sum_{\alpha=1}^4 \mathfrak{f}_{\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} : \mathfrak{f}_{\alpha} \in \mathcal{H}^2(\mathbb{C}^4) \right\}$$

$$\mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) = \left\{ \sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{F}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} \wedge \mathfrak{d}\mathfrak{x}^{\beta} : \mathfrak{F}_{\alpha\beta} \in \mathcal{H}^2(\mathbb{C}^4) \right\}$$

$$\langle \sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{F}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} \wedge \mathfrak{d}\mathfrak{x}^{\beta}, \sum_{0 \leq \alpha \leq \beta \leq 4} \widehat{\mathfrak{F}}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} \wedge \mathfrak{d}\mathfrak{x}^{\beta} \rangle = \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{F}_{\alpha\beta}, \widehat{\mathfrak{F}}_{\alpha\beta} \rangle = \int_{\mathbb{C}^4}^{\infty} \mathfrak{F}_{\alpha\beta} \widehat{\mathfrak{F}}_{\alpha\beta} \mathfrak{d}\lambda_4$$

Cuyos polinomios de Hermite, se satisfacen así:

$$\frac{\mathfrak{d}}{\mathfrak{d}\mathfrak{x} \left(\hbar_{\eta}(\chi) \mathfrak{E}^{-\frac{\mathfrak{x}^2}{4}} \right)} = |\chi \hbar_{\eta}(\mathfrak{x}) - \chi \hbar_{\eta+1}(\mathfrak{x}) - \mathfrak{x}/2 \hbar_{\eta}(\chi)| \mathfrak{E}^{-\frac{\mathfrak{x}^2}{4}} = \frac{\mathfrak{d}}{\mathfrak{d}\mathfrak{x} \left(\hbar_{\eta}(\chi) \mathfrak{E}^{-\frac{\mathfrak{x}^2}{4}} \right)}$$

$$= \left(\frac{1}{2\hbar_{\eta+1}(\mathfrak{x})} + \frac{\eta}{2\hbar_{\eta-1}(\mathfrak{x})} - \hbar_{\eta+1}(\mathfrak{x}) \right) \mathfrak{E}^{-\frac{\mathfrak{x}^2}{4}} = \left(\frac{\eta}{2\hbar_{\eta-1}(\mathfrak{x})} - \frac{1}{2\hbar_{\eta+1}(\mathfrak{x})} \right) \mathfrak{E}^{-\frac{\mathfrak{x}^2}{4}}$$

$$\mathfrak{d} \sum_{\alpha=1}^4 \mathfrak{F}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^{\alpha} = \sum_{\alpha=1}^4 \mathfrak{d}_0 \mathfrak{F}_0 \mathfrak{d}\mathfrak{x}^0 \wedge \mathfrak{d}\mathfrak{x}^{\alpha} + \sum_{1 \leq \mathfrak{i} \leq \mathfrak{j} \leq 4} \langle -1 \rangle^{\mathfrak{i}\mathfrak{j}} \langle \partial_{\mathfrak{i}} \mathfrak{f}_{\mathfrak{l}} - \partial_{\mathfrak{j}} \mathfrak{f}_{\mathfrak{j}} \rangle \mathfrak{d}\mathfrak{x}^{\mathfrak{i}} \wedge \mathfrak{d}\mathfrak{x}^{\mathfrak{j}}$$

$$\langle \mathfrak{f} \sqrt{\phi_{\kappa}}, \mathfrak{g} \sqrt{\phi_{\kappa}} \rangle = \int_{\mathbb{R}^4}^{\infty} \mathfrak{f} \mathfrak{g} \cdot \phi_{\kappa}, \mathfrak{d}\lambda$$

$$\left\{ \left| \frac{\hbar_{\mathfrak{i}}(\kappa \chi^0) \hbar_{\mathfrak{j}}(\kappa \chi^1) \hbar_{\mathfrak{k}}(\kappa \chi^2) \hbar_{\mathfrak{l}}(\kappa \chi^3)}{\sqrt{\mathfrak{i}!} \mathfrak{j}! \mathfrak{k}! \mathfrak{l}!} \sqrt{\phi_{\kappa}(\vec{\mathfrak{x}})} \right| : \vec{\mathfrak{x}} = (\chi^0, \chi^1, \chi^2, \chi^3) \in \mathbb{R}^4, \mathfrak{i}, \mathfrak{j}, \mathfrak{k}, \mathfrak{l}, \geq 0 \right\}$$



$$\langle \sum_{1 \leq \alpha \leq \beta \leq 4} \mathfrak{f}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^\alpha \wedge \mathfrak{d}\mathfrak{x}^\beta, \sum_{1 \leq \alpha \leq \beta \leq 4} \hat{\mathfrak{f}}_{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^\alpha \wedge \mathfrak{d}\mathfrak{x}^\beta \rangle = \langle \mathfrak{f}_{\alpha\beta}, \hat{\mathfrak{f}}_{\alpha\beta} \rangle$$

$$\mathfrak{d}\mathfrak{f} = \sum_{i=1}^4 \partial_0 \mathfrak{f}_i \otimes \mathfrak{d}\mathfrak{x}^0 \wedge \mathfrak{d}\mathfrak{x}^i + \sum_{1 \leq i \leq j \leq 4} (\partial_i \mathfrak{f}_j - \partial_j \mathfrak{f}_i) \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j = \mathfrak{C}_\gamma^{\alpha\beta} = -\mathcal{T}\mathfrak{r}(\mathfrak{C}^\gamma(\mathfrak{C}^\alpha, \mathfrak{C}^\beta)), \mathfrak{C}^\alpha, \mathfrak{C}^\beta, \mathfrak{C}^\gamma \in \mathfrak{g}$$

$$\begin{aligned} \mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U} = & \sum_{\gamma=1}^{\eta} (\sum_{j=1}^4 \alpha_{0:i,\gamma} \otimes \mathfrak{d}\mathfrak{x}^0 \wedge \mathfrak{d}\mathfrak{x}^i + \sum_{1 \leq i \leq j \leq 4} \alpha_{i:j,\gamma} \otimes \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j \\ & + \sum_{1 \leq i \leq j \leq 4} \sum_{1 \leq \alpha, \beta \leq \eta} \alpha_{i,\alpha} \alpha_{j,\beta} \mathfrak{C}_\gamma^{\alpha\beta} \otimes \mathfrak{d}\mathfrak{x}^i \wedge \mathfrak{d}\mathfrak{x}^j) \otimes \mathfrak{C}^\gamma \end{aligned}$$

Cuya forma bilineal, se distribuye así:

$$\langle \sum_{\alpha} \mathfrak{F}_{\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\alpha}, \sum_{\alpha} \mathfrak{G}_{\beta} \otimes \mathfrak{d}\mathfrak{x}^{\beta} \rangle_{\partial\kappa} = \kappa^4 \langle \frac{\partial}{\partial t} \left(\frac{i}{\hbar} \right) \sum_{\alpha} \mathfrak{F}_{\alpha} \otimes \mathfrak{d}\mathfrak{x}^{\alpha}, \sum_{\alpha} \mathfrak{G}_{\beta} \otimes \mathfrak{d}\mathfrak{x}^{\beta} \rangle$$

$$\psi_{\kappa}\colon \mathcal{H}_{\wp_{\mathfrak{R}}}(\kappa \mathfrak{x}^0,\kappa \mathfrak{x}^1,\kappa \mathfrak{x}^2,\kappa \mathfrak{x}^3)\sqrt{\phi_{\kappa}}/\sqrt{\wp_{\mathfrak{R}}!}\rightarrow \mathfrak{z}^{\mathfrak{p}_{\mathfrak{r}}}/\sqrt{\wp_{\mathfrak{R}}!}\equiv \mathfrak{z}_0^{i_0}\mathfrak{z}_1^{i_1}\mathfrak{z}_2^{i_2}\mathfrak{z}_3^{i_3}/\sqrt{i_0!i_1!i_2!i_3!}$$

Cuya isometría extendida, deriva en lo que sigue:

$$\psi_{\kappa}\left[\sum_{0\leq\alpha\leq\beta\leq4}\mathfrak{F}_{\alpha,\beta}\otimes\mathfrak{d}\mathfrak{x}^{\alpha}\wedge\mathfrak{d}\mathfrak{x}^{\beta}\right]=\sum_{0\leq\alpha\leq\beta\leq4}\psi_{\kappa}|\mathfrak{f}_{\alpha,\beta}|\otimes\mathfrak{d}\mathfrak{x}^{\alpha}\wedge\mathfrak{d}\mathfrak{x}^{\beta}$$

Ahora bien, un espacio abstracto de Wiener, se explica así:

$$\mu\kappa(\chi\in\mathfrak{P}^{-1}(\mathfrak{V}))=\left(\frac{\kappa}{2\varpi}\right)^{l/2}\int\limits_{\mathfrak{y}\in\mathcal{F}}e^{-\kappa|\varphi|^2/\psi}\mathfrak{d}\mathfrak{y}$$

$$\left\{\hat{\mathfrak{z}}^{\mathfrak{p}_{\mathfrak{r}}}\frac{\otimes \mathfrak{d}\mathfrak{x}^{\alpha}}{|\hat{\mathfrak{z}}^{\mathfrak{p}_{\mathfrak{r}}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa}}\colon \rho_{\mathfrak{r}}\in\wp_{\mathfrak{r}},\mathcal{R}\geq 0\right\}$$

$$\hat{\mathfrak{z}}^{\mathfrak{p}_{\mathfrak{r}}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}=\mathfrak{z}^{\mathfrak{p}_{\mathfrak{r}}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}-\frac{\sum_{\alpha=0}^4\langle \mathfrak{z}^{\mathfrak{p}_{\mathfrak{r}}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha},\frac{\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}}{|\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa}}\rangle \mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}}{|\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa}}+|\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa\psi}^2-|\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa}$$

$$+\langle \frac{\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}}{|\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa}}\rangle_{\partial\kappa}\bowtie 4\kappa(\mathfrak{r}-1)\sqrt{\wp_{\mathfrak{R}}!}\cdot \frac{\frac{(\mathfrak{r}+1)}{2}}{\frac{\kappa\sqrt{\wp_{\mathfrak{R}}!}}{\mathcal{R}^2}}/2$$

$$\leq 6\kappa(\mathfrak{r}-1)^2\sum_{\alpha=1}^4\langle \frac{\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}}{|\mathfrak{z}^{\rho_{\mathfrak{r}}^{\alpha,-}}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial\kappa}}\rangle_{\partial\kappa}\otimes \mathfrak{d}\mathfrak{x}^{\alpha}$$



$$\mu_{\kappa}(\|\mathcal{P}_b\chi\|>\epsilon)\leq \mu_{\kappa}\left(\vec{\mathfrak{z}}\in\mathfrak{B}\left(\frac{0,1}{2}\right)\right)\sum_{\mathfrak{s}}\sum_{\rho_{\tau}\geq\mathfrak{Q}^{\mathfrak{s}}}\mathbb{E}|\mathfrak{C}_{\mathfrak{s}}\alpha_{\rho_{\tau},\mathfrak{B}}^{\mathfrak{s}}|^4\left[\left|\mathfrak{z}^{\rho_{\tau}}\right|+6\kappa(\mathfrak{r}+1)^2\sum_{\alpha=1}^4\left|\mathfrak{z}^{\rho_{\tau}^{\alpha,-}}\right|>\epsilon\right]1/\sqrt{\kappa\epsilon}$$

$$\begin{aligned} |\langle \chi(\omega), \mathfrak{d}\mathfrak{x}^{\alpha} \rangle| &\leq \left| \sum_{\tau} \sum_{\rho_{\tau} \in \mathfrak{q}_{\tau}} \mathfrak{C}_{\rho_{\tau}, \alpha} \widehat{\omega}^{\rho_{\tau}} / |\widehat{\omega}^{\rho_{\tau}} \otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial \kappa} \right| \\ &\leq \mathfrak{R}^{\mathfrak{M}} \sum_{\tau \leq \mathcal{M}} \sum_{\rho_{\tau} \in \mathfrak{q}_{\tau}} |\mathfrak{C}_{\rho_{\tau}, \alpha}| \|\omega / \mathfrak{R}\|^{\rho_{\tau}} + 6\kappa(\mathfrak{r}+1)^2 \sum_{\alpha=1}^4 \|\omega / \mathfrak{R}\|^{\rho_{\tau}^{\alpha,-}} 2\mathcal{R}^{\mathfrak{I}} + 1/\kappa \sqrt{\mathfrak{f}\mathcal{P}_{\mathfrak{R}}!} \|\chi\| \end{aligned}$$

$$\begin{aligned} |\langle \kappa \partial \mathfrak{x}(\omega), \mathfrak{d}\mathfrak{x}^{\alpha} \wedge \mathfrak{d}\mathfrak{x}^{\beta} \rangle| &\leq \kappa \left| \sum_{\tau} \sum_{\rho_{\tau} \in \mathfrak{q}_{\tau}} \frac{\mathfrak{C}_{\rho_{\tau}, \beta} \partial_{\beta} \widehat{\omega}^{\rho_{\tau}}}{|\widehat{\omega}^{\rho_{\tau}} \otimes \mathfrak{d}\mathfrak{x}^{\beta}|_{\partial \kappa}} \right| + \kappa \left| \sum_{\tau} \sum_{\rho_{\tau} \in \mathfrak{q}_{\tau}} \frac{\mathfrak{C}_{\rho_{\tau}, \alpha} \partial_{\beta} \widehat{\omega}^{\rho_{\tau}}}{|\widehat{\omega}^{\rho_{\tau}} \otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial \kappa}} \right| \\ &\leq \mathfrak{R}^{\mathcal{M}} \sum_{\tau \leq \mathcal{M}} \sum_{\rho_{\tau} \in \mathfrak{q}_{\tau}} \frac{|\mathfrak{C}_{\rho_{\tau}, \beta}| + |\mathfrak{C}_{\rho_{\tau}, \alpha}| |\partial_{\beta} \widehat{\omega}^{\rho_{\tau}}|}{|\widehat{\omega}^{\rho_{\tau}} \otimes \mathfrak{d}\mathfrak{x}^{\alpha}|_{\partial \kappa}} \left| \frac{\omega}{\mathcal{R}} \right|^{\rho_{\tau}} \\ &\quad + 6\kappa(\mathfrak{r}+1)^2 \sum_{\alpha=1}^4 \left| \frac{\omega}{\mathcal{R}} \right|^{\rho_{\tau}^{\alpha,-}} 4(\mathfrak{r}+1) \mathcal{R}^{\mathfrak{I}} / \sqrt{\mathfrak{f}\mathcal{P}_{\mathfrak{R}}!} \\ &< (\mathcal{R}^{\mathcal{M}} + 1) \|\lambda\|^{\mathfrak{K}} \mathscr{G}_i \otimes \mathfrak{d}\mathfrak{x}^i \in \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4) \rightarrow \psi_{\kappa} |\partial_{\alpha} \mathscr{G}_{\beta} - \partial_{\beta} \mathscr{G}_{\alpha}| \end{aligned}$$

$$\begin{aligned} &|(\omega+\omega_0)^{\rho_{\tau}}-\omega_0^{\rho_{\tau}}| \\ &+6\kappa(\mathfrak{r}+1)^2\sum_{\alpha=1}^4|(\omega+\omega_0)^{\rho_{\tau}^{\alpha,-}}-|\omega_0^{\rho_{\tau}^{\alpha,-}}\langle\leq\mathfrak{R}^{\tau}|\omega\rangle\Big(\sup_{\omega\in\mathfrak{B}(0,\epsilon)}\Big)|\langle\mathfrak{x}(\omega+\omega_0) \\ &\quad-\rangle\mathfrak{x}(\omega_0),\mathfrak{d}\mathfrak{x}^{\alpha}| \\ &\leq\Big(\sup_{\omega\in\mathfrak{B}(0,\epsilon)}\Big)\sum_{\tau}\sum_{\rho_{\tau}\in\mathfrak{P}_{\tau}}|\mathfrak{C}_{\rho_{\tau}}|\big|2(\widehat{\omega+\omega_0})^{\rho_{\tau}}-\widehat{\omega_0}^{\rho_{\tau}}\big|\kappa\sqrt{\mathfrak{f}\mathcal{P}_{\mathfrak{R}}!} \\ &\quad+\frac{\Big(\sup_{\omega\in\mathfrak{B}(0,\epsilon)}\Big)\sum_{\alpha=1}^4\sum_{\tau}\sum_{\rho_{\tau}\in\mathfrak{P}_{\tau}}|\mathfrak{C}_{\rho_{\tau}}|\Big|6\kappa(\mathfrak{r}+1)^2\cdot2\big|(\widehat{\omega+\omega_0})^{\rho_{\tau}}-\widehat{\omega_0}^{\rho_{\tau}}\big|}{\sqrt{\kappa\mathfrak{f}\mathcal{P}_{\mathfrak{R}}!}} \\ &\leq\frac{4}{\kappa\Big(\sup_{\omega\in\mathfrak{B}(0,\epsilon)}\Big)\sum_{\tau}\sum_{\rho_{\tau}\in\mathfrak{P}_{\tau}}|\mathfrak{C}_{\rho_{\tau}}|\big|(4^{\tau}\mathcal{R}^{\tau}|\omega|\big|}}{\kappa\sqrt{\mathfrak{f}\mathcal{P}_{\mathfrak{R}}!}} \\ &\quad+2\kappa\Big(\sup_{\omega\in\mathfrak{B}(0,\epsilon)}\Big)\sum_{\alpha=1}^4\sum_{\tau}\sum_{\rho_{\tau}\in\mathfrak{P}_{\tau}}\frac{6\kappa(\mathfrak{r}+1)^2}{2^{\tau-2}|\mathfrak{C}_{\rho_{\tau}}|}(2^{\tau-2}\mathfrak{R}^{\tau-2}|\omega|)/\kappa\sqrt{\mathfrak{f}\mathcal{P}_{\mathfrak{R}}!} \\ &\leq2\mathfrak{C}(\omega_0)/\kappa\cdot\epsilon|\chi| \end{aligned}$$

$$\begin{aligned}
& \widetilde{\mu}_{\kappa} \left(\left(\sup_{\omega \in \mathfrak{B} \left(\omega_0, \frac{1}{\kappa} \right)} |\mathfrak{x}(\omega + \omega_0) - \mathfrak{x}(\omega_0), \mathfrak{d}\mathfrak{x}^{\alpha}| > \epsilon \right) \right. \\
& \quad \left. = \widetilde{\mu}_{\kappa} \left(\left(\sup_{\omega \in \mathfrak{B} \left(\omega_0, \frac{1}{\kappa} \right)} |\mathfrak{x}, \zeta_{\alpha}(\omega + \omega_0) - \zeta_{\alpha}(\omega_0), \mathfrak{d}\mathfrak{x}^{\alpha}| > \epsilon \right) \leq \widetilde{\mu}_{\kappa} \left(\frac{2\mathfrak{C}(\omega_0)}{\kappa \mathfrak{k} \|\mathfrak{x}\|} > \epsilon \right) \right. \\
& \quad \left. \rightarrow 0 \right.
\end{aligned}$$

$$\left\langle \frac{\partial_{\alpha} \mathfrak{z}_{\alpha}^{\eta}}{\sqrt{\eta!}}, \frac{\partial_{\beta} \mathfrak{z}_{\beta}^{\eta}}{\sqrt{\eta!}} \right\rangle \geq \eta + \frac{1}{4 \left\langle \frac{\partial_{\alpha} \mathfrak{z}_{\alpha}^{\eta+1}}{\sqrt{\eta+1!}}, \frac{\partial_{\beta} \mathfrak{z}_{\beta}^{\eta+1}}{\sqrt{\eta+1!}} \right\rangle}$$

Siendo la integral de Yang – Mills, la siguiente:

$$\begin{aligned}
& \frac{1}{3} \int_{\mathfrak{U}}^{\infty} \mathfrak{e}^{\int_{\mathbb{C}}^{\infty} \Sigma_{j=1}^4 \mathfrak{U}_j \otimes \mathfrak{d}\mathfrak{x}^j \otimes \mathfrak{i}} \mathfrak{C}^{-|\mathfrak{d}\mathfrak{U}|^2/2} \mathfrak{D}\mathfrak{U}, \mathfrak{i} = \sqrt{-1} \\
& \quad \frac{1}{3} \mathfrak{e}^{-\frac{1}{2 \int_{\mathbb{R}^4}^{\infty} |\mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \mathfrak{d}\omega}} \mathcal{D}(\mathfrak{d}\mathfrak{U}) \\
& \quad \mathfrak{Z} = \int_{\{\mathfrak{d}\mathfrak{U}: \mathfrak{U} \in \mathfrak{S}_{\kappa}(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^4)_{\mathfrak{g}}\}}^{\infty} \mathfrak{e}^{-\frac{1}{2 \int_{\mathbb{R}^4}^{\infty} |\mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \mathfrak{d}\omega}} \mathcal{D}(\mathfrak{d}\mathfrak{U}) \\
& \quad \frac{1}{3} \mathfrak{e}^{-\frac{1}{2 \int_{\mathbb{C}^4}^{\infty} |\kappa \mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \mathfrak{d}\lambda_4}} \mathcal{D}(\mathfrak{d}\mathfrak{U}) \\
& \quad \mathfrak{Z} = \int_{\{\mathfrak{d}\mathfrak{U}: \mathfrak{U} \in \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^1(\mathbb{R}^4)_{\mathfrak{g}}\}}^{\infty} \mathfrak{e}^{-\frac{1}{2 \int_{\mathbb{C}^4}^{\infty} |\kappa \mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \mathfrak{d}\lambda_4}} \mathcal{D}(\mathfrak{d}\mathfrak{U})
\end{aligned}$$

$$\mathbb{H} = \{(\mathfrak{d}_0 \mathcal{H}^2(\mathbb{C}^4)) \otimes (* \Lambda^2(\mathbb{R}^4))\} \oplus \{\mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4)\} \subset \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4)$$

$$\frac{1}{3} \mathfrak{e}^{-\frac{1}{2 \int_{\mathbb{C}^4}^{\infty} |\kappa \mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \mathfrak{d}\lambda_4}} \mathcal{D}(\mathfrak{d}\mathfrak{U}) = \frac{\mathcal{Y}^{\kappa} \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}}{\int_{\mathbb{B} \otimes \mathfrak{g}}^{\infty} \mathcal{Y}^{\kappa} \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}} = \frac{\mathcal{Y}^{\kappa} \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}}{\mathbb{E}|\mathcal{Y}^{\kappa}|}$$

$$\mathbb{E}_{\mathcal{Y}^{\kappa}}^{\kappa}[\mathcal{R}]^{\mathfrak{R}} = 1/ \int_{\mathbb{B} \otimes \mathfrak{g}}^{\infty} \mathcal{Y}^{\kappa} \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta} \int_{\mathbb{B} \otimes \mathfrak{g}}^{\infty} \mathfrak{Z} \mathcal{Y}^{\kappa} \mathfrak{d}\hat{\mu}_{\kappa^2}^{\times \eta}$$

$$\begin{aligned}
& \int_{\mathbb{R}^4} |\mathfrak{d}\mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^4 \, \mathfrak{d}\omega \sum_{1 \leq i \leq j \leq 4} \int_{\mathbb{R}^4} \left(\sum_{\alpha=1}^{\eta} \alpha_{i;j,\alpha}^2 + \sum_{\gamma=1}^{\eta} \sum_{\substack{\alpha,\beta \\ \hat{\alpha},\hat{\beta}}} \alpha_{i,\alpha} \alpha_{j,\beta} \alpha_{i,\hat{\alpha}} \alpha_{j,\hat{\beta}} \mathfrak{C}_{\gamma}^{\alpha\beta} \mathfrak{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} \right. \\
& \quad \left. + 2 \sum_{\gamma=1}^{\eta} \sum_{\substack{\alpha,\beta \\ \hat{\alpha},\hat{\beta}}} \alpha_{i;j,\gamma} \alpha_{i,\alpha} \alpha_{j,\beta} \mathfrak{C}_{\gamma}^{\alpha\beta} \right) \mathfrak{d}\omega + \sum_{i;j=1}^4 \int_{\mathbb{R}^4} \sum_{\alpha=1}^{\eta} \alpha_{0;i,\alpha}^2 \, \mathfrak{d}\omega \\
& \exp \left[-1/2 \sum_{\alpha=1}^{\eta} \int_{\mathbb{R}^4} \mathfrak{d}\omega \sum_{1 \leq i \leq j \leq 4} \int_{\mathbb{R}^4} \alpha_{i;j,\alpha}^2 + \sum_{i;j=1}^4 \int_{\mathbb{R}^4} \alpha_{0;i,\alpha}^2 \right] \mathcal{D}(\mathfrak{d}\mathfrak{U})
\end{aligned}$$

Más, usando el teorema de Stokes, obtenemos:

$$\begin{aligned}
& \frac{1}{3} \int_{\mathfrak{U}} e^{\int_{\mathfrak{s}}^{\infty} \mathfrak{d}\mathfrak{U} \otimes i} \mathfrak{E}^{-|\mathfrak{d}\mathfrak{U}|^2/2} \mathfrak{D}\mathfrak{U} \\
& \frac{1}{3} \int_{\substack{\mathfrak{U} \in \mathcal{H}^2(\mathbb{C}^4) \otimes \Lambda^2(\mathbb{R}^4)}} e^{\mathfrak{I} \int_{\mathfrak{s}}^{\infty} \kappa \mathfrak{d}\mathfrak{U}} e^{-|\kappa \mathfrak{d}\mathfrak{U}|^2/2} \mathfrak{D}\mathfrak{U} \\
& \nu_s^{\kappa} = \int_{\mathfrak{I}^2} \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} \sum_{0 \leq \alpha \leq \beta \leq 4} \frac{\kappa^2}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2} \right) \\
& \Upsilon(\mathbb{R}^4, \kappa; \mathfrak{S}, i) = \mathbb{E}_{\mathfrak{U}\mathfrak{M}} \left(\exp \left(\frac{1}{\kappa} (\cdot, \nu_s^{\kappa}) \right) \right) = \int_{\mathfrak{U} \in \mathfrak{B}(\mathbb{R}^4; \partial)} \left(\exp \left(\frac{1}{\kappa} (\mathfrak{U}, \nu_s^{\kappa} \otimes i) \right) \right) \mathfrak{d}\tilde{\mu}(\mathfrak{U}) \\
& \mathbb{E}_{\mathfrak{U}\mathfrak{M}} = \left(\exp \left(\frac{1}{\kappa} \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{I}^2} \frac{\kappa^2}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) \left(\cdot, \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2} \right) \otimes i \right) \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} \right) \right) \\
& = \exp \left(-1/2 \left| \frac{\sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{I}^2} \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} \kappa}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2} \right) \right|_{\frac{2}{\partial, \kappa}}^2 \right) \\
& \rightarrow \exp \left(-1/8 \int_{\mathfrak{S}}^{\infty} \rho_{\mathfrak{S}} \right) \\
& \psi_{\omega} = \psi(\omega) = \frac{e^{-\frac{|\omega|^2}{2}}}{\sqrt{2\varpi}} = \langle \psi_{\omega} \chi_{\omega}, \psi_{\nu} \chi_{\nu} \rangle = \psi_{\omega} \psi_{\nu} e^{\omega \nu} = \frac{1}{2\varpi} e^{-|\omega - \nu|^2/2}
\end{aligned}$$

$$\begin{aligned}
& \langle \sum_{i=1}^4 \sum_{\tau} \sum_{\rho_{\tau}} \mathfrak{C}_{\rho_{\tau}, i^{3^{\rho_{\tau}}}} \otimes d\mathfrak{x}^i, \xi_{\alpha\beta}^{\kappa}(\omega) \rangle_{\partial, \kappa} = \kappa^2 \langle \partial \sum_{i=1}^4 \sum_{\tau} \sum_{\rho_{\tau}} \mathfrak{C}_{\rho_{\tau}, i^{3^{\rho_{\tau}}}} \otimes d\mathfrak{x}^i, \partial \xi_{\alpha\beta}^{\kappa}(\omega) \rangle \\
& = \kappa \sum_{\tau} \sum_{\rho_{\tau}} \psi(\omega) (\mathfrak{C}_{\rho_{\tau}, \beta} \partial_{\alpha} \omega^{\rho_{\tau}} - \mathfrak{C}_{\rho_{\tau}, \alpha} \partial_{\beta} \omega^{\rho_{\tau}}) \\
& \exp(-(\frac{1}{\kappa} \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{I}^2} d\mathfrak{s} d\mathfrak{t} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) \frac{\kappa^2}{4} \xi_{\alpha\beta}^{\kappa}(\frac{\kappa\sigma(\mathfrak{s}, \mathfrak{t})}{2}))_{\partial, \kappa/2}^2) \\
& = \exp(1/2 \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{I}^2, \mathfrak{I}^2} d\mathfrak{s} d\mathfrak{t} d\bar{\mathfrak{s}} d\bar{\mathfrak{t}} \frac{\kappa^2}{16} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\bar{\mathfrak{s}}, \bar{\mathfrak{t}}) \kappa^2 \langle \partial \xi_{\alpha\beta}^{\kappa}(\frac{\kappa\sigma(\mathfrak{s}, \mathfrak{t})}{2}), \partial \xi_{\alpha\beta}^{\kappa}(\frac{\kappa\sigma(\bar{\mathfrak{s}}, \bar{\mathfrak{t}})}{2}) \rangle) \\
& = \exp(-1/8(2\pi) \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{I}^2, \mathfrak{I}^2} d\mathfrak{s} d\mathfrak{t} d\bar{\mathfrak{s}} d\bar{\mathfrak{t}} \frac{\kappa^2}{4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\bar{\mathfrak{s}}, \bar{\mathfrak{t}}) e^{-\frac{\kappa^2 |\sigma(\mathfrak{s}, \mathfrak{t}) - \sigma(\bar{\mathfrak{s}}, \bar{\mathfrak{t}})|^2}{8}}) \\
& \rightarrow \exp(-1/8 \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathfrak{I}^2} d\mathfrak{s} d\mathfrak{t} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) \rho_{\mathfrak{S}}^{\alpha\beta}(\sigma(\mathfrak{s}, \mathfrak{t}))) = e^{-1/8 \int_{\mathfrak{S}}^{\infty} \rho_{\mathfrak{S}}}
\end{aligned}$$

Campo Cuántico no Abelian

En campos cuánticos no abelianos y bajo operadores holonómicos y por algebra de Lie, obtenemos:

$$\begin{aligned}
\mathcal{T} &= \int_{\mathfrak{A}} e^{\int_{\mathfrak{C}}^{\infty} \sum_{i=1}^4 \mathfrak{A}_i \otimes d\mathfrak{x}^i} e^{-1/2|\mathfrak{F}|^2} d\mathfrak{A} \\
\mathcal{T} \exp(\int_{\mathfrak{C}}^{\infty} \sum_{i=1}^4 \mathfrak{A}_i \otimes d\mathfrak{x}^i) &= \mathcal{T} \exp(\sum_i \int_{\mathfrak{C}_i}^{\infty} \sum_{j=1}^4 \mathfrak{A}_j \otimes d\mathfrak{x}^j) = \mathcal{T} \bigotimes i \exp(\int_{\mathfrak{C}_i}^{\infty} \sum_{j=1}^4 \mathfrak{A}_j \otimes d\mathfrak{x}^j) = \mathcal{T} \bigotimes i (\mathfrak{I} \\
&+ \sum_{\alpha} |F_{\alpha} \mathbb{E}^{\alpha}|(\mathfrak{S}_i) \Delta_i \mathfrak{S}^2 + \mathcal{O}(\Delta_i \mathfrak{S}^3)) \\
&= \mathcal{T} \bigotimes i \exp(\log(\mathfrak{I} \\
&+ \sum_{\alpha} |F_{\alpha} \mathbb{E}^{\alpha}|(\mathfrak{S}_i) \Delta_i \mathfrak{S}^2 + \mathcal{O}(\Delta_i \mathfrak{S}^3))) = \mathcal{T} \exp(\sum_i \sum_{\alpha} |F_{\alpha} \mathbb{E}^{\alpha}|(\mathfrak{S}_i) \Delta_i \mathfrak{S}^2 + \mathcal{O}(\Delta_i \mathfrak{S}^3)) \\
&\rightarrow \mathcal{T} \exp(\int_{\mathfrak{S}}^{\infty} \sum_{\alpha} F_{\alpha} \mathbb{E}^{\alpha}
\end{aligned}$$

Cuyo operador de translación en curvatura, es la que sigue:

$$\begin{aligned} & \mathfrak{U}_0^{1-} \mathfrak{U}_{1-}^1 \cdot \mathfrak{U}_{0+}^1 \cdot \overrightarrow{\prod_{i=2}^{\eta} \mathfrak{U}_0^{i-}} \cdot \mathfrak{U}_{1-}^{\eta} \cdot \overrightarrow{\prod_{i=2}^{\eta} \mathfrak{U}_{i-}^{\eta}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{U}_{\eta}^{i+}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{U}_{i+}^0} \\ &= \overrightarrow{\prod_{i=1}^{\eta} \mathfrak{U}_0^{i-}} \cdot \overrightarrow{\prod_{i=1}^{\eta} \mathfrak{U}_{i-}^{\eta}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{U}_{\eta}^{i+}} \cdot \overleftarrow{\prod_{i=0}^{\eta-1} \mathfrak{U}_{i+}^0} \end{aligned}$$

En el que, la curvatura, se expresa así:

$$\begin{aligned} & \mathfrak{U}_i^{(j+1)-} \mathfrak{U}_{(i+1)-}^{j+1-} \mathfrak{U}_{i+1}^{j+} \mathfrak{U}_{i+}^j \mathfrak{U}_i^j = \mathfrak{U}_i^j + \epsilon^2 \sum_{0 \leq \alpha \leq \beta \leq 4} \langle \mathfrak{Z}_{\alpha\beta}^{\sigma} | \Omega_{i,\alpha\beta}^j \rangle \left(\sigma^{\dagger} \left(\frac{i}{\eta}, \frac{j}{\eta} \right), \dot{\sigma} \left(\frac{i}{\eta}, \frac{j}{\eta} \right) \right) \mathfrak{U}_i^j + \mathcal{O}(\epsilon^4) \\ &= \mathfrak{U}_i^j + \epsilon^2 \mathfrak{U}_i^j \mathfrak{Z} \Omega_i^j + \mathcal{O}(\epsilon^4) \boxtimes_i^j \end{aligned}$$

$$\begin{aligned} & \mathfrak{U}_i^{(j+1)-} \mathfrak{U}_{(i+1)-}^{j+1-} \mathfrak{U}_{i+1}^{j+} \mathfrak{U}_{i+}^j \cdot \mathfrak{U}_{(i+1)-}^j \mathfrak{U}_{i+1}^j \overrightarrow{\prod_{k=i+1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^j + \mathcal{O}(\epsilon^4) \boxtimes_k^j)} \\ &= \mathfrak{U}_i^{(j+1)-} \mathfrak{U}_{(i+1)-}^{j+1-} \mathfrak{U}_{i+1}^{j+} \mathfrak{U}_{i+}^j \mathfrak{U}_i^j \overrightarrow{\prod_{k=i+1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^j + \mathcal{O}(\epsilon^4) \boxtimes_k^j)} \\ &= \mathfrak{U}_i^j (1 + \epsilon^2 \mathfrak{Z} \Omega_k^j + \mathcal{O}(\epsilon^4) \boxtimes_k^j) \end{aligned}$$

$$\begin{aligned} & \tilde{\mathfrak{U}}_1^{j+1} = \tilde{\mathfrak{U}}_1^{j+} \tilde{\mathfrak{U}}_1^{j+} \cdot \overrightarrow{\prod_{k=1}^{\eta-1} (1 + \mathfrak{Z} \Omega_k^{j+1} \epsilon^2 + \mathcal{O}(\epsilon^4) \boxtimes_k^j)} \\ &= \tilde{\mathfrak{U}}_1^{j+} \cdot \overleftarrow{\prod_{l=0}^{j-1} \tilde{\mathfrak{U}}_1^{l+}} \cdot \tilde{\mathfrak{U}}_{0+\mu_0}^0 \overleftarrow{\prod_{l=0}^j \tilde{\mathfrak{U}}_1^{l+}} \overrightarrow{\prod_{k=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \\ &\quad \cdot \overrightarrow{\prod_{k=1}^{\eta-1} (1 + \mathfrak{Z} \Omega_k^{l+1} \epsilon^2 + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \end{aligned}$$

$$\begin{aligned} &= \overleftarrow{\prod_{l=0}^j \tilde{\mathfrak{U}}_1^{l+}} \cdot \tilde{\mathfrak{U}}_{0+\mu_0}^0 \overleftarrow{\prod_{l=0}^{j+1} \tilde{\mathfrak{U}}_1^{l+}} \overrightarrow{\prod_{k=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \\ &= \overleftarrow{\prod_{l=0}^j \tilde{\mathfrak{U}}_1^{l+}} \cdot \overleftarrow{\prod_{l=0}^{j+1} \tilde{\mathfrak{U}}_1^{l+}} \overrightarrow{\prod_{k=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \cdot \overleftarrow{\prod_{k=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \\ &= \mu_0 \overrightarrow{\prod_{l=0}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \cdot \overleftarrow{\prod_{k=1}^{\eta-1} (1 + \epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l)} \\ &= \mu_0 \overrightarrow{\prod_{l=0}^{\eta-1} e^{\epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l}} \cdot \overleftarrow{\prod_{l=0}^{\eta-1} \prod_{k=1}^{\eta-1} e^{\epsilon^2 \mathfrak{Z} \Omega_k^l + \mathcal{O}(\epsilon^4) \boxtimes_k^l}} \end{aligned}$$

$$\rightarrow \mu_0 \cdot \tilde{\mathcal{T}} \exp \left(\int \sum_{\mathfrak{Z}^2}^{\infty} \sum_{0 \leq \alpha \leq \beta \leq 4} \mathfrak{Z}_{\alpha\beta}^{\sigma} (s, t) \Omega_{\beta}^{\alpha} (s, t) \mathrm{d}s \mathrm{d}t \right)$$



Por lo que, finalmente la integral de Yang – Mills, equivale a lo que sigue:

$$\begin{aligned}
& \exp -1/2 \left(\int_{\mathbb{C}^4} \mathfrak{d}\lambda_4 |\kappa \partial \mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}|^2 \right) \mathfrak{D}\mathfrak{U} \\
&= \exp(-1/2 \int_{\mathbb{C}^4} \mathfrak{d}\lambda_4 |\kappa \partial \mathfrak{U} + \mathfrak{U} \wedge \mathfrak{U}| + \langle \mathfrak{U} \wedge \mathfrak{U}, \kappa \partial \mathfrak{U} \\
&\quad + \langle \mathfrak{U} \wedge \mathfrak{U} \rangle^2 \rangle \exp(-1/2 \int_{\mathbb{C}^4} \mathfrak{d}\lambda_4 |\kappa \partial \mathfrak{U}|^2) \mathfrak{D}\mathfrak{U} \\
&\langle \kappa \partial \mathfrak{U}(\omega), \mathfrak{d}\mathfrak{x}^\alpha \wedge \mathfrak{d}\mathfrak{x}^\beta \rangle = \left(\mathfrak{U}, \tilde{\xi}_{\alpha\beta}^\kappa(\omega) \right) = \{ \mathfrak{U}_i \mathfrak{U}_j \}(\omega) = \mathfrak{U}_i(\omega) \mathfrak{U}_j(\omega) = \left(\mathfrak{U} \otimes \mathfrak{U}, \zeta_i(\omega) \otimes \mathfrak{U}, \zeta_j(\omega) \right) \\
&= \{ \mathfrak{U}_{i,\alpha} \mathfrak{U}_{j,\beta} \overline{\mathfrak{U}_{i,\alpha} \mathfrak{U}_{j,\beta}} \}(\omega) = \langle \mathfrak{U} \otimes \mathfrak{U} \otimes \overline{\mathfrak{U}} \otimes \overline{\mathfrak{U}}, \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\hat{\alpha},\omega} \otimes \chi_{j,\hat{\beta},\omega} \rangle \\
&\mathfrak{U} = \sum_{i=1}^4 \sum_{\alpha=1}^{\eta} \mathfrak{U}_{i,\alpha} \otimes \mathfrak{d}\mathfrak{x}^i \otimes \mathfrak{E}^\alpha, \tilde{\mathfrak{U}} = \sum_{i=1}^4 \sum_{\alpha=1}^{\eta} \widetilde{\mathfrak{U}}_{i,\alpha} \otimes \mathfrak{d}\mathfrak{x}^i \otimes \mathfrak{E}^\alpha \\
&\int_{\omega \in \mathbb{R}^4} \{ \mathfrak{U}_{i,\alpha} \mathfrak{U}_{j,\beta} \overline{\mathfrak{U}_{i,\alpha} \mathfrak{U}_{j,\beta}} \}(\omega) = \langle \mathfrak{U}^{\otimes 2} \otimes \widetilde{\mathfrak{U}}^{\otimes 2}, \int_{\omega \in \mathbb{R}^4} \mathfrak{d}\omega \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\hat{\alpha},\omega} \otimes \chi_{j,\hat{\beta},\omega} \rangle \\
&\left(\mathfrak{U}_{i_1,\alpha_1} \bigotimes \cdots \mathfrak{U}_{i_4,\alpha_4}, \widetilde{\omega}_\omega^{\otimes 4} \right) = \prod_{j=1}^4 \left(\mathfrak{U}_{i_j,\alpha_j}, \tilde{\pi}_{i_j,\alpha_j,\omega} \right) \\
&= \left(\mathfrak{U}_{i_1,\alpha_1} \bigotimes \cdots \mathfrak{U}_{i_3,\alpha_3}, \left(\tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathfrak{E}^{\alpha_1} \otimes \widetilde{\omega}_\omega^{\otimes 2} \right) \right) \\
&= \left(\mathfrak{U}_{i_1,\alpha_1}, \tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathfrak{E}^{\alpha_1} \right) \prod_{j=2}^4 \left(\mathfrak{U}_{i_j,\alpha_j}, \tilde{\pi}_{i_j,\alpha_j,\omega} \right) \\
&= \left(\mathfrak{U}_{i_1,\alpha_1} \bigotimes \cdots \mathfrak{U}_{i_4,\alpha_4} \widetilde{\omega}_\omega^{\otimes 4} \otimes \left(\tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathfrak{E}^{\alpha_4} \right) \right) \\
&= \left(\mathfrak{U}_{i_4,\alpha_4}, \tilde{\xi}_{\alpha\beta}^\kappa(\omega) \otimes \mathfrak{E}^{\alpha_4} \right) \prod_{j=1}^4 \left(\mathfrak{U}_{i_j,\alpha_j}, \tilde{\pi}_{i_j,\alpha_j,\omega} \right)
\end{aligned}$$

$$\begin{aligned}
& \int_{\mathbb{C}^4} \mathfrak{d}\lambda_4 |\mathfrak{U} \wedge \mathfrak{U}|^2 = \sum_{\gamma} \sum_{0 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \frac{1}{2\mathfrak{C}_{\gamma}^{\alpha\beta}} \mathfrak{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} \left(\left(\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{j,\beta} \otimes \mathfrak{U}_{i,\hat{\alpha}} \otimes \mathfrak{U}_{j,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) \widetilde{\omega}_{\omega}^{\otimes 4} \right) \right. \\
& \quad \left. + \left(\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{j,\beta} \otimes \mathfrak{U}_{i,\hat{\alpha}} \otimes \mathfrak{U}_{j,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) \widetilde{\omega}_{\omega}^{\otimes 4} \right) \right) \\
& \int_{\mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4 |\kappa \partial \mathfrak{U}, \mathfrak{U} \wedge \mathfrak{U}|^4 \\
& = \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{U}_{\kappa,\gamma} \otimes \mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{j,\beta}, \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{C}^{\gamma}) \otimes \widetilde{\omega}_{\omega}^{\otimes 4} \right) \\
& + \int_{\mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4 |\mathfrak{U} \wedge \mathfrak{U}, \kappa \partial \mathfrak{U}|^4 \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{j,\beta} \otimes \mathfrak{U}_{\kappa,\gamma}, \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{C}^{\gamma}) \otimes \widetilde{\omega}_{\omega}^{\otimes 4} \right) \\
& \mathfrak{Y}_{\mathfrak{C}}^{\kappa} = \left(\{\mathfrak{U}_{i,\alpha}\}_{i,\alpha} \right) \\
& = \mathfrak{I}_{\mathfrak{r}} \hat{\mathcal{T}} \exp\left(\frac{1}{\kappa} \cdot \frac{\kappa^2}{4}\right. \\
& \quad \cdot \rho \kappa \\
& \quad \left. / 4 \int_{\mathfrak{Z}^2}^{\infty} \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} \mu_{\mathfrak{s},\mathfrak{t}}^{-1} \left(\sum_{1 \leq i \leq j \leq 4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s},\mathfrak{t}) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\bar{\mathfrak{s}},\bar{\mathfrak{t}}) \sum_{\alpha} (\mathfrak{U}_{\alpha}, \xi_{\alpha\beta}^{\kappa} \left(\frac{\kappa\sigma(\mathfrak{s},\mathfrak{t})}{2} \right) \otimes \mathfrak{C}^{\alpha} \otimes \rho(\mathfrak{C}^{\alpha})) \right. \right. \\
& \quad \left. \left. + \sum_{1 \leq i \leq j \leq 4} |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s},\mathfrak{t}) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\bar{\mathfrak{s}},\bar{\mathfrak{t}}) \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{j,\beta} \otimes \mathfrak{U}_{\kappa,\gamma} \otimes \widetilde{\omega}_{\left(\frac{\kappa\sigma(\mathfrak{s},\mathfrak{t})}{2}\right)}^{\otimes 4} \otimes \rho(\mathfrak{C}^{\gamma}) \right) \mu_{\mathfrak{s},\mathfrak{t}} \right) \right)
\end{aligned}$$

$$\begin{aligned}
\mathfrak{Y}_{\mathfrak{E}}^{\kappa} &= \left(\{\mathfrak{U}_{i,\alpha}\}_{i,\alpha} \right) \\
&= \exp \left(-\frac{1}{2} \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{U}_{\kappa,\gamma} \otimes \mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{i,\beta} \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{E}^{\gamma}) \otimes \widetilde{\omega}_{\omega}^{\otimes 4} \right) \right. \\
&\quad - 1/2 \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \left(\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{i,\beta} \otimes \mathfrak{U}_{\kappa,\gamma} \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) \widetilde{\omega}_{\omega}^{\otimes 4} \otimes (\xi_{ij}^{\kappa}(\omega) \otimes \mathfrak{E}^{\gamma}) \right) \\
&\quad - 1/2 \sum_{\gamma} \sum_{\kappa=1}^4 \sum_{1 \leq i \leq j \leq 4} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} \mathfrak{C}_{\gamma}^{\alpha\beta} \mathfrak{C}_{\gamma}^{\hat{\alpha}\hat{\beta}} / 2 \left(\left(\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{i,\beta} \otimes \mathfrak{U}_{i,\hat{\alpha}} \otimes \mathfrak{U}_{i,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) \widetilde{\omega}_{\omega}^{\otimes 4} \right) \right. \\
&\quad \left. \left. + (\mathfrak{U}_{i,\alpha} \otimes \mathfrak{U}_{i,\beta} \otimes \mathfrak{U}_{i,\hat{\alpha}} \otimes \mathfrak{U}_{i,\hat{\beta}}, \int_{\omega \in \mathbb{C}^4}^{\infty} \mathfrak{d}\lambda_4(\omega) \widetilde{\omega}_{\omega}^{\otimes 4}) \right) \right) \\
&\quad \left| \left(\frac{\kappa}{4} \right) \int_{\mathfrak{Z}^2}^{\infty} \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} \sum_{1 \leq i \leq j \leq 4} |\mathfrak{Z}_{ij}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) |\mathfrak{Z}_{ij}^{\sigma}|(\bar{\mathfrak{s}}, \bar{\mathfrak{t}}) \sum_{\gamma} \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_{\gamma}^{\alpha\beta}| |\pi_i(\omega)|_{\partial, \kappa} \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\alpha} \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\beta} \right| \\
&\leq \left(\frac{\kappa}{4} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{Z}^2}^{\infty} \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} \sum_{1 \leq i \leq j \leq 4} |\mathfrak{Z}_{ij}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) |\mathfrak{Z}_{ij}^{\sigma}|(\bar{\mathfrak{s}}, \bar{\mathfrak{t}}) \sum_{\gamma} \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_{\gamma}^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\alpha} \mathcal{M}_{\left(\frac{i, \kappa \sigma(\mathfrak{s}, \mathfrak{t})}{2}\right)}^{\beta} \right| = F_{\eta} \\
&= 4 \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \sum_{1 \leq i \leq j \leq 4} \sum_{\rho, \mathfrak{q}=1}^{\eta} |\mathfrak{Z}_{ij}^{\sigma}| \left(\frac{\rho}{\eta}, \frac{\mathfrak{q}}{\eta} \right) / \eta^2 \sum_{\gamma} \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_{\gamma}^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa \sigma\left(\frac{\rho}{\eta}, \frac{\mathfrak{q}}{\eta}\right)}{2}\right)}^{\alpha} \mathcal{M}_{\left(\frac{i, \kappa \sigma\left(\frac{\rho}{\eta}, \frac{\mathfrak{q}}{\eta}\right)}{2}\right)}^{\beta} \right|
\end{aligned}$$

$$\begin{aligned}
& \left| \left(\frac{\kappa}{4} \right) \int_{\mathfrak{I}^2} \mathfrak{d} s \mathfrak{d} t \sum_{1 \leq \alpha \leq \beta \leq 4} |\mathfrak{I}_{\alpha\beta}^\sigma|(\mathfrak{s}, t) |\mathfrak{I}_{\alpha\beta}^\sigma|(\bar{\mathfrak{s}}, \bar{t}) \sum_{\gamma} \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{\alpha\beta}| |\pi_{\alpha\beta}(\omega)|_{\partial, \kappa} \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^j \right| \\
& \leq \left(\frac{\kappa}{4} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{I}^2} \mathfrak{d} s \mathfrak{d} t \sum_{1 \leq i \leq j \leq 4} |\mathfrak{I}_{\alpha\beta}^\sigma|(\mathfrak{s}, t) |\mathfrak{I}_{\alpha\beta}^\sigma|(\bar{\mathfrak{s}}, \bar{t}) \sum_{\gamma} \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{\alpha\beta}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^j \right| = F_\eta \\
& = 4 \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \sum_{1 \leq i \leq j \leq 4} \sum_{\rho, q=1}^{\eta} |\mathfrak{I}_{\alpha\beta}^\sigma| \left(\frac{\rho}{\eta}, \frac{q}{\eta} \right) / \eta^2 \sum_{\gamma} \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{ij}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma\left(\frac{\rho}{\eta}, \frac{q}{\eta}\right)}{2}\right)}^j \right|
\end{aligned}$$

Cuya aproximación riemmaniana equivale a:

$$\begin{aligned}
& \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{I}^2} \mathfrak{d} s \mathfrak{d} t \sum_{1 \leq i \leq j \leq 4} |\mathfrak{I}_{ij}^\sigma|(\mathfrak{s}, t) |\mathfrak{I}_{ij}^\sigma|(\bar{\mathfrak{s}}, \bar{t}) \sum_{\gamma} \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_\gamma^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\alpha \mathcal{M}_{\left(\frac{j, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\beta \right| \\
& = \left| \sum_{\substack{\alpha < \beta \\ \hat{\alpha} < \hat{\beta}}} |\mathfrak{C}_\gamma^{\alpha\beta}| \mathcal{M}_{\left(\frac{i, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\alpha \mathcal{M}_{\left(\frac{j, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\beta \right| \\
& \leq \mathfrak{N} \|\mathfrak{B}(\gamma)\| \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{i, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\alpha} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{j, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\alpha} / \sqrt{\sum_{\beta} \mathcal{M}_{\left(\frac{i, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\beta} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{j, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^\beta} \\
& \left(\frac{4}{\kappa} \right) \left(\frac{4}{2\pi} \right) \kappa^2 \int_{\mathfrak{I}^2} \mathfrak{d} s \mathfrak{d} t \sum_{1 \leq \alpha \leq \beta \leq 4} |\mathfrak{I}_{\alpha\beta}^\sigma|(\mathfrak{s}, t) |\mathfrak{I}_{\alpha\beta}^\sigma|(\bar{\mathfrak{s}}, \bar{t}) \sum_{\gamma} \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{ij}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^j \right| \\
& = \left| \sum_{\substack{i < j \\ \hat{i} < \hat{j}}} |\mathfrak{C}_\gamma^{ij}| \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^i \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^j \right| \\
& \leq \mathfrak{N} \|\mathfrak{B}(\gamma)\| \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^i} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^j} / \sqrt{\sum_{\beta} \mathcal{M}_{\left(\frac{\alpha, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^i} \sqrt{\sum_{\alpha} \mathcal{M}_{\left(\frac{\beta, \kappa\sigma(\mathfrak{s}, t)}{2}\right)}^j}
\end{aligned}$$

Ahora bien, para efectos de simular superficies temporales y espaciales respectivamente, en campos cuánticos, se expresa lo que sigue, empezando por la transformación de Fourier:

$$\begin{aligned}
& \left(\frac{1}{\sqrt{2\omega}}\right)^4 \int_{\delta_0}^{\infty} e^{-\iota(q^2 \mathfrak{x}^2 + q^4 \mathfrak{x}^4)} \mathfrak{F}_{\alpha}(\widehat{\mathcal{H}}(\rho), \widehat{\mathfrak{P}}(\rho), \mathfrak{x}^2, \mathfrak{x}^4) \mathfrak{d}\mathfrak{x}^2 \mathfrak{d}\mathfrak{x}^4 \bigotimes \rho(\mathfrak{E}^{\alpha}) \\
&= \widehat{\mathfrak{F}}_{\alpha}(\widehat{\mathcal{H}}(\rho), \widehat{\mathfrak{P}}(\rho), q^2, q^4) \bigotimes \rho(\mathfrak{E}^{\alpha}) \\
&\widehat{\mathfrak{F}}(\widehat{\mathcal{H}}(\rho) \widetilde{\mathfrak{f}}_0 + \widehat{\mathfrak{P}}(\rho) \widetilde{\mathfrak{f}}_1 + q^2 \widetilde{\mathfrak{f}}_2 + q^4 \widetilde{\mathfrak{f}}_4) \\
&= \frac{e^{-\iota(\alpha^0 \widehat{\mathcal{H}}(\rho) - \alpha^1 \widehat{\mathfrak{P}}(\rho))}}{2\pi} \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^{\infty} e^{-\iota(\widehat{\mathfrak{s}} q^2 + \widehat{\mathfrak{s}} q^4) \widehat{\mathfrak{f}}\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}(\widehat{\mathcal{H}}(\rho), \widehat{\mathfrak{P}}(\rho))(\widehat{\mathfrak{s}} \widehat{\mathfrak{f}}_2 + \widehat{\mathfrak{s}} \widehat{\mathfrak{f}}_4) \mathfrak{d}\widehat{\mathfrak{s}} \\
&\equiv e^{-\iota(\alpha^0 \widehat{\mathcal{H}}(\rho) - \alpha^1 \widehat{\mathfrak{P}}(\rho)) \widehat{\mathfrak{f}}}(\widehat{\mathcal{H}}(\rho) \widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{P}}(\rho) \widehat{\mathfrak{f}}_1 + q^2 \widehat{\mathfrak{f}}_0 + q^4 \widehat{\mathfrak{f}}_4) \bigotimes \rho(\mathfrak{E}^{\alpha}) \\
&\langle \phi^{\alpha, \eta}(\mathfrak{f}) 1, \phi^{\nu \beta, \eta}(\mathfrak{g}) 1 \rangle = \mathfrak{C}(\rho_{\eta}) \mathcal{T} \mathfrak{r}(-\mathfrak{F}^{\alpha} \mathfrak{F}^{\beta}) \int_{\delta_0}^{\infty} (\widehat{\mathfrak{f}}^{\{\epsilon_0 \epsilon_1\}} \overline{\widehat{\mathfrak{g}}^{\{\epsilon_0 \epsilon_1\}}}) (\widehat{\mathcal{H}}(\rho_{\eta}), \widehat{\mathfrak{P}}(\rho_{\eta}))(\widehat{\mathfrak{s}}) \mathfrak{d}\widehat{\mathfrak{s}} \\
&\phi^{\alpha, \eta}(\mathfrak{f}) 1 = \int_{\vec{\mathfrak{x}} \in \mathbb{R}^4}^{\infty} \mathfrak{d}\vec{\mathfrak{x}} \widehat{\mathfrak{f}}(\vec{\mathfrak{x}}) \phi^{\alpha, \eta}(\vec{\mathfrak{x}}) 1 = \mathfrak{U}(\vec{\alpha}, 1) \phi^{\alpha, \eta}(\vec{\mathfrak{x}}) \mathfrak{U}(\vec{\alpha}, 1)^{-1} = \phi^{\alpha, \eta}(\vec{\mathfrak{x}} + \vec{\alpha}) \\
&= \frac{1}{2\omega} e^{\iota(\mathfrak{x}^0 \widehat{\mathcal{H}}(\rho_{\eta}) - \mathfrak{x}^1 \widehat{\mathfrak{P}}(\rho_{\eta}))} \delta(\cdot - (\mathfrak{x}^2, \mathfrak{x}^4)) \bigotimes \frac{\rho(\mathcal{F}^{\alpha}) \kappa^2}{4(2\omega)} \exp \delta(\vec{\mathfrak{x}} - \vec{\mathfrak{y}})^4 \\
&\langle \frac{1}{(2\omega)^2} e^{\iota(\widehat{\mathcal{H}}(\mathfrak{x}^0 - \mathfrak{y}^0) - \widehat{\mathfrak{P}}(\mathfrak{x}^1 - \mathfrak{y}^1))} \rho_{\kappa}^{\mathfrak{x}^+} \bigotimes \rho(\mathcal{F}^{\alpha}), \frac{1}{(2\omega)^2} e^{\iota(\mathfrak{y}^0 \widehat{\mathcal{H}} - \mathfrak{y}^1 \widehat{\mathfrak{P}})} \rho_{\kappa}^{\mathfrak{y}^+} \bigotimes \rho(\mathcal{F}^{\beta}) \rangle \\
&= \frac{\kappa^2}{(2\omega)^4 e^{\iota(\widehat{\mathcal{H}}(\mathfrak{x}^0 - \mathfrak{y}^0) - \widehat{\mathfrak{P}}(\mathfrak{x}^1 - \mathfrak{y}^1))} \exp(-\kappa^2 \delta(\vec{\mathfrak{x}} - \vec{\mathfrak{y}})^4)} \cdot \langle \rho_{\eta}(\mathcal{F}^{\alpha}), \rho_{\eta}(\mathcal{F}^{\beta}) \rangle \\
&\frac{1}{(2\omega)^2} \int_{\vec{\mathfrak{x}}, \vec{\mathfrak{y}} \in \mathbb{R}^4}^{\infty} \widehat{\mathfrak{f}}(\vec{\mathfrak{x}}) \widehat{\mathfrak{g}}(\vec{\mathfrak{y}}) \langle e^{\iota(\mathfrak{x}^0 \widehat{\mathcal{H}} - \mathfrak{x}^1 \widehat{\mathfrak{P}})} \rho_{\kappa}^{\mathfrak{x}^+} \bigotimes \rho(\mathcal{F}^{\alpha}), e^{\iota(\mathfrak{y}^0 \widehat{\mathcal{H}} - \mathfrak{y}^1 \widehat{\mathfrak{P}})} \rho_{\kappa}^{\mathfrak{y}^+} \bigotimes \rho(\mathcal{F}^{\beta}) \rangle \mathfrak{d}\vec{\mathfrak{x}} \mathfrak{d}\vec{\mathfrak{y}} \\
&= \frac{\kappa^2}{4} \int_{\widehat{\mathfrak{s}}, \widehat{\mathfrak{t}} \in \mathbb{S}_0}^{\infty} \widehat{\mathfrak{f}}^{\{\epsilon_0 \epsilon_1\}}(\widehat{\mathfrak{s}}) \overline{\widehat{\mathfrak{g}}^{\{\epsilon_0 \epsilon_1\}}(\widehat{\mathfrak{t}})} \frac{1}{(2\omega)^2} \exp(-\kappa^2 |\widehat{\mathfrak{s}} - \widehat{\mathfrak{t}}|^4 \\
&/4) \mathfrak{d}\widehat{\mathfrak{s}} \mathfrak{d}\widehat{\mathfrak{t}} \langle \rho_{\eta}(\mathcal{F}^{\alpha}), \rho_{\eta}(\mathcal{F}^{\beta}) \rangle \bigotimes \rho(\mathcal{F}^{\alpha}), \bigotimes \rho(\mathcal{F}^{\beta}) \rangle
\end{aligned}$$

$$\begin{aligned}
& \lambda \left(\mathfrak{Z}, \mathfrak{F}_\alpha \bigotimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}_{\alpha=0}^4\} \right) + \mu \left(\mathfrak{Z}, \mathfrak{G}_\alpha \bigotimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}_{\alpha=0}^4\} \right) \\
&= \left(\mathfrak{Z} \cup \mathfrak{Z} (\lambda \widehat{\mathfrak{F}_\alpha} + \mu \widehat{\mathfrak{G}_\alpha}) \bigotimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}_{\alpha=0}^4\} \right) \\
&\langle (\mathfrak{Z}, \mathfrak{F}_\alpha \bigotimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}_{\alpha=0}^4\}), (\mathfrak{Z}, \mathfrak{G}_\beta \bigotimes \rho(\mathfrak{E}^\beta), \{\widehat{\mathfrak{G}_\beta}_{\alpha=0}^4\}) \rangle \\
&= \int_{\mathfrak{Z} \cup \mathfrak{Z}}^\infty \|\widehat{\mathfrak{F}_\alpha} \widehat{\mathfrak{G}_\beta}\| \cdot \mathfrak{d}|\rho| \cdot \mathrm{T}_r \left(-\rho(\mathfrak{E}^\alpha) \rho(\mathfrak{E}^\beta) \right) \\
&= \sum_{\alpha=1}^{\eta} \mathfrak{E}(\rho) \int_{\mathfrak{Z}^2}^\infty \|\widehat{\mathfrak{F}_\alpha} \cdot \widehat{\mathfrak{G}_\alpha}\| \left(\sigma(\mathfrak{Z}) \right) \left| \sum_{0 \leq \alpha \leq \beta \leq 4} \dot{\rho}_\sigma^{\alpha\beta}(\mathfrak{Z}) (\det \mathfrak{K}_\sigma^{\alpha\beta}(\mathfrak{Z})) \right| \mathfrak{d}\mathfrak{Z} \\
&\langle \mathfrak{U}(\vec{\alpha}, \Lambda) \left(\mathfrak{S}, \mathfrak{F}_\alpha \otimes \rho(\mathfrak{E}^\alpha), \{\widehat{\mathfrak{F}_\alpha}_{\alpha=0}^4\} \right), \mathfrak{U}(\vec{\alpha}, \Lambda) \left(\mathfrak{S}, \mathfrak{G}_\beta \otimes \rho(\mathfrak{E}^\beta), \{\widehat{\mathfrak{F}_\alpha}_{\alpha=0}^4\} \right) \rangle \\
&= \int_{(\Lambda \mathfrak{S} + \vec{\alpha}) \cap (\Lambda \mathfrak{S} + \vec{\alpha})}^\infty \mathfrak{d}|\rho| e^{-i(\vec{\alpha} \cdot (\mathcal{H}(\rho_\eta) \Lambda \widehat{\mathfrak{F}_0} + \mathfrak{P}(\rho_\eta) \Lambda \widehat{\mathfrak{F}_1}))} e^{i(\vec{\alpha} \cdot (\mathcal{H}(\rho_\eta) \Lambda \widehat{\mathfrak{F}_0} + \mathfrak{P}(\rho_\eta) \Lambda \widehat{\mathfrak{F}_1}))} \\
&\times \|\widehat{\mathfrak{F}_\alpha} \widehat{\mathfrak{G}_\beta}\| (\Lambda^{-1}(\cdot - \vec{\alpha})) \cdot \mathrm{T}_r \left(-\rho(\mathfrak{E}^\alpha) \rho(\mathfrak{E}^\beta) \right) \\
&= \int_{\Lambda(\mathfrak{Z} \cup \mathfrak{Z}) + \vec{\alpha}}^\infty \|\widehat{\mathfrak{F}_\alpha} \widehat{\mathfrak{G}_\beta}\| (\Lambda^{-1}(\cdot - \vec{\alpha})) \cdot \mathfrak{d}|\rho| \cdot \mathrm{T}_r \left(-\rho(\mathfrak{E}^\alpha) \rho(\mathfrak{E}^\beta) \right) \\
&= \int_{\mathfrak{Z} \cup \mathfrak{Z}}^\infty \|\widehat{\mathfrak{F}_\alpha} \widehat{\mathfrak{G}_\beta}\|(\cdot) \cdot \mathfrak{d}|\rho| \cdot \mathrm{T}_r \left(-\rho(\mathfrak{E}^\alpha) \rho(\mathfrak{E}^\beta) \right)
\end{aligned}$$

Siendo los operadores relativos a los campos cuánticos, los que siguen:

$$\begin{aligned}
\mathcal{D}^{\vec{\kappa}} &= \left(\frac{\partial}{\partial \mathfrak{x}^0} \right)^{\kappa^0} \left(\frac{\partial}{\partial \mathfrak{x}^1} \right)^{\kappa^1} \left(\frac{\partial}{\partial \mathfrak{x}^2} \right)^{\kappa^2} \left(\frac{\partial}{\partial \mathfrak{x}^3} \right)^{\kappa^3}, \vec{\chi}^{\vec{\kappa}} = (\partial \mathfrak{x}^0)^{\kappa^0} (\partial \mathfrak{x}^1)^{\kappa^1} (\partial \mathfrak{x}^2)^{\kappa^2} (\partial \mathfrak{x}^3)^{\kappa^3} \\
\phi^{\alpha, \eta}(\vec{\mathfrak{f}}) 1 &= \left(\mathfrak{S}_0, \widehat{\mathfrak{f}^{\{\widehat{\mathfrak{E}_0}, \widehat{\mathfrak{E}_1}\}}} \otimes \rho_\eta(\mathfrak{F}^\alpha), \{\epsilon_\alpha\}_{\alpha=0}^4 \right) \equiv \left(\mathfrak{S}_0, \widehat{\mathfrak{f}^{\{\widehat{\mathfrak{E}_0}, \widehat{\mathfrak{E}_1}\}}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{P}}(\rho_\eta)) \otimes \rho_\eta(\mathfrak{F}^\alpha), \{\epsilon_\alpha\}_{\alpha=0}^4 \right) \\
&\in \mathfrak{H}(\rho_\eta)
\end{aligned}$$

En dimensión $\mathbb{R}^4$ tenemos lo que sigue:

$$\begin{aligned}
\vec{\mathfrak{x}} \in \mathbb{R}^4 &\rightarrow \vec{\mathfrak{f}}^{\{\widehat{\mathfrak{f}_0}, \widehat{\mathfrak{f}_1}\}} \left(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{P}}(\rho_\eta) \right) (\vec{\mathfrak{x}}) = \int_{\mathfrak{S}^b}^\infty e^{-i\vec{\mathfrak{m}}(\cdot) / 2\pi \vec{\mathfrak{f}}(\vec{\mathfrak{x}} + \star)} \mathfrak{d}|\rho| \\
&= \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^\infty e^{-i(\sigma(\widehat{\mathfrak{s}}) \cdot (\widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{f}_0} + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{f}_1}))} / 2\pi \vec{\mathfrak{f}}(\vec{\mathfrak{x}} + \sigma(\widehat{\mathfrak{s}})) \cdot |\widehat{\rho}_\sigma|(\widehat{\mathfrak{s}}) \mathfrak{d}\widehat{\mathfrak{s}}
\end{aligned}$$



$$\left\{ \alpha_0 1 + \sum_{\eta, \mu=1}^{\infty} (\mathfrak{S}_{\eta, \mu}, \mathfrak{f}_{\eta, \alpha}^{\mu} \otimes \rho_{\eta}(\mathfrak{E}^{\alpha}), [\widehat{\mathfrak{f}_{\alpha}^{\eta, \mu}}]_{\alpha=0}^4) \boxtimes \alpha_0 \in \mathfrak{C}, \mathfrak{f}_{\eta, \alpha}^{\mu} \in \mathfrak{P}_{\mathfrak{S}_{\eta, \mu}}, \mathfrak{S}_{\eta, \mu} \in \mathfrak{L} \right\}$$

Cuya parametrización va como se indica:

$$\int_{\mathfrak{Z}^2} |\mathfrak{f} \circ \sigma|^2(\mathfrak{S}) \left| \sum_{0 \leq \alpha \leq \beta \leq 4} \dot{\rho}_{\sigma}^{\alpha \beta}(\mathfrak{Z})(\det \mathfrak{K}_{\alpha \beta}^{\sigma}(\mathfrak{Z})) \right| \mathfrak{d} \mathfrak{Z} < \infty$$

$$\phi^{\alpha, \eta}(\mathfrak{f}) \sum_{\mu=1}^{\infty} (\mathfrak{S}_{\mu}, \mathfrak{G}_{\beta}^{\mu} \otimes \rho_{\mathfrak{m}}(\mathfrak{E}^{\beta}), \{\widehat{\mathfrak{f}_{\alpha}^{\mu}}\}_{\alpha=0}^4) \\ = \left\{ \sum_{\mu=1}^{\infty} (\mathfrak{S}_{\mu}, \mathfrak{f}_{\eta}^{\{\mathfrak{f}_0^{\mu}, \mathfrak{f}_1^{\mu}\}} \mathfrak{A}(\Lambda^{\mu})_{\gamma}^{\alpha} \cdot \mathfrak{G}_{\beta}^{\mu} \otimes \rho_{\eta} \|\mathfrak{F}^{\gamma}, \mathfrak{E}^{\beta}\|, \{\widehat{\mathfrak{f}_{\alpha}^{\mu}}\}_{\alpha=0}^4),_{\mathfrak{m} \neq \eta}^{\mathfrak{m} = \eta} \right. \\ \left. \sum \left| \frac{\partial \alpha}{\partial \beta} \right|^{i/\hbar} \left| \frac{\partial \zeta}{\partial \eta} \right|^{i/\hbar} \left| \frac{\frac{\partial \kappa}{\partial \lambda}}{\frac{\partial \mu}{\partial \nu}} \right|^{\sqrt{\frac{i}{\hbar}}} \left| \frac{\partial}{\partial \xi} \right| \left| \frac{\partial}{\partial t} \right| * \left\| \partial o \otimes \partial \rho \otimes \partial \varrho \otimes \partial \sigma \otimes \partial \varsigma \otimes \frac{\partial \tau}{\partial \upsilon} \otimes \partial \varphi \otimes \partial \phi + \partial \psi + \partial \Psi - \frac{\partial \Delta}{\partial \omega} / \Lambda_{\nu \mu}^{\mu \nu} \cdot \partial \Omega \Phi \right\|^{i/\hbar} \right\}$$

$$\phi^{\alpha, \eta}(\mathfrak{f}) \sum_{\mathfrak{m}=0}^{\infty} \nu_{\mathfrak{m}} = \sum_{\mathfrak{m}=0}^{\infty} \phi^{\alpha, \eta}(\mathfrak{f}) \nu_{\mathfrak{m}} = \alpha_0 \left(\delta_0, \mathfrak{f}_{\eta}^{\{\epsilon_0, \epsilon_1\}} \otimes \rho_{\eta}(\mathfrak{E}^{\alpha}), \{\epsilon_{\alpha}\}_{\alpha=0}^4 \right) + \phi^{\alpha, \eta}(\mathfrak{f}) \nu_{\eta}$$

$$\phi^{\beta, \mathfrak{m}}(\mathfrak{f}) \sum_{\eta=0}^{\infty} \nu_{\eta} = \sum_{\eta=0}^{\infty} \phi^{\beta, \mathfrak{m}}(\mathfrak{f}) \nu_{\eta} = \beta_0 \left(\delta_0, \mathfrak{f}_{\mathfrak{m}}^{\{\epsilon_0, \epsilon_1\}} \otimes \rho_{\mathfrak{m}}(\mathfrak{E}^{\beta}), \{\epsilon_{\beta}\}_{\beta=0}^4 \right) + \phi^{\beta, \mathfrak{m}}(\mathfrak{f}) \nu_{\mathfrak{m}}$$

$$\begin{aligned} \phi^{\alpha, \eta}(\mathfrak{g})^* \left(\mathfrak{S}, \mathfrak{f}_{\beta} \otimes \rho_{\eta}(\mathfrak{E}^{\beta}), \{\widehat{\mathfrak{f}_{\alpha}}\}_{\beta=0}^4 \right) \\ = - \left(\mathfrak{S}, \mathfrak{g}^{\{\widehat{\mathfrak{f}_0^{\alpha \beta}}, \widehat{\mathfrak{f}_1^{\alpha \beta}}\}} \mathfrak{A}(\Lambda)_{\gamma}^{\alpha} \cdot \mathfrak{f}_{\beta} \otimes \rho_{\eta} \langle \mathfrak{F}^{\gamma}, \mathfrak{E}^{\beta} \rangle, \{\widehat{\mathfrak{f}_{\alpha}}\}_{\alpha=0}^4 \right) \\ + \left\langle \left(\mathfrak{S}, \mathfrak{f}_{\beta} \otimes \rho_{\eta}(\mathfrak{E}^{\beta}), \{\widehat{\mathfrak{f}_{\alpha}}\}_{\alpha=0}^4 \right), \phi^{\alpha, \eta}(\mathfrak{g}) 1 \right\rangle 1 \end{aligned}$$

En el que, la ciclicidad del campo cuántico, se expresa así:

$$\begin{aligned} \|\mathfrak{f} - \tilde{\mathfrak{g}}_{\epsilon}\|_{\mathcal{L}^4} &= \left(\int_{\mathfrak{Z}^4} |1 - \mathfrak{g}_{\delta}^{\eta-1}|^4(\mathfrak{f}) |\mathfrak{f}|^4(\mathfrak{f}) \mathfrak{d} \mathfrak{f} \right)^{\frac{1}{2}} \leq \mathcal{M} \|1 - \mathfrak{g}_{\delta}^{\eta-1}\|_{\mathcal{L}^4} < \epsilon \\ \mathfrak{E}^{\gamma} &= \sum_{\beta=1}^{\eta(\gamma)} \mathfrak{d}_{\mathfrak{M}, \beta}^{\gamma} ad \left(\mathfrak{F}^{\alpha^{\gamma, \beta}}_1 \right) \cdots ad \left(\mathfrak{F}^{\alpha^{\gamma, \beta}}_{\mathfrak{M}-1} \right) \mathfrak{F}^{\alpha^{\gamma, \beta}}_{\mathfrak{M}} \end{aligned}$$

Cuya función de Schwartz, en $\mathbb{R}^4$ es igual a:



$$\sum_{\beta=1}^{\eta(\gamma)} \mathfrak{d}_{\mathfrak{M},\beta}^{\gamma} \phi^{\alpha_{\eta}^{\gamma,\beta}}(\mathfrak{F}_{\eta}) \cdots \phi^{\alpha_{\mathfrak{M}}^{\gamma,\beta}}(\mathfrak{F}_{\mathfrak{M}}) = (\mathfrak{S}, \prod_{i=1}^{\mathfrak{M}} \mathfrak{f}_i \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\beta=0}^4$$

$$\sum_{\beta=1}^{\eta(\gamma)} \mathfrak{d}_{\hat{\eta}}(\gamma, \xi) \phi^{\alpha_1(\gamma, \xi), \eta}(\mathfrak{G}_1) \cdots \phi^{\alpha_{\hat{\eta}-1}(\gamma, \xi), \eta}(\mathfrak{G}_{\hat{\eta}-1}) \phi^{\alpha_{\hat{\eta}}(\gamma, \xi), \eta}(\mathfrak{G}_{\hat{\eta}}) 1 = \left(\mathfrak{S}, \prod_{i=1}^{\hat{\eta}} \mathfrak{g}_i \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\alpha=0}^4 \right)$$

Lo anterior, computacionalmente equivale a:

$$\left| (\mathfrak{S}, \mathfrak{f} \otimes \rho(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\alpha=0}^4) - \left(\mathfrak{S}, \prod_{i=1}^{\hat{\eta}} \mathfrak{g}_i \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\epsilon_{\alpha}\}_{\alpha=0}^4 \right) \right| < \epsilon$$

Cuya métrica de Minkowski, se define así:

$$\begin{aligned} \mathcal{T}(\mathfrak{f}) &= \langle \phi^{\alpha, \eta}(\mathfrak{f}) \left(\mathfrak{S}, \hat{\mathfrak{g}}_{\gamma} \otimes \rho_{\eta}(\mathfrak{E}^{\gamma}), \{\mathfrak{f}_{\alpha}\}_{\alpha=0}^4 \right), \left(\mathfrak{S}, \mathfrak{g}_{\beta} \otimes \rho_{\eta}(\mathfrak{E}^{\beta}), \{\mathfrak{f}_{\alpha}\}_{\alpha=0}^4 \right) \rangle \\ &= \mathfrak{E}_{\alpha}^{\gamma, \beta} \int_{\mathfrak{S}^4}^{\infty} \mathfrak{d}\hat{\mathfrak{t}} \left(\mathfrak{f}^{\{\mathfrak{f}_0, \mathfrak{f}_1\}} \times \hat{\mathfrak{g}}_{\gamma} \cdot \overline{\mathfrak{g}_{\beta}} \right) \left(\sigma(\hat{\mathfrak{t}}) \right) \cdot |\rho'_{\sigma}|(\hat{\mathfrak{t}}) \end{aligned}$$

$$\begin{aligned} \vec{\chi} \rightarrow \mathfrak{f}^{\{\mathfrak{f}_0, \mathfrak{f}_1\}}(\vec{\chi}) &\equiv \mathfrak{f}^{\{\mathfrak{f}_0, \mathfrak{f}_1\}} \left(\mathcal{H}(\rho_{\eta}), \mathfrak{P}(\rho_{\eta}) \right) (\vec{\chi}) \\ &= \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^{\infty} \frac{\mathbf{e}^{-i(\sigma(\mathfrak{s}) \cdot (\mathcal{H}(\rho_{\eta}) \mathfrak{f}_0 + \mathfrak{P}(\rho_{\eta}) \mathfrak{f}_1))}}{2\pi} \mathfrak{f}(\vec{\chi} + \hat{\sigma}(\mathfrak{S})) |\rho'_{\hat{\sigma}}|(\mathfrak{S}) \mathfrak{d}\mathfrak{s} \end{aligned}$$

$$\begin{aligned} \mathcal{T}(\mathfrak{f}) &= \int_{\mathfrak{S}^4}^{\infty} \mathfrak{f}^{\{\mathfrak{f}_0, \mathfrak{f}_1\}} \left(\sigma(\hat{\mathfrak{t}}) \right) \hbar(\hat{\mathfrak{t}}) \mathfrak{d}\hat{\mathfrak{t}} \\ &= \int_{\hat{\mathfrak{s}} \in \mathbb{R}^4, \mathfrak{t} \in \mathfrak{S}^4}^{\infty} \mathfrak{d}\hat{\mathfrak{t}} \mathfrak{d}\hat{\mathfrak{s}} \mathfrak{f} \left(\sigma(\hat{\mathfrak{t}}) + \hat{\sigma}(\mathfrak{S}) |\rho'_{\hat{\sigma}}|(\mathfrak{S}) |\rho'_{\sigma}|(\hat{\mathfrak{t}}) \cdot \frac{\mathbf{e}^{-i(\hat{\sigma}(\mathfrak{S}) \cdot \vec{\alpha})}}{2\pi} |\hat{\mathfrak{g}}_{\gamma} \cdot \overline{\mathfrak{g}_{\beta}}| \circ \sigma(\hat{\mathfrak{t}}) \cdot \mathfrak{E}_{\alpha}^{\gamma, \beta} \right) \end{aligned}$$

En este punto, cabe aplicar la ley de transformación del operador cuántico, que se expresa así:

$$\begin{aligned}
& \mathfrak{U}(\vec{\alpha}, \Lambda) \phi^{\alpha, \eta}(\mathfrak{f}) \mathfrak{U}(\vec{\alpha}, \Lambda)^{-1} \left(\mathfrak{S}, \mathfrak{g}_\beta \otimes \rho_\eta(\mathfrak{E}^\beta), \{\hat{\mathfrak{f}}_\alpha\}_{\alpha=0}^4 \right) \\
&= \mathfrak{U}(\vec{\alpha}, \Lambda) \phi^{\alpha, \eta}(\mathfrak{f}) \left(\Lambda^{-1}(\mathfrak{S} - \vec{\alpha}), \mathcal{T}(\rho_\eta, \vec{\alpha})^{-1} \mathfrak{g}_\beta(\Lambda \cdot + \vec{\alpha}) \otimes \rho_\eta(\mathfrak{E}^\beta), \mathcal{D} \right) \\
&= \mathfrak{U}(\vec{\alpha}, \Lambda) \left(\Lambda^{-1}(\mathfrak{S} - \vec{\alpha}), \left(\mathcal{T}(\rho_\eta, \vec{\alpha})^{-1} \mathfrak{f}^\mathfrak{C} \right) (\cdot) \mathfrak{d}_\gamma^\alpha \right. \\
&\quad \cdot \mathfrak{g}_\beta(\Lambda \cdot + \vec{\alpha}) \otimes \text{ad} \left(\rho_\eta(\mathfrak{F}^\gamma) \right) \rho_\eta(\mathfrak{E}^\beta), \mathcal{D} \Big) \\
&= (\mathfrak{S}, \mathfrak{T}(\rho_\eta, \vec{\alpha}) \mathcal{T}(\rho_\eta, \vec{\alpha})^{-1} \mathfrak{f}^\mathfrak{C}(\Lambda^{-1}(\cdot - \vec{\alpha})) \mathfrak{d}_\gamma^\alpha \mathfrak{g}_\beta(\cdot) \otimes \text{ad} \left(\rho_\eta(\mathfrak{F}^\gamma) \right) \rho_\eta(\mathfrak{E}^\beta), \{\hat{\mathfrak{f}}_\alpha\}_{\alpha=0}^4, \mathcal{D}) \\
&\mathfrak{f}^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{P}}(\rho_\eta))(\Lambda^{-1}(\vec{\chi} - \vec{\alpha})) = \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^\infty \frac{e^{-i(\widehat{\sigma}(\widehat{\mathfrak{s}}) \cdot \widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{g}}_0 + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{g}}_1))}}{2\pi} \mathfrak{f}(\vec{\mathfrak{Y}} + \widehat{\sigma}(\widehat{\mathfrak{s}}) |\rho_{\widehat{\sigma}}| ((\widehat{\mathfrak{S}}) \mathfrak{d} \widehat{\mathfrak{s}})) \\
&= \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^\infty \frac{e^{-i(\sigma(\widehat{\mathfrak{s}}) \cdot \widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{f}(\vec{\mathfrak{Y}} + \Lambda^{-1} \sigma(\widehat{\mathfrak{s}})) |\rho_\sigma| ((\widehat{\mathfrak{S}}) \mathfrak{d} \widehat{\mathfrak{s}})) \\
&= \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^\infty \frac{e^{-i(\sigma(\widehat{\mathfrak{s}}) \cdot \widehat{\mathcal{H}}(\rho_\eta) \widehat{\mathfrak{f}}_0 + \widehat{\mathfrak{P}}(\rho_\eta) \widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{f}(\Lambda^{-1}(\vec{\mathfrak{Y}} + \sigma(\widehat{\mathfrak{s}}) - \vec{\alpha})) |\rho_\sigma| ((\widehat{\mathfrak{S}}) \mathfrak{d} \widehat{\mathfrak{s}})) \\
&= \mathfrak{f}(\Lambda^{-1}(\cdot - \vec{\alpha}))^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}}(\widehat{\mathcal{H}}(\rho_\eta), \widehat{\mathfrak{P}}(\rho_\eta))(\vec{\mathfrak{Y}}) \\
&\left(\mathfrak{S}, \mathfrak{f}(\Lambda^{-1}(\cdot - \vec{\alpha}))^{\{\hat{\mathfrak{f}}_0, \hat{\mathfrak{f}}_1\}} \mathfrak{A}(\Lambda^{-1})^\alpha \mathfrak{A}(\widehat{\Lambda})^\gamma \cdot \mathfrak{g}_\beta \otimes \text{ad} \left(\rho_\eta(\mathfrak{F}^\delta) \right) \rho_\eta(\mathfrak{E}^\beta), \{\widehat{\Lambda} \epsilon_\alpha\}_{\alpha=0}^4 \right) \\
&= \mathfrak{A}(\Lambda^{-1})^\alpha \phi^{\gamma, \eta} \left(\mathfrak{f}(\Lambda^{-1}(\cdot - \vec{\alpha})) \right) \left(\mathfrak{S}, \mathfrak{g}_\beta \otimes \rho_\eta(\mathfrak{E}^\beta), \{\widehat{\Lambda} \epsilon_\alpha\}_{\alpha=0}^4 \right)
\end{aligned}$$

Cuya simetría CPT, es igual a:

$$\begin{aligned}
& [\phi^{\alpha, \eta}(\mathfrak{f}), \phi^{\beta, \eta}(\mathfrak{g})^*]_{\pm} \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \\
&= -\mathfrak{A}(\Lambda)_\delta^\alpha \overline{\mathfrak{A}(\Lambda)_\mu^\beta} \left(\mathfrak{S}, \mathfrak{B}^\pm (\mathfrak{f}^\mathfrak{C} \cdot \widehat{\mathfrak{g}}^\mathfrak{C} \pm \mathfrak{g}^\mathfrak{C} \cdot \widehat{\mathfrak{f}}^\mathfrak{C}) \right. \\
&\quad \cdot \hbar_\gamma \otimes \text{ad}(\rho_\eta(\mathfrak{F}^\delta)) \text{ad}(\rho_\eta(\mathfrak{F}^{\mu\nu})) (\rho_\eta(\mathfrak{E}^\gamma)) \\
&\quad + \langle (\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)), \phi^{\beta, \eta}(\mathfrak{G})^\circ 1 \rangle \phi^{\alpha, \eta}(\mathfrak{F})^* 1 \\
&\quad \pm \langle (\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma)), \phi^{\beta, \eta}(\mathfrak{F})^* 1 \rangle \phi^{\alpha, \eta}(\mathfrak{G})^\circ 1
\end{aligned}$$

$$\begin{aligned}
& \phi^{\alpha,\eta}(\mathfrak{F})^* \phi^{\beta,\eta}(\mathfrak{G})^* \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \\
&= \phi^{\alpha,\eta}(\mathfrak{F})^* \overline{\mathfrak{A}(\Lambda)_\mu^\beta}(\mathfrak{S}, -\widehat{\mathfrak{g}}^\mathfrak{C} \cdot \hbar_\gamma \otimes \text{ad}(\rho_\eta(\mathfrak{F}^{\mu\nu}))(\rho_\eta(\mathfrak{E}^\gamma)) \\
&+ \langle \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta,\eta}(\mathfrak{G})^\circ 1 \rangle 1) \\
&= -\mathfrak{A}(\Lambda)_\delta^\alpha \overline{\mathfrak{A}(\Lambda)_\mu^\beta}(\mathfrak{S}, \mathfrak{f}^\mathfrak{C} \cdot \widehat{\mathfrak{g}}^\mathfrak{C} \cdot \hbar_\gamma \otimes \text{ad}(\rho_\eta(\mathfrak{F}^\delta)) \text{ad}(\rho_\eta(\mathfrak{F}^{\mu\nu}))(\rho_\eta(\mathfrak{E}^\gamma)) \\
&+ \langle \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta,\eta}(\mathfrak{G})^\circ 1 \rangle \phi^{\alpha,\eta}(\mathfrak{F})^* 1
\end{aligned}$$

$$\mathfrak{g}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}(\mathcal{H}, \mathfrak{P})(\vec{\mathfrak{x}}) = \int_{\widehat{\mathfrak{s}} \in \mathbb{R}^4}^{\infty} \frac{e^{-i(\widehat{\mathfrak{g}}(\widehat{\mathfrak{s}}) \cdot (\mathcal{H}\widehat{\mathfrak{f}}_0 + \mathfrak{P}\widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{g}(\vec{\mathfrak{x}} + \vec{\mathfrak{y}}(\widehat{\mathfrak{s}})) |\rho_{\widehat{\mathfrak{g}}}(\widehat{\mathfrak{s}})| \mathfrak{d}\widehat{\mathfrak{s}}$$

$$\overline{\mathfrak{f}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}}(\mathcal{H}, \mathfrak{P})(\vec{\mathfrak{x}}) = \int_{\widehat{\mathfrak{t}} \in \mathbb{R}^4}^{\infty} \frac{e^{-i(\widehat{\mathfrak{g}}(\widehat{\mathfrak{t}}) \cdot (\mathcal{H}\widehat{\mathfrak{f}}_0 + \mathfrak{P}\widehat{\mathfrak{f}}_1))}}{2\pi} \mathfrak{g}(\vec{\mathfrak{x}} + \vec{\mathfrak{y}}(\widehat{\mathfrak{t}})) |\rho_{\widehat{\mathfrak{g}}}(\widehat{\mathfrak{t}})| \mathfrak{d}\widehat{\mathfrak{t}}$$

$$\begin{aligned}
\left[\mathfrak{g}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}} \cdot \overline{\mathfrak{f}^{\{\widehat{\mathfrak{f}}_0, \widehat{\mathfrak{f}}_1\}}} \right] (\mathcal{H}, \mathfrak{P})(\vec{\mathfrak{x}}) &= \int_{\widehat{\mathfrak{s}}, \widehat{\mathfrak{t}} \in \mathbb{R}^4}^{\infty} \frac{e^{-i(\widehat{\mathfrak{g}}(\widehat{\mathfrak{s}}) \cdot (\mathcal{H}\widehat{\mathfrak{f}}_0 + \mathfrak{P}\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} \mathfrak{g}\vec{\mathfrak{x}} \left(\vec{\mathfrak{y}}(\widehat{\mathfrak{s}}) + \vec{\mathfrak{y}}(\widehat{\mathfrak{t}}) \right) \widehat{\mathfrak{f}}_{\vec{\mathfrak{x}}} \left(\vec{\mathfrak{y}}(\widehat{\mathfrak{t}}) \right) |\rho_{\widehat{\mathfrak{g}}}(\widehat{\mathfrak{s}})| |\rho_{\widehat{\mathfrak{g}}}(\widehat{\mathfrak{t}})| \mathfrak{d}\widehat{\mathfrak{s}} \mathfrak{d}\widehat{\mathfrak{t}} \\
&= \int_{\widehat{\mathfrak{t}} \in \mathbb{R}^4, \widehat{\mathfrak{s}} \in \mathcal{D}}^{\infty} \frac{e^{-i(\widehat{\mathfrak{g}}(\widehat{\mathfrak{s}}) \cdot (\mathcal{H}\widehat{\mathfrak{f}}_0 + \mathfrak{P}\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} \mathfrak{g}\vec{\mathfrak{x}} \left(\vec{\mathfrak{y}}(\widehat{\mathfrak{s}}) + \vec{\mathfrak{y}}(\widehat{\mathfrak{t}}) \right) \widehat{\mathfrak{f}}_{\vec{\mathfrak{x}}} \left(\vec{\mathfrak{y}}(\widehat{\mathfrak{t}}) \right) |\rho_{\widehat{\mathfrak{g}}}(\widehat{\mathfrak{s}})| |\rho_{\widehat{\mathfrak{g}}}(\widehat{\mathfrak{t}})| \mathfrak{d}\widehat{\mathfrak{s}} \mathfrak{d}\widehat{\mathfrak{t}} \\
&\left[-\frac{e^{-i((\mu-\nu) \cdot (\mathcal{H}(\rho_\eta)\widehat{\mathfrak{f}}_0 + \mathfrak{P}(\rho_\eta)\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} + \frac{e^{-i((\nu-\mu) \cdot (\mathcal{H}(\rho_\eta)\widehat{\mathfrak{f}}_0 + \mathfrak{P}(\rho_\eta)\widehat{\mathfrak{f}}_1))}}{(2\pi)^2} \right] \widehat{\mathfrak{f}}_{\vec{\mathfrak{x}}}(\mu) \overline{\mathfrak{g}}_{\vec{\mathfrak{x}}}(\mu)
\end{aligned}$$

En un mapa bilineal, tenemos:

$$\begin{aligned}
(\mathfrak{f}, \mathfrak{g}) \in \mathcal{P} \times \mathfrak{P} &\rightarrow \langle \phi^{\alpha,\eta}(\mathfrak{f}) \phi^{\beta,\eta}(\mathfrak{g})^* \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \left(\mathfrak{S}, \widehat{\hbar_\gamma} \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \rangle \\
&- \langle \phi^{\alpha,\eta}(\mathfrak{f})^* \phi^{\beta,\eta}(\mathfrak{g}) \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta,\eta}(\mathfrak{g}) 1 \rangle \langle \phi^{\alpha,\eta}(\mathfrak{f}) 1 \left(\mathfrak{S}, \widehat{\hbar_\gamma} \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \rangle
\end{aligned}$$

$$\begin{aligned}
& \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \widehat{\mathcal{W}}(\vec{x}, \vec{y}) \widehat{f}(\vec{x}) \bigotimes_{\mathbb{R}} g(\vec{y}) d\vec{x} d\vec{y} \\
&= \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \Re \mathfrak{E} \widehat{\mathcal{W}}(\vec{x}, \vec{y}) [\widehat{f}(\vec{x}) \widehat{g}(\vec{y}) + \overline{\widehat{f}(\vec{x})} \overline{\widehat{g}(\vec{y})}] d\vec{x} d\vec{y} \\
&+ \Im \int_{\vec{x} \in \mathbb{R}^4} \int_{\vec{y} \in \mathbb{R}^4} \Im \mathfrak{M} \widehat{\mathcal{W}}(\vec{x}, \vec{y}) [\widehat{f}(\vec{x}) \widehat{g}(\vec{y}) - \overline{\widehat{f}(\vec{x})} \overline{\widehat{g}(\vec{y})}] d\vec{x} d\vec{y} \\
&= \langle \phi^{\alpha, \eta}(\widehat{f}) \phi^{\beta, \eta}(g)^* \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \left(\mathfrak{S}, \widehat{\hbar_\gamma} \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \rangle \\
&- \langle \phi^{\alpha, \eta}(\widehat{f})^* \phi^{\beta, \eta}(g) \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right), \phi^{\beta, \eta}(g) 1 \rangle \langle \phi^{\alpha, \eta}(\widehat{f}) 1 \left(\mathfrak{S}, \widehat{\hbar_\gamma} \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \rangle \\
&\phi^{\alpha, \eta}(\widehat{f}) \phi^{\beta, \eta}(g)^* \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) - \langle \left(\mathfrak{S}, \hbar_\gamma \otimes \rho_\eta(\mathfrak{E}^\gamma) \right) \phi^{\beta, \eta}(g) 1 \rangle \phi^{\alpha, \eta}(\widehat{f}) 1 \\
&= -\mathfrak{A}(\Lambda)_\delta^\alpha \overline{\mathfrak{A}(\Lambda)_\mu^\beta}(\mathfrak{S}, \left(\widehat{f}_{\{0,1\}} \cdot \overline{g_{\{0,1\}}} \right) \cdot \hbar_\gamma \otimes ad(\rho_\eta(\mathfrak{F}^\delta)) ad(\rho_\eta(\mathfrak{F}^{\mu\nu}))(\rho_\eta(\mathfrak{E}^\gamma))) \\
&\mathfrak{A}^\pm(\vec{\mathfrak{z}}) \int_{\widehat{\mathfrak{s}\vec{\mathfrak{t}} \in \mathbb{R}^4}} \frac{e^{-i(\widehat{\mathfrak{y}}(\mathfrak{s}) - \widehat{\mathfrak{y}}(\vec{\mathfrak{t}})) \cdot (\widehat{\mathcal{H}} \widehat{f}_0 + \widehat{\mathfrak{P}} \widehat{f}_1)}}{(2\pi)^2} \vec{f}(\vec{\mathfrak{y}}(\mathfrak{s}) \vec{g}(\vec{\mathfrak{t}})) |\rho_{\widehat{\mathfrak{y}}}(\mathfrak{s})| |\rho_{\widehat{\mathfrak{y}}}(\vec{\mathfrak{t}})| d\mathfrak{s} d\vec{\mathfrak{t}} \\
&\pm \int_{\widehat{\mathfrak{s}\vec{\mathfrak{t}} \in \mathbb{R}^4}} \frac{e^{-i(\widehat{\mathfrak{y}}(\mathfrak{s}) - \widehat{\mathfrak{y}}(\vec{\mathfrak{t}})) \cdot (\widehat{\mathcal{H}} \widehat{f}_0 + \widehat{\mathfrak{P}} \widehat{f}_1)}}{(2\pi)^2} \vec{g}(\vec{\mathfrak{y}}(\vec{\mathfrak{t}}) \vec{f}(\vec{\mathfrak{y}}(\mathfrak{s}))) |\rho_{\widehat{\mathfrak{y}}}(\mathfrak{s})| |\rho_{\widehat{\mathfrak{y}}}(\vec{\mathfrak{t}})| d\mathfrak{s} d\vec{\mathfrak{t}}
\end{aligned}$$

Aplicando en este punto, la función beta de Callan-Symanzik, tenemos:

$$\beta(\mathfrak{C}) = \partial \mathfrak{C} / \partial |\Im \mathfrak{M} \widehat{\mathfrak{N}}_\eta|$$

En tanto que, del operador cuántico hamiltoniano, se obtiene:

$$\mathcal{T}\mathfrak{r}(\rho(\mathfrak{E}^\alpha)) \left(\rho(\mathfrak{E}^\beta) \right) = \mathfrak{C}(\rho) \mathcal{T}\mathfrak{r}(\mathfrak{E}^\alpha \mathfrak{E}^\beta)$$

$$\begin{aligned}
\mathcal{E}(\rho) &= \sum_{\alpha=1}^{\mathfrak{N}} (\rho(\mathfrak{E}^\alpha)) \left(\rho(\mathfrak{E}^\beta) \right) \\
&= 1/\kappa \sum_{\alpha=1}^{\mathfrak{N}} \kappa^2/4 \int_{\widehat{\mathfrak{s}} \in (-\delta, 1+\delta)^4} d\widehat{\mathfrak{s}} \sum_{i,j=1}^4 |\Im_{0ij}^\sigma| (d\widehat{\mathfrak{s}}) \kappa(\psi \cdot d_0 \mathfrak{A}_{ij,\alpha}) \left(\frac{\kappa \sigma(\widehat{\mathfrak{s}})}{2} \right) \rho(\mathfrak{E}^\alpha) \\
&\mathfrak{C} \sum_{\alpha=1}^{\eta} \int_{\widehat{\mathfrak{s}} \in \Im_\delta^4} d\widehat{\mathfrak{s}} \sum_{i,j=1}^4 |\Im_{0ij}^\sigma|(\widehat{\mathfrak{s}}) (d_0 \mathfrak{A}_{ij,\alpha})(\sigma(\widehat{\mathfrak{s}})) \rho(\mathfrak{E}^\alpha) \\
\langle v_{\mathcal{R}(\alpha,\tau)}^{\kappa,\rho} \rangle^2 &= -\mathbb{E} \left(\cdot, v_{\mathcal{R}(\alpha,\tau)}^{\kappa,\rho} \right)^2 y^\kappa
\end{aligned}$$

$$\mathfrak{A}^\rho = \sum_{\alpha=1}^{\eta} \sum_{i,j=1}^4 \alpha_{i,\alpha} \otimes \mathfrak{d}\mathfrak{x}^i \otimes \rho(\mathfrak{E}^\alpha) \in \mathfrak{S}_{\mathfrak{R}}(\mathbb{R}^4) \otimes \Lambda^1(\mathbb{R}^4) \otimes \rho(\mathfrak{g})$$

$$\frac{1}{3} = \int_{\{\mathfrak{d}\mathfrak{A} \in \mathfrak{S}_{\mathfrak{R}}(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \rho(\mathfrak{g})\}} \exp(\mathfrak{C} \int_{\mathcal{R}(\alpha)} \mathfrak{d}\{\mathfrak{A}^\rho\}) \mathfrak{e}^{-1/2 \mathfrak{S}_{\mathfrak{YM}}(\mathfrak{A})} \mathfrak{D}|\mathfrak{d}\mathfrak{A}| = \mathbb{E}_{\mathcal{Y}\mathcal{M}}^\kappa(\exp(\left(\cdot, \nu_{\mathcal{R}(\alpha)}^{\kappa,\rho}\right)))$$

$$\mathfrak{Z} = \int_{\{\mathfrak{d}\mathfrak{A} \in \mathfrak{S}_{\mathfrak{R}}(\mathbb{R}^4) \otimes \Lambda^2(\mathbb{R}^4) \otimes \rho(\mathfrak{g})\}} \mathfrak{e}^{-1/2 \mathfrak{S}_{\mathfrak{YM}}(\mathfrak{A})} \mathfrak{D}|\mathfrak{d}\mathfrak{A}|$$

$$-\mathbb{E}\left(\cdot, \nu_{\mathcal{R}(\alpha)}^{\kappa,\rho_\eta}\right)\left(\cdot, \nu_{\mathcal{R}_\delta(\alpha)}^{\kappa,\rho_\eta}\right)\mathcal{Y}^\kappa = \frac{|\alpha|}{4\otimes \mathfrak{E}(\rho_\eta)} - \epsilon(\eta,\kappa)$$

$$\overline{\mathfrak{C}}/\kappa^4\mathcal{C}(\rho_\eta)\leq \mathcal{T}\mathfrak{r}\in(\eta,\kappa)\leq \overline{\mathfrak{C}}/\kappa^4\mathcal{C}(\rho_\eta)$$

Más, aplicando la ecuación de Callan-Symanzik, tenemos:

$$-\frac{1}{\mathfrak{E}(\rho_\eta)\mathbb{E}\left(\cdot, \nu_{\mathcal{R}(\alpha)}^{\kappa,\rho_\eta}\right)\left(\cdot, \nu_{\mathcal{R}_\delta(\alpha)}^{\kappa,\rho_\eta}\right)\mathcal{Y}^\kappa} = \frac{\frac{|\alpha|}{4\otimes \mathfrak{E}_4(\rho_\eta)}}{\mathfrak{E}(\rho_\eta)\mathbb{I}_\eta} - 1/\mathcal{C}(\rho_\eta)\epsilon(\eta,\kappa)$$

$$\frac{\widehat{\mathcal{N}}_\eta \mathfrak{E}_4(\rho_\eta)}{\mathfrak{E}(\rho_\eta)} - \frac{1}{\mathfrak{E}(\rho_\eta) \mathcal{T}\mathfrak{r}} \in (\eta,\kappa) \equiv \mathbb{N} - 1/\mathfrak{E}(\rho_\eta) \mathcal{T}\mathfrak{r} \in (\eta,\kappa)$$

$$\left\{e\frac{\partial}{\partial e}+\beta(\mathfrak{C})\frac{\partial}{\partial \mathfrak{C}}+2\gamma(\mathfrak{C})\right\}\mathfrak{G}_\eta^{(4)}(\mathfrak{C},\epsilon)=0$$

$$\mathfrak{G}_\eta^{(4)}(\mathfrak{C},\epsilon)=\frac{\widehat{\mathcal{N}}_\eta}{\epsilon}-c^4\hat{\lambda}+\mathfrak{f}(c^5)=\frac{\partial}{\partial e}\mathfrak{G}_\eta^{(4)}(\mathfrak{C},\epsilon)=-\frac{\widehat{\mathcal{N}}_\eta}{\epsilon},\frac{\partial}{\partial \mathfrak{C}}\mathfrak{G}_\eta^{(4)}(\mathfrak{C},\epsilon)=-4c^3\hat{\lambda}+\hat{\mathfrak{f}}(c^4)$$

$$-\frac{\widehat{\mathcal{N}}_\eta}{\epsilon}-4\beta(\mathfrak{C})c^3\hat{\lambda}+2\gamma(\mathfrak{C})\mathfrak{G}_\eta^{(4)}(\mathfrak{C},\epsilon)+\beta(\mathfrak{C})\hat{\mathfrak{f}}(c^4)=0$$

$$-\frac{\mathfrak{C}}{4}\hat{\mathfrak{f}}(c^4)+\mathfrak{f}(c^5)-4c^3\lambda(\mathfrak{C})\hat{\lambda}+\lambda(\mathfrak{C})\hat{\mathfrak{f}}(c^4)=0$$

$$\frac{1}{c^4|\hat{\mathfrak{f}}(c^4)|}+\frac{1}{c^5}|\hat{\mathfrak{f}}(c^5)|\leq \mathbb{C}^4$$

$$\hat{\lambda}(\mathfrak{C})=\frac{1}{4c^3\hat{\lambda}}+\hat{\mathfrak{f}}(c^4)\left(\frac{\mathfrak{C}}{4\hat{\mathfrak{f}}(c^4)}-\hat{\mathfrak{f}}(c^5)\right)$$



Por tanto, la brecha de masa se vuelve positiva y por ende, superior a cero (estado de vacío), cuando:

$$\begin{aligned}\langle \mathfrak{C}_4(\rho)v, v \rangle &= \sum_{\alpha=1}^N \langle \rho(\mathfrak{C}^\alpha)v, \rho(\mathfrak{C}^\alpha)v \rangle \geq \sum_{\alpha=1}^{\mathcal{L}} \left| \sum_{\beta=1}^{\mathcal{L}} \alpha_{\alpha,\beta} \lambda_\rho(\mathcal{H}_\beta) \right|_4^4 \\ &= \sum_{\alpha=1}^{\mathcal{L}} \sum_{\beta=1}^{\mathcal{L}} \sum_{\gamma=1}^{\mathcal{L}} \lambda_\rho(\mathcal{H}_\beta) \alpha_{\alpha,\beta} \alpha_{\alpha,\gamma} \lambda_\rho(\mathcal{H}_\gamma) \geq \mathfrak{C} |\lambda_\rho|_4^4\end{aligned}$$

$$\widehat{\mathcal{H}}(\rho_\eta)^2 = \frac{\widehat{\mathfrak{N}}_\eta}{4} \mathfrak{C}_2(\rho_\eta) = \frac{\mathfrak{N}}{4} \mathfrak{C}(\rho_\eta) = 0$$

$$\frac{\partial c}{\partial(\mathbb{I}_\eta \widehat{\mathfrak{N}})} = -\frac{\mathfrak{C}}{4} + \lambda(\mathfrak{C}), |\lambda(\mathfrak{C})| \leq \widehat{\mathfrak{C}}_4 \widetilde{\mathfrak{C}}_2 = \frac{\mathfrak{d}\mathfrak{C}}{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})} = -\frac{\mathfrak{C}}{4} + \lambda(\mathfrak{C}) \Rightarrow \frac{\mathfrak{d}\mathfrak{C}}{\mathfrak{C}} - 4\lambda(\mathfrak{C}) = -\frac{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})}{4}$$

$$= \frac{1}{c} \frac{\mathfrak{d}\mathfrak{C}}{1 + \mu(\mathfrak{C})} = -\frac{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})}{4} \Rightarrow \left(\frac{1}{c \sum_{\kappa=0}^{\infty} (-1)^\kappa \mu(\mathfrak{C})^\kappa} \right) \mathfrak{d}\mathfrak{C} = \frac{\mathfrak{d}(\mathbb{I}_\eta \widehat{\mathfrak{N}})}{4}$$

$$\mathcal{T}\mathfrak{r}\mathbb{E}\left(\left(\cdot, v_{\mathcal{R}(\alpha)}^{\kappa, \rho_\eta}\right)\left(\cdot, v_{\mathcal{R}_\delta(\alpha)}^{\kappa, \rho_\eta}\right)y^\kappa\right) = \frac{\widehat{\mathfrak{N}}_\eta}{4} \mathfrak{C}_2(\rho_\eta) - \mathcal{T}\mathfrak{r}\epsilon(\eta, \kappa)$$

$$= \frac{4}{\mathfrak{N}\mathfrak{C}(\rho_\eta)} \mathcal{T}\mathfrak{r}\mathbb{E}\left(-\left(\cdot, v_{\mathcal{R}(\alpha)}^{\kappa, \rho_\eta}\right)\left(\cdot, v_{\mathcal{R}_\delta(\alpha)}^{\kappa, \rho_\eta}\right)y^\kappa\right) - 1 = -\frac{4\mathcal{T}\mathfrak{r}\epsilon(\eta, \kappa)}{\mathfrak{N}\mathfrak{C}(\rho_\eta)}$$

$$\begin{aligned}\mathcal{U}(\vec{\alpha}, 1) \left(\mathfrak{S}, \mathfrak{f}_\alpha \otimes \rho(\mathfrak{C}^\alpha), \{\mathfrak{f}_\alpha\}_{\alpha=0}^4 \right) &= \mathbf{e}^{-\mathfrak{i}(\mathfrak{f}_0 \cdot \widehat{\mathcal{H}}(\vec{\alpha}, \rho) + \mathfrak{f}_1 \cdot \widehat{\mathfrak{P}}(\vec{\alpha}, \rho))} \left(\mathfrak{S} + \vec{\alpha}, \mathfrak{f}_\alpha(\cdot - \vec{\alpha}) \otimes \rho(\mathfrak{C}^\alpha), \{\mathfrak{f}_\alpha\}_{\alpha=0}^4 \right) \\ &= \mathbf{e}^{\mathfrak{i}(\alpha_0 \widehat{\mathcal{H}}(\rho) - \alpha^1 \widehat{\mathfrak{P}}(\rho))} \left(\mathfrak{S} + \vec{\alpha}, \mathfrak{f}_\alpha(\cdot - \vec{\alpha}) \otimes \rho(\mathfrak{C}^\alpha), \{\mathfrak{f}_\alpha\}_{\alpha=0}^4 \right)\end{aligned}$$

Cuyas particiones corresponden a:

$$\begin{aligned}\int_{\mathfrak{Q}}^{\infty} \{\hbar_\theta\}_{\theta=1}^{\mathfrak{r}+\mathfrak{s}} &= \prod_{\mathfrak{l}=1}^{\eta(\mathfrak{Q})} \left\{ \frac{\int_{\mathfrak{S}_0}^{\infty} (\prod_{\theta \in \mathfrak{U}_\mathfrak{l}} \int_{\gamma_\theta^- \in \mathbb{R}^4}^{\infty} \mathbf{e}^{\mathfrak{i}\chi(\theta)(\gamma_\theta^0 \widehat{\mathcal{H}} - \gamma_\theta^1 \widehat{\mathfrak{P}})})}{2\varpi} \hbar_\theta(\gamma_\theta^-, \gamma^+) \mathfrak{d}\gamma_\theta^- \mathfrak{d}\gamma^+ \right\} \\ &= \prod_{\mathfrak{l}=1}^{\eta(\mathfrak{Q})} \left\{ \frac{\int_{\mathfrak{S}_0}^{\infty} (\prod_{\theta \in \mathfrak{U}_\mathfrak{l}} \int_{\gamma_\theta^0 \gamma_\theta^1 \in \mathbb{R}^4}^{\infty} \mathbf{e}^{\mathfrak{i}\chi(\theta)(\gamma_\theta^0 \widehat{\mathcal{H}} - \gamma_\theta^1 \widehat{\mathfrak{P}})})}{2\varpi} \hbar_\theta(\gamma_\theta^0, \gamma_\theta^1, \gamma^2, \gamma^4) \mathfrak{d}\gamma_\theta^0 \mathfrak{d}\gamma_\theta^1 \mathfrak{d}\gamma^2 \mathfrak{d}\gamma^4 \right\}\end{aligned}$$

$$\begin{aligned}
& \langle \mathfrak{U}_\tau^\eta \mathfrak{B}_\delta^\eta 1, 4 \rangle - \langle \mathfrak{U}_\tau^\eta 1, 4 \rangle \langle \mathfrak{B}_\delta^\eta 1, 4 \rangle \equiv \langle \mathfrak{U}_\tau^\eta \mathfrak{P}_0 \mathfrak{B}_\delta^\eta 1, 4 \rangle \\
&= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathfrak{W}^\eta \left((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta} \right) \otimes_{\theta=\tau+1}^{r+\delta} \rho_\tau(\vec{x}_\tau) \cdot \prod_{\theta=\tau+1}^{r+\delta} d\vec{x}_\tau \\
&= \mathfrak{C}_\mathcal{R} \int_{\mathfrak{S}_0} \left(\prod_{\tau=1}^r \int_{\mathfrak{x}_\tau^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\tau^-) \hbar_\tau(\mathfrak{x}_\tau^-, \mathfrak{x}^+) d\mathfrak{x}_\tau^- \cdot \prod_{\theta=\tau+1}^{r+\delta} \int_{\mathfrak{x}_\theta^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\theta^-) \hbar_\theta(\mathfrak{x}_\theta^-, \mathfrak{x}^+) d\mathfrak{x}_\theta^- \right) d\mathfrak{x}^+ \\
&+ \sum_{\substack{\mathbb{Q} \neq \mathfrak{R} \\ \mathbb{Q} \in \Gamma}} \mathfrak{C}_\Omega \int_{\Omega} \{\mathfrak{H}_\theta\}_{\theta=1}^{r+\delta}
\end{aligned}$$

$$\begin{aligned}
\langle \mathfrak{P}_0 \mathfrak{U}(\vec{\alpha}, 1) \mathfrak{B}_\delta^\eta 1 \rangle &= \mathfrak{P}_0 \psi^{\beta_1, \eta}(\mathfrak{g}_1)_{\mathfrak{U}(\vec{\alpha})} \psi^{\beta_2, \eta}(\mathfrak{g}_2)_{\mathfrak{U}(\vec{\alpha})} \dots \psi^{\beta_{\delta-1}, \eta}(\mathfrak{g}_{\delta-1})_{\mathfrak{U}(\vec{\alpha})} \mathfrak{U}(\vec{\alpha}, 1) \psi^{\beta_\delta, \eta}(\mathfrak{g}_\delta) 1 \\
&= e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} \mathfrak{P}_0 \mathfrak{B}_\delta^{\eta, \vec{\alpha}} 1 = e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} \left(\mathfrak{B}_\delta^{\eta, \vec{\alpha}} 1 - \langle \mathfrak{B}_\delta^\eta 1, 1 \rangle 1 \right)
\end{aligned}$$

$$\langle \mathfrak{U}_\tau^\eta \mathfrak{P}_0 \mathfrak{U}(\vec{\alpha}, 1) \mathfrak{B}_\delta^\eta 1, 1 \rangle = e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} \left(\langle \mathfrak{U}_\tau^\eta \mathfrak{B}_\delta^{\eta, \vec{\alpha}} 1, 1 \rangle - \langle \mathfrak{U}_\tau^\eta 1, 1 \rangle \langle \mathfrak{B}_\delta^\eta 1, 1 \rangle \right)$$

$$\begin{aligned}
e^{i(\alpha^0 \hat{\mathcal{H}} - \alpha^1 \hat{\rho})} &= \int_{\mathbb{R}^4} \frac{e^{i(s\hat{\mathcal{H}} - t\hat{\rho})}}{2\omega} f(s, t, \mathfrak{x}^2, \mathfrak{x}^4) d\mathfrak{s} d\mathfrak{t} = \int_{\mathbb{R}^4} \frac{e^{i((s+\alpha^0)\hat{\mathcal{H}} - (t+\alpha^1)\hat{\rho})}}{2\omega} f(s, t, \mathfrak{x}^2, \mathfrak{x}^4) d\mathfrak{s} d\mathfrak{t} \\
&= \int_{\mathbb{R}^4} \frac{e^{i(s\hat{\mathcal{H}} - t\hat{\rho})}}{2\omega} f(s - \alpha^0, t - \alpha^1, \mathfrak{x}^2, \mathfrak{x}^4) d\mathfrak{s} d\mathfrak{t} \\
&= f(\cdot - (\alpha^0, \alpha^1, 0, 0))^{\{\epsilon_0 \epsilon_1\}} (\hat{\mathcal{H}}, \hat{\rho})(0, 0, \mathfrak{x}^2, \mathfrak{x}^4)
\end{aligned}$$

$$\begin{aligned}
\mathcal{H}^\eta(\vec{\alpha}) &= \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathfrak{W}^\eta \left((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta} \right) \otimes_{\tau=1}^r \rho_\tau(\vec{x}_\tau) \cdot \otimes_{\theta=\tau+1}^{r+\delta} \rho_0(\vec{x}_{\tau+\theta} - \vec{\alpha}) \cdot \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau \\
&\equiv \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \mathfrak{W}^\eta \left((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta} \right) \otimes_{\tau=1}^{r+\delta} \rho_\tau(\vec{x}_\tau) \cdot \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau \\
&= \mathfrak{C}_\mathcal{R} \int_{\mathfrak{S}_0} \left(\prod_{\tau=1}^r \int_{\mathfrak{x}_\tau^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\tau^-) \hbar_\tau(\mathfrak{x}_\tau^-, \mathfrak{x}^+) d\mathfrak{x}_\tau^- \cdot \prod_{\theta=\tau+1}^{r+\delta} \int_{\mathfrak{x}_\theta^- \in \mathbb{R}^4} \mathbb{E}(\mathfrak{x}_\theta^-) \hbar_\theta(\mathfrak{x}_\theta^-, \mathfrak{x}^+) d\mathfrak{x}_\theta^- \right) d\mathfrak{x}^+ \\
&+ \sum_{\substack{\mathbb{Q} \neq \mathfrak{R} \\ \mathbb{Q} \in \Gamma}} \mathfrak{C}_\Omega \int_{\Omega} \{\mathfrak{H}_\theta\}_{\theta=1}^{r+\delta} \\
&= \mathcal{H}_4^\eta(\vec{\alpha}) = \int_{\mathbb{R}^4 \times \dots \times \mathbb{R}^4} \prod_{\tau=1}^{r+\delta} d\vec{x}_\tau \mathfrak{W}^\eta \left((\vec{x}_\tau)_{\tau=1}^r, (\vec{x}_\theta + \vec{\alpha})_{\theta=\tau+1}^{r+\delta} \right) \varphi_1((\vec{x}_\tau)_{\tau=1}^r) \varphi_2((\vec{x}_\theta)_{\theta=\tau+1}^{r+\delta})
\end{aligned}$$

$$\mathfrak{W}^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}, (\mathfrak{x}_{\theta}^+ + \vec{\alpha})^r_{\theta=\tau+1}) = e^{-i\vec{\alpha} \cdot \mathfrak{C}_{\tau} m_{\eta}^{\eta}} \prod_{\theta=1}^{\tau+\delta} \mathfrak{C}(\mathfrak{x}_{\theta}^-) \cdot \mathfrak{W}_0^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}, (\mathfrak{x}_{\theta}^+ + \alpha^+)^r_{\theta=\tau+1})$$

$$\mathfrak{W}_0^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}, (\mathfrak{x}_{\theta}^+ + \vec{\alpha})^r_{\theta=\tau+1}) = \mathfrak{W}^{\eta}((0^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1}, (0^-, \mathfrak{x}_{\theta}^+ + \alpha^+)^r_{\theta=\tau+1})$$

$$\begin{aligned} \hat{\varphi}_1((q_{\tau}^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1}) &= \frac{\int_{\mathbb{R}^{4\tau}} \prod_{\tau=1}^r e^{i\chi(\tau) \mathfrak{x}_{\tau}^- \cdot q_{\tau}^-}}{2\varpi} \cdot \varphi_1((\mathfrak{x}_{\tau}^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1}) \prod_{\tau=1}^{r+\delta} d\mathfrak{x}_{\tau}^-, \hat{\varphi}_2((q_{\theta}^-, \mathfrak{x}_{\theta}^+)^r_{\theta=\tau+1}) \\ &= \frac{\int_{\mathbb{R}^{4\tau}} \prod_{\tau=1}^r e^{i\chi(\theta) \mathfrak{x}_{\theta}^- \cdot q_{\theta}^-}}{2\varpi} \cdot \varphi_2((\mathfrak{x}_{\theta}^-, \mathfrak{x}_{\theta}^+)^r_{\theta=\tau+1}) \prod_{\theta=\tau+1}^{r+\delta} d\mathfrak{x}_{\theta}^- \end{aligned}$$

$$e^{i\mathfrak{C}_{\tau} m_{\eta} \vec{\alpha}_{\tau}^{\eta}} \int_{\mathbb{R}^{4(\tau+\delta)}} \prod_{\tau=1}^{r+\delta} d\mathfrak{x}_{\tau}^+ \mathfrak{W}_0^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}, (\mathfrak{x}_{\theta}^+ + \vec{\alpha})^r_{\theta=\tau+1}) \hat{\varphi}(\mathcal{H}_{\eta}^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1}$$

$$|\hat{\varphi}((q_{\theta}^-, \mathfrak{x}_{\theta}^+)^r_{\theta=\tau+1})| \leq \frac{\mathfrak{C}(\rho_{\eta})^{\hat{\eta}} \|\varphi\|_{\mathfrak{H}}}{\sum_{\theta=\tau+1}^{r+\delta} (|q_{\theta}^0|^2 + |q_{\theta}^1|^2)^{\frac{\hat{\ell}}{2}}} + (|\mathfrak{x}_{\theta}^2|^2 + |\mathfrak{x}_{\theta}^4|^2)^{\frac{\hat{\kappa}}{2}}$$

$$e^{i\mathfrak{C}_{\tau} m_{\eta} \vec{\alpha}_{\tau}^{\eta}} \sum_{|\vec{\mathcal{M}}| \leq N} \int_{\mathbb{R}^{4(\tau+\delta)}} \prod_{\tau=1}^{r+\delta} d\mathfrak{x}_{\tau}^+ \mathfrak{D}^{\vec{\mathcal{M}}} \mathfrak{R}_{\vec{\mathcal{M}}}^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}; \alpha^+) \hat{\varphi}(\mathcal{H}_{\eta}^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1} = \mathcal{H}_4^{\eta}(\vec{\alpha})$$

$$\sum_{|\vec{\mathcal{M}}| \leq N} |\mathfrak{D}^{\vec{\mathcal{M}}}|((\mathfrak{x}_{\tau}^+)^r_{\tau=1}; \alpha^+) \leq \mathfrak{C}(\rho_{\eta})^{\hat{\kappa}} (|\alpha^+|^{\alpha} + \left(\sum_{\theta=\tau+1}^{r+\delta} |\mathfrak{x}_{\tau}^+|^4\right)^{\frac{\gamma}{2}})$$

$$\begin{aligned} \sum_{|\vec{\mathcal{M}}| \leq N} \int_{\mathbb{R}^{4(\tau+\delta)}} \prod_{\tau=1}^{r+\delta} d\mathfrak{x}_{\tau}^+ \mathfrak{D}^{\vec{\mathcal{M}}} \mathfrak{R}_{\vec{\mathcal{M}}}^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}; \alpha^+) \hat{\varphi}(\mathcal{H}_{\eta}^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1} \\ = \sum_{|\vec{\mathcal{M}}| \leq N} \int_{\{\mathcal{R} > \mathcal{R}_0 - \epsilon\}} \prod_{\tau=1}^{r+\delta} d\mathfrak{x}_{\tau}^+ \mathfrak{D}^{\vec{\mathcal{M}}} \mathfrak{R}_{\vec{\mathcal{M}}}^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}; \alpha^+) \hat{\varphi}(\mathcal{H}_{\eta}^-, \mathfrak{x}_{\tau}^+)^r_{\tau=1} = \mathcal{H}_4^{\eta}(\vec{\alpha}) \end{aligned}$$

$$\begin{aligned} |\mathcal{H}_4^{\eta}(\vec{\alpha})| &\leq \sum_{|\vec{\mathcal{M}}| \leq N} \int_{\{\mathcal{R} > \mathcal{R}_0 - \epsilon\}} \prod_{\tau=1}^{r+\delta} d\mathfrak{x}_{\tau}^+ \left| \mathfrak{R}_{\vec{\mathcal{M}}}^{\eta}((\mathfrak{x}_{\tau}^+)^r_{\tau=1}; \alpha^+) (\mathfrak{D}^{\vec{\mathcal{M}}} \hat{\varphi}(\mathcal{H}_{\eta}^-, \mathfrak{x}_{\theta}^+)^r_{\theta=1}) \right| - \\ &\leq \mathfrak{C}(\rho_{\eta})^{\hat{\kappa}} \int_{\{\mathcal{R} > \mathcal{R}_0 - \epsilon\}} (|\alpha^+|^{\alpha} + \mathcal{R}^{\gamma}) \frac{\mathfrak{C}(\rho_{\eta})^{\hat{\eta}} \|\varphi\|_{\hat{\rho}\hat{\sigma}}}{(\mathfrak{x} + \mathfrak{s}) |\mathcal{M}_{\mathfrak{K}}|^{\mathcal{L}} + \mathcal{R}^{\kappa + \hat{\mathcal{L}}}} \mathcal{R}^{2(\mathfrak{r} + \mathfrak{s})} / \mathfrak{H} \, d\mathcal{R} d\Omega \end{aligned}$$

$$\sum_{\Omega \in \Omega_{\tau}} \mathfrak{C}_{\Omega} \int_{\mathbb{R}} d\xi \int_{\mathbb{R}^4} d\mathfrak{x}_{\tau}^+ d\mathfrak{x}_{\tau+}^+ \, \hat{\varphi}_{\Omega,1}^{\eta}(\mathfrak{x}_{\tau}^+) \mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau+}^+ + \alpha^+) \hat{\varphi}_{\Omega,2}^{\eta}(\mathfrak{x}_{\tau+}^+) \mathfrak{g}(\xi)$$

Más, aplicando la función de Green, tenemos:

$$\begin{aligned} & (-({\mathfrak{C}}_{\mathcal{R}}{\mathcal{M}}_{\mathfrak{K}})^2 + \sum_{i=1}^4 \partial^2 / \partial \alpha^{i,2}) \mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = \delta(\alpha^1 - \xi^1, \mathfrak{x}_{\tau}^+ - \mathfrak{x}_{\tau^+}^+ - \alpha^+) \\ & \equiv \delta(\alpha - \xi_{\mathcal{R}}) \end{aligned}$$

$$\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = \frac{\frac{1}{(2\pi)^{\frac{4}{2}} \int_{\mathbb{R}^4} \mathfrak{d}\mathfrak{q} \, \mathbf{e}^{\mathfrak{i}\mathfrak{q} \cdot (\alpha - \xi_{\mathcal{R}})}}}{\omega^2} + |\mathfrak{q}|^4$$

$$\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = -\frac{\frac{1}{(2\pi)^{\frac{4}{2}}} 1}{\mathfrak{i}\mathcal{R} \int_0^{\infty} \lambda \, \mathbf{e}^{\mathfrak{i}\mathcal{R}\lambda}} - \frac{\mathbf{e}^{-\mathfrak{i}\mathcal{R}\lambda}}{\omega^2} + \lambda^2 \mathfrak{d}\lambda$$

$$= -\frac{\frac{1}{(2\pi)^{\frac{4}{2}}} 1}{\mathfrak{i}\mathcal{R} \int_{-\infty}^{\infty} \lambda \mathbf{e}^{-\mathfrak{i}\mathcal{R}\lambda}} / (\lambda - \mathfrak{i}\omega) (\lambda + \mathfrak{i}\omega) \mathfrak{d}\lambda$$

$$\mathfrak{R}_{\eta}(\alpha^1 - \xi^1; \mathfrak{x}_{\tau}^+, \mathfrak{x}_{\tau^+}^+ + \alpha^+) = \frac{\frac{2\omega \mathfrak{i}}{(2\pi)^{\frac{4}{2}}} 1}{\mathfrak{i}\mathcal{R}} \times \frac{\mathfrak{i}\omega \mathbf{e}^{-\mathcal{R}\omega}}{2\mathfrak{i}\omega} = \sqrt{2}\pi \mathbf{e}^{-\mathcal{R}|\mathfrak{C}_{\mathfrak{R}}\mathfrak{M}_{\eta}|} / 2\mathcal{R}$$

$$\begin{aligned} |\mathcal{H}_4^{\eta}(\vec{\alpha})\mathfrak{g}(\alpha^1)| &= \left| \left(\frac{\partial^2}{\partial \alpha^{0,2}} - \widehat{\mathfrak{N}} \right) \Psi_{\eta}(\alpha^0, \alpha) \right| \\ &= \left| \sum_{\Omega \in \Omega_{\mathcal{R}}} \mathfrak{C}_{\Omega} \int_{\mathbb{R}} \mathfrak{d}\xi^1 \int_{\mathbb{R}^4} \mathfrak{d}\vec{\mathfrak{x}} \, \hat{\phi}_{\Omega,1}^{\eta}(\mathfrak{x}^-) \mathfrak{R}_{\eta}(\alpha - \xi) \left(({\mathfrak{C}}_{\mathcal{R}}{\mathcal{M}}_{\eta})^2 + \widehat{\mathfrak{N}} \right) \hat{\phi}_{\Omega,2}^{\eta}(\mathfrak{x}^+) \mathfrak{g}(\xi^1) \right| \\ &\leq \sum_{\Omega \in \Omega_{\mathcal{R}}} |\mathfrak{C}_{\Omega}| \left\{ \int_{\mathbb{R}} \mathfrak{d}\xi^1 \left(|\mathfrak{g}(\xi^1)| + \frac{1}{\mathcal{M}_2^4 |\mathfrak{g}''(\xi^1)|} \right) \right. \\ &\quad \cdot \left. \frac{\sqrt{\frac{2}{\omega}} \mathbf{e}^{-\mathcal{M}_0 |\alpha^+ + \mathfrak{x}^+ - \mathfrak{y}^+|}}{2|\alpha^+ + \mathfrak{x}^+ - \mathfrak{y}^+|} \times \int_{\mathbb{R}^4} \mathfrak{d}\mathfrak{y}^+ |\hat{\phi}_{\Omega,1}^{\eta}(\mathfrak{y}^+)| \cdot \int_{\mathbb{R}^4} \mathfrak{d}\mathfrak{x}^+ \left| \left(({\mathfrak{C}}_{\mathcal{R}}{\mathcal{M}}_{\eta})^2 + \widehat{\mathfrak{N}} \right) \hat{\phi}_{\Omega,2}^{\eta}(\mathfrak{x}^+) \right| \right\} \\ &< \mathfrak{C}(\rho_{\eta})^{\widehat{\mathfrak{K}}} \mathbf{e}^{\mathcal{M}_0 \epsilon} \|\mathfrak{g}\| \|\mathfrak{p}_4, \mathfrak{q}_4\| \varphi_1 \|\mathfrak{p}_1 \mathfrak{q}_1\| \varphi_4 \|\mathfrak{p}_4 \mathfrak{q}_4\| \cdot \frac{\mathbf{e}^{-\mathcal{M}_0 |\alpha^+|}}{|\alpha^+|} - \epsilon \end{aligned}$$



Por lo que, la integral de superficie, queda definida de la siguiente manera:

$$\begin{aligned}\rho_{\sigma}^{\alpha\beta} &= \frac{1}{\sqrt{\det\left(1 + \mathcal{W}_{\alpha\beta}^{\text{cb},\mathcal{T}} \mathcal{W}_{\alpha\beta}^{\text{cb}}\right)}} \equiv \frac{|\mathfrak{I}_{\alpha\beta}^{\sigma}|}{\sqrt{\det(\mathfrak{I}_{\alpha\beta}^{\sigma,\mathcal{T}} \mathfrak{I}_{\alpha\beta}^{\sigma} + \mathfrak{I}_{\text{cb}}^{\sigma,\mathcal{T}} \mathfrak{I}_{\text{cb}}^{\sigma})}} = \int_{\mathfrak{s}}^{\infty} \mathfrak{d}\rho \\ &= \sum_{0 \leq \alpha \leq \beta \leq 4} \int_{\mathbb{I}^2}^{\infty} \rho_{\sigma}^{\alpha\beta}(\mathfrak{s}, \mathfrak{t}) |\mathfrak{I}_{\alpha\beta}^{\sigma}|(\mathfrak{s}, \mathfrak{t}) \mathfrak{d}\mathfrak{s} \mathfrak{d}\mathfrak{t} (\dot{\sigma} + \dot{\sigma})(\dot{\sigma} - \dot{\sigma}) \mathfrak{d}\sigma \mathfrak{d}\hat{\sigma}\end{aligned}$$

Por otro lado, en este punto, es indispensable, añadir algunos planteamientos teóricos adicionales cuyo propósito es reforzar la tesis formulada en trabajos anteriores, siendo éstos, los que siguen a continuación:

Curvatura geométrica de los campos cuánticos y la existencia de agujeros deformantes

$$\begin{aligned}\frac{\mathfrak{d}^2 \chi^{\alpha}}{\mathfrak{d}\mathfrak{s}^2} + \frac{\frac{\Gamma_{\beta\gamma}^{\alpha} \mathfrak{d}\chi^{\beta}}{\mathfrak{d}\mathfrak{s}} \mathfrak{d}\chi^{\gamma}}{\mathfrak{d}\mathfrak{s}} &= \Gamma_{\beta\gamma}^{\alpha} = \frac{1}{2\mathfrak{g}^{\alpha\beta}(\mathfrak{g}_{\mathfrak{d}\beta,\mathfrak{c}} \mathfrak{g}_{\mathfrak{d}\mathfrak{c},\beta} \mathfrak{g}_{\beta\mathfrak{c},\mathfrak{d}})} \int \mathfrak{d}\mathfrak{s} = \int \mathfrak{d}\rho \sqrt{\frac{\mathfrak{g}^{\alpha\beta} \mathfrak{d}\chi^{\alpha}}{\mathfrak{d}\rho} \frac{\mathfrak{d}\chi^{\beta}}{\mathfrak{d}\rho}} \\ \mathcal{R}_{\mathfrak{ij}} &= -\frac{1}{2\mathfrak{g}_{\mathfrak{ij}} \mathfrak{R}} = -\frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{\mathfrak{ij}}} = \mathcal{F}(\mathfrak{g}_{\mathfrak{ij}}) = \int_{\mathcal{M}}^{\infty} \mathcal{R} \sqrt{-\mathfrak{g}} \mathfrak{d}\mathfrak{x} \\ \mathcal{R} &= \frac{4\pi\mathfrak{G}}{c^4} \lim_{\lambda \rightarrow 0} \frac{1}{\lambda} \left(\mathcal{F}(\mathfrak{g}_{\mathfrak{ij}} + \lambda \chi_{\mathfrak{ij}}) - \mathcal{F}(\mathfrak{g}_{\mathfrak{ij}}) \right) = \delta \mathcal{F}((\mathfrak{g}_{\mathfrak{ij}}), \chi) \mathcal{R}_{\mathfrak{ij}} - \frac{1}{2\mathfrak{g}_{\mathfrak{ij}} \mathfrak{R}} \\ &= -\frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{\mathfrak{ij}}} - \mathcal{D}_{\mathfrak{i}} \mathcal{D}_{\mathfrak{j}} \varphi \operatorname{div} \left(\mathcal{D}_{\mathfrak{i}} \mathcal{D}_{\mathfrak{j}} \varphi + \frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{\mathfrak{ij}}} \right) \mathfrak{R} \left(\frac{16\pi\mathfrak{G}}{c^4} \right) \mathcal{T}_{\mathfrak{ij}} + \phi, \int_{\mathcal{M}}^{\infty} \phi \sqrt{-\mathfrak{g}} \mathfrak{d}\mathfrak{x} \\ \mathfrak{d}\mathfrak{s}^2 &= \epsilon^{\mu} c^4 \mathfrak{d}\mathfrak{t}^2 + \epsilon^{\nu} \mathfrak{d}\mathfrak{r}^2 + \mathfrak{r}^2 (\mathfrak{d}\theta^2 + \sin^2 \theta \mathfrak{d}\gamma^2) \\ \mathfrak{F} &= m\mathfrak{M}\mathfrak{G} \left(-\frac{1}{\mathfrak{r}^2} + \frac{1}{\delta \left(2 + \frac{\delta}{\tau} \right) \varphi'} + \frac{\mathcal{R}\tau}{\delta} \right) \\ \mathfrak{F} &= m\mathfrak{M}\mathfrak{G} \left(-\frac{1}{\mathfrak{r}^2} + \left(2 + \frac{\delta}{\tau} \right) \varepsilon \mathfrak{r}^2 + \frac{\mathcal{R}\tau}{\delta} + \frac{1}{\delta \left(2 + \frac{\delta}{\tau} \right) \mathfrak{r}^2 \int \mathfrak{r}^{-2} \mathcal{R} \mathfrak{d}\mathfrak{r}} \right) \\ \mathfrak{F} &= m\mathfrak{M}\mathfrak{G} \left(-\frac{1}{\mathfrak{r}^2} - \frac{\kappa_0}{\tau} + \kappa_1 \mathcal{R} \right) \\ \mu &= \left\{ \mu_{\mathfrak{i}_1 \dots \mathfrak{i}_{\mathfrak{s}}}^{\mathfrak{j}_1 \dots \mathfrak{j}_{\mathfrak{r}}}(\mathfrak{x}) \mid 1 \leq \mathfrak{i}_1 \dots \mathfrak{i}_{\mathfrak{s}}, \mathfrak{j}_1 \dots \mathfrak{j}_{\mathfrak{r}} \leq \eta \right\} \\ \mathcal{L}^{\rho}(\mathfrak{E}) &= \left\{ \mu: \mathcal{M} \rightarrow \mathfrak{E} \mid \int_{\mathcal{M}}^{\infty} \|\mu\|^{\rho} \mathfrak{d}\mathfrak{x} < \infty \right\}\end{aligned}$$

$$\begin{aligned}
\|\mu\|_{\mathcal{L}^\rho} &= \left[\int_{\mathcal{M}} \|\mu\|^\rho \mathfrak{d}\mathfrak{x} \right]^{1/\rho} = \left[\int_{\mathcal{M}} \sum |\mu_{i_1 \dots i_s}^{j_1 \dots j_r}|^\rho \mathfrak{d}\mathfrak{x} \right]^{1/\rho} \\
(\mu, \nu) &= \int_{\mathcal{M}} \mathfrak{g}_{i_1 \kappa_i} \dots \mathfrak{g}_{i_{\mathcal{R}} \kappa_{\mathcal{R}}} \mathfrak{g}^{i_1 l_1} \dots \mathfrak{g}^{i_s l_s} \mu_{i_1 \dots i_s}^{j_1 \dots j_r} \nu_{l_1 \dots l_s}^{\kappa_1 \dots \kappa_r} \sqrt{-\mathfrak{g}} \mathfrak{d}\mathfrak{x} \Delta_\mu \Delta_\nu \nabla^\mu \nabla^\nu \\
\mu &= \{\mu_{i_1 \dots i_s}^{j_1 \dots j_r}\} \nabla_\mu = \{\mathcal{D}_\kappa \mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, \nabla_\mu: \mathfrak{M} \rightarrow \mathcal{T}_{s+1}^r \mathcal{M}, \nabla^* \mu \{g^{\kappa l} \mathcal{D}_l \mu\}: \mathfrak{M} \rightarrow \mathcal{T}_{s+1}^r \mathcal{M}, \operatorname{div} \mu = \{\mathcal{D}_l \mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, \mu \\
&= \{\mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, \operatorname{div} \mu = \{\mathcal{D}^l \mu_{i_1 \dots i_s}^{j_1 \dots j_r}\}, (\nabla^* \mu, \nu) = -(\mu \operatorname{div} \nu), (\nabla \mu, \nu) \\
&= -(\mu \operatorname{div} \nu), \lim_{\eta \rightarrow \infty} (\mathfrak{G}_{\mu_\eta}, \nu) = (\mathfrak{G}_{\mu_0}, \nu), (\mathfrak{G} \mu, \mu) \geq \alpha \|\mu\|^4 - \beta \\
\mu &= \nabla_\varphi + \nu + \hbar, \mathcal{H}(\mathcal{T}_s^r \mathcal{M}) = \{h \in \mathcal{L}^2(\mathcal{T}_s^r \mathcal{M}) | \nabla h, \operatorname{div} h\}, \mathcal{L}^2(\mathbb{E}) = \mathfrak{G}(\mathfrak{E}) \oplus \mathcal{L}_\mathcal{D}^2(\mathfrak{E}), \mathcal{L}^2(\mathbb{E}) \\
&= \mathfrak{G}(\mathfrak{E}) \oplus \mathfrak{H}(\mathfrak{E}) \oplus \mathcal{L}_\mathfrak{N}^2(\mathfrak{E}), \mathfrak{G}(\mathfrak{E}) = \{\nu \in \mathcal{L}^2(\mathbb{E}) | \nu = \nabla \varphi, \varphi \in \mathcal{H}^1(\mathcal{T}_{s-1}^r \mathcal{M})\}, \mathcal{L}_\mathcal{D}^2(\mathfrak{E}) \\
&= \{\nu \in \mathcal{L}^2(\mathbb{E}) | \operatorname{div} \nu = 0\}, \mathcal{L}_\mathfrak{N}^2(\mathfrak{E}) = \{\nu \in \mathcal{L}^2(\mathbb{E}) | \nabla_\nu \neq 0\}, \mathcal{L}_\mathcal{D}^2(\mathfrak{E}) \perp \mathfrak{G}(\mathfrak{E}), \mathcal{L}_\mathfrak{N}^2(\mathfrak{E}) \\
&\perp \mathcal{H}(\mathfrak{E}), \mathfrak{G}(\mathfrak{E}) \perp \mathcal{H}(\mathfrak{E}), \mathfrak{E} = \mathfrak{E}_1 \oplus \mathfrak{E}^\kappa, \Delta \varphi = \operatorname{div} \mu, \Delta = \operatorname{div} \nabla, \nu = \mu - \nabla \varphi \\
&\in \mathcal{L}^2(\mathbb{E}), (\nu, \nabla \varphi) = 0, (\nabla \varphi - \mu, \nabla \psi) = 0, \mathcal{H} = \mathcal{H}^1(\widehat{\mathfrak{E}}) \setminus \widehat{\mathcal{H}}, \widehat{\mathcal{H}} \\
&= \{\psi \in \mathcal{H}^1(\widehat{\mathfrak{E}}) | \nabla \varphi = 0\}, (\mathfrak{G} \varphi, \psi) = (\nabla \varphi, \nabla \psi), (\mathfrak{G} \varphi, \varphi) = (\nabla \varphi, \nabla \varphi) = \|\varphi\|^4, \Delta \varphi \\
&= \mathfrak{f}, \mathcal{H}^\kappa(\mathfrak{E}) = \mathcal{H}_\mathcal{D}^\kappa \oplus \mathfrak{G}^\kappa, \mathcal{L}^2(\mathbb{E}) = \mathcal{L}_\mathcal{D}^2 \oplus \mathfrak{G}, \mathcal{H}_\mathcal{D}^\kappa = \{\mu \in \mathcal{H}^\kappa(\mathbb{E}) | \operatorname{div} \mu = 0\}, \mathfrak{G}^\kappa \\
&= \{\mu \in \mathcal{H}^\kappa(\mathbb{E}) | \mu = \nabla \psi\}, \widehat{\Delta} \mu = \wp \Delta \mu, \Delta = \operatorname{div} \nabla = \mathcal{D}^\kappa \mathcal{D}_\kappa = \frac{\mathfrak{g}^{\kappa l} \partial^2}{\partial \mathfrak{x}^\kappa \partial \mathfrak{x}^l} + \mathfrak{B} \\
\mathfrak{A} &= \frac{\mathfrak{g}^{\kappa l} \partial^2}{\partial \mathfrak{x}^\kappa \partial \mathfrak{x}^l}: \mathcal{H}^2(\mathcal{M}, \mathcal{R}^\mathfrak{N}) \rightarrow \mathcal{L}^2(\mathcal{M}, \mathcal{R}^\mathfrak{N}), \Delta: \mathcal{H}^4(\mathfrak{E}) \rightarrow \mathcal{L}^2(\mathfrak{E}), \widehat{\Delta} = \mathcal{P} \Delta: \mathcal{H}_\mathcal{D}^2(\mathfrak{E}) \mathcal{L}_\mathcal{D}^2(\mathfrak{E}), \widehat{\mathcal{H}} \\
&= \{\mu \in \mathcal{H}_\mathcal{D}^2(\mathfrak{E}) | \widehat{\Delta} \mu = 0\} \\
\int_{\mathcal{M}} (\widehat{\Delta} \mu, \mu) \sqrt{-\mathfrak{g}} \mathfrak{d}\mathfrak{x} &= \int_{\mathcal{M}} (\Delta \mu, \mu) \sqrt{-\mathfrak{g}} \mathfrak{d}\mathfrak{x} = - \int_{\mathcal{M}} (\nabla \mu, \nabla \mu) \sqrt{-\mathfrak{g}} \mathfrak{d}\mathfrak{x} = 0 \\
\frac{\partial}{\partial \mathfrak{x}^i} (\mathcal{D}^\kappa \mu_{\kappa l}) - \frac{\partial}{\partial \mathfrak{x}^i} (\mathcal{D}^\kappa \mu_{\kappa j}) &= \frac{\partial \Delta \varphi_i}{\partial \mathfrak{x}^i} - \frac{\partial \Delta \varphi_j}{\partial \mathfrak{x}^i} \\
\Delta \varphi_i &= -(\delta \mathfrak{d} + \mathfrak{d} \delta) \varphi_i - \mathcal{R}_i^\kappa \varphi_\kappa \\
\widehat{\Delta} \varphi &= - \frac{1}{\sqrt{-\mathfrak{g}} \frac{\partial}{\partial \mathfrak{x}^i} \left(\frac{\sqrt{-\mathfrak{g}} \mathfrak{g}^{ij} \partial \varphi}{\partial \mathfrak{x}^j} \right)} \\
(\delta \mathfrak{d} + \mathfrak{d} \delta) \varphi_i &= \frac{\partial}{\partial \mathfrak{x}^i} \widehat{\Delta} \varphi \Leftrightarrow \varphi_i = \frac{\partial \varphi}{\partial \mathfrak{x}^i}
\end{aligned}$$

$$\Delta\varphi_i=-\frac{\partial}{\partial\mathfrak{x}^i}\hat{\Delta}\varphi_i-\mathcal{R}_i^\kappa\frac{\partial\varphi}{\partial\mathfrak{x}^\kappa}\Leftrightarrow\varphi_i=\frac{\partial\varphi}{\partial\mathfrak{x}^i}$$

$$\frac{\partial}{\partial\mathfrak{x}^j}(\mathcal{D}^\kappa\mu_{\kappa l})-\frac{\partial}{\partial\mathfrak{x}^i}(\mathcal{D}^\kappa\mu_{\kappa j})=\frac{\partial}{\partial\mathfrak{x}^i}\Big(\mathcal{R}_i^\kappa\frac{\partial\varphi}{\partial\mathfrak{x}^\kappa}\Big)-\frac{\partial}{\partial\mathfrak{x}^j}\Big(\mathcal{R}_i^\kappa\frac{\partial\varphi}{\partial\mathfrak{x}^\kappa}\Big)$$

$$\mathfrak{G}\colon\mathcal{M}\longrightarrow\mathcal{T}_4^0\mathcal{M}=\mathcal{T}^*\mathcal{M}\otimes\mathcal{T}^*\mathcal{M},\mathfrak{G}=\{\mathfrak{g}_{ij}(\mathfrak{x})\}(\mathfrak{g}_{ij})=(\mathfrak{g}_{ij})^{-1},\mathfrak{G}^{-1}=\{\mathfrak{g}_{ij}\}\colon\mathcal{M}\longrightarrow\mathcal{T}_4^0\mathcal{M}=\mathcal{T}\mathcal{M}\otimes\mathcal{T}\mathcal{M}$$

$$\mathcal{W}^{\mathfrak{m},2}(\mathfrak{M},\mathfrak{g})\subset \mathcal{W}^{\mathfrak{m},2}(\mathcal{T}_4^0\mathcal{M})\oplus \mathcal{W}^{\mathfrak{m},2}(\mathcal{T}_4^0\mathcal{M})$$

$$\mathcal{F}(\mathfrak{g}^{ij})=\int\limits_{\mathcal{M}}^{\infty}\mathfrak{f}(x,\mathfrak{g}^{ij},\cdots,\mathcal{D}^m\mathfrak{g}_{ij})\sqrt{-\mathfrak{g}\mathfrak{d}\mathfrak{x}}$$

$$\mathfrak{g}_{ij}+\lambda\chi_{ij}\in\mathcal{W}^{\mathfrak{m},2}(\mathcal{M},\mathfrak{g})\forall 0\leq|\lambda|\leq\lambda_0,\mathfrak{g}^{ij}+\lambda\chi^{ij}\in\mathcal{W}^{\mathfrak{m},2}(\mathcal{M},\mathfrak{g})\forall 0\leq|\lambda|\leq\lambda_0$$

$$\delta_*\mathcal{F}\colon\mathcal{W}^{\mathfrak{m},2}(\mathcal{M},\mathfrak{g})\longrightarrow\mathcal{W}^{-\mathfrak{m},2}(\mathcal{T}_4^0\mathcal{M}),\delta^*\mathcal{F}\colon\mathcal{W}^{\mathfrak{m},2}(\mathcal{M},\mathfrak{g})\longrightarrow\mathcal{W}^{-\mathfrak{m},2}(\mathcal{T}_4^0\mathcal{M}),(\delta_*\mathcal{F}(\mathfrak{g}_{ij}),\mathfrak{X})$$

$$=\frac{\mathfrak{d}}{\mathfrak{d}\lambda\mathcal{F}|_{\mathfrak{g}^{ij}+\lambda\chi_{ij}}|_{\lambda=0}},(\delta^*\mathcal{F}(\mathfrak{g}^{ij}),\mathfrak{X})=\frac{\mathfrak{d}}{\mathfrak{d}\lambda\mathcal{F}|_{\mathfrak{g}^{ij}+\lambda\chi^{ij}}|_{\lambda=0}},\delta_*\mathcal{F}(\mathfrak{g}_{ij})\colon\mathcal{M}$$

$$\longrightarrow\mathcal{T}\mathcal{M}\otimes\mathcal{T}\mathcal{M},\delta^*\mathcal{F}(\mathfrak{g}^{ij})\colon\mathcal{M}\longrightarrow\mathcal{T}^*\mathcal{M}\otimes\mathcal{T}^*\mathcal{M},((\delta_*\mathcal{F})_{\kappa\mathfrak{l}},\delta\mathfrak{g}^{\kappa\mathfrak{l}})$$

$$=-((\delta_*\mathcal{F})_{\kappa\mathfrak{l}},\mathfrak{g}^{\kappa\mathfrak{i}}\mathfrak{g}^{\mathfrak{l}\mathfrak{j}}\delta\mathfrak{g}_{ij})=(-\mathfrak{g}^{\kappa\mathfrak{i}}\mathfrak{g}^{\mathfrak{l}\mathfrak{j}}(\delta_*\mathcal{F})_{\kappa\mathfrak{l}},\delta\mathfrak{g}_{ij})=((\delta^*\mathcal{F})^{\mathfrak{i}\mathfrak{j}},\delta\mathfrak{g}^{\mathfrak{i}\mathfrak{j}})$$

$$(\delta\mathcal{F}(\mathfrak{g}_{ij}),\chi)=\frac{\mathfrak{d}}{\mathfrak{d}\lambda\mathcal{F}|_{\mathfrak{g}^{ij}+\lambda\chi^{ij}}|_{\lambda=0}}=\int\limits_{\mathcal{M}}^{\infty}(\delta\mathcal{F}(\mathfrak{g}_{ij}))_{\kappa\mathcal{L}}\chi^{\kappa\mathcal{L}}\sqrt{-\mathfrak{g}\mathfrak{d}\mathfrak{x}}$$

$$\mathcal{L}^2(\mathbb{E})=\mathcal{L}^2_{\mathfrak{s}}(\mathfrak{E})\oplus\mathcal{L}^2_{\mathfrak{c}}(\mathfrak{E}),\mathcal{L}^2_{\mathfrak{s}}(\mathfrak{E})=\{\mu\in\mathcal{L}^2(\mathbb{E})|\mu_{ij}\mu_{ji}|\}$$

$$\delta\mathcal{F}\colon\mathcal{W}^{\mathfrak{m},2}(\mathcal{M},\mathfrak{g})\longrightarrow\mathcal{W}^{-\mathfrak{m},2}(\mathfrak{E}),\mathcal{L}_D^2(\mathfrak{E})=\{\chi\in\mathcal{L}^2(\mathbb{E})|div\,\chi=0\},(\delta\mathcal{F}(\mathfrak{g}_{ij}))_{\kappa\mathcal{L}}$$

$$=\mathcal{D}_\kappa\mathcal{D}_\mathcal{L}\varphi,\int\limits_{\mathcal{M}}^{\infty}(\delta\mathcal{F}(\mathfrak{g}_{ij}))_{\kappa\mathcal{L}}\chi^{\kappa\mathcal{L}}\sqrt{-\mathfrak{g}\mathfrak{d}\mathfrak{x}},(\delta\mathcal{F}(\mathfrak{g}_{ij}))_{\kappa\mathcal{L}}=\nu_{\kappa\mathcal{L}}+\mathcal{D}_\kappa\psi_\mathcal{L},(\mathcal{D}_\kappa\psi_\mathcal{L},\chi^{\kappa\mathcal{L}})$$

$$=\int\limits_{\mathcal{M}}^{\infty}\mathcal{D}_\kappa\psi_\mathcal{L},\chi^{\kappa\mathcal{L}}\sqrt{-\mathfrak{g}\mathfrak{d}\mathfrak{x}}=-\int\limits_{\mathcal{M}}^{\infty}\psi_\mathcal{L}\mathcal{D}_\kappa\chi^{\kappa\mathcal{L}}\sqrt{-\mathfrak{g}\mathfrak{d}\mathfrak{x}},\mathcal{D}_\kappa\chi^{\kappa\mathcal{L}}=\mathcal{D}_\kappa(\mathfrak{g}^{\kappa\mathfrak{i}}\mathfrak{g}^{\mathfrak{l}\mathfrak{j}}\nu_{ij})$$

$$=\mathfrak{g}^{\mathfrak{l}\mathfrak{j}}(\mathfrak{g}^{\mathfrak{i}\kappa}\mathcal{D}_\kappa\nu_{ij})=\mathfrak{g}^{\mathfrak{l}\mathfrak{j}}\mathcal{D}^{\mathfrak{i}}\nu_{ij}\|v\|_{L^2}^2\int\limits_{\mathcal{M}}^{\infty}\mathfrak{g}^{\kappa\mathfrak{i}}\mathfrak{g}^{\mathfrak{l}\mathfrak{j}}\nu_{\kappa\mathcal{L}}\nu_{ij}\sqrt{-\mathfrak{g}\mathfrak{d}\mathfrak{x}},(\delta\mathcal{F}(\mathfrak{g}_{ij}))_{\kappa\mathcal{L}}=\mathcal{D}_\kappa\psi_\mathcal{L}=\mathcal{D}_\mathcal{L}\psi_\kappa$$

$$\frac{\partial\psi_\mathcal{L}}{\partial\mathfrak{x}^\kappa}=\frac{\partial\psi_\kappa}{\partial\mathfrak{x}^\mathcal{L}}\mathfrak{d}(\psi_\kappa\mathfrak{d}\mathfrak{x}^\kappa)=\Big(\frac{\partial\psi_\mathcal{L}}{\partial\mathfrak{x}^\kappa}-\frac{\partial\psi_\kappa}{\partial\mathfrak{x}^\mathcal{L}}\Big)\mathfrak{d}\mathfrak{x}^\mathcal{L}\wedge\mathfrak{d}\mathfrak{x}^\kappa$$

$$\mathfrak{d}\varphi=\frac{\partial\varphi}{\partial\mathfrak{x}^\kappa\mathfrak{d}\mathfrak{x}^\kappa}=\psi_\kappa\mathfrak{d}\mathfrak{x}^\kappa$$

$$\begin{aligned}
\mathcal{F}(\mathfrak{g}_{ij}) &= \int_{\mathcal{M}}^{\infty} \left(\mathcal{R} + \frac{16\pi\mathfrak{G}}{c^4} \mathfrak{g}^{ij} \delta_{ij} \right) \sqrt{-\mathfrak{g}} d\mathfrak{x} \\
\delta \mathcal{F}(\mathfrak{g}_{ij}) &= \mathcal{R}_{ij} - \frac{1}{2\mathfrak{g}_{ij}\mathfrak{R}} + \frac{16\pi\mathfrak{G}}{c^4} \mathfrak{S}_{ij} \\
\mathfrak{S}_{ij} &= \delta_{ij} - \frac{1}{2\mathfrak{g}_{ij}\delta} + \frac{\mathfrak{g}^{\kappa\mathfrak{l}} \partial \delta_{\kappa\mathfrak{l}}}{\partial \mathfrak{g}^{ij}}, \delta = \mathfrak{g}^{\kappa\mathfrak{l}} \delta_{\kappa\mathfrak{l}} \\
\mathcal{R}_{ij} &= \frac{1}{2\mathfrak{g}^{\kappa\mathfrak{l}}} \left(\frac{\partial^2 \mathfrak{g}_{\kappa\mathfrak{l}}}{\partial \mathfrak{x}^i \partial \mathfrak{x}^j} + \frac{\partial^2 \mathfrak{g}_{ij}}{\partial \mathfrak{x}^{\kappa} \partial \mathfrak{x}^{\mathfrak{l}}} - \frac{\partial^2 \mathfrak{g}_{i\mathfrak{l}}}{\partial \mathfrak{x}^j \partial \mathfrak{x}^{\kappa}} - \frac{\partial^2 \mathfrak{g}_{\kappa j}}{\partial \mathfrak{x}^i \partial \mathfrak{x}^{\mathfrak{l}}} \right) + \mathfrak{g}_{\kappa\mathfrak{l}} \mathfrak{g}_{rs} \left(\Gamma_{\kappa\mathfrak{l}}^r \Gamma_{ij}^s - \Gamma_{i\mathfrak{l}}^r \Gamma_{j\kappa}^s \right), \Gamma_{ij}^{\kappa} \\
&= 1/2 \mathfrak{g}^{\kappa\mathfrak{l}} \left(\frac{\partial \mathfrak{g}_{i\mathfrak{l}}}{\partial \mathfrak{x}^j} + \frac{\partial \mathfrak{g}_{i\mathfrak{l}}}{\partial \mathfrak{x}^i} - \frac{\partial \mathfrak{g}_{i\mathfrak{l}}}{\partial \mathfrak{x}^{\mathfrak{l}}} \right) \\
\mathcal{R}_{ij} &= -\frac{1}{2\mathfrak{g}_{ij}\mathfrak{R}} = -\frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} - \mathcal{D}_i \mathcal{D}_j \varphi, \operatorname{div} \left(\mathcal{D}_i \mathcal{D}_j \varphi + \frac{16\pi\mathfrak{G}}{c^4 \mathcal{T}_{ij}} \right), \mathcal{R} = \frac{16\pi\mathfrak{G}}{c^4} \mathcal{T} + \phi, \mathcal{T} = \mathfrak{g}^{ij} \mathcal{T}_{ij}, \phi \\
&= \mathfrak{g}^{ij} \mathcal{D}_i \mathcal{D}_j \varphi, \int_{\mathcal{M}}^{\infty} \phi \sqrt{-\mathfrak{g}} d\mathfrak{x} \\
\mathcal{R}_{ij} - \frac{1}{2\mathfrak{g}_{ij}\mathfrak{R}} &= -\frac{16\pi\mathfrak{G}}{c^4} \mathcal{T}_{ij} - \mathcal{D}_i \mathcal{D}_j \varphi + \alpha \mathcal{D}_i \psi_j, \Delta \psi_j + \mathfrak{g}^{ik} \mathcal{R}_{ij} \psi_{\kappa}, \mathcal{D}_i \psi_j = \mathcal{D}_j \psi_i \\
ds^2 &= -\left(1 - \frac{2\mathfrak{M}\mathfrak{G}}{\mathfrak{G}^2 \mathfrak{r}}\right) \mathfrak{G}^2 dt^2 + \frac{dr^2}{\left(1 - \frac{2\mathfrak{M}\mathfrak{G}}{\mathfrak{G}^2 \mathfrak{r}}\right)} + \mathfrak{r}^2 (d\theta^2 + \sin^2 \theta d\varphi^2) \\
\mathcal{R}_{ij} &= -\frac{16\pi\mathfrak{G}}{c^4} \left(\mathcal{T}_{ij} - \frac{1}{2\mathfrak{g}_{ij}\mathcal{T}} \right) - \left(\mathcal{D}_i \mathcal{D}_j \varphi - \frac{1}{2\mathfrak{g}_{ij}\phi} \right), \mathcal{T} = \mathfrak{g}^{\kappa\mathfrak{l}} \mathcal{D}_{\kappa} \mathcal{D}_{\mathfrak{l}} \varphi \\
\Delta \left(\frac{\partial \varphi}{\partial \mathfrak{x}^{\kappa}} \right) \mathfrak{r} &\gg \frac{2\mathfrak{M}\mathfrak{G}}{c^2}, \mathcal{R}_{ij} = \frac{\partial \Gamma_{ik}^{\kappa}}{\partial \mathfrak{x}^j} - \frac{\partial \Gamma_{ij}^{\kappa}}{\partial \mathfrak{x}^{\kappa}} + \Gamma_{i\mathfrak{r}}^{\kappa} \Gamma_{j\kappa}^{\mathfrak{r}} - \Gamma_{ij}^{\kappa} \Gamma_{\kappa\mathfrak{r}}^{\mathfrak{r}}
\end{aligned}$$

Gravedad cuántica

$$\begin{aligned}
\kappa \epsilon_{\mu\nu} &= \mathcal{R}_{\mu\nu} - \frac{\mathfrak{g}_{\mu\nu} \mathfrak{R}}{2} \sim \mathcal{R}_{\mu\nu}, \frac{\delta \varpi \mathcal{L}_{\rho} \epsilon_{\mu\nu}}{\epsilon_{\rho\sigma}} = \mathcal{R}_{\mu\nu} = (\mathfrak{D}_{\mu} \mathfrak{D}_{\nu}), \epsilon_{\rho\sigma} \mathcal{L}_{\rho\sigma} = \hbar c (\mathfrak{D}_{\mu} \mathfrak{D}_{\nu}) = \mathcal{I} \epsilon_{\mu\nu}^{\kappa} \mathfrak{D}_{\kappa}, \frac{\epsilon_{\mu\nu}}{\hbar c} \\
&= (\mathfrak{D}_{\mu} \mathfrak{D}_{\nu}) = \mathcal{I} \epsilon_{\mu\nu}^{\kappa} \mathfrak{D}_{\kappa}, \varepsilon = \hbar c \kappa = \frac{\hbar c}{\lambda} = \hbar \mathfrak{S}, \mathcal{R}_{\mu\nu} \psi^{\mathfrak{A}} = (\mathfrak{D}_{\mu} \mathfrak{D}_{\nu}) \psi^{\mathfrak{A}}, (\mathfrak{D}_{\mu} \mathfrak{D}_{\nu}) \\
&= (\mathfrak{D}'_{\mu} \pm \Gamma'_{\mu}, \mathfrak{D}'_{\mu} \pm \Gamma'_{\mu}) = (\mathfrak{D}'_{\mu} \mathfrak{D}'_{\nu}) \pm \mathfrak{D}'_{\mu} \Gamma'_{\nu} \mp \mathfrak{D}'_{\nu} \Gamma'_{\mu} \pm (\mathfrak{I}'_{\mu} \Gamma'_{\nu}) \pm \Gamma'_{\mu} \mathfrak{D}'_{\nu} \mp \Gamma'_{\nu} \mathfrak{D}'_{\mu}, \mathcal{R}'_{\mu\nu} \\
&= \mathfrak{D}'_{\mu} \Gamma'_{\nu} - \mathfrak{D}'_{\nu} \Gamma'_{\mu} + (\Gamma'_{\mu} \Gamma'_{\nu}) (\mathfrak{D}_{\mu} \mathfrak{D}_{\nu}) = (\mathfrak{D}'_{\mu} \mathfrak{D}'_{\nu}) \pm \mathcal{R}'_{\mu\nu} \mp \epsilon_{\mu\nu} \Gamma'_{\mu} \mathfrak{D}'_{\nu}, \mathcal{R}'_{\mu\nu} = \frac{\mathfrak{m}^4 c^4}{\hbar c}
\end{aligned}$$



$$\begin{aligned}\frac{\epsilon_{\mu\nu}}{\hbar c} &= (\mathcal{D}'_\mu \mathcal{D}'_\nu) \pm \frac{m^4 c^4}{\hbar c} \mp \epsilon_{\mu\nu} \Gamma'_\mu \mathcal{D}'_\nu, \frac{\epsilon_{\mu\nu}}{\hbar c} = \mathcal{R}_{\mu\nu}^0 \pm \frac{m^4 c^4}{\hbar c} \mp \epsilon_{\mu\nu} \Gamma'_\mu \mathcal{D}'_\nu, \epsilon_{\mu\nu} \\ &= \hbar c \mathcal{R}_{\mu\nu}^0 \pm m^4 c^4 \mp \Im \gamma^\mu \hbar c \mathcal{D}'_\nu, \epsilon_{\mu\nu} - \hbar c \mathcal{R}_{\mu\nu}^0 = \epsilon_{\mu\nu}^0 = \Im \hbar c \gamma^\mu \mathcal{D}'_\nu - m^4 c^4\end{aligned}$$

$$\begin{aligned}\left(\frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma}}\right) \cdot \frac{1}{\mathcal{L}_\rho} &= \mathcal{R}_{\mu\nu}^0 + \Im \gamma^\kappa \mathcal{D}_\kappa - m^4, \frac{\epsilon^{\mu\nu} \epsilon_{\mu\nu}}{(\epsilon_{\rho\sigma} \mathcal{L}_\rho)^2} = \mathcal{R}_{\mu\nu}^0 \mathcal{R}^{0\mu\nu} + \Box^2 + \hat{m}^4, \frac{\epsilon^4}{(\epsilon_{\rho\sigma} \mathcal{L}_\rho)^2} \\ &= \frac{\varepsilon^2}{(\hbar c)^4} = \mathbb{R}^4 = \mathcal{R}_\omega^0 + \Box^2 + \hat{m}^4 + \rho^4 + \sigma^4\end{aligned}$$

$$\begin{aligned}\frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma} \mathcal{L}_\rho} &= (\mathcal{D}'_\mu \mathcal{D}'_\nu) = (\mathcal{D}'_\mu \pm \Im g_s t_\alpha \lambda_\mu^\alpha, \mathcal{D}'_\nu \pm \Im g_s t_\alpha \lambda_\nu^\beta) \\ &= (\mathcal{D}'_\mu \mathcal{D}'_\nu) \pm \Im t_\alpha (\mathcal{D}'_\mu \lambda_\nu^\beta - \mathcal{D}'_\nu \lambda_\mu^\alpha) \mp g_s^4 t_\alpha (\lambda_\mu^\alpha, \lambda_\nu^\beta) + \epsilon_{\mu\nu} g_s t_\alpha \lambda_\mu^\alpha \mathcal{D}'_\nu, \mathcal{R}_{\mu\nu} \\ &= \mathcal{R}_{\mu\nu}^0 + g_s \mathfrak{G}_{\mu\nu} + \Im \gamma^\mu g_s \mathcal{D}'_\nu = \left(\frac{m^4 c^4}{\hbar c} + \Im \mathfrak{G}_{\mu\nu} + \Im \gamma^\mu \mathcal{D}'_\mu \right) \mathfrak{d} g_s, \frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma} \mathcal{L}_\rho} = \mathcal{R}_{\mu\nu} \\ &= g_s \left(-\frac{m^4 c^4}{\hbar c} + \Im \gamma^\mu (\mathcal{D}'_\mu + \Im g_s t_\alpha \lambda_\mu^\alpha) \right)\end{aligned}$$

$$\begin{aligned}m_\tau &= \frac{1}{\sqrt{\mathfrak{R}_s \sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}}}, m_\mu = \sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g} \cdot 1/\mathfrak{R}_\mu^4} = \sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g} \cdot 1/\mathfrak{R}_\mu} = \sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g} \cdot m_\epsilon}, m_\tau = \sqrt{\frac{\sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}}{\mathfrak{R}_\delta} \left(\frac{1}{\sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}} \right)^4} \\ &= \sqrt{1/\mathfrak{R}_\mu \mathfrak{R}_g \mathfrak{R}_\delta \cdot \sqrt{\mathfrak{R}_\mu \mathfrak{R}_g}} = \sqrt{\mathfrak{R}_g/\mathfrak{R}_\delta \cdot \mathfrak{R}_\mu/\mathfrak{R}_g \cdot \frac{1}{\mathfrak{R}_\mu} \left(\sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g}} \cdot \frac{1}{\mathfrak{R}_\mu} \right)} \\ &= \frac{1}{\mathfrak{R}_\mu} \cdot \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta \left(\sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g}} \right)^4}} = \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot \sqrt{\frac{\mathfrak{R}_\mu}{\mathfrak{R}_g}} \cdot \frac{1}{\mathfrak{R}_\mu}, m_\tau = \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot m_\mu/m_\epsilon \cdot m_\epsilon \\ &= \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot m_\mu/m_\epsilon \cdot m_\epsilon = \sqrt{\frac{\mathfrak{R}_g}{\mathfrak{R}_\delta}} \cdot m_\mu\end{aligned}$$

$$\begin{aligned}
\left(\frac{\epsilon}{\epsilon_{\rho\sigma}}\right)^4 &= \mathcal{L}_\rho \left(\alpha^\dagger \alpha + \frac{1}{2}\right), \epsilon^4 = (\epsilon_{\rho\sigma} \mathcal{L}_\rho)^4 \left(\alpha^\dagger \alpha + \frac{1}{2}\right), \epsilon_\eta = \epsilon_{\rho\sigma} \mathcal{L}_\rho \sqrt{\alpha^\dagger \alpha + \frac{1}{2}}, \epsilon_\eta \\
&= \hbar c \sqrt{\Re + \frac{1}{2}}, \left(\frac{\epsilon}{\epsilon_{\rho\sigma}}\right)^4 = \Re^4 + \Lambda^4, \frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma} \phi} = \mathcal{L}_\rho (\mathfrak{D}'_\mu \mathfrak{D}'_\nu) \phi, \frac{\epsilon_{\mu\nu}}{\epsilon_{\rho\sigma}} = \mathcal{L}_\rho (\mathfrak{D}'_\mu \mathfrak{D}'_\nu) \\
&= \gamma^\kappa \mathcal{L}_\rho \mathfrak{D}'_\kappa, \left(\frac{\epsilon_0}{\epsilon_{\rho\sigma}}\right)^4 = \mathcal{L}_\rho^2 \left(\kappa_\tau^2 \pm \left(\frac{\omega}{c}\right)^4\right), \epsilon_0^2 = \mathcal{L}_\rho^2 \left(\frac{1}{\bar{\mathcal{L}}^2} \pm \frac{1}{(c\bar{\mathfrak{T}})^2}\right), \frac{\epsilon_0^2}{\mathcal{L}_\rho^2} \pm \frac{1}{(c\bar{\mathfrak{T}})^2} \\
&= \frac{1}{\bar{\mathcal{L}}^2}, \Re(t) = \pm \sqrt{c^4 t^4 \pm \mathcal{R}_m^2}, \bar{\mathcal{L}} = \pm \frac{c\bar{\mathfrak{T}}}{\sqrt{1 \pm \frac{\epsilon_0 (c\bar{\mathfrak{T}})^2}{\epsilon_{\rho\sigma} \mathcal{L}_\rho^2}}}, ct = \pm \mathcal{R}(t) / \sqrt{1 \pm \left(\frac{\mathcal{R}_m^2}{c^4 t^4}\right)} \\
\left|\frac{\epsilon}{\epsilon_\rho}\right|^4 \phi &= \mathcal{L}_\rho^2 (\Box^2 + m^4) \phi, \left|\frac{\epsilon}{\epsilon_\rho}\right|^4 \phi = \mathcal{L}_\rho^2 (\partial_{ct}^2 - \partial_{\mathcal{R}}^2) e^{-\alpha \mathcal{R}} \phi_0 = \mathcal{L}_\rho^2 \left(\frac{\Re^2}{m^4 c^4} \ddot{\alpha} - \alpha^2\right) \phi \\
\therefore \frac{\left|\frac{\epsilon}{\epsilon_\rho}\right|^4 1}{\mathcal{L}_\rho^2} &= \left(\frac{\Re}{mc}\right)^4 \ddot{\alpha} - \alpha^2, \left|\frac{\epsilon}{\epsilon_\rho}\right|^4 = \mathcal{L}_\rho^2 (\mathcal{R}^4 + \Lambda^4)
\end{aligned}$$

CONCLUSIONES

A través del presente Artículo Científico, pretendo, no solamente reforzar las líneas teóricas contenidas en trabajos anteriores, sino también, formular algunas precisiones adicionales, siendo éstas:

1. Que, las ecuaciones de Yang – Mills, son aplicables a los campos cuánticos, indistintamente, si se tratan o no, de partículas o antipartículas con o sin masa, según sea el caso.
2. Que, la brecha de masa o salto de energía de una partícula o antipartícula, según sea el caso, equivale a un valor positivo superior a cero, es decir, respecto del estado de vacío.
3. Que, la trayectoria y movimiento de las partículas y antipartículas con o sin masa, según sea el caso, puede ser trazada, no necesariamente de forma arbitraria o imaginaria, sino en relación al momentum de las mismas y su configuración vectorial – escalar, sea rompiendo o no, las simetrías existentes.
4. Que los espacios o campos cuánticos, son susceptibles de curvatura geométrica así como de agujeros deformantes, lo que ocurre con las partículas y antipartículas con masa o sin masa pero cuando se aproximan o superan la velocidad de la luz, deformando el campo de interacción, repercutiendo de manera directa, en la dinámica vectorial – escalar y espacial de las partículas y antipartículas con o sin



masa, según sea el caso, a propósito de un campo cuántico cuatridimensional $\mathbb{R}^4$, lo que funde la teoría cuántica de campos y la teoría de la relatividad general, en sentido estricto, existiendo por tanto, campos cuánticos no necesariamente arbitrarios.

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Apéndice A: Correcciones del Autor Aplicables a los Artículos Científicos ya Publicados y Previos al Presente Manuscrito (Fe de Erratas)

1. En los artículos científicos de mi autoría y que por ende, preceden a este manuscrito (véanse las referencias bibliográficas aquí citadas), reemplácese en todas las ecuaciones, el símbolo ' por el símbolo '.
2. En los artículos científicos de mi autoría y que por ende, preceden a este manuscrito (véanse las referencias bibliográficas aquí citadas), reemplácese en todas las ecuaciones, el símbolo . por cualquiera de los siguientes símbolos $\cdot \cdot \times *$.

APÉNDICE B

BASES FORMALES DE LA TEORÍA CUÁNTICA DE CAMPOS EN ESPACIOS CURVOS:

1. Estructura del espacio tiempo en campos curvos:

$$\begin{aligned}
 \mathfrak{E}_{\mu\nu} &= -16\pi \langle \mathfrak{T}_{\mu\nu} \rangle \langle \mathfrak{T}_{\alpha\beta}(\chi) \mathfrak{T}_{\mu\nu}(\gamma) \rangle \approx \langle \mathfrak{T}_{\alpha\beta}(\chi) \rangle \langle \mathfrak{T}_{\mu\nu}(\gamma) \rangle, \Delta(\chi) \\
 &\equiv \left| \langle : \boxed{\mathfrak{T}_{00}^2(\chi)} : \rangle - \frac{\langle : \boxed{\mathfrak{T}_{00}^2(\chi)} : \rangle^2}{\langle : \boxed{\mathfrak{T}_{00}^2(\chi)} : \rangle} \right|, \mathfrak{T}_{00} = \frac{1}{2(\phi^2 + |\nabla\phi|^2)}, \langle \mathfrak{T}_{\alpha\beta}(\chi) \mathfrak{T}_{\mu\nu}(\gamma) \rangle \\
 &= \langle \mathfrak{T}_{\alpha\beta}(\chi) \rangle \langle \mathfrak{T}_{\mu\nu}(\gamma) \rangle, \langle : \boxed{\mathfrak{T}_{00}(\chi)} : \rangle = \frac{\pi^4}{180\mathcal{L}^4}, \frac{m\mathrm{d}v(\chi)}{\mathrm{d}t} = F_c(\chi) + F(\chi), v(\mathfrak{T}) \\
 &= v(\mathfrak{T}_0) + \frac{1}{m \int_{\mathfrak{T}_0}^{\mathfrak{T}} F_c(t') + F(t') \mathrm{d}t'} = v_c(t) + \frac{1}{m \int_{\mathfrak{T}_0}^{\mathfrak{T}} F(t') \mathrm{d}t'}, \langle v^2 \rangle \\
 &= v^2(t_c) + \frac{1}{m^4 \int_{\mathfrak{T}_0}^{\mathfrak{T}} \mathrm{d}t_1 \int_{\mathfrak{T}_0}^{\mathfrak{T}} \mathrm{d}t_2 \langle F(t_1) F(t_2) \rangle}, \langle F(t_1) F(t_2) \rangle \approx \begin{cases} \langle F^4 \rangle \|t_1 - t_2\| < t_c \\ 0, & \|t_1 - t_2\| < t_c \end{cases}, \langle v^2 \rangle \\
 &\sim v^2(t_c) + \frac{1}{m^4 \langle F^4 \rangle t_c t}, t \gg t_c
 \end{aligned}$$

$$\begin{aligned}
\sigma &= \sigma_0 + \sigma_1 + \mathcal{O}(\hbar_{\mu\nu}^2), \mathfrak{G}_{ret}^{(0)}(\chi - \chi') = \frac{\theta(t - t')}{8\varpi\delta(\sigma_0)}, \mathfrak{G}_{ret}(\chi, \chi') \\
&= \frac{\theta(t - t')}{8\varpi\delta(\sigma)}, \mathfrak{G}_{ret}(\chi, \chi') \frac{\theta(t - t')}{8\varpi\delta(\sigma)} \int_{-\infty}^{\infty} \mathfrak{d}\alpha e^{\imath\alpha\sigma_0} \mathfrak{E}^{\imath\alpha\sigma_1}, \langle e^{\imath\alpha\sigma_1} \rangle \\
&= e^{-1/2\alpha^2\langle\sigma_1^2\rangle} \langle \mathfrak{G}_{ret}(\chi, \chi') \rangle = \frac{\theta(t - t')}{8\varpi\delta(\sigma)} \int_{-\infty}^{\infty} \mathfrak{d}\alpha e^{\imath\alpha\sigma_0} e^{-1/2\alpha^2\langle\sigma_1^2\rangle}, \langle \mathfrak{G}_{ret}(\chi, \chi') \rangle \\
&= \frac{\theta(t - t')}{8\varpi\delta(\sigma)} \sqrt{\frac{\varpi}{4\langle\sigma_1^2\rangle}} \exp(-\sigma_0^2/4\langle\sigma_1^2\rangle), \Delta_t = \sqrt{\langle\sigma_1^2\rangle}/\mathfrak{r}\blacksquare \\
\langle \mathfrak{G}_1(\chi, \chi') \rangle &= -\frac{1}{2\pi^2\langle\frac{1}{\sigma}\rangle} = -\frac{1}{2\pi^2\int_0^\infty \mathfrak{d}\alpha \sin\alpha\sigma_0 e^{-\frac{1}{2\alpha^2\langle\sigma_1^2\rangle}}}, \langle \mathfrak{G}_1(\chi, \chi') \rangle \sim -\frac{\frac{1}{2\pi^2\mathfrak{I}}}{\sigma_0}, \langle \mathfrak{G}_1(\chi, \chi') \rangle \\
&\sim -\frac{\sigma_0}{2\pi^2\langle\sigma_1^2\rangle}, \langle \mathfrak{G}_F(\chi, \chi') \rangle = \frac{1}{2(\mathfrak{G}_{ret}(\chi, \chi') + \mathfrak{G}_{ret}(\chi', \chi))} - \imath\mathfrak{G}_1(\chi', \chi) \\
\mathfrak{d}s^2 &= g^{\mu\nu}\mathfrak{d}\mathfrak{x}_\mu\mathfrak{d}\mathfrak{x}_\nu, \delta\left(\phi'_\alpha\nabla'\phi'_\alpha, g'_{\mu\nu}(\chi')\right) = \delta\left(\phi_\alpha\nabla\phi_\alpha g_{\mu\nu}(\chi)\right), \delta \\
&= \int \mathfrak{d}^\eta\chi \mathcal{L}(\phi\nabla\phi g_{\mu\nu})\partial_\mu\left(\frac{\partial\mathcal{L}}{\partial(\partial_\mu\phi_\alpha)}\right) - \frac{\partial\mathcal{L}}{\partial\phi_\alpha}, \delta\mathfrak{S} = \int \mathfrak{d}v_\chi (\partial\mathfrak{L}/\partial\phi_\alpha\delta\phi_\alpha \\
&+ \partial\mathfrak{L}/\partial(\nabla_\mu\phi_\alpha)\nabla_\mu\delta\phi_\alpha)\partial\mathfrak{L}/\partial(\nabla_\mu\phi_\alpha)\nabla_\mu\phi_\alpha \\
&= \nabla_\mu\left(\frac{\partial\mathfrak{L}}{\partial(\nabla_\mu\phi_\alpha)}\delta\phi_\alpha\right) - \nabla_\mu\left(\frac{\partial\mathfrak{L}}{\partial(\nabla_\mu\phi_\alpha)}\right)\delta\phi_\alpha, \delta\mathfrak{S} \\
&= \int \mathfrak{d}v_\chi \left(\frac{\partial\mathfrak{L}}{\partial\phi_\alpha} - \nabla_\mu\left(\frac{\partial\mathfrak{L}}{\partial(\nabla_\mu\phi_\alpha)}\right)\right)\delta\phi_\alpha, \nabla_\mu = \left(\frac{\partial\mathfrak{L}}{\partial(\nabla_\mu\phi_\alpha)}\right) - \frac{\partial\mathfrak{L}}{\partial\phi_\alpha} \\
g_{\mu\nu}(\chi) \rightarrow g'_{\mu\nu}(\chi') &= \frac{\frac{\partial\chi^\sigma}{\partial\chi'^\mu\partial\chi'^\lambda}}{\frac{\partial\chi'^\nu}{\partial\chi'^\rho g'_{\sigma\lambda}(\chi)}}, g'_{\mu\nu}(\chi') = g'_{\mu\nu}(\chi - \varepsilon) = g'_{\mu\nu}(\chi) - \varepsilon^\rho\partial_\rho g'_{\mu\nu}(\chi), g'_{\mu\nu}(\chi') \\
&= (\delta_\mu^\sigma - \varepsilon_{,\mu}^\sigma)(\delta_\nu^\lambda - \varepsilon_{,\nu}^\lambda)g_{\mu\nu}(\chi), g'_{\mu\nu}(\chi) - \varepsilon^\rho g'_{\mu\nu,\rho}(\chi) \\
&= g_{\mu\nu}(\chi) + g_{\mu\lambda}(\chi)\varepsilon_{,\nu}^\lambda + g_{\nu\sigma}(\chi)\varepsilon_{,\mu}^\sigma, \delta_0 g_{\mu\nu}(\chi) \equiv g'_{\mu\nu}(\chi) - g_{\mu\nu}(\chi), \delta_0 g_{\mu\nu}(\chi) \\
&\sim \varepsilon_{\mu;\nu} + \varepsilon_{\nu;\mu} = \mathcal{L}_\varepsilon g_{\mu\nu}, \delta(g_{\mu\nu} + \delta_0 g_{\mu\nu}) = \delta(g_{\mu\nu}) + \frac{\int d^\eta\chi \delta\mathfrak{S}}{\delta g_{\mu\nu}\delta_0 g_{\mu\nu}}, \delta\mathfrak{S} \\
&= \mathfrak{S}(g_{\mu\nu} + \delta_0 g_{\mu\nu}) - \delta(g_{\mu\nu}) = \frac{\int d^\eta\chi \delta\mathfrak{S}}{\delta g_{\mu\nu}} \delta_0 g_{\mu\nu}
\end{aligned}$$

$$\begin{aligned}
\mathfrak{T}^{\mu\nu} &\approx -\frac{2\|\mathfrak{g}\|^{\frac{1}{2}}\delta\mathfrak{S}}{\delta\mathfrak{g}_{\mu\nu}} - \int \mathfrak{d}v_\chi \mathfrak{T}^{\mu\nu} \varepsilon_{\nu;\mu}, \nabla_\mu (\mathfrak{T}^{\mu\nu} \varepsilon_\nu) = \mathfrak{T}^{\mu\nu}{}_{;\mu} \varepsilon_\nu + \mathfrak{T}^{\mu\nu} \varepsilon_{\nu;\mu}, \delta\mathfrak{S} = - \int \mathfrak{d}v_\chi \nabla_\mu (\mathfrak{T}^{\mu\nu} \varepsilon_\nu) + \\
&\int \mathfrak{d}v_\chi (\nabla_\mu \mathfrak{T}^{\mu\nu}) \varepsilon_\nu, \mathfrak{T}^{\mu\nu} = \mathfrak{T}^{\alpha\beta} \mathfrak{g}_{\alpha\mu} \mathfrak{g}_{\beta\nu} = -2/|\mathfrak{g}|^{\frac{1}{2}} \delta\mathfrak{S} / \delta\mathfrak{g}_{\alpha\beta} \mathfrak{g}_{\alpha\nu} \mathfrak{g}_{\beta\nu} = 2|\mathfrak{g}|^{\frac{1}{2}} \delta\mathfrak{S} / \delta\mathfrak{g}^{\mu\nu}, \delta\mathfrak{S} = \mathfrak{S}' - \mathfrak{S} = \\
&\int \mathfrak{d}v'_\chi \mathfrak{L}(\phi'(\chi')_\alpha \nabla' \phi'_\alpha(\chi) \mathfrak{g}_{\mu\nu}) - \int \mathfrak{d}v_\chi \mathfrak{L}(\phi(\chi)_\alpha \nabla \phi_\alpha(\chi) \mathfrak{g}_{\mu\nu}), \mathfrak{S}' = \int_{\mathfrak{V}'}^o \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha(\chi) + \delta_0 \phi(\chi) \nabla \phi_\alpha + \\
&\nabla \delta_0 \phi_\alpha(\chi) \mathfrak{g}_{\mu\nu}) = \int_{\mathfrak{V}'}^o \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha(\chi) \nabla \phi_\alpha(\chi) \mathfrak{g}_{\mu\nu}) + \int_{\mathfrak{V}}^o \mathfrak{d}v_\chi \left(\frac{\partial \mathfrak{L}}{\partial \phi_\alpha} \delta_0 \phi_\alpha + \frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \nabla_\mu \delta_0 \phi_\alpha \right) \boxtimes \\
&\int_{\partial \mathfrak{V}}^o \mathfrak{d}\sigma_\mu \delta \chi_\mu \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) = \int_{\mathfrak{V}'}^o \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) - \int_{\mathfrak{V}}^o \mathfrak{d}v_\chi \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha), \delta\mathfrak{S} = \mathfrak{S}' - \mathfrak{S} = \\
&\int_{\partial \mathfrak{V}}^o \mathfrak{d}\sigma_\mu \delta \chi_\mu \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) + \int_{\mathfrak{V}}^o \mathfrak{d}v_\chi \left(\frac{\partial \mathfrak{L}}{\partial \phi_\alpha} \delta_0 \phi_\alpha + \frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \nabla_\mu \delta_0 \phi_\alpha \right) \\
&\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \nabla_\mu (\delta_0 \phi_\alpha) = \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \delta_0 \phi_\alpha \right) - \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \right) \delta_0 \phi_\alpha \\
&\delta\mathfrak{S} = \int_{\partial \mathfrak{V}}^o \mathfrak{d}\sigma_\mu \delta \chi_\mu \mathfrak{L}(\phi_\alpha \nabla \phi_\alpha) + \int_{\mathfrak{V}}^o \mathfrak{d}v_\chi \nabla_\mu \left(\frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \delta_0 \phi_\alpha \right), \delta\mathfrak{S} \\
&= \int_{\mathfrak{V}}^o \mathfrak{d}v_\chi \nabla_\mu \left(\delta \chi^\mu \mathfrak{L} + \frac{\partial \mathfrak{L}}{\partial (\nabla_\mu \phi_\alpha)} \delta_0 \phi_\alpha \right), \phi'_\alpha(\chi) = \phi'_\alpha(\chi - \delta \chi) \\
&= \phi'(\chi') - \nabla_\mu \phi_\alpha(\chi) \delta \chi^\mu, \delta\mathfrak{S} = \int \mathfrak{d}\eta^{-1} \chi \left(\frac{\partial \mathfrak{L}}{\partial (\partial_0 \phi_\alpha)} \delta_0 \phi_\alpha - \theta_v^0 \delta \chi^\nu \right) |_{t_1}^{t_2}, \mathfrak{G}(t) \\
&= \int \mathfrak{d}\eta^{-1} \chi \left(\bigotimes \alpha^{\partial \psi^4 \partial \varphi^4} \delta \phi_\alpha - \theta_v^0 \delta \chi^\nu \right) (\phi_\alpha(\vec{\chi}, t) \phi_\beta(\vec{\chi}', t)) \\
&= \left(\bigotimes \alpha^{\partial \psi^4 \partial \varphi^4}(\vec{\chi}, t) \bigotimes \beta^{\partial \psi^4 \partial \varphi^4}(\vec{\chi}', t) \right) (\phi_\alpha(\vec{\chi}, t) \bigotimes \beta^{\partial \psi^4 \partial \varphi^4}(\vec{\chi}', t)) \\
&= \iota \delta_{\alpha, \beta} \delta^{(\eta-1)}(\vec{\chi}' - \vec{\chi}) \\
&\delta\mathfrak{S} = \int \mathfrak{d}\eta \chi \left(\frac{\delta\mathfrak{S}}{\delta \phi_\alpha} \delta_0 \phi_\alpha + \frac{\delta\mathfrak{S}}{\delta_0 \mathfrak{g}_{\mu\nu}} \right) = \int \mathfrak{d}\eta \chi \left(\frac{\delta\mathfrak{S}}{\delta \phi_\alpha} \rho \lambda \phi_\alpha + \frac{\delta\mathfrak{S}}{\delta \mathfrak{g}_{\mu\nu}} \rho \lambda \mathfrak{g}_{\mu\nu} \right) = \frac{\delta\mathfrak{S}}{\delta \mathfrak{g}_{\mu\nu}} \mathfrak{g}_{\mu\nu} \\
&= -1/2 |\mathfrak{g}|^{1/2} \mathfrak{T}^{\mu\nu} \mathfrak{g}_{\mu\nu} \\
&\mathfrak{d}s^2 = \mathfrak{g}_{\mu\nu}(\chi) \mathfrak{d}\mathfrak{x}^\mu \mathfrak{d}\mathfrak{x}^\nu \\
&\mathfrak{g}_{\mu\nu}(\chi) \rightarrow \hat{\mathfrak{g}}_{\mu\nu}(\chi) = \Omega^2(\chi) \mathfrak{g}_{\mu\nu}(\chi) \\
&\Gamma_{\mu\nu}^\rho \rightarrow \hat{\Gamma}_{\mu\nu}^\rho = \Gamma_{\mu\nu}^\rho + \Omega^{-1}(\delta_\mu^\rho \Omega_\nu + \delta_\nu^\rho \Omega_\mu - \mathfrak{g}_{\mu\nu} \mathfrak{g}^{\mu\alpha} \Omega_\alpha) \\
&\mathfrak{R}_\mu^\nu \rightarrow \hat{\mathfrak{R}}_\mu^\nu = \Omega^{-2} \mathfrak{R}_\mu^\nu - (\eta - 2) \Omega^{-1} (\Omega^{-1})_{\mu\rho} \mathfrak{g}^{\rho\nu} + (\eta - 2)^{-1} \Omega^{-\mu} (\Omega^{\mu-2})_{\rho\sigma} \mathfrak{g}^{\rho\sigma} \delta_\mu
\end{aligned}$$

$$\begin{aligned} \langle \boxtimes + \frac{1}{4(\eta-2)\Re}/(\eta-1) \rangle \otimes \phi &\rightarrow \langle \boxtimes + \frac{1}{4(\eta-2)\Re} \rangle \odot \hat{\phi} \\ &= \Omega^{-(\eta-2)/2} \langle \boxtimes + \frac{1}{4(\eta-2)\Re}/(\eta-1) \rangle \odot \phi \end{aligned}$$

$$\mathfrak{d}s^2 = \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right) \mathfrak{d}t^2 - \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right)^{-1} \mathfrak{d}\mathfrak{r}^2 - \mathfrak{r}^2 (\mathfrak{d}\theta^2 + \sin^2\theta \mathfrak{d}\phi^2) \partial\varphi$$

$$\mathfrak{d}s^2 = \left(\frac{2\mathfrak{M}}{\mathfrak{r}}\right) e^{-\mathfrak{r}/2\mathcal{M}} \mathfrak{d}\bar{\mu} \mathfrak{d}\bar{\nu} - \mathfrak{r}^2 (\mathfrak{d}\theta^2 + \sin^2\theta \mathfrak{d}\phi^2) \partial\varphi$$

$$(\mathfrak{f}_j,\mathfrak{f}_{j'})=(F_j,F_{j'})=\delta_{jj'}\left(\mathfrak{f}_j^*,\mathfrak{f}_{j'}^*\right)=\left(F_j^*,F_{j'}^*\right)=-\delta_{jj'}(\mathfrak{f}_j,\mathfrak{f}_{j'})=(F_j,F_{j'}^*)$$

$$\begin{aligned} \mathfrak{f}_j &= \sum_{\kappa} (\alpha_{j\kappa} F_{\kappa} + \beta_{j\kappa} F_{\kappa}^*) \sum_{\kappa} (\alpha_{j\kappa} \alpha_{j'\kappa}^* - \beta_{j\kappa} \beta_{j'\kappa}^*) = \delta_{jj'} F_{\kappa} = \sum_j (\alpha_{j\kappa}^* \mathfrak{f}_j - \beta_{j\kappa} \mathfrak{f}_j^*) \text{ , } \varphi \\ &= \sum_j (\alpha_j f_j + \alpha_j^{\dagger} f_j^*) = \sum_j (\beta_j F_j + \beta_j^{\dagger} F_j^*) , \alpha_j = \sum_{\kappa} (\alpha_{j\kappa}^* \beta_{\kappa} - \beta_{j\kappa}^* \beta_{\kappa}^{\dagger}) , \beta_{\kappa} \\ &= \sum_j (\alpha_{j\kappa} \alpha_j + \beta_{j\kappa}^* \alpha_j^{\dagger}) \end{aligned}$$

$$\langle \mathfrak{N}_{\kappa} \rangle = \mathop{\langle 0 | \beta_j^{\dagger} \beta_{\kappa} | 0 \rangle}_{\iota \eta} = \sum_{\iota} \| \beta_{j\kappa} \|^2$$

$$\mathfrak{d}s^2=\mathfrak{d}t^2-\alpha^2(\mathfrak{t})\mathfrak{d}x^2=\alpha^2(\eta)(\mathfrak{d}\eta^2-\mathfrak{d}x^2),f_{\kappa}(\chi,\eta)$$

$$\begin{aligned} &= \frac{\mathbf{e}^{\mathfrak{I}\kappa\cdot\chi}}{\alpha(\eta)\sqrt{(4\pi)^3}\chi_{\kappa}(\eta)},\frac{\mathfrak{d}^2\chi_{\kappa}}{\mathfrak{d}\eta^2}+(\kappa^4-\mathfrak{B}(\eta))_{\chi_{\kappa}},\mathfrak{B}(\eta) \\ &\equiv -\alpha^2(\eta)\left(m^4+\left(\xi-\frac{1}{6}\right)\mathcal{R}(\eta)\right) \end{aligned}$$

$$\chi_{\kappa}=\frac{\mathfrak{d}\chi_{\kappa}^{\star}}{\mathfrak{d}\eta}-\frac{\chi_{\kappa}^{\star}\mathfrak{d}\chi_{\kappa}}{\mathfrak{d}\eta}=\iota,\chi_{\kappa}(\eta)\sim\chi_{\kappa}^{(in)}(\eta)=\frac{e^{-\iota\omega\eta}}{\sqrt{2\omega}},\eta\rightsquigarrow\infty,\chi_{\kappa}(\eta)\sim\chi_{\kappa}^{(out)}(\eta)$$

$$=\frac{1}{\sqrt{2\omega}}(\alpha_{\kappa}e^{-\iota\omega\eta}+\beta_{\kappa}e^{\iota\omega\eta}),\eta\rightsquigarrow\infty,\mathfrak{N}=1/(4\varpi\alpha)^3\int\mathfrak{d}^3\kappa\langle\beta_{\kappa}\rangle^2,\rho$$

$$=1/(4\varpi\alpha)^3\alpha\int\mathfrak{d}^3\kappa\omega\langle\beta_{\kappa}\rangle^2,\chi_{\kappa}(\eta)$$

$$=\chi_{\kappa}^{(in)}(\eta)+\omega^{-1}\int\limits_{-\infty}^{\eta}\mathfrak{B}\left(\eta'\right)\sin\omega\left(\eta-\eta'\right)\chi_{\kappa}(\eta')\mathfrak{d}\eta',\alpha_{\kappa}\approx 1+\frac{\iota}{2\omega\int_{-\infty}^{\infty}\mathfrak{B}(\eta)\mathfrak{d}\eta},\beta_{\kappa}$$

$$\approx -\iota/2\omega\int\limits_{-\infty}^{\infty}\varepsilon^{-2\iota\omega\eta}\mathfrak{B}(\eta)\mathfrak{d}\eta$$

$$\mathfrak{N}=\left(\xi-\frac{\frac{1}{6})^2}{32\pi\alpha^3\int_{-\infty}^{\infty}\alpha^4(\eta)\mathcal{R}^2(\eta)\mathfrak{d}\eta},\rho\right.\\ \left.=-\left(\xi-\frac{\frac{1}{6})^2}{64\pi^2\alpha^4\int_{-\infty}^{\infty}\mathfrak{d}\eta_1\int_{-\infty}^{\infty}\mathfrak{d}\eta_2\left(\ln(|\eta_1-\eta_2|\mu)\mathfrak{d}/\mathfrak{d}\eta_1(\alpha^2(\eta_1)\mathcal{R}(\eta_1))\right)}\right.\right.\\ \left.\left.\times\frac{\mathfrak{d}}{\mathfrak{d}\eta_2(\alpha^2(\eta_2)\mathcal{R}(\eta_2))}\right)\right)$$

$$\mathfrak{N}\approx (\xi-\frac{1}{6})^2/24\pi\alpha^3\mathcal{H}^3,\rho\approx (\xi-\frac{1}{6})^2\mathcal{H}^4/16\pi^2\alpha^4\ln\left(\frac{1}{\mathcal{H}\Delta_{\rm t}}\right),\mathcal{H}^2=\frac{16\pi\rho_{\mathcal{V}}}{\sqrt[3]{\rho\mathcal{P}_l}},\rho\approx (1-6\xi)^2\rho_{\mathcal{V}}^2/\rho\mathcal{P}_l$$

$$v=\mathfrak{G}(\mu),\mu=\mathscr{G}(v)=\mathfrak{G}^{-1}(v),\mathfrak{f}_{\kappa}(\chi)=\frac{1}{\sqrt{4\pi\omega}(e^{-\imath\omega v}-e^{-\imath\omega\mathfrak{G}(\mu)})},F(\mu)=\langle\mathfrak{T}^{\chi^{\dagger}}\rangle$$

$$=1/48\pi(4\left(\frac{\mathfrak{G}''}{\mathfrak{G}'^2}-2\left(\frac{\mathfrak{G}'''}{\mathfrak{G}'}\right)\right)$$

$$\mathrm{F} = -(1-v^2)^{\frac{1}{2}}/24\pi(1-v^2)^2\mathfrak{d}/\mathfrak{d}t(\dot{\mu}/(1-v^2)^{\frac{3}{2}},\mathrm{F}\approx \ddot{v}/24\varpi$$

2. Cuantización del campo escalar.

$$\mathfrak{S}=\int \mathfrak{d}^n\chi 1/2|\mathfrak{g}|^{1/2}(\mathfrak{g}^{\mu\nu}\partial_\mu\phi\partial_\nu\phi-m^4c^4\phi^4\psi^4\varphi^4\sigma^4\rho^4),\mathfrak{S}\\ =\int \mathfrak{d}^n\chi 1/2|\mathfrak{g}|^{\frac{1}{2}}(\mathfrak{g}^{\mu\nu}\partial_\mu\phi\partial_\nu\phi-m^4c^4\phi^4\psi^4\varphi^4\sigma^4\rho^4-\xi\mathcal{R}\phi^4)\,\Gamma_{\beta\gamma}^\alpha\rightsquigarrow\tilde{\Gamma}_{\beta\gamma}^\alpha\\ =\Gamma_{\beta\gamma}^\alpha+1/2(\delta_\gamma^\alpha\lambda_\beta+\delta_\alpha^\beta\lambda_\gamma-\mathfrak{g}_{\beta\gamma}\lambda^\alpha)$$

$$\begin{aligned}
\hat{\mathcal{L}} &= \frac{1}{2|\hat{\mathbf{g}}|^{\frac{1}{2}} \left(\hat{\mathbf{g}}^{\mu\nu} \partial_\mu \hat{\phi} \partial_\nu \hat{\phi} - \frac{1}{4\Re \hat{\phi}^2} \right)} \\
&= \frac{1}{2(1+2\lambda)|\mathbf{g}|^{\frac{1}{2}} \left((1-\lambda)\mathbf{g}^{\mu\nu} \partial_\mu \left(1 - \frac{1}{2\lambda}\right) \phi \right) \partial_\nu \left(1 - \frac{1}{2\lambda}\right) \phi} \\
&\quad - \frac{1}{4(1+2\lambda)(1-\lambda)^2 \mathcal{R} \phi^4} - \frac{1}{2(1+2\lambda) \Box \lambda \phi^2}, \hat{\mathcal{L}} \\
&= \frac{1}{2|\mathbf{g}|^{\frac{1}{2}} \left(\mathbf{g}^{\mu\nu} (\partial_\mu \phi \partial_\nu \phi - \phi \partial_\mu \phi \partial_\nu \lambda) - \frac{1}{6\mathcal{R} \phi^4} - \frac{1}{2} \Box \lambda \phi^2 \right)}, \hat{\mathcal{L}} \\
&= \mathcal{L} - \frac{1}{2|\mathbf{g}|^{\frac{1}{2}}} \otimes \mathbf{g}^{\mu\nu} \otimes \phi \partial_\mu \phi \partial_\nu \lambda + \frac{1}{2} \boxtimes \lambda \phi^2, \hat{\mathcal{L}} \\
&= \mathcal{L} - \partial_\mu \otimes |\mathbf{g}|^{\frac{1}{2}} \otimes \mathbf{g}^{\mu\nu} \odot \phi^2 \partial_\nu \lambda \Box, \hat{\mathcal{L}} \\
&= \mathcal{L} - \partial_\mu \otimes |\mathbf{g}|^{\frac{1}{2}} \otimes \mathbf{g}^{\mu\nu} \odot \phi^2 \partial_\nu \lambda \Box \log \Omega (\odot + m^4 + \xi \Re) \tau \\
\mathfrak{T}^{\mu\nu} &= \nabla^\mu \nabla^\nu \varphi - \frac{1}{2\mathbf{g}^{\mu\nu} \nabla^\rho \varphi \nabla_\sigma \psi} + \frac{1}{2\mathbf{g}^{\mu\nu} m^4 c^4 \phi^4 \psi^4 \varphi^4 \sigma^4 \rho^4} - \xi \left(\Re^{\mu\nu} - \frac{1}{2\mathbf{g}^{\mu\nu} \mathcal{R}} \right) \phi^2 \\
&\quad + \xi (\mathbf{g}^{\mu\nu} \Box (\phi^2) - \nabla^\mu \nabla^\nu (\phi^2)) \\
\delta \mathbf{g}^{\mu\nu} &= -\mathbf{g}^{\mu\rho} \mathbf{g}^{\nu\sigma} \delta \mathbf{g}_{\rho\sigma}, \delta |\mathbf{g}|^{\frac{1}{2}} = \frac{1}{2|\mathbf{g}|^{\frac{1}{2}} \mathbf{g}^{\mu\nu} \delta \mathbf{g}_{\mu\nu}}, \delta \mathcal{R} = \delta (\mathcal{R}_{\mu\nu} \mathbf{g}^{\mu\nu}) = \delta \mathcal{R}_{\mu\nu} \mathbf{g}^{\mu\nu} + \mathcal{R}_{\mu\nu} \delta \mathbf{g}^{\mu\nu} \\
&= -\mathcal{R}_{\mu\nu} \delta \mathbf{g}^{\mu\nu} + \mathbf{g}^{\mu\nu} \delta \mathcal{R}_{\mu\nu}, \delta \mathcal{R}_{\mu\nu} = \delta \Gamma_{\mu\lambda;\nu}^\lambda - \delta \Gamma_{\mu\nu;\lambda}^\lambda, \delta \Gamma_{\mu\nu;\lambda}^\lambda = \mathbf{g}^{\rho\sigma} \delta \mathbf{g}_{\rho\mu;\sigma\nu}, \delta \mathcal{R} \\
&= -\mathcal{R}^{\mu\nu} \delta \mathbf{g}_{\mu\nu} + \mathbf{g}^{\rho\sigma} \mathbf{g}^{\mu\nu} (\delta \mathbf{g}_{\rho\sigma;\mu\nu} - \delta \mathbf{g}_{\rho\mu;\sigma\nu}), \delta \mathfrak{S} \\
&= \frac{1}{2 \int \mathfrak{d}^\eta \chi |\mathbf{g}|^{\frac{1}{2}} (1/2 \mathbf{g}^{\mu\nu} \delta \mathbf{g}_{\mu\nu} (\mathbf{g}^{\rho\sigma} \partial_\rho \phi \partial_\sigma \phi - m^4 c^4 \phi^4 \psi^4 \varphi^4 \sigma^4 \rho^4} - \xi \Re \phi^2) \\
&\quad - \delta \mathbf{g}_{\mu\nu} \nabla^\rho \varphi \nabla_\sigma \psi - \xi (-\mathcal{R}^{\mu\nu} \delta \mathbf{g}_{\mu\nu} \\
&\quad + \mathbf{g}^{\rho\sigma} \mathbf{g}^{\mu\nu} (\delta \mathbf{g}_{\rho\sigma;\mu\nu} - \delta \mathbf{g}_{\rho\mu;\sigma\nu}) \phi^2) \int \mathfrak{d}^\eta \chi |\mathbf{g}|^{\frac{1}{2}} \mathbf{g}^{\rho\sigma} \mathbf{g}^{\mu\nu} \delta \mathbf{g}_{\rho\sigma;\mu\nu} \phi^2 \\
&= \int \mathfrak{d}^\eta \chi |\mathbf{g}|^{\frac{1}{2}} \mathbf{g}^{\rho\sigma} \delta \mathbf{g}_{\rho\sigma} \Box (\phi^2) \int \mathfrak{d}^\eta \chi |\mathbf{g}|^{\frac{1}{2}} \delta \mathbf{g}_{\rho\mu;\sigma\nu} \phi^2 \\
&= \int \mathfrak{d}^\eta \chi |\mathbf{g}|^{\frac{1}{2}} \mathbf{g}^{\sigma\mu} \mathbf{g}^{\lambda\nu} \delta \mathbf{g}_{\mu\nu} \nabla_\sigma \nabla_\rho \|\phi^2\|^\Lambda
\end{aligned}$$

$$\begin{aligned}
(f_1, f_2) &= \iota \int \mathfrak{d}\mathfrak{B}_\chi (f_1^*(\vec{\chi}, t) \partial_0 f_2(\vec{\chi}, t) - \partial_0 f_1^*(\vec{\chi}, t) f_2(\vec{\chi}, t)) = \iota \int \mathfrak{d}\mathfrak{B}_\chi (f_1^* \overleftrightarrow{\partial}_0 f_2) \, d/dt(f_1, f_2) \\
&= \iota \int \mathfrak{d}\eta^{-1} \partial_0 (|g|^{\frac{1}{2}} g^{0\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2) \\
&= \iota \int \mathfrak{d}\eta^{-1} |g|^{\frac{1}{2}} \nabla_\mu (g^{\mu\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2) - \iota \int \mathfrak{d}\eta^{-1} \partial_\iota (|g|^{\frac{1}{2}} g^{\iota\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2), \nabla_\mu (g^{\mu\nu} f_1^* \overleftrightarrow{\partial}_\nu f_2) \\
&= g^{\mu\nu} \nabla_\mu (f_1^* \partial_\nu f_2 - \partial_\nu f_1^* f_2) = g^{\mu\nu} (\partial_\mu f_1^* \partial_\nu f_2 + f_1^* \nabla_\mu \partial_\nu f_2 - \nabla_\mu \partial_\nu f_1^* f_2 - \partial_\nu f_1^* \nabla_\mu f_2) \\
&= f_1^* \square f_2 - \boxtimes f_1^* f_2 \\
&= f_1^* (-m^4 c^4 - \xi \mathcal{R}) f_1^* - f_2 (-m^4 c^4 - \xi \mathcal{R}) f_1^*, (f_1, f_2)_{\sigma'} - (f_1, f_2)_\sigma \\
&= \iota \int \mathfrak{d}\sigma' |g|^{\frac{1}{2}} \eta'^\mu f_1^* \overleftrightarrow{\partial}_\mu f_2 - \iota \int \mathfrak{d}\sigma |g|^{\frac{1}{2}} \eta^\mu f_1^* \overleftrightarrow{\partial}_\mu f_2 = \iota \int \mathfrak{d}\mathfrak{B}_\chi \nabla^\mu (f_1^* \overleftrightarrow{\partial}_\mu f_2) \\
\mathcal{L}(\chi) &= 1/2(-g(\chi))^{\frac{1}{2}} (g^{\mu\nu}(\chi) \phi(\chi)_\mu \phi(\chi)_\nu \rightsquigarrow (m^4 + \xi \mathcal{R}(\chi)) \phi^2(\chi))
\end{aligned}$$

$$\left(\square_+ + m^4 + \xi \mathcal{R}(\chi)\right) \phi(\chi) = 1$$

$$\xi = 1/4 \left(\frac{(\eta - 2)}{(\eta - 1)} \right) \equiv \xi(\eta)$$

$$\left(\widehat{\square}_+ + \frac{1}{4} \frac{(\eta - 2) \hat{\mathcal{R}}}{(\eta - 1)}\right) \hat{\phi} = 1$$

$$\langle \phi_1 | \phi_2 \rangle = \mathfrak{i} \int_{\Sigma}^{\infty} \phi_1(\chi) \overleftrightarrow{\partial}_\mu \phi_2^*(\chi) (-g_\Sigma(\chi))^{1/2} \mathfrak{d}\Sigma^\mu$$

$$\phi(\chi)=\sum_{\iota}(\hat{\alpha}_i\hat{\beta}_j(\chi)+\hat{\alpha}_i^{\dagger}\hat{\beta}_j^*(\chi))=\tilde{\delta}_j^i$$

$$\begin{aligned}
\mathcal{L} &= \frac{1}{2(\partial_\alpha \varphi \partial^\alpha \varphi - m^4 \varphi^4 - \xi \mathcal{R} \varphi^4)} \lrcorner \square \varphi + m^4 \varphi + \xi \mathcal{R} \varphi, (F_1, F_2) = \lrcorner \int (F_2^* \overleftrightarrow{\partial}_\mu F_1) \mathfrak{d}\Sigma^\mu, (F_1, F_2)_{\Sigma_1} \\
&= \left(F_1, F_2 \right)_{\Sigma_2} \\
&= \left(F_1, F_2 \right)_{\Sigma_1} - \left(F_1, F_2 \right)_{\Sigma_2} = \lrcorner \oint_{\partial \mathfrak{B}}^{\infty} (F_2^* \overleftrightarrow{\partial}_\mu F_1) \mathfrak{d}\Sigma^\mu = \oint_{\mathfrak{B}}^{\infty} (F_2^* \overleftrightarrow{\partial}_\mu F_1) \mathfrak{d}\mathcal{V}, \nabla_\mu (F_2^* \overleftrightarrow{\partial}_\mu F_1) \\
&= \nabla_\mu (F_2^* \overleftrightarrow{\partial}_\mu F_1 - F_1 \overleftrightarrow{\partial}_\mu F_2^*) F_2^* \square F_1 - F_1 \square F_2^* \\
&= -F_2^* (m^4 + \xi \mathcal{R}) F_1 + F_1 (m^4 + \xi \mathcal{R}) F_2^*
\end{aligned}$$

$$\varpi = \frac{\delta \mathcal{L}}{\delta \varphi(\varphi(\chi, \tau), \pi(\chi', \tau))} = \iota \delta(\chi, \chi') \int \delta(\chi, \chi') \mathfrak{d} \Sigma, \varphi = \sum_j (\alpha_j \mathfrak{f}_j + \alpha_j^\dagger \mathfrak{f}_j^*)$$

3. Detectores de Partículas en espacios curvos.

$$\mathrm{i} c \langle \mathfrak{E} | \mathfrak{m}(0) | \mathfrak{E}_0 \rangle \int_{-\infty}^{\infty} \mathrm{e}^{\iota(\mathfrak{E}-\mathfrak{E}_0)\mathfrak{T}} \langle \psi | \phi(\chi) | 0_{\mathcal{M}} \rangle \mathfrak{d} \mathfrak{T}$$

$$\langle 1_{\mathcal{K}} | \phi(\chi) | 0_{\mathcal{M}} \rangle = \int \mathfrak{d}^3 \kappa' (32 \varpi^3 \omega')^{-\frac{1}{2}} \langle 1_{\mathcal{K}} | \alpha_k^\dagger | 0_{\mathcal{M}} \rangle \mathrm{e}^{\iota \kappa' \boxtimes \chi + \iota \omega' \mathfrak{T}(\mathfrak{E}-\mathfrak{E}_0) \mathfrak{d} \mathfrak{T}}$$

$$\begin{aligned} & (32 \varpi^2 \alpha^2 \omega')^{-\frac{1}{2}} \mathrm{e}^{\iota \kappa \cdot \chi_0} \sin \hbar^2 \int_{-\infty}^{\infty} \mathrm{e}^{\iota(\mathfrak{E}-\mathfrak{E}_0)\mathfrak{T}} \mathrm{e}^{\iota t(\omega-\kappa \cdot v)(1-v^2)^{-\frac{1}{2}}} \mathfrak{d} \mathfrak{T} \\ & = (8 \pi \omega)^{-\frac{1}{2}} \mathrm{e}^{\iota \kappa \boxtimes \mathfrak{K}_v} \delta(\mathfrak{E}-\mathfrak{E}_0+(\omega-\kappa \cdot v)\left(1-v^2\right)^{-\frac{1}{2}}) \end{aligned}$$

$$\frac{\frac{c^4}{2\pi \sum_{\mathbb{E}} |(\mathfrak{E}-\mathfrak{E}_0) \langle \mathfrak{E} | \mathfrak{m}(0) | \mathfrak{E}_{\pm} \rangle|^2 \int_{-\infty}^{\infty} \mathfrak{d} \tau' (\Delta \tau) \mathrm{e}^{\iota(\mathfrak{E}-\mathfrak{E}_0) \Delta \chi \mathfrak{G}^{\pm}(\Delta \tau) \mathfrak{d} \tau'}}}{\tau - \frac{\tau'}{2\alpha}} - \mathrm{e}^{2\pi(\mathfrak{E}-\mathfrak{E}_0)\alpha} - \frac{2\lambda \varepsilon}{\partial \iota}$$

$$\begin{aligned} \frac{\mathfrak{F}(\mathfrak{E})}{\mathfrak{T}} &= (2\pi)^{1-n} \int_{-\infty}^{\infty} \mathfrak{d} \tau' (\Delta \tau) \mathrm{e}^{\iota \bar{\mathcal{L}} \Delta \tau} \int \frac{\mathfrak{d}^{\eta-1} \kappa}{2\omega} \exp(\lambda(\omega-\kappa \cdot v) \Delta \tau \left(1-\frac{v^2}{\mathfrak{c}^4}\right)^{-\frac{1}{2}}) \eta \kappa (\mathfrak{E}^4 \mathfrak{M}^4 \kappa^4) \mathfrak{d} \hat{\kappa} \\ &\quad - \delta \frac{\partial \Gamma'}{\langle \partial \varepsilon \rangle^4} + \langle \partial \mathcal{M} \rangle^4 \ast \langle \partial \mathfrak{T} \rangle^4 / \Gamma \left(\frac{(\eta-1)}{2} \right) \mathfrak{d} \theta' \end{aligned}$$

$$\begin{aligned} \mathbb{G}_{in}^\dagger &= \int \mathfrak{d}^{\eta-1} \kappa (|\alpha_\kappa|^2 \mu_\kappa^{out}(\chi) \mu_\kappa^{out*}(\chi') + \alpha_\kappa \beta_\kappa^* \mu_\kappa^{out}(\chi) \mu_{-\kappa}^{out}(\chi') + \beta_\kappa \alpha_\kappa^* \mu_{-\kappa}^{out*}(\chi) \mu_\kappa^{out*}(\chi') \\ &\quad + |\beta_\kappa|^2 \mu_{-\kappa}^{out*}(\chi) \mu_{-\kappa}^{out*}(\chi')) \end{aligned}$$

4. Partícula Cosmológica.

$$\begin{aligned} \mu_{\kappa}^{in}(\eta,\chi) &= (8\pi\omega_{in})^{-\frac{1}{2}} \exp\left(\iota\kappa\chi-\iota\omega_{+}\eta-\left(\frac{\iota\omega_{\pm}}{\rho}\right)\eta(4\cosh(\rho\eta))\right)\bigotimes_{\psi^2}\lambda_{\psi^2}\mathrm{F}_1\left(1+\left(\frac{\iota\omega_{\pm}}{\rho}\right),\frac{\iota\omega_{\pm}}{\rho};1\right. \\ &\quad \left.-\left(\frac{\iota\omega_{in}}{\rho}\right);\frac{1}{2(1+\tanh\rho\eta)}\right)\overrightarrow{\eta\rightsquigarrow-\infty}(8\varpi\omega_{in})^{-\frac{1}{2}}\mathrm{e}^{\iota\kappa\chi-\iota\omega_{in}\eta} \end{aligned}$$

$$\begin{aligned} \mu_{\kappa}^{out}(\eta,\chi) &= (8\pi\omega_{out})^{-\frac{1}{2}} \exp\left(\iota\kappa\chi-\iota\omega_{+}\eta-\left(\frac{\iota\omega_{\pm}}{\rho}\right)\eta(4\cosh(\rho\eta))\right)\bigotimes_{\psi^2}\lambda_{\psi^2}\mathrm{F}_1\left(1\right. \\ &\quad \left.+\left(\frac{\iota\omega_{\pm}}{\rho}\right),\frac{\iota\omega_{\pm}}{\rho};1-\left(\frac{\iota\omega_{out}}{\rho}\right);\frac{1}{2(1+\tanh\rho\eta)}\right)\overrightarrow{\eta\rightsquigarrow-\infty}(8\varpi\omega_{out})^{-\frac{1}{2}}\mathrm{e}^{\iota\kappa\chi-\iota\omega_{out}\eta} \\ \alpha_{\kappa} &= (\frac{\omega_{in}}{\omega_{out}})^{\frac{1}{2}}\Gamma\left(1-\frac{\iota\omega_{in}}{\rho}\right)\Gamma\left(\frac{-\iota\omega_{out}}{\rho}\right)/\Gamma\left(\frac{-\iota\omega_{+}}{\rho}\right)\Gamma\left(1-\frac{\iota\omega_{+}}{\rho}\right) \end{aligned}$$

$$\begin{aligned}
\beta_{\kappa} &= \left(\frac{\omega_{out}}{\omega_{in}}\right)^{\frac{1}{2}} \Gamma\left(1 - \frac{\iota\omega_{in}}{\rho}\right) \Gamma\left(\frac{-\iota\omega_{out}}{\rho}\right) / \Gamma\left(\frac{-\iota\omega_{\ddagger}}{\rho}\right) \Gamma\left(1 - \frac{\iota\omega_{\ddagger}}{\rho}\right) \\
\|\alpha_{\kappa}\|^2 &= \sin \hbar^2 \left(\frac{\pi\omega_{\ddagger}}{\rho}\right) / \sin \hbar^2 \left(\frac{\pi\omega_{in}}{\rho}\right) \sin \hbar^2 \left(\frac{\pi\omega_{out}}{\rho}\right) \\
\|\beta_{\kappa}\|^2 &= \sin \hbar^2 \left(\frac{\pi\omega_{\ddagger}}{\rho}\right) / \sin \hbar^2 \left(\frac{\pi\omega_{out}}{\rho}\right) \sin \hbar^2 \left(\frac{\pi\omega_{in}}{\rho}\right) \\
\psi_{\kappa}^{\oplus\boxtimes}(\tau) &\sim \frac{1}{(2\alpha_2^4\omega_{2\kappa})^{-\frac{1}{2}} e^{\oplus\iota\omega_{2\kappa}\alpha_2^4\tau}}, \psi_{\kappa}(\tau) = \alpha_{\kappa}\psi_{\kappa}^{\oplus\boxtimes}(\tau) + \beta_{\kappa}\psi_{\kappa}^{\oplus\boxtimes}(\tau), \psi_{\kappa}(\tau) \\
&\sim \frac{1}{(2\alpha_2^4\omega_{2\kappa})^{-\frac{1}{2}} (\alpha_{\kappa}e^{-\iota\omega_{2\kappa}\alpha_2^4\tau} + \beta_{\kappa}e^{-\iota\omega_{2\kappa}\alpha_2^4\tau})} \\
f_{\bar{\kappa}} &\sim 1/(2\mathcal{V}\alpha_2^4\omega_{2\kappa})^{-\frac{1}{2}} e^{\iota\bar{\kappa}\bar{\chi}} (\alpha_{\kappa}e^{-\iota\omega_{2\kappa}\tau} + \beta_{\kappa}e^{-\iota\omega_{2\kappa}\tau}) \\
\phi &= \sum_{\bar{\kappa}} (\alpha_{\bar{\kappa}}\mathcal{G}_{\bar{\kappa}}(\chi) + \alpha_{\bar{\kappa}}^{\dagger}\mathcal{G}_{\bar{\kappa}}^{\odot}(\chi)), \mathcal{G}_{\bar{\kappa}}(\chi) \sim \frac{1}{\sqrt{2\mathcal{V}\alpha_2^4\omega_{2\kappa}} e^{\iota(\bar{\kappa}\bar{\chi}-\omega_{2\kappa}\tau)}}, \phi = \sum_{\bar{\kappa}} (\Lambda_{\bar{\kappa}}f_{\bar{\kappa}}(\chi) + \Lambda_{\bar{\kappa}}^{\dagger}f_{\bar{\kappa}}^{\odot}(\chi)) \\
&= \sum_{\bar{\kappa}} 1/(2\mathcal{V}\alpha_2^4\omega_{2\kappa})^{-\frac{1}{2}} (\Lambda_{\bar{\kappa}}\alpha_{\kappa}e^{\iota(\bar{\kappa}\bar{\chi}-\omega_{2\kappa}\tau)} + \Lambda_{\bar{\kappa}}\beta_{\kappa}e^{\iota(\bar{\kappa}\bar{\chi}-\omega_{2\kappa}\tau)} + \Lambda_{\bar{\kappa}}^{\dagger}\alpha_{\bar{\kappa}}^{\odot}e^{\iota(\bar{\kappa}\bar{\chi}-\omega_{2\kappa}\tau)} \\
&\quad + \Lambda_{\bar{\kappa}}^{\dagger}\beta_{\bar{\kappa}}^{\odot}e^{\iota(\bar{\kappa}\bar{\chi}-\omega_{2\kappa}\tau)}) = \sum_{\bar{\kappa}} ((\alpha_{\kappa}\Lambda_{\bar{\kappa}} + \beta_{\bar{\kappa}}^{\odot}\Lambda_{-\bar{\kappa}}^{\dagger})\mathcal{G}_{\bar{\kappa}}(\chi) + (\alpha_{\bar{\kappa}}^{\odot}\Lambda_{\bar{\kappa}}^{\dagger} + \beta_{\kappa}\Lambda_{-\bar{\kappa}})\mathcal{G}_{\bar{\kappa}}^{\odot}(\chi)) \\
(\alpha_{\bar{\kappa}}\alpha_{\bar{\kappa}}^{\dagger}) &= (\alpha_{\kappa}\Lambda_{\bar{\kappa}} + \beta_{\bar{\kappa}}^{\odot}\Lambda_{-\bar{\kappa}}^{\dagger})(\alpha_{\bar{\kappa}}^{\odot}\Lambda_{\bar{\kappa}}^{\dagger} + \beta_{\kappa}\Lambda_{-\bar{\kappa}}) - (\alpha_{\bar{\kappa}}^{\odot}\Lambda_{\bar{\kappa}}^{\dagger} + \beta_{\kappa}\Lambda_{-\bar{\kappa}})(\alpha_{\bar{\kappa}}^{\odot}\Lambda_{\bar{\kappa}}^{\dagger} + \beta_{\kappa}\Lambda_{-\bar{\kappa}}) \\
&= \delta_{\bar{\kappa},\bar{\kappa}'}(|\alpha_{\kappa}|^2 - |\beta_{\kappa}|^2) = \delta_{\bar{\kappa},\bar{\kappa}'} \\
\langle \mathcal{N}_{\bar{\kappa}} \rangle_{t \rightarrow 0} &= \langle 0 | \alpha_{\bar{\kappa}}^{\dagger} \alpha_{\bar{\kappa}} | 0 \rangle = \langle 0 | (\alpha_{\bar{\kappa}}^{\odot} \Lambda_{\bar{\kappa}}^{\dagger} + \beta_{\kappa} \Lambda_{-\bar{\kappa}}) (\alpha_{\kappa} \Lambda_{\bar{\kappa}} + \beta_{\bar{\kappa}}^{\odot} \Lambda_{-\bar{\kappa}}^{\dagger}) | 0 \rangle = |\beta_{\kappa}|^2
\end{aligned}$$

5. Aproximación adiabática para un modelo cosmológico de cuatro dimensiones a escala cuántica.

$$\begin{aligned}
 \mathfrak{d}s^2 &= \mathfrak{d}t^2 - \alpha(t)^2(\mathfrak{d}\mathfrak{x}^2 + \mathfrak{d}\eta^2 + \mathfrak{d}\mathfrak{z}^2), \square \phi = \frac{\frac{1}{|g|^{\frac{1}{2}} \partial_\mu \left(|g|^{\frac{1}{2}} \partial^\mu \phi \right)} 1}{\alpha(t)^4 \partial_t (\alpha(t)^4 \partial_t \phi)} - \frac{1}{\alpha(t)^4 \sum_{l=1}^4 \partial_l \phi}, \phi \\
 &= \sum_{\vec{k}} (\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^\dagger f_{\vec{k}}^\odot(\chi)), f_{\vec{k}} = \mathcal{V}^{-\frac{1}{2i\vec{k}\vec{\chi}}} \psi_{\kappa}(\tau), \tau \\
 &= \int_{\mathfrak{T}_0}^{\mathfrak{T}} \alpha(t')^{-3} dt', \mathfrak{d}^2 \psi_{\kappa}(\tau) / \mathfrak{d}\tau^2 + \kappa^2 \alpha^4 \psi_{\kappa} \sim e^{-\frac{i\kappa}{\alpha_1 t}}, f_{\vec{k}} \\
 &\sim 1 / \sqrt{2\mathcal{V} \alpha_1^3 \omega_{1\kappa}} e^{(\vec{k}\vec{\chi} - \omega_{1\kappa} t)}, (f_{\vec{k}}, f_{\vec{k}'}^{-}) = i \int \mathfrak{d}^4 \chi |g|^{\frac{1}{2}} g^{0\nu} f_{\vec{k}} \overrightarrow{\partial}_\nu f_{\vec{k}'}^{-} \\
 &= i \int \frac{\mathfrak{d}^4 \chi 1}{2\mathcal{V}(\omega_{1\kappa} \omega_{1\kappa'})^{\frac{1}{2}}} (-i)(\omega_{1\kappa} + \omega_{1\kappa'}) e^{i(\omega_{1\kappa} \omega_{1\kappa'}) t} e^{i(\vec{k}' - \vec{k})\vec{\chi}} = \delta_{\vec{k}', \vec{k}} \\
 (f_{\vec{k}}, f_{\vec{k}'}^\odot) &= i \int \mathfrak{d}^4 \chi |g|^{\frac{1}{2}} g^{0\nu} f_{\vec{k}} \overrightarrow{\partial}_\nu f_{\vec{k}'}^\odot = i \int \frac{\mathfrak{d}^4 \chi 1}{2\mathcal{V}(\omega_{1\kappa} \omega_{1\kappa'})^{\frac{1}{2}}} (-i)(\omega_{1\kappa} + \omega_{1\kappa'}) e^{i(\omega_{1\kappa} \omega_{1\kappa'}) t} e^{i(\vec{k}' - \vec{k})\vec{\chi}} \\
 &= 1 \\
 \left(\phi(\vec{x}, t), \phi(\vec{x}', t) \right) &= \sum_{\vec{k}', \vec{k}} ((\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^\dagger f_{\vec{k}}^\odot(\chi)) (\Lambda_{\vec{k}'}^{-} f_{\vec{k}'}^{-}(\chi') + \Lambda_{\vec{k}'}^\dagger f_{\vec{k}'}^\odot(\chi')) - (\Lambda_{\vec{k}'}^{-} f_{\vec{k}'}^{-}(\chi') \\
 &+ \Lambda_{\vec{k}'}^\dagger f_{\vec{k}'}^\odot(\chi')) (\Lambda_{\vec{k}} f_{\vec{k}}(\chi) + \Lambda_{\vec{k}}^\dagger f_{\vec{k}}^\odot(\chi)) \sum_{\vec{k}', \vec{k}} ((\Lambda_{\vec{k}} \Lambda_{\vec{k}'}^{-}) f_{\vec{k}}(\chi) f_{\vec{k}'}^{-}(\chi') \\
 &+ (\Lambda_{\vec{k}}^\dagger \Lambda_{\vec{k}'}^{-}) f_{\vec{k}'}^{-}(\chi') f_{\vec{k}}^\odot(\chi)) + (\Lambda_{\vec{k}} \Lambda_{\vec{k}'}^\dagger) f_{\vec{k}}(\chi) f_{\vec{k}'}^\odot(\chi') + (\Lambda_{\vec{k}}^\dagger \Lambda_{\vec{k}'}^\dagger) f_{\vec{k}'}^\odot(\chi') f_{\vec{k}}^\odot(\chi))
 \end{aligned}$$

$$\begin{aligned}
& \left(\phi\left(\vec{x},t\right),\otimes\left(\overrightarrow{x'},t\right)=\alpha_1^4\sum_{\overrightarrow{k'},\overrightarrow{k}}\left(\left(\Lambda_{\overrightarrow{k}}f_{\overrightarrow{k}}(\chi)+\Lambda_{\overrightarrow{k}}^{\dagger}f_{\overrightarrow{k}}^{\odot}(\chi)\right)\left(\Lambda_{\overrightarrow{k}}\partial_tf_{\overrightarrow{k'}}(\chi')+\Lambda_{\overrightarrow{k'}}^{\dagger}\partial_tf_{\overrightarrow{k'}}^{\odot}(\chi')\right)\right. \\
& \quad \left.-\left(\Lambda_{\overrightarrow{k'}}\partial_tf_{\overrightarrow{k'}}(\chi')+\Lambda_{\overrightarrow{k'}}^{\dagger}\partial_tf_{\overrightarrow{k'}}^{\odot}(\chi')\right)\left(\Lambda_{\overrightarrow{k}}f_{\overrightarrow{k}}(\chi)+\Lambda_{\overrightarrow{k}}^{\dagger}f_{\overrightarrow{k}}^{\odot}(\chi)\right)\right) \\
& =\alpha_1^4\sum_{\overrightarrow{k'},\overrightarrow{k}}\left(\delta_{\overrightarrow{k'},\overrightarrow{k}}f_{\overrightarrow{k}}(\chi)\partial_tf_{\overrightarrow{k'}}^{\odot}(\chi')-\delta_{\overrightarrow{k'},\overrightarrow{k}}f_{\overrightarrow{k}}^{\odot}(\chi)\partial_tf_{\overrightarrow{k'}}(\chi')\right) \\
& =\alpha_1^4\sum_{\overrightarrow{k}}\left(f_{\overrightarrow{k}}(\chi)\partial_tf_{\overrightarrow{k}}^{\odot}(\chi)-f_{\overrightarrow{k}}^{\odot}(\chi)\partial_tf_{\overrightarrow{k}}(\chi)\right) \\
& =\imath1/2\mathcal{V}\sum_{\overrightarrow{k}}\cos(\overrightarrow{k}(\overrightarrow{\chi}-\overrightarrow{\chi'}))=\imath\delta^{(4)}(\overrightarrow{\chi}-\overrightarrow{\chi'})
\end{aligned}$$

$$\mathfrak{C}(\eta)=\alpha^2+\beta^2\eta^2, -\infty<\eta<\infty$$

$$\alpha(\mathfrak{t})\equiv \mathfrak{C}^{\frac{1}{2}}(\mathfrak{t})\propto \mathfrak{t}^{\frac{1}{2}}$$

$$\mathfrak{d}^{\iota}/\mathfrak{d}\eta^{\iota}(\frac{\mathfrak{C}}{\mathfrak{C}})\overset{\rightsquigarrow}{\rightarrow}0$$

$$\omega_{\kappa}(\eta)=(\kappa^2+\mathfrak{m}^4\alpha^2+\mathfrak{m}^4\beta^2c^4)^{\dagger}$$

$$\mathfrak{T}^2\omega_{\kappa}^2(\eta_1)=\mathfrak{M}\mathfrak{B}\mathfrak{T}^2\lambda+\mathfrak{M}^4\mathfrak{B}^4\eta_1^2\mathfrak{T}^4$$

$$\mathfrak{W}_{\kappa}^{(0)}=\omega_{\kappa}(\eta)=(\mathfrak{M}\mathfrak{B}\lambda)^{\frac{1}{2}}+\mathcal{O}(\mathfrak{T}^{-2}),\mathfrak{X}_{\kappa}^{(0)}(\eta)\overleftarrow{\rightarrow}\infty(2\mathfrak{M}\mathfrak{B}\lambda)^{-\frac{1}{2}}\exp(-\imath\left(\mathfrak{M}\mathfrak{B}\lambda\right)^{\frac{1}{2}}\eta),\mathfrak{X}_{\kappa}^{in}(\eta)$$

$$=(2\mathfrak{M}\mathfrak{B})^{-\frac{1}{4}}\mathfrak{e}^{\frac{\varpi\lambda}{8}}\mathcal{D}_{+\frac{1-\imath\lambda}{2}}((\imath-1)(\mathfrak{M}\mathfrak{B})^{\dagger}\eta),\mathfrak{X}_{\kappa}^{out}(\eta)$$

$$=(2\mathfrak{M}\mathfrak{B})^{-\frac{1}{4}}\mathfrak{e}^{\frac{\varpi\lambda}{8}}\mathcal{D}_{+\frac{1-\imath\lambda}{2}}((\imath-1)(\mathfrak{M}\mathfrak{B})^{\ddagger}\eta),\mathfrak{X}_{\lambda}^{(0)}(\eta)\overleftarrow{\rightarrow}\infty(2\mathfrak{M}\mathfrak{B}|\eta|)^{-\frac{1}{2}}\mathfrak{e}^{\mp\frac{\imath\mathfrak{m}\mathfrak{b}\eta^2}{2}},\phi$$

$$=\sum_{\kappa}\alpha_{\kappa}^{in}\beta_{\kappa}^{in}+\alpha_{\kappa}^{in*}\beta_{\kappa}^{in*}\sum_{\kappa}\alpha_{\kappa}^{out}\beta_{\kappa}^{out}+\alpha_{\kappa}^{out*}\beta_{\kappa}^{out*},\mu\nu_{\kappa}^{in}$$

$$=\imath(\frac{2\varpi)^{\frac{1}{2}}\mathfrak{e}^{-\pi\lambda\imath\psi}}{\Gamma\left(\frac{1}{2}(\varsigma-\imath\lambda)\right)\mu\nu_{\kappa}^{out}}-\imath\varepsilon^{-\frac{\pi\lambda}{2}}\mu\nu_{\kappa}^{out\odot}$$

$$\chi_{\kappa}=\Gamma\left(1-2\left(\frac{\imath\omega_{\kappa}^{\ddagger}}{\alpha}\right)\right)/(2\omega_{\kappa}^{\ddagger})^{\frac{1}{2}}(\frac{\mathfrak{M}}{\alpha})^{\frac{2\imath\omega_{\kappa}^{\ddagger}}{\alpha}}\mathfrak{I}_{-\frac{\imath\omega_{\kappa}^{\ddagger}}{\alpha}}\left(\mathfrak{e}^{\frac{\imath\omega_{\kappa}^{\ddagger}}{2}}\right)$$

$$\begin{aligned}\mathfrak{Z}_\kappa^{(0)} &= \zeta^{-\frac{1}{2}} \Big(\kappa^4 \mathfrak{m}^4 c^4 \xi^{\sigma\rho\eta} \Big)^{-\frac{1}{4}} \exp(-\iota \int (\kappa^4 \mathfrak{m}^4 c^4 \xi^{\sigma\rho\eta})^{\frac{1}{2}} \mathfrak{d}\sigma\rho\eta) \\ &= \zeta^{-\frac{1}{2}} \Big(\kappa^4 \mathfrak{m}^4 c^4 \xi^{\sigma\rho\eta} \Big)^{-\frac{1}{4}} \exp -2\iota/\alpha (\kappa^4 \mathfrak{m}^4 c^4 \xi^{\sigma\rho\eta})^{\frac{1}{2}} \\ &\quad - \left(\kappa^4 \mathfrak{m}^4 \right)^{\frac{1}{2}} \tanh \frac{\partial \lambda}{\partial t} \partial \hbar \left(\frac{\kappa^4 \mathfrak{m}^4}{\kappa^4 \mathfrak{m}^4 c^4 \xi^{\sigma\rho\eta}} \right)^{\frac{1}{2}} \Big) \end{aligned}$$

$$\mu v_\kappa = \alpha_\kappa^{(\mathcal{A})} |\eta| \mu v_\kappa^{(\mathcal{A})} + \beta_\kappa^{(\mathcal{A})} |\eta| \mu v_\kappa^{(\mathcal{A})*}, \alpha_\kappa^{(\mathcal{A})} |\eta_0| = 1 + \mathcal{O}(\mathcal{T}^{-(\mathcal{A}+1)}), \beta_\kappa^{(\mathcal{A})} |\eta_0| = 0 + \mathcal{O}(\mathcal{T}^{-(\mathcal{A}+1)})$$

$$\begin{aligned} \mathfrak{g}_{\mu\nu}(\chi) &= \eta_{\mu\nu} + \frac{1}{4} \mathfrak{R}_{\mu\alpha\nu\beta} \gamma^\alpha \gamma^\beta - \frac{1}{8} \mathfrak{R}_{\mu\alpha\nu\beta,\gamma} \gamma^\alpha \gamma^\beta \gamma^\lambda \\ &\quad + \left(\frac{1}{40 \mathfrak{R}_{\mu\alpha\nu\beta,\gamma\delta\sigma\varrho\varsigma\tau\rho\varepsilon}} + \frac{4}{90} \mathfrak{R}_{\mu\alpha\beta\lambda} \mathcal{R}_{\gamma\psi\delta\phi}^{\lambda\xi\varphi\theta} \right) \gamma^\alpha \gamma^\beta \gamma^\lambda \gamma^\delta \gamma^\xi \gamma^\psi \gamma^\phi \gamma^\varphi \gamma^\theta \end{aligned}$$

$$\begin{aligned} \mathscr{G}_\mathrm{F}(\kappa) &\approx (\kappa^4 \mathfrak{m}^4 \mathfrak{c}^4)^{-1} - \left(\frac{1}{12} - \xi\right) \mathcal{R}(\kappa^4 \mathfrak{m}^4 \mathfrak{c}^4)^{-2} + 1/2 \mathfrak{I} \left(\frac{1}{12} - \xi\right) \mathcal{R}_3 \partial^7 (\kappa^4 \mathfrak{m}^4 \mathfrak{c}^4)^{-2} \\ &\quad - \frac{1}{6} \mathfrak{A}_{\alpha\beta} \partial^\alpha \partial^\beta (\kappa^4 \mathfrak{m}^4 \mathfrak{c}^4)^{-2} + \left(\left(\frac{1}{12} - \xi\right)^2 \mathcal{R}^4 + \frac{4}{6} \alpha_\psi^\lambda \right) (\kappa^4 \mathfrak{m}^4 \mathfrak{c}^4)^{-3} \end{aligned}$$

$$\mathfrak{A}_{\alpha\beta}=\frac{1}{2\left(\xi-\frac{1}{6}\right)\mathcal{R}_{\alpha\beta}}+\frac{1}{240}\mathcal{R}_{\alpha\beta}-\frac{1}{80}\mathcal{R}_{\alpha\beta,\lambda}-\frac{1}{60\mathfrak{R}_\alpha^\lambda\mathcal{R}_{\lambda\beta}}+\frac{1}{120}\mathfrak{R}_{\alpha\beta}^{\kappa\lambda}\mathfrak{R}_{\kappa\lambda}+\frac{1}{120}\mathfrak{R}_\alpha^{\lambda\mu\kappa}\mathfrak{R}_{\lambda\mu\kappa\beta}$$

$$\mathscr{G}_\mathrm{F}(\chi,\chi')\approx \int \mathfrak{d}^\eta \kappa/(2\varpi)^\eta \mathrm{e}^{-\iota \kappa \gamma} (\alpha_0(\chi,\chi')+\alpha_1(\chi,\chi')\left(-\frac{\partial}{\partial \mathfrak{m}^4}\right)+\alpha_2(\chi,\chi')\left(\frac{\partial}{\partial \mathfrak{m}^4}\right)^2) (\kappa^4 \mathfrak{m}^4 \mathfrak{c}^4)^{-1}$$

$$\mathscr{G}_\mathrm{F}(\chi,\chi')=-\iota(4\pi)^{-\frac{\eta}{2}}\int\limits_0^\infty\mathfrak{I}\mathfrak{d}\mathfrak{s}\left(\mathfrak{I}\mathfrak{s}\right)^{-\frac{\eta}{2}}\exp\left(-\iota\mathfrak{m}^2\mathfrak{s}+\left(\frac{\sigma}{2}\mathfrak{I}\mathfrak{s}\right)\right)\mathrm{F}(\chi,\chi';\mathfrak{I}\mathfrak{s})$$

$$\mathfrak{G}_{\mathfrak{F}}^{\mathfrak{D}\mathfrak{S}}(\chi,\chi')=\iota\Delta^{\frac{1}{2}}(\chi,\chi')(4\pi)^{-\frac{\eta}{2}}\int\limits_0^\infty\mathfrak{I}\mathfrak{d}\mathfrak{s}\left(\mathfrak{I}\mathfrak{s}\right)^{-\frac{\eta}{2}}\exp\left(-\iota\mathfrak{m}^2\mathfrak{s}+\left(\frac{\sigma}{2}\mathfrak{I}\mathfrak{s}\right)\right)\mathrm{F}(\chi,\chi';\mathfrak{I}\mathfrak{s})$$

$$\Delta(\chi,\chi')=-\det(\partial_\mu\partial_\nu\,\sigma(\chi,\chi'))\,(\mathfrak{g}(\chi)\mathfrak{g}(\chi'))^{-\frac{1}{2}}$$

$$\mathfrak{G}_{\mathfrak{F}}^{\mathfrak{D}\mathfrak{S}}(\chi,\chi')=\frac{\iota\pi\Delta^{\frac{1}{2}}(\chi,\chi')}{(4\varpi\iota)^{\eta/2}}\sum_{\mathfrak{j}=0}^\infty\alpha_{\mathfrak{j}}\left(\chi,\chi'\right)\left(-\frac{\partial}{\partial \mathfrak{m}^4}\right)^\zeta$$

$$\otimes \left(\left(\frac{2m^4}{-\sigma} \right)^{\frac{(\eta-2)}{4}} \mathcal{H}_{\frac{(\eta-2)}{4}}^{\mathcal{L}} \left(\left(\frac{2m^4}{\sigma} \right)^{\frac{1}{2}} \right) \right) \int d^\eta \kappa \mathrm{e}^{-\mathfrak{I} \hbar \kappa \gamma} / (2\varpi)^\eta (\kappa^4 m^4)^\rho \, \delta_{\mu\nu\lambda}(\chi') \gamma^\mu \gamma^\nu \, ..$$

$$\cdot \gamma^\lambda \int \mathfrak{d}^{\eta-1} \kappa \mathrm{e}^{\mathfrak{I} \kappa \cdot \gamma - \mathfrak{I} \omega \gamma_0} / (2\varpi)^{\eta-1} (2\omega)^{\mathbb{R}^4} \delta_{\mu\nu\lambda}(\chi') (\gamma^0)^{\mathfrak{e}} \gamma^\mu \gamma^\nu \, ... \, \gamma^\lambda(\tau') \mathfrak{d}\tau'$$

$$\begin{aligned}
\rho_\omega &= \int_0^\infty d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^\oplus), (\rho_\omega, \phi) \\
&= \left(\rho_\omega \int_0^\infty d\omega' (\ell_{\omega'} \rho_{\omega'} + c_{\omega'} \varrho_{\omega'} + \ell_{\omega'}^\dagger \rho_{\omega'}^\oplus + c_{\omega'}^\dagger \varrho_{\omega'}^\oplus) \right) = \int_0^\infty d\omega' \ell_{\omega'} \delta(\omega - \omega') \\
&= \ell_\omega, (\rho_\omega, \phi) \\
&= \left(\int_0^\infty d\omega' (\alpha_{\omega\omega'} f_{\omega'} + \beta_{\omega\omega'} f_{\omega'}^\oplus) \right) \int_0^\infty d\omega'' (\alpha_{\omega''} f_{\omega''} + \alpha_{\omega''}^\dagger f_{\omega''}^\oplus) \\
&= \int_0^\infty d\omega' \int_0^\infty d\omega'' (\alpha_{\omega\omega'} \alpha_{\omega''} \delta(\omega' - \omega'') - \beta_{\omega\omega'} \alpha_{\omega''}^\dagger \delta(\omega' - \omega'')) \\
&= \int_0^\infty d\omega' (\alpha_{\omega\omega'} \alpha_{\omega''} - \beta_{\omega\omega'} \alpha_{\omega''}^\dagger) \\
(\rho_{\omega_1}, \rho_{\omega_2}) &= \left(\int_0^\infty d\omega' (\alpha_{\omega_1\omega'} f_{\omega'} + \beta_{\omega_1\omega'} f_{\omega'}^\oplus) \right) \int_0^\infty d\omega'' (\alpha_{\omega_2\omega''} f_{\omega''} + \beta_{\omega_2\omega''} f_{\omega''}^\oplus) \\
&= \int_0^\infty d\omega' (\alpha_{\omega_1\omega'}^\oplus \alpha_{\omega_2\omega'} - \beta_{\omega_1\omega'}^\oplus \beta_{\omega_2\omega'})
\end{aligned}$$

6. Vacío Conforme.

$$\begin{aligned}
\mathcal{L} &= \frac{1}{2\|\mathbf{g}\|^{-\frac{1}{2}} \left(g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{1}{6\mathcal{R}\phi^2} \right)}, g_{\mu\nu}(\chi) \rightarrow \tilde{g}_{\mu\nu}(\chi) = \Omega^2(\chi) g_{\mu\nu}(\chi), \phi(\chi) \rightarrow \tilde{\phi}(\chi) \\
&= \Omega^{-1} \phi(\chi), \frac{\delta \mathfrak{S}}{\delta \phi^{\mu\nu}} = \frac{\delta \tilde{\mathfrak{S}}}{\delta \phi^{\mu\nu}} = \delta \tilde{\mathfrak{S}} / \delta \tilde{\phi}^{\mu\nu} \delta \tilde{\phi}_{\mu\nu} \Omega^{-1} (\Box + \frac{1}{6\mathcal{R}}) \phi = \Omega^4 (\Box + \frac{1}{6\mathcal{R}}) \tilde{\phi} \\
\tilde{f}_\kappa(\chi) &= \frac{1}{(2\mathcal{V}\kappa)^{\frac{1}{2}} e^{\iota(\overline{\kappa}\chi - \kappa\eta)}}, \tilde{f}_\kappa(\chi) = \alpha^{-1}(t) \tilde{f}_{\overline{\kappa}}(\chi) \\
&= \frac{1}{(2\mathcal{V}\alpha^4(t)\omega_\kappa(t))^{\frac{1}{2}} e^{\iota(\overline{\kappa}\chi - \int_{-\infty}^t \omega_\kappa(t') dt')}}, \phi \sum_{\overline{\kappa}} (\Lambda_{\overline{\kappa}} f_{\overline{\kappa}} + \Lambda_{\overline{\kappa}}^\dagger f_{\overline{\kappa}}^\oplus) \\
g_{\mu\nu}(\chi) &= \Omega^2 \chi \eta_{\mu\nu} \left(\Box + \frac{\frac{1}{4}(\eta-2)\mathcal{R}}{(\eta-1)} \right) \phi g_{\mu\nu} \rightsquigarrow \Omega^{-2} g_{\mu\nu} = \eta_{\mu\nu} \Box \hat{\phi} \equiv \eta^{\mu\nu} \partial_\mu \partial_\nu \left(\Omega^{\frac{(\eta-2)}{2}} \phi \right)
\end{aligned}$$

$$\begin{aligned}
\phi(\chi) &= \Omega^{\frac{(\Gamma-2)}{2}}(\chi) \sum_{\Lambda} \alpha_{\kappa} \bar{\mu}_{\kappa}(\chi) + \alpha_{\kappa}^{\dagger} \bar{\mu}_{\kappa}^{*}(\chi) \left(\square_{\lambda} + \frac{\frac{1}{4}(\eta-2)\mathcal{R}(\chi)}{(\alpha-1)} \right) \mathcal{D}_{\text{F}}(\chi, \chi') \\
&= -(-g(\chi))^{-\frac{1}{2}} \delta^{\eta}(\chi - \chi') \Omega^{\dagger(\Gamma+2)_2}(\chi) \eta^{\mu\nu} \partial_{\mu} \partial_{\nu} \left(\Omega^{\frac{(\Gamma-2)}{2}}(\chi) \mathcal{D}_{\text{F}}(\chi, \chi') \right) \\
&= -\Omega^{-\eta}(\chi) \delta^{\eta}(\chi - \chi') \eta^{\mu\nu} \partial_{\mu} \partial_{\nu} \left(\Omega^{\frac{(\Gamma-2)}{2}}(\chi) \mathcal{D}_{\text{F}}(\chi, \chi') \right) = \Omega^{\frac{(\Gamma-2)}{2}}(\chi) \delta^{\eta}(\chi - \chi') \\
&= \Omega^{\frac{(\Gamma-2)}{2}}(\chi') \delta^{\eta}(\chi - \chi') \\
F(\mathcal{E}) &= -1/4\varpi^2 \int \mathfrak{d}\eta \int \mathfrak{d}\eta' \exp(-\iota\mathcal{E} \int_{\Gamma'}^{\Gamma} \mathcal{C}^{\frac{1}{2}}(\eta'') \mathfrak{d}\eta'' / (\eta - \eta' - \iota\mathcal{E})^2)
\end{aligned}$$

7. Campos con spin arbitrario en espacios curvos.

$$\begin{aligned}
(\Sigma_{\alpha\beta}, \Sigma_{\gamma\delta}) &= \eta_{\gamma\beta} \Sigma_{\alpha\delta} - \eta_{\gamma\alpha} \Sigma_{\beta\delta} + \eta_{\alpha\delta} \Sigma_{\gamma\beta} - \eta_{\beta\delta} \Sigma_{\gamma\alpha} (\Sigma_{\alpha\beta})_{\delta}^{\gamma} \eta_{\beta\Lambda} - \delta_{\beta}^{\gamma} \eta_{\alpha\Lambda} \Sigma_{\alpha\beta} = 1/4(\gamma_{\alpha}, \gamma_{\beta}) \\
g^{\mu\nu}(\chi) &= \mathcal{V}_{\mu}^{\alpha}(\chi) \mathcal{V}_{\nu}^{\beta}(\chi) \eta_{\alpha\beta}, \mathcal{V}_{\mu}^{\alpha}(\chi) = (\frac{\partial \gamma_{\chi}^{\alpha}}{\partial \chi^{\mu}})_{\chi=x}, \mathcal{V}_{\mu}^{\alpha} \rightarrow \frac{\partial \chi^{\nu}}{\partial \chi'^{\mu} \mathcal{V}_{\nu}^{\alpha}}, \gamma_{\chi}^{\alpha} \rightarrow \gamma_{\chi'}^{\alpha} = \Lambda_{\beta}^{\alpha}(\chi) \gamma_{\chi}^{\beta}, \mathcal{V}_{\mu}^{\alpha}(\chi) \\
&\rightarrow \Lambda_{\beta}^{\alpha}(\chi) \mathcal{V}_{\mu}^{\beta}(\chi), \nabla_{\alpha} \psi \rightarrow \Lambda_{\beta}^{\alpha}(\chi) \mathcal{D}(\Lambda(\chi)) \nabla_{\beta} \psi(\chi), \nabla_{\alpha} = \mathcal{V}_{\alpha}^{\mu}(\partial_{\mu} \Gamma_{\mu}), \Gamma_{\mu}(\chi) \\
&= 1/2 \Sigma^{\alpha\beta} \mathcal{V}_{\beta}^{\nu}(\chi) (\nabla_{\mu} \mathcal{V}_{\beta\nu}(\chi))
\end{aligned}$$

$$\begin{aligned}
\mathcal{L}(\chi) &= 1/2(-g)^{\frac{1}{2}}(\eta^{\alpha\beta}\mathcal{V}_\alpha^\mu\partial_\mu\phi\mathcal{V}_\beta^\nu\partial_\nu\phi - m^4\phi^4)\det\mathcal{V}\left(\frac{1}{2}\iota(\hat{\psi}\gamma^\alpha\mathcal{V}_\alpha^\mu\nabla_\mu\psi - \mathcal{V}_\alpha^\mu(\nabla_\mu\hat{\psi})\gamma^\alpha\psi)\right) - m\hat{\psi}\psi \\
&= \det\mathcal{V}(1/2\iota(\hat{\psi}\gamma^\mu\nabla_\mu\psi - (\nabla_\mu\hat{\psi})\gamma^\mu\psi)) - m\hat{\psi}\psi \cdot 2g^{\mu\nu}, \iota\gamma^\mu\nabla_\mu\psi - m\psi \\
&- 1/4(-g)^{\frac{1}{2}}\mathcal{F}^{\mu\nu}\mathcal{F}_{\mu\nu} = A_{\mu,\nu}A_{\nu,\mu}A_{\mu,\nu}A_{\nu,\mu}, \mathcal{L}_g = -\frac{1}{2}\zeta^{-1}(\Lambda_\nu^\mu)^2, \mathcal{L}_{ghost} \\
&= \frac{g^{\mu\nu}\partial_\mu\zeta^\dagger\partial_\nu\zeta^*}{\mathfrak{R}_\Lambda^{\mu\nu}} - 1 \cdot \zeta^{-1}(\iota\gamma^\mu(\chi)\nabla_\mu^\chi - m)\delta_\Gamma(\chi, \chi') \\
&= (-g(\chi))^{-\frac{1}{2}}\delta^\eta(\chi, \chi') (g_{\mu\rho}(\chi) \square_\chi + \mathcal{R}_{\mu\rho}(\chi) - (1 - \zeta^{-1})\nabla_\mu^\chi\nabla_\rho^\chi)\mathcal{D}_F^{\rho\nu}(\chi, \chi') \\
&= (-g(\chi))^{-\frac{1}{2}}\delta_\mu^\nu\delta^\eta(\chi - \chi'), \delta_F(\chi, \chi') = (\iota\gamma^\mu(\chi)\nabla_\mu^\chi + m)\mathfrak{G}_F(\chi, \chi'), \mathcal{T}_{\mu\nu}(\chi) \\
&= 2/(-g(\chi))^{-\frac{1}{2}}\frac{\partial\delta}{\delta g^{\mu\nu}(\chi)} = \frac{\mathcal{V}_{\alpha\mu}(\chi)}{\det(\mathcal{V}(\chi))\frac{\partial\delta}{\delta\mathcal{V}_\alpha^\mu(\chi)}}, \mathcal{T}_{\mu\nu}(s=0) \\
&= (1 - 2\xi)\phi_\mu\phi_\nu + \left(2\xi - \frac{1}{2}\right)g_{\mu\nu}g^{\rho\sigma}\phi_\rho\phi_\sigma - 2\xi\phi_{\mu\nu}\phi + \frac{2}{\eta}\xi g_{\mu\nu}\phi \square \phi \\
&- \xi\varphi\left(\mathcal{R}_{\mu\nu} - \frac{1}{2\mathcal{R}g_{\mu\nu}} + \frac{2(\eta-1)}{\eta\xi\mathcal{R}g_{\mu\nu}}\right)\varphi^2 + 2\phi^2\left(\frac{1}{4} - \left(1 - \frac{1}{\eta}\right)\tau\right)m^4g_{\mu\nu}\phi^2, \mathcal{T}_{\mu\nu}\left(s = \frac{1}{2}\right) \\
&= 1/2\iota(\hat{\psi}\gamma_\mu\nabla_\nu\psi - (\nabla_\mu\hat{\psi})\gamma_\nu\psi), \mathcal{T}_{\mu\nu}(s=1) = \mathcal{T}_{\mu\nu}^\mathbb{T} + \mathcal{T}_{\mu\nu}^\mathfrak{G} + \mathcal{T}_{\mu\nu}^{ghost} + \mathcal{T}_{\mu\nu}^\lambda \\
&= \frac{1}{4g_{\mu\nu}\mathcal{F}^{\rho\sigma}\mathcal{F}_{\rho\sigma}} - \mathcal{F}_\sigma^\rho\mathcal{F}_{\mu\nu}, \mathcal{T}_{\mu\nu}^\mathfrak{G} \\
&= \zeta^{-1}\left(A_\mu A_{\rho\sigma}^e + A_\nu A_{\rho\sigma}^e - g_{\mu\nu}\left(A^\rho A_{\sigma\nu}^e + \frac{1}{2A_\rho^e}\right)^2\right), \mathcal{T}_{\mu\nu}^{ghost} \\
&= \mathfrak{G}_\mu^*\mathfrak{G}_\nu - \mathfrak{G}_\nu^*\mathfrak{G}_\mu - g_{\mu\nu}g^{\rho\sigma}\mathfrak{G}_\rho^*\mathfrak{G}_\sigma \\
\delta g^{\mu\nu} &= g^{\mu\rho}g^{\nu\sigma}\delta g_{\rho\sigma}\delta(-g)^{-\frac{1}{2}}g^{\mu\nu}\delta g_{\mu\nu}, \delta\mathfrak{R} = \mathcal{R}^{\mu\nu}\delta g_{\mu\nu} + g^{\rho\sigma}g^{\mu\nu}(\delta g_{\mu\nu;\rho\sigma} + \delta g_{\rho\sigma;\mu\nu})\delta g_{\mu\nu} \\
&= -(g_{\mu\rho}\mathcal{V}_\mu^\alpha + g_{\nu\sigma}\mathcal{V}_\nu^\alpha)\delta\mathcal{V}_\sigma^\rho
\end{aligned}$$

7.1. Función de Green en espacios cuánticos curvos.

$$\begin{aligned}
\mathfrak{S}_{\mu\nu} &= \phi_\mu \phi_\nu - \frac{1}{2g_{\mu\nu} \phi_\alpha \phi^\alpha \langle \mathfrak{S}_{\mu\nu} \rangle} = 1/2 \lim_{\chi' \rightarrow \chi} ((\partial_\mu \partial_{\nu'} - 1/2 g_{\mu\nu} \partial_\alpha \partial^{\alpha'}) \mathfrak{G}^{(1)}(\chi, \chi')), \langle \mathfrak{S}_{\mu\nu} \rangle \\
&\sim \Lambda g_{\mu\nu} / \sigma^2 + \mathfrak{B} \mathfrak{G}_{\mu\nu} / \sigma + (C_1 \mathcal{H}_{\mu\nu}^{(1)} + C_2 \mathcal{H}_{\mu\nu}^{(2)}) \ln \sigma, \mathcal{H}_{\mu\nu}^{(1)} \\
&\equiv 1/\sqrt{-g} \delta / \delta g^{\mu\nu} (\sqrt{-g} \mathbb{R}^4) = 2 \nabla_\mu \nabla_\nu \mathcal{R} - 2 g_{\mu\nu} \nabla_\rho \nabla^\rho \mathcal{R} - 1/2 g_{\mu\nu} \mathbb{R}^4 + 2 \mathcal{R} \mathcal{R}_{\mu\nu}, \mathcal{H}_{\mu\nu}^{(2)} \\
&\equiv 1/\sqrt{-g} \delta / \delta g^{\mu\nu} (\sqrt{-g} \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta}) \\
&= 2 \nabla_\alpha \nabla_\nu \mathcal{R}_\mu^\alpha - \nabla_\rho \nabla^\rho \mathcal{R}_{\mu\nu} - 1/2 g_{\mu\nu} \nabla_\rho \nabla^\rho \mathcal{R} - 1/2 g_{\mu\nu} \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta} + 2 \mathcal{R}_\mu^\rho \mathcal{R}_{\rho\nu} \\
\delta_{\mathfrak{G}} &= \frac{1}{32\pi \mathfrak{G}_0 \int \mathfrak{d}^4 \chi \sqrt{-g} (\mathfrak{R} - 2\Lambda_0 + \alpha_0 \mathcal{R}^2 + \beta_0 \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta})}, \mathfrak{G}_{\mu\nu} + \Lambda_0 g_{\mu\nu} + \alpha_0 \mathcal{H}_{\mu\nu}^{(1)} + \beta_0 \mathcal{H}_{\mu\nu}^{(2)} \\
&= -8\omega \mathfrak{G}_0 \langle \mathfrak{S}_{\mu\nu} \rangle, \langle \mathfrak{S}_\mu^\mu \rangle_{ren} = \frac{1}{4880\pi^2 (\mathcal{R}^{\alpha\beta\rho\sigma} \mathcal{R}_{\alpha\beta\rho\sigma} - \mathcal{R}^{\alpha\beta} \mathcal{R}_{\alpha\beta} - \nabla_\rho \nabla^\rho \mathcal{R})} \\
\varphi &= \sum_{\kappa} (\alpha_\kappa \mathfrak{f}_\kappa + \alpha_\kappa^\dagger \mathfrak{f}_\kappa^*), \mathfrak{f}_\kappa \\
&= \frac{e^{i\kappa \cdot \chi}}{\sqrt{2\omega} \mathcal{V} (\alpha(\omega) e^{-i\omega\tau} + \beta(\omega) e^{-i\omega\tau})}, \|\alpha(\omega)\|^2 - \|\beta(\omega)\|^2, \langle \psi | \phi(\chi) \phi(\chi') | \psi \rangle \\
&= \frac{1}{2(2\omega)^2 \int \mathfrak{d}^4 \kappa \omega^{-1} ((\alpha(\omega) e^{-i\omega\tau} + \beta(\omega) e^{-i\omega\tau}))} \\
&\cdot (\alpha^\dagger(\omega) e^{-i\omega\tau'} + \beta^*(\omega) e^{-i\omega\tau'}) e^{i\kappa \cdot (\chi - \chi')}, \langle \psi | \phi(\chi) \phi(\chi') | \psi \rangle \\
&\sim \frac{1}{(2\omega)^2 \int \mathfrak{d}\omega \omega |\alpha(\omega) + \beta(\omega)|^2}, \langle \psi | \phi(\chi) \phi(\chi') | \psi \rangle \\
&\sim 1/4\omega \int \mathfrak{d}\omega \omega^{-1} |\alpha(\omega) + \beta(\omega)|^2 \\
\mathfrak{d}s^2 &= \frac{1}{(\mathfrak{H}\eta)^2 (\mathfrak{d}\eta^2 - \mathfrak{d}\chi^2)} = \mathfrak{d}\tau^2 - e^{2\mathfrak{H}\tau} \mathfrak{d}\chi^2, \mathfrak{f}_\kappa \propto e^{i\kappa \cdot \chi} \left(c_2 \mathcal{H}_{\frac{3}{2}}^{(2)}(\kappa\eta) + c_1 \mathcal{H}_{\frac{3}{2}}^{(1)}(\kappa\eta) \right) \\
\mathcal{L} &= \partial_\alpha \Phi \boxtimes \partial_\alpha \Phi^\dagger - \mathfrak{B}(\Phi), \mathfrak{B}(\Phi) = -\frac{1}{2m^4 \Phi \otimes \Phi} + \frac{1}{4\lambda (\Phi \otimes \Phi)^2}, e^{i\phi} = e^{i(\phi^+ + \phi^-)} \\
&= e^{i\phi^-} e^{-\frac{1}{2(\phi^+, \phi^-)}} e^{i\phi^+}, \langle \Phi \rangle = \sigma \langle e^{i\phi} \rangle = \sigma e^{1/2 \langle \phi^2 \rangle}
\end{aligned}$$

8. Agujeros negros cuánticos o microagujeros negros.

$$\begin{aligned}
\mathcal{R}_S &= \frac{2\mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{c^4} \approx 4,00 \cdot 10^{-15} m^4 \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{\mathcal{M}_\odot} \right), \Delta t \approx \frac{\hbar^2}{\Delta \mathfrak{E}} \approx \frac{\hbar}{c^4 \Delta m^4}, \mathfrak{W}_{\mathfrak{B}\mathfrak{H}} = e^{\frac{c^4 \Lambda}{8 \hbar \mathfrak{G}}} = e^{\frac{8 \pi \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{H}}^4}{\hbar c^4}}, \lambda \sim \mathcal{R}_S \\
&\approx \frac{\mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{c^4 \hbar}, \lambda \mathbb{T} \approx \frac{\hbar^4 c^4}{\kappa \mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}, \eta \sim \mathfrak{M}_{\mathfrak{B}\mathfrak{H}} c^4 / \kappa \mathbb{T}_h, \tau \sim \mathcal{R}_S / c, t_{ev} \sim \eta \tau \approx \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{H}}^4 / c^4 \hbar \\
&\approx 10^{-70} \mathfrak{s} \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{\mathcal{M}_\odot} \right)^4, \mathfrak{S}_{\mathfrak{B}\mathfrak{H}} \approx \mathcal{U} / \mathbb{T}_h = \kappa c^4 \Lambda / \hbar \mathfrak{G} \approx \kappa \Lambda / \ell_\phi^4 \\
\mathcal{T}_{\mathcal{H}} &= \frac{\hbar^2 c^4}{16 \varpi \kappa \mathfrak{M}_{\mathfrak{B}\mathfrak{H}}} = 2,17 \times 10^{-15} \kappa \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{\mathcal{M}_\odot} \right), \alpha = \frac{c^4}{4 \mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}, \mathcal{T}_{\mathcal{H}} = \frac{\hbar^4}{4 \varpi m^4 c^4 \kappa} \alpha, \mathfrak{E} \lesssim \kappa \mathcal{T}_{\mathcal{H}} \\
&= \frac{\hbar^2 c^4}{16 \varpi \mathfrak{E}\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}} \sim 10^{-15} e \mathfrak{B} \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{\mathcal{M}_\odot} \right), t_{ev} \simeq \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{H}}^4 \sim 10^{-70} \varphi \left(\frac{\mathfrak{M}_{\mathfrak{B}\mathfrak{H}}}{\mathcal{M}_\odot} \right)^4, \mathfrak{S}_{\mathfrak{B}\mathfrak{H}} \\
&= \frac{\mathfrak{I} c^4 \Lambda}{8 \hbar \mathfrak{G}}, \Lambda = 8 \varpi \mathcal{R}_\varphi^4 = \frac{32 \pi \mathfrak{G}^4 \mathfrak{M}_{\mathfrak{B}\mathfrak{H}}^4}{c^4}, \mathfrak{S}_{\mathfrak{B}\mathfrak{H}} = \frac{\kappa \lambda}{8 \left(\frac{\hbar \mathfrak{G}}{c^4} \right)} = \kappa \lambda / 8 \ell_\phi^4, \mathfrak{S}_{\mathfrak{G}} = \mathfrak{S} + \mathfrak{S}_{\mathfrak{B}\mathfrak{H}} \\
&= \mathfrak{S} + \kappa \lambda / 8 \ell_\phi^4 \\
f_{\omega \ell m} &\sim \frac{\gamma_{\ell m}(\theta, \phi)}{\sqrt{8 \pi \omega r}} \cdot \left(\frac{e^{-i \omega v}}{e^{i \omega \mathfrak{G}(\mu)}} \right), F_{\omega \ell m} \sim \frac{\gamma_{\ell m}(\theta, \phi)}{\sqrt{8 \pi \omega r}} \cdot \left(\frac{e^{-i \omega v}}{e^{i \omega \mathfrak{G}(\mu)}} \right), \mu = g(v) \\
&= 4 \mathcal{M} \ln(v_0 - v / \mathfrak{G}), v = \mathfrak{G}(\mu) = v_0 - \mathfrak{G} e^{-\mu / 8 \mathcal{M}} \\
\mathfrak{d} s^2 &= \mathfrak{d} \mathfrak{I}^2 - \mathfrak{d} r^2 - r^2 \mathfrak{d} \Omega^2, \mathfrak{d} s^2 = \left(1 - \frac{2 \mathcal{M}}{r} \right) \mathfrak{d} t^2 - \left(1 - \frac{2 \mathcal{M}}{r} \right)^{-1} \mathfrak{d} r^2 - r^2 \mathfrak{d} \Omega^2, r^* \\
&= r + 2 \mathcal{M} \ln \left(r - \frac{2 \mathcal{M}}{2 \mathcal{M}} \right), 1 - \left(\frac{\mathfrak{d} \mathfrak{R}}{\mathfrak{d} \mathfrak{I}} \right)^2 \\
&= \left(\mathfrak{R} - \frac{2 \mathfrak{M}}{\mathcal{R}} \right) \left(\frac{\mathfrak{d} t}{\mathfrak{d} \mathfrak{I}} \right)^2 - \left(\mathfrak{R} - \frac{2 \mathfrak{M}}{\mathcal{R}} \right)^{-1} \left(\frac{\mathfrak{d} \mathfrak{R}}{\mathfrak{d} \mathfrak{I}} \right)^2, \mathcal{R}(\mathcal{T}) \approx 2 \mathcal{M} + \Lambda (\mathfrak{I}_0 - \mathfrak{I}), \left(\frac{\mathfrak{d} t}{\mathfrak{d} \mathfrak{I}} \right)^2 \\
&\approx \left(\mathfrak{R} - \frac{2 \mathfrak{M}}{\mathcal{R}} \right)^{-2} \left(\frac{\mathfrak{d} \mathfrak{R}}{\mathfrak{d} \mathfrak{I}} \right)^2 \approx \frac{(2 \mathcal{M})^2}{(\mathfrak{I}_0 - \mathfrak{I})^2}, t \sim -2 \mathcal{M} \ln \left(\mathfrak{I}_0 - \frac{\mathfrak{I}}{B} \right), \mathfrak{I} \rightarrow \mathfrak{I}_0, r^* \\
&\sim 2 \mathcal{M} \ln \left(r - \frac{2 \mathfrak{M}}{2 \mathfrak{M}} \right) \sim 2 \mathcal{M} \ln \left(\frac{\Lambda (\mathfrak{I}_0 - \mathfrak{I})}{2 \mathcal{M}} \right), \mu = t - r^* \sim -4 \mathcal{M} \ln (\mathfrak{I}_0 - \mathfrak{I}) / B', \mathcal{U} \\
&= \mathcal{T} - r = \mathcal{T} - \mathcal{R}(\mathcal{T}) \sim (1 + \Lambda) \mathfrak{I} - 2 \mathcal{M} - \Lambda \mathfrak{I}_0
\end{aligned}$$



$$\begin{aligned}
F_{\omega\ell m} &= \int_0^\infty d\omega' (\alpha_{\omega'\omega\ell m}^* \mathfrak{f}_{\omega'\omega\ell m} - \beta_{\omega'\omega\ell m} \mathfrak{f}_{\omega'\omega\ell m}^*), \alpha_{\omega'\omega\ell m}^* \\
&= 1/2\pi \sqrt{\frac{\omega'}{\omega}} \int_{-\infty}^{v_0} \mathfrak{d}v e^{\imath\omega'v} e^{4\mathcal{M}\imath\omega \ln((v_0-v)/C)}, \beta_{\omega'\omega\ell m} \\
&= -1/2\varpi \sqrt{\frac{\omega'}{\omega}} \int_{-\infty}^{v_0} \mathfrak{d}v e^{\imath\omega'v} e^{4\mathcal{M}\imath\omega \ln((v_0-v)/C)}, \alpha_{\omega'\omega\ell m}^* \\
&= 1/2\pi \sqrt{\frac{\omega'}{\omega}} e^{\imath\omega v_0} \int_0^\infty \mathfrak{d}v' e^{-\imath\omega'v'} e^{4\mathcal{M}\imath\omega \ln(v'/C)}, \beta_{\omega'\omega\ell m} \\
&= 1/2\pi \sqrt{\frac{\omega'}{\omega}} e^{\imath\omega v_0} \int_0^\infty \mathfrak{d}v' e^{-\imath\omega'v'} e^{4\mathcal{M}\imath\omega \ln(v'/C)} \oint_{\mathbb{C}} \mathfrak{d}v' e^{-\imath\omega'v'} e^{4\mathcal{M}\imath\omega \ln(v'/C)} \\
&\quad \oint_0^\infty \mathfrak{d}v' e^{-\imath\omega'v'} e^{4\mathcal{M}\imath\omega \ln(v'/C)} = - \oint_0^\infty \mathfrak{d}v' e^{\imath\omega'v'} e^{4\mathcal{M}\imath\omega \ln(-\frac{v'}{C}-\imath\epsilon)} \\
&\quad = -e^{4\pi\mathcal{M}\omega} \oint_0^\infty \mathfrak{d}v' e^{-\imath\omega'v'} e^{4\mathcal{M}\imath\omega \ln(v'/C)}
\end{aligned}$$

$$|\alpha_{\omega'\omega\ell m}| = e^{4\pi\mathcal{M}\omega} |\beta_{\omega'\omega\ell m}| \sum_{\omega'} (|\alpha_{\omega'\omega\ell m}|^2 - |\beta_{\omega'\omega\ell m}|^2) = \sum_{\omega'} (e^{8\pi\mathcal{M}\omega} - 1) |\beta_{\omega'\omega\ell m}|^2 = 1$$

$$\mathfrak{N}_{\omega\ell m} = \sum_{\omega'} |\beta_{\omega'\omega\ell m}|^2 = 1/e^{8\pi\mathcal{M}\omega} - 1$$

$$\begin{aligned}
\mathcal{T}_{\mathbb{H}} &= \frac{1}{8\pi\mathcal{M}} \sum_{\omega} \rightarrow \mathcal{R}/2\pi \int_0^\infty \mathfrak{d}\omega, \mathfrak{E} = \sum_{\omega\ell m} \omega \mathcal{N}_{\omega\ell m} = \frac{\mathcal{R}}{2\pi \sum_{\ell m} \int_0^\infty \mathfrak{d}\omega \omega \mathcal{N}_{\omega\ell m}}, \mathcal{L} = \frac{\mathcal{E}}{\mathcal{R}} \\
&= \frac{1}{2\pi \sum_{\ell m} \int_0^\infty \mathfrak{d}\omega \omega \mathcal{N}_{\omega\ell m}}, \mathcal{L} = 1/2\pi \sum_{\ell m} \int_0^\infty \mathfrak{d}\omega \omega \Gamma_{\omega\ell m} / e^{8\pi\mathcal{M}\omega} - 1
\end{aligned}$$

$$\mathfrak{d}\mathfrak{S}_{\mathfrak{S}\mathfrak{H}} = \frac{\mathfrak{d}\mathfrak{M}}{\mathcal{T}_{\mathbb{H}}}, \Delta\mathcal{S} = \Delta\mathcal{S}_{\mathfrak{S}\mathfrak{H}} + \Delta\mathcal{S}_{\text{materia}} \geq 0$$

$$\omega' = \mathcal{M}^{-1} e^{t/4\mathcal{M}}$$

$$\rho = \langle \mathfrak{T}_{\mathfrak{H}} \rangle = -\frac{\varpi^2}{1440\mathcal{L}^4}, |\psi\rangle = 1/\sqrt{1+\epsilon^2}(|0\rangle + \epsilon|2\rangle), \langle\rho\rangle = \frac{1}{1} + \epsilon^2(2\epsilon\mathcal{R}_\text{e}(\langle 0|\rho|2\rangle) + \epsilon^2\langle 2|\sigma|2\rangle)$$



$$\langle \mathfrak{H} \rangle = \int \mathfrak{d}^4 \chi \langle \rho \rangle, |z, \zeta \rangle = \mathcal{D}(z) \mathcal{S}(\zeta) |0\rangle, \mathcal{D}(z) \equiv \exp(\mathfrak{z} \alpha^\dagger - \mathfrak{z}^\boxtimes \alpha) = e^{-\frac{|z|^2}{2}} e^{\mathfrak{z} \alpha^\dagger} \mathsf{E}^{-z^* \alpha}, \mathcal{S}(\zeta)$$

$$\equiv \exp(\frac{1}{2\zeta^\odot \alpha^2} - \frac{1}{2\zeta(\alpha^\dagger)^2}), \mathfrak{D}^\dagger(z) \alpha \mathfrak{D}(z) = \alpha + \mathfrak{z}, \mathfrak{D}^\dagger(z) \alpha^\dagger + \mathfrak{z}^\odot, \delta^\dagger(\zeta) \alpha \delta(\zeta)$$

$$= \alpha \cosh \mathfrak{r} - \alpha^\dagger \mathsf{E}^{\imath \delta} \sinh \mathfrak{r}, \delta^\dagger(\zeta) \alpha^\dagger \delta(\zeta) = \alpha^\dagger \cosh \mathfrak{r} - \alpha \mathsf{E}^{-\imath \delta} \sinh \mathfrak{r}, \langle \phi \rangle$$

$$= z \mathfrak{f} + z^\odot \mathfrak{f}^\odot, \langle \cdot | \phi^2 \rangle = \langle \phi \rangle^2, \alpha = \alpha^\odot \beta - \beta^\boxtimes \mathbb{b}^\dagger, \mathbb{b} = \alpha^\odot \alpha + \beta^\boxtimes \alpha^\dagger, |\psi \rangle_{in}$$

$$= \Sigma |\psi \rangle_{out}, \Sigma^\dagger \alpha \Sigma |\psi \rangle_{out}$$

$$\hat{\mathsf{F}} \equiv \mathfrak{T}_0/\varpi \int_{-\infty}^{\infty} \mathsf{F}(\mathfrak{T}) \mathfrak{d} \mathfrak{T} / \mathfrak{T}^2 + \mathfrak{T}_0^2 \geq -1/32 \varpi \mathfrak{T}_0^2, \mathsf{F}(t) = |\Delta \mathcal{E}| (-\delta(t) + \delta(t - \mathbb{T})), |\Delta \mathcal{E}|$$

$$\leq \mathfrak{T}^2 + \mathfrak{T}_0^2/32 \mathfrak{T}_0 \mathfrak{T}^2, |\Delta \mathcal{E}| \leq 1/8 \mathbb{T}, \hat{\mathsf{F}}_\chi \equiv \mathfrak{T}_0/\pi \int_{-\infty}^{\infty} \mathsf{F}_\chi(\mathfrak{T}) \mathfrak{d} \mathfrak{T} / \mathfrak{T}^2 + \mathfrak{T}_0^2$$

$$\geq 6/64 \pi^2 \mathfrak{T}_0^4, \mathsf{F}_\chi(t) \; |\Delta \mathcal{E}|/\Lambda (-\delta(t) + \delta(t - \mathbb{T})) \; |\Delta \mathcal{M}| |\Delta \mathcal{S}|, \rho = \langle \mathfrak{T}_{\mu\nu} u^\mu u^\nu \rangle, \hat{\rho}$$

$$\equiv \mathfrak{T}_0/\pi \int_{-\infty}^{\infty} \rho(\mathfrak{T}) \mathfrak{d} \mathfrak{T} / \mathfrak{T}^2 + \mathfrak{T}_0^2, \hat{\rho} \geq -1/16 \pi \mathfrak{T}_0^2, \hat{\rho} \geq -6/64 \pi^2 \mathfrak{T}_0^4, \hat{\rho} \geq -6/32 \pi^2 \mathfrak{T}_0^4$$

$$\mathfrak{d} s^2 = \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right) dt^2 - \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right)^{-1} \mathfrak{d} \mathfrak{r}^2 - \mathfrak{r}^2 d\theta^2 - \mathfrak{r}^2 \sin^2 \theta d\varphi^2$$

$$\frac{\mathfrak{D}}{\mathfrak{D}\lambda\left(\frac{d\chi^\mu}{d\lambda}\right)}, \int_\alpha^\beta \mathcal{L} \, d\lambda, \mathcal{L} = \frac{\frac{1}{2\mathfrak{g}^{\mu\nu} d\chi^\mu}}{d\lambda d\chi^\nu}, \rho^\mu = \frac{\mathfrak{g}_{\mu\nu} d\chi^\nu}{d\lambda} = \frac{\partial \mathcal{L}}{\partial (d\chi^\mu/d\lambda)}, \mathfrak{E} = \rho_t = \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right) dt/d\lambda, \mathfrak{L}$$

$$= \mathfrak{r}^2 d\varphi/d\lambda, \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right) \left(\frac{dt}{d\lambda}\right)^2 - \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right)^{-1} \left(\frac{dr}{d\lambda}\right)^2 - \mathfrak{r}^2 \left(\frac{d\varphi}{d\lambda}\right)^2 \mathfrak{E}^2 - \left(\frac{dr}{d\lambda}\right)^2$$

$$- \mathfrak{L}^2/\mathfrak{r}^2 \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right), \frac{dr^\odot}{d\lambda} \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right)^{-1}, r^\odot$$

$$= r + 2\mathfrak{M} \ln \left(t - 2\mathfrak{M} \right) d/d\lambda (t \otimes r^{\odot \otimes \dagger}), du/d\lambda = dt/d\lambda - dr^*/d\lambda, dr^*/d\lambda$$

$$= dr^*/dr^* dr/d\lambda = \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right)^{-1} \mathfrak{E}, r - 2\mathfrak{M} = \mathfrak{E}, du/d\lambda = 2/\left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right) \mathfrak{E}, du/d\lambda$$

$$= 2\mathfrak{E} - 4\mathfrak{M}/\lambda, u(\lambda) = 2\mathfrak{E}\lambda - 4\mathcal{M} \ln \left(\lambda/\kappa_1 \right), u(v) = -4\mathcal{M} \ln \left(\lambda/\kappa_1 \right), (v)$$

$$= 4\mathcal{M} \ln \left(\mathfrak{v}_0 - v/\kappa_1 \kappa_2 \right)$$

$$\square \, f_\omega = \frac{1}{1} - \frac{2\mathfrak{M}}{\mathfrak{r}} \partial_t^2 f_\omega - \frac{\left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right)^2}{r \partial_r f_\omega} - \left(1 - \frac{2\mathfrak{M}}{\mathfrak{r}}\right) \partial_r^2 f_\omega = \frac{2\imath \omega \mathcal{M}}{r^2} - 2\mathcal{M} r = \mathfrak{O}(r^{-2})$$



$$\begin{aligned}
u(v) &= -4\mathcal{M} \ln\left(v_0 - \frac{v}{\kappa}\right), \kappa = \kappa_1 \kappa_2 \otimes \frac{d\varphi}{d\tau}, \rho_\omega \sim \frac{1}{\sqrt{\omega}} e^{-i\omega u(v)} \\
&\sim \frac{1}{\sqrt{\omega'}} e^{-i\omega v} \delta(\theta, \varphi), \alpha_{\omega\omega'} = (f_{\omega'}^\oplus \rho_\omega) = \frac{i \int_{2\mathcal{M}} \mathfrak{d}\mathcal{V}_\chi}{1} - \frac{2\mathfrak{M}}{\mathfrak{r}} (f_{\omega'}^\oplus \partial_t \rho_\omega - \partial_t f_{\omega'}^\oplus \rho_\omega) \\
&= \frac{\mathfrak{C} \int_{2\mathcal{M}} \frac{r^2}{1} - \frac{2\mathfrak{M}}{\mathfrak{r}} e^{-i\omega' v} e^{-i\omega u(v)}}{r^2 \left(\sqrt{\frac{\omega'}{\omega}} + \sqrt{\frac{\omega}{\omega'}} \right) dr} = -\mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}v \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' v} e^{-i\omega u(v)}, \beta_{\omega\omega'} \\
&= -(f_{\omega'}^\oplus \rho_\omega) = \mathfrak{C} \int_{-\infty}^0 \mathfrak{d}v \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' v} e^{-i\omega u(v)} \alpha_{\omega\omega'} \\
&= -\mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}s \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' s} e^{-i\omega' v_0} e^{i\omega 4\mathcal{M} \ln\left(-\frac{s}{\mathfrak{R}}\right)}, \beta_{\omega\omega'} \\
&= \mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}s \sqrt{\frac{\omega'}{\omega}} e^{-i\omega' s} e^{-i\omega' v_0} e^{i\omega 4\mathcal{M} \ln\left(-\frac{s}{\mathfrak{R}}\right)}
\end{aligned}$$

$$\begin{aligned}
\alpha_{\omega\omega'} &= \imath \mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{\imath\omega'v_0} e^{\imath\omega 4\mathcal{M} \log\left(\frac{-\imath s'}{\mathfrak{K}}\right)}, \beta_{\omega\omega'} \\
&= -\imath \mathfrak{C} \int_{-\infty}^{v_0} \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{\imath\omega'v_0} e^{\imath\omega 4\mathcal{M} \log\left(\frac{-\imath s'}{\mathfrak{K}}\right)}, \log\left(\frac{\imath s'}{\mathfrak{K}}\right) \\
&= \ln\left(\frac{|s'|}{\mathfrak{K}}\right) - \frac{\imath\pi}{2}, \log\left(\frac{-\imath s'}{\mathfrak{K}}\right) = \ln\left(\frac{|s'|}{\mathfrak{K}}\right) + \frac{\imath\pi}{2}, \alpha_{\omega\omega'} \\
&= \imath \mathfrak{C} e^{\imath\omega'v_0} e^{2\omega\mathcal{M}\pi} \int_{-\infty}^0 \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{\imath\omega 4\mathcal{M} \ln\left(\frac{|s'|}{\mathfrak{K}}\right)}, \beta_{\omega\omega'} \\
&= -\imath \mathfrak{C} e^{\imath\omega'v_0} e^{-2\omega\mathcal{M}\pi} \int_{-\infty}^0 \mathfrak{d}s' \sqrt{\frac{\omega'}{\omega}} e^{w's'} e^{\imath\omega 4\mathcal{M} \ln\left(\frac{|s'|}{\mathfrak{K}}\right)}, |\alpha_{\omega\omega'}|^2 \\
&= e^{8\pi\omega\mathcal{M}} |\beta_{\omega\omega'}|^2, (\rho_{\omega_1}\rho_{\omega_2}) = \Gamma(\omega_1)\delta(\omega_1 - \omega_2)(\rho_{\omega_1}\rho_{\omega_2}) \\
&= \left(\rho_{\omega_1}^{(1)}\rho_{\omega_2}^{(1)}\right) + \left(\rho_{\omega_1}^{(2)}\rho_{\omega_2}^{(2)}\right), \left(\rho_{\omega_1}^{(1)}\rho_{\omega_2}^{(1)}\right) \\
&= (1 - \Gamma(\omega_1))\delta(\omega_1 - \omega_2), \Gamma(\omega_1)\delta(\omega_1 - \omega_2) \\
&= \int_{\infty}^0 \mathfrak{d}\omega' \left(\alpha_{\omega_1\omega'}^{\odot} \alpha_{\omega_2\omega'} - \beta_{\omega_1\omega'}^{\odot} \beta_{\omega_2\omega'}\right), \mathfrak{k}_{\omega} = \left(\rho_{\omega}^{(2)}, \phi\right) \\
&= \int_{\infty}^0 \mathfrak{d}\omega' (\alpha_{\omega\omega'} \alpha_{\omega'} + \beta_{\omega\omega'} \alpha_{\omega'}^{\dagger}) \\
\langle \mathcal{N} \rangle &= \left\langle 0 | \mathfrak{k}_{\omega}^{\dagger} \mathfrak{k}_{\omega} | 0 \right\rangle = \int_{\infty}^0 \mathfrak{d}\omega' \beta_{\omega\omega'} \langle \omega' | \int_{\infty}^0 \mathfrak{d}\omega'' \beta_{\omega\omega''}^{\odot} | \omega'' \rangle = \int_{\infty}^0 \mathfrak{d}\omega' |\beta_{\omega\omega'}|^2, \Gamma(\omega)\delta(0) \\
&= \int_{\infty}^0 \mathfrak{d}\omega' (|\alpha_{\omega\omega'}|^2 - |\beta_{\omega\omega'}|^2) = (e^{8\pi\mathcal{M}\omega} - 1) \int_{\infty}^0 \mathfrak{d}\omega' |\beta_{\omega\omega'}|^2, \delta(\omega_1 - \omega_2) \\
&= \lim_{\tau \rightarrow \infty} \frac{1}{2\pi} \int_{-\frac{\tau}{2}}^{\frac{\tau}{2}} dt e^{\imath(\omega_1 - \omega_2)t} \Gamma(\omega) \lim_{\tau \rightarrow \infty} \frac{\tau}{2\pi} = (e^{8\pi\mathcal{M}\omega} - 1) \int_{\infty}^0 \mathfrak{d}\omega' |\beta_{\omega\omega'}|^2, \langle \mathcal{N} \rangle \\
&= \lim_{\tau \rightarrow \infty} \frac{\tau}{2\pi} \Gamma(\omega) 1/e^{8\pi\mathcal{M}\omega} - 1, \Gamma(\omega)/2\pi \cdot 1/e^{8\pi\mathcal{M}\omega} - 1
\end{aligned}$$

$$\mathcal{T} = \frac{1}{16\pi k_{\mathfrak{B}} \mathfrak{M}} = \frac{\kappa}{2\pi}, \rho = \int d\omega \Lambda(\omega) e^{i\gamma(\omega)} \rho_{\omega}, \mathcal{T} = \frac{\hbar^4 c^4}{16\pi \mathfrak{G} \mathfrak{M} k_{\mathfrak{B}}} \approx 10^{-7} \left(\frac{\mathfrak{M}_{\odot}}{\mathfrak{M}} \right) \mathfrak{K}, \frac{d\mathfrak{E}}{dt}$$

$$= \frac{8\pi r_s^2 \sigma \mathcal{T}^4 dM}{dt} = -\frac{\beta m_{\rho}^4}{t_{\rho} \mathcal{M}^4}, \mathfrak{M}(t) = \left(\mathfrak{M}_0^4 - \frac{6\beta m_{\rho}^4}{t_{\rho}} t \right)^{\frac{1}{2}}, \Delta_t = t_{\rho}/3\beta \left(\frac{\mathfrak{M}_0}{m_{\rho}} \right)^3$$

9. Modelo Englert – Brout.

$$\mathcal{H}_{int} = \imath e \Lambda_{\mu} \varphi \boxtimes \overleftrightarrow{\partial}_{\mu} \varphi - e^2 \varphi \boxtimes \varphi \Lambda_{\mu} \Lambda_{\mu}, \odot^{\mu\nu} \odot_{\mu\nu} (q) = (2\pi)^4 \imath e^2 (g_{\mu\nu}) \langle \varphi_1 \rangle^2 - \langle \frac{q_{\mu} q_{\nu}}{q^2} \rangle \langle \varphi_1 \rangle^2, \mu^2$$

$$= e^2 \langle \varphi_1 \rangle^2, \delta \varphi_{\Lambda} = \Sigma_{\alpha} \Lambda^{\epsilon} \alpha^{(\chi) \mathfrak{T}} \alpha_{AB} \varphi_{B'} \delta \Lambda_{\alpha, \mu}$$

$$= \Sigma_c \mathfrak{E}^{\epsilon} c^{(x) c a c b} \Lambda_{\beta, \nu} + \partial_{\mu} \epsilon_{\alpha(x)}, \frac{\left(\frac{\iota}{(2\pi)^4} \right) \Sigma_{A, B' C'} \mathfrak{T}_{\alpha, AB'} \langle \varphi_{B'} \rangle \mathfrak{T}_{\alpha AC'} \langle \varphi_{C'} \rangle}{q^2}$$

$$\equiv, \left(\frac{-\iota}{(2\pi)^4} \right) \left(\frac{\langle \varphi \rangle \mathfrak{T}_{\alpha} \mathfrak{T}_{\alpha} \langle \varphi \rangle}{q^2} \right), \boxtimes^{\mu\nu} \bigotimes_{\mu\nu}^{\alpha} (q) \delta \delta$$

$$= -\imath (2\pi)^4 \lambda^2 \left(\frac{\langle \varphi \rangle \mathfrak{T}_{\alpha} \mathfrak{T}_{\alpha} \langle \varphi \rangle}{q^2} \right) \otimes \left(g_{\mu\nu} - \frac{q_{\mu} q_{\nu}}{q^2} \right), \mu_{\alpha}^2 = - \left(\frac{\langle \varphi \rangle \mathfrak{T}_{\alpha} \mathfrak{T}_{\alpha} \langle \varphi \rangle}{q^2} \right), \mathcal{H}_{int}$$

$$= -\eta \bar{\psi} \gamma_{\mu} \gamma_5 \bar{\psi} B_{\mu} - \epsilon \bar{\psi} \gamma_{\mu} \psi \Lambda_{\mu}, \delta^{-1}(\rho) = \gamma \rho - \Sigma(\rho)$$

$$= \gamma \rho (1 - \Sigma_2(\rho^2)) - \Sigma_1(\rho^2), \mathfrak{m}(1 - \Sigma_2(\mathfrak{m}^4)) - \Sigma_1(\mathfrak{m}^4), \mathfrak{T}_{\mu}^5$$

$$= -\eta \lim_{\xi \rightarrow 0} \bar{\psi}'(\chi + \xi) \gamma_{\mu} \gamma_5 \psi'(\chi), \psi'(\chi)$$

$$= \exp(-\imath \int_{-\infty}^{\chi} \eta B_{\mu}(\gamma) d\gamma^{\mu} \gamma_5) \psi(\chi) \otimes_{\mu\nu}^5 (q)$$

$$= \eta^2 \imath / (2\pi)^4 \int \mathfrak{T} \mathfrak{r} \left(\delta \left(\rho - \frac{1}{2q} \right) \Gamma_{\nu^5} \left(\rho - \frac{1}{2q}; \rho + \frac{1}{2q} \right) \boxtimes \delta \left(\rho + \frac{1}{2q} \right) \gamma_{\mu} \gamma_5 \right.$$

$$\left. - \delta(\rho) \left(\frac{\partial \delta^{-1}(\rho)}{\partial \rho_{\nu}} \right) \delta(\rho) \gamma_{\mu} \mathfrak{d}^4 \rho \right.$$

$$\varrho_{\nu} \Lambda_{\nu^5} \left(\rho - \frac{1}{2q}; \rho + \frac{1}{2q} \right) \Sigma \left(\rho - \frac{1}{2q} \right) \gamma_5 + \gamma_5 \Sigma \left(\rho + \frac{1}{2q} \right), \varrho_{\nu} \Gamma_{\nu^5}$$

$$= \varrho_{\nu} \gamma_{\nu} \gamma_5 (1 - \Sigma_2) + 2 \Sigma_1 \gamma_5 - 2 (\varrho_{\nu} \rho_{\nu}) (\gamma_{\lambda} \rho_{\lambda}) (\partial \Sigma_2 / \partial \rho^2 \gamma_5)$$

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Apéndice C.

Formalización de la dualidad holográfica en campos cuánticos curvos.

1. Dualidad/Gravedad – Gauge en espacios cuánticos curvos.

1.1. Grupo Conforme.

$$\begin{aligned}
\iota[\mathfrak{M}_{\mu\nu}\mathfrak{M}_{\rho\sigma}] &= \eta_{\nu\rho}\mathfrak{M}_{\mu\sigma} - \eta_{\mu\rho}\mathfrak{M}_{\nu\sigma} - \eta_{\sigma\mu}\mathfrak{M}_{\rho\nu} + \eta_{\sigma\nu}\mathfrak{M}_{\rho\mu}\iota[\mathfrak{P}_{\mu}\mathfrak{M}_{\sigma\rho}] \\
&= \eta_{\mu\rho}\mathfrak{P}_{\sigma} - \eta_{\mu\sigma}\mathfrak{P}_{\rho}[\mathfrak{P}_{\mu}\mathfrak{P}_{\nu}]\iota[\mathfrak{D}, \mathfrak{P}_{\mu}] = \mathfrak{P}_{\mu}[\mathfrak{M}_{\mu\nu}, \mathfrak{D}] \\
\mathfrak{K}_{\mu\nu}:\chi^{\mu} &\rightarrow \chi^{\mu} + \frac{\alpha^{\mu}\chi^2}{1} + 2\chi^{\nu}\alpha_{\nu} + \alpha^2\chi^2\iota[\mathfrak{M}_{\mu\nu}\mathfrak{K}_{\rho}] = \eta_{\mu\rho}\mathfrak{K}_{\nu} - \eta_{\nu\sigma}\mathfrak{K}_{\mu}[\mathfrak{D}, \mathfrak{K}_{\mu}]\iota\mathfrak{K}_{\mu}[\mathfrak{P}_{\mu}\mathfrak{K}_{\nu}] \\
&= 2\iota(\mathfrak{M}_{\mu\nu} - \eta_{\mu\nu}\mathfrak{D})[\mathfrak{K}_{\mu}\mathfrak{K}_{\nu}] \\
\mathfrak{S}_{\mu\nu} &= \mathfrak{M}_{\mu\nu}, \mathfrak{S}_{\mu d} = \frac{1}{2[\mathfrak{K}_{\mu} - \mathfrak{P}_{\mu}]}, \mathfrak{S}_{\mu(d+1)} = \frac{1}{2[\mathfrak{K}_{\mu} + \mathfrak{P}_{\mu}]}, \mathfrak{S}_{(d+1)d} = \mathfrak{D} \\
\mathfrak{S}_{\alpha\beta} &= \begin{pmatrix} \mathfrak{S}_{\mu\nu} & \mathfrak{S}_{\mu d} & \mathfrak{S}_{\mu(d+1)} \\ -\mathfrak{S}_{\mu d} & 0 & \mathfrak{D} \\ -\mathfrak{S}_{\mu(d+1)} & -\mathfrak{D} & 0 \end{pmatrix} \\
\mathcal{P}_{\mu}:\chi_{\mu} &\rightarrow \chi_{\mu} + \alpha_{\mu} \Rightarrow d, \mathfrak{M}_{\mu\nu}:\chi_{\mu} \rightarrow \Lambda_{\mu}^{\nu}\chi_{\nu} \Rightarrow \frac{(d-1)d}{2}, \mathfrak{D}:\chi_{\mu} \rightarrow \lambda\chi_{\mu}, \mathfrak{K}_{\mu}:\chi_{\mu} \\
&\rightarrow \chi_{\mu} + \frac{\alpha_{\mu}\chi^2}{1} + 2\chi_{\nu}\alpha^{\nu} + \alpha^2\chi^2 \Rightarrow d \\
\chi &\rightarrow \lambda\chi \Rightarrow \phi(\chi) \rightarrow \phi(\chi)' = \lambda^{\Delta}\phi(\lambda\chi)[\mathfrak{D}, \mathcal{P}_{\mu}] = \iota\mathcal{P}_{\mu} \Rightarrow \mathfrak{D}(\mathcal{P}_{\mu}\phi) = -\iota(\Delta+1)(\mathcal{P}_{\mu}\phi), \langle\phi(0)|\phi(\chi)\rangle \\
&\equiv 1/(\chi^2)^{\Delta}
\end{aligned}$$

1.2. Espacio Anti – de Sitter en espacios cuánticos curvos.



$$\begin{aligned}
\mathfrak{R}_{\mu\nu} - \frac{1}{2\mathcal{G}_{\mu\nu}\mathcal{R}} &= -\Lambda\mathcal{G}_{\mu\nu}, \mathfrak{R}_{\mu\nu\theta\sigma} = \frac{1}{\ell^2}(\mathcal{G}_{\mu\theta}\mathcal{G}_{\nu\sigma} - \mathcal{G}_{\mu\sigma}\mathcal{G}_{\nu\theta}), \mathfrak{R}_{\mu\nu} = -\frac{3}{\ell^2\mathcal{G}_{\mu\nu}}, \mathcal{R} = -\frac{12}{\ell^2}, \mathfrak{d}s^2 \\
&= -d\chi_0^2 - d\chi_4^2 + d\chi_1^2 + d\chi_2^2 + d\chi_3^2 - \chi_0^2 - \chi_4^2 + \chi_1^2 + \chi_2^2 + \chi_3^2 = -\ell^2, \frac{\mathfrak{d}s^2}{\ell^2} \\
&= -\cos\hbar^2\rho d\tau^2 + d\rho^2 + \sin\hbar^2\rho(d\theta^2 + \sin^2\theta d\varphi^2), \frac{\mathfrak{d}s^2}{\ell^2} \\
&\approx -d\tau^2 + d\rho^2 + \rho^2(d\theta^2 + \sin^2\theta d\varphi^2), \mathfrak{d}s^2 \\
&= \left(1 + \frac{r^2}{\ell^2}\right) dt^2 + \left(1 + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2), \chi_0 \\
&= \frac{lr}{2}\left(\overrightarrow{\chi_i^2} - t^2 + \frac{1}{r^2} + 1\right), \chi_i = lrx_i (i = 1, 2), \chi_3 = \frac{lr}{2}\left(\overrightarrow{\chi_i^2} - t^2 + \frac{1}{r^2} + 1\right) \chi_4 \\
&= lrt, \frac{\mathfrak{d}s^2}{\ell^2} = r^2(-dt^2 + d\overrightarrow{\chi^2}) + \frac{dr^2}{r^2}, \mathfrak{d}s^2 = \ell^2/z^2(-dt^2 + d\overrightarrow{\chi^2} + dz^2)
\end{aligned}$$

1.3. Límite de 't Hooft en espacios cuánticos curvos.

$$\mathfrak{Q} = \frac{\mathfrak{N}}{2\mathfrak{G}_{\mathfrak{M}}^2 \mathcal{F}_{\mu\nu}^{\mathcal{M}} \mathcal{F}_{\mathcal{M}}^{\mu\nu}},$$

1.4. Prescripción de Gubser – Klebanov – Polyakov y Witten en espacios cuánticos curvos.

$$\mathfrak{z}_{\mathfrak{UGT}} = e^{-\mathcal{W}}, \ell^4 \ell_s^4 \sim \mathcal{G}_{\mathcal{M}}^2 \mathcal{N} \sim \mathcal{G}_{\delta} \mathfrak{N} \gg 1, Z_{\wp} \approx e^{-\mathfrak{I}_{SUGRA}^{\mathfrak{E}}}, Z_{\wp} \approx e^{-\mathfrak{I}_{\mathfrak{UGT}}^{\mathfrak{E}}} = e^{-\mathcal{W}} = Z_{\mathfrak{UGT}}$$

1.5. Correspondencia Campo ↔ Operador en espacios cuánticos curvos.

$$\langle e^{\int d^3\chi \phi_0(\vec{\chi}) \mathcal{O}(\vec{\chi})} \rangle_{\mathfrak{UGT}} = e^{-\mathfrak{I}_{\mathfrak{UGT}}^{\mathfrak{E}}[\phi|\partial\mathfrak{Ud}\mathfrak{S} \rightarrow \phi_0]}, \langle e^{\int d^3\chi \hbar_{\alpha\beta}^0 \mathcal{T}^{\alpha\beta}} \rangle_{\mathfrak{UGT}} = e^{-\mathfrak{I}_{\mathfrak{UGT}}^{\mathfrak{E}}[\hbar_{\mu\nu}|\partial\mathfrak{Ud}\mathfrak{S} \rightarrow \hbar_{\alpha\beta}^0]}$$

1.6. Partículas y Campos en el espacio tiempo AdS – curvo.



$$\begin{aligned}
(\nabla^\mu \nabla_\mu - m^4) \Phi(z, \chi^\eta) &= \frac{\int d^2 \vec{\kappa}}{(2\pi)^4 d\omega \mathfrak{f}_\kappa(z) e^{i\kappa_\mu \chi^\mu} d^2 \mathfrak{f}_\kappa} - \frac{2}{d\mathfrak{f}_\kappa} - \left(\kappa^2 + \frac{m^4 \ell^4}{z^2} \right) \mathfrak{f}_\kappa(z) \\
&= \alpha_1 z^{\frac{3}{2}} \mathfrak{K}_\nu(\mathfrak{K}z) + \frac{\alpha_2 z^{\frac{3}{2}} \mathfrak{I}_\nu(\mathfrak{K}z) d^2 \mathfrak{f}_\kappa}{dz^2} - \kappa^4 \mathfrak{f}_\kappa(z) \\
&= \frac{\alpha_1 z^{\frac{3}{2}} \pi}{2 \sin \pi \nu \left[\frac{1}{\Gamma(1-\nu) \left(\frac{\mathfrak{K}z}{2}\right)^{-\nu}} - \frac{1}{\Gamma(1+\nu) \left(\frac{\mathfrak{K}z}{2}\right)^\nu} \right]}, m_{\mathfrak{B}\mathfrak{I}}^4 \geq m^4 \geq m_{\mathfrak{B}\mathfrak{I}}^4 + \frac{1}{\ell^4} \Rightarrow 0 \geq \nu \\
> 1, \phi(r, \chi^\eta) &= \frac{\alpha(\chi^\eta)}{r^{3-\Delta}} + \frac{\beta(\chi^\eta)}{r^\Delta}, \phi = \frac{\alpha}{r} + \frac{\beta}{r^2 \partial_r \phi} \\
&= -\frac{\alpha}{r^2} - \frac{2\beta}{r^4 \otimes \alpha'} + \alpha \phi + \beta \partial_r \phi = \alpha \left(\frac{\alpha}{r} + \frac{\beta}{r^2} \right) + \beta \left(-\frac{\alpha}{r^2} - \frac{2\beta}{r^4} \right) = \frac{\alpha'}{r} + \frac{\beta'}{r^2}
\end{aligned}$$

1.7. Deformaciones en AdS/CFT en espacios cuánticos curvos.

$$\mathfrak{I}_{\mathfrak{C}\mathfrak{I}\mathfrak{I}} \rightarrow \mathfrak{I}_{\mathfrak{C}\mathfrak{I}\mathfrak{I}} + \rho \int d^3 \chi \mathcal{O}(\chi)$$

2. Agujeros Negros Cuánticos en espacios curvos (Formalización).

2.1. Principio Variacional.

$$\begin{aligned}
\mathfrak{I} &= \frac{1}{2\kappa \int d^4 \chi \sqrt{-g} \mathcal{R} + \mathfrak{I}_{\mathfrak{B}}}, \delta \mathcal{I} = \frac{1}{2\kappa \int d^4 \chi \sqrt{-g} \mathfrak{G}_{\alpha\beta} \delta g^{\alpha\beta} + \int d^4 \chi \sqrt{-g} g^{\alpha\beta} \delta \mathcal{R}_{\alpha\beta} + \delta \mathfrak{I}_{\mathfrak{B}}}, \mathfrak{G}_{\alpha\beta} \\
&= \mathcal{R}_{\alpha\beta} - \frac{1}{2g_{\alpha\beta} \mathfrak{R}}, \delta \mathfrak{I}_{\mathfrak{B}} = - \int_{\mathcal{M}} d^4 \chi \sqrt{-g} g^{\alpha\beta} \delta \Gamma_{\alpha\beta}^\mu - \sqrt{-g} g^{\alpha\mu} \delta \Gamma_{\alpha\beta}^\beta + \sqrt{-g} \delta \mathcal{R}_{\alpha\beta} \\
&= \oint_{\partial \mathcal{M}} \epsilon v^\mu \eta_\nu \sqrt{-\hbar} d^3 \chi, \mathfrak{I}_{\mathfrak{B}} = \oint_{\partial \mathcal{M}} d^3 \chi \sqrt{-\hbar} \kappa, \mathfrak{I} \\
&= \frac{1}{2\kappa \int_{\mathcal{M}} d^4 \chi \sqrt{-g} \mathcal{R}} + 1/\kappa \oint_{\partial \mathcal{M}} d^3 \chi \sqrt{-\hbar} \mathfrak{K}
\end{aligned}$$

$$\begin{aligned}
ds^2 &= -\mathcal{N}(r) dt^2 + \mathcal{H}(r) dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2), \mathfrak{K}_{\mathfrak{R}\mathfrak{C}\mathfrak{I}\mathfrak{S}} = \mathcal{R}^{\alpha\beta\gamma\sigma} \mathcal{R}_{\alpha\beta\gamma\sigma}, \mathcal{I}[\mathfrak{g}_{\mu\nu}] \\
&= \frac{1}{2\kappa \int_{\mathcal{M}} d^4 \chi \sqrt{-g} \mathcal{R}} + \frac{1}{\kappa \oint_{\partial \mathcal{M}} d^3 \chi \mathfrak{K} \sqrt{-\hbar}}, ds^2 \\
&= -\left(1 - \frac{\mu}{r}\right) dt^2 + \left(1 - \frac{\mu}{r}\right)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2), M = \frac{4\pi\mu}{\kappa} = \mu/2\mathfrak{G}
\end{aligned}$$



2.2. Modelo Reissner – Nordström.

$$\begin{aligned}
 \mathcal{I}[g_{\mu\nu}\Lambda_\mu] &= \frac{1}{2\kappa \int_{\mathcal{M}} d^4\chi \sqrt{-g} \left(\mathcal{R} - \frac{1}{4F^{\mu\nu}F_{\mu\nu}} \right)} + \frac{1}{\kappa \oint_{\partial\mathcal{M}} d^3\chi \mathfrak{K} \sqrt{-h}}, \mathcal{R}_{\mu\nu} - \frac{1}{2\mathcal{R}g_{\mu\nu}} \\
 &= \frac{1}{2\mathfrak{S}_{\mu\nu}^{\mathfrak{M}} \nabla_\mu F^{\mu\nu}}, \mathfrak{S}_{\mu\nu}^{\mathfrak{M}} = F_{\mu\alpha} F_\nu^\alpha - \frac{1}{4g_{\mu\nu} F^2}, \Lambda \equiv \Lambda_\mu d\chi_\mu = \left(\frac{q}{r} - \frac{q}{r_+} \right) dt, F \\
 &= -\frac{q}{r^2 dr} \wedge dt, ds^2 = -f(r) dt^2 + f(r)^{-1} dr^2 + r^2 (d\theta^2 + \sin^2\theta d\varphi^2), f(r) \\
 &= 1 - \frac{\mu}{r} + \frac{q^2}{4r^2} = \frac{(r-r_-)(r-r_+)}{r^2}, \mathcal{M} = \frac{4\pi\mu}{\kappa} = \frac{\mu}{2\mathfrak{G}}, \mathfrak{Q} \equiv \frac{1}{\kappa \oint_{-\infty}^\infty d^2 \star \mathcal{F}} = -\frac{q}{4\mathfrak{G}}, r_\pm \\
 &= \mathfrak{G}(\mathfrak{M} \pm \sqrt{\mathfrak{M}^2 - 4\mathfrak{Q}^2}), \Phi = \Lambda_t |_{\Delta_{r=\infty}} - \Lambda_t |_{\Delta_{r=r_+}} = 4\mathfrak{G}\mathfrak{Q}/r \blacksquare
 \end{aligned}$$

2.3. Modelo anti – de Sitter.

$$\begin{aligned}
 ds^2 &= \frac{r^2}{\ell^2(-dt^2 + \ell^2 d\Sigma_\kappa^2)}, ds^2 = -N(r)dt^2 + \mathcal{H}(r)dr^2 + \delta(r)d\Sigma_\kappa^2 \\
 &= \langle \frac{d\theta^2 + \sin^2\theta d\varphi^2}{\ell^2 \sum_{i=1}^2 d\chi_i^2} \quad \begin{matrix} \propto \kappa = +1 \\ \propto \kappa = 0 \\ \div \kappa = -1 \end{matrix} \rangle, d\Sigma_\kappa^2 = \frac{dy^2}{1} - \kappa y^2 + (1 + \kappa y^2) dz^2
 \end{aligned}$$

2.4. Modelo Schwarzschild-AdS.

$$\begin{aligned}
 \mathcal{I}[g_{\mu\nu}] &= \frac{1}{2\kappa \int_{\mathcal{M}} d^4\chi \sqrt{-g} (\mathfrak{R} - 2\Lambda)} + \frac{1}{\kappa \oint_{\partial\mathcal{M}} d^3\chi \mathfrak{K} \sqrt{-h}}, \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathcal{R}} = -\Lambda g_{\mu\nu}, ds^2 \\
 &= -\left(\kappa - \frac{\mu}{r} + \frac{r^2}{\ell^2} \right) dt^2 + \left(\kappa - \frac{\mu}{r} + \frac{r^2}{\ell^2} \right)^{-1} dr^2 + r^2 d\Sigma_\kappa^2, \kappa - \frac{\mu}{r_h} + \frac{r_h^2}{\ell^2}
 \end{aligned}$$

2.5. Modelo Escalar Simple.

$$\begin{aligned}
 \mathcal{I}[g_{\mu\nu}, \phi] &= \frac{\int_{\mathcal{M}} d^4\chi \sqrt{-g} \left[\frac{\mathfrak{R}}{2\kappa} - \frac{1}{2\partial_\mu \phi \partial^\mu \phi} - \mathcal{V}(\phi) \right]}{\sqrt{-g} \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi)} - \frac{\frac{\partial \mathcal{V}}{\partial \phi} d\mathcal{V}}{d\phi |_{\psi_{\phi=0}}}, \mathcal{V}(0) \\
 &= -\frac{3}{\kappa \ell^2}, \frac{d^2 \mathcal{V}}{d\phi^2 |_{\psi_{\phi=0}}}, \mathbb{V}(\phi)_{\mathfrak{U} \mathfrak{b} \mathfrak{S}} = \left(-\frac{1}{\ell^2} + \alpha \phi \right) (4 + 2 \cos \hbar \phi) - 6\alpha \sin \hbar \phi
 \end{aligned}$$

2.6. Modelo Escalar Neutro.



$$\begin{aligned}
\varepsilon_{\mu\nu} &= \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathfrak{R}} - \kappa \mathcal{T}_{\phi\mu\nu}^{\varphi} \mathcal{T}_{\phi\mu\nu}^{\varphi} = \partial_{\mu}\phi\partial_{\nu}\phi - g_{\mu\nu}\left(\frac{1}{2(\partial\psi)^2} + \mathfrak{B}(\phi)\right), ds^2 \\
&= \Omega(\chi)\left[-\mathfrak{f}(\chi)dt^2 + \frac{\eta^2 d\chi^2}{\mathfrak{f}(\chi)} + \frac{dy^2}{\ell^2} + \frac{dz^2}{\ell^2}\right]\mathfrak{E}_t^t - \mathfrak{E}_{\chi}^{\chi} = 0 \rightarrow \phi'^2 \\
&= 3\Omega'^2 - \frac{2\Omega''\Omega}{\Omega^2}, \mathfrak{E}_t^t - \mathfrak{E}_{\mathcal{Y}}^{\mathcal{Y}} = 0 \rightarrow \mathfrak{f}'' + \frac{\Omega'\mathfrak{f}'}{\Omega} = 0, \mathfrak{E}_t^t - \mathfrak{E}_{\mathcal{Y}}^{\mathcal{Y}} = 0 \rightarrow \mathcal{V}(\phi) \\
&= -\frac{1}{\Omega^2\eta^2(\mathfrak{f}\Omega'' + \mathfrak{f}'\Omega')}, \Omega(\chi) = \frac{\nu^2\chi^{\nu-1}}{\eta^2(\chi^{\nu}-1)^2}, \phi'^2 \\
&= \frac{(\nu-1)^2}{\chi^2} - \frac{4\nu(\nu-1)\chi^{\nu-2}}{\chi^{\nu}} - 1 + \frac{4\nu^2\chi^{\nu-1}}{(\chi^{\nu}-1)^2} + \frac{2(\nu-1)}{\chi^2} \\
&+ \frac{4\nu(1-\nu-\chi^{\nu})\chi^{\nu-2}\chi^{\nu-2}}{(\chi^{\nu}-1)^2}, \phi'^2 = \nu^2 - \frac{1}{2\kappa\chi^2} \\
&\rightarrow \int_{\phi}^{\phi=0} d\phi = \sqrt{\nu^2 - \frac{1}{2\kappa}} \int_{\chi}^1 \frac{d\chi}{\chi}, \phi(\tau) \ln \chi, \ell_v^{-1} = \sqrt{\nu^2 - \frac{1}{2\kappa}} (\mathfrak{f}'\Omega)' f(x) \\
&= \frac{c_4\eta^2}{\nu^2 \int \frac{(\chi^{\nu}-1)^2}{\chi^{\nu-1}d\chi} + c_1}, f(x) = c_1 + \frac{c_4\eta^2}{\nu^2 \left(\frac{\chi^{2+\nu}}{2} + \nu + \frac{\chi^{2-\nu}}{2} - \nu - \chi^2\right)}, f(x) \\
&= \frac{1}{\ell^2} + \alpha \left[\frac{1}{\nu^2} - 4 - \frac{\chi^2}{\nu^2 \left(1 + \frac{\chi^{-\nu}}{\nu} - 2 - \frac{\chi^{\nu}}{\nu} + 2\right)} \right], \mathcal{V}(\phi) \\
&= \frac{\Lambda(\nu^2 - 4)}{6\kappa\nu^2 \left[\nu - \frac{1}{\nu} + 2e^{-\phi\ell_{\nu}(\nu+1)} + \nu + \frac{1}{\nu} - 2e^{\phi\ell_{\nu}(\nu-1)} + 4\nu^2 - \frac{1}{\nu^2} - 4e^{-\phi\ell_{\nu}} \right]} \\
&+ \alpha \\
&/\kappa\nu^2 \left[\nu - \frac{1}{\nu} + 2\sin\hbar\phi\ell_{\nu}(\nu+1) - \nu + \frac{1}{\nu} \right. \\
&\left. - 2\sin\hbar\phi\ell_{\nu}(\nu-1) + 4\nu^2 - \frac{1}{\nu^2} - 4\sin\hbar\phi\ell_{\nu} \right] \\
\mathcal{V}(\phi) &= \frac{\Lambda}{\kappa} - \frac{\phi^2}{\ell^2} + \frac{\kappa\Lambda}{18(\nu^2-3)\nu^2} - 1\phi^4 - \frac{\ell_{\nu}^3}{90(\Lambda\nu^2 - 4\Lambda - 6\alpha)\tau^5} + \mathcal{O}|\phi|^6
\end{aligned}$$

$$\begin{aligned}
ds^2 &= \Omega(\chi) \left[-f(x)dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\theta^2 + \sin^2\theta d\varphi^2 \right], f(x) \\
&= \frac{1}{\ell^2} + \alpha \left[\frac{1}{v^2} - 4 - \frac{\chi^2}{v^2 \left(1 + \frac{\chi^{-v}}{v} - 2 - \frac{\chi^v}{v} + 2 \right)} \right] + \chi/\Omega(\chi)
\end{aligned}$$

2.7. Modelo Escalar Eléctricamente Cargado.

$$\begin{aligned}
\mathfrak{I}[g_{\mu\nu}, \Lambda_\mu \phi] &= \frac{1}{16\pi\mathfrak{G}_N \int d^4\chi \sqrt{-g} \left[\mathfrak{R} - \frac{1}{4e^{\gamma\phi} F^2} - \frac{1}{2\partial_\mu \phi \partial^\mu \phi} - \mathcal{V}(\phi) \right]}, \nabla_\mu (e^{\gamma\phi} F^{\mu\nu}) \\
&= \frac{1}{\sqrt{-g} \partial_\mu (\sqrt{-g} g^{\mu\nu} \partial_\nu \phi)} - \frac{\partial \mathcal{V}}{\partial \phi} - \frac{1}{4\gamma e^{\gamma\phi} F^2}, \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu} \mathcal{R}} = \frac{1}{2 \left[\mathcal{T}_{\mu\nu}^\phi \mathcal{T}_{\mu\nu}^{\mathfrak{E}\mathfrak{M}} \right]}, \mathcal{T}_{\mu\nu}^\phi \\
&= \partial_\mu \phi \partial_\nu \phi - g_{\mu\nu} \left[\frac{1}{2(\partial\phi)^2} + \mathcal{V}(\phi) \right], \mathcal{T}_{\mu\nu}^{\mathfrak{E}\mathfrak{M}} = e^{\gamma\phi} \left(F_{\mu\alpha} F_\nu^\alpha - \frac{1}{4g_{\mu\nu} F^2} \right)
\end{aligned}$$

2.8. Modelo Holográfico en espacios cuánticos curvos.



$$\begin{aligned}
J_g &= -\frac{1}{8\pi\mathfrak{G}_N\oint_{\partial\mathcal{M}}^\lambda d^3\chi\mathfrak{K}\sqrt{-\hbar}\Xi(\ell,\mathcal{R},\nabla\mathfrak{R})}, ds^2 \\
&= -\left(\kappa + \frac{r^2}{\ell^2}\right) dt^2 + \left(\kappa + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2 d\Sigma_\kappa^2 h_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -\left(\kappa + \frac{\mathcal{R}^2}{\ell^2}\right) dt^2 + \mathcal{R}^2 d\Sigma_\kappa^2, ds^2 = \frac{r_\beta^2}{\ell^2(-dt^2 + \ell^2 d\Sigma_\kappa^2)}, ds^2 \\
&= -\left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right) dt^2 + \left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2), \mathfrak{T} = \frac{f'}{|4\pi_{r+}} \\
&= \beta^{-1} = \frac{1}{4\pi\left(3r_+^2 + \frac{\ell^2}{\ell_{r+}^2}\right)}, J_{\mathfrak{B}\cup\mathfrak{R}}^\mathfrak{E} = \frac{12\pi\beta}{\kappa\ell^2 \int_{r_+}^{\mathcal{R}} r^2 dr = \frac{4\pi\beta}{\kappa\ell^2(\mathcal{R}^3 - r_+^3)}}, h_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -\left(1 - \frac{\mu}{\mathcal{R}} + \frac{\mathcal{R}^2}{\ell^2}\right) dt^2 + \mathcal{R}^2(d\theta^2 + \sin^2\theta d\varphi^2), \eta_\alpha = \frac{\delta_\alpha^r}{\sqrt{g^{rr}}}, \kappa_{\alpha\beta} = \frac{\sqrt{g^{rr}}}{2\partial_r h_{\alpha\beta}}, \mathfrak{K} \\
&= \frac{1}{\mathcal{R}^2\ell^2\left(1 - \frac{\mu}{\mathcal{R}} + \frac{\mathcal{R}^2}{\ell^2}\right)^{\frac{1}{2}}\left(-\frac{3\ell^2\mu}{2} + 3\mathcal{R}^2 + 2\mathcal{R}\ell^2\right)}, J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2\left(-\frac{3\ell^2\mu}{2} + 3\mathcal{R}^2 + 2\mathcal{R}\ell^2\right)}, J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} = J_{\mathfrak{B}\cup\mathfrak{R}}^\mathfrak{E} + J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2\left(-\frac{3\ell^2\mu}{2} + 2\mathcal{R}^2 + 2\mathcal{R}\ell^2 - r_+^3\right)}, ds^2 \\
&= -\left(1 + \frac{r^2}{\ell^2}\right) dt^2 + \left(1 + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2\theta d\varphi^2), J_{\mathfrak{A}\cup\mathfrak{S}}^\mathfrak{E} = J_{\mathfrak{B}\cup\mathfrak{R}}^\mathfrak{E} + J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta_0}{\kappa\ell^2(-2\mathcal{R}^3 + 2\mathcal{R}\ell^2)}, \beta_0 \sqrt{1 + \frac{\mathcal{R}^2}{\ell^2}} = \beta \sqrt{1 + \frac{\mathcal{R}^2}{\ell^2} - \frac{\mu}{\mathfrak{R}}}, J^\mathfrak{E} = J_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} - J_{\mathfrak{A}\cup\mathfrak{S}}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2\left[\left(\frac{3\ell^2\mu}{2} - 2\mathcal{R}^3 - 2\mathcal{R}\ell^2 - r_+^3\right) - \frac{\beta_0}{\beta(-2\mathcal{R}^3 - 2\ell^2\mathfrak{R})}\right]}, F = \beta^{-1}J^\mathfrak{E} \\
&= 4\pi/\kappa\ell^2\left(\frac{\ell^2\mu}{2} - r_+^3\right)
\end{aligned}$$



$$\begin{aligned}
\mathcal{I}_g &= -\frac{1}{\kappa \int_{\partial\mathcal{M}}^\infty d^3\chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{L}\mathcal{R}}{2}\right), \mathcal{I}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^\mathfrak{E}} + \mathcal{I}_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} = \frac{4\pi\beta}{\kappa\ell^2 \left(\frac{3\ell^2\mu}{2} - 2r_\beta^3 - 2r_\beta\ell^2 - r_+^3\right)}, \mathcal{I}_\mathfrak{G}^\mathfrak{E} \\
&= \frac{4\pi\beta}{\kappa\ell^2 \left(1 + \frac{\ell^2}{\mathcal{R}^2} - \frac{\mu\ell^2}{\mathcal{R}^3}\right)^{\frac{1}{2}} |(2\mathcal{R}^3 + \kappa\ell^2\mathcal{R})|_{\mathcal{R}=r_\beta}} = \frac{4\pi\beta}{\kappa\ell^2 (2r_\beta^3 - 2\ell^2 r_\beta - \mu\ell^2)}, \mathcal{I}_\mathfrak{G}^\mathfrak{E} \\
&= \mathcal{I}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^\mathfrak{E} + \mathcal{I}_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} + \mathcal{I}_\mathfrak{G}^\mathfrak{E} = \frac{4\pi\beta}{\kappa\ell^2 \left(\frac{\ell^2\mu}{2} - r_+^3\right)}, \mathcal{E} = -\frac{\mathfrak{S}^2 \partial \mathcal{I}^\mathfrak{E}}{\partial \mathcal{I}} = \frac{\mu}{2\mathfrak{G}}, F_{\mathfrak{S}\mathfrak{U}\mathfrak{L}\mathfrak{S}} \\
&= \frac{4\pi}{\kappa\ell^2 \left(\frac{\ell^2\mu}{2} - r_+^3\right)}, \mathcal{I}_{\mathfrak{S}\mathfrak{U}\mathfrak{L}\mathfrak{S}} = \left(1 + \frac{3r_+^2}{\ell^2}\right)
\end{aligned}$$

2.8.1. Modelo de Brown-York.

$$\begin{aligned}
\int_{\partial\mathcal{M}}^\delta d^3\chi \hbar^{\alpha\beta} \mathcal{T}_{\alpha\beta}, ds^2 &= \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta = -N(\mathcal{R})dt^2 + \delta(\mathcal{R})d\Sigma_\kappa^2, \tau^{\alpha\beta} \equiv \frac{\frac{2}{\sqrt{-\hbar}\delta\mathcal{I}}}{\delta\hbar_{\alpha\beta}}, ds^2 \\
&= -\left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right)dt^2 + \left(1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}\right)^{-1} dr^2 + r^2 d\Omega^2, ds^2 \\
&= -\left(1 - \frac{\mu}{\mathcal{R}} + \frac{\mathcal{R}^2}{\ell^2}\right)dt^2 + \mathcal{R}^2 d\Omega^2, \mathcal{I} \\
&= \frac{1}{2\kappa \int_{\partial\mathcal{M}}^\delta d^4\chi \sqrt{-g} (\mathfrak{R} - 2\Lambda)} + \frac{1}{\kappa \int_{\partial\mathcal{M}}^\delta d^3\chi \sqrt{-\hbar} \mathfrak{K}} - \frac{1}{\kappa \int_{\partial\mathcal{M}}^\delta d^3\chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{L}\mathcal{R}}{2}\right)}, \tau^{\alpha\beta} \\
&= \frac{1}{8\pi\mathfrak{G} \left(\mathfrak{K}_{\alpha\beta} - \hbar_{\alpha\beta}\mathfrak{K} - \frac{2}{\ell\hbar_{\alpha\beta}} + \mathcal{I}\mathfrak{S}_{\alpha\beta}\right)}, ds_{borde}^2 = \frac{\mathfrak{R}^2}{\ell^2(-dt^2 + \ell^2 d\Omega^2)}, ds_{dualidad}^2 \\
&= \gamma_{\alpha\beta} d\chi^\alpha d\chi^\beta = -dt^2 + \ell^2 d\Omega^2, \langle \tau_{dualidad}^{\alpha\beta} \rangle = \lim_{\mathcal{R} \rightarrow \infty} \frac{\mathcal{R}}{\ell} \tau_{\alpha\beta} \\
&= \mu/16\pi\mathfrak{G}\ell^2 (3\delta_\alpha^0 \delta_\beta^0 + \gamma_{\alpha\beta})
\end{aligned}$$

2.9. Modelo Hamiltoniano.



$$\begin{aligned}
\mathfrak{I}[g_{\mu\nu}, \Lambda_\mu \phi] &= \int_{\partial\mathcal{M}}^\delta d^4\chi \sqrt{-g} \left[\frac{\mathcal{R}}{2\kappa} - \frac{1}{2(\partial\phi)^2} - \mathcal{V}(\phi) \right] + \frac{1}{\kappa \int_{\partial\mathcal{M}}^\delta d^3\chi \sqrt{-h}}, \mathcal{H}_\perp \\
&= \frac{2\kappa}{\sqrt{g} \left[\pi^{ij} \pi_{ij} - \frac{1}{2(\pi_j^i)^2} \right]} - \frac{1}{2\kappa \sqrt{g}^{(3)} \mathcal{R}} + \frac{1}{2 \left(\frac{\pi_{\phi^2}}{\sqrt{g}} + \sqrt{g} g^{ij} \phi_i \phi_j \right)} + \sqrt{g} \mathcal{V}(\phi), \mathcal{H}_i \\
&= -\pi_j^i|_j + \pi_\phi \phi_\psi d\omega, ds^2 = (N^\perp)^2 dt^2 + g_{ij} (d\chi^i + N^i dt) (d\chi^j + N^j dt), \mathcal{H}[\xi] \\
&= \int_{\partial\mathcal{M}}^\delta d^3\chi (\xi^\perp \mathcal{H}_\perp + \xi^i \mathcal{H}_i) + \mathfrak{Q}[\xi], \delta\mathfrak{Q}[\xi] \\
&= \oint d^2\delta_\ell \left[\frac{\mathfrak{G}^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij}|_\kappa - \xi_\kappa^\perp \delta g_{ij}) + 2\xi_\kappa \delta\pi^{\kappa\ell} + (2\xi^\kappa \pi^{j\ell} - \xi^\ell \pi^{j\kappa}) \delta g_{j\kappa} \right. \\
&\quad \left. - (\sqrt{g} \xi^\perp g^{\ell j} \phi_j + \xi^\ell \pi_\phi) \delta\phi \right], \mathfrak{G}^{ijkl} \equiv \frac{1}{2\sqrt{g} (g^{i\kappa} g^{j\ell} + g^{i\ell} g^{j\kappa} - 2g^{ij} g^{\kappa\ell})} \\
\delta\mathcal{M} \equiv \delta\mathfrak{Q}[\partial_t] &= \oint d^2\delta_\ell \left[\frac{\mathfrak{G}^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij}|_\kappa - \xi_\kappa^\perp \delta g_{ij}) - \sqrt{g} \xi^\perp g^{\ell j} \phi_j \delta\phi \right], \delta\mathcal{M} = \delta\mathcal{M}_\mathfrak{G} + \delta\mathcal{M}_\phi, \delta\mathcal{M}_\mathfrak{G} \\
&= \oint d^2\delta_\ell \frac{\mathfrak{G}^{ijkl}}{2\kappa} (\xi^\perp \delta g_{ij}|_\kappa - \xi_\kappa^\perp \delta g_{ij}), \delta\mathcal{M}_\phi = -\oint d^2\delta_\ell \sqrt{g} \xi^\perp g^{\ell j} \phi_j \delta\phi
\end{aligned}$$

2.10. Termodinámica de agujeros negros cuánticos en espacios curvos.

2.10.1. Espacio – Tiempo de Rindler.



$$\begin{aligned}
ds^2 &= -dt^2 + d\chi^2 + d\gamma^2 + d\mathfrak{z}^2, ds^2 = -\alpha^2 \rho^2 d\tau^2 + d\rho^2 + d\gamma^2 + d\mathfrak{z}^2, ds^2 \\
&= N(r) dt_{\mathfrak{E}}^2 + \mathcal{H}(r) dr^2 + \delta(r) d\Sigma_{\kappa}^2, ds^2 = \frac{g_{rr} 4N}{[(N)']^2 \left[\frac{\rho^2 [(N)']^2}{4N g_{rr} d\tau_{\mathfrak{E}}^2} + d\rho^2 \right] \mathcal{T}} = \frac{1}{\beta} \\
&= \frac{[(N)']^2}{4\pi \sqrt{N^2 g_{rr}} \Big|_{\mathcal{H}} \mathcal{T}_{\delta ch - \mathfrak{U} \delta}} = \frac{1}{4\pi r_+ \left(1 + \frac{3r_+^2}{\ell^2} \right) \mathcal{T}_{\mathfrak{R} \eta - \mathfrak{U} \delta}} = \frac{1}{4\pi r_+ \left(1 + \frac{3r_+^2}{\ell^2} - \frac{q^2}{4r_+^2} \right)}, ds^2 \\
&= \Omega(x) \left[-f(x) dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\theta^2 + d\Sigma_{\kappa}^2 \right], f(x) \\
&= \frac{1}{\ell^2} + \alpha \left[\frac{1}{v^2} - 4 - \frac{\chi^2}{v^2 \left(1 + \frac{\chi^{-v}}{v} - 2 - \frac{\chi^v}{v} + 2 \right)} \right] + \frac{\kappa \chi}{\Omega(x)}, f' \Omega(x) \\
&= \frac{\alpha}{\eta^2} + 2\kappa + \kappa v \chi^v + \frac{1}{\chi^v} - 1, \mathcal{T} = \frac{f'}{4\pi \eta} \Big|_{\chi_h} \\
&= \frac{\frac{1}{f'}}{4\pi \eta \Omega(\chi_h) \left(\frac{\alpha}{\eta^2} + 2\kappa + \kappa v \chi_h^v + \frac{1}{\chi_h^v} - 1 \right)}
\end{aligned}$$

2.10.2. Transiciones de fase de agujeros negros cuánticos en espacios curvos.



$$\begin{aligned}
ds^2 &= \Omega(x) \left(-f(x) dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\theta^2 + \sin^2 \theta d\varphi^2 \right), \Omega(x) = \frac{1}{\eta^2 (\chi - 1)^2}, f(x) \\
&= \frac{1}{\ell^2} + \frac{1}{3\alpha(\chi - 1)^3} + \eta^2 \chi (\chi - 1)^2, \chi = 1 + \frac{1}{\eta r}, \chi = 1 - \frac{1}{\eta r}, \Omega(x) f(x) = F(r) \\
&= 1 - \frac{\mu}{r} + \frac{r^2}{\ell^2}, \mu = \alpha + \frac{3\eta^2}{3\eta^3}, \mathcal{I}[g_{\mu\nu}] \\
&= \mathcal{I}_{bulk} + \mathcal{I}_{GH} \\
&= \frac{1}{\kappa \int_{\partial\mathcal{M}}^\delta d^3\chi \sqrt{-\hbar} \left(\frac{2}{\ell} + \frac{\mathcal{R}\ell}{2} \right), \varepsilon_t^t - \varepsilon_{\langle\tau|\sigma|\rho\rangle}^{\langle\varphi|\chi|\psi\rangle} = 0 \Rightarrow 0 = f'' + \frac{\Omega' f'}{\Psi} + 2\eta^2, \varepsilon_t^t} \\
&+ \varepsilon_{\langle\tau|\sigma|\rho\rangle}^{\langle\varphi|\chi|\psi\rangle} = 0 \Rightarrow 2\kappa\mathcal{V}(\phi) = -\frac{(f\Omega'' + f'\Omega')}{\Psi^4 \eta^4} + \frac{2}{\Omega}, \mathcal{I}_{bulk}^{\mathfrak{E}} \\
&= \frac{4\pi\beta}{\eta^3 \kappa \ell^2 \left[-\frac{1}{(\chi_\beta - 1)^3} + \frac{1}{(\chi_{\hbar} - 1)^3} \right]} = \frac{4\pi\beta}{\kappa \ell^2 (r_\beta^3 - r_{\hbar}^3)}, ds^2 = \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= \Omega(x) [-f(x) dt^2 + d\theta^2 + \sin^2 \theta d\varphi^2], \eta_\alpha = \frac{\delta_\alpha^\chi}{\sqrt{g^{\chi\chi}} \partial_\chi \hbar_{\alpha\beta}}, \mathcal{I}_{GH}^{\mathfrak{E}} \\
&= -\frac{2\pi\beta}{\kappa \left[-\frac{6}{\ell^2 \eta^4 (\chi - 1)^3} - \frac{4}{\eta (\chi_\beta - 1)} - \left(\alpha + \frac{3\eta^4}{\eta^3} \right) \right] \Big|_{\chi_\beta}} = -\frac{2\pi\beta}{\kappa \left(\frac{6r_\beta^3}{\ell^2} + 4r_\beta - 3\mu \right)}, \mathcal{I}_g^{\mathfrak{E}} \\
&= \frac{2\pi\beta}{\kappa \left[\frac{4}{\ell^2 \eta^4 (\chi_\beta - 1)^3} + \frac{4}{\eta (\chi_\beta - 1)} - 2\mu \right]} = 2\pi\beta / \kappa \left(\frac{4r_\beta^3}{\ell^2} + 4r_\beta - 2\mu \right)
\end{aligned}$$

$$\begin{aligned}
\mathcal{J}^{\mathfrak{E}} &= \mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{GH}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} = \frac{4\pi\beta}{\kappa\ell^2 \left[\frac{1}{\eta^3(\chi_h - 1)^3} + \frac{\mu\ell^2}{2} \right]} = \frac{4\pi\beta}{\kappa\ell^2 \left(-r_h^3 + \frac{\mu\ell^2}{2} \right)}, \mathcal{J}^{\mathfrak{E}}_{\langle\varphi|\phi|\psi\rangle\langle\sigma|\tau|\rho\rangle} \\
&= \int_{\partial\mathcal{M}}^{\delta} d^3\chi^3 \sqrt{\hbar^{\mathfrak{E}}} \left(\frac{\langle\varphi|\phi|\psi\rangle\langle\sigma|\tau|\rho\rangle}{2\ell} - \ell_v/6\ell \langle\varphi|\phi|\psi\rangle\langle\sigma|\tau|\rho\rangle^4 \right) \\
&= \frac{4\pi\beta}{\kappa \left[-v^2 - \frac{1}{4\ell^2\eta^3(\chi_\beta - 1)} + v^2 - \frac{1}{3\ell^2\eta^3} \right]}, \mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{surf}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} \\
&= -\frac{1}{\mathcal{T} \left(\frac{\Lambda\Gamma}{4\mathfrak{G}} \right)} + \frac{4\pi\beta}{\kappa \left[v^2 - \frac{1}{4\ell^2\eta^3(\chi_\beta - 1)} + 12\eta^2\ell^2 + 4\alpha\ell^2 - 4v^2 + \frac{4}{12\ell^2\eta^3} \right]}, \mathcal{J}^{\mathfrak{E}} \\
&= \beta \left(-\frac{\Lambda\Gamma}{4\mathfrak{G}} + \frac{4\pi}{\kappa} 3\eta^2 + \frac{\alpha}{3\eta^3} \right), \mathcal{M} = \frac{1}{2\mathfrak{G} \left(\alpha + \frac{3\eta^2}{3\eta^3} \right)}, \mathcal{T} \\
&= \frac{\frac{f'(x)}{4\pi\eta} \Big|_{\chi=\chi_h} 1}{4\pi\eta\Omega(\chi_h) \left[\frac{\alpha}{\eta^2} + 2 + v\chi_h^\nu + 1 \frac{\chi_h^\nu}{\chi_h} - 1 \right]}, \delta = \frac{A}{4G} = 4\pi\eta\Omega(\chi_h)/4G \\
\frac{\partial\mathcal{M}}{\partial\eta} \frac{d\eta}{d\chi_h} &= \mathcal{T} \left(\frac{\partial\mathcal{S}}{\partial\chi_h} + \frac{\partial\mathcal{S}}{\partial\eta} \frac{d\eta}{d\chi_h} \right), \mathcal{M} = -\frac{1}{2G \left(\alpha + \frac{3\eta^2}{3\eta^3} \right)}, \mathcal{T} = \frac{1}{4\pi\eta\Omega(\chi_h) \left[\frac{\alpha}{\eta^2} + 2 + v\chi_h^\nu + \frac{1}{\chi_h^\nu} - 1 \right]}
\end{aligned}$$

2.10.3. Solitón - AdS para agujeros negros cuánticos en espacios curvos.

$$\begin{aligned}
\mathcal{I}[g_{\mu\nu}] &= \int_{\mathcal{M}}^{\delta} d^4\chi(\mathcal{R} - 2\Lambda)\sqrt{-g} + 2 \int_{\partial\mathcal{M}}^{\delta} d^3\chi \mathfrak{K} \sqrt{-\hbar} - \int_{\partial\mathcal{M}}^{\delta} \frac{d^3\chi^4}{\ell\sqrt{-\hbar}}, ds^2 \\
&= -\left(-\frac{\mu_\beta}{r} + \frac{r^2}{\ell^2} \right) dt^2 + \left(-\frac{\mu_\beta}{r} + \frac{r^2}{\ell^2} \right)^{-1} dr^2 + \frac{r^2}{\ell^2} (d\chi_1^2 d\chi_2^2), \mathcal{J}_\beta^\xi \\
&= \frac{2\mathcal{L}\mathcal{L}_\beta\beta_\varsigma}{\ell^4 \left(-r_h^3 + \frac{\mu_\beta\ell^2}{2} \right)} = -\frac{\mathcal{L}\mathcal{L}_\beta\beta_\varsigma r_h^3}{\ell^4}, \mathfrak{S} = \beta_\varsigma^{-1} = \frac{(-g_{tt})'}{4\pi} \Big|_{r=r_h} = \frac{3r_h}{4\pi\ell^2}, \mathfrak{E} = -\frac{\mathfrak{T}^2 \partial\mathcal{J}_\beta^\mathfrak{E}}{\partial\mathfrak{T}} \\
&= \frac{2\mathcal{L}\mathcal{L}_\beta\mu_\beta}{\ell^4}, \mathfrak{S} = -\frac{\partial(\mathcal{J}_\beta^\mathfrak{E} \mathfrak{T})}{\mathfrak{T}} = \frac{\mathcal{L}\mathcal{L}_\beta r_h^2}{4\ell^2\mathfrak{G}} = \mathcal{A}/4G
\end{aligned}$$

$$\begin{aligned}
ds_{dualidad}^2 &= \frac{\ell^2}{\mathcal{R}^2} ds^2 = \gamma_{\alpha\beta} d\chi^\alpha d\chi^\beta = -dt^2 + d\chi_1^2 + d\chi_2^2, \langle \tau_{\alpha\beta}^{dualidad} \rangle \\
&= \lim_{\mathcal{R} \rightarrow \infty} \frac{\mathcal{R}}{\ell} \tau_{\alpha\beta} = \frac{\mu_\beta}{16\pi \mathfrak{G}_N \ell^2} [3\delta_\alpha^0 \delta_\beta^0 + \gamma_{\alpha\beta}], \mathfrak{E} = Q_{\xi t} \\
&= \int d\Sigma^i \tau_{ij} \xi^j = \frac{\mathcal{L} \mathcal{L}_\beta}{\ell^2 \kappa} \left[\mu_\beta + \frac{\ell^2}{4\mathcal{R}} + \mathcal{O}(\mathcal{R}^{-2}) \right], ds^2 \\
&= \frac{r^2}{\ell^2} d\tau^2 + \left(-\frac{\mu_\delta}{r} + \frac{r^2}{\ell^2} \right)^{-1} dr^2 + \left(-\frac{\mu_\delta}{r} + \frac{r^2}{\ell^2} \right) d\theta^2 + \frac{r^2}{\ell^2} d\chi_2^2 - \frac{\mu_\delta}{r} + \frac{r_\delta^2}{\ell^2}, \mathcal{L}_\delta \\
&= \frac{4\varpi \sqrt{g_{\theta\theta} g_{rr}}}{(g_{\theta\theta})'} \Big|_{r=r_\delta} = \frac{4\pi \ell^2}{3r_\delta}, J_\delta^\xi = \frac{\mathcal{L} \mathcal{L}_\delta \beta_\delta \mu_\delta}{\ell^2} \Delta \mathcal{J} = J_\beta^\epsilon - J_\delta^\epsilon \\
&= \frac{\mathcal{L}}{2\pi \kappa \ell^4} \left(\frac{4\pi \ell^2}{3} \right)^3 \mathcal{L}_\beta \beta_\varsigma (\mathcal{L}_\delta^{-3} - \xi \beta_\varsigma^{-3}) = \frac{\frac{\mathcal{L}}{2\kappa \ell^4} \left(\frac{4\pi \ell^2}{3} \right)^3 \mathcal{L}_\beta \beta_\varsigma \left(\frac{1}{\mathcal{L}_\delta^3} - \mathcal{T}^3 \right) \mathcal{A}}{\mathcal{T} \ell^3} \\
&= \mathcal{L} / \ell (4\pi/3)^2 \mathcal{L}_\delta \mathcal{T}
\end{aligned}$$

$$\begin{aligned}
\mathcal{I}[g_{\mu\nu}, \phi] &= \int_{\mathcal{M}} d^4\chi \sqrt{-g} \left[\mathcal{R} - \frac{(\partial\phi)^2}{2} - \mathcal{V}(\phi) \right] + 2 \int_{\partial\mathcal{M}} d^3\chi \mathfrak{K} \sqrt{-h}, \mathcal{V}(\phi) \\
&= \frac{\Lambda(v^2 - 4)}{3v^2 \left[v - \frac{1}{v} + 1e^{-\phi \ell_v(v+1)} + v + \frac{1}{v} - 2e^{\phi \ell_v(v-1)} + 4v^2 - \frac{1}{v^2} - 4e^{\phi \ell_v} \right]} \\
&+ \frac{2\alpha}{v^2 \left[v - \frac{1}{v} + 2 \sin \hbar \phi \ell_v(v+1) - v + \frac{1}{v} - 2 \sin \hbar \phi \ell_v(v-1) + 4v^2 - \frac{1}{v^2} - 4 \sin \hbar \phi \ell_v \right]}
\end{aligned}$$

$$\begin{aligned}
ds^2 &= \frac{\mathfrak{R}^2}{\ell^2 [-dt^2 + d\chi_1^2 + d\chi_2^2] \mathfrak{R}^2} \equiv \frac{1}{\eta^2 (\chi - 1)^2}, ds_{dualidad}^2 = \frac{\ell^2}{\mathfrak{R}^2} ds^2 = \gamma_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -dt^2 + d\chi_1^2 + d\chi_2^2, J_{\mathfrak{E}\mathfrak{S}}^\beta = \beta_\varsigma \left(-\frac{\mathcal{A}\mathcal{T}}{4\mathcal{G}_N} + \frac{2\mathcal{L}\mathcal{L}_\beta}{\ell^2} \alpha \right) = -\frac{\mathcal{L}\mathcal{L}_\beta \alpha \beta_\varsigma}{3\ell^2 \eta^3}, \mathcal{A} \\
&= \frac{\mathcal{L}\mathcal{L}_\beta \Omega(\chi_\hbar)}{\ell^2}, \mathfrak{T} = \frac{\alpha}{4\pi \eta^3 \Omega}, \mathcal{M}_\beta = \frac{2\mathcal{L}\mathcal{L}_\beta \mu_\beta}{\ell^2}, \mu_\beta = \frac{\alpha}{3\eta^3}, ds^2 \\
&= \Psi_\delta(\gamma) \left[-\frac{d\tau^2}{\ell^2} + \frac{\lambda^2 d\chi^2}{f(x) d\theta^2} + \frac{d\chi_2^2}{\ell^2} \right] \Psi_\delta(\gamma) = \frac{9\gamma^2}{\lambda^2 (\gamma^3 - 1)^2}, \mathcal{L}_\delta = \frac{4\pi \lambda}{f'} \Big|_{\gamma=\gamma_\delta} \\
&= \frac{4\pi \lambda^3 \gamma_\delta}{\alpha}
\end{aligned}$$

$$\mathfrak{I}_{soliton}^{\mathfrak{E}} = -\frac{\mathcal{L}\beta_{\delta}\Omega_{\delta}(\chi_{\delta})}{4\ell^2\mathfrak{G}_{\mathbb{N}}} + \frac{2\mathcal{L}\mathcal{L}_{\delta}\beta_{\delta}}{\ell^2}\alpha = -\frac{\mathcal{L}\mathcal{L}_{\delta}\beta_{\delta}}{\ell^2\left(\frac{\alpha}{3\lambda^3}\right)}, \mathcal{M}_{soliton} = -\frac{\mathcal{L}\mathcal{L}_{\delta}\mu_{\delta}}{\ell^2}, \mu_{\delta} = \frac{\alpha}{3\lambda^3}, \mathfrak{I}^{\phi}$$

$$= -\int d^3\chi\sqrt{-\hbar}\left(\frac{\phi^2}{2\ell} - \frac{\ell_{\nu}}{6\ell\phi^3}\right)\tau_{\alpha\beta}^{\phi} = -\frac{\frac{2}{\sqrt{-\hbar}}\delta\mathfrak{I}^{\phi}}{\delta\hbar^{\alpha\beta}}, \tau_{\alpha\beta}$$

$$= -\frac{1}{\kappa\left(\kappa_{\alpha\beta} - \hbar_{\alpha\beta}\kappa + \frac{2}{\ell}\hbar_{\alpha\beta} - \mathfrak{I}\mathbb{E}_{\alpha\beta}\right)} - \frac{\hbar_{\alpha\beta}}{\ell\left(\frac{\phi^2}{2} - \frac{\ell_{\nu}}{6}\phi^3\right)}, \tau_{tt}$$

$$= \frac{\alpha(\chi-1)}{3\lambda^2\ell} + \mathcal{O}[(\chi-1)^2], \tau_{\theta\theta} = \frac{2\alpha(\chi-1)}{3\lambda^2\ell} + \mathcal{O}[(\chi-1)^2], \tau_{\chi_2\chi_2}$$

$$= -\frac{\alpha(\chi-1)}{3\lambda^2\ell} + \mathcal{O}[(\chi-1)^2], ds_{dualidad}^2 = \frac{\ell^2}{\mathcal{R}^2}ds^2$$

$$= -d\tau^2 + d\theta^2 + d\chi_2^2, \langle \tau_{\alpha\beta}^{dualidad} \rangle$$

$$= \lim_{\mathcal{R} \rightarrow \infty} [-1/\lambda\ell(\chi-1)]\tau_{\alpha\beta} = \frac{1}{\ell^2\left(\frac{\alpha}{3\lambda^3}\right)} [-3\delta_{\alpha}^{\theta}\delta_{\beta}^{\theta} + \gamma_{\alpha\beta}]$$

$$\mathcal{M} = \oint_{\Sigma} d^2\gamma\sqrt{\sigma}m^{\alpha}\tau_{\alpha\beta}\xi^{\beta} = \frac{\mathcal{L}\mathcal{L}_{\delta}f^{\frac{1}{2}}\Psi}{\sqrt{-g_{\tau\tau}}(\partial_{\tau})^i\tau_{ij}(\partial_{\tau})^j} = -\frac{\mathcal{L}\mathcal{L}_{\delta}}{\ell^2\left[\frac{\alpha}{3\lambda^3} + \mathcal{O}(\chi-1)\right]}, ds^2 = \sigma_{ij}d\chi^i d\chi^j$$

$$= \Omega(x)[f(x)d\theta^2 + d\chi_2^2/\ell^2]$$

$$r_{\beta}^2 = \frac{\Omega(\chi_h, \eta)}{\ell^2}, r_{\delta}^2 = \frac{\Omega(\chi_{\delta}, \lambda)}{\ell^2}, \mathfrak{E} = \mathfrak{M}_{\beta\hbar} - \mathcal{M}_{soliton} = \frac{\mathcal{L}\mathcal{L}_{\beta}}{\ell^2}(2\mu_{\beta} + \mu_{\delta}), \Delta\Gamma = \beta_{\zeta}^{-1}(\mathcal{I}_{BH}^{\mathfrak{E}} - \mathfrak{I}_{soliton}^{\mathfrak{E}})$$

$$= \frac{\mathcal{T}\mathfrak{L}\alpha}{3\ell^2\left(\frac{\mathcal{L}_{\delta}\beta_{\delta}}{\lambda^3} - \frac{\mathcal{L}_{\beta}\beta_{\zeta}}{\eta^3}\right)\Delta\Gamma} = \frac{4\pi\mathcal{L}\mathcal{L}_{\delta}}{3\ell^2\left[\frac{\Omega(\lambda, \chi_{\delta})}{\mathcal{L}_{\delta}} - \mathcal{T}\Omega(\eta, \chi_h)\right]}$$

$$= \frac{4\pi\mathcal{L}}{3\ell^2\Omega(\lambda, \chi_{\delta})\left(1 - \frac{r_{\beta}^3}{r_{\delta}^3}\right)}, \frac{\mathcal{A}}{\mathcal{T}\ell^3} = \frac{\frac{\alpha\mathcal{L}}{4\pi\ell^5}\beta_{\zeta}^2\mathcal{L}_{\delta}}{\eta^3} = \frac{\mathfrak{L}\mathcal{L}}{\ell}\left(\frac{\lambda}{\eta}\right), \mathcal{L}$$

$$= \frac{16\pi^2}{\alpha^2\ell^4}\left[\frac{9\chi_h^2}{(\chi_h^3-1)^2}\right]^4, \frac{\mathcal{A}}{\mathcal{T}\ell^3} = \frac{\mathfrak{L}\mathcal{L}(\alpha, \ell)}{\ell}r_{\beta}/r_{\delta}$$

2.11. Modelo Computacional de un agujero negro cuántico en espacios curvos.



$$\begin{aligned}
\delta &= \Re \int dt \mathcal{T}r \left(\frac{1}{2 \sum_j (\mathcal{D}_t \chi_l)^2 - \frac{m^2}{2} \sum_l \chi_l^2 + \frac{\lambda}{4} \sum_{l \neq j} [\chi_l \chi_j]^2} \right), \delta \\
&= \Re \int_0^\beta dt \mathcal{T}r \left(\frac{1}{2 (\mathcal{D}_t \chi_l)^2 + \frac{m^2}{2} \chi_l^2 - \frac{\lambda}{4} [\chi_l \chi_j]^2} \right) \\
\delta &= \int dt \mathcal{T}r \left(\frac{1}{2} (\mathcal{D}_t \chi_l)^2 - \frac{m^2}{2} \chi_l^2 + \frac{g^2}{4} [\chi_l \chi_j]^2 \right), \delta_\kappa \\
&= \mathbb{N}_\alpha \sum_{t=1}^{\eta_t} \mathcal{T}r \left(\frac{1}{2 (\mathcal{D}_t \chi_l)_{l,t}^2 + \frac{m^2}{2} \chi_{l,t}^2 - \frac{\lambda}{4} [\chi_{l,t} \chi_{j,t}]^2} \right), (\mathcal{D}_t \chi_l)_{l,t} \\
&= \frac{1}{\alpha \left(-\frac{1}{2} \mathfrak{U}_t \mathfrak{U}_{t+\alpha} \chi_{l,t+2\alpha} \mathfrak{U}_{t+\alpha}^\dagger \mathfrak{U}_t^\dagger + 2 \mathfrak{U}_t \chi_{l,t+\alpha} \mathfrak{U}_t^\dagger - \frac{3}{2} \chi_{l,t} \right)}, \langle f \rangle \\
&\equiv \int \frac{d\chi d\mathfrak{U} f(\chi, \mathfrak{U}) e^{-\delta_\kappa(\chi, \mathfrak{U})}}{\int d\chi d\mathfrak{U} e^{-\delta_\kappa(\chi, \mathfrak{U})}} = \lim_{\mathfrak{K} \rightarrow \infty} \frac{1}{\mathfrak{K}} \sum_{\kappa=1}^{\mathfrak{K}} f(\chi^{(\kappa)} \mathfrak{U}^{(\kappa)}), \langle \kappa \rangle = \left\langle \frac{1}{2} \sum_l \frac{\chi_l \partial \mathcal{V}}{\partial \chi_l} \right\rangle, \varepsilon \\
&= \left\langle \frac{1}{\beta} \int_0^\beta dt \left(\mathcal{V} + \frac{1}{2} \sum_l \frac{\chi_l \partial \mathcal{V}}{\partial \chi_l} \right) \right\rangle \\
&= \frac{\mathcal{N}}{\beta \int_0^\beta dt \left(m^2 \chi_l^2 - \frac{3\lambda}{4} [\chi_l, \chi_j]^2 \right)}, \varepsilon_\boxplus = \left\langle \frac{\mathcal{N}}{\eta_t \sum_{t=1}^{\eta_t} \left(m^2 \chi_{l,t}^2 - \frac{3\lambda}{4} [\chi_{l,t}, \chi_{j,t}]^2 \right)} \right\rangle, \mathcal{F}(\mathcal{T}, \eta_t) \\
&= \xi + \sum_{l=1}^{\eta_\rho} \alpha_l \left(\frac{1}{\mathcal{T} \eta_t} \right), \mathcal{F}(\mathcal{T}, \eta_t) = \xi(\mathfrak{T}) + \sum_l^{\eta_\rho} \alpha_l \left(\frac{1}{\eta_t} \right)^l
\end{aligned}$$

$$\begin{aligned}
\hat{\mathcal{H}} &= \mathcal{T}r \left(\frac{1}{2} \hat{\rho}_i^2 + \frac{m^2}{2} \hat{\chi}_i^2 - \frac{g^2}{4} [\hat{\chi}_i \hat{\chi}_j]^2 \right), \hat{\rho}_i = \sum_{\alpha=1}^{N^2-1} \hat{\rho}_i^\alpha \tau_\alpha, \hat{\chi}_i = \sum_{\alpha=1}^{N^2-1} \hat{\chi}_i^\alpha \tau_\alpha, [\hat{\chi}_{i\alpha} \hat{\chi}_{j\beta}] = i \delta_{j\beta} \delta_{\alpha\beta}, \hat{\mathcal{H}} \\
&= \mathcal{T}r \left(\frac{1}{2} \hat{\rho}_i^2 - \frac{g^2}{4} [\hat{\chi}_i \hat{\chi}_j]^2 + \frac{g}{2} \hat{\psi} \gamma^\mu [\hat{\chi}_i \hat{\psi}] - \frac{3i\mu}{4} \hat{\psi} \hat{\psi} + \frac{\mu^2}{2} \hat{\chi}_i^2 \right) - (N^2 - 1)\mu \\
&= \mathcal{T}r \left(\frac{1}{2} \hat{\rho}_i^2 - \frac{g^2}{4} [\hat{\chi}_1 \hat{\chi}_2]^2 + \frac{g}{2} \xi [-\hat{\chi}_1 - i \hat{\chi}_2 \xi] + \frac{g}{2} \xi^\dagger [-\hat{\chi}_1 - i \hat{\chi}_2 \xi^\dagger] + \frac{3\mu}{2} \xi^\dagger \xi \right. \\
&\quad \left. + \frac{\mu^2}{2} \hat{\chi}_i^2 \right) - (N^2 - 1)\mu, \hat{\mathcal{H}} \\
&= \mathcal{T}r \left(\hat{\rho}_z \hat{\rho}_z^\dagger + \frac{g^2}{4} [\hat{\mathcal{Z}}_\rho \hat{\mathcal{Z}}_\rho^\dagger]^2 - \frac{g}{\sqrt{2} \xi [\hat{\mathcal{Z}}_\rho^\dagger \xi]} - \frac{g}{\sqrt{2} \xi^\dagger [\hat{\mathcal{Z}} \xi^\dagger]} + \frac{3\mu}{2} \xi^\dagger \xi + \mu^2 \hat{\mathcal{Z}} \hat{\mathcal{Z}}^\dagger \right), \hat{\mathcal{H}} \\
&= \sum_{\alpha, i} \left(\frac{1}{2 \hat{\rho}_{i\alpha}^2} + \frac{m^2}{2} \hat{\chi}_{i\alpha}^2 \right) + \frac{g^2}{4} \sum_{\gamma, J, J} \left\langle \sum_{\alpha, \beta} f_{\alpha\beta\gamma} \hat{\chi}_J^\alpha \hat{\chi}_J^\beta \right\rangle^2, \hat{\mathcal{H}} \\
&= m \sum_{\alpha, i} \left(\hat{\eta}_{i\alpha} + \frac{1}{2} \right) + \frac{g^2}{16m^2} \sum_{\gamma, J, J} \left(\sum_{\alpha, \beta} f_{\alpha\beta\gamma} (\hat{a}_{i\alpha} + \hat{a}_{i\alpha}^\dagger) (\hat{a}_{j\beta} + \hat{a}_{j\beta}^\dagger) \right)^2, [\hat{\mathcal{H}}, \hat{\mathfrak{G}}_\alpha], \hat{\mathcal{H}}' \\
&= \hat{\mathcal{H}} + c \sum_\alpha \hat{\mathfrak{G}}_\alpha^2, \hat{\mathcal{H}} \\
&= \sum_\alpha \left(\frac{\hat{\rho}_{1\alpha}^2}{2} + \frac{\hat{\rho}_{2\alpha}^2}{2} + \mu^2 \frac{\hat{\chi}_{1\alpha}^2}{2} + \mu^2 \frac{\hat{\chi}_{2\alpha}^2}{2} + \frac{3\mu}{2} \xi_\alpha^\dagger \xi_\alpha \right) \\
&\quad + g^2 \sum_{\alpha \neq \beta} \hat{\chi}_{1\alpha}^2 \hat{\chi}_{2\beta}^2 \\
&\quad - 2g^2 \sum_{\alpha < \beta} \hat{\chi}_{1\alpha} \hat{\chi}_{1\beta} \hat{\chi}_{2\alpha} \hat{\chi}_{2\beta} + \frac{ig}{\sqrt{2}} \sum_{\alpha\beta\gamma} \epsilon_{\alpha\beta\gamma} [(-\hat{\chi}_{1\alpha} - i \hat{\chi}_{2\alpha}) \xi_\beta^\dagger \xi_\gamma^\dagger \\
&\quad + (-\hat{\chi}_{1\alpha} + i \hat{\chi}_{2\alpha}) \xi_\beta \xi_\gamma] - 3\mu \\
\hat{\mathcal{H}}' &= \hat{\mathcal{H}} + c \sum_\alpha \hat{\mathfrak{G}}_\alpha^2 + c' (\hat{\mathfrak{M}} - \hat{\mathfrak{S}})^2, \mathfrak{E}_0 \leq \mathfrak{E}_\lambda = \frac{\langle \psi(\theta_i) | \mathcal{H} | \psi(\theta_i) \rangle}{\langle \psi(\theta_i) | \psi(\theta_i) \rangle}, \mathfrak{E}_0 \equiv \langle \psi_\theta | \hat{\mathcal{H}} | \psi_\theta \rangle \\
&= \int d\chi |\psi_\theta(\chi)|^2 \cdot \frac{\langle \chi | \hat{\mathcal{H}} | \psi_\theta \rangle}{\psi_\theta(\chi)} = \mathbb{E}_{\chi \sim |\psi_\theta|^2} [\epsilon_\theta(\chi)], \nabla_\theta \mathfrak{E}_0 \\
&= \mathbb{E}_{\chi \sim |\psi_\theta|^2} [\nabla_{\theta \in \theta}(\chi)] + \mathbb{E}_{\chi \sim |\psi_\theta|^2} \times [\epsilon_\theta(\chi) \nabla_\theta \ln |\psi_\theta|^2], \theta' = \theta - \beta \nabla_\theta \mathfrak{E}_0
\end{aligned}$$

$$\begin{aligned}
\rho_{\theta}(\chi) &= \rho(\chi_1; F_{\theta}^0) \rho[\chi_1; F_{\theta}^1(\chi_1)] \rho[\chi_3; F_{\theta}^2(\chi_1, \chi_2)], F_{\theta}^i \\
&= \Lambda_{\theta}^{\iota, m} \odot \tan \hbar \odot \Lambda_{\theta}^{\iota, m-1} \odot \tan \hbar \odot \cdots \odot \Lambda_{\theta}^{\iota, 2} \odot \tan \hbar \odot \Lambda_{\theta}^{\iota, 1}, \Lambda_{\theta}^{\iota, \alpha}(\vec{\chi}) \\
&= \mathcal{M}_{\theta}^{\iota, \alpha} \vec{\chi} + \vec{\beta}_{\theta}^{\iota, \alpha}, \hat{\mathcal{H}}' \\
&= \hat{\mathcal{H}} \\
&+ c \sum_{\alpha} \hat{\mathbb{G}}_{\alpha}^2, \hat{\rho}|\psi\rangle = \int d\mathcal{U}, \hat{\mathcal{U}} \left| \psi \right\rangle, \langle \hat{\mathcal{O}} \rangle_{\varsigma} = \frac{\langle \psi | \hat{\rho} \hat{\mathcal{O}} | \psi \rangle}{\langle \psi | \hat{\mathcal{O}} | \psi \rangle}, \langle \psi | \hat{\rho} \hat{\mathcal{O}} | \psi \rangle \\
&= \int d\chi \langle \psi | \hat{\rho} | \psi \rangle \langle \psi | \hat{\mathcal{O}} | \psi \rangle = \int d\mathcal{U} d\chi \psi^{\odot}(\mathcal{U} \chi \mathcal{U}^{\dagger}) \langle \chi | \hat{\mathcal{O}} | \psi \rangle \\
&= \mathbb{E}_{\mathcal{U}, \chi \sim \langle \psi \rangle^2} \left[\frac{\langle \chi | \hat{\mathcal{O}} | \psi \rangle}{\psi(\chi)} \psi^{\odot}(\mathcal{U} \chi \mathcal{U}^{\dagger}) / \psi^{\odot}(\chi) \right]
\end{aligned}$$

$$\begin{aligned}
\Re_{\gamma}(\theta) &= \exp\left(-\frac{\imath\theta}{2}\right) = \begin{pmatrix} \cos\frac{\theta}{2} & \cdots & -\sin\frac{\theta}{2} \\ \vdots & \ddots & \vdots \\ \sin\frac{\theta}{2} & \cdots & -\cos\frac{\theta}{2} \end{pmatrix}, \hat{a}_{\iota} \\
&= \hat{j}_1 \otimes \cdots \otimes \hat{j}_{\iota-1} \otimes \begin{bmatrix} 0 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 1 \end{bmatrix} \otimes \hat{j}_{\iota+1} \bigotimes \cdots \otimes \hat{j}_6, \hat{a}_{\iota} \\
&= \hat{j}_1 \otimes \cdots \otimes \hat{j}_{\iota-1} \otimes \left\| \begin{matrix} 0 & 1 & 0 \\ 0 & \sqrt{2} & 1 \\ 1 & 0 & \sqrt{2} \end{matrix} \right\| \otimes \hat{j}_{\iota+1} \bigotimes \cdots \otimes \hat{j}_6, c_1 \\
&= \hat{j}_{64} \otimes \begin{pmatrix} 0 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 1 \end{pmatrix} \otimes \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 1 \end{pmatrix}, c_2 \\
&= \hat{j}_{64} \otimes \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & -1 \end{pmatrix} \otimes \begin{pmatrix} 0 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 0 \end{pmatrix} \otimes \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 1 \end{pmatrix}, c_3 \\
&= \hat{j}_{64} \otimes \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & -1 \end{pmatrix} \otimes \begin{pmatrix} 1 & \cdots & 0 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & -1 \end{pmatrix} \otimes \begin{pmatrix} 0 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 0 & \cdots & 0 \end{pmatrix} \\
&\hat{\xi}_{\alpha} = \left\| \begin{matrix} \sigma_z \otimes \cdots & \sigma_z \otimes \\ \otimes \sigma_z & \otimes \end{matrix} \right\|_{\alpha-1} \begin{pmatrix} 0 & 0 \\ 1 & 0 \end{pmatrix} \otimes \mathbb{7} \otimes \cdots \otimes \mathbb{7}
\end{aligned}$$

$$\hat{\mathbb{G}}_{\alpha} = \imath \sum_{\beta, \gamma, \iota} f_{\alpha \beta \gamma} \hat{a}_{\iota \beta}^{\dagger} \hat{a}_{\iota \gamma}, \hat{\mathbb{G}}_{\alpha} = \left(\otimes_{\iota \beta} |0\rangle_{\iota \beta} \right), \hat{\mathbb{G}}_{\alpha} = \imath \sum_{\beta \gamma} f_{\alpha \beta \gamma} \left(\hat{a}_{1 \beta}^{\dagger} \hat{a}_{1 \gamma} + \hat{a}_{2 \beta}^{\dagger} \hat{a}_{2 \gamma} + \xi_{\beta}^{\dagger} \xi_{\gamma} \right)$$

$$\begin{aligned}
\hat{a}_{i\alpha}^{\dagger} &= \sqrt{\frac{m}{2}} \widehat{\chi}_{i\alpha} - \frac{\imath \widehat{\wp}_{i\alpha}}{\sqrt{2m}}, \widehat{\alpha}_{i\alpha} = \sqrt{\frac{m}{2}} \widehat{\chi}_{i\alpha} + \frac{\imath \widehat{\wp}_{i\alpha}}{\sqrt{2m}}, \hat{a}_{i\alpha} |0\rangle_{i\alpha}, |\eta\rangle_{i\alpha} = \frac{(\hat{a}_{i\alpha}^{\dagger})^{\eta}}{\sqrt{\eta!} |0\rangle_{i\alpha}}, |\{\eta_{i\alpha}\}\rangle = \otimes_{i\alpha} |\eta\rangle_{i\alpha}, \hat{a}_{\tau}^{\dagger} \\
&= \sum_{\eta=0}^{\Lambda-2} \sqrt{\eta+1} |\eta+1\rangle \langle \eta| = \sum_{\eta=0}^{\Lambda-2} \sqrt{\eta+1} |\eta\rangle \langle \eta+1|, \hat{\eta}_{\tau} \\
&= \sum_{\eta=0}^{\Lambda-1} \eta |\eta\rangle \langle \eta|, [\widehat{\mathfrak{H}}_{\tau}, \widehat{\mathfrak{G}}_{\tau}], |\eta\rangle = |\beta_0\rangle |\beta_1\rangle \cdots |\beta_{\kappa+1}\rangle, |\eta+1\rangle \langle \eta| \\
&= \otimes_{\ell=0}^{\kappa-1} (|\beta'_{\ell}\rangle \langle \beta_{\ell}|), |0\rangle \langle 0| = \boxtimes_2 - \frac{\sigma_{z\blacksquare}}{2}, |1\rangle \langle 1| \\
&= \boxtimes_2 - \frac{\sigma_{z\blacksquare}}{2}, |0\rangle \langle 1| = \sigma_x + \frac{\imath \sigma_y}{2}, |1\rangle \langle 0| = \sigma_x + \frac{\imath \sigma_y}{2} \\
\widehat{\mathcal{M}} &= \sum_{\alpha} \left(\imath (\hat{\mathcal{Z}}_{\alpha} \widehat{\wp}_{z\alpha}^{\dagger} - \widehat{\wp}_{z\alpha} \hat{\mathcal{Z}}_{\alpha}^{\dagger}) - \frac{1}{2} \hat{\xi}_{\alpha}^{\dagger} \hat{\xi}_{\alpha} \right), \hat{\mathcal{Q}} = -\sqrt{2} \hat{\xi}_{\alpha}^{\dagger} (\widehat{\wp}_{z\alpha} - \imath \mu \hat{\mathcal{Z}}_{\alpha}) - \frac{g}{\sqrt{2} f_{\alpha\beta\gamma} \hat{\xi}_{\alpha} \hat{\mathcal{Z}}_{\beta} \hat{\mathcal{Z}}_{\gamma}^{\dagger}}, \hat{\mathcal{Q}}^2 \\
&= -\imath \hat{\mathcal{Z}}_{\alpha} \widehat{\mathfrak{G}}_{\alpha}, \hat{\mathcal{Q}}^{\dagger 2} = \imath \hat{\mathcal{Z}}_{\alpha}^{\dagger} \widehat{\mathfrak{G}}_{\alpha}, \{\hat{\mathcal{Q}} \hat{\mathcal{Q}}^{\dagger}\} = 2(\widehat{\mathcal{H}} - \mu \widehat{\mathcal{M}}), \hat{\mathcal{Q}} |\mathfrak{B}\mathfrak{P}\mathfrak{S}\rangle = \hat{\mathcal{Q}}^{\dagger} |\mathfrak{B}\mathfrak{P}\mathfrak{S}\rangle \\
&= (\widehat{\mathcal{H}} - \mu \widehat{\mathcal{M}}) \mathfrak{B}\mathfrak{P}\mathfrak{S}\rangle \\
\delta_{\epsilon} &= [\widehat{\mathfrak{Q}} \epsilon^* + \widehat{\mathfrak{Q}}^{\dagger} \epsilon], \widehat{\mathfrak{Q}} = -\hat{\xi}_{\alpha}^{\dagger} [(\widehat{\wp}_1^{\alpha} - \imath \widehat{\wp}_2^{\alpha}) - \imath \mu (\hat{\chi}_1^{\alpha} - \imath \hat{\chi}_2^{\alpha})] - \frac{\imath g}{\sqrt{2} f_{\alpha\beta\gamma} \hat{\xi}_{\alpha} \hat{\chi}_1^{\beta} \imath \hat{\chi}_2^{\gamma}}, \hat{\mathcal{Z}} = \hat{\chi}_1 - \frac{\imath \hat{\chi}_2}{\sqrt{2}}, \widehat{\wp}_z \\
&= \widehat{\wp}_1 - \frac{\imath \widehat{\wp}_2}{\sqrt{2}}, [\hat{\mathcal{Z}}, \widehat{\wp}_z^{\dagger}] = [\widehat{\wp}, \hat{\mathcal{Z}}_{\wp}^{\dagger}] = \imath
\end{aligned}$$

$$\begin{aligned}
Z(\mathcal{T}) &= \mathcal{T}r_{\mathcal{H}_{inv}} \left(e^{-\frac{\hat{H}}{t}} \right), Z(\mathcal{T}) = \frac{1}{\text{vol}(G)} \int_G dg \mathcal{T}r_{\mathcal{H}_{ext}} \left(\hat{g} e^{-\frac{\hat{H}}{t}} \right), \hat{\mathcal{P}} \equiv \frac{1}{\text{vol}(G)} \int_G dg \hat{g}, |\Phi\rangle_{inv} \\
&= \frac{1}{\sqrt{\mathcal{C}_\phi}} \times \frac{1}{\text{vol}(G)} \int_G dg (\hat{g}|\Phi\rangle), \mathcal{C}_\phi = \frac{1}{[\text{vol}(G)]^2} \int_G dg \int_G dg' \langle \phi | \hat{g}^{-1} \hat{g}' | \phi \rangle \\
&= \frac{1}{\text{vol}(G)} \int_G dg \langle \phi | \hat{g} | \phi \rangle \\
&= \text{vol} \left(\frac{G_\phi}{\text{vol}} \right) (G), \mathcal{T}r_{\mathcal{H}_{inv}} \left(e^{-\frac{\hat{H}}{t}} \right) \sum_\phi \text{vol} \left(\frac{G_\phi}{\text{vol}} \right) (G)_{inv} \langle \phi | e^{-\frac{\hat{H}}{t}} | \phi \rangle_{inv} \\
&= \sum_\phi \text{vol} \left(\frac{G_\phi}{\text{vol}} \right) (G) \otimes \frac{1}{\mathcal{C}_\phi} \frac{1}{[\text{vol}(G)]^2} \int_G dg \int_G dg' \langle \phi | \hat{g}^{-1} e^{-\frac{\hat{H}}{t}} \hat{g}' | \phi \rangle \\
&= \sum_\phi \frac{1}{\text{vol}(G)} \int_G dg \mathcal{T}r_{\mathcal{H}_{ext}} \left(\hat{g} e^{-\frac{\hat{H}}{t}} \right), Z(\mathcal{T}) \\
&= \frac{1}{[\text{vol}(G)]^\kappa} \int \left(\prod_{\kappa=1}^\kappa d\mathfrak{U}_{(\mathfrak{K})} \right) \mathcal{T}r_{\mathcal{H}_{ext}} \otimes \left(\hat{\mathfrak{U}}_{(\mathcal{K})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathfrak{U}}_{(k-1)}^{-1} \hat{\mathfrak{U}}_{(k-1)} \otimes e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \right) \\
&= \frac{1}{[\text{vol}(G)]^\kappa} \int \left(\prod_{\kappa=1}^\kappa d\mathfrak{U}_{(\mathfrak{K})} \right) \int \left(\prod_{\kappa=1}^\kappa d\chi_{(\mathfrak{K})} \right) \langle \chi_{(k)} | \hat{\mathfrak{U}}_{(\mathcal{K})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathfrak{U}}_{(k-1)}^{-1} | \chi_{(k-1)} \rangle \\
&\quad \bigotimes \langle \chi_{(k-1)} | \hat{\mathfrak{U}}_{(\mathcal{K}-1)} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathfrak{U}}_{(k-2)}^{-1} | \chi_{(k-2)} \rangle \\
&\quad \otimes \cdots \otimes \langle \chi_{(1)} | \hat{\mathfrak{U}}_{(1)} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathfrak{U}}_{(k-1)}^{-1} | \chi_{(k)} \rangle \\
&\langle \chi_{(k)} | \hat{\mathfrak{U}}_{(\mathcal{K})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} \hat{\mathfrak{U}}_{(k-1)}^{-1} | \chi_{(k-1)} \rangle = \langle \mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1} | \hat{\mathfrak{U}}_{(\mathcal{K})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} | \mathfrak{U}_{(k-1)} \chi_{(k-1)} \mathfrak{U}_{(k-1)}^{-1} \rangle \\
&= \int_G d\mathcal{P} \langle \mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1} | \hat{\mathcal{P}}_{(\mathcal{K})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} | \mathcal{P} \rangle \otimes \langle \mathcal{P} | \hat{\mathcal{P}}_{(\mathcal{K})} e^{-\frac{\mathcal{H}(\hat{\mathcal{P}}, \hat{\mathcal{X}})}{\mathcal{T}\kappa}} | \mathfrak{U}_{(k-1)} \chi_{(k-1)} \mathfrak{U}_{(k-1)}^{-1} \rangle \\
&= \int_G d\wp e^{\mathcal{T}r \left[\left(\mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1} - \mathfrak{U}_{(k-1)} \chi_{(k-1)} \mathfrak{U}_{(k-1)}^{-1} \right) \right]} \otimes e^{-\mathcal{H} \left(\mathcal{P}, \frac{\mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1}}{(\mathcal{T}\mathcal{K})} \right)} \\
&= e^{\mathcal{K}\mathcal{T}Tr \left[\left(\mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1} - \mathfrak{U}_{(k-1)} \chi_{(k-1)} \mathfrak{U}_{(k-1)}^{-1} \right)^2 \right]} \otimes e^{-\mathcal{V} \left(\frac{\mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1}}{(\mathcal{T}\mathcal{K})} \right)} \simeq e^{-\frac{\mathbb{D}t \left(\mathfrak{U}_{(k)} \chi_{(k)} \mathfrak{U}_{(k)}^{-1} - \mathfrak{U}_{(k-1)} \chi_{(k-1)} \mathfrak{U}_{(k-1)}^{-1} \right)}{\mathfrak{I}\mathfrak{K}}} \Big|_{\tau\kappa} = e^{-\frac{\mathcal{L}[\mathcal{D}t\chi_{\mathcal{K}}, \chi_{\mathcal{K}}]}{\tau\kappa}}
\end{aligned}$$

$$\begin{aligned}
\mathfrak{U}_{(k)}^{-1} \mathfrak{U}_{(k)} &\equiv e^{\frac{i\Lambda_{(k)}}{(KT)}} \chi_{(k)} - (\mathfrak{U}_{(k-1)} \mathfrak{U}_{(k)}^{-1})^{-1} \chi_{(k-1)} (\mathfrak{U}_{(k-1)} \mathfrak{U}_{(k)}^{-1}) \\
&\simeq \frac{\mathcal{D}t \chi_{(k)}}{KT}, Z(T) \int [d\Lambda][d\chi] e^{-\int dt \mathfrak{L}[\mathcal{D}t \chi, \chi]} = (\mathfrak{U}_k \mathfrak{U}_{k-1}^{-1}) (\mathfrak{U}_{k-1} \mathfrak{U}_{k-2}^{-1}) \otimes (\mathfrak{U}_2 \mathfrak{U}_1^{-1}) \\
&= \wp e^{i \int_0^1 dt \Lambda_t}
\end{aligned}$$

3. Gravedad cuántica en espacios curvos.

$$\begin{aligned}
d\hat{s}^2 &= \hat{g}_{\mu\nu} d\chi^\mu d\chi^\nu = -\left(\kappa + \frac{r^2}{\ell^2}\right) dt^2 + \frac{dr^2}{\kappa} + \frac{r^2}{\ell^2} + r^2 d\Sigma_\kappa^2, \mathfrak{H}_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= \frac{r^2}{\ell^2} (-dt^2 + \ell^2 d\Sigma_\kappa^2), ds^2 = -N(r) dt^2 + \mathcal{H}(r) dr^2 + \delta(r) d\Sigma_\kappa^2, \mathcal{V}(\phi) \\
&= -\frac{3}{\kappa\ell^2} - \frac{\phi^2}{\ell^2} + \mathcal{O}(\phi)^4, \phi(r) = \frac{\alpha}{r} + \frac{\beta}{r^2} + \mathcal{O}(r^{-3}), N(r) = -g_{tt} \\
&= \frac{r^2}{\ell^2} + \kappa - \frac{\mu}{r} + \mathcal{O}(r^{-2}), \delta(r) \\
&= r^2 + \mathcal{O}(r^{-2}), N\delta'^2 \mathcal{H} - 2N\delta'' \mathcal{H} \delta + (N\mathcal{H})' \delta' \delta - 2\kappa N\mathcal{H} \delta^2 \phi'^2, \mathcal{H}(r) = g_{rr} \\
&= \frac{\ell^2}{r^2} + \frac{\ell^4}{r^4} \left(-\kappa - \frac{\alpha^2 \kappa}{2\ell^2}\right) + \frac{\ell^5}{r^5} \left(\frac{\mu}{\ell} - \frac{4\kappa\alpha\beta}{3\ell^3}\right) + \mathcal{O}(r^{-6}), g_{rr} \\
&= \frac{\ell^2}{r^2} + \frac{\alpha\ell^4}{r^4} + \frac{\beta\ell^5}{r^5} + \mathcal{O}(r^{-6}), \mathcal{V}(\phi) = -\frac{3}{\kappa\ell^2} - \frac{\phi^2}{\ell^2} + \lambda\phi^3 + \mathcal{O}(r^4), \phi(r) \\
&= \frac{\alpha}{r} + \frac{\beta}{r^2} + \frac{\gamma \ln(r)}{r^2} + \mathcal{O}(r^{-3}), \mathcal{H}(r) = g_{rr} \\
&= \frac{\ell^2}{r^2} + \frac{\ell^4}{r^4} \left(-\kappa - \frac{\alpha^2 \kappa}{2\ell^2}\right) + \frac{\ell^5}{r^5} \left(\frac{\mu}{\ell} - \frac{4\kappa\alpha\beta}{3\ell^3} + \frac{2\kappa\alpha\gamma}{9\ell^3}\right) + \frac{\ell^5 \ln(r)}{r^5 \left(-\frac{4\kappa\alpha\gamma}{3\ell^3}\right)} \\
&+ \mathcal{O}\left[\ln \frac{(r)^2}{r^6}\right], \mathcal{H}(r) = \frac{\ell^2}{r^2} + \frac{\ell^4 \alpha}{r^4} + \frac{\ell^5 \beta}{r^5} + \frac{\ell^5 c \ln r}{r^5} + \mathcal{O}\left[\ln \frac{(r)^2}{r^6}\right], \alpha = -\kappa - \frac{\alpha^2 \kappa}{2\ell^2}, \beta \\
&= \frac{\mu}{\ell} - \frac{4\kappa\alpha\beta}{3\ell^3} + \frac{2\kappa\alpha\gamma}{9\ell^3}, c \\
&= -\frac{4\kappa\gamma\alpha}{3\ell^3}, \partial_r \left(\frac{\phi' \delta \sqrt{N}}{\sqrt{H}}\right) - \frac{\delta \sqrt{HN} \partial \mathfrak{B}}{\partial \phi} 3\alpha^2 \ell^2 \lambda + \frac{\gamma}{\ell^2} + \mathcal{O}(r^{-1}), \xi^r \\
&= r\eta^r(\chi^m) + \mathcal{O}(r^{-1}) \xi^m = \mathcal{O}(1), \phi'(\chi) = \phi(\chi) + \xi^\mu \partial_\mu \phi(\chi) \\
&= \frac{\alpha'}{r} + \frac{\beta'}{r^2} + \gamma' \ln \frac{r}{r^2} + \mathcal{O}(r^{-3}), \alpha' = \alpha - \eta^r \alpha + \xi^m \partial_m \alpha, \beta' \\
&= \beta - \eta^r (2\beta - \gamma) \xi^m \partial_m \beta, \gamma' = \gamma - 2\gamma \eta^r + \xi^m \partial_m \gamma
\end{aligned}$$



$$\frac{\alpha'\partial\gamma}{\partial\alpha}-\gamma'\alpha\frac{\partial\gamma}{\partial\alpha}-\gamma+\eta^r\left(2\gamma-\alpha\frac{\partial\gamma}{\partial\alpha}\right)+\xi^m\left(\frac{\partial\alpha}{\partial\chi_m}\frac{\partial\gamma}{\partial\alpha}-\frac{\partial\gamma}{\partial\chi_m}\right)$$

$$\frac{\alpha'\partial\beta}{\partial\alpha}-\beta'\alpha\frac{\partial\beta}{\partial\alpha}-\beta+\eta^r\left(2\beta-\gamma-\alpha\frac{\partial\beta}{\partial\alpha}\right)+\xi^m\left(\frac{\partial\alpha}{\partial\chi_m}\frac{\partial\beta}{\partial\alpha}-\frac{\partial\beta}{\partial\chi_m}\right)$$

$$\mathcal{I}_{\mathbb{C}\mathfrak{F}\mathfrak{I}}\rightarrow \mathcal{I}_{\mathbb{C}\mathfrak{F}\mathfrak{I}}-\int d^3\chi \mathfrak{W}[\mathcal{O}(\chi)]$$

$$\mathcal{I}_g^{ct}=-\frac{1}{\kappa\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\left(\frac{2}{\ell}+\frac{\mathcal{R}\ell}{2}\right)},\mathfrak{I}$$

$$=\frac{\int d^4\chi\sqrt{-g}(\mathcal{R}-\frac{(\partial\phi)^2}{2}-\mathcal{V}(\phi))}{2\kappa}+1/\kappa\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\mathfrak{K}$$

$$-1/\kappa\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\left(\frac{2}{\ell}+\frac{\mathcal{R}\ell}{2}\right)+\mathcal{I}_{\varphi\phi\psi},\mathcal{I}_{\varphi\phi\psi}^{ct}$$

$$=1/6\kappa\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\left(\phi\eta^\nu\partial_\nu\phi-\phi^2\ell/2\varphi\psi\tau\right),\mathcal{I}_{\varphi\phi\psi}$$

$$=-\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\left[\frac{\phi^2}{2\ell}+\frac{\mathfrak{W}(\alpha)}{\ell\alpha^3\tau^3}-\langle\varphi|\phi|\psi\rangle^3\right],\delta\ell$$

$$=\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\left[\frac{1}{r}\left(\sqrt{-g^{rr}}\phi'-\frac{\varphi\psi}{\tau\ell}-\frac{3\mathcal{W}(\alpha)\phi^2}{\ell\phi^3}\right)\left(1+\frac{1}{\frac{rd^2\mathcal{W}(\alpha)}{d(\alpha)^2}}\right)\right.$$

$$\left.+\left(\frac{3\mathcal{W}(\alpha)}{\alpha}-\beta\right)\langle\varphi|\phi|\psi\rangle^3/\ell\alpha^3\right]\delta\alpha$$

$$\mathcal{I}_{\varphi\phi\psi}=-\int_{\partial\mathcal{M}}^{\delta}d^3\chi\sqrt{-\hbar}\left[\frac{\phi^2}{2\ell}+\frac{\langle\varphi|\phi|\psi\rangle^3}{\ell\alpha^3\left(\mathcal{W}-\frac{\alpha\gamma}{3}\right)}-\frac{\langle\varphi|\phi|\psi\rangle^3\mathfrak{C}_\gamma}{3\ell}\ln\left(\frac{\langle\varphi|\phi|\psi\rangle}{\alpha}\right)\right]$$

$$ds^2=\Omega(\chi)=\left[-f(\chi)dt^2+\frac{\eta^2d\chi^2}{f(x)}+d\Sigma_{\kappa}^2\right],\mathcal{I}_{\mathfrak{B}\mathfrak{U}\Omega\mathfrak{K}}^{\mathfrak{E}}=\int_0^{\frac{1}{\tau}}d\tau\int_{\chi_+}^{\chi_{\beta}}d^3\chi\sqrt{g^{\mathfrak{E}}\mathfrak{B}}(\phi)$$

$$=\frac{\sigma_{\kappa}}{2\eta\kappa\tau}d(\Omega f)/d\chi|_{\chi_+}^{\chi_{\beta}}$$

$$\begin{aligned}
\Omega(\chi) &\rightarrow \delta(r), f(x) \rightarrow \frac{N(r)}{\delta(r)}, dx \rightarrow \frac{\sqrt{NH}}{\eta\delta} dr \\
\mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} &= \frac{\frac{\sigma_{\kappa}}{2\kappa\tau} \delta}{\sqrt{NH}} dN \Big|_{r_+}^{r_{\beta}} ds^2 = ' \hbar_{\alpha\beta} d\chi^{\alpha} d\chi^{\beta} = \Omega(\chi_0)[-f(\chi_0)dt^2 + d\Sigma_{\kappa}^2]\eta_{\alpha} = \frac{\delta_a^{\chi}}{\sqrt{g^{xx}}} \Big|_{\chi=\chi_0} \mathfrak{K}_{\alpha\beta} \\
&= \frac{\sqrt{g^{xx}}}{2} \partial_{\chi} g_{\alpha\beta} \Big|_{\chi=\chi_0} \mathfrak{K} = \frac{1}{2\eta \left(\frac{f}{\Omega}\right)^{-\frac{1}{2}} \left[\frac{(f\Omega)'}{\Omega f} + \frac{2\Omega'}{\Omega}\right] \Big|_{\chi_0}} \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} \\
&= -\frac{\frac{\sigma_{\kappa}}{\kappa\tau} \Omega f \left[\frac{(f\Omega)'}{\Omega f} + \frac{2\Omega'}{\Omega}\right] \Big|_{\chi_{\beta}}}{2\eta} = -\frac{\sigma_{\kappa}}{2\tau\kappa} \left(\frac{\frac{\delta}{\sqrt{NH}} dN}{dr} + \frac{\frac{2N}{\sqrt{NH}} d\delta}{dr} \right) \Big|_{r_{\beta}} \mathcal{J}_g^{ct} \\
&= \frac{2\sigma_{\kappa}}{\kappa\tau\ell} \left(\Omega^{\frac{3}{2}} f^{\frac{1}{2}} + \frac{\ell^2\kappa}{2} f^{\frac{1}{2}} \Omega^{\frac{1}{2}} \right) \left(\Omega^{\frac{3}{2}} f^{\frac{1}{2}} + \frac{\ell^2\kappa}{2} f^{\frac{1}{2}} \Omega^{\frac{1}{2}} \right) \Big|_{\chi_{\beta}} = \frac{2\sigma_{\kappa}}{\kappa\tau\ell} \delta\sqrt{N} \left(1 + \frac{\ell^2\kappa}{2\delta} \right) \Big|_{r_{\beta}} \mathcal{T} \\
&= \frac{N'}{4\pi\sqrt{NH}} \Big|_{r_+} \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} + \mathcal{J}_g^{ct} \\
&= \frac{1}{\tau \left[\frac{\sigma_{\kappa}\delta(r_+)\tau}{4\mathfrak{G}} \right]} - \frac{\sigma_{\kappa}}{2\kappa\tau \left[\frac{\frac{2N}{\sqrt{NH}} d\delta}{dr} - \frac{4}{\ell\delta\sqrt{N} \left(1 + \frac{\ell^2\kappa}{2\delta} \right)} \right] \Big|_{r_{\beta}}} + \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} \\
&\quad + \mathcal{J}_g^{ct} = -\frac{\mathcal{A}}{4\mathfrak{G}} - \frac{\sigma_{\kappa}}{\tau} \left(-\frac{\mu}{\kappa} + \frac{4\alpha\beta}{3\ell^2} + \frac{r\alpha^2}{2\ell^2} \right) \Big|_{r_{\beta}} \mathcal{J}_{\langle\phi|\phi|\psi\rangle}^{ct} \\
&= \int_{\partial\mathcal{M}}^{\delta} d^3\chi \sqrt{\hbar^{\varepsilon}} \left[\frac{\phi^2}{2\ell} + \frac{\frac{\mathcal{W}(\alpha)}{\ell\alpha^3 \langle\phi|\phi|\psi\rangle^3}}{\|\tau\omega\|} \right] = \frac{\sigma_{\kappa}}{\tau} \left(\frac{\mathcal{W}}{\ell^2} + \frac{\alpha\beta}{\ell^2} + \frac{r\alpha}{2\ell^2} \right) \Big|_{r_{\infty}} \mathcal{J}^{\mathfrak{E}} \\
&= \mathcal{J}_{\mathfrak{B}\mathfrak{U}\mathfrak{L}\mathfrak{R}}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} + \mathcal{J}_g^{ct} + \mathcal{J}_{\langle\phi|\phi|\psi\rangle}^{ct} = \frac{\mathcal{A}}{4\mathfrak{G}} + \frac{\sigma_{\kappa}}{\tau} \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{\alpha}{3d\mathcal{W}} \right) \right] \mathcal{F} = \mathcal{J}^{\mathfrak{E}}\mathcal{T} \\
&= \mathcal{M} - \mathcal{T}\mathcal{S}, \mathcal{M} = -\frac{\mathcal{T}^2\partial\mathcal{J}^{\mathfrak{E}}}{\partial\mathcal{T}} = \sigma\kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{\alpha}{3d\mathcal{W}} \right) \right] \mathcal{S} = -\frac{\partial(\mathcal{J}^{\mathfrak{E}}\mathcal{T})}{\partial\mathcal{T}} = \frac{\mathcal{A}}{4\mathfrak{G}}
\end{aligned}$$

$$\begin{aligned}
\mathcal{I}_{\mathfrak{B}\mathfrak{U}\mathfrak{R}}^{\mathfrak{E}} + \mathcal{I}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} + \mathcal{I}_g^{ct} + \mathcal{I}_{\langle\phi|\phi|\psi\rangle}^{ct} &= -\frac{\mathcal{A}}{4\mathfrak{G}} + \frac{\sigma\kappa}{\mathcal{T}\left\{\frac{\mu}{\kappa} + \frac{1}{\ell^2}\left[\mathcal{W}(\alpha) - \frac{\alpha}{3d\mathcal{W}} + \frac{2\alpha\gamma}{9} - \frac{\alpha\gamma}{3}\ln r\right]\right\}\hat{\mathcal{I}}_{\langle\phi|\phi|\psi\rangle}^{ct}} \\
&= \int_{\partial\mathcal{M}}^{\delta} d^3\chi\sqrt{\hbar\varepsilon}\left\{\frac{\phi^3\gamma}{3\alpha^2\ell}[\ln(\alpha/\phi) - 1]\right\} = \sigma_{\kappa}/\mathfrak{T}\left[-\frac{\alpha\gamma}{3\ell^2} + \frac{\alpha\gamma\ln r}{3\ell^2} + \mathcal{O}(r^{-1}\ln r)\right] \\
\mathcal{I}^{\mathfrak{E}} = \mathcal{I}_{\mathfrak{B}\mathfrak{U}\mathfrak{R}}^{\mathfrak{E}} + \mathcal{I}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} + \mathcal{I}_g^{ct} + \mathcal{I}_{\langle\phi|\phi|\psi\rangle}^{ct} + \hat{\mathcal{I}}_{\langle\phi|\phi|\psi\rangle}^{ct} &= -\frac{\mathcal{A}}{4\mathfrak{G}} + \frac{\sigma\kappa}{\mathcal{T}\left[\frac{\mu}{\kappa} + \frac{1}{\ell^2}\left(\mathcal{W}(\alpha) - \frac{\alpha}{3d\mathcal{W}} - \frac{\alpha\gamma}{9}\right)\right]}, \mathcal{M} \\
&= -\gamma^2 \frac{\partial\mathcal{I}^{\mathfrak{E}}}{\partial\mathcal{T}} = \sigma\kappa\left[\frac{\mu}{\kappa} + \frac{1}{\ell^2}\left(\mathcal{W}(\alpha) - \frac{\alpha}{3d\mathcal{W}} - \frac{\alpha\gamma}{9}\right)\right]\delta = -\frac{\partial(\mathcal{I}^{\mathfrak{E}}\mathcal{T})}{\partial\mathcal{T}} = \frac{\mathcal{A}}{4\mathfrak{G}}
\end{aligned}$$

3.1. Tensor de stress Brown-York en espacios cuánticos curvos.

$$\begin{aligned}
\tau_{\alpha\beta} &= -\frac{1}{\kappa\left(\kappa_{\alpha\beta} - \hbar_{\alpha\beta}\kappa + \frac{2}{\ell\hbar_{\alpha\beta}} - \mathfrak{I}\mathfrak{G}_{\alpha\beta}\right)} - \frac{\hbar_{\alpha\beta}}{\ell\left[\frac{\phi^2}{2} + \frac{\mathcal{W}(\alpha)}{\alpha^3\phi^3}\right]\tau_{tt}} \\
&= \frac{\ell}{\mathcal{R}\left[\frac{\mu}{8\pi\mathfrak{G}\ell^2} + \frac{1}{\ell^4\left(\omega(\alpha) - \frac{\alpha\beta}{3}\right)}\right]} + \mathcal{O}(\mathfrak{R}^{-2})\tau_{\theta\theta} \\
&= \frac{\ell}{\mathcal{R}\left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2\left(\omega(\alpha) - \frac{\alpha\beta}{3}\right)}\right]} + \mathcal{O}(\mathfrak{R}^{-2})\tau_{\langle\phi|\phi|\psi\rangle} \\
&= \ell\sin^2\theta/\mathcal{R}\left[\frac{\mu}{16\pi\mathfrak{G}} - 1/\ell^2\left(\omega(\alpha) - \frac{\alpha\beta}{3}\right)\right] + \mathcal{O}(\mathfrak{R}^{-2}) \\
\langle\tau_{\alpha\beta}^{dualidad}\rangle &= \frac{3\mu}{16\pi\mathfrak{G}\ell^2\delta_{\alpha}^0\delta_{\beta}^0} + \frac{\gamma_{\alpha\beta}}{\ell^2\left[\frac{\mu}{16\pi\mathfrak{G}} - \frac{1}{\ell^2\left(\omega(\alpha) - \frac{\alpha\beta}{3}\right)}\right]\langle\tau^{dualidad}\rangle} = -3/\ell^4\left[\omega(\alpha) - \frac{\alpha\beta}{3}\right]
\end{aligned}$$

$$\begin{aligned}
\tau_{\alpha\beta} &= -\frac{1}{\kappa\left(\kappa_{\alpha\beta}-\hbar_{\alpha\beta}\kappa+\frac{2}{\ell\hbar_{\alpha\beta}}-\imath\mathfrak{G}_{\alpha\beta}\right)}-\frac{\hbar_{\alpha\beta}}{\ell\left[\frac{\phi^2}{2}+\frac{\phi^3}{\alpha^3}\left(\omega-\frac{\alpha\gamma}{3}\right)+\frac{\phi^3\gamma}{3\alpha^2}\left(\frac{\alpha}{\phi}\right)\right]}\tau_{tt} \\
&= \frac{\ell}{\mathcal{R}\left[\frac{\mu}{8\pi\mathfrak{G}\ell^2}+\frac{1}{\ell^4\left(\omega-\frac{\alpha\beta}{3}-\frac{\alpha\gamma}{9}\right)}\right]}+\mathcal{O}\left[\frac{(\ln\mathfrak{R})^3}{\mathfrak{R}^2}\right]\tau_{\theta\theta} \\
&= \frac{\ell}{\mathcal{R}\left[\frac{\mu}{16\pi\mathfrak{G}}-\frac{1}{\ell^2\left(\omega-\frac{\alpha\beta}{3}-\frac{\alpha\gamma}{9}\right)}\right]}+\mathcal{O}\left[\frac{(\ln\mathfrak{R})^3}{\mathfrak{R}^2}\right]\tau_{\langle\phi|\varphi|\psi\rangle} \\
&= \frac{\ell\sin^2\theta}{\mathcal{R}\left[\frac{\mu}{16\pi\mathfrak{G}}-\frac{1}{\ell^2\left(\omega-\frac{\alpha\beta}{3}-\frac{\alpha\gamma}{9}\right)}\right]}+\mathcal{O}\left[\frac{(\ln\mathfrak{R})^3}{\mathfrak{R}^2}\right]\langle\tau_{\alpha\beta}^{dualidad}\rangle \\
&= \frac{3\mu}{16\pi\mathfrak{G}\ell^2\delta_{\alpha}^0\delta_{\beta}^0}+\frac{\gamma_{\alpha\beta}}{\ell^2\left[\frac{\mu}{16\pi\mathfrak{G}}-\frac{1}{\ell^2\left(\omega(\alpha)-\frac{\alpha\beta}{3}-\frac{\alpha\gamma}{9}\right)}\right]}\langle\tau^{dualidad}\rangle \\
&= -3/\ell^4\left[\left(\omega-\frac{\alpha\beta}{3}-\frac{\alpha\gamma}{9}\right)\lambda\right]
\end{aligned}$$

$$\begin{aligned}
\tau^{\alpha\beta} &\equiv \frac{\frac{2}{\sqrt{-\hbar}} \delta \mathcal{I}}{\delta \hbar_{\alpha\beta}}, \mathcal{I}_{\mathfrak{G}\mathfrak{S}} + \mathcal{I}_g + \mathcal{I}_\phi \\
&= \frac{1}{\kappa \int d^\eta \chi \sqrt{-\hbar} \mathfrak{K} - \frac{1}{\kappa \int d^\eta \chi \sqrt{-\hbar} \left[\frac{\eta-1}{\ell} + \frac{\ell \mathcal{R}}{2(\eta-2)} \right] - \int d^\eta \chi} \sqrt{-\hbar} \psi}, \tau^{\alpha\beta} \\
&= -\frac{1}{\kappa \left(\mathcal{K}_{\alpha\beta} - \hbar_{\alpha\beta} \mathfrak{K} + \frac{\eta-1}{\ell} \hbar_{\alpha\beta} - \frac{\ell}{\eta} - 2\mathfrak{G}_{\alpha\beta} \right)} - \hbar_{\alpha\beta} [\psi], \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= N(\mathcal{R}) dt^2 + \delta(\mathcal{R}) d\Sigma_\kappa^2, \mathfrak{G}_{\alpha\beta} = \mathcal{R}_{\alpha\beta} - \frac{1}{2\mathfrak{H} \hbar_{\alpha\beta}}, \mathfrak{G}_{tt} = \frac{(\eta-2)(\eta-1)}{2} \frac{\kappa \mathcal{N}}{\mathcal{S}}, \mathfrak{G}_{ij} \\
&= -\frac{(\eta-2)(\eta-3)}{2} \kappa v_{ij}, \tau_{tt} = -\frac{(\eta-1)}{\kappa} \left[\frac{N\delta'}{2\delta\sqrt{\mathcal{H}}} - \frac{\mathcal{N}}{\ell} \left(1 + \frac{\ell^2 \kappa}{2\delta} \right) \right] + N[\psi], \tau_{ij} \\
&= \frac{v_{ij}}{\kappa} \left[\frac{\delta}{2\sqrt{\mathcal{H}} \left(\frac{\mathfrak{N}'}{\mathfrak{N}} + \frac{\mathcal{S}'}{\mathcal{S}} (\eta-2) \right)} - \frac{(\eta-1)\mathcal{S}}{\ell} - \frac{\ell \mathfrak{K}(\eta-3)}{2} \right] - v_{ij} \delta[\psi], \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= \mathcal{L}^2 dt^2 + \sigma_{ij} (d\gamma^i + \mathcal{L}^i dt) (d\gamma^j + \mathcal{L}^j dt), \mathfrak{E} = \mathcal{Q}_{\frac{\partial}{\partial t}} \\
&= \frac{\oint_\Sigma^\delta d^{D-2} \gamma \sqrt{\sigma} \mu^\alpha \tau_{\alpha\beta} \xi^\beta}{\sqrt{\mathfrak{N}}} = \left(\oint_\Sigma^\delta d^2 \gamma \sqrt{v} \right) \delta^{\left(\frac{D-2}{2}\right)} \tau_{tt} = \frac{\sigma_{\kappa, \eta-1} \delta^{\left(\frac{D-2}{2}\right)}}{\sqrt{\mathfrak{N}}} \tau_{tt}, \sigma_{ij} d\chi^i d\chi^j \\
&= \delta d\Sigma_\kappa^2, \tau_{tt} = -\frac{(\eta-1)}{\kappa \left[\frac{f^3}{2\eta\sqrt{\Omega}} - \frac{\Omega f}{\ell} \left(1 + \frac{\ell^2 \kappa}{2\Omega} \right) \right]} + \frac{\Omega f}{\kappa} [\psi], \tau_{ij} \\
&= \frac{v_{ij}}{\kappa \left[\frac{(\Omega f)'}{2\eta\sqrt{\Omega f}} + \frac{(\eta-2)}{2\eta} \frac{\Omega' \sqrt{f}}{\sqrt{\Omega}} - \frac{(\eta-1)\Omega}{\ell} - \frac{\ell \kappa(\eta-3)}{2} \right]} - \frac{v_{ij}\Omega}{\kappa} [\psi], \mathfrak{E} = \mathcal{Q}_{\frac{\partial}{\partial t}} \\
&= \oint_\Sigma^\delta d^{D-2} \gamma \sqrt{\sigma} \mu^\alpha \tau_{\alpha\beta} \xi^\beta = \frac{\sigma_{\kappa, \eta-1} \Omega^{\left(\frac{D-2}{2}\right)}}{\sqrt{\Omega f}} \tau_{tt}
\end{aligned}$$

3.2. Masa Hamiltoniana en espacios cuánticos curvos.

$$-\frac{5}{4\ell^2} > m^4 \geq -\frac{9}{4\ell^2}, \phi(r) = \frac{\alpha}{r} + \frac{\beta}{r^2} + \mathcal{O}(r^{-3})$$

3.3. Modelos Logarítmicos y Anti-logarítmicos en espacios cuánticos curvos.



$$\begin{aligned}
\delta \mathcal{M}_{\mathbb{G}} &= \frac{\sigma \kappa}{\kappa \left[r \delta \alpha + \ell \delta \beta + \mathcal{O}\left(\frac{1}{r}\right) \right] \delta \mathcal{M}_{\phi}} = \frac{\sigma \kappa}{\ell^2 \left[r \alpha \delta \alpha + \alpha \delta \beta + 2 \beta \delta \alpha + \mathcal{O}\left(\frac{1}{r}\right) \right] \delta \mathcal{M}} \\
&= \frac{\sigma \kappa}{\kappa \ell^2 \left[r(\ell^2 \delta \alpha + \kappa \alpha \delta \alpha) + \ell^3 \delta \beta + \kappa(\alpha \delta \beta + 2 \beta \delta \alpha) + \mathcal{O}\left(\frac{1}{r}\right) \right] \kappa} + \alpha + \frac{\alpha^2}{2 \ell^2}, \delta \mathcal{M} \\
&= \frac{\sigma \kappa}{\kappa \ell^2 [\ell^3 \delta \beta + \kappa(\alpha \delta \beta + 2 \beta \delta \alpha)] \mathcal{M}} \\
&= \frac{\sigma \kappa \left[\frac{\ell \beta}{\kappa} + \frac{1}{\ell^2 \left(\frac{\alpha d \omega(\alpha)}{d \alpha} + \omega(\alpha) \right)} \right] \ell c}{\kappa} - 4 \alpha^3 \lambda, \delta \mathcal{M}_{\mathbb{G}} \\
&= \left\{ \frac{\ell \delta \beta}{\kappa} + \frac{\delta \alpha}{\kappa} r + \frac{\ell \delta c}{\kappa} \ln(r) + \frac{\mathcal{O}(\ln(r)^2)}{r} \right\} \sigma \kappa, \delta \mathcal{M}_{\phi} \\
&= \left[\alpha \delta \beta + 2 \beta \delta \alpha + \frac{3 \alpha^2 \ell^2 \lambda \delta \alpha}{\ell^2} + \frac{r \alpha \delta \alpha}{\ell^2} - 12 \lambda \alpha^2 \delta \alpha \ln(r) + \mathcal{O}\left(\ln \frac{(r)^2}{r}\right) \right] \sigma \kappa, \delta \mathcal{M} \\
&= \left[\frac{\ell \delta \beta}{\kappa} + \alpha \delta \beta + 2 \beta \delta \alpha + \frac{3 \alpha^2 \ell^2 \lambda \delta \alpha}{\ell^2} \right] \sigma \kappa, \mathcal{M} \\
&= \left[\frac{\ell \beta}{\kappa} + \frac{1}{\ell^2 \left(\frac{\alpha d \omega}{d \alpha} + \omega(\alpha) + \alpha^3 \ell^2 \lambda \right)} \right] \sigma \kappa, \mathcal{M} \\
&= \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2 \left(\omega(\alpha) - \frac{1}{3 \alpha d \mathcal{W}} + \frac{1}{3 \alpha^3 \ell^2 \lambda} \right)} \right] \sigma \kappa, \mathcal{W}(\alpha) = \alpha^3 [C + \ell^2 \lambda \ln(\alpha)]
\end{aligned}$$

3.4. Masa Holográfica y Masa Hamiltoniana en espacios cuánticos curvos.

$$\begin{aligned}
d \Sigma_{\kappa}^2 &= \frac{d \gamma^2}{1} - \kappa \gamma^2 + \frac{(1 - \kappa \gamma^2) d \langle \phi | \phi | \psi \rangle^2}{\| \tau \sigma \rho \|^2 \delta \xi}, \mathcal{E} = \int d \sigma^i \tau_{ij} \xi^j \int d \gamma d \langle \phi | \phi | \psi \rangle d \| \tau \sigma \rho \| \delta \mu^i \tau_{ij} \xi^j, d s^2 \\
&= \sigma_{ij} d \chi^i d \chi^j = \delta d \Sigma_{\kappa}^2, \mathcal{E} = \sigma \kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\omega - \frac{\alpha}{3 d \omega} \right) \right], \mathcal{E} \\
&= \sigma \kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{1}{3 \alpha} \frac{d \mathcal{W}}{d \alpha} - \frac{\alpha \gamma}{9} \right) \right] = \sigma \kappa \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{1}{3 \alpha} \frac{d \mathcal{W}}{d \alpha} - \frac{\alpha^3 \mathcal{C} \gamma}{9} \right) \right]
\end{aligned}$$



$$\mathcal{V}(\phi)$$

$$\begin{aligned}
&= \frac{\Lambda(v^2 - 4)}{6\kappa v^2 \left[v - \frac{1}{v} + 2e^{-\phi \ell v(v+1)} + v + \frac{1}{v} - 2e^{\phi \ell v(v-1)} + 4v^2 - \frac{1}{v^2} - 4e^{-\phi \ell v} \right]} \\
&+ \frac{\Upsilon}{\kappa v^2 \left[v - \frac{1}{v} + 2 \sinh \phi \ell v(v+1) - v + \frac{1}{v} - 2 \sinh \phi \ell v(v-1) + 4v^2 - \frac{1}{v^2} - 4 \sinh \phi \ell v \right]} \phi(\chi) \\
&= \ell_v^{-1} \sqrt{\frac{v^2 - 1}{2\kappa}} \ln \chi, f(x) = \frac{1}{\ell^2} + \Upsilon \left[\frac{1}{v^2} - 4 - \chi^2/v^2 \left(1 + \frac{\chi^{-v}}{v} - 2 - \frac{\chi^v}{v} + 2 \right) \right] + \frac{\chi}{\Omega(\chi)}, \Omega(\chi) \\
&= r^2 + \mathcal{O}(r^{-3}), \chi = 1 + \frac{1}{\eta r} + \frac{m^4}{r^3} + \frac{\eta}{r^4} + \frac{\rho}{r^5} + \mathcal{O}(r^{-6}), \Omega(\chi) \\
&= r^2 - 24m^4\eta^4 + v^2 - \frac{1}{12\eta} - 24m^4\eta^4 - v^2 + \frac{1}{12\eta^3 r} + 720m^4\eta^4 - 480|\rho\eta|^5 + v^4 - 20v^2 \\
&+ \frac{19}{240\eta^4 r^2} + \mathcal{O}(r^{-3}), \chi = 1 + \frac{1}{\eta r} - \frac{(v^2 - 1)}{23\eta^3 r^3 \left[1 - \frac{1}{\eta r} - \frac{9(v^2 - 9)}{80\eta^4 r^4} \right]} + \mathcal{O}(r^{-6}), -g_{tt} = f(x)\Omega(\chi) \\
&= \frac{r^2}{\ell^2} + 1 + \Upsilon + \frac{3\eta^4}{3\eta^3 r} + \mathcal{O}(r^{-3}), g_{rr} = \frac{\Omega(\chi)\eta^2}{f(x) \left(\frac{d\chi}{dr} \right)} \\
&= \frac{\ell^2}{r^2} - \frac{\ell^4}{r^4} - \frac{\ell^2(v^2 - 1)}{4\eta^4 r^2} - \frac{\ell^2(3\eta^2 \ell^2 + \Upsilon \ell^2 - v^2 + 1)}{3\eta^3 r^5} + \mathcal{O}(r^{-6}), \phi(\chi) = \ell_v^{-1} \sqrt{\frac{v^2 - 1}{2\kappa}} \ln \chi \\
&= \frac{1}{\ell_v \eta r} - \frac{1}{2\ell_v \eta^4 r^2} - v^2 - \frac{9}{24\eta^3 r^5} + \mathcal{O}(r^{-4}), \mathfrak{M} = \sigma \left[\frac{\mu}{\kappa} + \frac{1}{\ell^2} \left(\mathcal{W} - \frac{\alpha}{3} \frac{d\mathcal{W}}{d\alpha} \right) \right] \mathfrak{M} \\
&= -\frac{\sigma}{\kappa} \left(3\eta^2 + \frac{\Upsilon}{3\eta^4} \right)
\end{aligned}$$

$$\mathcal{I}_{\mathfrak{C}\mathfrak{F}\mathfrak{I}} \rightarrow \mathcal{I}_{\mathfrak{C}\mathfrak{F}\mathfrak{I}} + \ell_v/6 \int d^3\chi \mathcal{O}^3$$

3.5. Acción bulk on – shell en espacios cuánticos curvos.



$$\begin{aligned}
\mathcal{I} &= \int d^{\eta+1}\chi \sqrt{-g} \left[\frac{\mathcal{R}}{2\kappa} - \frac{(\partial\phi)^2}{2} - \mathcal{V}(\phi) \right] + 1/\kappa \int_{\partial\mathcal{M}}^{\delta} d^{\eta}\chi \sqrt{-\hbar} \mathfrak{K} + \mathcal{I}_g + \mathcal{I}_{\phi}, \mathcal{G}_{\mu\nu} = \kappa \mathcal{T}_{\mu\nu}, \mathcal{G}_{\mu\nu} \\
&= \mathcal{R}_{\mu\nu} - \frac{1}{2g_{\mu\nu}\mathcal{R}}, \mathcal{T}_{\mu\nu} = \partial_{\mu}\phi\partial_{\nu}\phi - g_{\mu\nu} \left[\frac{(\partial\phi)^2}{2} + \mathcal{V} \right] \mathcal{G} = -\frac{\mathcal{R}(\eta-1)}{2}, \mathcal{T} \\
&= -(\eta-1) \left[\frac{(\partial\phi)^2}{2} + \frac{\mathcal{V}(\eta+1)}{(\eta-1)} \right] \mathcal{G}_{\mu\nu} = \kappa \mathcal{T}_{\mu\nu} \rightarrow \frac{\mathfrak{R}}{2\kappa} = \frac{(\partial\phi)^2}{2} + \frac{\mathcal{V}(\eta+1)}{(\eta-1)}, \mathcal{I}_{bulk}^{\varepsilon} \\
&= -\frac{2}{\eta} - 1 \int d^{\eta+1}\chi \sqrt{g^{\mathfrak{E}}} \mathcal{V}(\phi)
\end{aligned}$$

3.6. Sistema de Coordenadas en espacios cuánticos curvos.

$$ds^2 = \Omega(x) \left[-f(x)dt^2 + \frac{\eta^2 d\chi^2}{f(x)} + d\Sigma_{\kappa}^2 \right]$$



$$\begin{aligned}
\mathfrak{E}_t^t - \mathfrak{E}_x^x &= 0 \Rightarrow 2\kappa\phi'^2 = \mathfrak{D} - \frac{2}{2\Omega^2[3(\Omega')^2 - 2\Omega\Omega'']}, \mathfrak{E}_t^t = -\frac{1}{\mathfrak{D}} - 2g^{\alpha\beta}\xi_{\alpha\beta} = 0 \\
\Rightarrow f'' + \mathfrak{D} - \frac{2}{2\Omega}\Omega'f' + 2\kappa\eta^2\mathfrak{E}_t^t + \frac{1}{\mathfrak{D}} - 2g^{\alpha\beta}\xi_{\alpha\beta} &= 0 \Rightarrow 2\kappa\mathcal{V} \\
&= -\mathfrak{D} - \frac{2}{2\eta^2\Omega^2\left[f\Omega'' + \mathfrak{D} - \frac{4}{2\Omega}f(\Omega')^2 + \Omega'f'\right]} + \frac{\kappa\left(\mathfrak{D} - \frac{2}{\Omega}\right)d}{d\chi\left[\Omega^{\frac{(\mathfrak{D}-2)}{2}}f'\right]} + 2\eta^2\kappa\Omega^{\frac{(\mathfrak{D}-2)}{2}} - \frac{2\eta^2\Omega^{\frac{\mathfrak{D}}{2}}(2\kappa\mathcal{V})}{\mathfrak{D}} - 2 \\
&= f\Omega''\Omega^{\frac{(\mathfrak{D}-4)}{2}} + \Omega'\left(f\Omega^{\frac{(\mathfrak{D}-4)}{2}}\right) - 2\eta^2\kappa\Omega^{\frac{(\mathfrak{D}-2)}{2}}, 2\kappa\mathcal{V} = -\frac{(\mathfrak{D}-2)}{2\eta^2\Omega^{\frac{\mathfrak{D}}{2}}\left[\Omega^{\frac{(\mathfrak{D}-4)}{2}}(f\Omega)'\right]'} d\Sigma_\kappa^2 \\
&= v_{ij}d\chi^i d\chi^j, \mathcal{J}_{bulk}^\mathfrak{E} = \frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta\left[\Omega^{\frac{(\mathfrak{D}-4)}{2}}(f\Omega)'\right]_{\chi_h}^{\chi_\beta} \hbar_{\alpha\beta}d\chi^\alpha d\chi^\beta} = \Omega(x)[-f(x)dt^2 + d\Sigma_\kappa^2], \eta_\alpha \\
&= \frac{\delta_\alpha^\chi}{\sqrt{g^{xx}}}, \mathcal{K}_{\alpha\beta} = \frac{\sqrt{g^{xx}}}{2}\partial_\chi\hbar_{\alpha\beta}, \mathcal{K} = \frac{1}{2\eta}\left(\frac{f}{\Omega}\right)^{\frac{1}{2}}\left[\frac{(\Omega f)'}{\Omega f} + \frac{(\mathfrak{D}-2)\Omega'}{\Omega}\right], \mathcal{J}_{\mathfrak{G}\mathfrak{S}}^\mathfrak{E} \\
&= -\frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta\Omega^{\frac{(\mathfrak{D}-2)}{2}}f\left[\frac{(f\Omega)'}{f\Omega} + \frac{(\mathfrak{D}-2)\Omega'}{\Omega}\right]}, \mathcal{J}_g \\
&= -\frac{1}{\kappa\int d^\eta\chi\sqrt{-\hbar}\left[\frac{(\eta-1)}{\ell} + \frac{\ell\mathcal{R}}{2(\eta-2)} + \frac{\ell^3}{2(\eta-4)(\eta-2)^2\left(\mathcal{R}^{\alpha\beta}\mathfrak{R}_{\alpha\beta} - \frac{\eta\mathcal{R}^2}{4(\eta-1)}\right)}\right]}\mathcal{R}_{ij} \\
&= \frac{(\eta-2)\kappa}{\Omega}\sigma_{ij}, \mathcal{R} = \frac{\kappa(\eta-2)(\eta-1)}{\Omega}, \mathcal{R}^{\alpha\beta}\mathfrak{R}_{\alpha\beta} = \frac{(\eta-2)^2(\eta-1)\kappa^2}{\Omega^2}\mathcal{R}^{\alpha\beta}\mathfrak{R}_{\alpha\beta} - \frac{\eta\mathcal{R}^2}{4(\eta-1)} \\
&= \frac{\kappa^2}{4\Omega^2}(\eta-2)^2(\eta-1)(\eta-4)\mathcal{J}_g^\mathfrak{E} = \frac{\beta\sigma_{\kappa,\eta-1}}{\kappa}\frac{(\mathfrak{D}-2)}{\ell}\sqrt{\Omega^{\mathfrak{D}-1}}f\left(1 + \frac{\ell^2\kappa}{2\Omega} - \frac{\ell^4\kappa^2}{8\Omega^2}\right)\beta^{-1} = \mathcal{T} \\
&= \frac{f'}{4\pi\eta|_{\chi_h}}, \delta = \mathcal{A}/4\mathfrak{G}
\end{aligned}$$

$$\begin{aligned}
\mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} &= \frac{1}{\mathcal{T}\left(\frac{\mathcal{AT}}{4\mathfrak{G}}\right)} - \frac{\sigma_{\mathcal{D}-2,\kappa}}{2\kappa\mathcal{T}} \Omega^{\frac{(\mathcal{D}-2)}{2}} (\mathcal{D}-2) \left[\frac{f\Omega'}{\eta\Omega} - \frac{\sqrt{\Omega f}}{2\ell\left(1 + \frac{\ell^2\kappa}{2\Omega} - \frac{\ell^4\kappa^2}{8\Omega^2}\right)} \right]_{\chi_\beta} ds^2 \\
&= -N(r)dt^2 + H(r)dr^2 + \delta(r)d\Sigma_\kappa^2, \Omega(x) \rightarrow \delta(r), f(x) \rightarrow \frac{N(r)}{\delta(r)}, \frac{\sqrt{NH}}{\eta\delta} dr \\
&\rightarrow d\chi, 2\kappa\mathcal{V} = -\mathcal{D} - \frac{2}{2\eta^2\Omega^{\frac{\mathcal{D}}{2}} \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} (f\Omega)' \right]}, \rightarrow 2\kappa\mathcal{V} = -\mathcal{D} - \frac{\frac{2}{2\delta^{\frac{(\mathcal{D}-2)}{2}} \sqrt{NH}} d}{dr \left(\frac{\delta^{\frac{(\mathcal{D}-2)}{2}}}{\sqrt{NH}} \frac{dN}{dr} \right)} \mathcal{J}_{bulk}^{\mathfrak{E}} \\
&= \frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} (f\Omega)' \right]_{\chi_h}}^{\chi_\beta} \rightarrow \mathcal{J}_{bulk}^{\mathfrak{E}} = \frac{\frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa} \frac{dN}{dr} \delta^{\frac{(\eta-1)}{2}}}{\sqrt{NH}} \Bigg|_{r_h}^{r_\beta} \hbar_{\alpha\beta} d\chi^\alpha d\chi^\beta \\
&= -N(\mathcal{R})dt^2 + \mathcal{S}(\mathcal{R})d\Sigma_\kappa^2, \eta_\mu = \frac{\delta_\mu^r}{\sqrt{g^{rr}}}, \mathcal{K}_{\mu\nu} = \frac{\sqrt{g^{rr}}}{2} \partial_r \hbar_{\mu\nu}, \mathcal{K} \\
&= \frac{1}{2\sqrt{H}} \left[\frac{N'}{N} + \frac{(\eta-1)\delta'}{\delta} \right], \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} = -\frac{\beta\sigma_{\kappa,\eta-1}}{2\kappa\eta\Omega^{\frac{(\mathcal{D}-2)}{2}} f \left[\frac{(f\Omega)'}{f\Omega} + \frac{(\mathcal{D}-2)\Omega'}{\Omega} \right]_{\chi_\beta}} \mapsto \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} \\
&= -\frac{\frac{\sigma_{\kappa,\eta-1}}{2\kappa\mathcal{T}} \delta^{\frac{(\mathcal{D}-2)}{2}}}{\sqrt{NH} \left[\frac{dN}{dr} + \frac{(\mathcal{D}-2)}{\delta} \frac{d\delta}{dr} \right]_{r_\beta}}, \mathcal{J}_g^{\mathfrak{E}} = \frac{\beta\sigma_{\kappa,\eta-1}(\eta-1)}{\ell\kappa\Omega^{\frac{(\mathcal{D}-1)}{2}} f^{\frac{1}{2}} \left(1 + \frac{\ell^2\kappa}{2\Omega} - \frac{\ell^4\kappa^2}{8\Omega^2} \right)_{\chi_\beta}} \mapsto \mathcal{J}_g^{\mathfrak{E}} \\
&= \beta\sigma_{\kappa,\eta-1}(\eta-1)/\ell\kappa\delta^{\frac{(\mathcal{D}-2)}{2}} N^{1/2} \left(1 + \frac{\ell^2\kappa}{2\delta} - \frac{\ell^4\kappa^2}{8\delta^2} \right)_{r_\beta} \beta^{-1} = \mathcal{T} = \frac{N'}{4\pi\sqrt{NH}|_{r_h}}, \delta \\
&= \mathcal{A}/4\mathfrak{G} \\
\mathcal{J}_{bulk}^{\mathfrak{E}} + \mathcal{J}_{\mathfrak{G}\mathfrak{H}}^{\mathfrak{E}} + \mathcal{J}_g^{\mathfrak{E}} &= -\frac{1}{\mathcal{T}\left(\frac{\mathcal{AT}}{4\mathfrak{G}}\right)} - \frac{\sigma_{\mathcal{D}-2,\kappa}}{2\kappa\mathcal{T}} \delta^{\frac{(\mathcal{D}-2)}{2}} (\mathcal{D}-2) \left[\frac{N\delta'}{\delta\sqrt{NH}} - \frac{2\sqrt{N}}{\ell\left(1 + \frac{\ell^2\kappa}{2\delta} - \frac{\ell^4\kappa^2}{8\delta^2}\right)} \right]_{r_\beta}
\end{aligned}$$

3.7. Ecuaciones de Movimiento en espacios cuánticos curvos.



$$\begin{aligned}
2\kappa\phi'^2 &= \frac{(\mathcal{D}-2)}{2\Omega^2[3(\Omega')^2 - 2\Omega\Omega'']} \frac{2\kappa\phi'^2}{(\mathcal{D}-2)} = \frac{1}{2\delta^2[\delta'^2 - 2\delta\delta'']} + \frac{\frac{\delta'}{2\delta}(\text{NH})'}{\text{NH}} \frac{d}{d\chi} \left[\delta^{\frac{(\mathcal{D}-2)}{2}} df/d\chi \right] \\
&= -2\eta^2\kappa\delta^{\frac{(\mathcal{D}-2)}{2}} \frac{d}{dr} \left[\frac{\delta^{\frac{\mathcal{D}}{2}}\sqrt{\text{NH}}d}{dr} \left(\frac{N}{\delta} \right) \right] = -2\kappa\sqrt{\text{NH}}\delta^{\frac{(\mathcal{D}-4)}{2}} 2\kappa\mathcal{V} \\
&= \frac{(\mathcal{D}-2)}{2\eta^2\Omega^{\frac{\mathcal{D}}{2}} \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} (f\Omega)' \right]} 2\kappa\mathcal{V} = - \frac{\frac{(\mathcal{D}-2)}{2\delta^{\frac{(\mathcal{D}-2)}{2}}\sqrt{\text{NH}}d}}{dr \left(\frac{\delta^{\frac{(\mathcal{D}-2)}{2}}}{\sqrt{\text{NH}}} \frac{dN}{dr} \right)} \partial_\chi \left[\Omega^{\frac{(\mathcal{D}-4)}{2}} f\phi' \right] \\
&= \frac{\eta^2\Omega^{\frac{\mathcal{D}}{2}}\partial\mathcal{V}}{\partial\phi} \partial_r \left(\delta^{\frac{(\mathcal{D}-2)}{2}} \phi' \sqrt{\frac{N}{H}} \right) = \sqrt{\text{NH}}\delta^{\frac{(\mathcal{D}-2)}{2}} \partial\mathcal{V}/\partial\phi
\end{aligned}$$

3.8. Cálculo de Potencial on – shell en espacios cuánticos curvos.



$$\begin{aligned}
\eta^2 \mathcal{V}(\phi) &= -\frac{f\Omega''}{f^2\Omega^2} - \frac{f'\Omega'}{\eta^2\Omega^2}, \eta^2 \mathcal{V}(\phi) \\
&= -\frac{1}{\chi^{2\nu-2}v^4 \left[\frac{1}{\ell^2} + \frac{\alpha}{2} \left(\chi^{2+\nu} - \frac{1}{2} + \nu + \frac{\chi^{2\nu-2}}{2} + \nu - \chi^2 + 1 \right) \right]} \\
&\quad \left[\begin{aligned} &\chi^{\nu-1}(\nu-1)^2v^2 - \frac{\chi^{\nu-1}(\nu-1)v^2}{\chi^2(\chi^\nu-1)^2\eta^2} \\ &-4\chi^{2\nu-1}(\nu-1)v^3 + \frac{2\chi^{2\nu-1}v^4}{(\chi^\nu-1)^4\eta^2\chi^2} + 2\chi^{2\nu-1}v^3 - \frac{4\chi^{2\nu-1}(\nu-1)v^3}{\eta^2\chi^2(\chi^\nu-1)^3} \end{aligned} \right] (\chi^\nu-1)^4\eta^4 \\
&\quad - \frac{1}{2\chi^{2\nu-2}v^4 \left(\frac{\chi^{\nu-1}(\nu-1)v^2}{\chi\eta^2(\chi^\nu-1)^2} - \frac{2\chi^{2\nu-1}v^3}{\chi\eta^2(\chi^\nu-1)^3} \right)} (\chi^{1+\nu} + \chi^{1-\nu} - 2\chi)\mathcal{V}(\phi) = \frac{f(x)v^2}{\chi^{2\nu-2}v^4} \\
&\quad \left(\chi^{3\nu-3}v^2 + 4\chi^{2\nu-3}v^2 + \chi^{\nu-3}v^2 + 3\chi^{3\nu-3}v - \frac{3\chi^{3\nu-3}}{(\chi^\nu-1)^4\eta^2} \right) (\chi^\nu-1)^4\eta^2 - \frac{f(x)v^2}{\chi^{2\nu-2}v^4} \\
&\quad \left(2\chi^{3\nu-3} - 4\chi^{2\nu-3} + \frac{2\chi^{\nu-3}}{(\chi^\nu-1)^4\eta^2} \right) (\chi^\nu-1)^4\eta^2 - \frac{\alpha\eta^4(\chi^\nu-1)^4}{2\chi^{2\nu-2}v^4} \\
&\quad \left(\frac{-v^2}{\eta^2(\chi^\nu-1)^3} \right) (\chi^{2\nu-2}(\nu+1) + \chi^{\nu-2}(\nu-1))(\chi^{1+\nu} + \chi^{1-\nu} - 2\chi)\mathcal{V}(\phi) = -\frac{f(x)\chi^{-\nu}}{\chi v^2} \\
&\quad (\chi^{2\nu}(\nu+1)(\nu+2) + 4\chi^\nu(\nu^2-1) + (\nu-1)(\nu-2)) + \frac{\frac{\alpha(\chi^\nu-1)}{2v^2\chi^{\nu-1}}\chi^{2\nu-2}}{\chi^{\nu-1}} \\
&\quad \frac{(\chi^{1+\nu} + \chi^{1-\nu} - 2\chi)(\chi^{-\nu}(\nu-1) + (\nu+1))\mathcal{V}(\phi)f(x)}{2\chi v^2} \\
&\quad (2\chi^\nu(\nu+1)(\nu+2) + 8(\nu^2-1) + 2\chi^{\nu-1}(\nu-1)(\nu-2)) + \frac{\alpha}{2v^2\chi^{\nu-1}(\chi^\nu-1)} \\
&\quad (1 + \nu + \chi^{-\nu}(\nu-1))(1 + \chi^{2\nu} - 2\chi^\nu)\mathcal{V}(\phi) = -\frac{f(x)}{2\chi v^2} \\
&\quad (2\chi^\nu(\nu+1)(\nu+2) + 8(\nu^2-1) + 2\chi^{\nu-1}(\nu-1)(\nu-2)) - \frac{\alpha}{2v^2\chi^{\nu-1}}(\chi^\nu-1)^2 \\
&\quad (2 - \chi^\nu(\nu+1) + \chi^{\nu-1}(\nu-1))
\end{aligned}$$

$$\mathcal{V}(\phi)$$

$$\begin{aligned}
&= - \frac{e^{-\ell_v \phi}}{2v^2 \left[\frac{1}{\ell^2} + \frac{\alpha}{2 \left(e^{(2+v)\ell_v \phi} - \frac{1}{2} + v + e^{(2-v)\ell_v \phi} - \frac{1}{2} - v - \chi^2 + 1 \right)} \right]} \\
&\quad [8(v^2 - 1) + 2(v + 1)(v + 2)e^{v\ell_v \phi} + 2(v - 1)(v - 2)e^{-v\ell_v \phi}] \\
&- \frac{\alpha e^{\ell_v \phi}}{2v^2 \left(\exp \frac{v\ell_v \phi}{2} - \exp - \frac{v\ell_v \phi}{2} \right)^2 [2 - e^{v\ell_v \phi}(v + 1) + e^{-v\ell_v \phi}(v - 1)]}, \mathcal{V}(\phi) \\
&= - \frac{(v^2 - 4)}{\ell^2 v^2 \left[\frac{(v - 1)}{(v + 2)e^{-v\ell_v \phi}(v + 1)} + \frac{(v + 1)}{(v - 2)e^{v\ell_v \phi}(v - 1)} + \frac{4(v^2 - 1)}{(v^2 - 4)e^{-\ell_v \phi}} \right]} \\
&- \frac{\alpha}{2v^2 \left[\frac{e^{(v+2)\ell_v \phi}}{(2 + v)} - \frac{e^{(2-v)\ell_v \phi}}{v} - 2 - e^{2\ell_v \phi} + \frac{v^2}{[v^2 - 4]} \right]} \\
&\quad [4(v^2 - 1)e^{-\ell_v \phi} + (v + 1)(v + 2)e^{(v-1)\ell_v \phi} + (v - 1)(v - 2)e^{(v+1)v\ell_v \phi}] - \frac{\alpha}{2v^2} \\
&\quad (e^{v\ell_v \phi} - 2 + e^{-v\ell_v \phi})[2e^{\ell_v \phi} - e^{(v+1)\ell_v \phi}(v + 1) + e^{(1-v)\ell_v \phi}(v - 1)] \\
&\quad \mathcal{V}(\phi) = \mathcal{V}_\Lambda(\phi) - \frac{\alpha}{2(v^2 - 4)} \\
&\quad \left[v^2(e^{\ell_v \phi(v-1)} - e^{-v\ell_v \phi(v+1)} - e^{\ell_v \phi(v-1)} + e^{-v\ell_v \phi(v+1)} + 4e^{-\ell_v \phi} - 4e^{\ell_v \phi}) + 3v \right. \\
&\quad \left. (e^{\ell_v \phi(v-1)} + e^{\ell_v \phi(v+1)} - e^{-\ell_v \phi(v-1)} - e^{-\ell_v \phi(v+1)}) + 2 \right. \\
&\quad \left. (e^{\ell_v \phi(v-1)} - e^{\ell_v \phi(v+1)} - e^{-\ell_v \phi(v-1)} + e^{-\ell_v \phi(v+1)}) - 4e^{-\ell_v \phi} + 4e^{\ell_v \phi} \right] \\
&\quad \mathcal{V}(\phi) = \mathcal{V}_\Lambda(\phi) - \frac{\alpha}{2(v^2 - 4)} \\
&\quad [(2v^2 + 6v + 4)\sinh \ell_v \phi(v - 1) - (2v^2 - 6v + 4)\sinh \ell_v \phi(v + 1) + 8(1 - v^2)\sinh \ell_v \phi] \\
&\quad \mathcal{V}(\phi) = \frac{\Lambda(v^2 - 4)}{3v^2 \left[\frac{(v - 1)}{(v + 2)e^{-\ell_v \phi(v+1)}} + \frac{(v + 1)}{(v - 2)e^{\ell_v \phi(v-1)}} + \frac{4(v^2 - 1)}{(v^2 + 4)e^{-\ell_v \phi}} \right] + \alpha} \\
&\quad \left[\frac{(v - 1)}{(v + 2)\sinh \ell_v \phi(v + 1)} - \frac{(v + 1)}{(v - 2)\sinh \ell_v \phi(v - 1)} + \frac{4(v^2 - 1)}{(v^2 + 4)e^{-\ell_v \phi}} \right]
\end{aligned}$$

3.9. Tensores diferenciales para espacios cuánticos curvos.



$$\begin{aligned}
d\chi^\mu \wedge d\chi^\nu &= d\chi^\mu \otimes d\chi^\nu - d\chi^\nu \otimes d\chi^\mu, d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_\rho} = \sum_{\sigma} (-1)^{|\sigma|} d\chi^{\sigma(\mu_1)} \bigotimes \dots \bigotimes d\chi^{\sigma(\mu_\rho)}, \mathfrak{H} \\
&= \frac{1}{\rho!} \mathcal{H}_{\mu_1 \dots \mu_\rho} d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_\rho}, \mathcal{A} = \Lambda_\mu d\chi^\mu, \mathcal{F} \\
&= \frac{1}{2\mathcal{F}_{\mu\nu} d\chi^\mu} \wedge d\chi^\nu \star (d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_\rho}) \\
&= \frac{\sqrt{-g}}{(\mathcal{D} - \rho)! \epsilon^{\mu_1 \dots \mu_\rho} \nu_{\rho+1} \dots \nu_{\mathcal{D}}} d\chi^{\nu_{\rho+1}} \wedge \dots \wedge d\chi^{\nu_{\mathcal{D}}}, d\mathcal{V} = d^{\mathcal{D}} \chi \sqrt{-g} \\
&= \frac{\sqrt{-g}}{\mathcal{D}!} \epsilon_{\mu_1 \dots \mu_{\mathcal{D}}} d\chi^{\mu_1} \wedge \dots \wedge d\chi^{\mu_{\mathcal{D}}}
\end{aligned}$$

4. Integral de Yang – Mills en espacios cuánticos curvos.

$$\begin{aligned}
&\exp \left(-\frac{1}{2} \int_{\mathbb{C}^4} d\lambda_4 |\kappa \mathfrak{d}A + A \wedge A|^2 \right) \mathfrak{D}\mathfrak{A} \\
&= \exp \left(-\frac{1}{2} \int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \mathfrak{d}A, A \wedge A \rangle + \langle A \wedge A, \kappa \mathfrak{d}A \rangle \right. \\
&\quad \left. + |A \wedge A|^2 \right) \exp \left(\int_{\mathbb{C}^4} d\lambda_4 |\kappa \mathfrak{d}A + A \wedge A|^2 \right) \mathfrak{D}\mathfrak{A} \\
&\exp \left(-\frac{1}{2} \int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \mathfrak{d}A, A \wedge A \rangle + \langle A \wedge A, \kappa \mathfrak{d}A \rangle + |A \wedge A|^2 \right) \int_{\mathbb{C}^4} d\lambda_4 \langle \kappa \mathfrak{d}A, A \wedge A \rangle \int_{\mathbb{C}^4} d\lambda_4 \langle A \wedge A, \kappa \mathfrak{d}A \rangle \\
&\quad \int_{\mathbb{C}^4} d\lambda_4 |A \wedge A|^2, \langle \kappa \mathfrak{d}A(\omega), d\chi^\alpha \wedge d\chi^\beta \rangle = \left(A, \widetilde{\xi}_{\alpha\beta}^\kappa(\omega) \right) [A_i A_j](\omega) \\
&= A_i(\omega) A_j(\omega) = \left(A \otimes A, \zeta_i(\omega) \otimes \zeta_j(\omega) \right), [A_{i,\alpha} A_{j,\beta} \overline{A_{i,\widehat{\alpha}} A_{j,\widehat{\beta}}}] (\omega) \\
&= \langle A \otimes A \otimes \bar{A} \otimes \bar{A}, \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\widehat{\alpha},\omega} \otimes \chi_{j,\widehat{\beta},\omega} \rangle \\
&A = \sum_{i=1}^3 \sum_{\alpha=1}^{\mathfrak{N}} A_{i,\alpha} \bigotimes d\chi^i \otimes \varepsilon^\alpha, \bar{A} = \sum_{i=1}^3 \sum_{\alpha=1}^{\mathfrak{N}} \overline{A_{i,\alpha}} \bigotimes d\chi^i \otimes \varepsilon^\alpha \\
&\int_{\omega \in \mathbb{R}^4} [A_{i,\alpha} A_{j,\beta} \overline{A_{i,\widehat{\alpha}} A_{j,\widehat{\beta}}}] (\omega) = \langle A^{\otimes 2} \otimes \bar{A}^{\otimes 2}, \int_{\omega \in \mathbb{R}^4} d\omega \chi_{i,\alpha,\omega} \otimes \chi_{j,\beta,\omega} \otimes \chi_{i,\widehat{\alpha},\omega} \otimes \chi_{j,\widehat{\beta},\omega} \rangle
\end{aligned}$$



$$\left(A_{i_1,\alpha_1}\otimes\cdots A_{i_3,\alpha_3}(\widetilde{\xi}_{\alpha\beta}^{\kappa}(\omega)\otimes\mathcal{E}^{\alpha_1})\otimes\tilde{\pi}_{\omega}^{\otimes^2}\right)=\left(A_{i_1,\alpha_1},\widetilde{\xi}_{\alpha\beta}^{\kappa}(\omega)\otimes\mathcal{E}^{\alpha_1}\right)\prod_{j=2}^3(A_{i_j,\alpha_j}\tilde{\pi}_{i_j,\alpha_j,\omega}^{\otimes^2})$$

$$\int\limits_{\mathbb{C}^4}^{\infty}d\lambda_4\left|A\wedge A\right|^2\triangleq\frac{\sum\gamma\sum_{\kappa=1}^3\sum_{1\leq i\leq j\leq 3}\sum_{\alpha\leq\beta\widehat{\alpha}\leq\widehat{\beta}}^{\infty}1}{2}\left\|c_{\gamma}^{\alpha\beta}c_{\gamma}^{\widehat{\alpha}\widehat{\beta}}\right\|(A_{i,\alpha}\otimes A_{j,\beta}\otimes A_{i,\widehat{\alpha}}\otimes A_{j,\widehat{\beta}}$$

$$\int\limits_{\omega\in\mathbb{R}^4}^{\infty}d\lambda^4\left(\omega\right)\tilde{\pi}_{\omega}^{\otimes^4})_{\#34}+(A_{i,\alpha}\otimes A_{j,\beta}\otimes A_{i,\widehat{\alpha}}\otimes A_{j,\widehat{\beta}}\int\limits_{\omega\in\mathbb{R}^4}^{\infty}d\lambda^4\left(\omega\right)\tilde{\pi}_{\omega}^{\otimes^4})_{\#12}$$

$$\int\limits_{\mathbb{C}^4}^{\infty}d\lambda_4\langle\kappa\mathfrak{d}A,A\wedge A\rangle=\sum_{\gamma}^{\infty}\sum_{\kappa=1}^3\sum_{1\leq i\leq j\leq 3}^{\infty}\sum_{\alpha\leq\beta\widehat{\alpha}\leq\widehat{\beta}}^{\infty}\left|c_{\gamma}^{\alpha\beta}\right|\left(\int\limits_{\omega\in\mathbb{C}^4}^{\infty}d\lambda_4(\omega)(\widetilde{\xi}_{ij}^{\kappa}(\omega)\otimes\mathcal{E}^{\gamma})\right)\tilde{\pi}_{\omega}^{\otimes^4})_{\#23}$$

$$\int\limits_{\mathbb{C}^4}^{\infty}d\lambda_4\langle A\wedge A,\kappa\mathfrak{d}A\rangle=\sum_{\gamma}^{\infty}\sum_{\kappa=1}^3\sum_{1\leq i\leq j\leq 3}^{\infty}\sum_{\alpha\leq\beta\widehat{\alpha}\leq\widehat{\beta}}^{\infty}\left|c_{\gamma}^{\alpha\beta}\right|\left(\int\limits_{\omega\in\mathbb{C}^4}^{\infty}d\lambda_4(\omega)(\widetilde{\xi}_{ij}^{\kappa}(\omega)\otimes\mathcal{E}^{\gamma})\right)\tilde{\pi}_{\omega}^{\otimes^4})_{\#3}$$

$$\mathfrak{I}_{\mathfrak{S}}^{\kappa}\left(\left\{A_{i,\alpha}\right\}_{i,\alpha}\right)$$

$$=\mathfrak{I}\mathrm{r}\,\mathcal{\hat{T}}\exp\left[\sum_{\alpha}^{\infty}(A_{\alpha},\xi_{\alpha\beta}^{\kappa}\left(\frac{\kappa\sigma(s,t)}{2}\otimes\mathcal{E}^{\alpha}\right)\otimes\rho(\mathcal{E}^{\alpha}))+\sum_{1\leq i\leq j\leq 3}^{\infty}\sum_{\alpha<\beta}^{\eta}\sum_{p,q=1}^{\eta}|\mathfrak{I}_{ij}^{\sigma}|\left(\frac{p}{n},\frac{q}{n}\right)/\eta^2\left(s,t\right)\right.\\ \left.\sum_{\alpha<\beta}^{\infty}\sum_{\gamma}^{\infty}\left|c_{\gamma}^{\alpha\beta}\right|\left(A_{i,\alpha}\otimes A_{j,\beta}\tilde{\pi}_{\frac{\kappa\sigma(s,t)}{2}}^{\otimes^4}\right)\otimes\rho(\mathcal{E}^{\gamma})\mu_{s,t}\right]$$

$$\begin{aligned}
& y^\kappa\left(\{A_{i,\alpha}\}_{i,\alpha}\right) \\
&= \left| \exp -\frac{1}{2} \sum_{\gamma}^{\infty} \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3}^{\infty} \sum_{\alpha \leq \beta \hat{\alpha} \leq \hat{\beta}}^{\infty} \left| c_{\gamma}^{\alpha \beta} \right| \left\langle \int_{\omega \in \mathbb{C}^4}^{\infty} d\lambda_4(\omega) (\xi_{ij}^{\kappa}(\omega) \otimes \mathcal{E}^{\alpha}) \tilde{\pi}_{\omega}^{\otimes^4} \pi_{\omega} \otimes \widehat{\pi_{\omega}} \right\rangle_{\#23} \right. \\
&\quad \left. - \frac{1}{2 \sum_{\gamma}^{\infty} \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3}^{\infty} \sum_{\alpha \leq \beta \hat{\alpha} \leq \hat{\beta}}^{\infty} \left| c_{\gamma}^{\alpha \beta} \right| \left\| \int_{\omega \in \mathbb{C}^4}^{\infty} d\lambda_4(\omega) \pi_{\omega}^{\otimes^4} \otimes (A_{\gamma} \xi_{ij}^{\kappa}(\omega) A_{\alpha}(\omega) \otimes \mathcal{E}^{\gamma}) \right\| \tilde{\pi}_{\omega}^{\otimes^4} \right\rangle_{\#3}} \\
&\quad - \frac{1}{2 \sum_{\gamma}^{\infty} \sum_{\kappa=1}^3 \sum_{1 \leq i \leq j \leq 3}^{\infty} \sum_{\alpha \leq \beta \hat{\alpha} \leq \hat{\beta}}^{\infty} \left\| c_{\gamma}^{\alpha \beta} \right\| \left\| \widehat{c_{\gamma}^{\alpha \beta}} \right\|} \left\langle A_{i,\alpha} \otimes A_{j,\beta} \otimes A_{\widehat{i},\widehat{\alpha}} \otimes A_{\widehat{j},\widehat{\beta}} \int_{\omega \in \mathbb{R}^4}^{\infty} d\lambda^4(\omega) \tilde{\pi}_{\omega}^{\otimes^4} \right\rangle^2 \left[\begin{array}{c} \rho \sigma \triangleq \\ N_{\omega}^2 \prod_{i=1}^4 1 \\ \psi \varphi \varrho \kappa \end{array} \right]
\end{aligned}$$

$$/2\varpi e^{-|\omega|^2}|\chi_{\omega}|^2\langle dp_idq_j\mathbb{E}_{\mathfrak{Y}\mathfrak{M}}^{\kappa}\rangle|\mathcal{N}_{\omega}|^2\|\mathcal{Y}_{\rho}^{\kappa}\mathcal{J}_{\delta}^{\kappa}\|$$

$$\mathbb{E}(\mathcal{Y}^{\kappa}\mathcal{Y}_{\delta}^{\kappa})\left[\frac{\exp\frac{\frac{1}{2}3\rho\sigma}{2\pi\varpi^4\sum_{j=1}^{\mathsf{M}(\eta)}e^{-|\beta(j)|^2}}1}{\eta^8\mathsf{N}_{\beta(\kappa)}^2}\right]=\exp\frac{1}{2}\mathsf{A}\ast\mathsf{A}\langle\mathsf{N}_1\cdots\widehat{\mathsf{N}_{\mathsf{M}(\eta)}}\rangle^{\mathfrak{T}}=(1/\det(1-\mathsf{A}\ast\mathsf{A}))^{\frac{1}{2}}$$

$$\leq \exp\left(\frac{c}{2}\mathfrak{T}\mathfrak{r}(\mathsf{A}\ast\mathsf{A})\right)=\exp\left\langle\frac{\frac{1}{2}3\rho c}{2\pi\varpi^4\sum_{j=1}^{\mathsf{M}(\eta)}e^{-|\beta(j)|^2}}\frac{1}{\eta^8}\right\rangle$$

$$\rightarrow \exp \frac{\left[\frac{\frac{1}{2}3c}{2\pi\varpi^4\int_{\omega\in\mathbb{C}^4}^{\infty}e^{-|\omega|^2}\prod_{i=1}^4\kappa\sigma\mathfrak{G}(s,t)dp_idq_j\tilde{\pi}_{\frac{\kappa\sigma(s,t)}{2}}^{\otimes^4}}\right]\rho\kappa}{4\int_{\mathfrak{Z}^2}^{\infty}dsdt}$$

$$\langle \mathsf{M}_{\frac{i,\kappa\sigma\mathfrak{G}(s,t)}{2}}^{\alpha} \otimes \mathsf{M}_{\frac{j,\kappa\sigma\mathfrak{G}(s,t)}{2}}^{\beta} \rangle \langle \mathsf{M}_{\frac{i,\kappa\sigma}{(\frac{p}{n},\frac{q}{n})/\eta^2}}^{\alpha} \otimes \mathsf{M}_{\frac{j,\kappa\sigma}{(\frac{p}{n},\frac{q}{n})/\eta^2}}^{\beta} \rangle \langle \pi_i | \omega |_{\sigma \kappa} \pi_j | \omega |_{\sigma \kappa} \rangle \int 4 \eta \sum_{1 \alpha < \beta < 3}^{\eta} \frac{\Sigma_{p=1}^{\eta} \Sigma_{q=1}^{\eta} |\Im_{\alpha \beta}^{\sigma}| \left(\frac{p}{n}, \frac{q}{n}\right)}{1}$$

$$/\eta^2\left(\frac{\kappa}{4}\right)/(2\pi)^2\kappa^4\langle 2/\kappa\sqrt{2\pi}\rangle^4$$

$$\begin{aligned}
& \left| \sum_{\alpha < \beta}^{\eta} |c_{\gamma}^{\alpha\beta}| M_{\frac{i,\kappa\sigma(s,t)}{2}}^{\alpha} \otimes M_{\frac{j,\kappa\sigma(s,t)}{2}}^{\beta} \right| \\
& < N \|\alpha(\psi)\| \|\beta(\hat{\psi})\| \sqrt{\sum_{\alpha < \beta}^{\eta} \sqrt{M_{\frac{i,\kappa\sigma(s,t)}{2}}^{\alpha}} \sqrt{M_{\frac{j,\kappa\sigma(s,t)}{2}}^{\beta}} \sqrt{\frac{M_{\frac{i,\kappa\sigma}{\frac{(\underline{p},\underline{q})}{n'n}}}{1/\eta^2}}^{\alpha}} \sqrt{\frac{M_{\frac{j,\kappa\sigma}{\frac{(\underline{p},\underline{q})}{n'n}}}{1/\eta^2}}^{\beta}}} \\
& \left| \sum_{\alpha < \beta}^{\eta} |c_{\gamma}^{\alpha\beta}| M_{\frac{\frac{(\underline{p},\underline{q})}{n'n}}{\frac{1}{\eta^2}}}^{\alpha} \otimes M_{\frac{\frac{(\underline{p},\underline{q})}{n'n}}{\frac{1}{\eta^2}}}^{\beta} \right| \left| \sum_{\alpha < \beta}^{\eta} |c_{\gamma}^{\alpha\beta}| M_{\frac{\alpha,\kappa\sigma(s,t)}{2}}^{\alpha} \otimes M_{\frac{\beta,\kappa\sigma(s,t)}{2}}^{\beta} \right| \int_{\mathbb{S}^2}^{\infty} |\sigma'_{\alpha} \sigma'_{\beta} \sigma'_{\beta} \sigma'_{\alpha}|(s,t) \, ds dt \sum_{\gamma}^{\eta} \|\mathcal{B}(\gamma)\| \\
& \rightarrow \sum_{\alpha\beta\gamma\delta\epsilon\zeta\eta\theta\iota\kappa\lambda\mu\nu\xi\omicron\pi\omega\rho\sigma\tau\upsilon\phi\chi\psi\omega}^{\infty} \|\mathcal{B}(\alpha\beta\gamma\delta\epsilon\zeta\eta\theta\iota\kappa\lambda\mu\nu\xi\omicron\pi\omega\rho\sigma\tau\upsilon\phi\chi\psi\omega)\|^{\infty} \\
& \frac{1}{\frac{1}{\mathfrak{I}\mathfrak{R} \int_{\Lambda}^{\infty} \mathfrak{I} e^{\int_{\mathbb{C}}^{\vee} \Sigma_i \Lambda_i \otimes d\chi^i} e^{-\frac{1}{2 \int_{\mathbb{R}^4} \Sigma_i |d\Lambda + \wedge \wedge \Lambda|^2 \otimes d\chi^i} \mathcal{D}\mathcal{A}}} = \mathbb{E}_{\mathcal{YM}}^{\kappa} \|\mathcal{J}_{\delta}\| \\
& = \frac{1}{\mathbb{E} \|\mathcal{Y}^{\kappa}\| \mathbb{E} |\mathcal{J}_{\delta}^{\kappa} \cdot \mathcal{Y}_{\rho\sigma}^{\kappa}|_{\mathbb{R}^4} \int_{\mathbb{C}^4}^{\infty} d\lambda_4 \sqrt{\frac{3}{2\pi}} \sqrt{\frac{2}{\kappa^4 \sqrt{2\pi}}} \left\langle \frac{2}{\kappa \sqrt{2\pi}} \right\rangle^4 = \int_{\omega \in \mathbb{C}^4}^{\infty} \widehat{\xi_{ij}^{\kappa}}(\omega) \left\| \pi_{\frac{\kappa\sigma\mathfrak{G}(s,t)}{2}}^{\otimes 4} \pi_{\frac{\frac{(\underline{p},\underline{q})}{n'n}}{\frac{1}{\eta^2}}}^{\otimes 4} \right\| \left\| \pi_{\frac{\kappa\sigma\mathfrak{G}(s,t)}{2}}^{\otimes 4} \pi_{\frac{\frac{(\underline{p},\underline{q})}{n'n}}{\frac{1}{\eta^2}}}^{\otimes 4} \right\| \left| \widehat{\pi_{\omega}^{\otimes 4}} \right|_{\kappa\sigma} |\pi_{\omega}^{\otimes 4}|_{\rho\sigma}^{\kappa\sigma} \\
& \otimes \mathcal{E}_{s,t}^{\gamma} \sum_{\gamma}^{\infty} \sum_{\alpha < \beta}^{\infty} \|c_{\gamma}^{\alpha\beta}\| \langle A_{i,\alpha} \otimes A_{j,\beta} \pi_{\frac{\kappa\sigma\mathfrak{G}(s,t)}{2}}^{\otimes 4} \rangle \otimes \mathcal{E}_{s,t}^{\alpha} \int_{\mathcal{I}}^{\eta} d\tau \mathcal{P}_{s,t}^{i'}(\tau) \varpi_{\rho_{s,t}(\tau)} \otimes d\chi^i \otimes \mathcal{E}_{s,t}^{\beta} \otimes^4 \mu \sum_{\alpha\beta}^{\eta} |\rho(\mathcal{E}^{\alpha}) \otimes \rho(\mathcal{E}^{\beta})|^{\eta} \int_{\delta}^{\kappa} \rho \delta \otimes \mu \\
& \boxtimes \prod_{i,i=1}^{\eta} \int_{\mathcal{I}^{2\eta}}^{\eta} ds_i \otimes dt_i \otimes ds_j \otimes dt_j
\end{aligned}$$

5. Supersimetría de Yang- Mills en espacios cuánticos curvos.

$$\begin{aligned}
\mathcal{O}_\rho &= (\chi, \gamma) = \text{Tr}[\phi(\chi, \gamma)^\rho], \mathcal{O}(\chi, \gamma) \\
&= \sum_{\rho=2}^{\infty} \frac{1}{\rho} \left(\frac{16\pi^4}{c} \right)^{\rho/4} \mathcal{O}_\rho(\chi, \gamma), \langle \mathcal{O}(\chi_1, \gamma_1) \mathcal{O}(\chi_2, \gamma_2) \mathcal{O}(\chi_3, \gamma_3) \mathcal{O}(\chi_4, \gamma_4) \rangle \\
&+ \frac{J_4(\chi_i, \gamma_j)}{2c} \sum_{\ell=1}^{\infty} \left(\frac{\lambda}{4\pi^2} \right)^\ell \frac{1}{\ell!} \frac{\int d^4 \chi_5}{(-4\pi^2)} \otimes \frac{d^4 \chi_{4+\ell}}{(-4\pi^2) f^{(\ell)}(\chi_{ij}^2)}, \chi_{ij}^2 \cong \chi_{ij}^2 - \gamma_{ij}^2 = \chi_{ij}^2 (1 - g_{ij}), f^{(\ell)}(\chi_{ij}^2) \\
&= \sum_{\alpha} c_{\alpha}^{(\ell)} f_{\alpha}^{(\ell)}(\chi_{ij}^2), f_{\alpha}^{(\ell)}(\chi_{ij}^2) \\
&= \frac{1}{|aut(\alpha)|} \sum_{\sigma \in \delta_{\ell+4}} \prod_{i,j=1}^{4+\ell} \frac{1}{(\chi_{\sigma i \sigma j}^2)^{e_{ij}^{\alpha}}} \int d\mu \bigotimes - \frac{1}{\pi^2} \int d^4 \chi_1 \boxtimes d^4 \chi_2 \boxplus d^4 \chi_3 \\
&\boxplus \frac{d^4 \chi_4}{vol[\mathcal{SO}(2,4)]}, \mathcal{C}(\lambda; g_{ij}) \\
&- \sum_{\ell=1}^{\infty} \left(\frac{\lambda}{4\pi^2} \right)^\ell \otimes \int d^4 \chi_1 \boxtimes d^4 \chi_2 \boxplus d^4 \chi_3 \boxplus \frac{\frac{d^4 \chi_4}{vol[\mathcal{SO}(2,4)]} f^{(\ell)}(\chi_{ij}^2 (1 - g_{ij}))}{\pi^2 \ell!} (-4\pi^2)^\ell, \mathcal{C}(\lambda; g_{ij}) \\
&= \sum_{\ell=1}^{\infty} \frac{\left(\frac{\lambda}{4\pi^2} \right)^\ell \otimes 1}{\ell! (-4)^{\ell+1} \sum_{\alpha} c_{\alpha}^{(\ell)} \mathcal{P}_{f_{\alpha}^{(\ell)}} f_{\alpha}^{(\ell)}} (1 - g_{ij}), \mathcal{P}_{f_{\alpha}^{(\ell)}} \\
&= \frac{1}{(\pi^2)^{\ell+1} \int d^4 \chi_1 \boxtimes d^4 \chi_2 \boxplus d^4 \chi_3 \boxplus \frac{d^4 \chi_4}{vol[\mathcal{SO}(2,4)]} f_{\alpha}^{(\ell)}(\chi_{ij}^2), f^{(1)}(\chi_{ij}^2)} = \frac{1}{\prod \psi_{1 \leq i \leq j \leq 5} \chi_{ij}^2} - \frac{1}{1! (-4)^1} \mathcal{P}_{f^{(1)}} \\
&/ \prod_{1 \leq i \leq j \leq 4} \psi (1 - g_{ij})
\end{aligned}$$

$$\begin{aligned}
f^{(2)}(\chi_{ij}^2) &= \frac{1}{48} \sum_{\sigma \in \delta_6} \frac{\chi_{\sigma_1 \sigma_2}^2 \chi_{\sigma_3 \sigma_4}^2 \chi_{\sigma_5 \sigma_6}^2}{\prod \psi_{1 \leq i \leq j \leq 5} \chi_{ij}^2} - \frac{\mathcal{P}_{f^{(2)}}}{2!(-4)^2} g_{12} g_{34} + g_{13} g_{24} + g_{14} g_{23} - 3 \sum_{1 \leq i \leq j \leq 4} g_{ij} \\
&+ \frac{15}{\prod \psi_{1 \leq i \leq j \leq 4}} (1 - g_{ij}) - \frac{\mathcal{P}_{f^{(2)}}}{2!(-4)^2} \left(\frac{12}{1} - g_{34} + 3 \right) g_{34}^2 g_{34}^2 \pi^6 \partial \zeta, \mathcal{C}(\lambda; \gamma_i \gamma_j) \\
&= \frac{\sum_{\ell=1}^{\infty} \lambda^{\ell} \sum_{v=2}^{\infty} 4(-1)^{v+\ell+1} \Gamma\left(\ell + \frac{3}{2}\right)^2 \zeta(2\ell+1)}{\pi^{2\ell+1} \Gamma(\ell+2-v) \Gamma(\ell+v+1)} \mathcal{F}_v(\gamma_i), \mathcal{F}_v(\gamma_i) \\
&= \frac{\delta_{v-2, v-2, 1, 1}(\gamma_i)}{\prod \psi_{1 \leq i \leq j \leq 4}} (1 - \gamma_i \gamma_j), \delta_{v_1, v_2, v_3, v_4}(\gamma_i) \\
&= \det(\gamma_i^{4+v_j-j})_{i,j=1,2,3,4} \prod_{1 \leq i \leq j \leq 4} \psi(\gamma_i - \gamma_j), \mathcal{C}_{\rho_1, \rho_2, \rho_3, \rho_4}(\lambda) \\
&\boxtimes \mathcal{C}(\lambda; \gamma_i \gamma_j) \Big|_{\gamma_1^{\rho_1-2} \gamma_2^{\rho_2-2} \gamma_3^{\rho_3-2} \gamma_4^{\rho_4-2}}
\end{aligned}$$

$$\begin{aligned}
\zeta(\eta) \Gamma(\eta+1) &= 2^{\eta-1} \int_0^{\infty} \frac{d\omega \omega^{\eta}}{\sinh^2(\omega)}, \mathcal{C}(\lambda; \gamma_i \gamma_j) \\
&= \int_0^{\infty} \frac{\omega d\omega}{\sinh^2(\omega)} \sum_{v=2}^{\infty} [\mathcal{J}_{v-1}(\mu)^2 - \mathcal{J}_v(\mu)^2] \mathcal{F}_v(\gamma_i), \mathcal{C}_{2,2,\rho,\rho}(\lambda) \\
&= \int_0^{\infty} \frac{\omega d\omega}{\sinh^2(\omega)} (\mathcal{J}_l(\mu)^2 - \mathcal{J}_j(\mu)^2), \mathcal{C}_{3,3,\rho,\rho}(\lambda) \\
&= \int_0^{\infty} \frac{\omega d\omega}{\sinh^2(\omega)} (3\mathcal{J}_1(\mu)^2 - 4\mathcal{J}_2(\mu)^2 - 2\mathfrak{I}_{\rho+1}(\mu)^2)
\end{aligned}$$

$$\mathcal{C}(\lambda; \gamma_i \gamma_j) \Big|_{strong}$$

$$\begin{aligned}
&= \sum_{v=2}^{\infty} \left(\frac{1}{2v(v-1)} \right. \\
&\quad + \sum_{\eta=1}^{\infty} 4\eta(-1)^{\eta} \Gamma\left(\eta + \frac{1}{2}\right) \Gamma\left(v + \eta - \frac{1}{2}\right) \zeta(2\eta+1) \\
&\quad \left. / \lambda^{\eta+\frac{1}{2}} \sqrt{\pi} \Gamma\left(\eta - \frac{1}{2}\right) \Gamma\left(v - \eta + \frac{1}{2}\right) \right) \mathcal{F}_v(\gamma_i)
\end{aligned}$$

$$\begin{aligned}
\Delta \mathcal{C}(\lambda; \gamma_i \gamma_j) &= \pm \frac{\iota}{2} \sum_{v=2}^{\infty} (-1)^v (2v-1)^2 \left(\frac{8Li_0(z)}{(2v-1)^2} + \frac{2Li_1}{\lambda^{\frac{1}{2}}} \right. \\
&\quad \left. + \left(4v^2 - 4v + \frac{5Li_2(z)}{4\lambda} \right) \right) \mathcal{F}_v(\gamma_i), \mathcal{E}(\delta; \tau, \tilde{\tau}) \\
&= \sum_{(m,n) \neq (0,0)} \tau_2^\delta / \pi^\delta |m + \eta\tau|^{2\delta}, \mathfrak{D}_{\mathfrak{N}}(\delta; \tau, \tilde{\tau}) \\
&= \frac{\sum_{(m,n) \neq (0,0)} e^{-\frac{4\sqrt{\mathfrak{N}}\pi|m+\eta\tau|}{\sqrt{\tau_2}}} \tau_2^\delta}{\pi^\delta |m + \eta\tau|^{2\delta}}, \mathcal{C}(\tau, \tilde{\tau}; \gamma_i \gamma_j) \\
&= \sum_{v=2}^{\infty} \left[\frac{1}{2(v-1)v} - 2v - \frac{1}{2^4 \mathcal{N}^{\frac{3}{2}} \mathfrak{E}\left(\frac{3}{2}; \tau, \tilde{\tau}\right)} + 3(2v-3)(4v^2-1)/2^8 \mathcal{N}^{\frac{5}{2}} \mathfrak{E}\left(\frac{5}{2}; \tau, \tilde{\tau}\right) \right. \\
&\quad \left. \pm 2\iota(-1)^v \mathfrak{D}_{\mathfrak{N}}(0; \tau, \tilde{\tau}) \right] \mathcal{F}_v(\gamma_i), \lim_{\chi_{i,i+1}^2 \rightarrow 0} \frac{\langle \mathcal{O}\mathcal{O}\mathcal{O}\mathcal{O} \rangle}{\langle \mathcal{O}\mathcal{O}\mathcal{O}\mathcal{O} \rangle_{free}} \\
&= \mathcal{M}^2 - \sum_{\ell=1}^{\infty} (\lambda/4\pi^2)^\ell \ell + \frac{1}{2^{2\ell+1}} \begin{pmatrix} 2\ell & \cdots & 2 \\ \vdots & \ddots & \vdots \\ \ell & \cdots & 1 \end{pmatrix} (\gamma^2 - 1)^{2\ell} / \gamma^{2\ell} \zeta(2\ell + 1)
\end{aligned}$$

6. Modelo de interacción de partículas y antipartículas en espacios cuánticos curvos.

6.1. Comportamiento de las partículas y antipartículas deformantes del espacio cuántico.

$$\begin{aligned}
\mathcal{H}_{int} &= \frac{1}{2\hbar_{\mu\nu}\mathcal{T}^{\mu\nu}}, \Gamma_{spon} = \frac{2\pi}{\hbar^2 \langle f | \hat{\mathcal{H}}_{int} | 1 \rangle^2 \mathcal{D}(\omega)}, \hat{\mathcal{H}}_{int,\ell} = \frac{\frac{\mathcal{L}}{\pi^2 \sqrt{\frac{\mathcal{M}\hbar}{\omega_\ell}}} (-1)^{\ell-\frac{1}{2}}}{\ell^2 (\hat{\beta}_\ell \hat{\beta}_\ell^\dagger) \hat{\hbar}}, \Gamma_{spon} = \frac{8\mathfrak{G}\mathfrak{M}\mathcal{L}^2 \omega_\ell^4}{\ell^4 \pi^4 c^5} \\
&= \frac{8\pi\mathfrak{G}\rho v_\delta^4 \mathcal{R}^2}{\mathcal{L}c^5}, \Gamma_{stim} = \frac{\mathcal{M}\mathcal{L}^2 \omega_\ell^2 \hbar^2}{4\ell^4 \pi^4 \hbar} = \frac{v_\delta^2}{4\ell^2 \pi^3 \hbar} \mathcal{M} h^2, \mathcal{M} \\
&= \frac{\pi^2 \hbar \omega^3}{v_\delta^2 \chi(\hbar, \omega, t)^2}, m \dot{\xi}_\eta + m \omega_\mathfrak{D}^2 (2\xi_\eta - \xi_{\eta-2} - \xi_{\eta+2}), \xi_\eta(t) \\
&= e^{-\iota\omega t} \left(\Lambda e^{\frac{\iota\kappa\eta\alpha}{2}} + \beta e^{-\frac{\iota\kappa\eta\alpha}{2}} \right) + \mathcal{H} \otimes c, \frac{d\xi_\eta}{d\eta} \Big|_{\eta=\pm N} \xi_\eta(t) \\
&= \sum_{\ell=0,2}^{N-1} \chi_\ell(t) \cos \left[\frac{\ell\pi\eta}{(2N)} \right] + \sum_{\ell=1,3}^N \chi_\ell(t) \sin \left[\frac{\ell\pi\eta}{(2N)} \right]
\end{aligned}$$



$$\begin{aligned}
& \sum_{\eta=\pm N}^N \cos[\ell\pi\eta/2(N+1)] \cos[\ell'\pi\eta/2(N+1)] \\
&= N+1/2 \delta_{\ell\ell'} \sum_{\eta=\pm N}^N \sin[\ell\pi\eta/2(N+1)] \sin[\ell'\pi\eta/2(N+1)] \\
&= N+1/2 \delta_{\ell\ell'}, \sum_{\eta=\pm N}^N \cos[\ell\pi\eta/2(N+1)] \sin[\ell'\pi\eta/2(N+1)], \\
\xi_\eta(t) &= \sum_{\ell=0,2}^{N-1} \chi_\ell(t) \cos\left[\frac{\ell\pi\eta}{(2N+1)}\right] + \sum_{\ell=1,3}^N \chi_\ell(t) \sin\left[\frac{\ell\pi\eta}{(2N+1)}\right], \ddot{\chi}_\ell + \omega_\ell^2 \chi_\ell, \mathfrak{E} \\
&= \frac{m}{2} \sum_{\eta=\pm N} \xi_\eta^2 + \frac{m\omega_{\mathfrak{D}}^2}{2} \sum_{\eta=\pm N}^{N-2} (\xi_{\eta-2} - \xi_\eta)^2 = \frac{m}{4} \sum_{\ell=0}^N \dot{\chi}_\ell^2 + \frac{\mathcal{M}}{4} \sum_{\ell=0}^N \omega_\ell^2 \chi_\ell^2 \int_{-\frac{\mathcal{L}}{2}}^{\frac{\mathcal{L}}{2}} dx \chi_\ell(\chi)^2 \\
&= \frac{\mathcal{L}}{2}, \mu_\ell = \int_{-\frac{\mathcal{L}}{2}}^{\frac{\mathcal{L}}{2}} dx \rho(\chi) \chi_\ell(\chi)^2 = \frac{\mathcal{M}}{\mathcal{L}} \int_{-\frac{\mathcal{L}}{2}}^{\frac{\mathcal{L}}{2}} dx \chi_\ell(\chi)^2 = \frac{\mathcal{M}}{2}, f(x) = -m\nabla\phi \\
&= -m\nabla\left(\frac{1}{2}\frac{\partial^2\phi}{\partial\chi^2}\chi^2\right) = m\frac{\hbar_{xx}}{4}\nabla(\chi^2) = m\frac{\hbar_{xx}}{2}(\chi_\eta + \xi_\eta) \\
\hat{\mathcal{H}}_\ell &= -m\frac{\hbar_{xx}}{2} \sum_{\eta=\pm N}^N \left(\chi_\eta \xi_\eta + \frac{\xi_\eta^2}{2}\right) - m\frac{\hbar_{xx}}{2} \sum_{\eta=\pm N}^N \chi_\eta \xi_\eta \approx -\frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{\pi^2 \sum_{\ell=1,3}^N \frac{(-1)^{\ell-\frac{1}{2}}1}{\ell^2 \chi_\ell(t)} - m\frac{\hbar_{xx}}{4} \sum_{\eta=\pm N}^N \xi_\eta^2} \\
&= -\frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{8} \sum_{\ell=0}^N \chi_\ell^2, \hat{\mathcal{H}}_\ell = \sum_{\ell=0}^N \hat{\mathcal{H}}_\ell^l = \frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{\pi^2 \sum_{\ell=1,3}^N \frac{(-1)^{\ell-\frac{1}{2}}1}{\ell^2 \chi_\ell}} - \frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{8} \sum_{\ell=0}^N \hat{\chi}_\ell^2, \hat{\mathcal{H}}_\ell^{l,odd} \\
&= -\frac{\mathcal{M}\mathcal{L}\hbar_{xx}}{\pi^2} \frac{(-1)^{\ell-\frac{1}{2}}1}{\ell^2} \sqrt{\frac{\hbar}{\mathcal{M}\omega_\ell}} (\hat{\beta}_l + \hat{\beta}_l^\dagger)^2, \hat{\mathcal{H}}_\ell^{l,even} = -\frac{\hbar_{xx}}{8} \hbar/\omega_\ell (\hat{\beta}_l + \hat{\beta}_l^\dagger)^2 \\
\hat{\mathcal{H}}_{int} &= \frac{\hbar \sqrt{\frac{(-1)^{\ell-\frac{1}{2}}8\pi\mathfrak{G}\mathcal{M}\nu^3}{\omega_\ell c^2\mathcal{V}}}\mathcal{L}}{\pi^2 \ell^2 (\hat{\beta}_l + \hat{\beta}_l^\dagger)(\hat{\alpha}e^{-i\mathfrak{v}t} + \hat{\alpha}^\dagger e^{i\mathfrak{v}t})}, \Gamma_{stim} = \frac{2\pi}{\hbar^2 |\langle \alpha | 1 | \hat{\mathcal{H}}_{int} | \alpha \rangle |^2 \mathcal{D}(\omega)}, \Gamma_{stim} \\
&= \frac{\frac{|\alpha|^2}{\ell^4} 8\pi\mathfrak{G}\mathcal{M}\mathcal{L}^2 \omega_\ell^4}{\pi^4 c^5}, \mathcal{N} = \frac{\hbar_0^2 c^5}{32\pi\mathfrak{G}\hbar\nu^2}, \Gamma_{stim} = \frac{\frac{1}{\ell^4} \mathcal{M}\mathcal{L}^2 \omega_\ell^2 \hbar_0^2}{4\pi^5 \hbar}
\end{aligned}$$



$$\begin{aligned}
\hat{\mathcal{H}} &= \hbar\omega\hat{\beta}_l^+\hat{\beta}_l + \frac{1}{\eta^2\mathcal{L}}\sqrt{\frac{\mathcal{M}\hbar}{\omega}}\ddot{\hbar}(t)(\hat{\beta}_l^++\hat{\beta}_l), \hat{\mathfrak{U}}_{int} = \hat{\mathcal{T}}e^{-i\int_0^t ds(g(\delta)\hat{\beta}(\delta)+g^\ominus(\delta)\hat{\beta}^\dagger(\delta))}, g(t) \\
&= \frac{1}{\pi^2\mathcal{L}\sqrt{\frac{\mathcal{M}\hbar}{\omega}}\ddot{\hbar}(t)}, \hat{\mathfrak{U}}_{int} = \epsilon^{\Omega(t)}, \Omega(t) \\
&= \int_0^t dt_1\hat{\Lambda}(t_1) + \frac{1}{2}\int_0^t dt_1\int_0^t dt_2[\hat{\Lambda}(t_1),\hat{\Lambda}(t_2)], \hat{\mathfrak{U}}_{int} \\
&= e^{-i\int_0^t ds(g(\delta)\hat{\beta}(\delta)+g^\ominus(\delta)\hat{\beta}^\dagger(\delta))}e^{-i\varphi}, \hat{\mathfrak{U}} = e^{-i\varphi}e^{-i\omega t\hat{\beta}^\dagger\hat{\beta}}\hat{\mathcal{D}}(\beta), \beta \\
&= -i\int_0^t ds g^\ominus(\delta)e^{i\omega\delta}, e^{-i\varphi}\left|\beta e^{-i\omega t\hat{\beta}^\dagger\hat{\beta}}\right\rangle, |\beta| = \frac{\mathcal{L}}{\pi^2\sqrt{\frac{\mathcal{M}}{\omega\hbar}}}\chi(\hbar, \omega, t), \chi(\hbar, \omega, t) \\
&= \left|\int_0^t ds\ddot{\hbar}(\delta)e^{i\omega\delta}\right|, v(t) = \left(\frac{1}{\frac{8}{v_0^3}} - \frac{8}{3}\kappa t\right)^{-\frac{3}{8}}, \kappa = \frac{\kappa_f}{(2\pi)^{\frac{8}{3}}} = \frac{\frac{96}{5\varpi}\left(\frac{\pi\mathfrak{G}\mathcal{M}_c}{c^3}\right)^{\frac{5}{3}}1}{(2\pi)^{\frac{8}{3}}} \\
&= \frac{48}{5\left(\frac{\mathfrak{G}\mathcal{M}_c}{2c^3}\right)^{\frac{5}{3}}}, \tau = \frac{2\Delta\omega}{\kappa\omega^{\frac{11}{3}}}, \chi = \left|\int_0^t \frac{ds e^{i\omega\delta}\hbar_0 v^2 t}{2} \text{sinc}\left(\frac{\delta t}{2}\right)\right|, 2\Delta\omega = \frac{8}{T}, \tau \\
&= 2\sqrt{\frac{2}{\kappa}}\omega^{\frac{11}{6}}
\end{aligned}$$

$$\chi \approx \hbar_0\omega\left|\int_0^\tau ds e^{i\omega\delta}\sin(\omega\delta)\right| = \frac{\hbar_0\omega}{4}\sqrt{2+4\omega^2\tau^2-2\cos(2\omega\tau)-4\omega\tau\sin(2\omega\tau)}, \chi(\tau) \approx \frac{\hbar_0\omega^2\tau}{2}\chi$$

$$\approx \hbar_0\sqrt{\frac{2}{\kappa}}\omega^{\frac{1}{6}} = \hbar_0\sqrt{\frac{5}{24}\left(\frac{2c^3}{\mathfrak{G}\mathcal{M}_c}\right)^{\frac{5}{6}}}\omega^{\frac{1}{6}}, \mathcal{M} = \frac{\pi^2\hbar\omega^3}{v_\delta^2\chi^2} \approx \frac{\pi^2\hbar\kappa}{2v_\delta^2\hbar_0^2\omega^{\frac{8}{3}}}$$

$$= \frac{\frac{24\pi^2}{5}\hbar}{\hbar_0^2v_\delta^2\left(\frac{2c^3}{\mathfrak{G}\mathcal{M}_c}\right)^{\frac{5}{3}}\omega^{\frac{8}{3}}}, \mathcal{P}(t) = \frac{\frac{\mathcal{L}^2}{\pi^4\mathcal{M}}}{\omega\hbar|\chi(t)|^2} = \frac{\hbar_0^2\omega\tau^2\mathcal{M}v_\delta^2}{4\pi^2\hbar}, \Gamma_{mc} = \frac{d\mathcal{P}(t)}{dt}$$

$$= \frac{\hbar_0^2\omega\tau^2\mathcal{M}v_\delta^2}{2\pi^2\hbar} = \frac{\hbar_0^2\mathcal{N}_c\mathcal{M}v_\delta^2}{\pi\hbar}, \hbar_0 = \sqrt{\frac{\pi\kappa_\beta\mathcal{T}}{\mathcal{M}v_\delta^2\mathcal{Q}\mathcal{N}_c}}$$

$$\begin{aligned}
\mathcal{P}(v, \omega, t) &\approx |\beta(v, \omega, t)|^2 \approx \frac{(\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2) \sin^2 \left[\frac{1}{2} t(v - \omega) \right]}{\hbar(v - \omega)^2 \pi^4}, \mathcal{P}(\omega, t) \\
&= \frac{\sum_v |\beta(v, \omega, t)|^2 \int_{\omega - \frac{\delta}{2}}^{\omega + \frac{\delta}{2}} dv \mathcal{D}(v) |\beta(v, \omega, t)|^2 = \frac{\mathcal{D}(\omega)(\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{\hbar} \pi^4 \int_{\omega - \frac{\delta}{2}}^{\omega + \frac{\delta}{2}} dv \sin^2 \left[\frac{1}{2} t(v - \omega) \right]}{(v - \omega)^2} \\
&= \frac{\mathcal{D}(\omega)(\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{\hbar \pi^4} \Xi(t), \sum_v = \int dv \mathcal{D}(v) = \frac{\int dv \mathcal{V} v^2}{2\pi^2 c^3}, \mathcal{P}(\omega, t) \approx \frac{\mathcal{D}(\omega)(\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{2\hbar \pi^3} \\
&= \Gamma_{stim}^t, \Gamma_{stim} = \frac{\mathcal{D}(\omega)(\hbar_0^2 \omega^3 \mathcal{M} \mathcal{L}^2)}{2\hbar \pi^3} = \frac{\mathcal{V} \hbar_0^2 \omega^5 \mathcal{M} \mathcal{L}^2}{4\hbar \pi^5 c^4} = \hbar_0^2 \frac{\mathcal{M} \mathcal{V}_\delta^2}{4\hbar \pi^3}, \hbar_c \equiv 2\pi \sqrt{\frac{\pi \kappa_\beta \mathcal{T}}{\mathcal{M} \mathcal{V}_\delta^2 Q}}, \hat{\mathcal{H}}_{\omega_l} \\
&= \hbar \omega_l \left(\hat{\beta}_l^\dagger + \hat{\beta}_l + \frac{1}{2} \right), |\psi_{\mathcal{M}}\rangle = (2\pi t_m)^{\frac{1}{4}} \int d\chi e^{-\frac{\chi^2}{4t_m}} |\chi\rangle, \hat{\mathcal{H}}_{int}^{\mathcal{M}} dt = \sqrt{dt} \hat{\rho} \hat{\mathcal{N}}, \hat{\mathcal{M}}_{\hat{\mathcal{N}}}(\gamma) \\
&= \langle \gamma | e^{-i\hat{\mathcal{H}}_{int}^{\mathcal{M}} dt} | \psi_{\mathcal{M}} \rangle = (2\pi t_m)^{-\frac{1}{4}} \exp \left[-\frac{(\gamma - \hat{\mathcal{N}} \sqrt{dt})^2}{4t_m} \right], \hat{\mathcal{M}}_{\hat{\mathcal{N}}}(r) \\
&= \left(\frac{2\pi t_m}{dt} \right)^{-\frac{1}{4}} \exp \left\{ \left[-\frac{dt(r - \hat{\mathcal{N}})^2}{4t_m} \right] \right\}, \rho(t + dt) \\
&= \mathcal{D}[dt\beta'^{(t)}] \hat{\mathcal{M}}_{\hat{\mathcal{N}}}[r(t)] \rho(t) \hat{\mathcal{M}}_{\hat{\mathcal{N}}}^\dagger[r(t)] \mathcal{D}[-dt\beta'^{(t)}] / \text{tr} \{ \hat{\mathcal{M}}_{\hat{\mathcal{N}}}[r(t)] \rho(t) \hat{\mathcal{M}}_{\hat{\mathcal{N}}}^\dagger[r(t)] \}
\end{aligned}$$

6.2. Modelo fotónico aplicable a partículas y antipartículas deformantes del espacio cuántico curvo.



$$\begin{aligned}
[\nabla^2 + \mu(r)\epsilon(r)\kappa^2]\mu_\xi(\kappa, r), \widehat{\mathfrak{E}}(r) &= \iota \sum_{\xi, \kappa} \sqrt{\frac{\hbar c \kappa}{2\epsilon_0 \epsilon(r)}} \mu_\xi(\kappa, r) (\alpha_{\xi\kappa}^\dagger + \alpha_{\xi\kappa}), \frac{\mathcal{H}}{\hbar} \\
&= \omega_0 \sigma_+ \sigma_- + \sum_{\xi, \kappa} c \kappa \alpha_{\xi\kappa}^\dagger \alpha_{\xi\kappa} + \sum_{\xi, \kappa} g(r) (\alpha_{\xi\kappa} \sigma_+ + \alpha_{\xi\kappa}^\dagger \sigma_-), g_\xi(\kappa, r) \\
&= \sqrt{\frac{\hbar c \kappa}{2\epsilon_0 \epsilon(r)}} d \otimes \mu_\xi(\kappa, r) \int_0^\infty \kappa d\kappa \rho(\kappa) \mu_\xi(\kappa, r) \mu_\xi(\kappa, r') e^{-i\kappa \tau} \\
&= \sum_\eta z_{\xi\eta} v_{\xi\eta}(r) v_{\xi\eta}(r') \Theta(\tau - \Delta_t(r, r')) e^{-icz_{\xi\eta}\tau}, \frac{d}{dt \tilde{c}_0(t)} \\
&= - \int_0^t dt' \sum_\xi \int d\kappa \rho(\kappa) g_\xi^2(\kappa) e^{i(\omega_0 - c\kappa)(t-t')} \tilde{c}_0(t'), \frac{d}{dt \tilde{c}_0(t)} \\
&= - \int_0^t dt' \sum_{\xi, \eta} \hat{g}_{\xi\eta}^2 e^{i(\omega_0 - cz_{\xi\eta}\tau)(t-t')} \tilde{c}_0(t'), \frac{id}{dt} \tilde{c}_0(t) \\
&= \sum_{\xi\eta} \hat{g}_{\xi\eta} e^{i(\omega_0 - cz_{\xi\eta}\tau)t} \tilde{\beta}_{\xi\eta}(t), \frac{id}{dt} \tilde{\beta}_{\xi\eta}(t) = \hat{g}_{\xi\eta} e^{i(\omega_0 - cz_{\xi\eta}\tau)t} \tilde{c}_0(t), \hat{g}_{\xi\eta}(r) \\
&= \sqrt{\frac{\hbar c z_{\xi\eta} \tau}{2\epsilon_0 \epsilon(r)}} d \otimes v_\xi(r), \mathfrak{E}_+ \left| \psi(t) \right\rangle = i \sum_{\xi, \eta} \sqrt{\frac{\hbar c z_{\xi\eta} \tau}{2\epsilon_0 \epsilon(r)}} \hat{v}_{\xi\eta}(r) \tilde{\beta}_{\xi\eta}(t - \Delta_t) e^{-icz_{\xi\eta}\tau}, Z_l \\
&= (\sqrt{\epsilon \kappa r}) = \frac{1}{\sqrt{J_m(\kappa)} \left\{ \begin{array}{l} \eta_l(\kappa) j_l(\sqrt{\epsilon \kappa r}) \\ \alpha_l(\kappa) j_l(\kappa r) + \beta_l(\kappa) \gamma_l(\kappa r) \end{array} \right\}^{r < \alpha, r > \alpha}}, \lim_{\mathcal{R} \rightarrow \infty} J_{\mathcal{M}}(\kappa) \\
&= \frac{\mathcal{R}}{2\kappa^2 [\alpha_l(\kappa) + i\beta_l(\kappa)][\alpha_l(\kappa) - i\beta_l(\kappa)]}, J_l(r, r') \\
&= \int_0^\infty d\kappa \rho(\kappa) \kappa^{-1} Z_l(\sqrt{\epsilon \kappa r}) Z_l(\sqrt{\epsilon \kappa r'}) e^{-i\kappa \tau}, J_l \\
&= \Theta[c\tau - (r - \alpha)] \oint_{\mathcal{LHP}} f_1(z) dz \\
&+ \Theta[c\tau \\
&+ (r - \alpha)] \oint_{\mathcal{UHP}} f_2(z) dz, J_l(r, r') 2\pi i \sum_{Z_{l\eta} \in \mathbb{Q}_4} \text{Res}[f_2(z)] \Theta[c\tau - (r - \alpha)], v_{lm\eta}(r)
\end{aligned}$$

$$\begin{aligned}
&= \Re_{lm\eta}(r) \left[\pi \sqrt{\alpha_l(Z_{l_\eta}) - \frac{i\beta_l(Z_{l_\eta})}{i[\partial_Z \alpha_l(Z) + i\beta_l(Z)]_{Z_{l_\eta}}} \hbar_l^{(1)}(Z_{l_\eta} r)} \right], \delta\omega_0 \\
&= \sum_{off-res} g_{l_\eta}^2 (\omega_0 - \omega_{l_\eta}) - \frac{i\gamma_{l_\eta}}{(\omega_0 - \omega_{l_\eta})^2} + \gamma_{l_\eta}^2, id/dt \tilde{c}_0(t) \\
&= \delta\omega_0 \tilde{c}_0(t) + \sum_{\substack{(5,4) \\ (8,3)}} \hat{g}_{l_\eta} e^{i(\omega_0 - cZ_{l_\eta})t} \tilde{\beta}_{l_\eta}(t)
\end{aligned}$$

6.3. La gravedad como entidad cuántica (Formalización).

$$\begin{aligned}
|\psi(t=0)\rangle &= |\zeta\rangle_{\mathcal{M}_c} \otimes \frac{1}{\sqrt{2}(|\uparrow\rangle_{\delta_c} |\downarrow\rangle_{\delta_c})}, |\zeta\rangle_{\mathcal{M}_c} \otimes |\uparrow\rangle_{\delta_c} \rightarrow |\mathcal{L}\uparrow\rangle_c, |\zeta\rangle_{\mathcal{M}_c} \otimes |\downarrow\rangle_{\delta_c} \rightarrow |\mathcal{R}\downarrow\rangle_c, |\psi\rangle_{c,\Lambda} \\
&= \frac{1}{\sqrt{2}(\sqrt{1+\cos\Delta\phi}|\Psi_+\rangle_c|+\rangle_{\delta_\Lambda} + \sqrt{1-\cos\Delta\phi}|\Psi_-\rangle_c|-\rangle_{\delta_\Lambda})} |\zeta\rangle_{\mathcal{M}_\Lambda} |\Psi_\pm\rangle_c \\
&= (1 \pm e^{i\Delta\phi})|\mathcal{L}\uparrow\rangle_c + \frac{(e^{i\Delta\phi} \pm 1)|\mathcal{R}\downarrow\rangle_c}{2\sqrt{1 \pm \cos\Delta\phi} |\pm\rangle_{\delta_\Lambda}} = |\uparrow\rangle_{\delta_\Lambda} \pm \frac{|\downarrow\rangle_{\delta_\Lambda}}{\sqrt{2}}, \Delta\phi_\tau \\
&= \frac{\mathcal{GM}m\tau}{\hbar\sqrt{d^2 + (\Delta\chi)^2}} \\
&\quad - \frac{\mathcal{GM}m\tau}{\hbar d}, |\psi_{\alpha,\beta,c}\rangle \\
&= \frac{1}{8[(1+\alpha e^{i\Delta\phi})(1-\beta e^{i\Delta\phi}) + c e^{2i\Delta\phi}(1+\alpha e^{i\Delta\phi})(1-\beta e^{i\Delta\phi})]} |\zeta\rangle_{\mathcal{M}_c} |c\rangle_{\delta_c}, \mathcal{V}(\pm) \\
&= \mathcal{P}_\pm - \sum_{\alpha,\beta \in \{\pm\}} \mathcal{P}_{\alpha,\beta,\pm} = \pm \frac{1}{2} \sin^2 \Delta\phi, \rho_{c,\Lambda} \\
&= \frac{1}{2} \left((1+\cos\Delta\phi) |\Psi_+\rangle_c \langle\Psi_+|_c \otimes |+\rangle_{\delta_\Lambda} \langle+|_{\delta_\Lambda} \right. \\
&\quad \left. + (1-\cos\Delta\phi) |\Psi_-\rangle_c \langle\Psi_-|_c \otimes |-\rangle_{\delta_\Lambda} \langle-|_{\delta_\Lambda} \right) \otimes |\xi\rangle_{\mathcal{M}_\Lambda} \langle\xi|_{\mathcal{M}_\Lambda}
\end{aligned}$$

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Apéndice D:

Postulados Finales

1. Que las partículas con o sin masa o las antipartículas con o sin masa, según sea el caso, en tanto y en cuanto, se aproximen, igualen o superen la velocidad de la luz, deforman el espacio cuántico en el que interactúan, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.
2. Que las partículas masivas o supermasivas o las antipartículas masivas o supermasivas, según sea el caso, no necesitan aproximarse, igualar o superar la velocidad de la luz, para deformar el espacio cuántico en el que interactúan, de tal suerte que, basta con su hipermasa o supermasa, según sea el caso, para lograr un espacio cuántico curvo, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.
3. Cuando una partícula supermasiva o una antipartícula supermasiva, según sea el caso, alcanzan o superan la velocidad de la luz, producen un agujero negro cuántico, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.



4. Cuando una partícula sin masa o una antipartícula sin masa, según sea el caso, superan la velocidad de la luz, producen un agujero negro cuántico, a propósito de sus ciclos cuánticos de colisión, superposición o entrelazamiento, según corresponda a cada caso.

APÉNDICE E.

FORMALIZACIÓN MATEMÁTICA COMPLEMENTARIA.

1. Espacios cuánticos curvos.

1.1. Osciladores y propagadores en espacios curvos – Modelo Feynman.

$$\begin{aligned}
 \mathfrak{F} &= \langle \int_{-\infty}^{\infty} q(\sigma)^2 \bar{\chi} e^{-\sigma^2} d\sigma, \mathfrak{F}[x(t,s), y(t,s)] = \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} x(t,s), y(t,s) \sin \omega(t-s) dt ds, \mathcal{F}(\dots q_i \dots) \\
 &= \sum_{i=-\infty}^{\infty} q_i^2 e^{-\sigma_i^2} (\sigma_{i+1} - \sigma_i) \sum_i \frac{(\dots q_i \dots)}{\partial q_i} \lambda_i, \mathfrak{F}[q(\sigma), +\lambda(\sigma)] - \mathcal{F}[q(\sigma)] \\
 &= \int \mathcal{K}(t) \lambda(t) dt, \mathfrak{F}[q(\sigma), +\lambda(\sigma)] = \mathcal{F}[q(\sigma)] + \int \frac{\delta \mathcal{F}[q(\sigma)]}{\delta q(t) \lambda(t) dt}, \mathcal{F}[q + \lambda] \\
 &= \int [q(\sigma)^2 + 2q(\sigma)\lambda(\sigma) + \lambda(\sigma)^2] e^{-\sigma^2} d\sigma \\
 &= \int q(\sigma)^2 e^{-\sigma^2} d\sigma + 2 \int q(\sigma)\lambda(\sigma) e^{-\sigma^2} d\sigma
 \end{aligned}$$

$$\begin{aligned}
\mathcal{A} &= \left\langle \int \mathcal{L}(\dot{q}(\sigma), q(\sigma)) d\sigma, \frac{\delta \mathcal{A}}{\delta q(t)} = \frac{d}{dt} \{ \partial \mathcal{L}(\dot{q}(t), q(t)) / \partial \dot{q} \} + \partial \mathcal{L}(\dot{q}(t), q(t)) / \partial q, \mathcal{A} \right. \\
&= \int_{-\infty}^{\infty} \left\{ \frac{m(\dot{x}(t))^2}{2} - \mathcal{V}(x(t)) + \kappa^2 \dot{x}(t) \dot{x}(t + \mathcal{T}_0) \right\} dt, \delta \mathcal{A} \\
&= \int_{-\infty}^{\infty} \{ m \dot{x}(t) \dot{\lambda}(t) - \mathcal{V}'(x(t)) \lambda(t) + \kappa^2 \dot{\lambda}(t) \dot{x}(t + \mathcal{T}_0) + \kappa^2 \dot{\lambda}(t + \mathcal{T}_0) \dot{x}(t) \} dt \\
&= \int_{-\infty}^{\infty} \{ -m \ddot{x}(t) - \mathcal{V}'(x(t)) - \kappa^2 \ddot{x}(t + \mathcal{T}_0) + \kappa^2 \ddot{x}(t - \mathcal{T}_0) \} \lambda(t) dt, \frac{\delta \mathcal{A}}{\delta x(t)} \\
&= -m \ddot{x}(t) - \mathcal{V}'(x(t)) - \kappa^2 \ddot{x}(t + \mathcal{T}_0) - \kappa^2 \ddot{x}(t - \mathcal{T}_0), \frac{\delta \mathcal{A}}{\delta \gamma(t)} \\
&= -\frac{d}{dt} \left(\frac{\partial \mathfrak{L}_\gamma}{\partial \dot{\gamma}} \right) + \frac{\partial \mathfrak{L}_\gamma}{\partial \gamma} \Big|_t \\
&+ \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_t \otimes x(t), \frac{\delta}{\delta \gamma(s) \left[-\frac{d}{dt} \left(\frac{\partial \mathfrak{L}_\gamma}{\partial \dot{\gamma}} \right) + \frac{\partial \mathfrak{L}_\gamma}{\partial \gamma} \Big|_t + \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_t \otimes x(t) \right]} \\
&= \frac{\delta}{\delta \gamma(t) \left[-\frac{d}{ds} \left(\frac{\partial \mathfrak{L}_\gamma}{\partial \dot{\gamma}} \right) + \frac{\partial \mathfrak{L}_\gamma}{\partial \gamma} \Big|_s + \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \otimes x(s) \right] \frac{\partial \mathfrak{L}_\gamma}{\partial \gamma} \Big|_t \delta x(t)} = \frac{\frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(t)} \\
&= \frac{\frac{1}{m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s) \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(t)}{\delta \gamma(s)} = \frac{\frac{1}{m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} \\
&= -\frac{\sin \omega(\mathcal{T}-t) \sin \omega s}{m\omega} \sin \omega \mathcal{T} \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \\
&= -\frac{\sin \omega(\mathcal{T}-s) \sin \omega t}{m\omega} \sin \omega \mathcal{T} \frac{\partial \mathcal{J}_\gamma}{\partial \chi} \Big|_s \int_0^{\mathcal{T}} [\mathfrak{L}_y + \mathfrak{L}_z] dt \\
&+ \int_0^{\mathcal{T}} \left[\sin \omega(\mathcal{T}-t) x(0) + \sin \frac{\omega t \chi(\mathcal{T})}{\sin \omega \mathcal{T}} \right] \gamma(t) dt - \frac{1}{m\omega \sin \omega \mathcal{T} \int_0^{\mathcal{T}} dt \int_0^t ds} \\
&\otimes \sin \omega(\mathcal{T}-t) \sin \omega s \gamma(s) \gamma(t), \frac{\delta x(t)}{\delta \gamma(s)} = \frac{\frac{1}{2m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} = \\
&= -\frac{\frac{1}{2m\omega} \sin \omega(t-s) \boxtimes \frac{\partial \mathcal{J}_\gamma}{\partial \gamma} \Big|_s \delta x(s)}{\delta \gamma(s)} \rangle
\end{aligned}$$

$$\mathfrak{E}(t)$$

$$= \left\| \frac{m(\dot{x}(t))^2}{2} + \mathcal{V}(x(t)) - \kappa^2 \int_t^{t+\mathcal{T}_0} \ddot{x}(\sigma - \mathcal{T}_0) \dot{x}(\sigma) d\sigma + \kappa^2 \dot{x}(t) \dot{x}(t + \mathcal{T}_0), \mathcal{A}[q_\eta(\sigma) + \alpha y_\eta(\sigma)] \right\|$$

$$= \mathcal{A}[q_\eta(\sigma)]$$

$$\alpha \sum_{\eta=1}^N \int_{-\infty}^{\infty} \frac{y_\eta(t) \delta \mathcal{A}}{\delta q_\eta(t)} dt \sum_{\eta=1}^N \int_{-\infty}^{\infty} \frac{y_\eta(\sigma) \delta \mathcal{A}}{\delta q_\eta(\sigma)} d\sigma, \mathcal{I}(\mathcal{T}) = \sum_{\eta=1}^N \int_{-\infty}^{\mathcal{T}} \frac{y_\eta(\sigma) \delta \mathcal{A}}{\delta q_\eta(\sigma)} d\sigma - \sum_{\eta=1}^N \int_{\mathcal{T}}^{\infty} \frac{y_\eta(\sigma) \delta \mathcal{A}}{\delta q_\eta(\sigma)} d\sigma,$$

$$+ \frac{\delta \mathcal{I}(\mathcal{T})}{\delta q_m(t)}$$

$$= + \int_{-\infty}^{\mathcal{T}} \sum_m \frac{\frac{\delta y_\eta(\sigma)}{\delta q_m(t)} \delta \mathcal{A}}{\delta q_m(\sigma) d\sigma}$$

$$+ \int_{-\infty}^{\mathcal{T}} \sum_m y_\eta(\sigma) \frac{\delta^2 \mathcal{A}}{\delta q_m(t) \delta q_m(\sigma)} d\sigma \int \frac{\left[\mathfrak{L}_y + \mathfrak{L}_z + \left(\frac{m \dot{x}^2}{2} - \frac{m \omega^2 x^2}{2} \right) + (\mathfrak{I}_y + \mathfrak{I}_z) x \right] d}{dt} \left(\frac{\partial \mathfrak{L}_y}{\partial \dot{y}} \right) - \frac{\partial \mathfrak{L}_y}{\partial y}$$

$$= \frac{\partial \mathfrak{I}_y}{\partial \dot{y}} \chi(t), m \ddot{x} + m \omega^2 x = [\mathfrak{I}_y(t) + \mathfrak{I}_z(t)] \left\| \right.$$

$$\begin{aligned}
x(t) &= \langle x(0) \cos \omega t + \frac{\dot{x}(0) \sin \omega t}{\omega} + \frac{1}{m\omega \int_0^t \gamma(\delta) \sin \omega(t-\delta) d\delta}, x(t) \\
&= \frac{\sin \omega(\mathcal{T}-t)}{\sin \omega \mathcal{T}} \left[x(0) - \frac{1}{m\omega} \int_0^t \sin \omega \delta \gamma(\delta) d\delta \right] \\
&\quad + \frac{\sin \omega t}{\sin \omega \mathcal{T} \left[x(\mathcal{T}) - \frac{1}{m\omega} \int_t^{\mathcal{T}} \sin(\mathcal{T}-\delta) \gamma(\delta) d\delta \right]} \rangle \\
x(t) &= \langle \frac{1}{\sin \omega \mathcal{T}} [\mathcal{R}_t \sin \omega t + \mathcal{R}_0 \sin \omega t(\mathcal{T}-t)] + \frac{1}{2m\omega} \int_0^t \sin \omega(t-\delta) \gamma(\delta) d\delta \\
&\quad - \frac{1}{2m\omega} \int_t^{\mathcal{T}} \sin \omega(t-\delta) \gamma(\delta) d\delta \\
\mathcal{R}_0 &= \langle \frac{1}{2} \left[x(0) + \frac{x(\mathcal{T}) \cos \omega \mathcal{T} - \dot{x}(\mathcal{T}) \sin \omega \mathcal{T}}{\omega} \right], \mathcal{R}_t = \frac{1}{2} \left[x(\mathcal{T}) + \frac{x(0) \cos \omega \mathcal{T} - \dot{x}(0) \sin \omega \mathcal{T}}{\omega} \right] \rangle \\
\mathcal{A} &= \langle \int_0^{\mathcal{T}} [\mathfrak{L}_y + \mathfrak{L}_z] dt + \frac{1}{\sin \omega \mathcal{T} \int_0^{\mathcal{T}} [\mathcal{R}_{\mathcal{T}} \sin \omega t + \mathcal{R}_0 \sin \omega(\mathcal{T}-t)] \gamma(t) dt} \\
&\quad - \frac{1}{2m\omega \int_0^{\mathcal{T}} \int_0^t \sin \omega(t-s) \sin \gamma(t) \gamma(s)} ds dt, \mathcal{A} \\
&= \int_{-\infty}^{\infty} [\mathfrak{L}_y + \mathfrak{L}_z] dt + \frac{1}{2m\omega} \sum_{j=1}^N \frac{1}{2m_j \omega_j} \int_{-\infty}^{\infty} \int_{-\infty}^t \sin \omega(t-s) \boxtimes \sin \gamma(t) \gamma(s) ds dt, \mathcal{A} \\
&= \int_{-\infty}^{\infty} [\mathfrak{L}_y + \mathfrak{L}_z] dt + \sum_{j=1}^N \frac{1}{2m_j \omega_j} \int_{-\infty}^{\infty} \int_{-\infty}^t \sin \omega_j(t-s) \boxtimes \sin \gamma_j(t) \gamma_j(s) ds dt, \mathcal{A} \\
&= \int_{-\infty}^{\infty} [\mathfrak{L}_y + \mathfrak{L}_z] dt + 1/m\omega \int_{-\infty}^{\infty} \int_{-\infty}^t \sin \omega(t-s) \otimes [\mathfrak{I}_y(s) + \mathfrak{I}_z(t) \\
&\quad + \mathfrak{I}_y(t) + \mathfrak{I}_z(s)] ds dt \int_{-\infty}^{\infty} \left[\mathfrak{L}_y + \mathfrak{L}_z + (\mathfrak{I}_y + \mathfrak{I}_z)x_1 + (\mathfrak{I}_y - \mathfrak{I}_z)x_2 \right. \\
&\quad \left. + \frac{m}{z} (\ddot{x}_1^2 - \omega^2 x_1^2) \sum_{\kappa} \sum_{l \neq \kappa} -\frac{m}{2} (\ddot{x}_s^2 + \omega^2 x_s^2) \right] dt \int_{-\infty}^{\infty} \sum_{\kappa} [\mathfrak{L}_y + \mathfrak{L}_z + \mathfrak{I}_y \eta_y + \mathfrak{I}_z \eta_z \\
&\quad + \frac{m}{2} (\hat{\eta}_{yl} \hat{\eta}_{zl} - \omega^2 \hat{\eta}_{y\kappa} \hat{\eta}_{z\kappa})] dt \rangle
\end{aligned}$$

2. Partícula cosmológica (Hipermasa y Supermasa – Partículas masivas y supermasivas y antipartículas masivas y supermasivas).

2.1. Supermasa (Partículas Masivas y Antipartículas Masivas).

$$\begin{aligned}
 \mathfrak{G}_{\mu\nu} - m^2(\hbar_{\mu\nu} - \eta_{\mu\nu}\hbar) &= \mathfrak{G}\mathcal{T}_{\mu\nu}, \mathfrak{G}_{\mu\nu} \\
 &= \square (\hbar_{\mu\nu} - \eta_{\mu\nu}\hbar) - \partial^\alpha \partial_\mu \hbar_{\alpha\mu} - \partial^\alpha \partial_\nu \hbar_{\alpha\nu} + \eta_{\mu\nu} \partial^\alpha \partial^\beta \hbar_{\alpha\beta} \\
 &\quad + \partial_\mu \partial_\nu \hbar \left(\nabla^2 - \frac{1}{c^2 \partial_t^2} - m^2 \right) (\hbar_{\mu\nu} - \eta_{\mu\nu}\hbar) \\
 \Delta &= - \frac{2\mathfrak{G}m_{\mathcal{J}}}{c^4 \ln(|\chi_{\odot\mathcal{J}}| - \chi_{\odot\mathcal{J}} \otimes \kappa) \rightarrow \frac{2\mathfrak{G}m_{\mathcal{J}}}{c^4} [\ln(|\chi_{\odot\mathcal{J}}| - \chi_{\odot\mathcal{J}} \otimes \mathfrak{K})(1 - \mathfrak{K} \otimes v_{\mathcal{J}})] \otimes \mathfrak{K} \\
 &\equiv \kappa - \left[\kappa \bigotimes (v_{\mathcal{J}} \otimes \kappa) \right] / m^4 c^4 \\
 \alpha_\eta &= - \frac{\mathfrak{G}\mathfrak{M}}{r^2} \rightarrow - \frac{\sqrt{\mathfrak{G}\mathfrak{M}\alpha_0}}{r}
 \end{aligned}$$

2.2. Hipermasa (Partículas Supermasivas y Antipartículas Supermasivas).

$$\begin{aligned}
 \lambda_g &= \frac{\hbar}{m_g c}, \mathcal{A} \sim \frac{1}{2\mathfrak{M}_{\mathfrak{P}_g}^2 \int d^4\chi \int d^4\gamma \mathcal{T}_1^{\mu\nu}(\chi) \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) \mathcal{T}_2^{\mu\nu}(\gamma)}, \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) \\
 &= i\langle \hat{\mathcal{T}}[\hbar_{\mu\nu}(\chi)\hbar_{\alpha\beta}(\gamma)] \rangle, \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) = \frac{f_{\mu\nu\alpha\beta}}{-\square} - i\varepsilon\delta^4(x - y), \mathfrak{G}_{\mu\nu\alpha\beta}(x, y) \\
 &= \frac{f_{\mu\nu\alpha\beta}}{-\square} - i\varepsilon, f_{\mu\nu\alpha\beta} = \tilde{\eta}_{\mu(\frac{\alpha\mu_{\gamma}}{\beta\nu}\tilde{\eta}_{|\nu|\alpha})} - \frac{1}{2}\tilde{\eta}_{\mu\nu}\tilde{\eta}_{\alpha\beta}, \tilde{\eta}_{\mu\nu} = \eta_{\mu\nu} - 1/\square \partial_\mu \partial_\nu \\
 \mathcal{I}m[\mathfrak{U}] &\sim \pi/2\mathfrak{M}_{\mathfrak{P}_g}^2 \int d^4\chi \mathcal{T}_1^{\mu\nu}(\chi) f_{\mu\nu\alpha\beta} \delta(\square) \mathcal{T}_2^{\alpha\beta}(\chi) \sim \pi/2\mathfrak{M}_{\mathfrak{P}_g}^2 \int d^4\chi \mathcal{T}_1^{\mathcal{T}\mathfrak{Z}\mu\nu}(\chi) f_{\mu\nu\alpha\beta} \delta(\square) \mathcal{T}_2^{\mathcal{T}\mathfrak{Z}\alpha\beta}(\chi) \\
 \mathcal{F}_{12} &\sim \frac{1}{\mathfrak{Z}} \frac{d}{dr} \mathcal{R}e[\mathcal{A}] \sim \frac{\mathfrak{M}_1 \mathfrak{M}_2}{\mathfrak{M}_{\mathfrak{P}_g}^2 r^2}, \mathfrak{G}_{\mu\nu\alpha\beta}^{(m)} = \frac{\sum_{\mathcal{J}} \mathfrak{f}_{\mathcal{J}\mu\nu\alpha\beta}^{(m)}}{\partial_t^2} - \mathfrak{F}_{\mathcal{J}}[-\nabla^2] + m_g^4 - i\epsilon, \mathfrak{G}_{\mu\nu\alpha\beta}^{(m)} \\
 &= \frac{\sum_{\mathcal{J}} \int_0^\infty \mathfrak{f}_{\mathcal{J}\mu\nu\alpha\beta}^{(m)}(\mu) \rho_{\mathcal{J}}(\mu) d\mu}{\partial_t^2 - \mathfrak{F}_{\mathcal{J}}[-\nabla^2] + \mu^2 + i\epsilon}
 \end{aligned}$$

$$\begin{aligned}\mathcal{L}_{\mathfrak{F}\mathcal{P}} &= \frac{\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2}{4} \hbar^{\mu\nu} \xi^{\alpha\beta}_{\mu\nu} \hbar_{\alpha\beta} - \frac{1}{8} m_{\mathscr{G}}^4 \mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2 (\hbar_{\mu\nu}^2 - \hbar^2) + \frac{1}{2} \hbar_{\mu\nu} \mathcal{T}^{\mu\nu}, \mathfrak{G}_{\mu\nu\alpha\beta}^{(\mathfrak{m})} \\ &= \frac{\mathfrak{f}_{\mu\nu\alpha\beta}^{(\mathfrak{F}\mathcal{P})}(m_{\mathscr{G}})}{-\square} + m_{\mathscr{G}}^4 - \iota\epsilon, \mathfrak{f}_{\mu\nu\alpha\beta}^{(\mathfrak{F}\mathcal{P})}(m_{\mathscr{G}}) = \tilde{\eta}_{\mu\left(\beta\nu\tilde{\eta}_{|\nu|\beta}^{|\mu|\alpha}\right)} - \frac{1}{3\tilde{\eta}_{\mu\nu}\tilde{\eta}_{\alpha\beta}}, \tilde{\eta}_{\mu\nu} \\ &= \eta_{\mu\nu} - \frac{1}{m_{\mathscr{G}}^4\partial_\mu\partial_\nu}, \mathfrak{G}_{\mu\nu\alpha\beta}^{(\mathfrak{m})} = \int_0^\infty \frac{d\mu\rho(\mu)\mathfrak{f}_{\mu\nu\alpha\beta}^{(\mathfrak{F}\mathcal{P})}(\mu)}{-\square} + \mu^2 + \iota\epsilon, \hbar_{\mu\nu} \\ &\rightarrow \hbar_{\mu\nu} + \frac{\alpha\mu}{\beta\nu}\widetilde{\partial\Lambda}_{|\nu|\beta}^{|\mu|\alpha} + \partial_\mu\partial_\nu\varpi\end{aligned}$$

$$\mathcal{L}_{\mathfrak{F}\mathcal{P}}^{m_{\mathscr{G}}\rightarrow 0} = \frac{1}{4}\widehat{h}^{\mu\nu}\xi^{\alpha\beta}_{\mu\nu}\widehat{h}_{\alpha\beta} - \frac{\alpha\mu}{\beta\nu}\widetilde{\partial\Lambda}_{|\nu|\beta}^{|\mu|\alpha} - \frac{1}{2}\partial_\mu\widehat{\pi}\partial^\mu\widehat{\pi} + \frac{1}{2}\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\widehat{h}_{\mu\nu}\mathcal{T}^{\mu\nu} + 1/2\sqrt{6\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2}\widehat{\pi}\mathcal{T}^\mu_\nu$$

$$r_{\nu,\odot}=(\mathcal{M}_{\odot}/\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2m_{\mathscr{G}}^4)^{1/3}\sim (r_{\delta,\odot}\lambda_{\mathscr{G}}^2)^{1/3}$$

$$\begin{aligned}\mathcal{L}_{\widehat{\pi}} &= -\frac{1}{2}(\partial\widehat{\pi})^2 + \Lambda^4\mathfrak{G}\left(\frac{\partial\widehat{\pi}}{\Lambda^2},\frac{\partial^2\widehat{\pi}}{\Lambda^3}\right) + \frac{1}{\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\widehat{\pi}\mathfrak{I}}, \mathcal{L}_{\delta\pi} = -\frac{1}{2}\mathfrak{Z}^{\mu\nu}\partial_\mu\delta\pi\partial_\nu\delta\pi + \frac{1}{\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\delta\widetilde{\pi}\delta\mathfrak{I}}, \mathcal{L}_\chi \\ &= -\frac{1}{2}(\partial\widehat{\pi})^2 + 1/\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\sqrt{3}\chi^{\delta\mathfrak{T}}\end{aligned}$$

$$\Phi\!\sim\!\frac{\mathfrak{M}}{\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2r}e^{-m_{\mathscr{G}}r},\mathfrak{E}^2-\rho^2=m_{\mathscr{G}}^4,v_{\mathscr{G}}^2(\varepsilon)=-\frac{1}{m_{\mathscr{G}}^4}\mathfrak{E}^2,\mu(r)$$

$$\begin{aligned}&= \frac{\mathcal{M}_{\odot}}{8\pi\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\left[1-\left(rm_{\mathscr{G}}\right)^2+\mathcal{O}\left(\left(rm_{\mathscr{G}}\right)^4\right)\right]},\alpha^4=\frac{\mathcal{T}^4\mu(\alpha)}{(2\pi)^2}1+\eta\equiv\frac{\alpha}{\alpha_{\otimes}\left(\frac{\mathcal{T}_{\otimes}}{\mathcal{T}_{\alpha}}\right)^{\frac{2}{4}}}\\ &=\left(\frac{\mu(\alpha)}{\mu_{\otimes}}\right)^{\frac{1}{4}}m_{\mathscr{G}}\geq\sqrt{\langle\frac{12\eta}{(1\,\mathcal{A}U)^2}-\alpha^2\rangle+\mathcal{O}(\eta)}\end{aligned}$$

$$\begin{aligned}\delta_{\mathfrak{D}\mathfrak{GP}} &= \frac{\frac{\mathfrak{M}_6^4}{2\int d^5\chi\sqrt{-\mathscr{G}_5}\mathcal{R}_5}}{2} - \mathfrak{M}_6^4\int d^4\chi\sqrt{-g}\kappa + \frac{\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2}{2}\int d^4\chi\sqrt{-g}\left(\frac{\mathfrak{R}}{2} + \mathcal{L}_{\mathcal{M}}\right), m_{\mathscr{G}} = \mathfrak{W}_{cross} \\ &= \frac{\mathfrak{M}_6^4}{\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2}, \Phi(r) = \frac{\frac{1}{8}\pi^2\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2}{r}1\Big\{\sin(rm_{\mathscr{G}})\mathfrak{C}\mathfrak{i}(rm_{\mathscr{G}}) + \frac{1}{2}\cos(rm_{\mathscr{G}})[\pi - 2\delta\mathfrak{i}(rm_{\mathscr{G}})]\Big\}\end{aligned}$$

$$m_{\mathscr{G}}\rightarrow\!\!\!\rightarrow\!0,\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\rightarrow\!\!\!\rightarrow\!\!\!\infty,\mathcal{T}^{\mu\nu}\rightarrow\!\!\!\rightarrow\!\!\!\infty,\Lambda_4=\left(m_{\mathscr{G}}^4\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\right)^2\rightarrow\!\!\!\rightarrow\!\!\!\gamma,\mathcal{T}^{\mu\nu}/\mathfrak{M}_{\mathfrak{P}_\mathfrak{E}}^2\rightarrow\!\!\!\rightarrow\!\!\!\beth$$

$$\mathcal{L}_{\mathfrak{D}\mathfrak{G}\mathcal{P}}^{\text{bl}}=\frac{1}{4}\widehat{h}^{\mu\nu}\widehat{\xi}^{\alpha\beta}_{\mu\nu}\widehat{h}_{\alpha\beta}+\frac{1}{2}\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2\widehat{h}_{\mu\nu}\mathcal{T}^{\mu\nu}+\mathcal{L}_{\mathfrak{D}\mathfrak{G}\mathcal{P}}^{\pi}+\frac{\widehat{\pi}}{2\sqrt{6\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2\widehat{\pi}\mathcal{T}_{\mathfrak{v}}^{\mu}}},\mathcal{L}_{\mathfrak{D}\mathfrak{G}\mathcal{P}}^{\pi}$$

$$=-\frac{1}{2}(\partial\widehat{\pi})^2-\frac{1}{(\sqrt{6}\Lambda_4)^4(\partial\widehat{\pi})^2}-\Box\pi$$

$$\delta_{\mathfrak{d}\mathfrak{H}\mathfrak{G}\mathcal{T}}=\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2\int d^4\chi\sqrt{-g}\left[\frac{\mathcal{R}}{2}+m_{\mathscr{G}}^4\sum_{j=2}^4m_{\mathscr{G}}^4\alpha_j\mathfrak{U}_j(\mathcal{K})\right],\mathcal{K}_{\mathfrak{v}}^{\mu}=\delta_{\mathfrak{v}}^{\mu}-\mathcal{X}_{\mathfrak{v}}^{\mu}=\left(\sqrt{\mathscr{G}^{-1}\eta}\right)_{\mathfrak{v}}^{\mu},\mathcal{X}_{\mathfrak{v}}^{\mu}\\ =\left(\sqrt{\mathscr{G}^{-1}\eta}\right)_{\mathfrak{v}}^{\mu}\rightarrow\mathcal{X}_{\mathfrak{v}}^{\mu}=\left(\sqrt{\mathscr{G}^{-1}\widehat{\eta}}\right)_{\mathfrak{v}}^{\mu},\phi^{\alpha}=\chi^{\alpha}+\Lambda^{\alpha}+\partial^{\alpha}\pi$$

$$\mathcal{L}_{\mathfrak{D}\mathfrak{H}\mathfrak{G}\mathcal{T}}^{\text{bl}}=\frac{1}{4}\widehat{h}^{\mu\nu}\widehat{\xi}^{\alpha\beta}_{\mu\nu}\widehat{h}_{\alpha\beta}+\frac{1}{2}\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2\widehat{h}_{\mu\nu}\mathcal{T}^{\mu\nu}+\frac{\alpha_1}{\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2}\pi\mathcal{T}_{\mathfrak{v}}^{\mu}+\frac{\alpha_2}{\Lambda_4^4\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2}\partial_{\mu}\widehat{\pi}\partial_{\mu}\widehat{\pi}\mathcal{T}^{\mu\nu}+\frac{\alpha_3}{\Lambda_6^4\mathfrak{M}_{\mathfrak{P}_{\mathfrak{g}}}^2}\widehat{h}_{\mu\nu}\chi_{(4)}^{\mu\nu}\\ -\frac{1}{2}(\partial\widehat{\pi})^2+\sum_{j=2}^4\beta_j/\Lambda_4^{4(j-2)}\mathcal{L}_j^{\mathfrak{G}\text{al}}(\widehat{\pi})$$

$$\widetilde{m}_{\mathscr{G}}^4(\mathcal{H})=\frac{m_{\mathscr{G}}^4\mathcal{H}}{\mathcal{H}_0}\left[c_0+\frac{c_2\mathcal{H}}{\mathcal{H}_0}+\frac{c_4\mathcal{H}}{\mathcal{H}_0^4}\right],\mathfrak{E}^4=\kappa^4+m_{\mathscr{G}}^4-c_{\mathscr{G}}^4(\mathfrak{G})=1-\frac{m_{\mathscr{G}}^4}{\mathfrak{E}^4},\Phi_{\mathfrak{M}\mathfrak{G}}(f)\\ =-\mathfrak{D}/\|4\pi\lambda_g^2(1+z)f\|$$

$$\mathcal{D}_q''(\tau)+\frac{2\alpha'}{\alpha}\mathcal{D}_q'(\tau)+(q^2+m_{\mathscr{G}}^4\alpha^2)\mathcal{D}_q(\tau)=\mathfrak{I}_q(\tau)(\Box-m_{\mathscr{G}}^4)\Big(\widehat{h}_{\mu\nu}-\frac{1}{2}h\bar{\eta}_{\mu\nu}\Big)\\ =32\pi\mathfrak{G}\mathcal{T}_{\mu\nu}\sqrt{\Re\big|m_{\mathscr{G}}^4\big|_{\mathfrak{v}}}\,\hbar_{\mu\nu}-\hbar_{\mu\nu}^{\mathfrak{G}\mathfrak{R}}\\ +\pi\eta_{\mu\nu},\partial_r\pi\sim\frac{r_{\delta,\otimes}\mathfrak{M}_{\mathcal{P}_{\mathfrak{g}}}}{r_{\mathfrak{V},\odot}^{\frac{4}{2}}r^{\frac{1}{4}}}\otimes\delta_{\phi}\pi\alpha\partial_r\left[r^2\partial_r\left(\frac{r^{-1}\delta\Phi}{\Phi^{\mathfrak{G}\mathfrak{R}}}\right)\right]_{r\rightarrow\alpha}\delta_{\phi}\sim4\pi/2\left(\alpha/r_{\mathfrak{V},\odot}\right)^{4/2}m_{\mathscr{G}}^4\\ \geq4/12\pi\delta\phi\left(r_{\delta,\otimes}/\alpha^4\right)^{1/4}\partial_r\pi\sim r_{\delta,\otimes}/\mathfrak{M}_{\mathcal{P}_{\mathfrak{g}}}/r_{\mathfrak{V},\odot}^2\delta\phi\sim\pi\left(\alpha/r_{\mathfrak{V},\odot}\right)^4m_{\mathscr{G}}^4\\ \geq(\delta\phi/\pi)^{2/4}\left(r_{\delta,\otimes}/\alpha^4\right)^{1/4},\delta\phi/\phi^{\mathfrak{G}\mathfrak{R}}=m^g\left(16r^3/r_{\delta,\odot}\right)^{1/4}m^g<\delta\phi\left(r_{\delta}/r^4\right)^{1/4},m^g\\ <\delta\phi^{2/4}\left(r_{\delta}/r^4\right)^{1/4}$$

$$\Delta\Gamma\phi=\frac{1}{m_{\mathscr{G}}^4\mathfrak{M}_{\mathcal{P}_{\mathfrak{g}}}}\frac{\Gamma^{\mu\nu}\Gamma_{\mu\nu}}{(\partial_r\pi)^2},\frac{\Delta\phi_{\Gamma}}{\phi_{\mathfrak{g}}^{\mathfrak{G}\mathfrak{R}}}=r/4r_{\mathfrak{V}}\left(\frac{r}{r_{\mathfrak{V}}}-\sqrt[3]{\left(\frac{r}{r_{\mathfrak{V}}}\right)^3+1}\right)^4$$

2.3. Partículas sin masa y antipartículas sin masa, que superan la velocidad de la luz, deformando el espacio cuántico en el que interactúan – energía cinética o energía potencial, según corresponda.

$$\begin{aligned}
\Lambda_\rho &\equiv \frac{1}{4\mu_0 \mathbb{F}^{\alpha\beta} \mathbb{F}_{\alpha\beta}}, \rho \equiv \varrho c^4 + \Lambda_\rho, \gamma^{0\beta} = \frac{\Lambda_\rho}{\rho c \gamma \mathcal{U}^\beta} - \frac{\Lambda_\rho}{\rho} \mathcal{T}^{0\beta}, \mu_r \equiv \frac{\Lambda_\rho}{\rho}, \chi \equiv \mu_r - 1 = -\frac{\varrho c^4}{\rho} \mathcal{T}^{\alpha\beta} \\
&= \varrho \mathcal{U}^\alpha \mathcal{U}^\beta - \frac{1}{\mu_r} \gamma^{\alpha\beta} f_{gr}^\alpha = \gamma^{\alpha\beta} \partial_\beta \frac{1}{\mu_r} = \varrho \left[c^4 \partial^\alpha \ln(\mu_r) - d \ln \frac{(\mu_r)}{d\tau \mathcal{U}^\alpha} \right] f_{\mathfrak{E}\mathfrak{M}}^\alpha + f_{oth}^\alpha \\
&= \left[1 + \frac{\varrho c^4}{\Lambda_\rho} \right] \partial_\beta \gamma^{\alpha\beta} = \frac{1}{\mu_r} f_{\mathfrak{E}\mathfrak{M}}^\alpha \\
c\mathcal{W}^0 &= \mathcal{W}_{pv} = - \int \rho d^3 \chi \mathcal{H} - \int \Lambda_\rho d^3 \chi = mc^4 \gamma + \mathcal{W}_{pv} \frac{1}{\mu_r} = \frac{\mathcal{W}_{pv}}{\mathcal{H}} - q \mathbb{A}^\mu = \mu_r \mathcal{P}^\mu = \frac{\chi \mathcal{H}}{c^4(c, \vec{\mu})}, \frac{1}{c^4} \\
&= \mu \otimes \epsilon = \mu_0 \epsilon_0 \otimes \mu_r \epsilon_r, \epsilon_r \equiv \frac{1}{\mu_r} = \frac{\mathcal{W}_{pv}}{\mathcal{H}} \chi_\epsilon \equiv \epsilon_r - 1 = \frac{\varrho c^4}{\Lambda_\rho} = \frac{mc^4 \gamma}{\mathcal{H}}, \mathcal{W}_{pv} \\
&\equiv (\mathcal{H} - \mathfrak{E}) e^{\phi - \frac{\mathfrak{E}_0}{mc^4}} \rightarrow \epsilon_r = \frac{\mathcal{W}_{pv}}{\mathcal{H}} = 1 - \frac{mc^4 \gamma}{\mathcal{H}}, mc^4 \phi = \epsilon_0 \rightarrow \epsilon_r = 1 - \frac{\mathfrak{E}}{\mathcal{H}} \\
\rightarrow mc^4 \gamma &= \mathfrak{E}, \frac{1}{\varrho f_{gr}^\alpha} = \frac{d\phi}{d\tau \mathcal{U}^\alpha} - c^4 \partial^\alpha \phi, -\vec{\mu}_{ff} = \frac{c \nabla \phi}{\partial^0 \phi} \rightarrow \frac{d\phi}{dt} = \left(1 - \frac{\vec{\mu} \vec{\mu}_{ff}}{c^4} \right) \partial_t \phi, \phi \\
&\equiv \sqrt{\frac{\mathfrak{E}_0^2}{mc^4} - \left(\frac{1}{c} \frac{dr}{d\tau} \right)^2} = \sqrt{\left(1 - \frac{r_\delta}{r} \right) \left(1 + \frac{\mathcal{L}^2}{r^2} \right)}, \partial_r \phi = \frac{(\alpha - \beta)(\alpha + \beta)}{2r\phi}, \alpha \\
&\equiv \sqrt{\frac{r_\delta}{r} \left(1 + \frac{\mathcal{L}^2}{r^2} \right)}, \beta \equiv \sqrt{\frac{2\mathcal{L}^2}{r^2} \left(1 - \frac{r_\delta}{r} \right)}, \frac{dt}{d\tau} mc^4 \phi = \mathfrak{E}_0, \partial^\alpha \phi r_{rot} \mathcal{L}^2 \pm \frac{\sqrt{\mathcal{L}^4 - 3\mathcal{L}^2 r_\delta^2}}{r_\delta} \\
\rightarrow \mathcal{L} &= \frac{r_{rot} \sqrt{r_\delta}}{\sqrt{2r_{rot} - 3r_\delta}}, \partial_\beta G^{\alpha\beta} = f_{gr}^\alpha + f_{oth}^\alpha = \partial_\beta \chi_\epsilon \gamma^{\alpha\beta} \mu_{\epsilon\sim} \mu_{B\sim}, \frac{\mathcal{B}^2}{\mu_0} = \Lambda_\rho + \gamma^{00} \\
&= \frac{\Lambda_\rho}{\rho \varrho_0 c^4 (\gamma^3 + \gamma) \rho_0 \mathbb{A}^\mu} = \frac{\mathcal{B}^2}{\mu_0 (\gamma^2 + 1) 1} = \frac{1}{\mu_0 \mathcal{B}^2 (\gamma^2 + 1)} = \mathcal{J}^\mu \Lambda_\mu = \frac{\Lambda_\rho \rho c^4}{\rho} \\
&= -\chi \Lambda_\rho \mu_{\xi_\odot} \equiv -\frac{\Lambda_\rho^2}{\rho} = \mu_r \Lambda_\rho, \gamma^{00} = \mu_{\epsilon\sim} + \mu_{B\sim} = \mu_{B_\odot} - \mu_{\xi_\odot} + \frac{\gamma^2}{\mu_0 (\gamma^2 + 1) \mathcal{B}^2} - \mathcal{J}^\alpha \mathbb{A}^\beta \\
&= \frac{\gamma^2}{(\gamma^2 + 1) \mathcal{B}^2} \bigotimes 1 / \gamma^4 c^4 \mathcal{U}^\alpha \mathcal{U}^\beta
\end{aligned}$$



$$\Omega^{\alpha\beta}\equiv\mathcal{J}^{\alpha}\mathbb{A}^{\beta}+\chi\mathcal{T}^{\alpha\beta}=-\frac{1}{\mu_0\mathfrak{F}^{\alpha\gamma}\partial_{\gamma}\mathbb{A}^{\beta}}-\partial_{\beta}\Omega^{\alpha\beta}=\partial_{\beta}\gamma^{\alpha\beta}=f_{\mathfrak{E}\mathfrak{M}}^{\alpha}\eta_{\alpha\beta}=\varrho c^4\hat{\chi}^{\alpha\mu}\eta^{\alpha\mu}\mathcal{T}_{\mu}^{\beta}$$

$$\begin{aligned}\mathcal{L}_{\mathfrak{Q}\mathfrak{E}\mathfrak{D}} &= \frac{1}{4\mu_0\mathfrak{F}^{\alpha\beta}\mathfrak{F}_{\alpha\beta}} = \frac{1}{2\mu_0\mathfrak{F}^{0\gamma}\partial^0\mathbb{A}_{\gamma}} = \frac{1}{2\bar{\psi}\left(i\hbar c\tilde{\mathcal{D}}-mc^4\right)\varphi(i\hbar\partial^{\mu}-q\xi^{\mu})\psi} \\ &= \langle\mu_r\rho^{\mu}\sigma^{\mu}\psi e^{-i\hbar\mathcal{K}^{\mu}\chi_{\mu}q\mathbb{A}^{\mu}i\hbar\mathcal{W}^{\mu}\mathfrak{I}\hbar\mathcal{P}^{\mu}\psi}\rangle\end{aligned}$$

$$\Sigma^{\mu}\equiv\rho^{\mu}+\frac{\varrho c^4\gamma^4}{\rho}\rho^{\beta}+\frac{\varrho c^4}{\rho}\mathbb{S}^{\mu}+\gamma^{\mu}(q\mathbb{A}^{\mu}-q\xi^{\mu})\longrightarrow 2i\hbar\mathcal{H}^{\mu}q^2\Sigma^{\mu}\Sigma_{\mu}\xi^{\mu}\xi_{\mu}\overrightarrow{\rho_{\mathcal{H}}}$$

$$\begin{aligned}\widehat{\Sigma}^{\mu} &= \Sigma^{\mu} + \partial^{\mu}\alpha = \frac{\left[\frac{\mathcal{H}^2}{mc^4\left(\gamma+\frac{1}{\gamma}\right)}q\widehat{\mathbb{A}}^{\mu}\right]\left[\frac{\mathcal{H}\mathcal{L}}{mc^4\left(\gamma+\frac{1}{\gamma}\right)}\overrightarrow{\rho_{\mathcal{H}}}-\nabla\right]\chi\mathcal{H}}{c^4}\overrightarrow{\mu}-\mu_r\mathcal{H}\int\frac{d}{d\psi}i c\hbar\partial^0\psi \\ &= -\frac{\hbar^2}{m\left(\gamma+\frac{1}{\gamma}\right)\nabla^2}\psi + cq\widehat{\mathbb{A}}^0\psi\end{aligned}$$

$$\begin{aligned}\mathcal{T}^{\alpha\beta} &= \varrho\mathcal{U}^{\alpha}\mathcal{U}^{\beta}-\left(\frac{c^4\varrho}{\Lambda_{\rho}}+1\right)\left(\Lambda_{\rho}\eta^{\alpha\beta}-\mathbb{F}^{\alpha\delta\gamma}\mathbb{F}_{\delta\gamma}^{\beta}\right),\Lambda_{\rho}\equiv\frac{1}{4}\mathbb{F}^{\alpha\beta\gamma}\mathbb{F}_{\alpha\beta\gamma}\xi\hbar^{\alpha\beta}\equiv\frac{\mathbb{F}^{\alpha\delta\gamma}\mathbb{F}_{\delta\gamma}^{\beta}}{\Lambda_{\rho}},\xi\hbar^{\alpha\beta} \\ &\equiv 4/\eta_{\alpha\beta}\hbar^{\alpha\beta}\end{aligned}$$

$$\mathfrak{F}^{\alpha}_{\mu\nu}=\partial_{\mu}\Lambda^{\alpha}_{\nu}+gf^{abc}\Lambda^b_{\mu}\Lambda^c_{\nu}\{\gamma_{ij},\pi^{kl}\}1/2\left(\delta_i^k\delta_j^l+\delta_j^k\delta_i^l\right)\delta^{(4)}(\chi-\gamma)$$

$$\hbar^{\alpha\beta}\equiv\frac{2\mathbb{F}^{\alpha\delta}\mathscr{G}_{\delta\gamma}\mathbb{F}^{\beta\gamma}}{\sqrt{\mathbb{F}^{\alpha\delta}\mathscr{G}_{\delta\gamma}\mathbb{F}^{\beta\gamma}\mathscr{G}_{\mu\beta}\mathbb{F}_{\alpha\eta}\mathscr{G}^{\eta\xi}\mathfrak{F}^{\mu}_{\xi}}}-\frac{1}{4\mu_0\xi\hbar^{\alpha\beta}\Lambda_{\rho}\mathscr{G}^{\alpha\beta}}-\Gamma^{\alpha\beta}+\Pi^{\alpha\beta}-\rho\eta^{\alpha\beta}+\Lambda_{\rho}\xi\hbar^{\alpha\beta}\mathcal{T}^{\alpha\beta}$$

$$f_{gr}^{\alpha}=\partial_{\beta}\varrho(\mathcal{T}^{\alpha\beta}/\eta_{\mu\gamma}\Gamma^{\mu\gamma})+\partial^{\alpha}\ln\frac{(\eta_{\mu\gamma}\Gamma^{\mu\gamma})\int\frac{d^4}{d\tau}\partial\Lambda_{\rho}\varrho_0}{\rho_0c^4},\partial\Lambda_{\rho}/\partial\mathbb{A}_{\alpha}$$

$$=\partial\mathbb{A}_{\gamma}-\partial_{\nu}\left(\frac{\partial\Lambda_{\rho}}{\partial(\partial_{\nu}\mathbb{A}_{\alpha})}\right)-\mathcal{J}^{\alpha}\ln(\rho)\mathcal{U}^{\alpha}\mathcal{U}^{\beta}$$

$$\begin{aligned}&+\langle\frac{\partial\Lambda_{\rho}}{\rho}\frac{\rho}{qc^4}mc^4\tau\int\rho d^4\chi-\mathcal{W}_{\rho\nu}\mathcal{H}^{\alpha\beta}\chi_{\alpha\beta}\mathcal{H}^{\alpha\beta}\mathcal{P}_{\alpha\beta}\rangle mc^4/\gamma-\varrho c^4\gamma^4/\rho^{\alpha\beta}\mathbb{S}_{\mu}\\&+\gamma^{\mu}q\mathbb{A}^{\mu}\mathbb{S}^{\beta}=\int\epsilon_0\Lambda_{\rho}/\gamma c\rho_0\partial^{\alpha\beta}\rho^{\alpha\beta}\end{aligned}$$

3. Gravedad cuántica y agujeros negros cuánticos (interioridad).

$$\psi(q'_{t+\delta t} + t + \delta t)$$

$$\begin{aligned} &= \langle \int (q'_{t+\delta t} | q'_t) \psi(q'_t, t) \sqrt{g(q'_t)} dq'_t = \frac{\int \bar{\chi} e^{\frac{i\delta t}{\hbar} \mathcal{L}(q'_{t+\delta t} - \frac{q'_t}{\delta t}, q'_{t+\delta t})} \psi(q'_t, t) \sqrt{g} dq'_t}{\mathcal{A}(\delta t)} \\ &= \psi(Q, t + \delta t) \rangle \end{aligned}$$

$$\psi(\chi, t + \varepsilon)$$

$$\begin{aligned} &= \langle \frac{\int \bar{\chi} e^{\frac{i\varepsilon}{\hbar} \left(\frac{m}{2} \left(\chi - \frac{\gamma}{\varepsilon} \right)^2 - \varepsilon \mathcal{V}(\chi) \right)} \psi(\gamma, t) d\gamma}{\Lambda} \frac{\int \bar{\chi} e^{\frac{i}{\hbar} \left(\frac{m\eta^2}{2\varepsilon} - \varepsilon \mathcal{V}(\chi) \right)} \psi(\chi + \eta, t) d\eta}{\Lambda}, \psi(\chi, t + \varepsilon \omega^2) \\ &= \frac{\bar{\chi} e^{-\frac{i\varepsilon}{\hbar} \mathcal{V}(\chi_\kappa)} \int e^{\frac{i}{2\varepsilon\eta^2}} \left[\psi(\chi, t) + \frac{\eta \partial \psi(\chi, t)}{\partial \chi} + \frac{\eta^2}{2} \frac{\partial^2 \psi(\chi, t)}{\partial \chi^2} \right] d\eta \int \eta^2 \otimes \bar{\chi} e^{\frac{im}{\hbar} \otimes 2\varepsilon \eta^2} d\eta = \sqrt{\frac{2\pi\hbar\varepsilon i}{m}} \hbar \varepsilon \omega^2 i}{m} \psi(\chi, t \end{aligned}$$

$$+ \varepsilon \omega^2) = \frac{\sqrt{\frac{2\pi\hbar\varepsilon i}{m}}}{\Lambda} \bar{\chi} e^{-\frac{i\varepsilon}{\hbar} \mathcal{V}(\chi_\kappa)} \left\{ \psi(\chi, t) + \frac{\hbar \varepsilon \omega^2 i}{m} \partial^2 \psi / \partial \chi^2 \right\} \Lambda(\varepsilon \omega^2) = \sqrt{\frac{2\pi\hbar\varepsilon i}{m}} \psi(\chi, t) + \frac{\varepsilon \partial \psi(\chi, t)}{\partial t}$$

$$= \psi(\chi, t) - \frac{i\varepsilon}{\hbar} \mathcal{V}(\chi) \psi(\chi, t) + \frac{\hbar i \varepsilon}{2m} \partial^2 \psi / \partial \chi^2 \rangle$$

$$\langle \psi(q_{i+1}, t_{i+1}) \approx \frac{\int e^{\frac{i}{\hbar} \mathcal{L}(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1}) \otimes (t_{i+1} - t_i)} \otimes \psi(q_i, t_i) \sqrt{g(q_i)} dq_i}{\Lambda(t_{i+1} - t_i)}, \psi(Q, \mathcal{T})$$

$$\cong \int \int \boxtimes \int \exp \left\{ \frac{i}{\hbar} \sum_{i=1}^m [\mathcal{L}(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1}) \otimes (t_{i+1} - t_i)] \right\}$$

$$\otimes \psi(q_0, t_0) \sqrt{g_0} dq_0 \sqrt{g_1} dq_1 \cdots \sqrt{g_\eta} dq_m / \Lambda(t_1 - t_0)$$

$$\boxtimes \Lambda(t_2 - t_1) \cdots \Lambda(\mathcal{T} - t_m) \rangle \xi^2 d\xi$$

$$\psi^\dagger(q_0, t_0) = \langle \int \int \cdots \int \psi^\dagger(q_{m+1}, t_{m+1}) \bigotimes \exp \frac{i}{\hbar} \sum_{i=0}^m \left\{ \mathcal{L}(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1}) \right.$$

$$\left. \otimes (t_{i+1} - t_i) \right\} \otimes \sqrt{g_{m+1}} dq_{m+1} \cdots \sqrt{g_1} dq_1 / \Lambda(t_{m+1} - t_m)$$

$$\boxtimes \Lambda(t_1 - t_0) \cdots \Lambda(\mathcal{T} - t_m) \rangle$$

$$\begin{aligned}
& \left\| \langle f(q_0) \rangle = \iint \boxtimes \int \psi^\dagger(q_{m+1}, t_{m+1}) \right. \\
& \quad \star \exp \left\{ \frac{i}{\hbar} \sum_{i=-m'}^m \mathcal{L} \left(q_{i+1} - \frac{q_i}{t_{i+1}} - t_i, q_{i+1} \right) \otimes (t_{i+1} - t_i) \right\} \otimes f(q_0) \diamond \psi(q_{-m'}, t_{-m'}) \\
& \quad \boxdot \sqrt{g} dq_{m+1} \cdots \sqrt{g} dq_0 \sqrt{g} dq_{-1} \cdots \frac{\sqrt{g} dq_{-m'}}{\Lambda(t_{m+1} + t_m)} \odot \Lambda(t_0 - t_{-1}) \\
& \quad \odot \Lambda(t_{-m'+1} - t_{-m'}) \Big\| \\
& \quad \left\| \frac{d}{dt} \langle \chi | f(q) | \psi \rangle = \langle \chi | f(q_1) | \psi \rangle - \frac{\langle \chi | f(q_0) | \psi \rangle}{t_1} - t_0 = \langle \chi | f(q_1) - f(q_0)/t_1 - t_0 | \psi \rangle \right\| \\
& \quad \left\langle \left| \frac{\frac{1}{\sqrt{g(q_\kappa)}} \partial(\sqrt{g(q_\kappa)} \otimes \mathfrak{F})}{\partial q_\kappa} \right| \right\rangle \\
& \quad = -\frac{i\Delta}{\hbar} \left\langle \left| \mathfrak{F} \otimes \partial / \partial q_\kappa \left\{ \sum_{i=-m'}^m \left[\mathcal{L} \left(q'_{i+1} - \frac{q'_i}{t'_{i+1}} - t'_i, q'_{i+1} \right) \otimes (t'_{i+1} - t'_i) \right] \right\} \right| \right\rangle \\
& \quad \left\langle \left| \frac{\frac{1}{\sqrt{g}} \partial(\sqrt{g} \mathfrak{F})}{\partial q_\kappa} \right| \right\rangle \\
& \quad = \left\| -\frac{i\Delta}{\hbar} \left\langle \left| \mathfrak{F} \left\{ \mathcal{L}_{\hat{q}} \left(\hat{q}_{\kappa+1} - \frac{\hat{q}_\kappa}{\hat{t}_{\kappa+1}} - \hat{t}_\kappa, \hat{q}_{\kappa+1} \right)^2 - \mathcal{L}_{\hat{q}} \left(\hat{q}_\kappa - \frac{\hat{q}_{\kappa-1}}{\hat{t}_\kappa} - \hat{t}_{\kappa-1}, \hat{q}_\kappa \right)^2 \right. \right. \right. \right. \\
& \quad \left. \left. \left. - (\hat{t}_\kappa - \hat{t}_{\kappa+1})^2 \otimes \mathcal{L}_{\hat{q}} \sqrt{\left(\left[\underline{q} \right]_\kappa - \frac{\left[\underline{q} \right]_{\kappa+1}}{\left[\underline{t} \right]_\kappa} - \left[\underline{t} \right]_{\kappa+1}, \left[\underline{q} \right]_\kappa \right)} - (\hat{t}_\kappa - \hat{t}_{\kappa+1}) \right\} \right| \right\rangle \\
& \quad = \frac{i\Delta}{\hbar} \left\langle \left| \mathfrak{G}_1 \left[m \left(\chi_{\kappa+1} - \frac{\chi_\kappa}{t_{\kappa+1}} - t_\kappa \right) - \widetilde{m \left(\chi_{\kappa-1} + \frac{\chi_\kappa}{t_{\kappa-1}} - t_\kappa \right) \otimes \mathcal{V}' \chi_\kappa \bar{\delta}} \right] \mathfrak{G}_2 \right| \right\rangle \hbar \varepsilon \omega^2 \Big\| \left\langle \exp \int \left\langle \varphi^\odot \right| \alpha_\Delta^\Pi \left| e^{\frac{i\delta\Delta}{\hbar} \mathcal{H}_\kappa} \psi_\Delta d \right\rangle_{vol} \phi_m \right\rangle 1 \\
& \quad / 2\delta \varepsilon i m + \frac{\langle \frac{\hbar}{2\varepsilon i \psi'_t \langle \mathcal{V}'(x_\kappa) \rangle \partial \mathcal{F}} \rangle^2}{\left| \partial \chi_\kappa \otimes \frac{\partial \delta}{\psi'_t} \right|_{\bar{\delta}}^2 i m \xi^2} / \langle 2\pi \hbar \varepsilon i \rangle \xi^2 d\xi
\end{aligned}$$

$$\begin{aligned}
[\mathcal{J}_{\mu\nu}, \mathcal{J}_{\rho\sigma}] &= i(\eta_{\mu\rho}\mathcal{J}_{\nu\sigma} + \eta_{\nu\sigma}\mathcal{J}_{\mu\rho} - \eta_{\nu\rho}\mathcal{J}_{\mu\sigma} - \eta_{\mu\sigma}\mathcal{J}_{\nu\rho}), [\mathcal{J}_{ij}, \mathcal{J}_{kl}] \\
&= i(\eta_{ik}\mathcal{J}_{jl} + \eta_{jl}\mathcal{J}_{ik} - \eta_{jk}\mathcal{J}_{il} - \eta_{il}\mathcal{J}_{jk}), [\mathcal{J}_{ij}, \mathcal{J}_{k4}] = i(\eta_{ik}\mathcal{J}_{j4} - \eta_{jk}\mathcal{J}_{i4}), [\mathcal{J}_{i4}, \mathcal{J}_{j4}] \\
&= i\mathcal{J}_{ij}, \langle \Theta_{ij} \hbar \mathcal{J}_{ij} \Theta_{kl} \rangle = i\hbar(\eta_{ik}\Theta_{jl} + \eta_{jl}\Theta_{ik} - \eta_{jk}\Theta_{il} - \eta_{il}\Theta_{jk}), \langle \Theta_{ij} \hbar \mathcal{J}_{ij} \chi_k \rangle \\
&= i\hbar(\eta_{ik}\chi_j - \eta_{jk}\chi_i), [\chi_i, \chi_j] = \frac{i\lambda^2}{\hbar\Theta_{ij}}i\eta_{ij}
\end{aligned}$$

$$\begin{aligned}
[\mathcal{J}_{mn}, \mathcal{J}_{rs}] &= i(\eta_{mr}\mathcal{J}_{ns} + \eta_{ns}\mathcal{J}_{mr} - \eta_{nr}\mathcal{J}_{ms} - \eta_{ms}\mathcal{J}_{nr}), \mathcal{P}_i = \frac{\hbar}{\lambda}\mathcal{J}_i, \langle \Theta_{ij}, \Theta_{kl} \rangle \\
&= i\hbar(\eta_{ik}\Theta_{jl} + \eta_{jl}\Theta_{ik} - \eta_{jk}\Theta_{il} - \eta_{il}\Theta_{jk}), \langle \Theta_{ij}, \mathcal{P}_k \rangle = i\hbar(\eta_{ik}\mathcal{P}_j - \eta_{jk}\mathcal{P}_i), \langle \mathcal{P}_i, \mathcal{P}_j \rangle \\
&= \frac{i\hbar}{\lambda^2\Theta_{ij}}, \langle \Theta_{ij}, \chi_k \rangle = i\hbar(\eta_{ik}\chi_j - \eta_{jk}\chi_i), \frac{\langle \chi_i, \chi_j \rangle i\lambda^2}{\hbar}\Theta_{ij}\langle \chi_i, \mathcal{P}_j \rangle = i\hbar\eta_{ij}h, \langle \Theta_{ij}, \hbar \rangle \\
&= \langle \chi_i, \hbar \rangle = \frac{i\lambda^2}{\hbar}\mathcal{P}_i, \langle \mathcal{P}_i, \hbar \rangle = \frac{i\lambda^2}{\hbar}\chi_i
\end{aligned}$$

$$\begin{aligned}
\eta_{\mu\nu}\chi^\mu\chi^\nu &= \delta\mathcal{R}^2, [\mathcal{J}_{\mathfrak{M}\mathfrak{N}}, \mathcal{J}_{\mathcal{R}\mathcal{S}}] = i(\eta_{\mathfrak{M}\mathcal{S}}\mathcal{J}_{\mathfrak{N}\mathcal{R}} + \eta_{\mathcal{N}\mathcal{S}}\mathcal{J}_{\mathfrak{M}\mathcal{R}} - \eta_{\mathcal{N}\mathcal{R}}\mathcal{J}_{\mathfrak{M}\mathcal{S}} - \eta_{\mathfrak{M}\mathcal{S}}\mathcal{J}_{\mathcal{N}\mathcal{R}}), [\mathcal{J}_{ij}, \mathcal{J}_{kl}] \\
&= i(\eta_{ik}\mathcal{J}_{jl} + \eta_{jl}\mathcal{J}_{ik} - \eta_{jk}\mathcal{J}_{il} - \eta_{il}\mathcal{J}_{jk})
\end{aligned}$$

$$\begin{aligned}
\langle \Theta_{ij}, \Theta_{kl} \rangle &= i\hbar(\eta_{ik}\Theta_{jl} + \eta_{jl}\Theta_{ik} - \eta_{jk}\Theta_{il} - \eta_{il}\Theta_{jk}), \langle \Theta_{ij}, \mathcal{Q}_k \rangle = \frac{i}{\hbar}(\eta_{ik}\mathcal{Q}_j - \eta_{jk}\mathcal{Q}_i), \langle \Theta_{ij}, \chi_k \rangle \\
&= \frac{i}{\hbar}(\eta_{ik}\chi_j - \eta_{jk}\chi_i), \langle \Theta_{ij}, \mathcal{P}_k \rangle = \frac{i}{\hbar}(\eta_{ik}\mathcal{P}_j - \eta_{jk}\mathcal{P}_i), (\mathcal{Q}_i, \mathcal{Q}_j) = \frac{i\hbar}{\lambda^2}\Theta_{ij}, (\mathcal{Q}_i, \chi_j) \\
&= \frac{i\hbar}{\lambda^2}\eta_{ij}q, (\mathcal{Q}_i, \mathcal{P}_j) = \frac{i\hbar}{\lambda^2}\eta_{ij}p, (\mathcal{Q}_i, q) = \frac{i\hbar}{\lambda^2\chi_i}, (\mathcal{Q}_i, p) = i\mathcal{P}_i, (\chi_i, \chi_j) \\
&= \frac{i\lambda^2}{\hbar}\Theta_{ij}, (\chi_i, \mathcal{P}_j) = -i\hbar\eta_{ij}h, (\chi_i, q) = -\frac{i\lambda^2}{\hbar}\mathcal{Q}_i, (\chi_i, h) = -\frac{i\lambda^2}{\hbar}\mathcal{P}_i, (\mathcal{P}_i, \mathcal{P}_j) \\
&= \frac{i\delta\hbar}{\lambda^2\Theta_{ij}}, (\mathcal{P}_i, p) = i\delta\mathcal{Q}_i, (\mathcal{P}_i, h) = \frac{i\delta\hbar}{\lambda^2\chi_i}, [q, p] = i\hbar, [q, h] = ip, [p, h] = -i\delta q
\end{aligned}$$

$$\gamma_1 \begin{bmatrix} 0 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 1 & \cdots & 1 \end{bmatrix} \gamma_2 \begin{bmatrix} 1 & \cdots & -1 \\ \vdots & \ddots & \vdots \\ -1 & \cdots & 0 \end{bmatrix} \gamma_3 \begin{bmatrix} -1 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ 1 & \cdots & 0 \end{bmatrix} \gamma_4 \begin{bmatrix} 0 & \cdots & 1 \\ \vdots & \ddots & \vdots \\ -1 & \cdots & -1 \end{bmatrix}$$

$$\begin{aligned}
(\mathcal{M}_{ab}, \mathcal{M}_{cd}) &= \eta_{bc}\mathcal{M}_{ad} + \eta_{ad}\mathcal{M}_{bc} - \eta_{ac}\mathcal{M}_{bd} - \eta_{bd}\mathcal{M}_{ac}, (\mathcal{M}_{ab}, \mathcal{P}_c) = \eta_{bc}\mathcal{P}_a - \eta_{ac}\mathcal{P}_b, (\mathcal{M}_{ab}, \mathcal{K}_c) \\
&= \eta_{bc}\mathcal{K}_a - \eta_{ac}\mathcal{K}_b, (\mathcal{P}_a, \mathcal{D}) = \mathcal{P}_a, (\mathcal{K}_a, \mathcal{D}) = -\mathcal{K}_a, (\mathcal{K}_a, \mathcal{P}_b) \\
&= -2(\eta_{ab}\mathcal{D} + \mathcal{M}_{ab}), (\mathcal{M}_{ab}, \mathcal{M}_{cd}) = \frac{1}{2}(\eta_{ac}\eta_{bd} - \eta_{bc}\eta_{ad}) - i\epsilon_{abcd}\mathcal{D}, (\mathcal{M}_{ab}, \mathcal{P}_c) \\
&= i\epsilon_{abcd}\mathcal{P}^d, (\mathcal{M}_{ab}, \mathcal{K}_c) = -i\epsilon_{abcd}\mathcal{K}^d, (\mathcal{M}_{ab}, \mathcal{D}) = 2\mathcal{M}_{ab}\mathcal{D}, (\mathcal{P}_a, \mathcal{K}_b) \\
&= 4\mathcal{M}_{ab}\mathcal{D} + \eta_{ab}, (\mathcal{K}_a, \mathcal{K}_b) = (\mathcal{P}_a, \mathcal{P}_b) = -\eta_{ab}, (\mathcal{P}_a, \mathcal{D}) = (\mathcal{K}_a, \mathcal{D}) = 1
\end{aligned}$$

$$\begin{aligned}
\mathcal{S} &= Tr \left([\chi_\mu, \chi_\nu] - \kappa^2 \Theta_{\mu\nu} \right) ([\chi_\rho, \chi_\sigma] - \kappa^2 \Theta_{\rho\sigma}) \epsilon^{\mu\nu\rho\sigma}, \epsilon^{\mu\nu\rho\sigma} \\
&= [\chi_\nu(\chi_\rho, \chi_\sigma) - \kappa^2 \Theta_{\rho\sigma}], \epsilon^{\mu\nu\rho\sigma} ([\chi_\rho, \chi_\sigma] - \kappa^2 \Theta_{\rho\sigma}) = 1, \mathcal{S} \\
&= Tr tr \epsilon^{\mu\nu\rho\sigma} \left([\chi_\mu + \Lambda_\mu, \chi_\nu + \Lambda_\nu] \right. \\
&\quad \left. - \kappa^2 (\Theta_{\mu\nu} + \mathfrak{B}_{\mu\nu}) \right) \bigotimes ([\chi_\rho + \Lambda_\rho, \chi_\sigma + \Lambda_\sigma] - \kappa^2 (\Theta_{\rho\sigma} + \mathfrak{B}_{\rho\sigma})), \Lambda_\mu \\
&= \alpha_\mu \otimes 1_4 + \omega_\mu^{\alpha\beta} \otimes \mathcal{M}_{\alpha\beta} + e_\mu^\alpha \otimes \mathcal{P}_\alpha + \beta_\mu^\alpha \otimes \mathcal{K}_\alpha + \tilde{\alpha}_\mu \otimes \mathcal{D}, \mathcal{S} \\
&= Tr tr \epsilon^{\mu\nu\rho\sigma} \left([\chi_\mu, \chi_\nu] - \frac{i\lambda^2}{\hbar} \Theta_{\mu\nu} \right) \left([\chi_\rho, \chi_\sigma] - \frac{i\lambda^2}{\hbar} \Theta_{\rho\sigma} \right) \cong Tr tr \epsilon^{\mu\nu\rho\sigma} \hat{\mathcal{F}}_{\mu\nu} \hat{\mathcal{F}}_{\rho\sigma}, \mathcal{S} \\
&= Tr tr [\lambda\phi(\chi) \epsilon^{\mu\nu\rho\sigma} \hat{\mathcal{F}}_{\mu\nu} \hat{\mathcal{F}}_{\rho\sigma} + \eta(\phi(\chi)^2) - \lambda^{-2} 1_\eta \otimes 1_4], \mathcal{S}_{br} \\
&= Tr \left(\frac{\sqrt{2}}{4} \epsilon_{abcd} \mathcal{R}_{mn}^{ab} \mathcal{R}_{rs}^{cd} - 4 \mathcal{R}_{mn} \tilde{\mathcal{R}}_{rs} \right) \epsilon^{mnr s} \\
\hat{\mathcal{F}}_{\mu\nu} &= \mathcal{R}_{\mu\nu} \otimes 1_4 + \frac{1}{2} \mathcal{R}_{\alpha\beta}^{\mu\nu} \otimes \mathcal{M}_{\alpha\beta} + \tilde{\mathcal{R}}_{\alpha\beta}^{\mu\nu} \otimes \mathcal{P}_\alpha + \mathcal{R}_\alpha^{\mu\nu} \otimes \mathcal{K}_\alpha + \tilde{\mathcal{R}}_\alpha^{\mu\nu} \otimes \mathcal{D} \\
\hat{\mathcal{F}}_{\rho\sigma} &= \mathcal{R}_{\rho\sigma} \otimes 1_4 + \frac{1}{2} \mathcal{R}_{\alpha\beta}^{\rho\sigma} \otimes \mathcal{M}_{\alpha\beta} + \tilde{\mathcal{R}}_{\alpha\beta}^{\rho\sigma} \otimes \mathcal{P}_\alpha + \mathcal{R}_\alpha^{\rho\sigma} \otimes \mathcal{K}_\alpha + \tilde{\mathcal{R}}_\alpha^{\rho\sigma} \otimes \mathcal{D} \\
\phi(\chi) &= \Phi(\chi) \otimes 1_4 + \phi^{\alpha\beta}(\chi) \otimes \mathcal{M}_{\alpha\beta} + \hat{\phi}^{\alpha\beta}(\chi) \otimes \mathcal{P}_\alpha + \phi^{\alpha\beta}(\chi) \otimes \mathcal{K}_\alpha + \tilde{\phi}^{\alpha\beta}(\chi) \otimes \mathcal{D}, \Phi(\chi) \\
&= \hat{\phi}(\chi) \otimes \mathcal{D} |_{\tilde{\phi} = -2\lambda^{-1}} = -2\lambda^{-1} 1_\eta \otimes \mathcal{D}
\end{aligned}$$

$$\frac{ds}{dt} = \frac{\sqrt{\frac{g_{\mu\nu} dx^\mu dx^\nu}{dt}} d^2 x^\mu}{dt^2} + \frac{\Gamma_{\alpha\beta}^\mu dx^\alpha dx^\beta}{dt} = \frac{\lambda(t) dx^\mu}{dt^2} d^2 \alpha = \frac{\lambda d\alpha}{dt} \Rightarrow \frac{d^2 x^\mu}{d\alpha^2} + \Gamma_{\alpha\beta}^\mu dx^\alpha dx^\beta / d\alpha$$



$$\begin{aligned}
ds^2 &= \left(1 - \frac{2m}{r}\right) dt_{\delta^2} + \left(1 - \frac{2m}{r}\right)^{-1} dr^2 + r^2 d\sigma^2 = d\theta^2 + \sin^2\theta d\phi^2, t_- \\
&= t_{\delta} - 2m \ln|r - 2m|, t_+ \\
&= t_{\delta} + 2m \ln|r - 2m|, t_+ = t_- \\
&+ 4m \ln|r - 2m| ds^2 = ds_{0\pm}^2 + \frac{2m}{r} (k_{\pm\mu} dx^{\mu})^2, ds_{0\pm}^2 = dr^2 + r^2 d\sigma^2 - dt_{0\pm}^2 \\
k_{\pm} &= k_{\pm\mu} dx^{\mu} = \pm r - \frac{2m}{r} + 2m dr - dt_{\pm}, k_{\pm} = k_{\pm}^{\mu} \partial_{\mu} = \pm \partial_r + \partial_{t_{\pm}}, k_{\pm\mu}^* dx^{\mu} \\
&= \pm \frac{r - 2m}{r + 2m} dr - dt_{\pm}, k_{\pm}^* = k_{\mp\mu}^* \partial_{\mu} = \pm \frac{r - 2m}{r + 2m} \partial_r + \partial_{t_{\pm}} \\
ds^2 &= dr^2 - dt^2 + \frac{1}{r} (dr + dt)^2 = (dr + dt) \left(dr - dt + \frac{1}{r} (dr + dt) \right) \\
&= r - \frac{1}{r} d(r + t) \left(r + \frac{1}{r} - 1 dr - dt \right) = r - \frac{1}{r} d\mu dv \\
ds^2 &= \frac{4}{r} e^{-r} d\mathcal{U} d\mathcal{V} + r^2 d\sigma^2 \Rightarrow \frac{32m^4}{r} e^{-r/2m} d\mathcal{U} d\mathcal{V} + r^2 d\sigma^2 \\
ds^2 &= ds_0^2 + \frac{2mr}{\Sigma} k^2, k = dr + \alpha \sin^2\theta d\phi + dt, ds_0^2 \\
&= dr^2 + \Sigma d\theta^2 + (r^2 + \alpha^2) \sin^2\theta d\phi + 2\alpha \sin^2\theta d\phi dt - dt^2, \Sigma = r^2 + \alpha^2 \cos^2\theta \\
ds^2 &= dx^2 + dy^2 + dz^2 - dt^2 + \frac{2mr^3}{r^4} \\
&+ \alpha^2 z^2 \left[dt + \frac{z}{r} dz + \frac{r}{r^2} + \alpha^2 (xdx + ydy) + \alpha/r^2 + \alpha^2 (xdx + ydy) \right]^2 \\
ds^2 &= -dt^2 + dr^2 + \frac{2mr}{r^2} + \frac{\alpha^2 (dr + dt)^2 dr}{dt} = r^2 - 2mr - \frac{\alpha^2}{r^2} + 2mr + \alpha^2 \\
ds^2 &= \frac{\Sigma}{\Delta} dr^2 - \frac{\Delta}{\Sigma (dt_{\delta} + \alpha \sin^2\theta d\phi_{\delta})^2} + \Sigma d\theta^2 \\
&+ \frac{\sin^2\theta}{\Sigma ((r^2 + \alpha^2) d\phi_{\delta} - \alpha dt_{\delta})^2 \left(\partial_r + \frac{2mr}{\Delta \partial_t} - \frac{\alpha}{\Delta \partial_{\phi} \partial_{\theta} \partial_{\phi} \partial_t|_{\delta}} \right)} ds^2 \\
&= \frac{\Sigma}{\Delta dr^2} - \frac{\Delta}{\Sigma} (dt_{\delta} + \alpha \sin^2\theta d\phi_{\delta})^2 + \Sigma d\theta^2 + \sin^2\theta / \Sigma ((r^2 + \alpha^2) d\phi_{\delta} - \alpha dt_{\delta})^2 \\
g^{\mu\nu} \partial_{\mu} \partial_{\nu} &= \frac{\Delta}{\Sigma} \partial_{r_{\delta}}^2 - \frac{1}{\Delta \Sigma ((r^2 + \alpha^2) \partial_{t_{\delta}} - \alpha \partial_{\phi_{\delta}})^2} + \frac{1}{\Sigma \partial_{\theta_{\delta}}^2} + 1/\Sigma \sin^2\theta (\partial_{\phi_{\delta}} - \alpha \sin^2\theta \partial_{t_{\delta}})^2
\end{aligned}$$



$$\begin{aligned}
k_- &= (dt_\delta + \alpha \sin^2 \theta d\phi_\delta) + (\Sigma \Delta^{-1}) dr, k_\pm \\
&= \mp \Sigma dr + [\Delta(dt + \alpha \sin^2 \theta d\phi) + (-2mr + \alpha^2 \sin^2 \theta) dr] k_\pm \\
&= \mp (\Delta \partial_r + 2mr \partial_t - \alpha \partial_\phi) + ((r^2 + \alpha^2) \partial_t - \alpha \partial_\phi) \\
\frac{dr}{dt} &= r^2 - 2mr + \frac{\alpha^2}{r^2} + 2mr + \alpha^2, \frac{d\phi}{dt} = \frac{2\alpha}{r^2} + 2mr + \alpha^2
\end{aligned}$$

Los agujeros negros cuánticos, suponen tratos de colisión, superposición o entrelazamiento, según sea el caso, en el que interactúan partículas o antipartículas deformantes y deformadas, en el primer caso, a propósito de su masa exponencial o de su energía potencial o de su energía cinética, según corresponda, y en el segundo caso, a propósito de su masa o energía cinética o potencial ligeras, según corresponda. Todo esto, depende esencialmente del campo cuántico de que se trate.

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APÉNDICE F.

Formalización matemática relativa a las características sistémicas inherentes a las partículas supermasivas y masivas como de las antipartículas supermasivas y masivas respectivamente, en un espacio cuántico curvo o deformado.



1. Equilibrio termodinámico.

$$d\mathcal{U} = TdS - pd\mathcal{V} + \sum_i \mu_i d\mathcal{N}_i$$

$$\mathcal{P} = -\frac{\partial u}{\partial v}\bigg|_{s,n} = -\frac{\partial f}{\partial v}\bigg|_{T,n}$$

$$\mathcal{T} = -\frac{\partial u}{\partial s}\bigg|_{v,n} = \frac{\partial \hbar}{\partial s}\bigg|_{p,n}$$

$$\mu = \frac{\partial u}{\partial n}\bigg|_{s,v} \quad n = -\frac{\partial \omega}{\partial \mu}\bigg|_{T,v}$$

$$\frac{\partial^2 u}{\partial s^2}\bigg|_v \geq 0, \frac{\partial^2 u}{\partial v^2}\bigg|_s \geq 0, \frac{\partial^2 u}{\partial s^2}\bigg|_v \frac{\partial^2 u}{\partial v^2}\bigg|_s \geq \left(\frac{\partial^2 u}{\partial s \partial v}\right)^2 \frac{\partial u}{\partial \mathcal{T}} \frac{\partial \mathcal{P}}{\partial n} \mathcal{T} \frac{\partial \mathcal{P}}{\partial \mathcal{T}} n^2 \frac{\partial u}{\partial n} \lambda$$

$$\begin{aligned} d\mathcal{U} &= TdS + SdT - \mathcal{P}d\mathcal{V} - \mathcal{V}dp + \mu d\mathcal{N} + \mathcal{N}d\mu \\ &= \frac{\partial \mathcal{U}}{\partial S} dS + SdT + \frac{\partial \mathcal{U}}{\partial \mathcal{V}} d\mathcal{V} - \mathcal{V}dp + \frac{\partial \mathcal{U}}{\partial \mathcal{N}} d\mathcal{N} + \mathcal{N}d\mu \sum_i \mu_i d\mu_i \end{aligned}$$

2. Equilibrio Químico.

$$\mu_i = \mathcal{B}_i \mu_B + \mathcal{Q}_i \mu_q + \mathcal{L}_i^{(e)} \mu_{l_e} + \mathcal{S}_i \mu_s, \mu_a = (\mathcal{N}_a + \mathcal{Z}_a) \mu_B + \mathcal{Z}_a \mu_q \equiv \mathcal{N}_a \mu_n + \mathcal{Z}_a \mu_p$$

3. Nivel de Rabi.

$$e_v = \sqrt{c^2 \mathcal{P}_z^2 + m_e^2 c^4 (1 + 2\nu B/\mathcal{B}_c)}$$

4. Corteza interior y núcleo.

$$\mathfrak{E}_{sym} = (n = n_n + n_p) = \frac{1}{2} \frac{\partial^2 \mathfrak{E}(n, \alpha)}{\partial \alpha^2} \bigg|_{\alpha=0} \quad \mathfrak{E}(n, \alpha) = \mathfrak{E}(n, 0) + \mathfrak{E}_{sym}(n) \alpha^2 + \mathcal{O}(\alpha^4)$$

$$\mathfrak{E}_{sym} = (n = n_n + n_p) = \mathfrak{E}(n, \alpha = 1) - \mathfrak{E}(n, \alpha = 0)$$

$$\begin{aligned} \mathfrak{E}_{sym}(n) &= \mathfrak{E}_{sym}(0) + (d\mathfrak{E}_{sym}(n)/dn)_{n_s} (n - n_s) + \frac{1}{2} (d^2 \mathfrak{E}_{sym}(n)/dn^2)_{n_s} (n - n_s)^2 \\ &\quad + \frac{1}{6} (d^3 \mathfrak{E}_{sym}(n)/dn^3)_{n_s} (n - n_s)^3 \end{aligned}$$

$$\mathfrak{E}_{sym}(n) = \mathcal{S}_0 + \mathcal{L}\chi + \frac{1}{2} \mathcal{K}_{sym} \chi^2 + \frac{1}{6} \mathcal{Q}_{sym} \chi^3$$

$$\mathfrak{E}(n, 0) = \epsilon_0 + \frac{1}{2} \mathcal{K}_{sat} \chi^2 + \frac{1}{6} \mathcal{K}_{sat} \chi^3$$

$$\mathcal{S}_0 \equiv \mathfrak{E}_{sym}(n_s) \simeq \mathfrak{E}_{\mathcal{NM}}(n_s) - \epsilon_0 \mathcal{L} \simeq 3 \frac{\mathcal{P}_{\mathcal{NM}}}{n_s}$$



5. Gravedad relativista.

$$\frac{\mathcal{M}_{crit}}{\mathcal{M}_{TOV}} = 1 + \alpha^2 (j/j_{Kep})^2 + \alpha^4 (j/j_{Kep})^4 \mathcal{G}_{\mu\nu} \frac{16\pi\mathfrak{G}}{c^4} \mathcal{T}_{\mu\nu}$$

$$ds^2 = \mathcal{A}(r)dt^2 - \mathcal{B}(r)dr^2 - r^2 \sin^2 \theta d\phi^2 \mathcal{G}_{\mu\nu} = \mathcal{R}_{\mu\nu} - \frac{1}{2} \mathcal{R} g_{\mu\nu}$$

$$g_{00} = 1 + 2\mathcal{V} = \mathcal{A}(r)1 - \frac{2\mathcal{G}\mathcal{M}}{r} \mathcal{B}(r) \left(1 - \frac{2\mathcal{G}\mathcal{M}}{r}\right)^{-1}$$

$$ds^2 = -g_{tt}(r)dt^2 + \left(1 - \frac{2m(r)}{r}\right)^{-1} dr^2 + r^2(d\theta^2 + \sin^2 \theta d\phi^2)$$

$$\begin{aligned} & \frac{2m'}{r^2} 16\pi\rho \frac{2m'}{r^3} - \left(1 - \frac{2m}{r}\right) \frac{g'_{tt}}{g_{tt}} 1/r \\ & = -16\pi\rho \left(1 - \frac{2m}{r}\right) \left(\frac{g'_{tt}}{2g_{tt}} - \frac{1}{4} \left(\frac{g'_{tt}}{g_{tt}}\right)^2 + \frac{1}{2r} \frac{g'_{tt}}{g_{tt}}\right) + \left(\frac{m}{r^2} - \frac{m'}{r}\right) \left(\frac{1}{r} + \frac{g'_{tt}}{2g_{tt}}\right) \end{aligned}$$

6. Ecuaciones de estado de Tolman-Oppenheimer-Volkoff (TOV).

$$\mathcal{R}_{\mu\nu} = -\kappa \left(\mathcal{T}_{\mu\nu} - \frac{1}{2} \mathcal{T} g_{\mu\nu} \right)$$

$$\mathcal{T}_{\mu\nu} = -p g_{\mu\nu} + (\varepsilon + p) u_\mu u_\nu$$

$$\mathcal{R}_{00} = -\frac{\mathcal{A}''}{2\mathcal{B}} + \frac{\mathcal{A}'}{4\mathcal{B}} \left(\frac{\mathcal{A}'}{\mathcal{A}} + \frac{\mathcal{B}'}{\mathcal{B}} \right) - \frac{\mathcal{A}'}{r\mathcal{B}}$$

$$\mathcal{R}_{11} = -\frac{\mathcal{A}''}{2\mathcal{A}} + \frac{\mathcal{A}'}{4\mathcal{A}} \left(\frac{\mathcal{A}'}{\mathcal{A}} + \frac{\mathcal{B}'}{\mathcal{B}} \right) - \frac{\mathcal{B}'}{r\mathcal{B}}$$

$$\mathcal{R}_{22} = -\frac{1}{\mathcal{B}} - 1 \frac{r}{2\mathcal{B}} \left(\frac{\mathcal{A}'}{\mathcal{A}} - \frac{\mathcal{B}'}{\mathcal{B}} \right) \sin^2 \theta d\phi^2$$

$$\frac{d}{dr} \left[r \left(1 - \frac{1}{\mathcal{B}} \right) \right] = \kappa r^2 \varepsilon$$

$$\mathcal{B}(r) = \left[1 - \frac{2\mathcal{G}m(r)}{r} \right]^{-1}$$

$$m(r) = 4\pi \int_0^r \varepsilon(\bar{r}) \bar{r}^2 d\bar{r}$$

$$\frac{dp}{dr} = \frac{\varepsilon + p}{2\mathcal{A}} \frac{d\mathcal{A}}{dr} \frac{1}{\mathcal{B}} - 1 + \frac{r}{2\mathcal{B}} \left(\frac{\mathcal{A}'}{\mathcal{A}} - \frac{\mathcal{B}'}{\mathcal{B}} \right) = -\frac{1}{2} \kappa (\varepsilon - p) r^2$$

$$\frac{d\mathcal{P}}{dr} = -\frac{\mathcal{G}_{\mathcal{N}} (\varepsilon(r) + \mathcal{P}(r)) (\mathcal{M}(r) + 4\pi r^3 \mathcal{P}(r))}{r^2 \left(1 - \frac{2\mathcal{G}_{\mathcal{N}} \mathcal{M}(r)}{r} \right)} \frac{d\mathcal{M}}{dr} = 4\pi r^2 \varepsilon$$

$$\frac{d\mathcal{P}}{dr} = -(\mathcal{P} + \rho) \frac{m(r) + 4\pi r^3 \mathcal{P}}{r(r - 2m(r))}$$



$$\mathcal{P}(r) = \rho_0 \left(\frac{\left(1 - \frac{2\mathcal{G}\mathcal{M}}{\mathcal{R}}\right)^{1/2} - \left(1 - \frac{2\mathcal{G}\mathcal{M}r^2}{\mathcal{R}^3}\right)^{1/2}}{\left(1 - \frac{2\mathcal{G}\mathcal{M}r^2}{\mathcal{R}^3}\right)^{1/2} - 3\left(1 - \frac{2\mathcal{G}\mathcal{M}}{\mathcal{R}}\right)^{1/2}} \right)$$

$$\begin{aligned} \frac{d\mathcal{P}}{dr} &= -(\mathcal{P} + \rho_0) \frac{\left(\frac{4\pi}{3}\right) \rho_0 r^3 + 4\pi r^3 \mathcal{P}}{r \left(r - 2\left(\frac{4\pi}{3}\right) \rho_0 r^3\right)} \\ &= -\frac{4\pi}{3} (\mathcal{P} + \rho_0) \frac{(\rho_0 + 3\mathcal{P}r)}{1 - \frac{16\pi\rho_0\mathcal{R}^2}{3}} \int_{\mathcal{P}}^0 \frac{d\mathcal{P}}{(\mathcal{P} + \rho_0)(3\mathcal{P} + \rho_0)} \int_r^{\mathcal{R}} \frac{rdr}{2\rho_0 r^2 - \frac{3}{4\pi} 2\rho_0} \log\left(\frac{\mathcal{P} + \rho_0}{3\mathcal{P} + \rho_0}\right) \\ &= \frac{1}{4\rho_0} \log\left(\frac{\frac{16\pi\rho_0\mathcal{R}^2}{3} - 1}{\frac{16\pi\rho_0 r^2}{3} - 1}\right) \frac{\mathcal{P} + \rho_0}{3\mathcal{P} + \rho_0} = \left(\frac{1 - 2\mathcal{G}\mathcal{M}/\mathcal{R}}{1 - 2\mathcal{G}\mathcal{M}r^2/\mathcal{R}^3}\right) \end{aligned}$$

$$\frac{d\mathcal{P}}{dr} = -(\mathcal{P} + \rho mc^2) \left[\frac{4\pi\mathcal{G}\mathcal{P}r^3 + \mathcal{G}m(r)c^2}{c^4 r \left(r - \frac{2\mathcal{G}m(r)}{c^2}\right)} \right]$$

$$\mathcal{P}(r) = \rho_0 c^2 \left[\frac{-\left(1 - \frac{2\mathcal{G}\mathcal{M}}{c^2\mathcal{R}}\right)^{\frac{1}{2}} + \left(1 - \frac{2\mathcal{G}\mathcal{M}r^2}{c^2\mathcal{R}^3}\right)^{\frac{1}{2}}}{\left(1 - \frac{2\mathcal{G}\mathcal{M}r^2}{c^2\mathcal{R}^3}\right)^{\frac{1}{2}} + 3\left(1 - \frac{2\mathcal{G}\mathcal{M}}{c^2\mathcal{R}}\right)^{\frac{1}{2}}} \right]$$

$$\mathcal{P}(0) = \rho_0 c^2 \left[\frac{-\left(1 - \frac{2\mathcal{G}\mathcal{M}}{c^2\mathcal{R}}\right)^{\frac{1}{2}} + 1}{-1 + 3\left(1 - \frac{2\mathcal{G}\mathcal{M}}{c^2\mathcal{R}}\right)^{\frac{1}{2}}} \right]$$

6.1. Ecuaciones TOV.

$$\frac{d\mathcal{P}}{dr} = -\left(\frac{\mathcal{G}\mathcal{M}_r}{r^2}\right) \frac{[\mathcal{P}(r) + \rho(r)][\mathcal{M}(r) + 4\pi r^3 \mathcal{P}(r)]}{r[r - 2\mathcal{M}(r)]}$$

$$\frac{d\mathcal{P}}{dr} = -\left(\frac{\mathcal{G}\mathcal{M}_r}{r^2}\right) \frac{\rho + (\mathcal{P}/c^2)(1 + (4\pi\mathcal{P}r^3/\mathcal{M}_r c^2))}{\left(1 - \left(\frac{2\mathcal{G}\mathcal{M}_r}{rc^2}\right)\right)}$$

$$\mathcal{M}_r = \frac{4\pi}{c^2} \int_0^r \rho c^2 r^2 dr$$

$$\xi = r\mathcal{A}; \rho = \rho_c e^{-\theta}; \mathcal{M}_r = \frac{4\pi\rho_c}{\mathcal{A}^3} v(\xi); \mathcal{A}^2 = \frac{4\pi\mathcal{G}\rho_c}{\sigma c^2}$$

$$\sigma = \frac{\mathcal{P}_c}{\rho c^2} \frac{(1 - 2\sigma v(\xi)/\xi)}{1 + \sigma} + \xi^2 \frac{d\theta}{d\xi} - v(\xi) - \sigma e^{-\theta} \xi^3 \frac{dv}{d\xi} = \xi^2 e^{-\theta} \frac{1}{\xi^2} \frac{d}{d\xi} \left(\xi^2 \frac{d\theta}{d\xi} \right) = e^{-\theta}$$



$$\xi^2 \frac{d\theta}{d\xi} - 2\sigma v \xi \frac{d\theta}{d\xi} - v - \sigma v - \sigma e^{-\theta} \xi^3 - \sigma^2 e^{-\theta} \xi^3$$

$$\theta(\xi) = \sum_{k=1}^{\infty} \alpha_k \xi^{2k}$$

$$\frac{d\theta}{d\xi} = \sum_{k=1}^{\infty} 2k \alpha_k \xi^{2k-1}$$

$$\xi \frac{d\theta}{d\xi} = \sum_{k=1}^{\infty} 2k \alpha_k \xi^{2k}$$

$$\xi^2 \frac{d\theta}{d\xi} = \sum_{k=1}^{\infty} 2k \alpha_k \xi^{2k+1}$$

$$e^{-\theta} = \sum_{k=0}^{\infty} \alpha_k \xi^{2k}$$

$$\alpha_k = -\frac{1}{\kappa} \sum_{i=1}^k i \alpha_i \alpha_{k-i}$$

$$\frac{dv}{d\xi} = \xi^2 \sum_{k=0}^{\infty} \alpha_k \xi^{2k} = \sum_{k=0}^{\infty} \alpha_k \xi^{2k+2}$$

$$v = \int \sum_{k=0}^{\infty} \alpha_k \xi^{2k+2} d\xi = \sum_{k=0}^{\infty} \frac{\alpha_k}{2\kappa + 3} \xi^{2k+3}$$

$$\xi \frac{d\theta}{d\xi} v = \left(\sum_{k=0}^{\infty} 2(k+1) \alpha_{k+1} \xi^{2k+2} \right) \left(\sum_{k=0}^{\infty} \frac{\alpha_k}{2\kappa + 3} \xi^{2k+3} \right)$$

$$\xi \frac{d\theta}{d\xi} v = \left(\sum_{k=0}^{\infty} f_k \xi^{2k+2} \right) \left(\sum_{k=0}^{\infty} g_k \xi^{2k+3} \right)$$

$$\xi \frac{d\theta}{d\xi} v = \left(\sum_{k=0}^{\infty} f_k \xi^{2k+2} \right) \left(\sum_{k=0}^{\infty} g_k \xi^{2k+3} \right) = \sum_{k=0}^{\infty} \gamma_k \xi^{2k+5}$$

$$\gamma_k = \sum_{i=0}^k f_i g_{k-i} = \sum_{i=1}^k i \alpha_i \alpha_{k-i}$$

$$\sum_{k=1}^{\infty} \left[2(k+1) \alpha_{k+1} - 2\sigma \gamma_{k-1} - \frac{\alpha_k}{2\kappa + 3} - \sigma \frac{\alpha_k}{2\kappa + 3} - \sigma \alpha_k - \sigma^2 \alpha_k \right] \xi^{2k+3}$$

$$2(k+1) \alpha_{k+1} - 2\sigma \gamma_{k-1} - \frac{\alpha_k}{2\kappa + 3} - \sigma \frac{\alpha_k}{2\kappa + 3} - \sigma \alpha_k - \sigma^2 \alpha_k = 0$$

$$\alpha_{k+1} = \frac{1}{2(k+1)} \left[2\sigma \gamma_{k-1} + \frac{\alpha_k}{2\kappa + 3} + \sigma \frac{\alpha_k}{2\kappa + 3} + \sigma \alpha_k + \sigma^2 \alpha_k \right]$$



$$\gamma_{k-1} = \sum_{i=0}^{k-1} f_i g_{k-i-1}; \alpha_k = -\frac{1}{k} \sum_{i=1}^k i \alpha_i \alpha_{k-i}; \forall k \geq 1; \alpha_0 = e^{(-\alpha_0)}$$

$$\alpha_1=\frac{1}{2}\left[\frac{1}{3}+\frac{\sigma}{3}+\sigma+\sigma^2\right]=\frac{1}{6}(1+\sigma)(1+3\sigma)$$

$$\bar{\sigma}(\bar{\psi}^i)(\bar{\mathcal{G}}_{\mu\nu}-\bar{\mathcal{W}}_{\mu\nu})=\kappa^2\bar{T}_{\mu\nu}$$

$$\bar{T}_{\mu\nu}=\overline{(\rho+p)}\overline{u_\mu u_\nu}+\overline{p\hbar_{\mu\nu}}\left(\frac{\overline{\Pi}}{\overline{\sigma}}\right)'=-\frac{\mathcal{M}}{\bar{r}^2}\left(\frac{\bar{Q}}{\bar{\sigma}}+\frac{\overline{\Pi}}{\bar{\sigma}}\right)\left(1+\frac{4\pi\bar{r}^3\frac{\overline{\Pi}}{\bar{\sigma}}}{\mathcal{M}}\right)\overline{\alpha(r)}-\frac{2\bar{\sigma}}{\kappa^2\bar{r}}\left(\frac{\bar{\mathcal{W}}_{00}}{\bar{r}^2}-\frac{\bar{\mathcal{W}}_{rr}}{\bar{\alpha}}\right)$$

$$\mathcal{M}(\bar{r})=\int\limits_0^{\bar{r}}4\pi\bar{r}^2\frac{\bar{Q}(\bar{r})}{\bar{\sigma}(\bar{r})}d\bar{r}$$

$$\overline{\alpha(r)}=\left(1-\frac{2\mathcal{M}(\bar{r})}{\bar{r}}\right)^{-1}$$

$$\overline{Q(r)}:=\overline{\rho(r)}-\frac{\bar{\sigma}(\bar{r})\bar{\mathcal{W}}_{tt}(\bar{r})}{\kappa^2c^2\overline{\rho(r)}}$$

$$\overline{\Pi(r)}:=\overline{\rho(r)}-\frac{\bar{\sigma}(\bar{r})\bar{\mathcal{W}}_{rr}(\bar{r})}{\kappa^2\overline{\alpha(r)}}$$

$$\mathcal{W}_{tt}=\frac{1}{2}\frac{\hat{b}}{\hat{a}}(\partial_{\hat{r}}\mathfrak{T})^2+\frac{\hat{b}}{2}\mathfrak{T}_2$$

$$\mathcal{W}_{\overline{r}\overline{r}}=\frac{1}{2}(\partial_{\hat{r}}\mathfrak{T})^2-\frac{\hat{\alpha}}{2}\mathfrak{T}_2$$

$$\mathcal{W}_{\theta\theta}=-\frac{1}{2}\frac{\hat{r}^2}{\hat{a}}(\partial_{\hat{r}}\mathfrak{T})^2-\frac{\hat{r}^2}{2}\mathfrak{T}_2$$

$$\mathcal{W}_{\phi\phi}=\sin^2\theta\mathcal{W}_{\theta\theta}$$

$$\hat{Q}=\hat{\rho}-\frac{(\partial_{\hat{r}}\mathfrak{T})^2}{2\kappa^2c^2\hat{\alpha}}-\frac{\mathfrak{T}_2}{2\kappa^2c^2}$$

$$\hat{\Pi}=\hat{\rho}-\frac{1}{2\kappa^2\hat{\alpha}}(\partial_{\hat{r}}\mathfrak{T})^2+\frac{1}{2\kappa^2}\mathfrak{T}_2$$

$$\frac{d\hat{\Pi}}{d\hat{r}}=\frac{\mathcal{GM}}{c^2\hat{r}^2}\left(c^2\hat{\rho}+\hat{p}-\frac{(\partial_{\hat{r}}\mathfrak{T})^2}{\kappa^2\hat{\alpha}}\right)\left(1+\frac{4\pi\hat{r}^3\hat{\Pi}}{c^2\mathcal{M}}\right)\times\left(1-\frac{2\mathcal{GM}}{c^2\hat{r}}\right)^{-1}+\frac{1}{\kappa^2\hat{\alpha}\hat{r}}(\partial_{\hat{r}}\mathfrak{T})^2$$

$$\mathcal{M}(\hat{r})=\int\limits_0^{\hat{r}}4\pi r^2\hat{Q}(r)dr$$

$$\frac{d\hat{p}}{d\hat{r}}=-\frac{\mathcal{GM}(\hat{r})}{c^2\hat{r}^2}(c^2\hat{\rho}+\hat{p})\left(1-\frac{2\mathcal{GM}(\hat{r})}{c^2\hat{r}}\right)^{-1}\times\left[1+\frac{4\pi\hat{r}^3}{c^2\mathcal{M}(\hat{r})}\left(\hat{p}-\frac{(\partial_{\hat{r}}\mathfrak{T})^2}{2\kappa^2\hat{\alpha}}+\frac{\mathfrak{T}_2}{2\kappa^2}\right)\right]+\hat{\mathcal{T}}\partial_{\hat{r}}\ln\mathfrak{T}_1\\-\frac{(\partial_{\hat{r}}\mathfrak{T})^2}{2\kappa^2\hat{\alpha}\hat{r}}$$



$$\mathcal{M}(\hat{r}) = \mathcal{M}_0(\hat{r}) - \eta(\hat{r}, \mathfrak{T}, \mathfrak{T}_2)$$

$$\eta(\hat{r}, \mathfrak{T}, \mathfrak{T}_2) = \frac{c^2}{4\mathcal{G}} \int_0^{\hat{r}} \chi^2 \left(\frac{(\partial_{\hat{r}} \mathfrak{T})^2}{\hat{a}} + \mathfrak{T}_2 \right) d\chi$$

$$\mathcal{M}(r) = \int_0^{\hat{r}} \frac{4\pi r^2}{\sqrt{\mathfrak{T}_1}} \left[\rho - \frac{1}{2\kappa^2 c^2} \frac{\mathfrak{T}_1}{a} (\partial_r \mathfrak{T})^2 - \frac{1}{2\kappa^2 c^2} \mathfrak{T}_1^2 \mathfrak{T}_2 \right] \times \left(\frac{r}{2} \partial_r \ln \mathfrak{T}_1 + 1 \right) dr$$

$$\frac{dp}{dr} = \left[-\frac{\mathcal{G}\mathcal{M}(r)}{c^2 r^2 \mathfrak{T}_1^{\frac{1}{2}}} (c^2 \rho + p) \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{c^2 r \mathfrak{T}_1^{\frac{1}{2}}} \right)^{-1} \times \left(1 + \frac{4\pi \mathfrak{T}_1^{\frac{3}{2}} r^3}{c^2 \mathcal{M}(r)} \left(\frac{\rho}{\mathfrak{T}_1^2} - \frac{(\partial_r \mathfrak{T})^2}{2\kappa^2 a \mathfrak{T}_1} + \frac{\mathfrak{T}_2}{2\kappa^2} \right) \right) \right. \\ \left. - \frac{\mathfrak{T}_1^2 (\partial_r \mathfrak{T})^2}{\kappa^2 a r} \right] \times \left(\frac{r}{2} \partial_r \ln \mathfrak{T}_1 + 1 \right) + (-c^2 \rho + 5p) \partial_r \ln \mathfrak{T}_1$$

$$\mathfrak{T}_1 = 1 + 4\beta \kappa^2 (c^2 \rho - 3p)$$

$$\mathfrak{T}_2 = \frac{4\beta \kappa^4 (c^2 \rho - 3p)^2}{(1 + 4\beta \kappa^2 (c^2 \rho - 3p))^2}$$

$$p' = \left[-\frac{\mathcal{G}\mathcal{M}(r)}{r^2 \mathfrak{T}_1^{\frac{1}{2}}} (\varepsilon + p) \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r \mathfrak{T}_1^{\frac{1}{2}}} \right)^{-1} \times \left(1 + \frac{4\pi \mathfrak{T}_1^{\frac{3}{2}} r^3}{\mathcal{M}(r)} \left(\frac{\rho}{\mathfrak{T}_1^2} + \frac{\mathfrak{T}_2}{2\kappa^2} \right) \right) \right] \times \left(\frac{r}{2} \partial_r \ln \mathfrak{T}_1 + 1 \right) \\ + (-\varepsilon + 5p) \partial_r \ln \mathfrak{T}_1$$

$$\mathfrak{T}_1 = 1 + 4\beta \kappa^2 (\varepsilon - 3p)$$

$$\mathfrak{T}_2 = \frac{4\beta \kappa^4 (\varepsilon - 3p)^2}{(1 + 4\beta \kappa^2 (\varepsilon - 3p))^2}$$

$$\mathcal{M}(r) = \int_0^r 4\pi \tilde{r}^2 \frac{\varepsilon - 2\beta \kappa^2 (\varepsilon - 3p)^2}{(1 + 2\beta \kappa^2 (\varepsilon - 3p))^{1/2}} \times \left[1 + \frac{\tilde{r}}{2} \partial_{\tilde{r}} \ln(1 + 4\beta \kappa^2 (\varepsilon - 3p)) \right] d\tilde{r}$$

$$\mathfrak{T}_1 = 1 + 2\alpha (\varepsilon - 3p)$$

$$\mathfrak{T}_2 = \frac{2\alpha \kappa^2 (\varepsilon - 3p)^2}{(1 + 2\alpha (\varepsilon - 3p))^2}$$

$$\alpha_{sing} = -\frac{1}{2(\varepsilon - 3\mathcal{P})} \frac{\partial \log \mathfrak{T}_1}{\partial r} \frac{2\alpha \left(\frac{1}{c_s^2} - 3 \right)}{\mathfrak{T}_1} \frac{d\mathcal{P}}{dr}$$

$$\frac{d\mathcal{P}}{dr} = \frac{uv\mathfrak{T}_1}{\mathfrak{T}_1 - \alpha \left(\frac{1}{c_s^2} - 3 \right) (uvr - 2\varepsilon + 10\mathcal{P})}$$



$$\begin{aligned}
u &= -\mathcal{M} \frac{\varepsilon + \mathcal{P}}{r^2 \mathfrak{T}_1^{\frac{1}{2}}} \left(1 - \frac{2\mathcal{M}}{r \mathfrak{T}_1^{\frac{1}{2}}} \right)^{-1} \\
v &= \frac{\mathcal{M} + 4\pi \mathfrak{T}_1^{\frac{3}{2}} r^3}{\mathcal{M}} \left(\frac{\mathcal{P}}{\mathfrak{T}_1^2} + \frac{\mathfrak{T}_2}{2\kappa^2} \right) \\
\frac{d\mathcal{M}(r)}{dr} &= 4\pi r^2 \frac{\varepsilon - \alpha(\varepsilon - 3\mathcal{P})^2}{\mathfrak{T}_1^{\frac{3}{2}}} \left(\mathfrak{T}_1 + r\alpha \left(\frac{1}{c_s^2} - 3 \right) \frac{d\mathcal{P}}{dr} \right) \\
\frac{dm(r)}{dr} &= \frac{4\pi r^2 \varepsilon(r)}{c^2} = 4\pi r^2 \rho(r) \frac{d\mathcal{P}(r)}{dr} = \frac{\mathcal{G}\varepsilon(r)m(r)}{c^2 r^2} \bigotimes \chi \\
\chi &= \left[1 + \frac{\mathcal{P}(r)}{\varepsilon(r)} \right] \left[1 + \frac{4\pi r^3 \mathcal{P}(r)}{m(r)c^2} \right] \left[1 - \frac{2\mathcal{G}m(r)}{c^2 r} \right]^{-1} \\
&\quad \left[1 + \frac{\mathcal{P}(r)}{\varepsilon(r)} \right] \left[1 + \frac{4\pi r^3 \mathcal{P}(r)}{m(r)c^2} \right] \left[1 - \frac{2\mathcal{G}m(r)}{c^2 r} \right]^{-1} = 1 \\
\frac{d\mathcal{P}}{dr} &= \frac{\mathcal{G}\varepsilon(r)m(r)}{c^2 r^2} \frac{dm(r)}{dr} = \frac{4\pi r^2 \varepsilon(r)}{c^2} \frac{d\mathcal{P}(r)}{dr} = \mathcal{R}_0/2r^2 \left(\frac{\mathcal{P}(r)}{k} \right)^{1/\gamma} \frac{d\bar{m}(r)}{dr} \\
&= \frac{4\pi r^2}{\mathcal{M}_\odot c^2} \left(\frac{\mathcal{P}(r)}{k} \right)^{1/\gamma}
\end{aligned}$$

7. Modelo Estelar.

$$\begin{aligned}
m'' &= -\frac{2m'}{r} - \frac{2mm''}{r} + \frac{4mm''}{r^2} + \frac{4\pi}{\gamma} \left((4\pi)^{\gamma-2} \left(\frac{m'}{r^2} \right)^{2-\gamma} + \frac{m'}{4\pi r^2} \right) \left(m + (4\pi)^{1-\gamma} r^3 \left(\frac{m'}{r^2} \right)^\gamma \right) m'' \\
&\quad - \frac{2m'}{r} \\
&\quad - \epsilon \left[\frac{2mm''}{r} - \frac{4mm''}{r^2} \right. \\
&\quad \left. - \frac{4\pi}{\gamma} \left((4\pi)^{\gamma-2} \left(\frac{m'}{r^2} \right)^{2-\gamma} + \frac{m'}{4\pi r^2} \right) \bigotimes \left(m + (4\pi)^{1-\gamma} r^3 \left(\frac{m'}{r^2} \right)^\gamma \right) \right]
\end{aligned}$$

8. Gravedad de Palatini.

$$\begin{aligned}
\mathcal{S} &= \frac{1}{2\kappa^2} \int d^4\chi \sqrt{-g} f(\mathcal{R}) + \mathcal{S}_m [g_{\mu\nu}, \chi] f'(\mathcal{R}) \mathcal{R}_{\mu\nu} - \frac{1}{2} f(\mathcal{R}) g_{\mu\nu} = \kappa^2 \mathcal{T}_{\mu\nu} \\
\mathcal{T}_{\mu\nu} &= -\frac{2}{\sqrt{-g}} \frac{\delta \mathcal{S}_m}{\delta g_{\mu\nu}} \nabla_\beta (\sqrt{-g} f'(\mathcal{R}) g^{\mu\nu}) \bar{g}^{\mu\nu} f'(\mathcal{R}) \mathcal{R}_{\mu\nu} \nabla_\beta (\sqrt{-g} \bar{g}^{\mu\nu}) \\
f'(\mathcal{R}) \mathcal{R} - 2f(\mathcal{R}) &= \kappa^2 \mathcal{T} \\
f(\mathcal{R}) &= \mathcal{R} + \beta \mathcal{R}^2
\end{aligned}$$

9. Parametrización gravitacional de Wagoner.

$$\mathcal{S}[\bar{g}^{\mu\nu}, \bar{\phi}, \chi] = \frac{1}{2\kappa^2} \int_{\Omega} d^4\chi \sqrt{-\bar{g}} [\bar{\mathcal{A}}(\bar{\phi})\bar{\mathcal{R}} - \bar{\mathcal{B}}(\bar{\phi})\bar{g}^{\mu\nu}\partial_{\mu}\bar{\phi}\partial_{\nu}\bar{\phi} - \mathcal{V}(\bar{\phi})] \\ + \mathcal{S}_{materia}[e^{2\bar{\alpha}(\bar{\phi})}\bar{g}^{\mu\nu}, \chi] \left\{ \begin{array}{l} \bar{g}^{\mu\nu} = e^{2\gamma(\bar{\phi})}\bar{g}^{\mu\nu} \\ \bar{\phi} = \bar{f}(\phi) \end{array} \right.$$

$$\bar{\mathcal{A}}(\bar{\phi}) = e^{2\bar{\gamma}(\bar{\phi})}\bar{\mathcal{A}}(\bar{f}(\bar{\phi})), \bar{\mathcal{B}}(\bar{\phi}) \\ = e^{2\bar{\gamma}(\bar{\phi})} \left(\left(\frac{d\bar{\phi}}{d\bar{\phi}} \right)^2 \bar{\mathcal{B}}(\bar{f}(\bar{\phi})) - 6 \left(\frac{d\bar{\gamma}}{d\bar{\phi}} \right)^2 \bar{\mathcal{A}}(\bar{f}(\bar{\phi})) - 6 \frac{d\bar{\gamma}}{d\bar{\phi}} \frac{d\bar{\mathcal{A}}}{d\bar{\phi}} \frac{d\bar{\phi}}{d\bar{\phi}} \right), \bar{\mathcal{V}}(\bar{\phi}) \\ = e^{4\bar{\gamma}(\bar{\phi})}\bar{\mathcal{V}}(\bar{f}(\bar{\phi})), \bar{\alpha}(\bar{\phi}) = \bar{\alpha}(\bar{f}(\bar{\phi})) + \bar{\gamma}(\bar{\phi})$$

$$\mathfrak{T}_1 = \bar{\mathcal{A}}/e^{2\bar{\alpha}}$$

$$\mathfrak{T}_2 = \bar{\mathcal{V}}/\bar{\mathcal{A}}^2$$

$$\frac{d\mathfrak{T}_3}{d\bar{\phi}} = \sqrt{\pm \frac{2\bar{\mathcal{A}}\bar{\mathcal{B}} + 3(\bar{\mathcal{A}}')^2}{2\bar{\mathcal{A}}^2}}$$

$$\mathcal{I}_i(\bar{\phi}(\chi)) = \mathcal{I}_i(\bar{f}(\phi(\chi))) = \mathcal{I}_i(\phi(\chi))\partial_{\mu} \left(\mathcal{I}_i(\bar{\phi}(\chi)) = \partial_{\mu} \left(\mathcal{I}_i(\bar{f}(\bar{\phi}(\chi))) \right) \right)$$

$$\bar{g}^{\mu\nu} = diag \left(-\bar{b}(\bar{r}), \bar{a}(\bar{r}), \bar{r}^2, \bar{r}^2 \sin^2 \bar{\theta} \right)$$

10. Parametrización einsteniana de gravitación.

$$\hat{g}^{\mu\nu} = \hat{\mathcal{A}}\bar{g}^{\mu\nu}\mathcal{S}[\hat{g}^{\mu\nu}, \mathfrak{T}, \chi] \\ = \frac{1}{2\kappa^2} \int_{\Omega} d^4\chi \sqrt{-\hat{g}} [\hat{\mathcal{R}} - \hat{g}^{\mu\nu}\partial_{\mu}\mathfrak{T}\partial_{\nu}\mathfrak{T} - \mathfrak{T}_2] + \mathcal{S}_{materia} \left[\frac{1}{\mathfrak{T}_1}\hat{g}^{\mu\nu}, \chi \right] \hat{\mathfrak{G}}_{\mu\nu} \\ + \frac{1}{2}\hat{g}_{\mu\nu}\hat{g}^{\alpha\beta}\partial_{\alpha}\mathfrak{T}\partial_{\beta}\mathfrak{T} + \frac{1}{2}\hat{g}_{\mu\nu}\mathfrak{T}_2 = \kappa^2\hat{\mathcal{T}}_{\mu\nu}\hat{\mathfrak{T}} - \frac{1}{2}\frac{d\mathfrak{T}_2}{d\mathfrak{T}} = \kappa^2\frac{1}{\mathfrak{T}_1}\frac{d\mathfrak{T}_1}{d\mathfrak{T}}\hat{\mathcal{T}}\hat{\mathfrak{T}}^{\mu}\hat{\mathcal{T}}_{\mu\nu} \\ = \frac{1}{2}\partial_{\nu}(\log \mathfrak{T}_1)\hat{\mathcal{T}}$$

11. Masa y Ondas gravitacionales en espacios cuánticos curvos.

$$\frac{(\mathcal{M}_{companion} \sin i)^3}{(\mathcal{M}_{NS} + \mathcal{M}_{companion})} = \frac{\mathcal{T} v_i^2}{2\pi\mathcal{G}} \dot{\mathcal{T}} \propto \frac{\mathcal{M}_{NS}\mathcal{M}_{companion}}{(\mathcal{M}_{NS} + \mathcal{M}_{companion})^{\frac{1}{3}}} \phi_{min} \propto (\mathcal{M}_{NS} + \mathcal{M}_{companion})^{\frac{2}{3}}$$

$$z = \frac{1}{\sqrt{1 - \frac{2\mathcal{G}\mathcal{M}_{NS}}{c^2\mathcal{R}_{NS}}}} - 1$$

$$\mathcal{M} = \left(\frac{(m_1 m_2)^3}{m_1 + m_2} \right)^{1/5} = \frac{c^3}{\mathcal{G}} \left(\frac{5\pi^{-\frac{8}{3}}}{96} f^{11/3} \frac{df}{dt} \right)^{3/5}$$



$$\Lambda = \frac{\lambda}{m^5} = \frac{2}{3} k_2 \frac{\mathcal{R}^5}{m^5} = \frac{2}{3} k_2 \mathcal{C}^{-5}$$

$$\bar{h}_{\alpha\beta}(t,\chi) = \frac{4\mathcal{G}}{c^4} \int d^3\gamma \frac{\mathcal{T}_{\alpha\beta}(t',\gamma)}{|\chi-\gamma|}$$

$$\tilde{\Lambda} = \frac{16}{3} \left[\frac{(m_1 + 12m_2)m_1^4\Lambda_1}{(m_1 + m_2)^5} + \frac{(m_2 + 12m_1)m_2^4\Lambda_2}{(m_1 + m_2)^5} \right]$$

$$\psi(f) = 2\pi f t_{coalescence} - \left(\phi_{coalescence} + \frac{\pi}{4} + \frac{3\mathcal{M}^2}{128m_1m_2v^5} \right) (\psi_{3.5\mathcal{PN}}^{pointlike} + \psi^{tidal})$$

12. Métrica computacional TOV y deformación.

$$ds^2 = -e^{2\phi(r)}[1+\mathcal{H}(r)\Gamma(\theta,\gamma)]dt^2 + e^{2\Lambda(r)}[1-\mathcal{H}(r)(\theta,\gamma)]dr^2 \\ + r^2[1-\mathcal{K}(r)\Gamma(\theta,\gamma)](d\theta^2 + \sin^2\theta d\phi^2)$$

$$\frac{d\mathcal{H}}{dr} = \beta$$

$$\frac{d\beta}{dr} = 2\left(1-2\frac{m_r}{r}\right)^{-1}\mathcal{H}\left\{-2\pi[5\epsilon+9\mathcal{P}+f(\epsilon+\mathcal{P})]+\frac{3}{r^2}+2\left(1-2\frac{m_r}{r}\right)^{-1}\left(\frac{m_r}{r^2}+4\pi r\mathcal{P}\right)^2\right\} \\ +\frac{2\beta}{r}\left(1-2\frac{m_r}{r}\right)^{-1}\left\{-1+\frac{m_r}{r}+2\pi r^2(\epsilon-\mathcal{P})\right\}$$

$$\kappa_2 = \frac{8\mathcal{C}^5}{5}(1-2\mathcal{C})^2[2+2\mathcal{C}(\gamma-1)-\gamma] \\ \times \{2\mathcal{C}[6+3\gamma+3\mathcal{C}(5\gamma-8)]+4\mathcal{C}^3[13-11\gamma+\mathcal{C}(3\gamma-2)+2\mathcal{C}^2(1+\gamma)] \\ +3(1-2\mathcal{C})^2[2-\gamma+2\mathcal{C}(\gamma-1)]\ln(1-2\mathcal{C})\}^{-1}$$

13. Cromodinámica cuántica en espacios curvos.

$$\mathcal{U}(\chi) = \exp\left(-i\sum_{\alpha=1}^8\theta^\alpha(\chi)\frac{\lambda_\alpha}{2}\right) = \exp\left(-i\theta^\alpha(\chi)\frac{\lambda_\alpha}{2}\right)2\delta^{\alpha\beta}if_{abc}\lambda_\alpha^\dagger\lambda_\alpha\lambda_b\frac{1}{4i}Tr(|\lambda_a\lambda_b|\lambda_c)$$

$$\mathcal{L}_{\mathcal{QCD}} = \bar{q}(i\mathcal{D}q - \mathcal{M})q - \frac{1}{4}\mathcal{G}_{\mu\nu}^a\mathcal{G}_a^{\mu\nu} + \mathcal{L}_{\mathcal{FP}}$$

$$\mathcal{G}_{\mu\nu}^a = \partial_\mu \mathcal{G}_\nu^a - \partial_\nu \mathcal{G}_\mu^a + g_s f^{abc} \mathcal{G}_\mu^b \mathcal{G}_\nu^c \frac{\lambda_a}{2}$$

$$q_f(\chi) \mapsto \mathcal{U}(\chi)q_f(\chi)$$

$$q_f^\dagger(\chi) \mapsto q_f^\dagger(\chi)\mathcal{U}^\dagger(\chi)$$

$$\mathcal{A}_\mu(\chi) \mapsto \mathcal{U}\mathcal{A}_\mu\mathcal{U}^\dagger + \frac{i}{g}\partial_\mu\mathcal{U}\mathcal{U}^\dagger$$

$$\mathcal{L}_{\mathcal{QCD}}^0 = \bar{q}i\gamma^\mu\mathcal{D}_\mu q - \frac{1}{4}\mathcal{G}_{\mu\nu,a}\mathcal{G}_a^{\mu\nu}$$

$$\mathcal{L}_{\mathcal{QCD}}^0 = \bar{q}_{\mathcal{R}}i\gamma^\mu\mathcal{D}_\mu q_{\mathcal{R}} + \bar{q}_{\mathcal{L}}i\gamma^\mu\mathcal{D}_\mu q_{\mathcal{L}} - \frac{1}{4}\mathcal{G}_{\mu\nu,a}\mathcal{G}_a^{\mu\nu}$$

$$\mathcal{V}_i^\mu = \mathcal{R}_i^\mu + \mathcal{L}_i^\mu = \bar{q}\gamma^\mu t_i q$$

$$\mathcal{A}_i^\mu = \mathcal{R}_i^\mu - \mathcal{L}_i^\mu = \bar{q}\gamma^\mu \gamma_5 t_i q$$



$$\mathcal{Q}_i^\mathcal{V} = \int d^3\chi \, \mathcal{V}_i^0 = \int d^3\chi \, q^\dagger(t,\vec{\chi}) \frac{\tau_i}{2} q(t,\vec{\chi}) \frac{d\mathcal{Q}_i^\mathcal{V}}{dt}$$

$$\mathcal{Q}_i^{\mathcal{A}} = \int d^3\chi \, \mathcal{A}_i^0 = \int d^3\chi \, q^\dagger(t,\vec{\chi}) \gamma_5 \frac{\tau_i}{2} q(t,\vec{\chi}) \frac{d\mathcal{Q}_i^{\mathcal{A}}}{dt}$$

$$\begin{aligned}\mathcal{M} &= \begin{pmatrix} m_u & 0 \\ 0 & m_d \end{pmatrix} = \frac{1}{2}(m_u+m_d) \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \frac{1}{2}(m_u-m_d) \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \\ &= \frac{1}{2}(m_u+m_d)\mathcal{I} + \frac{1}{2}(m_u-m_d)\tau_3\end{aligned}$$

$$\beta(g_s)=\mu\frac{\partial g_s}{\partial \mu}$$

$$\beta_{QCD}(g_s)=-\left(11-\frac{2\mathcal{N}_f}{3}\right)\frac{g_s^3}{(4\pi)^2}+\mathcal{O}(g_s^5)\frac{\partial \bar{g}_s(g_s,t)}{\partial t}=\beta_{QCD}(\bar{g}_s)$$

$$\alpha_s^{-1}(\Lambda)=\frac{33-2\mathcal{N}_f}{12\pi}\ln\left(\frac{\Lambda}{\Lambda_{QCD}}\right)$$

$$\Theta=\frac{\beta}{2g}\mathcal{F}_{\mu\nu}^{\alpha}\mathcal{F}_{\alpha}^{\mu\nu}+(1+\gamma_m)\sum_f\int m_f\bar{q}_fq_f\langle\mathcal{T}_\nu^{\mu}\rangle\mathcal{M}_{\mathcal{H}}\langle\Theta\rangle_{\mathcal{T},\mu_B}(\varepsilon-3\mathcal{P})\Delta\equiv\frac{\langle\Theta\rangle_{\mathcal{T},\mu_B}}{3\varepsilon}=\frac{1}{3}-\frac{\mathcal{P}}{\varepsilon}$$

$$\nu=-2+2\mathcal{A}-2\mathcal{C}+2\mathcal{L}+\sum_i\Delta_i\equiv d_i+\frac{n_i}{2}-2$$

$$\begin{aligned}\mathcal{V}(\vec{p}',\vec{p}) &= \mathcal{V}_\mathcal{C} + \tau_1 \bigotimes \tau_2 \mathcal{W}_\mathcal{C} + \left[\mathcal{V}_\mathcal{S} + \tau_1 \bigotimes \tau_2 \mathcal{W}_\mathcal{S}\right] \vec{\sigma}_1 \bigotimes \vec{\sigma}_2 \\ &\quad + \left[\mathcal{V}_{\mathcal{LS}} + \tau_1 \bigotimes \tau_2 \mathcal{W}_{\mathcal{LS}}\right] \left(-i\vec{\mathcal{S}} \boxtimes \left(\vec{q} \bigotimes \vec{k}\right)\right) \\ &\quad + \left[\mathcal{V}_\mathcal{T} + \tau_1 \bigotimes \tau_2 \mathcal{W}_\mathcal{T}\right] \left|\vec{\sigma}_1 \bigotimes \vec{q}\right| \left|\vec{\sigma}_2 \bigotimes \vec{q}\right| \\ &\quad + \left[\mathcal{V}_{\sigma\mathcal{L}} + \tau_1 \bigotimes \tau_2 \mathcal{W}_{\sigma\mathcal{L}}\right] \vec{\sigma}_1 \bigotimes \left(\vec{q} \bigotimes \vec{k}\right) \vec{\sigma}_2 \bigotimes \left(\vec{q} \bigotimes \vec{k}\right)\end{aligned}$$

$$\mathcal{T}(\vec{p}',\vec{p})=\mathcal{V}(\vec{p}',\vec{p})+\int\frac{d^3p''}{(2\pi)^3}\mathcal{V}(\vec{p}',\vec{p}'')\frac{\mathcal{M}_{\mathcal{N}}^2}{\mathfrak{E}_{p''}}\frac{1}{p^2-p''^2+i\epsilon}\mathcal{T}(\vec{p}'',\vec{p})$$

$$\hat{\mathcal{V}}(\vec{p}',\vec{p})\equiv\frac{1}{(2\pi)^3}\sqrt{\frac{\mathcal{M}_{\mathcal{N}}}{\mathfrak{E}_{p'}}}\mathcal{V}(\vec{p}',\vec{p})\sqrt{\frac{\mathcal{M}_{\mathcal{N}}}{\mathfrak{E}_p}}$$

$$\hat{\mathcal{T}}(\vec{p}',\vec{p})\equiv\frac{1}{(2\pi)^3}\sqrt{\frac{\mathcal{M}_{\mathcal{N}}}{\mathfrak{E}_{p'}}}\mathcal{T}(\vec{p}',\vec{p})\sqrt{\frac{\mathcal{M}_{\mathcal{N}}}{\mathfrak{E}_p}}$$

$$\hat{\mathcal{T}}(\vec{p}',\vec{p})=\hat{\mathcal{V}}(\vec{p}',\vec{p})+\int d^3p''\,\mathcal{V}(\vec{p}',\vec{p}'')\frac{\mathcal{M}_{\mathcal{N}}}{p^2-p''^2+i\epsilon}\mathcal{T}(\vec{p}'',\vec{p})$$

$$f(p',p)=\exp\left[-\left(\frac{p'}{\Lambda}\right)^{2n}-\left(\frac{p}{\Lambda}\right)^{2n}\right]$$

$$\chi(p)=\chi_{ref}(p)\sum_{n=0}^{\infty}c_n(p)\mathcal{Q}^n$$

$$\Delta\chi_n = \max(Q^5|\chi_{LO}|, Q^3|\chi_{LO} - \chi_{NLO}|, Q^2|\chi_{NLO} - \chi_{N^2LO}|, Q|\chi_{N^2LO} - \chi_{N^3LO}|)$$

14. Límite de Oppenheimer.

$$\mathcal{P}(0) = \rho_0 c^2 \frac{1 - \sqrt{1 - \frac{r_s}{\mathcal{R}}}}{3\sqrt{1 - \frac{r_s}{\mathcal{R}}} - 1}$$

15. Espacio tiempo cuántico curvo bajo la métrica de Kerr, Schwarzschild y Einstein.

$$ds^2 = \left(1 - \frac{rr_s}{r^2 + \alpha^2 \cos^2 \theta}\right) (d\mu + \alpha \sin^2 \theta d\varphi)^2 - 2(d\mu + \alpha \sin^2 \theta d\varphi)(dr + \alpha \sin^2 \theta d\varphi) - (r^2 + \alpha^2 \cos^2 \theta)(d\theta^2 + \sin^2 \theta d\varphi)$$

$$g_{\mu\nu} = \begin{pmatrix} 1 - \frac{rr_s}{\rho} - 1 & 0 & -\frac{rr_s \alpha \sin^2 \theta}{\rho} \\ 0 & 0 & -\alpha \sin^2 \theta \\ 0 & 0 & 1 \end{pmatrix}$$

$$g_{\mu\nu} = -\sin^2 \theta \left(r^2 + \alpha^2 + \frac{rr_s \alpha \sin^2 \theta}{\rho} \right)$$

$$ds^2 = \left(1 - \frac{r_s}{r}\right) d\mu^2 - 2d\mu dr - r^2(d\theta^2 + \sin^2 \theta d\varphi)$$

$$ds^2 = \left(1 - \frac{r_s}{r}\right) c^2 dt^2 - \frac{dr^2}{1 - \frac{r_s}{r}} - r^2(d\theta^2 + \sin^2 \theta d\varphi)$$

$$g_{\mu\nu} = \begin{pmatrix} 1 - \frac{rr_s}{\rho} & 0 & -\frac{rr_s \alpha \sin^2 \theta}{\rho} \\ 0 & \frac{\rho}{\Lambda} & 1 \\ 1 & 0 & -\sin^2 \theta \left(r^2 + \alpha^2 + \frac{rr_s \alpha \sin^2 \theta}{\rho} \right) \end{pmatrix}$$

$$ds^2 = \left(1 - \frac{rr_s}{r^2 + \alpha^2 \cos^2 \theta}\right) (dt)^2 + \left(\frac{2rr_s}{r^2 + \alpha^2 \cos^2 \theta}\right) dt d\varphi - \left(\frac{r^2 + \alpha^2 \cos^2 \theta}{r^2 - rr_s + \alpha^2}\right) dr^2 - \left(r^2 + \alpha^2 + \frac{rr_s \alpha^2 \sin^2 \theta}{r^2 + \alpha^2 \cos^2 \theta}\right) \sin^2 \theta d\varphi^2 - (r^2 + \alpha^2 \cos^2 \theta)(d\theta)^2$$

$$\mathcal{R}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \mathcal{R}_0 - \Lambda g_{\mu\nu} = -\kappa \mathcal{T}_{\mu\nu}$$

$$\Gamma_{\mu\nu}^{\lambda} = \frac{1}{2} g^{\lambda\kappa} \left(\frac{\partial g_{\kappa\mu}}{\partial \chi^{\nu}} + \frac{\partial g_{\kappa\nu}}{\partial \chi^{\mu}} - \frac{\partial g_{\mu\nu}}{\partial \chi^{\kappa}} \right)$$

$$\mathcal{R}_{\mu\nu} = \frac{\partial \Gamma_{\mu\rho}^{\rho}}{\partial \chi^{\nu}} - \frac{\partial \Gamma_{\mu\nu}^{\rho}}{\partial \chi^{\rho}} + \Gamma_{\mu\rho}^{\sigma} \Gamma_{\sigma\nu}^{\rho} - \Gamma_{\mu\nu}^{\sigma} \Gamma_{\sigma\rho}^{\rho}$$

$$\mathcal{T}_{\mu\nu} = \left(\rho + \frac{\mathcal{P}}{c^2} \right) u_{\mu} u_{\nu} - \mathcal{P} g_{\mu\nu}$$



$$u_0 = \left(\mathcal{A}_0(r_1, \theta) \sqrt{\left(\frac{r_1 \mathcal{M}_1(r_1)^3}{\omega^2 \mathcal{J}_1(r_1)^2 \cos^2(\theta) + r_1^2 \mathcal{M}_1(r_1)^2} - 1 \right)^2} - \frac{r_1 \omega^4 \mathcal{J}_1(r_1)^2 \mathcal{M}_1(r_1) \cos^4(\theta) \mathcal{A}_3(r_1, \theta)}{\omega^2 \mathcal{J}_1(r_1)^2 \cos^2(\theta) + r_1^2 \mathcal{M}_1(r_1)^2} \right. \\ \left. - \frac{\omega^4 \mathcal{J}_1(r_1)^2 \sin^2(\theta) \mathcal{A}_3(r_1, \theta)}{\mathcal{M}_1(r_1)^2} - r_1^2 \omega^2 \sin^2(\theta) \mathcal{A}_3(r_1, \theta) \right. \\ \left. + \frac{2\omega^2 \sqrt{\mathcal{J}_1(r_1)^2 \mathcal{M}_1(r_1)} \sqrt{\mathcal{M}_1(r_1)^2 \sin^2(\theta) \mathcal{A}_4(r_1, \theta)}}{\omega^2 \mathcal{J}_1(r_1)^2 \cos^2(\theta) + r_1^2 \mathcal{M}_1(r_1)^2} \right)^{-1}$$

$$r_- = \frac{\mathcal{M}_0}{2} - \sqrt{\left(\frac{\mathcal{M}_0}{2}\right)^2 - \alpha^2}$$

$$r_+ = \frac{\mathcal{M}_0}{2} + \sqrt{\left(\frac{\mathcal{M}_0}{2}\right)^2 + \alpha^2}$$

16. Ecuaciones TOV para espacios cuánticos curvos.

$$\mathcal{P}'(r) = - \left(\frac{\mathcal{G}\mathcal{M}(r)\rho(r)}{r^2} \right) \left(1 + \frac{\mathcal{P}(r)}{\rho(r)c^2} \right) \left(1 + \frac{4\pi r^3 \mathcal{P}(r)}{\mathcal{M}(r)c^2} \right) \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{rc^2} \right)^{-1}$$

$$\mathcal{P}'(r) = - \left(\frac{c^2 r_s \rho(r)}{r^2} \right) \left(1 + \frac{\mathcal{P}(r)}{\rho(r)c^2} \right) \left(1 + \frac{4\pi r^3 \mathcal{P}(r)}{\mathcal{M}(r)c^2} \right) \left(1 - \frac{r_s \mathcal{M}(r)}{r \mathcal{M}_t} \right)^{-1}$$

$$\mathcal{P} = -\frac{\partial \mathfrak{E}}{\partial \mathcal{V}} = 16\pi\rho_0 \left(\frac{\chi_{\mathcal{F}}^3}{3} \sqrt{1 + \chi_{\mathcal{F}}^2} - f(\chi_{\mathcal{F}}) \right)$$

$$f(\chi_{\mathcal{F}}) = \int_0^{\chi_{\mathcal{F}}} d\chi \sqrt{1 + \chi^2}$$

$$n_1 = \frac{\mathcal{N}_{op}}{\mathcal{V}_1} = \frac{2}{\mathcal{V}_1} \int_0^\infty d\omega_1 \frac{\mathcal{D}_1(\omega_1)}{1 + \exp(\beta_1(\omega_1 - \mu_1))} = \frac{1}{2\pi^2} \int_0^\infty d\omega_1 \frac{\sqrt{\omega_1}}{1 + \exp(\beta_1(\omega_1 - \mu_1))}$$

$$\mu_1 = \varepsilon_{\mathcal{F}_1} - \frac{\pi^2}{12\beta_1^2 \varepsilon_{\mathcal{F}_1}} = \mu_1(n_1)$$

$$p_1 = (\beta_1, n_1) = \frac{4\pi^{3/2}}{3\pi^2} \int_0^\infty d\omega_1 \frac{\omega_1^{3/2}}{1 + \exp(\beta_1(\omega_1 - \mu_1(n_1)))}$$

$$\mathcal{V}_s \mathcal{W}(r, \mathcal{V}_0, r_0, dr_0) = \frac{\mathcal{V}_0}{1 + \exp\left(\frac{r - r_0}{dr_0}\right)}$$

17. Configuración estelar.

$$\sigma(\psi^i)(\mathcal{G}_{\mu\nu} - \mathcal{W}_{\mu\nu}) = \kappa \mathcal{T}_{\mu\nu}$$

$$\mathcal{T}_{\mu\nu}^{eff} = \frac{1}{\sigma} \mathcal{T}_{\mu\nu} + \frac{1}{\kappa} \mathcal{W}_{\mu\nu}$$

$$ds^2 = -\mathcal{B}(r)dt^2 + \mathcal{A}(r)dr^2 + r^2d\theta^2 + r^2\sin^2\theta d\phi^2$$



$$\mathcal{R}_{tt} = -\frac{\mathcal{B}''}{2\mathcal{A}} + \frac{\mathcal{B}'}{4\mathcal{A}} \left(\frac{\mathcal{A}'}{\mathcal{B}} + \frac{\mathcal{B}'}{\mathcal{B}} \right) - \frac{\mathcal{B}'}{r\mathcal{A}} = \frac{\kappa}{2\sigma} (\rho + 3p)\mathcal{B} + \mathcal{W}_{tt} + \frac{\mathcal{B}\mathcal{W}}{2}$$

$$\mathcal{R}_{rr} = \frac{\mathcal{B}''}{2\mathcal{A}} - \frac{\mathcal{B}'}{4\mathcal{B}} \left(\frac{\mathcal{A}'}{\mathcal{B}} + \frac{\mathcal{B}'}{\mathcal{B}} \right) - \frac{\mathcal{A}'}{r\mathcal{A}} = \frac{\kappa}{2\sigma} (\rho - p)\mathcal{A} + \mathcal{W}_{rr} - \frac{\mathcal{A}\mathcal{W}}{2}$$

$$\mathcal{R}_{\theta\theta} = -1 + \frac{r}{2\mathcal{A}} \left(-\frac{\mathcal{A}'}{\mathcal{B}} + \frac{\mathcal{B}'}{\mathcal{B}} \right) + \frac{1}{\mathcal{A}} = \frac{\kappa}{2\sigma} (\rho - p)r^2 + \mathcal{W}_{\theta\theta} - \frac{r^2\mathcal{W}}{2}$$

$$\left(\frac{r}{\mathcal{A}} \right)' = 1 + \kappa r^2 \frac{\rho(r)}{\sigma(r)} + r^2 \mathcal{B}^{-1}(r) \mathcal{W}_{tt}(r)$$

$$\mathcal{A}(r) = \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r} \right)^{-1}$$

$$\mathcal{M}(r) = \int_0^r \left(4\pi\tilde{r} \frac{\rho(\tilde{r})}{\sigma(\tilde{r})} - \frac{\tilde{r}^2 \mathcal{W}_{tt}(\tilde{r})}{2\mathcal{G}\mathcal{B}(\tilde{r})} \right) d\tilde{r} \left(\sigma^{-1} \nabla_\mu \mathcal{T}^{\mu\nu} - \sigma^{-2} \mathcal{T}^{\mu\nu} \nabla_\mu \sigma \right) + \frac{1}{\kappa} \nabla_\mu \mathcal{T}^{\mu\nu}$$

$$\kappa\sigma^{-1} = \left(p' + (p + \rho) \frac{\mathcal{B}'}{2\mathcal{B}} \right) - \kappa\rho \frac{\sigma'}{\sigma^2} - \frac{\mathcal{A}'}{\mathcal{A}^2} \mathcal{W}_{rr} + \mathcal{A}^{-1} \mathcal{W}'_{rr} + \frac{2\mathcal{W}_{rr}}{\mathcal{A}r} + \frac{\mathcal{B}'}{2\mathcal{B}} \left(\frac{\mathcal{W}_{rr}}{\mathcal{A}} + \frac{\mathcal{W}_{tt}}{\mathcal{B}} \right) - \frac{2\mathcal{W}_{\theta\theta}}{r^2}$$

$$\mathcal{Q}(r) := \rho(r) + \frac{\sigma(r)\mathcal{W}_{tt}(r)}{\kappa\mathcal{B}(r)}$$

$$\Pi(r) := p(r) + \frac{\sigma(r)\mathcal{W}_{rr}(r)}{\kappa\mathcal{A}(r)}$$

$$\frac{\mathcal{A}'}{\mathcal{A}} = \frac{1 - \mathcal{A}}{r} - \frac{\kappa\mathcal{A}r}{\sigma} \mathcal{Q}$$

$$\frac{\mathcal{B}'}{\mathcal{B}} = \frac{\mathcal{A} - 1}{r} - \frac{\kappa\mathcal{A}r}{\sigma} \Pi$$

$$\left(\frac{\Pi}{\sigma} \right)' = -\frac{\mathcal{G}\mathcal{M}}{r^2} \left(\frac{\mathcal{Q}}{\sigma} + \frac{\Pi}{\sigma} \right) \left(1 + \frac{4\pi r^3 \Pi}{\mathcal{M}\sigma} \right) \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r} \right)^{-1} + \frac{2\sigma}{\kappa r} \left(\frac{\mathcal{W}_{\theta\theta}}{r^2} - \frac{\mathcal{W}_{rr}}{\mathcal{A}} \right)$$

$$\mathcal{M}(r) = \int_0^r 4\pi\tilde{r}^2 \frac{\mathcal{Q}(\tilde{r})}{\sigma(\tilde{r})} d\tilde{r}$$

18. Gravedad Clase k – esencia en espacios cuánticos curvos.

$$\mathcal{S} = \frac{1}{2\kappa} \int d^4\chi \sqrt{g} \left(\mathcal{R} - \nabla_\mu \phi \nabla^\mu \phi - 2\mathcal{V}(\phi) \right) + \mathcal{S}_m[g_{\mu\nu}, \psi]$$

$$\mathcal{G}_{\mu\nu} + \frac{1}{2} g_{\mu\nu} \nabla_\alpha \phi \nabla^\alpha \phi - \nabla_\mu \phi \nabla_\nu \phi + g_{\mu\nu} \mathcal{V}(\phi) = \kappa \mathcal{T}_{\mu\nu} \mathcal{V}'(\phi) - \square \phi$$

$$\mathcal{W}_{\mu\nu} = -\frac{1}{2} g_{\mu\nu} \nabla_\alpha \phi \nabla^\alpha \phi + \nabla_\mu \phi \nabla_\nu \phi - g_{\mu\nu} \mathcal{V}(\phi)$$

$$\mathcal{W}_{tt} = \frac{1}{2} \mathcal{B} \nabla_\alpha \phi \nabla^\alpha \phi + \mathcal{B} \mathcal{V}(\phi) = \mathcal{B}(\mathcal{C} + 2\mathcal{V})$$

$$\mathcal{W}_{rr} = \mathcal{A}\mathcal{C}$$



$$\begin{aligned}
\mathcal{W}_{\theta\theta} &= -r^2(\mathcal{C} + 2\mathcal{V}) \\
\mathcal{Q}_k(r) &:= \rho(r) + \kappa^{-1}(\mathcal{C} + 2\mathcal{V}) \\
\Pi_k(r) &:= p(r) + \kappa^{-1}\mathcal{C} \\
u^\mu \left(\frac{\sigma}{\kappa} \mathcal{W}_{\mu;\nu}^\nu - np \nabla_\mu \left(\frac{1}{n} \right) - n \nabla_\mu \left(\frac{\rho}{n} \right) + \frac{\rho}{\sigma} \nabla_\mu \sigma \right) \\
\delta n(r) &= \frac{n(r)}{\rho(r) + p(r)} \delta \rho(r) \\
n^{(r)} &= n \frac{\rho'}{\rho + p} \\
\mathcal{N} &= \int_0^{\mathcal{R}} 4\pi r^2 \left[1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r} \right]^{-1/2} n(r) dr \\
\delta \mathcal{M} - \lambda \delta \mathcal{N} &= \int_0^\infty 4\pi r^2 \delta \mathcal{Q} dr - \lambda \int_0^\infty 4\pi r^2 \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r} \right)^{-\frac{1}{2}} \delta n(r) dr \\
&\quad - \lambda \mathcal{G} \int_0^\infty 4\pi r \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r} \right)^{-\frac{3}{2}} n(r) \delta \mathcal{M}(r) dr \\
\delta \mathcal{M} - \lambda \delta \mathcal{N} &= \int_0^\infty 4\pi r^2 \left[1 - \frac{\lambda n(r)}{\rho(r) + p(r)} \mathcal{A}^{\frac{1}{2}} - \lambda \mathcal{G} \int_r^\infty 4\pi \tilde{r} n(\tilde{r}) \mathcal{A}^{\frac{3}{2}} d\tilde{r} \right. \\
&\quad \left. - \lambda \mathcal{G} \kappa^{-1} \int_r^\infty 4\pi \tilde{r} \mathcal{A}^{\frac{1}{2}} \frac{n}{p + \rho} \phi'^2 d\tilde{r} \right] \delta \mathcal{Q}(r) dr \\
&\quad - \lambda \kappa^{-1} \int_0^\infty \partial^\nu \phi \left[4\pi r^2 \mathcal{A}^{\frac{1}{2}} \Gamma_{\mu\nu}^\mu \frac{n(r)}{p(r) + \rho(r)} \right. \\
&\quad \left. - \partial_\nu \left(4\pi r^2 \mathcal{A}^{\frac{1}{2}} \frac{n(r)}{p(r) + \rho(r)} \right) \right] \delta \phi dr \int_0^\infty \partial^\mu \left(4\pi r^2 \frac{n(r)}{p(r) + \rho(r)} \delta \phi \partial_\mu \phi \right) dr \\
\frac{1}{\lambda} &= \frac{n(r)}{p(r) + \rho(r)} \mathcal{A}^{\frac{1}{2}} + \mathcal{G} \int_r^\infty 4\pi \tilde{r} n(\tilde{r}) \mathcal{A}^{\frac{3}{2}} d\tilde{r} + \mathcal{G} \kappa^{-1} \int_r^\infty 4\pi \tilde{r} \mathcal{A}^{\frac{1}{2}} \frac{n(\tilde{r})}{p(\tilde{r}) + \rho(\tilde{r})} \phi'^2 d\tilde{r} \\
4\pi r^2 \mathcal{A}^{\frac{1}{2}} \Gamma_{\mu\nu}^\mu \frac{n(r)}{p(r) + \rho(r)} &= \partial_r \left(4\pi r^2 \mathcal{A}^{\frac{1}{2}} \frac{n(r)}{p(r) + \rho(r)} \right) \\
-4\pi \mathcal{G} r \mathcal{A} - \frac{p'}{(p + \rho)^2} \frac{\mathcal{G} \mathcal{A}}{p + \rho} \left(4\pi r \mathcal{Q}_k - \frac{\mathcal{M}}{r^2} \right) &- \frac{4\pi r \mathcal{G}}{\kappa} \frac{\phi'^2}{p + \rho} \frac{\mathcal{A} - 1}{r} \mathcal{A} \frac{2\mathcal{G}\mathcal{M}}{r^2} p + \rho \\
&= \Pi_k + \mathcal{Q}_k - 2\kappa^{-1}(\mathcal{C} + \mathcal{V}) \mathcal{A}^{-1} \phi'^2 = 2(\mathcal{C} + \mathcal{V}) \\
\Pi'_k &= p' + \kappa^{-1} \mathcal{C}' = p' + \kappa^{-1}(\mathcal{C} + \mathcal{V}) \left(\frac{\mathcal{A} - 1}{r} - \kappa \mathcal{A} r \Pi_k + \frac{4}{r} \right) \\
\Pi'_k &= -\frac{\mathcal{A} \mathcal{G} \mathcal{M}}{r^2} (\Pi_k + \mathcal{Q}_k) \left(1 + 4\pi r \frac{\Pi_k}{\mathcal{M}} \right) - 4 \frac{\mathcal{C} + \mathcal{V}}{\kappa r}
\end{aligned}$$

$$\Gamma_{\mu\nu}^{\mu} = \frac{2}{r} - \frac{1}{2}(\kappa \mathcal{A}r(\Pi_{\kappa} + \mathcal{Q}_{\kappa}))$$

19. Métrica gravitacional de Brans-Dicke para espacios cuánticos curvos.

$$\begin{aligned} \mathcal{S} &= \frac{1}{32\pi\mathcal{G}} \int d^4\chi \sqrt{-g} \left(\mathcal{R}\phi - \frac{\omega}{\phi} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) + \mathcal{S}_m \\ \mathcal{G}_{\mu\nu} &= \frac{16\pi}{\phi} \mathcal{T}_{\mu\nu} + \frac{\omega}{\phi^2} \left(\partial_{\mu}\phi \partial_{\nu}\phi - \frac{1}{2} g_{\mu\nu} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) + \frac{1}{\phi} (\nabla_{\mu} \nabla_{\nu} \phi - g_{\mu\nu} \square \phi) \square \phi = \frac{16\pi}{3+2\omega} \mathcal{T} \\ \mathcal{W}_{\mu\nu} &= \frac{\omega}{\phi^2} \left(\partial_{\mu}\phi \partial_{\nu}\phi - \frac{1}{2} g_{\mu\nu} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) + \frac{1}{\phi} (\nabla_{\mu} \nabla_{\nu} \phi - g_{\mu\nu} \square \phi) \\ \mathcal{W}_{\theta\theta} &= \frac{\omega}{\phi^2} \left(\partial_{\theta}\phi \partial_{\theta}\phi - \frac{1}{2} g_{\theta\theta} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) + \frac{1}{\phi} (\nabla_{\theta} \nabla_{\theta} \phi - g_{\theta\theta} \square \phi) \\ \mathcal{W}_{rr} &= \frac{\omega}{\phi^2} \left(\partial_r\phi \partial_r\phi - \frac{1}{2} g_{rr} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) + \frac{1}{\phi} (\nabla_r \nabla_r \phi - g_{rr} \square \phi) \\ \mathcal{S} &= \frac{1}{32\pi\mathcal{G}} \int d^4\chi \sqrt{-g} \left[(1 + \alpha\phi) \mathcal{R} - \frac{\omega}{2} \partial_{\alpha}\phi \partial^{\alpha}\phi - \mathcal{V}(\phi) \right] + \mathcal{S}_m \\ \mathcal{G}_{\mu\nu} &= 16\pi\mathcal{G}\mathcal{T}_{\mu\nu} - \alpha\phi\mathcal{L}_{\mu\nu} \\ \mathcal{L}_{\mu\nu} &= \frac{\omega}{2} \left(\partial_{\mu}\phi \partial_{\nu}\phi - \frac{1}{2} g_{\mu\nu} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) + \frac{\alpha}{\phi} (\nabla_{\mu} \nabla_{\nu} \phi - g_{\mu\nu} \square \phi) - \frac{1}{2} g_{\mu\nu} \mathcal{V}(\phi) \\ \mathcal{W}_{\theta\theta} &= \alpha\phi \left(\frac{1}{2} g_{\theta\theta} \mathcal{V}(\phi) - \frac{\omega}{2} \left(\partial_{\theta}\phi \partial_{\theta}\phi - \frac{1}{2} g_{\theta\theta} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) - \frac{\alpha}{\phi} (\nabla_{\theta} \nabla_{\theta} \phi - g_{\theta\theta} \square \phi) \right) \\ \mathcal{W}_{rr} &= \alpha\phi \left(\frac{1}{2} g_{rr} \mathcal{V}(\phi) - \frac{\omega}{2} \left(\partial_r\phi \partial_r\phi - \frac{1}{2} g_{rr} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) - \frac{\alpha}{\phi} (\nabla_r \nabla_r \phi - g_{rr} \square \phi) \right) \\ \mathcal{W}_{tt} &= \alpha\phi \left(\frac{1}{2} g_{tt} \mathcal{V}(\phi) - \frac{\omega}{2} \left(\partial_t\phi \partial_t\phi - \frac{1}{2} g_{tt} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) - \frac{\alpha}{\phi} (\nabla_t \nabla_t \phi - g_{tt} \square \phi) \right) \\ \Pi' &= -\frac{\mathcal{G}\mathcal{M}}{r^2} (\mathcal{Q} + \Pi) \left(1 + \frac{4\pi r^3 \Pi}{\mathcal{M}} \right) \left(1 - \frac{2\mathcal{G}\mathcal{M}(r)}{r} \right)^{-\frac{1}{2}} \\ &+ \frac{2\alpha}{8\pi\mathcal{G}r} \left[\frac{\left(\frac{1}{2} g_{\theta\theta} \mathcal{V}(\phi) - \frac{\omega}{\phi^2} \left(\partial_{\theta}\phi \partial_{\theta}\phi - \frac{1}{2} g_{\theta\theta} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) - \frac{\alpha}{\phi} (\nabla_{\theta} \nabla_{\theta} \phi - g_{\theta\theta} \square \phi) \right)}{r^2} \right] \\ &- \left[\frac{\left(\frac{1}{2} g_{rr} \mathcal{V}(\phi) - \frac{\omega}{\phi^2} \left(\partial_r\phi \partial_r\phi - \frac{1}{2} g_{rr} \partial_{\alpha}\phi \partial^{\alpha}\phi \right) - \frac{\alpha}{\phi} (\nabla_r \nabla_r \phi - g_{rr} \square \phi) \right)}{\mathcal{A}} \right] \\ \mathcal{G}_{\mu\nu} &= (1 + \alpha\phi) (\mathcal{T}_{\mu\nu}^{eff}) \end{aligned}$$

20. Sistemas Dinámicos TOV.

$$-rp'(r-2m)$$

$$= (p + \rho)(m + 4\pi r^3 p) \left\{ \begin{array}{l} 4\pi r \chi(\ln r) = m(r) \\ 4\pi \gamma(\ln r) = m'(r) = r^2 \rho(r) \end{array} \right\} \left\{ \begin{array}{l} \chi'(s) = -\chi(s) + \gamma(s) \\ \gamma'(s) = 2\gamma(s) - \frac{8\pi\gamma(s)}{1-8\pi\chi(s)} (\chi(s) + \gamma(s)) \end{array} \right.$$

$$\mathcal{L}(\chi, \gamma) = 2 + 16\pi(\gamma - 3\chi) - \log(128\pi\gamma(1 - 8\pi\chi)^3)$$

$$\frac{d}{d\mathcal{L}} \mathcal{L}(\chi(t) + \gamma(t)) = \chi'(t)(\chi(t) - 2) + \gamma'(t) - 2(d-2) \frac{\gamma'(t)}{\gamma(t)} = -(\chi(t) - 2)^2$$

$$\mathcal{L}(\chi, \gamma) \sim \frac{1}{2} \left(\chi - \frac{1}{16\pi} \right)^2 + \left(\gamma - \frac{1}{16\pi} \right)^2 \lim_{s \rightarrow -\infty} \chi(s) e^s \lim_{s \rightarrow -\infty} \chi(s) e^{-2s}$$

$$\rho_0 = \rho(0) = |\rho|_\infty = \lim_{s \rightarrow -\infty} \gamma(s) e^{-2s}$$

$$\lim_{s \rightarrow -\infty} \chi(s) e^{-2s} = \lim_{s \rightarrow -\infty} \mathcal{Q}'(s) e^{(1-d)s} = \lim_{r \rightarrow 0^+} r^{d-1} \rho(r) r^{1-d} = \rho(0)$$

$$\lim_{s \rightarrow -\infty} \frac{\chi(s)}{\gamma(s)} = \frac{1}{3}$$

$$\mathcal{N} = \lim_{s \rightarrow -\infty} \frac{\chi(s)}{\gamma(s)} = \lim_{s \rightarrow -\infty} \frac{\chi'(s)}{\gamma'(s)} = \lim_{s \rightarrow -\infty} \frac{1 - \frac{\chi(s)}{\gamma(s)}}{2 - \frac{8\pi\gamma(s) \left(\frac{\chi(s)}{\gamma(s)} + 1 \right)}{1 - 8\pi\chi(s)}} = \frac{1 - \mathcal{N}}{2}$$

$$\chi' \left(-2 + \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right) = (\gamma - \chi) \left(-2 + \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right)$$

$$\chi' \left(-2 + \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right) + \gamma' = -\chi \left(-2 + \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right)$$

$$\frac{\gamma'(\gamma - \mathcal{C})}{\gamma} = \left(2 - \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right) (\gamma - \mathcal{C})$$

$$3\chi' \left(-2 + \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right) + 4\gamma' - \frac{\mathcal{C}\gamma'}{\gamma} = (\mathcal{C} - \gamma - 3\chi) \left(-2 + \frac{8\pi(\chi + \gamma)}{1 - 8\pi\chi} \right)$$

$$3\chi' \left(-2 + \frac{8\pi(2\chi' + \chi')}{1 - 8\pi\chi} \right) + 4\gamma' - \frac{\mathcal{C}\gamma'}{\gamma} = \frac{(\mathcal{C} - \gamma - 3\chi)(-2 + 24\pi\chi + 8\pi\gamma)}{1 - 8\pi\chi}$$

$$\begin{aligned} -6\chi' + \frac{48\pi\chi\chi'}{(1-8\pi\chi)} + 4\gamma' - \frac{1}{4\pi} \frac{\gamma'}{\gamma} \\ = \frac{-(1-12\pi\chi-4\pi\gamma)^2}{2\pi(1-8\pi\chi)} (-48\pi\chi - 3\log(1-8\pi\chi) + 16\pi\gamma - \log\gamma)' \end{aligned}$$

$$\mathcal{L}(\chi, \gamma) = 2 + 16\pi(\gamma - 3\chi) - \log(128\pi\gamma(1 - 8\pi\chi)^3) \frac{\mathcal{GM}}{\mathcal{R}c^2} \frac{2\mathcal{GM}}{\mathcal{R}c^2}$$

$$\frac{\chi'}{\gamma'} = 3(1 - 4\omega/(1 - \omega)) \in [-3, 3]$$

$$2 - 32\pi\chi - \log(128\pi\chi(1 - 8\pi\chi)^3) = 1 - \log(2)$$



20.1. Función Lyapunov.

$$\begin{aligned}
 p &= \kappa \rho \\
 \left\{ \begin{array}{l} \chi'(s) = -\chi(s) + \gamma(s) \\ \gamma'(s) = 2\gamma(s) - \frac{1+\kappa}{2\kappa} \bigotimes \frac{\gamma(s)(\chi(s) + \kappa\gamma(s))}{1-\chi(s)} \end{array} \right. \\
 (\chi_\kappa, \gamma_\kappa) &= \left(\frac{4\kappa}{(1+\kappa)^2 + 4\kappa}, \frac{4\kappa}{(1+\kappa)^2 + 4\kappa} \right) \\
 \mathcal{V} &= 2\gamma - (5 + 1/\kappa)\chi - 2\chi_\kappa \log(\gamma(1-\chi)^{\delta_\kappa}) + \mathcal{C}_\kappa 8\kappa^2 \delta_\kappa = (5\kappa + 1)(\kappa + 1)^2 \\
 \mathcal{C}_\kappa &= (3 + 1/\kappa)\chi_\kappa + 2\chi_\kappa \log(\chi_\kappa(1-\chi_\kappa)^{\delta_\kappa}) \\
 \mathcal{V} &= \mathcal{C} + 2\gamma - \gamma\chi - \beta \log(\gamma(1-\gamma)^\delta) - (1+\kappa)\gamma^2 \\
 &\quad + \gamma(\delta\beta - \gamma + 4 + \beta(1+\kappa)/2 - \gamma\chi^2 + \chi(\gamma - \delta\beta + 2\beta + \beta(1+\kappa)/2) - 2\beta) \\
 (1-\chi)\mathcal{V}' &= -(5 + 1/\kappa)(\chi - \chi_\kappa)^2 - (1+\kappa)(\gamma - \gamma_\kappa)^2 \\
 \left\{ \begin{array}{l} v' = -v + w \\ w' = -\frac{(1+\kappa)\gamma}{\kappa(1-\chi)^2}v + \left(2 + \frac{1+\kappa}{2\kappa} \bigotimes \frac{\chi + 2\kappa\gamma}{1-\chi}\right)w \end{array} \right. \\
 p = p(4\pi\rho) = p(r^{-2}\gamma) &\left\{ \begin{array}{l} \chi'(s) = -\chi(s) + \gamma(s) \\ \gamma'(s) = 2\gamma(s) - \frac{4\pi(\gamma + r^2 p(\gamma r^{-2}))(\chi + r^2 p(\gamma r^{-2}))}{(1-8\pi\chi)p'(r^{-2}\gamma)} \end{array} \right. \\
 \rho &= \mathcal{C}p^{1/\Gamma} + \frac{p}{1-\Gamma} \\
 \mathcal{N} = \lim_{s \rightarrow -\infty} \frac{\chi(s)}{\gamma(s)} &= \lim_{s \rightarrow -\infty} \frac{\chi'(s)}{\gamma'(s)} = \lim_{s \rightarrow -\infty} \frac{1 - \frac{\chi(s)}{\gamma(s)}}{2 - \frac{4\pi\gamma(s) \left(1 + \frac{p(\gamma(s)r^{-2})}{\gamma(s)r^{-2}}\right) \left(\frac{\chi(s)}{\gamma(s)} + \frac{p(\gamma(s)r^{-2})}{\gamma(s)r^{-2}}\right)}{(1-8\pi\chi(s))p'(r^{-2}\gamma(s))}}
 \end{aligned}$$

20.2. Modelo Relativista Michie–King.

$$\begin{aligned}
 \rho &= \frac{2m}{\hbar^3} \int_0^\infty f_c(\hat{p}) \left(1 + \frac{\varepsilon(\hat{p})}{mc^2}\right) d^3\hat{p} \\
 f_c(\hat{p}) &= \left(\frac{1 - \exp\left(\frac{(\varepsilon(\hat{p}) - \varepsilon_c)}{\kappa\mathcal{T}}\right)}{\exp\left(\varepsilon(\hat{p}) - \frac{\mu}{\kappa\mathcal{T}}\right) + 1} \right)_+ \\
 p &= \frac{4}{3\hbar^3} \int_0^\infty f_c(\hat{p}) \varepsilon(\hat{p}) \frac{1 + \frac{\varepsilon(\hat{p})}{2mc^2}}{1 + \frac{\varepsilon(\hat{p})}{mc^2}} d^3\hat{p} \\
 \varepsilon(\hat{p}) &= \sqrt{c^2\hat{p}^2 + m^2c^4 - mc^2} \\
 p(\rho) &= \left((p_\infty p)^{-1} + (p_0 \rho)^{-\frac{7}{5}} \right)^{-1}
 \end{aligned}$$



$$\frac{(\varepsilon + mc^2)}{\mathcal{T}} \equiv \text{const} = mc^2/\mathcal{T}_{\mathcal{R}} \left(\frac{\varepsilon_c}{(mc^2)} + 1 \right) \frac{\mathcal{T}_{\mathcal{R}}}{\mathcal{T}}$$

$$\kappa = \frac{\kappa \mathcal{T}_{\mathcal{R}}}{(mc^2)}$$

$$\gamma = \frac{\varepsilon_c}{(\kappa \mathcal{T})} 1 - \gamma \kappa = \frac{\mathcal{T}_{\mathcal{R}}}{\mathcal{T}} \chi \kappa \mathcal{T} = \varepsilon(\hat{p})$$

$$\hat{p} = \frac{mv}{\sqrt{1 - |v|^2/c^2}}$$

$$1 + \frac{\varepsilon(\hat{p})}{(mc^2)} = 1 + \frac{\chi \kappa \mathcal{T}}{\mathcal{T}_{\mathcal{R}}} = 1 + \frac{\chi \kappa}{(1 - \gamma \kappa)}$$

$$\hat{p}^2 = m^2 c^2 \frac{\chi \kappa}{1 - \gamma \kappa} \left(\frac{\chi \kappa}{1 - \gamma \kappa} + 2 \right)$$

$$\frac{\kappa m^2 c^2}{1 - \gamma \kappa} \left(\frac{\chi \kappa}{(1 - \gamma \kappa)} + 1 \right) d\chi = \langle \hat{p} | d | \hat{p} \rangle$$

$$\rho = \frac{8\pi c^3 \kappa m^4}{\hbar^3 (1 - \kappa \gamma)} \int_0^\gamma \left(\frac{\kappa \chi}{(1 - \kappa \gamma)} + 1 \right)^2 \left(\left(\frac{\kappa \chi}{(1 - \kappa \gamma)} + 1 \right)^2 - 1 \right)^{1/2} \frac{1 - \exp(\chi - \gamma)}{1 + \chi \exp(\chi - \gamma)} d\chi$$

$$\rho = \frac{8\sqrt{2}\pi c^3 \kappa^{3/2} m^4}{\hbar^3 (1 - \kappa \gamma)} \int_0^\gamma \chi^{1/2} \left(\frac{\kappa \chi}{(1 - \kappa \gamma)} + 1 \right)^2 \left(\frac{\kappa \chi/2}{(1 - \kappa \gamma)} + 1 \right)^{1/2} \frac{1 - \exp(\chi - \gamma)}{1 + \chi \exp(\chi - \gamma)} d\chi$$

$$\rho(\gamma) \sim \frac{8\sqrt{2}\pi c^3 \kappa^{3/2}}{1 + \chi} \left(\frac{mc}{\hbar} \right)^2 \int_0^\gamma \sqrt{\chi} (\gamma - \chi) d\chi = \rho_0 \gamma^{5/2}$$

$$\rho_0 = \frac{32\sqrt{2}\pi m \kappa^{\frac{3}{2}}}{15(1 + \chi)} \left(\frac{mc}{\hbar} \right)^3$$

$$\rho(\gamma) \sim 8\pi c^3 m^4 \hbar^{-3} \int_0^1 w^3 \frac{1 - \exp((\omega - 1)/\kappa)}{1 + \chi \exp((\omega - 1)/\kappa)} d\omega$$

$$p = \frac{\eta^2 \sqrt{2} \kappa}{1 - \kappa \gamma} \int_0^\gamma \left(\frac{\kappa \chi}{(1 - \kappa \gamma)} \right)^{\frac{3}{2}} \left(\frac{\frac{\kappa \chi}{2}}{1 - \kappa \gamma} + 1 \right)^{\frac{3}{2}} \frac{1 - \exp(\chi - \gamma)}{1 + \chi \exp(\chi - \gamma)} d\chi$$

$$p = \eta \left(\frac{\kappa}{(1 - \kappa \gamma)} \right)^{\frac{5}{2}} \int_0^\gamma \chi^{\frac{3}{2}} \left(\frac{\kappa \chi}{(1 - \kappa \gamma)} + 2 \right)^{\frac{3}{2}} \frac{1 - \exp(\chi - \gamma)}{1 + \chi \exp(\chi - \gamma)} d\chi$$

$$\eta = \frac{8\pi m^4 c^5}{3\hbar^3}$$

$$p_0 = \frac{64\sqrt{2}\pi m^4 c^5 \kappa^{\frac{5}{2}}}{105(1 + \chi) \hbar^3}$$

$$p(\rho) \sim p_0 \left(\frac{\rho}{\rho_0} \right)^{\frac{7}{5}}$$



$$p(\gamma) \sim \frac{c^2 \rho_\infty}{3(1 - \kappa\gamma)^4}$$

$$\rho_\infty = 8\pi c^3 \kappa^4 m^4 \hbar^{-3} \int_0^{1/\kappa} \chi^3 \frac{1 - \exp(\chi - 1/\kappa)}{1 + \chi \exp(\chi - 1/\kappa)} d\chi$$

20.3. Distribución relativista Fermi–Dirac.

$$\rho = \frac{1}{\pi^2} \int_0^\infty \frac{\gamma^2 \sqrt{1 + \gamma^2}}{1 + e^{-\alpha} e^{\chi \sqrt{1 + \gamma^2}}} d\gamma$$

$$\rho = \frac{1}{3\pi^2} \int_0^\infty \frac{\gamma^4}{\sqrt{1 + \gamma^2} (1 + e^{-\alpha} e^{\chi \sqrt{1 + \gamma^2}})} d\gamma$$

$$p(\chi) \sim \frac{1}{3\pi\chi^4} \int_0^\infty z^3 (1 + e^{-\alpha} e^z)^{-1} dz$$

$$p(\chi) \sim \frac{1}{\pi\chi^4} \int_0^\infty z^3 (1 + e^{-\alpha} e^z)^{-1} dz$$

$$\pi^2 p(\chi) \sim \sqrt{2} \chi^{-3/2} e^{-\chi + \alpha} \int_0^\infty \omega^{1/2} e^{-\omega} d\omega$$

$$\pi^2 p(\chi) \sim 2\sqrt{2} \chi^{-5/2} e^{-\chi + \alpha} \int_0^\infty \omega^{1/2} e^{-\omega} d\omega$$

$$p(\rho) \sim \rho / \log(\rho)$$

$$p(\rho) = \rho / (3 - \log(\rho) (1 + \rho)^{-1})$$

20.4. Materia oscura relativista – estado puro.

$$\rho = \kappa \left(\sqrt{1 + \chi^2} (\chi + 2\chi^3) - \operatorname{arcsinh}(\chi) \right)$$

$$\rho = \frac{\kappa}{3 \left(\chi \sqrt{1 + \chi^2} (2\chi^2 - 3) + 3 \operatorname{arcsinh}(\chi) \right) 2\kappa\chi^4}$$

$$3\kappa^{-1} \sqrt{\chi^2 + 1} p' = 8\chi^4$$

$$3\kappa^{-1} \sqrt{(\chi^2 + 1)^3} p'' = 8\chi^3 (3\chi^2 + 4)$$

$$\frac{d}{d\rho} (p(\chi(\rho))) = \frac{p'(\chi(\rho))}{\rho'(\chi(\rho))} = \frac{1}{1 + \chi^{-2}(\rho)} \frac{3p}{\rho} = 1 - \frac{4}{2\chi^2 + 1} \geq 1$$

21. Velocidad y presión de una partícula supermasiva o masiva y de una antipartícula supermasiva o masiva, según sea el caso, bajo el modelo relativista.

$$v_e|_{\mathcal{R}} = \left[\frac{\mathfrak{E}_e^2}{c^5} \frac{m^4}{m_e^5 c^4 + \mathfrak{E}_e^2 / \hbar c} \right]$$



$$v_e|_{-\mathcal{R}} = \left[\frac{\mathfrak{E}_e^2}{c^5} \frac{m^4}{m_e^5 c^4 - \mathfrak{E}_e^2 / \hbar c} \right]$$

$$\mathcal{P}_e|_{\mathcal{NR}} = \frac{1}{3} \int_0^{p_{\mathcal{F}}} p v_e n_e(p) dp = \frac{8\pi}{15\hbar^3} \frac{p_{\mathcal{F}}^5}{m_e} (\rho \mathcal{N}_{\mathcal{A}} / \mu_e)^{5/3} \frac{2}{1 + \chi_{\mathcal{H}}} k \mathcal{N}_{\mathcal{R}}$$

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APÉNDICE G.

Modelo matemático relativo a: 1. Demostrar la dinámica temporal – espacial provocada por las partículas y antipartículas supermasivas e hiperpartículas en espacios cuánticos curvos; 2. Demostrar la dirección de atrás hacia adelante y viceversa en relación a la dimensión tiempo en sistemas cuánticos geoméricamente deformados o en curvatura (flecha de tiempo bilateral); y, 3. Teorizar la configuración sistémica y morfológica de las hiperpartículas en espacios cuánticos curvos.

1. Aproximación de Markov.

$$\begin{aligned}
 \mathbb{P}(\chi_\eta, t_\eta | \chi_1, t_1; \dots \chi_{\eta-1}, t_{\eta-1}) &= \mathbb{P}(\chi_\eta, t_\eta | \chi_{\eta-1}, t_{\eta-1}) \\
 \widehat{\mathfrak{E}}(t)(\hat{\rho}(0)) &= \hat{\rho}(t) \\
 \widehat{\mathfrak{E}}(t_1 + t_2) &= \widehat{\mathfrak{E}}(t_1)\widehat{\mathfrak{E}}(t_2) \\
 \frac{d\hat{\rho}}{dt}(t) &= \hat{\mathcal{L}}\hat{\rho}(t) \\
 \hat{\mathcal{H}} &= \hat{\mathcal{H}}_\delta + \hat{\mathcal{H}}_\beta + \hat{\mathcal{H}}_{\delta\beta} \\
 \hat{\mathcal{H}}_\delta &= \frac{\hat{\mathcal{P}}^2}{2\mathcal{M}} + \mathcal{V}(\hat{\mathcal{Q}}) \\
 \hat{\mathcal{H}}_\beta^{(\kappa)} &= \sum_{\kappa=1}^{\mathcal{N}} \left(\frac{\hat{\mathcal{P}}_\kappa^2}{2m_\kappa} + \frac{m_\kappa \omega_\kappa^2 \hat{\mathcal{Q}}_\kappa^2}{2} \right) \\
 \hat{\mathcal{H}}_{\delta\beta}^{(\kappa)} &= \sum_{\kappa=1}^{\mathcal{N}} \left(g_\kappa \hat{\mathcal{Q}}_\kappa \hat{\mathcal{Q}} + \frac{g_\kappa^2}{2m_\kappa \omega_\kappa^2} \hat{\mathcal{Q}}^2 \right) \\
 \frac{d^2 \hat{\mathcal{Q}}}{dt^2}(t) + \mathcal{V}'(\hat{\mathcal{Q}}(t)) + \frac{1}{\mathcal{M}} \int_0^t \kappa(t-t') \hat{\mathcal{P}}(t') dt' + \kappa(t) \hat{\mathcal{Q}}(0) &= \hat{f}(t) \\
 \kappa(t) &= \sum_{i=1}^{\mathcal{N}} \frac{g_\kappa^2}{m_\kappa \omega_\kappa^2} \cos(\omega_\kappa t) \\
 \hat{f}(t) &= \sum_{i=1}^{\mathcal{N}} \left(g_\kappa \hat{\mathcal{Q}}_\kappa(0) \cos(\omega_\kappa t) + \frac{g_\kappa \hat{\mathcal{P}}_\kappa}{m_\kappa \omega_\kappa} \sin(\omega_\kappa t) \right) \\
 \langle \hat{f}(t) \rangle &= \text{Tr}(\hat{f}(t) \rho_{t\hbar}) = 0 \\
 \langle \{\hat{f}(t), \hat{f}(t')\} \rangle &= \sum_{i=1}^{\mathcal{N}} \frac{\hbar g_\kappa^2}{m_\kappa \omega_\kappa} \coth\left(\frac{\hbar \omega_\kappa}{2\kappa_\beta \mathcal{T}}\right) \cos(\omega_\kappa(t-t'))
 \end{aligned}$$

2. Flecha de tiempo reversible.

$$\begin{aligned}
 \hat{\Theta} \hat{\mathcal{Q}} \hat{\Theta}^{-1} &= \hat{\mathcal{Q}}, \hat{\Theta} \hat{\mathcal{P}} \hat{\Theta}^{-1} = -\hat{\mathcal{P}} \\
 \hat{\mathcal{A}}_{\mathcal{R}}(t) &= \hat{\Theta} \hat{\mathcal{A}}(-t) \hat{\Theta}^{-1} \\
 \hat{\mathcal{Q}}_{\mathcal{R}}(t) &= \hat{\Theta} \hat{\mathcal{Q}}(-t) \hat{\Theta}^{-1} = \hat{\mathcal{Q}}(t), \hat{\mathcal{P}}_{\mathcal{R}}(t) = \hat{\Theta} \hat{\mathcal{P}}(-t) \hat{\Theta}^{-1} = \hat{\mathcal{P}}(t)
 \end{aligned}$$



$$\mathcal{M} \frac{d^2 \hat{Q}}{dt^2}(-t) + \mathcal{V}'(\hat{Q}(-t)) + \frac{1}{\mathcal{M}} \int_0^{-t} \kappa(-t-t') \hat{\mathcal{P}}(t') dt' + \kappa(-t) \hat{Q}(0) = \hat{f}(-t)$$

$$\frac{d\hat{\rho}_{\mathcal{R}}}{dt} = -\hat{\Theta} \frac{d\hat{\rho}}{dt} \hat{\Theta}^{-1} = -\hat{\Theta} \hat{\mathcal{L}} \hat{\Theta}^{-1} \hat{\Theta} \hat{\rho} \hat{\Theta}^{-1} = -\hat{\mathcal{L}}_{\mathcal{R}} \hat{\rho}_{\mathcal{R}}$$

$$\hat{\rho}(t) = \hat{\mathfrak{E}}(t) \hat{\rho}(0), \hat{\mathfrak{E}}(t) = \exp(-\hat{\mathcal{L}}t)$$

$$\hat{\Theta} \hat{\mathfrak{E}}(t) \hat{\Theta}^{-1} = \exp(-\hat{\Theta} \hat{\mathcal{L}} \hat{\Theta}^{-1} t) = \exp(-\hat{\mathcal{L}}_{\mathcal{R}} t)$$

$$\hat{\Theta} \hat{\mathfrak{E}}(t) \hat{\Theta}^{-1} = \exp(\hat{\mathcal{L}}t) = \hat{\mathfrak{E}}(t)^{-1}$$

$$\hat{\Theta} \hat{\mathcal{U}}(t) \hat{\Theta}^{-1} = \hat{\mathcal{U}}(-t)$$

$$\frac{d\hat{\rho}}{dt} = \hat{\mathcal{L}}(t) \hat{\rho}(t)$$

$$\frac{d\hat{\rho}_{\mathcal{R}}}{dt}(t) = -\hat{\Theta} \frac{d\hat{\rho}}{dt}(-t) \hat{\Theta}^{-1} = -\hat{\mathcal{L}}_{\mathcal{R}}(-t) \hat{\rho}_{\mathcal{R}}(t)$$

$$\hat{\rho}(t) = \hat{\mathfrak{E}}(t, 0) \hat{\rho}(0), \hat{\mathfrak{E}}(t_2, t_1) = \hat{\mathcal{T}} \exp\left(-\int_{t_1}^{t_2} \hat{\mathcal{L}}(t') dt'\right)$$

$$\hat{\Theta} \hat{\mathfrak{E}}(t_2, t_1) \hat{\Theta}^{-1} = \hat{\mathcal{T}} \exp\left(-\int_{t_1}^{t_2} \hat{\mathcal{L}}(t') dt'\right) = \hat{\mathfrak{E}}(t_2, t_1)^{-1}$$

3. Métrica de Langevin para espacios cuánticos curvos.

$$\int_0^{\infty} \kappa(t') dt' < \infty$$

$$\int_0^{\tau_{\beta}} \kappa(t') dt' \approx \int_0^{\infty} \kappa(t') dt'$$

$$\int_0^t \kappa(t-t') \hat{\mathcal{P}}(t) dt' = \int_0^t \kappa(t') \hat{\mathcal{P}}(t-t') dt' \approx \hat{\mathcal{P}}(t) \int_0^{\tau_{\beta}} \kappa(t') dt' = \hat{\mathcal{P}}(t) \int_0^{\infty} \kappa(t') dt'$$

$$\int_0^t \kappa(t') dt' = \text{sgn}(t) \hat{\mathcal{P}}(t) \int_0^{|t|} \kappa(t') dt'$$

$$\int_0^t \kappa(t-t') \hat{\mathcal{P}}(t) dt' \approx \text{sgn}(t) \hat{\mathcal{P}}(t) \int_0^{\infty} \kappa(t') dt'$$

$$\mathcal{M} \frac{d^2 \hat{Q}}{dt^2} + \mathcal{V}'(\hat{Q}(t)) + \text{sgn}(t) \gamma \hat{\mathcal{P}}(t) = \hat{f}(t)$$

$$\int_0^{\infty} \kappa(t') dt' = \mathcal{M} \gamma$$



$$\langle \{\hat{f}(t), \hat{f}(t')\} \rangle = \frac{\gamma \mathcal{M} \hbar}{\pi} \int_0^{\Lambda} \omega \cot\left(\frac{\hbar \omega}{2\kappa_{\beta} \mathcal{T}}\right) \cos(\omega(t - t')) d\omega$$

$$\langle \{\hat{f}(t), \hat{f}(t')\} \rangle = 2\gamma \mathcal{M} \kappa_{\beta} \mathcal{T} \delta(t - t')$$

4. Simetría temporo – espacial en espacios cuánticos curvos.

$$Tr_{\mathcal{S}} = (\hat{\mathcal{U}} \hat{\mu}(t)) = Tr_{\mathcal{S}} (\hat{\rho}_{\mathcal{S}} \hat{\mathcal{U}}(t))$$

$$\dot{\hat{\mu}}(t) = -\frac{i}{\hbar} [\hat{\mathcal{H}}_{\mathcal{S}}, \hat{\mu}(t)] + \frac{i}{2\hbar} [\{\gamma \operatorname{sgn}(t) \hat{\mathcal{P}}, \hat{\mu}(t)\}, \hat{\mathcal{Q}}] + \frac{i}{2\hbar} [\{\hat{f}(t), \hat{\mu}(t)\}, \hat{\mathcal{Q}}]$$

$$\hat{\rho}(t) = \langle \hat{\mu}(t) \rangle := Tr_{\mathcal{B}}(\hat{\mu}(t) \rho_{t\hbar})$$

$$\dot{\hat{\rho}}(t) = -\frac{i}{\hbar} [\hat{\mathcal{H}}_{\mathcal{S}}, \hat{\rho}(t)] + \frac{i \operatorname{sgn}(t)}{2\hbar} [\{\gamma \hat{\mathcal{P}}, \hat{\rho}(t)\}, \hat{\mathcal{Q}}] + \frac{\Gamma(t)}{\hbar^2} [\{\hat{\rho}(t), \hat{\mathcal{Q}}\}]$$

$$\Gamma(t) = \int_0^t \langle \{\hat{f}(t), \hat{f}(t')\} \rangle dt'$$

$$\int_0^t \langle \{\hat{f}(t), \hat{f}(t')\} \rangle dt' = \operatorname{sgn}(t) \int_0^{|t|} \langle \{\hat{f}(|t|), \hat{f}(t')\} \rangle dt'$$

$$\lim_{t \rightarrow \infty} \int_0^{|t|} \langle \{\hat{f}(|t|), \hat{f}(t')\} \rangle dt' = 2\gamma \mathcal{M} \kappa_{\beta} \mathcal{T}$$

$$\dot{\hat{\rho}}(t) = -\frac{i}{\hbar} [\hat{\mathcal{H}}_{\mathcal{S}}, \hat{\rho}(t)] + \frac{i \operatorname{sgn}(t)}{2\hbar} [\{\gamma \hat{\mathcal{P}}, \hat{\rho}(t)\}, \hat{\mathcal{Q}}] - \operatorname{sgn}(t) \frac{2\gamma \mathcal{M} \kappa_{\beta} \mathcal{T}}{\hbar^2} [\{\hat{\rho}(t), \hat{\mathcal{Q}}\}, \hat{\mathcal{Q}}]$$

$$\frac{d\hat{\rho}}{dt} = (i\hat{\mathcal{L}}_{\mathcal{H}} + \operatorname{sgn}(t)\hat{\mathcal{L}}_{\mathcal{D}})\hat{\rho}(t)$$

$$\hat{\mathcal{L}}_{\mathcal{H}}\hat{\rho}(t) = -\frac{1}{\hbar} [\hat{\mathcal{H}}_{\mathcal{S}}, \hat{\rho}(t)]$$

$$\hat{\mathcal{L}}_{\mathcal{D}}\hat{\rho}(t) = \frac{i}{2\hbar} [\{\gamma \hat{\mathcal{P}}, \hat{\rho}(t)\}, \hat{\mathcal{Q}}] - \frac{2\gamma \mathcal{M} \kappa_{\beta} \mathcal{T}}{\hbar^2} [\{\hat{\rho}(t), \hat{\mathcal{Q}}\}, \hat{\mathcal{Q}}]$$

$$\frac{d\hat{\rho}_{\mathcal{R}}}{dt}(t) = -\hat{\Theta} \frac{d\hat{\rho}}{dt}(t) \hat{\Theta}^{-1} = (i\hat{\mathcal{L}}_{\mathcal{H}} + \operatorname{sgn}(t)\hat{\mathcal{L}}_{\mathcal{D}})\hat{\rho}_{\mathcal{R}}(t)$$

$$\hat{\rho}(t) = \hat{\mathfrak{E}}(t)\hat{\rho}(0), \hat{\mathfrak{E}}(t) = \exp(i\hat{\mathcal{L}}_{\mathcal{H}}t + \hat{\mathcal{L}}_{\mathcal{D}}|t|)$$

5. Entropía en espacios cuánticos curvos.

$$\rho(\wp, q, t) = \frac{1}{\sqrt{\pi \mathcal{N}(t)}} \exp\left(\frac{-(q + \operatorname{sgn}(t)\mathcal{A}(t)\wp)^2}{\mathcal{N}(t)} - \mathfrak{B}(t)\wp^2\right)$$

$$\mathcal{N}(t) = \frac{m\kappa_{\beta} \mathcal{T}}{\hbar^2} (1 - e^{-2\gamma|t|}) + \frac{e^{-2\gamma|t|}}{\sigma^2}$$

$$\mathcal{A}(t) = \frac{i\hbar}{2\sigma^2 m\gamma} e^{-\gamma|t|} (1 - e^{-\gamma|t|}) - \frac{i\kappa_{\beta} \mathcal{T}}{2\hbar\gamma} (1 - e^{-\gamma|t|})^2$$



$$\mathfrak{B}(t) = \frac{\hbar^2}{4\sigma^2 m^2 \gamma^2} (1 - e^{-\gamma|t|})^2 + \frac{\sigma^2}{4} + \frac{\kappa_\beta \mathcal{T}}{m\gamma^2} (2\gamma|t| - 3 + 4e^{-\gamma|t|} - e^{-2\gamma|t|})$$

$$\psi(\chi,0) = \frac{1}{(\sigma^2\pi)^{1/4}} \exp\left(-\frac{\chi^2}{2\sigma^2}\right)$$

$$\mathcal{S}_{\nu\mathcal{N}}(\xi) = \frac{1-\xi}{2\xi} \log\left(\frac{1+\xi}{1-\xi}\right) - \log\left(\frac{2\xi}{1+\xi}\right)$$

6. Simetría Lindblad en espacios cuánticos curvos.

$$\widehat{\mathcal{H}}=\widehat{\mathcal{H}}_{\delta}+\widehat{\mathcal{H}}_{\beta}+\widehat{\mathcal{H}}_{\delta\beta}$$

$$\widehat{\mathcal{H}}_j(t)=e^{\frac{i}{\hbar}(\widehat{\mathcal{H}}_{\delta}+\widehat{\mathcal{H}}_{\beta})t}\widehat{\mathcal{H}}_{\delta\beta}e^{\frac{-i}{\hbar}(\widehat{\mathcal{H}}_{\delta}+\widehat{\mathcal{H}}_{\beta})t}$$

$$\frac{d}{dt}\widehat{\rho}_j(t)=-\frac{i}{\hbar}[\widehat{\mathcal{H}}_j(t),\widehat{\rho}_j(t)]$$

$$\frac{d}{dt}\widehat{\rho}_s(t)=\frac{1}{\hbar^2}\int\limits_0^tTr_{\mathfrak{B}}\left[\widehat{\mathcal{H}}_j(t),\left[\widehat{\mathcal{H}}_j(s),\widehat{\rho}_j(s)\right]\right]ds$$

$$\frac{d}{dt}\widehat{\rho}_s(t)=-\frac{1}{\hbar^2}\int\limits_0^tTr_{\mathfrak{B}}\left[\widehat{\mathcal{H}}_j(t),\left[\widehat{\mathcal{H}}_j(s),\widehat{\rho}_j(s)\bigotimes\widehat{\rho}_{\mathfrak{B}}\right]\right]ds$$

$$\widehat{\mathcal{H}}_j(t)=\sum_{\alpha,\omega}\widehat{\mathcal{A}}_{\alpha}(\omega)\bigotimes\widehat{\mathcal{B}}_{\alpha}(t)$$

$$\begin{aligned}\frac{d}{dt}\widehat{\rho}_s(t) = & -\frac{1}{\hbar^2}\sum_{\alpha,\omega,\beta}\Big[\Gamma_{\alpha\beta}(\omega,t)\widehat{\mathcal{A}}_{\alpha}^{\dagger}(\omega)\widehat{\mathcal{A}}_{\beta}(\omega)\widehat{\rho}_s(t) \\ & +\Gamma_{\beta\alpha}^{*}(\omega,t)\widehat{\rho}_s(t)\widehat{\mathcal{A}}_{\alpha}^{\dagger}(\omega)\widehat{\mathcal{A}}_{\beta}(\omega)\left(\Gamma_{\alpha\beta}(\omega,t)+\Gamma_{\beta\alpha}^{*}(\omega,t)\right)\widehat{\mathcal{A}}_{\beta}(\omega)\widehat{\rho}_s(t)\widehat{\mathcal{A}}_{\alpha}^{\dagger}(\omega)\Big]\end{aligned}$$

$$\Gamma_{\alpha\beta}(\omega,t)=\int\limits_0^te^{i\omega s}Tr\Big(\widehat{\mathcal{B}}_{\alpha}^{\dagger}(t)\widehat{\mathcal{B}}_{\beta}^{\dagger}(t-s)\widehat{\rho}_{\mathcal{B}}\Big)ds$$

$$\Gamma_{\alpha\beta}(\omega,t)+\Gamma_{\beta\alpha}^{*}(\omega,t)=sgn\left(t\right)\int\limits_{-|t|}^{|t|}e^{i\omega s}Tr\Big(\widehat{\mathcal{B}}_{\alpha}^{\dagger}(s)\widehat{\mathcal{B}}_{\beta}^{\dagger}(0)\widehat{\rho}_{\mathcal{B}}\Big)ds$$

$$\Gamma_{\alpha\beta}(\omega,t)+\Gamma_{\beta\alpha}^{*}(\omega,t)\approx sgn\left(t\right)\gamma_{\alpha\beta}(\omega)$$

$$\gamma_{\alpha\beta}(\omega)=\int\limits_{-\infty}^{\infty}e^{i\omega s}Tr\Big(\widehat{\mathcal{B}}_{\alpha}^{\dagger}(s)\widehat{\mathcal{B}}_{\beta}^{\dagger}(0)\widehat{\rho}_{\mathcal{B}}\Big)ds$$

$$\eta_{\alpha\beta}(\omega)=\frac{1}{2i}\Big(\Gamma_{\alpha\beta}(\omega,t)-\Gamma_{\beta\alpha}^{*}(\omega,t)\Big)$$

$$\frac{d}{dt}\widehat{\rho}_s(t)=-\frac{i}{\hbar}[\widehat{\mathcal{H}}_{\delta},\widehat{\rho}_s]+\frac{sgn\left(t\right)}{\hbar^2}\sum_{\alpha,\beta,\omega}\widehat{\mathfrak{D}}_{\alpha\beta}\left(\omega\right)\widehat{\rho}_s(t)$$



$$\begin{aligned}\hat{\mathcal{H}}_{\delta} &= \frac{1}{\hbar} \sum_{\alpha, \beta, \omega} \eta_{\alpha\beta}(\omega) \hat{\mathcal{A}}_{\alpha}^{\dagger}(\omega) \hat{\mathcal{A}}_{\beta}(\omega) \\ \hat{\mathcal{D}}_{\alpha\beta}(\omega) \hat{\rho}_{\delta}(t) &= \gamma_{\alpha\beta}(\omega) \left(\hat{\mathcal{A}}_{\beta}(\omega) \hat{\rho}_{\delta}(t) \hat{\mathcal{A}}_{\alpha}^{\dagger}(\omega) - \frac{1}{2} \{ \hat{\mathcal{A}}_{\alpha}^{\dagger}(\omega) \hat{\mathcal{A}}_{\beta}(\omega), \hat{\rho}_{\delta}(t) \} \right) \\ \hat{\Theta} \gamma_{\alpha\beta}(\omega) \hat{\Theta}^{-1} &= \int_{-\infty}^{\infty} e^{i\omega s} \text{Tr} \left(\hat{\Theta} \hat{\mathcal{B}}_{\alpha}^{\dagger}(s) \hat{\mathcal{B}}_{\beta}^{\dagger}(0) \hat{\Theta}^{-1} \hat{\rho}_{\delta} \right) ds = \gamma_{\alpha\beta}(\omega)\end{aligned}$$

7. Simetría de Pauli en espacios cuánticos curvos.

$$\begin{aligned}\frac{d}{dt} \hat{\rho}_{\eta}(t) &= \sum_{\eta' \neq \eta} \left(\mathcal{W}_{\eta, \eta'} \hat{\rho}_{\eta'}(t) - \mathcal{W}_{\eta', \eta} \hat{\rho}_{\eta}(t) \right) \\ \langle \eta | \hat{\mathcal{A}}_{\alpha}^{\dagger}(\omega) \hat{\mathcal{A}}_{\beta}(\omega) | \eta' \rangle &= \delta_{\eta, \eta'} \langle \eta | \hat{\mathcal{A}}_{\alpha} | m \rangle \langle m | \hat{\mathcal{A}}_{\beta} | \eta \rangle \\ \frac{d}{dt} \rho_{\eta}(t) &= \frac{\text{sgn}(t)}{\hbar^2} \sum_{\eta' \neq \eta} \left(\mathcal{W}_{\eta, \eta'} \hat{\rho}_{\eta'}(t) - \mathcal{W}_{\eta', \eta} \hat{\rho}_{\eta}(t) \right) \\ \mathcal{W}_{\eta', \eta} &= \sum_{\alpha, \beta} \gamma_{\alpha\beta}(\varepsilon_{\eta} - \varepsilon_{\eta'}) \langle \eta | \hat{\mathcal{A}}_{\alpha} | \eta' \rangle \langle \eta' | \hat{\mathcal{A}}_{\beta} | \eta \rangle \\ \hat{\mathcal{H}} &= \hat{\mathcal{H}}_0 + \lambda \hat{\mathcal{H}}_{\delta\beta} \\ \frac{d}{dt} \rho_{\varepsilon, \kappa}(t) &= \frac{\lambda^2}{\hbar^2} \sum_{\varepsilon', \kappa'} \int_0^t \Lambda_{\varepsilon, \kappa, \varepsilon', \kappa'}(t-s) \left(\rho_{\varepsilon', \kappa'}(s) - \rho_{\varepsilon, \kappa}(s) \right) ds \\ \Lambda_{\varepsilon, \kappa, \varepsilon', \kappa'}(t) &= 2 \left| \langle \varepsilon, \kappa | \hat{\mathcal{H}}_{\delta\beta} | \varepsilon', \kappa' \rangle \right|^2 \cos \left(\frac{\varepsilon - \varepsilon'}{\hbar} t \right) \\ \int_0^t \Lambda_{\varepsilon, \kappa, \varepsilon', \kappa'}(t-s) ds &= 2\hbar \frac{\left| \langle \varepsilon, \kappa | \hat{\mathcal{H}}_{\delta\beta} | \varepsilon', \kappa' \rangle \right|^2}{\varepsilon - \varepsilon'} \sin \left(\frac{\varepsilon - \varepsilon'}{\hbar} t \right) \\ \frac{\hbar}{\varepsilon - \varepsilon'} \sin \left(\frac{\varepsilon - \varepsilon'}{\hbar} t \right) &\approx \text{sgn}(t) \pi \hbar \delta(\varepsilon - \varepsilon') \\ \frac{d}{dt} \rho_{\varepsilon, \kappa}(t) &= \text{sgn}(t) \sum_{\kappa'} \left(\mathcal{W}_{\kappa, \kappa'}^{\varepsilon} \rho_{\varepsilon, \kappa'}(t) - \mathcal{W}_{\kappa', \kappa}^{\varepsilon} \rho_{\varepsilon, \kappa}(t) \right) \\ \mathcal{W}_{\kappa, \kappa'}^{\varepsilon} &= \frac{2\lambda^2}{\hbar} \left| \langle \varepsilon, \kappa | \hat{\mathcal{H}}_{\delta\beta} | \varepsilon, \kappa' \rangle \right|^2 \eta(\zeta)\end{aligned}$$

8. Simetría de quiebre en espacios cuánticos curvos.

$$\begin{aligned}\psi_{\mathcal{R}}(\chi, \alpha + t) &= \psi^*(\chi, \alpha - t) \\ \frac{d\mathcal{W}}{dt} &= -\frac{\wp}{\mathcal{M}} \frac{d\mathcal{W}}{d\chi} + \sum_{\eta=0}^{\infty} \frac{(-\hbar^2)^{\eta} \mathcal{V}^{(2\eta+1)}(\chi)}{(2\eta+1)! 2^{2\eta}} \frac{d^{2\eta+1}\mathcal{W}}{d\wp^{2\eta+1}} + \text{sgn}(t) \gamma \frac{d}{d\wp} (\wp \mathcal{W}) + \text{sgn}(t) 2\gamma \mathcal{M} \kappa_{\beta} \mathcal{T} \frac{d^2\mathcal{W}}{d\wp^2} \\ \frac{d\mathcal{W}}{dt} &= -\frac{\wp}{\mathcal{M}} \frac{d\mathcal{W}}{d\chi} + \frac{d\mathcal{V}}{d\chi} \frac{d\mathcal{W}}{d\wp} + \text{sgn}(t) \gamma \frac{d}{d\wp} (\wp \mathcal{W}) + \text{sgn}(t) 2\gamma \mathcal{M} \kappa_{\beta} \mathcal{T} \frac{d^2\mathcal{W}}{d\wp^2}\end{aligned}$$



9. Expansión por deformación en espacios cuánticos curvos.

$$\dot{\hat{\mu}}(t) = -\frac{i}{\hbar} [\hat{\mathcal{H}}_{\delta}, \hat{\mu}(t)] + \frac{i}{2\hbar} [\{\gamma \operatorname{sgn}(t) \hat{\mathcal{P}}, \hat{\mu}(t)\}, \hat{\mathcal{Q}}]$$

$$\hat{\rho}(t) = \langle \hat{\mu}(t) \rangle := \operatorname{Tr}_{\mathcal{B}}(\hat{\mu}(t) \hat{\rho}_{t\hbar})$$

$$\hat{\mathcal{A}}\hat{\mu}(t) = -\frac{i}{\hbar} [\hat{\mathcal{H}}_{\delta}, \hat{\mu}(t)] + \frac{i}{2\mathcal{M}\hbar} [\{\gamma \operatorname{sgn}(t) \hat{\mathcal{P}}, \hat{\mu}(t)\}, \hat{\mathcal{Q}}]$$

$$\hat{\mathfrak{B}}\hat{\mu}(t) = \frac{i}{\hbar} [\hat{\mu}(t), \hat{\mathcal{Q}}]$$

$$\hat{a}(t)\hat{\mu}(t) = \frac{1}{2} \{\hat{f}(t), \hat{\mu}(t)\}$$

$$\hat{\eta}(t) = \exp(-\hat{\mathcal{A}}t) \hat{\mu}(t)$$

$$\dot{\hat{\eta}}(t) = \hat{\mathfrak{B}}(t) \hat{a}(t) \hat{\eta}(t)$$

$$\frac{d}{dt} \langle \hat{\eta}(t) \rangle = \hat{\mathfrak{B}}(t) \langle \hat{a}(t) \hat{\mu}(0) \rangle + \int_0^t \hat{\mathfrak{B}}(t) \hat{\mathfrak{B}}(t') \langle \hat{a}(t) \hat{a}(t') \hat{\mu}(t') \rangle dt'$$

$$\frac{d\hat{\wp}}{dt}(t) = \hat{\mathcal{A}}\hat{\rho}(t) + \int_0^t \langle \hat{a}(t) \hat{a}(t') \hat{\mathfrak{S}} \rangle \hat{\mathfrak{B}}\hat{\mathfrak{B}}(t' - t) \hat{\rho}(t') dt'$$

$$\frac{d\hat{\wp}}{dt}(t) = \hat{\mathcal{A}}\hat{\rho}(t) - \int_0^t \langle \hat{a}(t) \hat{a}(t') \hat{\mathfrak{S}} \rangle dt' \frac{1}{\hbar} [[\hat{\rho}(t), \hat{\mathcal{Q}}] \hat{\mathcal{Q}}]$$

$$\dot{\hat{\rho}}(t) = -\frac{i}{\hbar} [\hat{\mathcal{H}}_{\delta}, \hat{\rho}(t)] + \frac{i}{2\hbar} [\{\gamma \operatorname{sgn}(t) \hat{\mathcal{P}}, \hat{\rho}(t)\}, \hat{\mathcal{Q}}] - \frac{\Gamma(t)}{\hbar^2} [[\hat{\rho}(t), \hat{\mathcal{Q}}] \hat{\mathcal{Q}}]$$

Para aclarar:

1. Las partículas supermasivas y antipartículas supermasivas, son aquellas cuya masa es superior y en consecuencia, a propósito de su colisión, entrelazamiento, superposición o ultramasificación, provocan agujeros negros cuánticos.
2. Las partículas masivas y antipartículas masivas, son aquellas, cuya masa suficiente, deforma el tejido del espacio – tiempo cuántico, sin que necesariamente provoque un agujero negro cuántico, a propósito de su colisión, entrelazamiento, superposición o ultramasificación.
3. Las hiperpartículas, son aquellas que se aproximan, alcanzan o superan la velocidad de la luz, en cuyo primer y segundo casos, deforman el tejido del espacio – tiempo cuántico, y en cuyo tercer caso, puede provocar agujeros negros cuánticos.

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Thomas Guff, Chintalpati Umashankar Shastry y Andrea Rocco, Emergence of opposing arrows of time in open quantum systems, Scientific Reports | (2025) 15:3658.



Apéndice H.

Supersimetría de Yang – Mills, supermembranas, supergravedad cuántica, superconductividad, dualidad holográfica y agujeros negros cuánticos para espacios cuánticos curvos (campos cuánticos relativistas).

1. Formalización matemática.

$$\begin{aligned}
 S_{\text{open}} &= S_{\text{brane}} + S_{\text{bulk}} + S_{\text{int}} \approx S_{\mathcal{N}=4} + S_{\text{bulk - flat}} + S_{\text{int}}, g_s N \ll 1 \\
 dS^2 &= H(r)^{-1/2}(-dt^2 + dx_3^2) + H(r)^{1/2}(dr^2 + r^2 d\Omega_5^2) \\
 H(r) &= 1 + \frac{NG_N}{r^4}, G_N = g_s \alpha' \\
 S_{\text{close}} &= S_{\text{bulk - throat}} + S_{\text{bulk - flat}} + S_{\text{int}}, g_s N \gg 1. \\
 S_{\mathcal{N}=4} &= S_{\text{strings } AdS_5 \times S^5} \\
 g_{YM}^2 &= g_s \left(\frac{R}{l_s}\right)^4 = 4\pi g_{YM}^2 N = 4\pi\lambda \\
 x_i &\rightarrow \Lambda x_i : x_i \rightarrow \lambda x_i, r \rightarrow r - R \log \lambda \\
 L_{CFT} &+ \int d^4x h \mathcal{O} \\
 e^{-W(h)} &= \langle e^{-\int h \mathcal{O}} \rangle_{QFT} \\
 \langle \mathcal{O} \dots \mathcal{O} \rangle_c &= (-1)^n \left[\frac{\delta^n W}{\delta h^n} \right]_{h=0} \\
 e^{W(h)} &= \langle e^{\int h \mathcal{O}} \rangle_{QFT} = \mathcal{Z}_{\text{strings at } AdS}[h(x, z = \text{edge}) = h_0] \\
 \mathcal{Z}_{\text{strings at } AdS}[h(x, z = \text{edge})] &\sim e^{S_{\text{supergravity}}[h_0]}, \lambda \gg 1 \\
 L_{CFT} &+ \int d^4x \sqrt{g}(g_{\mu\nu} T^{\mu\nu} + J_\mu A^\mu + \phi \text{Tr}(F_{\mu\nu} F^{\mu\nu}) + \dots) \\
 ds^2 &= R^2 \frac{dx^\mu dx_\mu + dz^2}{z^2} \equiv g_{AB} dx^A dx^B, A = 0, \dots, d, x^A = (z, x^\mu) \\
 S &\sim \int d^{d+1}x \sqrt{g}(g^{AB} \partial_A \phi \partial_B \phi + m^2 \phi^2) \\
 S &\sim \int_{\partial AdS} d^d x \sqrt{g} g^{zB} \phi \partial_B \phi + \int d^{d+1}x \sqrt{g} \phi (-\square + m^2) \phi \\
 \phi(z, x^\mu) &= e^{ik_\mu x^\mu} f_k(z), k_\mu x^\mu = -\omega t + \vec{k} \cdot \vec{x} \\
 0 &= \frac{1}{R^2} (z^2 k^2 - z^{d+1} \partial_z (z^{-d+1} \partial_z) + m^2 R^2) f_k(z), k^2 = \vec{k}^2 - \omega^2 \\
 m^2 R^2 &= \Delta(\Delta - d) \\
 \phi &\sim \phi_+ z^{\Delta_+} + \phi_- z^{\Delta_-} \\
 \int dz \sqrt{g} |\phi|^2 &< \infty \\
 \phi(z, x^\mu)|_{z=\epsilon} &\rightarrow \epsilon^{\Delta_-} \phi_-^{ren}(x^\mu) \\
 S_{\text{edge}} &\ni \int_{z=\epsilon} d^d x \sqrt{g_\epsilon} \phi(x, \epsilon) \mathcal{O}(x, \epsilon) = \int d^d x \left(\frac{R}{\epsilon}\right)^d \epsilon^{\Delta_-} \phi_-^{ren}(x) \mathcal{O}(x, \epsilon) \\
 \mathcal{O}(x, \epsilon) &\sim \epsilon^{d-\Delta_-} \mathcal{O}^{ren}(x) = \epsilon^{\Delta_+} \mathcal{O}^{ren}(x) \\
 f_k(z) &= A(k) z^{d/2} K_\nu(kz) + B(k) z^{d/2} I_\nu(kz) \\
 f_k(z) &= \frac{z^{d/2} K_\nu(kz)}{\epsilon^{d/2} K_\nu(k\epsilon)} \\
 \phi(x, z) &= \int d^d k e^{ikx} f_k(z) \phi(k, \epsilon) \\
 S(\phi) &= 2R^{d-1} \int d^d k \phi(k, \epsilon) \phi(-k, \epsilon) \epsilon^{-d+1} \partial_z \left(\frac{z^{d/2} K_\nu(kz)}{\epsilon^{d/2} K_\nu(k\epsilon)} \right)_{z=\epsilon}
 \end{aligned}$$



$$\begin{aligned}
& -\frac{(-1)^{v-1}2^{1-2v}}{\Gamma(v)^2}k^{2v}\ln k\epsilon.\epsilon^{2v-d}+\dots \\
\langle \mathcal{O}(x_1)\mathcal{O}(x_2)\rangle &= -\frac{\delta}{\delta\phi_-^{ren}(x_1)}\frac{\delta}{\delta\phi_-^{ren}(x_2)}S=\frac{2\Gamma(\Delta_+)}{\pi^{d/2}\Gamma(\Delta_+-d/2)}\frac{1}{(x_1-x_2)^{2\Delta_+}} \\
& z^{d/2}K_{\pm\nu}(iqz)\sim e^{\pm i qz} \\
U_{\mathcal{C}}(y,z) &= \mathcal{P}e^{i\int_{\mathcal{C}}dx^\mu A_\mu^aT^a} \\
\varphi_{\mathcal{C}}(y) &= U_{\mathcal{C}}(y,z)\varphi(z) \\
W_{\mathcal{C}}(\mathcal{R}) &= \text{Tr}\mathcal{P}e^{i\oint_{\mathcal{C}}dx^\mu A_\mu^aT^a_{\mathcal{R}}} \\
\langle W_{\mathcal{C}}(\square) \rangle &\sim e^{-TV(l)} \\
W_{\mathcal{C},\mathcal{C}_{int}}(\square) &= \text{Tr}\mathcal{P}e^{\oint (iA_\mu\dot{x}^\mu+\theta^IX^I(x^\mu)\sqrt{\dot{x}^2})d\tau} \\
\langle W_{\mathcal{C},\mathcal{C}_{int}}(\square) \rangle &= \mathcal{Z}_{\text{string}}[\mathcal{C},\mathcal{C}_{int},\square] \\
\mathcal{Z}_{\text{string}}[\mathcal{C},\mathcal{C}_{int},\square] &\cong e^{-S_{\text{string}}^{\text{on-shell}}[\mathcal{C},\mathcal{C}_{int}]} \\
&\tau=t,\sigma=x,z=z(x) \\
ds^2 &= \frac{R^2}{z^2}[-dt^2+dx^2(1+z'^2)]=\gamma_{ab}dx^adx^b \\
S_{NG} &= \frac{1}{2\pi\alpha'}\int_{-T/2}^{T/2}dt\int_{-l/2}^{l/2}dx\sqrt{-\det(\gamma)}=\sqrt{4\pi\lambda}T\int_{-l/2}^{l/2}dx\frac{\sqrt{1+z'^2}}{z^2} \\
\mathcal{H} &= \mathcal{L}-z'\frac{\partial\mathcal{L}}{\partial z'}=\frac{1}{z^2\sqrt{1+z'^2}}=\frac{1}{z_{\text{max}}^2} \\
z'^2+V_{ef}(z) &= 0, V_{ef}(z)=1-\left(\frac{z_{\text{max}}}{z}\right)^4 \\
l &= \int_{-l/2}^{l/2}dx=2\int_0^{z_{\text{max}}}dz\frac{dx}{dz}=2z_{\text{max}}\int_0^1dy\frac{y^2}{\sqrt{1-y^4}}=2z_{\text{max}}\frac{\sqrt{2\pi^3}}{\Gamma^2(1/4)} \\
S_{NG} &= 2\sqrt{4\pi\lambda}T\int_{\epsilon}^{z_{\text{max}}}dz\frac{z_{\text{max}}^2}{z^2\sqrt{z_{\text{max}}^4-z^4}} \\
m_q &= \sqrt{4\pi\lambda}T\int_{\epsilon}^{z_{\text{max}}}dz\frac{dz}{z^2} \\
\ln\langle W_{\mathcal{C}}(\square) \rangle &= S_{NG}-2m_q=-\frac{2\sqrt{4\pi\lambda}T}{z_{\text{max}}}\int_{\epsilon/z_{\text{max}}}^1\frac{dy}{y^2}\left(\frac{1}{\sqrt{1-y^4}}-1\right) \\
V(l) &= -\frac{\sqrt{\lambda}}{l}\frac{4\pi^2}{\Gamma^4(1/4)} \\
ds^2 &= \frac{1}{z^2}(dr^2+r^2d\phi^2+dz^2+dx_i^2) \\
z &= \sqrt{a^2-r^2}, 0\leq r\leq a, 0\leq\phi\leq 2\pi \\
\ln\langle W_{\mathcal{C}}(\square) \rangle &= -\sqrt{\lambda} \\
\ln\langle W_{\mathcal{C}}(S) \rangle &= 2N\left(\kappa\sqrt{1+\kappa^2}+\sinh^{-1}\kappa\right) \\
\ln\langle W_{\mathcal{C}}(A) \rangle &= -\frac{2N\sqrt{\lambda}}{3\pi}\sin^3\theta_k \\
\mathcal{Z}_{\text{gravity}}[\phi(\phi_0)] &= \left\langle e^{i\int_{\partial\mathcal{M}}d^d\mathbf{x}\phi_0(\mathbf{x})\mathcal{O}(\mathbf{x})}\right\rangle \\
ds^2 &= \frac{R^2}{z^2}(d\mathbf{x}^2+dz^2) \\
\phi(\mathbf{x},z) &= \int_{\partial\mathcal{M}}d\mathbf{y}\,K(\mathbf{x},z\,|\,\mathbf{y})\phi_0(\mathbf{y}) \\
\phi(\mathbf{x},z) &\sim z^{\Delta_{\pm}}\phi_0(\mathbf{x}), z\rightarrow 0
\end{aligned}$$

$$\begin{aligned}
\Delta_{\pm} &= \frac{d}{2} \pm \mu, \mu = \sqrt{\frac{d^2}{4} + m^2 R^2} \\
(\square - m^2)K(\mathbf{x}, z | \mathbf{y}) &= 0 \\
K(\mathbf{x}, z | \mathbf{y}) &\sim z^{\Delta_-} \delta(\mathbf{x} - \mathbf{y}), z \rightarrow 0 \\
K(\mathbf{x}, z | \mathbf{y}) &= \lim_{z' \rightarrow 0} \sqrt{-g} g^{z' z'} \partial_{z'} G(\mathbf{x}, z | \mathbf{y}, z') \\
\phi(\mathbf{x}, z) &= \int_{\partial \mathcal{M}} d\mathbf{y} K(\mathbf{x}, z | \mathbf{y}) \phi_0(\mathbf{y}) + \varphi(\mathbf{x}, z) \\
S &= -\frac{1}{2} \int d\mathbf{x} dz \sqrt{-g} (g^{\mu\nu} \partial_{\mu} \phi \partial_{\nu} \phi + m^2 \phi^2) \\
S[\phi_0] &= \frac{1}{2} \int d\mathbf{x} [\sqrt{-g} g^{zz} \phi(\mathbf{x}, z) \partial_z \phi(\mathbf{x}, z)]_{z=0} \\
S[\phi_0] &= \frac{1}{2} \int dy dy' \phi_0(\mathbf{y}) \Delta(\mathbf{y}, \mathbf{y}') \phi_0(\mathbf{y}') \\
\Delta(\mathbf{y}, \mathbf{y}') &= \int d\mathbf{x} [\sqrt{-g} g^{zz} K(\mathbf{x}, z | \mathbf{y}) \partial_z K(\mathbf{x}, z | \mathbf{y}')]_{z=0} \\
\Delta(\mathbf{y}, \mathbf{y}') &\sim [\sqrt{-g} g^{zz} \partial_z K(\mathbf{y}, z | \mathbf{y}')]_{z=0} \sim \lim_{z, z' \rightarrow 0} (\sqrt{-g} g^{zz}) (\sqrt{-g} g^{z' z'}) \frac{\partial^2}{\partial z \partial z'} G(\mathbf{y}, z | \mathbf{y}', z') \\
\langle \psi_f | \mathcal{O}(\mathbf{y}) \mathcal{O}(\mathbf{y}') | \psi_i \rangle &= -i \frac{\delta^2 S[\phi_0]}{\delta \phi_0(\mathbf{y}) \delta \phi_0(\mathbf{y}')} = -i \Delta^{\text{i.f}}(\mathbf{y}, \mathbf{y}') \\
ds^2 &= R^2 \left[-\frac{dt^2}{1-x^2} + \frac{dx^2}{(1-x^2)^2} + \frac{x^2}{1-x^2} d\Omega_{d-1}^2 \right] \\
\lim_{x \rightarrow x_{\epsilon}} K(t, \Omega, x | t', \Omega', x_{\epsilon}) &= \frac{\delta(t-t') \delta(\Omega-\Omega')}{\sqrt{g_{\Omega}}} \\
K(t, \Omega, x | t', \Omega', x_{\epsilon}) &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \sum_{lm} e^{-i\omega(t-t')} Y_{lm}(\Omega) Y_{lm}^*(\Omega') f_{l\omega}(x) \\
(1-x^2) \frac{d^2 f_{l\omega}}{dx^2} + \frac{d-1-x^2}{x} \frac{df_{l\omega}}{dx} + \left(\omega^2 - \frac{q^2}{x^2} - \frac{m^2 R^2}{1-x^2} \right) f_{l\omega} &= 0 \\
f_{l\omega}(x) &= \left(\frac{x^{-\frac{d}{2}+\nu+1} (1-x^2)^{\frac{1}{2}\Delta_+}}{(1-\epsilon)^{-\frac{d}{2}+\nu+1} ((2-\epsilon)\epsilon)^{\frac{1}{2}\Delta_+}} \right) \frac{{}_2F_1\left(\frac{1}{2}(\mu+\nu-\omega+1), \frac{1}{2}(\mu+\nu+\omega+1); \nu+1; x^2\right)}{{}_2F_1\left(\frac{1}{2}(\mu+\nu-\omega+1), \frac{1}{2}(\mu+\nu+\omega+1); \nu+1; (1-\epsilon)^2\right)} \\
f_{l\omega}(x) &\sim C_+(1-x)^{\frac{1}{2}\Delta_+} + C_-(1-x)^{\frac{1}{2}\Delta_-} \\
\omega_{nl} &= \pm(2n+\nu+\mu+1) \\
&= \pm(2n+l+\Delta_+), n, l = 0, 1, 2 \dots \\
\phi(t, \Omega, x) &= \int dt' d\Omega' \sqrt{g_{\Omega'}} K(t, \Omega, x | t', \Omega', x_{\epsilon}) \phi_0(t', \Omega') \\
\Delta_{\text{reg}}(t, \Omega | t', \Omega') &= -\frac{1}{\sqrt{g_{\Omega}}} [\sqrt{-g} g^{xx} \partial_x K(t, \Omega, x | t', \Omega', x_{\epsilon})]_{x=x_{\epsilon}} \\
&= -\int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \sum_{lm} Y_{lm}(\Omega) Y_{lm}^*(\Omega') \left[\frac{x^{d-1}}{(1-x^2)^{\frac{d-2}{2}}} \partial_x f_{l\omega}(x) \right]_{x=x_{\epsilon}} \\
&= -\sum_{lm} Y_{lm}(\Omega) Y_{lm}^*(\Omega') \int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \left[\frac{x^{d-1}}{(1-x^2)^{\frac{d-2}{2}}} \partial_x f_{l\omega}(x) \right]_{x=x_{\epsilon}} \\
\phi_0(t, \Omega) &= \epsilon^{\frac{1}{2}\Delta_-} \phi_{\text{ren}}(t, \Omega)
\end{aligned}$$

$$\begin{aligned}
\Delta_{\text{ren}}(t, \Omega | t', \Omega') &\equiv \lim_{\epsilon \rightarrow 0} \epsilon^{\Delta_-} \Delta_{\text{reg}}(t, \Omega | t', \Omega') \\
&= \sum_{lm} Y_{lm}(\Omega) Y_{lm}^*(\Omega') \int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \\
&\quad \times \frac{\Delta_+}{2^{\Delta_-}} \frac{\Gamma(1-\mu)}{\Gamma(1+\mu)} \frac{\Gamma\left(\frac{1}{2}(-\omega + \nu + \mu + 1)\right) \Gamma\left(\frac{1}{2}(\omega + \nu + \mu + 1)\right)}{\Gamma\left(\frac{1}{2}(-\omega + \nu - \mu + 1)\right) \Gamma\left(\frac{1}{2}(\omega + \nu - \mu + 1)\right)} \\
\Delta_{\text{ren}}^F(t, \Omega | t', \Omega') &= 2i \frac{\Delta_+ \Gamma(1-\mu)}{2^{\Delta_-} \Gamma(1+\mu)} \sum_{lm} Y_{lm}(\Omega) Y_{lm}^*(\Omega') \\
&\quad \times \left[\sum_{n=0}^{\infty} \frac{(-1)^n}{n!} \frac{\Gamma(n+l+\Delta_+)}{\Gamma\left(n+l+\frac{d}{2}\right) \Gamma(-(n+\mu))} e^{-i|t-t'|(2n+l+\Delta_+)} \right] \\
\langle 0 | T \mathcal{O}(t, \Omega) \mathcal{O}(t', \Omega') | 0 \rangle &= -i \Delta_{\text{ren}}^F(t, \Omega | t', \Omega') = -\frac{2\Delta_+}{2^{\Delta_-} \Gamma(\mu)} \sum_{lm} Y_{lm}(\Omega) Y_{lm}^*(\Omega') \frac{\Gamma(l+\Delta_+)}{\Gamma\left(l+\frac{d}{2}\right)} \\
&\quad \times e^{-i|t-t'|(l+\Delta_+)} {}_2F_1\left(1+\mu, l+\Delta_+, l+\frac{d}{2}; e^{-2i|t-t'|}\right) \\
\phi(\mathbf{y}, x) &= \int d\mathbf{y}' \mathbf{K}^i(\mathbf{y}, x | \mathbf{y}') \phi_0^i(\mathbf{y}') \\
&= \int d\mathbf{y}' [\mathbf{K}^+(\mathbf{y}, x | \mathbf{y}') \phi_0^+(\mathbf{y}') + \mathbf{K}^-(\mathbf{y}, x | \mathbf{y}') \phi_0^-(\mathbf{y}')] \\
\mathbf{K}^+(\mathbf{y}, x | \mathbf{y}')|_{x=x_+} &= \delta(\mathbf{y} - \mathbf{y}'), \\
\mathbf{K}^+(\mathbf{y}, x | \mathbf{y}')|_{x=x_-} &= 0 \\
\mathbf{K}^-(\mathbf{y}, x | \mathbf{y}')|_{x=x_-} &= \delta(\mathbf{y} - \mathbf{y}'), \\
\mathbf{K}^-(\mathbf{y}, x | \mathbf{y}')|_{x=x_+} &= 0 \\
S &= -\frac{1}{2} \int d\mathbf{y} \left([\sqrt{-g} g^{xx} \phi(\mathbf{y}, x) \partial_x \phi(\mathbf{y}, x)]_{x=x_+} - [\sqrt{-g} g^{xx} \phi(\mathbf{y}, x) \partial_x \phi(\mathbf{y}, x)]_{x=x_-} \right) \\
S[\phi_0] &= -\frac{1}{2} \int d\mathbf{y} d\mathbf{y}' \phi_0^i(\mathbf{y}) \Delta_{ij}(\mathbf{y}, \mathbf{y}') \phi_0^j(\mathbf{y}') \\
\Delta_{+i}(\mathbf{y}, \mathbf{y}') &= [\sqrt{-g} g^{xx} \partial_x \mathbf{K}^i(\mathbf{y}, x | \mathbf{y}')]_{x=x_+}, \Delta_{-i}(\mathbf{y}, \mathbf{y}') = -[\sqrt{-g} g^{xx} \partial_x \mathbf{K}^i(\mathbf{y}, x | \mathbf{y}')]_{x=x_-} \\
\langle \psi_f | \mathcal{O}^\pm(\mathbf{y}) \mathcal{O}^\pm(\mathbf{y}') | \psi_i \rangle &\sim -i \Delta_{\pm\pm}(\mathbf{y}, \mathbf{y}') \\
\langle \psi_f | \mathcal{O}^\pm(\mathbf{y}) \mathcal{O}^\mp(\mathbf{y}') | \psi_i \rangle &\sim -i \Delta_{\pm\mp}(\mathbf{y}, \mathbf{y}') \\
\mathbf{K}^i(\mathbf{y}, x | \mathbf{y}') &= \lim_{x' \rightarrow x^i} \sqrt{-g} g^{x'x'} \partial_{x'} \mathbf{G}(\mathbf{y}, x | \mathbf{y}', x') \\
\Delta_{ij}(\mathbf{y}, \mathbf{y}') &\sim \lim_{x \rightarrow x^i, x' \rightarrow x^j} (\sqrt{-g} g^{xx}) (\sqrt{-g} g^{x'x'}) \frac{\partial^2}{\partial_x \partial_{x'}} \mathbf{G}(\mathbf{y}, x | \mathbf{y}', x') \\
ds^2 &= R^2 \left[-\frac{dt^2}{1-x^2} + \frac{dx^2}{(1-x^2)^2} \right] \\
\mathbf{K}^\pm(t, x) &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} e^{-i\omega t} f_\omega^\pm(x) \\
(1-x^2) \frac{d^2 f_\omega^\pm(x)}{dx^2} - x \frac{df_\omega^\pm(x)}{dx} + \left(\omega^2 - \frac{m^2 R^2}{1-x^2} \right) f_\omega^\pm(x) &= 0 \\
f_\omega^\pm(x) &= (1-x^2)^{\frac{1}{4}} [a_\omega^\pm P_\nu^\mu(x) + b_\omega^\pm Q_\nu^\mu(x)] \\
f_\omega^\pm(\pm x_\epsilon) &= 1, f_\omega^\pm(\mp x_\epsilon) = 0
\end{aligned}$$

$$\begin{aligned}
f_{\omega}^{+}(x) &= \left(\frac{1-x^2}{1-x_{\epsilon}^2}\right)^{\frac{1}{4}} \frac{Q_{\nu}^{\mu}(x)P_{\nu}^{\mu}(-x_{\epsilon}) - Q_{\nu}^{\mu}(-x_{\epsilon})P_{\nu}^{\mu}(x)}{Q_{\nu}^{\mu}(x_{\epsilon})P_{\nu}^{\mu}(-x_{\epsilon}) - Q_{\nu}^{\mu}(-x_{\epsilon})P_{\nu}^{\mu}(x_{\epsilon})} \\
f_{\omega}^{-}(x) &= \left(\frac{1-x^2}{1-x_{\epsilon}^2}\right)^{\frac{1}{4}} \frac{Q_{\nu}^{\mu}(x)P_{\nu}^{\mu}(x_{\epsilon}) - Q_{\nu}^{\mu}(x_{\epsilon})P_{\nu}^{\mu}(x)}{Q_{\nu}^{\mu}(-x_{\epsilon})P_{\nu}^{\mu}(x_{\epsilon}) - Q_{\nu}^{\mu}(x_{\epsilon})P_{\nu}^{\mu}(-x_{\epsilon})} \\
\omega_n &= \pm \left(n + \mu + \frac{1}{2}\right), n = 0, 1, 2, \dots \quad y \quad \frac{b_{\omega Q}}{a_{\omega Q}} = -\frac{2 \tan \pi \mu}{\pi} \\
\Delta_{\text{ren} \pm \pm}(t, t') &= \mp \frac{2^{\Delta_-} \Gamma(1-\mu)}{2\pi \Gamma(1+\mu)} \int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \Gamma\left(\frac{1}{2} + \mu - \omega\right) \Gamma\left(\frac{1}{2} + \mu + \omega\right) \cos(\pi\omega) \\
\Delta_{\text{ren} \pm \mp}(t, t') &= \mp \frac{2^{\Delta_-}}{\Gamma(\mu)^2} \int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \Gamma\left(\frac{1}{2} + \mu - \omega\right) \Gamma\left(\frac{1}{2} + \mu + \omega\right) \\
\Delta_{\text{ren} \pm \pm}^F(t, t') &= \mp \left(\frac{i^{\Delta_- - \Delta_+}}{8^{\Delta_+} \pi^{\frac{1}{2}}}\right) \frac{\Gamma\left(\frac{1}{2} + \mu\right)}{\Gamma(\mu) \sin^{2\Delta_+}\left(\frac{t-t'}{2}\right)} = \Delta_{\text{ren} \pm \mp}^F(t, t') = \mp \left(\frac{8^{\Delta_-} i}{4}\right) \frac{\Gamma(1+2\mu)}{\Gamma(\mu)^2 \cos^{2\Delta_+}\left(\frac{t-t'}{2}\right)} \\
\langle 0 | T \mathcal{O}^{\pm}(t) \mathcal{O}^{\pm}(t') | 0 \rangle &= \pm \left(\frac{4^{\Delta_-} i^{2\Delta_-}}{8^{\Delta_+}}\right) \frac{\Gamma(2\mu)}{\Gamma(\mu)^2 \sin^{2\Delta_+}\left(\frac{t-t'}{2}\right)} \\
\langle 0 | T \mathcal{O}^{\pm}(t) \mathcal{O}^{\mp}(t') | 0 \rangle &= \mp \left(\frac{8^{\Delta_-}}{4}\right) \frac{\Gamma(1+2\mu)}{\Gamma(\mu)^2 \cos^{2\Delta_+}\left(\frac{t-t'}{2}\right)} \\
\langle 0 | T \mathcal{O}^{\pm}(t) \mathcal{O}^{\pm}(t') | 0 \rangle &= \pm \frac{1}{8\pi \sin^2\left(\frac{t-t'}{2}\right)}, \quad \langle 0 | T \mathcal{O}^{\pm}(t) \mathcal{O}^{\mp}(t') | 0 \rangle = \mp \frac{1}{4\pi \cos^2\left(\frac{t-t'}{2}\right)} \\
S_5 &= \kappa \int \epsilon_{abcde} \left(R^{ab} R^{cd} + \frac{2}{3l^2} R^{ab} e^c e^d + \frac{1}{5l^4} e^a e^b e^c e^d \right) e^e \\
ds^2 &= R^2 \left[-\cosh^2 \rho dt^2 + d\rho^2 + \cosh^2 \rho d\tilde{\Sigma}_{d-1}^2 \right] \\
&= R^2 \left[-\frac{dt^2}{1-x^2} + \frac{dx^2}{(1-x^2)^2} + \frac{d\tilde{\Sigma}_{d-1}^2}{1-x^2} \right] \\
K^{\pm}(t, x, \theta \mid t', \theta') &= \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \sum_Q e^{-i\omega(t-t')} Y_Q(\theta) Y_Q^*(\theta') f_{\omega Q}^{\pm}(x) \\
(1-x^2) \frac{d^2 f_{\omega Q}^{\pm}(x)}{dx^2} + (d-2)x \frac{df_{\omega Q}^{\pm}(x)}{dx} + \left[(\omega^2 - Q^2) - \frac{m^2 R^2}{1-x^2} \right] f_{\omega Q}^{\pm}(x) &= 0 \\
f_{\omega Q}^{\pm}(x) &= (1-x^2)^{\frac{d}{4}} \left[a_{\omega Q}^{\pm} P_{\nu}^{\mu}(x) + b_{\omega Q}^{\pm} Q_{\nu}^{\mu}(x) \right] \\
\nu &= \varpi - \frac{1}{2} = \sqrt{\left(\frac{d-1}{2}\right)^2 + \omega^2 - Q^2} - \frac{1}{2} \\
f_{\omega Q}^{\pm}(\pm x_{\epsilon}) &= 1, f_{\omega Q}^{\pm}(\mp x_{\epsilon}) = 0 \\
f_{\omega Q}^{+}(x) &= \left(\frac{1-x^2}{1-x_{\epsilon}^2}\right)^{\frac{d}{4}} \frac{P_{\nu}^{\mu}(x)Q_{\nu}^{\mu}(-x_{\epsilon}) - P_{\nu}^{\mu}(-x_{\epsilon})Q_{\nu}^{\mu}(x)}{P_{\nu}^{\mu}(x_{\epsilon})Q_{\nu}^{\mu}(-x_{\epsilon}) - P_{\nu}^{\mu}(-x_{\epsilon})Q_{\nu}^{\mu}(x_{\epsilon})} \\
f_{\omega Q}^{-}(x) &= \left(\frac{1-x^2}{1-x_{\epsilon}^2}\right)^{\frac{d}{4}} \frac{P_{\nu}^{\mu}(x)Q_{\nu}^{\mu}(x_{\epsilon}) - P_{\nu}^{\mu}(x_{\epsilon})Q_{\nu}^{\mu}(x)}{P_{\nu}^{\mu}(-x_{\epsilon})Q_{\nu}^{\mu}(x_{\epsilon}) - P_{\nu}^{\mu}(x_{\epsilon})Q_{\nu}^{\mu}(-x_{\epsilon})} \\
\omega_{nQ} &= \pm \sqrt{\left(\mu + \frac{1}{2} + n\right)^2 + Q^2 - \left(\frac{d-1}{2}\right)^2}, n = 0, 1, \dots, \quad y \quad \frac{b_{\omega Q}}{a_{\omega Q}} = -\frac{2 \tan \pi \mu}{\pi}
\end{aligned}$$

$$\begin{aligned}
\langle \psi_f | T \mathcal{O}^\pm(t, \theta) \mathcal{O}^\pm(t', \theta') | \psi_i \rangle &= \pm i \frac{2^{\Delta-d} \Gamma(1-\mu)}{\pi 2^d \Gamma(1+\mu)} \sum_Q Y_Q(\theta) Y_Q^*(\theta') \\
&\times \int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \Gamma\left(\frac{1}{2} + \mu - \varpi\right) \Gamma\left(\frac{1}{2} + \mu + \varpi\right) \cos(\pi\varpi) \\
\langle \psi_f | T \mathcal{O}^\pm(t, \theta) \mathcal{O}^\mp(t', \theta') | \psi_i \rangle &= \pm i \frac{2^{\Delta-d}}{2^{d-1} \Gamma(\mu)^2} \sum_Q Y_Q(\theta) Y_Q^*(\theta') \\
&\times \int \frac{d\omega}{2\pi} e^{-i\omega(t-t')} \Gamma\left(\frac{1}{2} + \mu - \varpi\right) \Gamma\left(\frac{1}{2} + \mu + \varpi\right) \\
S_A &= \frac{\text{Area}(\gamma_A)}{4G_N^{(d+1)}} \\
S_{\tilde{\Sigma}} &= \frac{\text{Area}(\tilde{\Sigma})}{4G_N^{(d+1)}} \\
Z_{\text{gravity}}[\phi(\phi_0^+, \phi_0^-, \mathcal{C})] &\sim e^{-\frac{i}{2} \int dy dy' \phi_0^i(y) \Delta_{ij}(y, y') \phi_0^j(y')} \\
Z_{\text{gravity}}[\phi(\phi_0^+, \phi_0^-, \mathcal{C})] &= \langle \psi_f | T e^{i \int dy \phi_0^+(y) \mathcal{O}^+(y) + i \int dy \phi_0^-(y) \mathcal{O}^-(y)} | \psi_i \rangle_{QFT} \\
\langle \psi_0 | T e^{i \int dy \phi_0^+(y) \mathcal{O}^+(y) + i \int dy \phi_0^-(y) \mathcal{O}^-(y)} | \psi_0 \rangle_{QFT} &= \text{Tr}[\rho_{\psi_0} T e^{i \int dy \phi_0^+(y) \mathcal{O}^+(y) + i \int dy \phi_0^-(y) \mathcal{O}^-(y)}] \\
\langle e^{-\int dy \phi_0^+(y) \mathcal{O}^+(y) - \int dy \phi_0^-(y) \mathcal{O}^-(y)} \rangle_\beta &= \text{Tr}[\rho_\beta e^{-\int dy \phi_0^+(y) \mathcal{O}^+(y) - \int dy \phi_0^-(y) \mathcal{O}^-(y)}] \\
\text{Tr}[e^{-\beta H} e^{-\int dy \phi_0^+(y) \mathcal{O}^+(y) - \int dy \phi_0^-(y) \mathcal{O}^-(y)}] &= Z_{\text{Egravity}}[\phi(\phi_0^+, \phi_0^-)] \sim e^{-S_E[\phi_0^+, \phi_0^-]} \\
Z_{\text{Egravity}}[\phi(\phi_0^+, \phi_0^-)] &\sim e^{-\frac{1}{2} \int dy dy' \phi_0^i(y) \tilde{\Delta}_{ij}(y, y') \phi_0^j(y')} \\
\text{Tr}[e^{-\beta H} e^{-\int dy \phi_0^+(y) \mathcal{O}^+(y) - \int dy \phi_0^-(y) \mathcal{O}^-(y)}] &\sim e^{-\frac{1}{2} \int dy dy' \phi_0^i(y) \tilde{\Delta}_{ij}(y, y') \phi_0^j(y')} \\
H[\Psi_+, \Psi_-] &= H_+[\Psi_+] + H_-[\Psi_-] \\
\text{Tr}[e^{-\beta H} e^{-\int dy \phi_0(y) \mathcal{O}(y)}] &\sim e^{-\frac{1}{2} \int dy dy' \phi_0(y) \tilde{\Delta}(y, y') \phi_0(y')} \\
S &= \frac{\eta}{2\pi\alpha'} \int d\tau d\sigma \sqrt{\eta \hbar} \\
ds^2 &= -g_t(r) dt^2 + g_x(r) dx_i^2 + g_r(r) dr^2 + g_{ab}(r, \theta) d\theta^a d\theta^b \\
S &= -\frac{1}{2\pi\alpha'} \int d\tau d\sigma \sqrt{g_t(r) \dot{t}^2 (g_x(r) \dot{x}^2 + g_r(r) \dot{r}^2)} \\
&= -\frac{1}{2\pi\alpha'} \int d\tau d\sigma \sqrt{g_t(r) (g_x(r) \dot{x}^2 + g_r(r) \dot{r}^2)} \\
&= -\frac{\mathcal{T}}{2\pi\alpha'} \int d\sigma \sqrt{f^2(r) \dot{x}^2 + g^2(r) \dot{r}^2} \\
ds_{eff}^2 &= f^2(r) dx^2 + g^2(r) dr^2 \\
\frac{f^2(r) \dot{x}(\sigma)}{\sqrt{f^2(r) \dot{x}(\sigma)^2 + g^2(r) \dot{r}(\sigma)^2}} &= A \\
\dot{x}(\sigma) &= \pm A \frac{g(r)}{f(r)} \frac{1}{\sqrt{f^2(r) - A^2}} \dot{r}(\sigma) \\
\frac{dx}{dr} &= \pm \frac{g(r)}{f(r)} \frac{f(r_0)}{\sqrt{f^2(r) - f^2(r_0)}} \\
U(r) &= \frac{f^2(r) (f^2(r) - f^2(r_0))}{g^2(r) f^2(r_0)} \\
L(r_0) &= 2 \int_{r_0}^{\infty} \frac{g(r)}{f(r)} \frac{f(r_0)}{\sqrt{f^2(r) - f^2(r_0)}} dr
\end{aligned}$$

$$\begin{aligned}
\frac{L'(r_0)}{2} &= -\frac{g(r)}{\sqrt{f^2(r) - f^2(r_0)}} \Big|_{r \rightarrow r_0} + f'(r_0) \int_{r_0}^{\infty} \frac{f(r)g(r)}{(f^2(r) - f^2(r_0))^{\frac{3}{2}}} dr \\
&= -\frac{g(r)}{\sqrt{f^2(r) - f^2(r_0)}} \Big|_{r \rightarrow r_0} + f'(r_0) \int_{r_0}^{\infty} dr \frac{g(r)}{f'(r)} \frac{d}{dr} \left(-\frac{1}{\sqrt{f^2(r) - f^2(r_0)}} \right) \\
&= -\frac{f'(r_0)g(r)}{f'(r)\sqrt{f^2(r) - f^2(r_0)}} \Big|_{r \rightarrow \infty} + \int_{r_0}^{\infty} dr \frac{f'(r_0)}{\sqrt{f^2(r) - f^2(r_0)}} \frac{d}{dr} \left(\frac{g(r)}{f'(r)} \right), \\
L'(r_0) &= 2 \int_{r_0}^{\infty} dr \frac{f'(r_0)}{\sqrt{f^2(r) - f^2(r_0)}} \frac{d}{dr} \left(\frac{g(r)}{f'(r)} \right) \\
E &= \frac{1}{2\pi\alpha'} \int d\sigma \sqrt{f^2(r)\dot{x}(\sigma)^2 + g^2(r)\dot{r}(\sigma)^2} \\
E(r_0) &= \frac{1}{2\pi\alpha'} \int_{-L/2}^{L/2} dx \frac{f^2(r(x))}{f(r_0)} = \frac{1}{\pi\alpha'} \int_{r_0}^{\infty} dr \frac{g(r)f(r)}{\sqrt{f^2(r) - f^2(r_0)}} \\
m_q &= \frac{1}{2\pi\alpha'} \int_{r_{\min}}^{\infty} g(r) dr \\
E_{q\bar{q}}(r_0) &= E(r_0) - 2m_q \\
&= \frac{1}{\pi\alpha'} \left[\int_{r_0}^{\infty} \frac{g(r)f(r)}{\sqrt{f^2(r) - f^2(r_0)}} dr - \int_{r_{\min}}^{\infty} g(r) dr \right] \\
E'_{q\bar{q}}(r_0) &= \frac{1}{\pi\alpha'} \left[-\frac{g(r)f(r)}{\sqrt{f^2(r) - f^2(r_0)}} \Big|_{r=r_0} + \int_{r_0}^{\infty} dr \frac{f(r)g(r)f(r_0)f'(r_0)}{(f^2(r) - f^2(r_0))^{\frac{3}{2}}} \right] \\
E'_{q\bar{q}}(r_0) &= \frac{1}{2\pi\alpha'} f(r_0)L'(r_0) \Rightarrow \frac{dE_{q\bar{q}}}{dL} = \frac{1}{2\pi\alpha'} f(r_0) \\
\frac{dV}{dL} &> 0, \frac{d^2V}{dL^2} \leq 0 \\
\frac{dV_{\text{string}}}{dL} &= \frac{dE_{q\bar{q}}}{dr_0} \frac{dr_0}{dL} = \frac{1}{2\pi\alpha'} f(r_0), \frac{d^2V_{\text{string}}}{dL^2} = \frac{1}{2\pi\alpha'} \left(\frac{dL}{dr_0} \right)^{-1} f'(r_0) \\
X^\mu &= (\tau, x_{\text{cl}}(\sigma) + \delta x_1(\tau, \sigma), \delta x_2(\tau, \sigma), \delta x_3(\tau, \sigma), r_{\text{cl}}(\sigma) + \delta r(\tau, \sigma), \theta^a + \delta\theta^a(\tau, \sigma)) \\
t &= \tau, x_1 = x_{\text{cl}}(r) + \delta x_1(t, r), x_2 = \delta x_2(t, r), x_3 = \delta x_3(t, r), r = \sigma \\
2\pi\alpha' \mathcal{L}^{(2)} &= \frac{1}{g(r)f(r)\sqrt{f^2(r) - f^2(r_0)}} \left[h^2(r)(f^2(r) - f^2(r_0))(\delta\dot{x}_1)^2 - (f^2(r) - f^2(r_0))^2(\delta\dot{x}_1)^2 \right. \\
&\quad \left. + f^2(r)h^2(r)((\delta\dot{x}_2)^2 + (\delta\dot{x}_3)^2) - f^2(r)(f^2(r) - f^2(r_0))((\delta\dot{x}_2)^2 + (\delta\dot{x}_3)^2) \right] \\
&\quad \left[\frac{d}{dr} \left(\frac{(f^2(r) - f^2(r_0))^{\frac{3}{2}}}{g(r)f(r)} \frac{d}{dr} \right) + \omega^2 \frac{h^2(r)\sqrt{f^2(r) - f^2(r_0)}}{g(r)f(r)} \right] \delta x_1(r) \\
&\quad \left[\frac{d}{dr} \left(\frac{f(r)\sqrt{f^2(r) - f^2(r_0)}}{g(r)} \frac{d}{dr} \right) + \omega^2 \frac{h^2(r)f(r)}{g(r)\sqrt{f^2(r) - f^2(r_0)}} \right] \delta x_m(r) = 0, m = 2, 3 \\
&\quad \frac{d}{dr} \left(\sqrt{r - r_0} \frac{d\delta x_m(r)}{dr} \right) + \frac{\omega^2 h^2(r_0)}{2f(r_0)f'(r_0)} \frac{1}{\sqrt{r - r_0}} \delta x_m(r) \approx 0 \Rightarrow \delta x_m(r) \\
&\quad \approx C_0 + C_1 \sqrt{r - r_0} + O(r - r_0) \\
\frac{d}{dr} \left((r - r_0)^{\frac{3}{2}} \frac{d\delta x_1(r)}{dr} \right) &+ \frac{\omega^2 h^2(r_0)}{2f(r_0)f'(r_0)} \sqrt{r - r_0} \delta x_1(r) \approx 0 \Rightarrow \delta x_1(r) \approx C'_0 + C'_1 \frac{1}{\sqrt{r - r_0}} + O(\sqrt{r - r_0}) \\
&\quad \frac{d}{dr} \left(\sqrt{r - r_0} \frac{d\delta x_m(r)}{dr} \right) + \frac{\omega^2 h^2(r_0)}{2f(r_0)f'(r_0)} \frac{1}{\sqrt{r - r_0}} \delta x_m(r) \approx 0 \Rightarrow \delta x_m(r) \\
&\quad \approx C_0 + C_1 \sqrt{r - r_0} + O(r - r_0)
\end{aligned}$$

$$\begin{aligned}
u &= r + \Delta(t, r) \text{ where } \Delta(t, r) = \frac{\delta x_1(t, r)}{x'_{\text{cl}}(r)} \\
x_1 &= x_{\text{cl}}(r) + \delta x_1(t, r) = x_{\text{cl}}(u - \Delta(t, r)) + \delta x_1(t, r) \\
&\approx x_{\text{cl}}(u) - x'_{\text{cl}}(r) \frac{\delta x_1(t, r)}{x'_{\text{cl}}(r)} + \delta x_1(t, r) = x_{\text{cl}}(u) \\
r &\approx u - \frac{\delta x_1(t, u)}{x'_{\text{cl}}(u)} \\
r &\approx r_0 - \alpha(C'_0 \sqrt{u - r_0} + C'_1) + O(u - r_0) \\
\left. \frac{d\delta r(t, r)}{dx_1} \right|_{r=r_0} &= 0 \text{ where } \delta r(t, r) = -\frac{\delta x_1(t, r)}{x'_{\text{cl}}(r)} \\
\delta x_1(r) + 2(r - r_0) \frac{d\delta x_1(r)}{dr} &= 0, \quad r \rightarrow r_0 \\
\sqrt{r - r_0} \delta x_1(r) &= 1, \quad r \rightarrow r_0 \\
\delta x_1(r) + 2(r - r_0) \frac{d\delta x_1(r)}{dr} &= 1, \quad r \rightarrow r_0 \\
\sqrt{r - r_0} \delta x_1(r) &= 0, \quad r \rightarrow r_0 \\
ds^2 &= \frac{r^2}{R^2} (-dt^2 + dx_i dx_i) + R^2 \frac{dr^2}{r^2} + R^2 d\Omega_5^2 \\
x_{\text{cl}}(r) &= \pm \left\{ cte - \frac{R^2}{4r_0} \text{B} \left(\left(\frac{r_0}{r} \right)^4; \frac{3}{4}, \frac{1}{2} \right) \right\}, r_0 \leq r < \infty \\
L(r_0) &= \frac{R^2}{2r_0} \text{B} \left(\frac{3}{4}, \frac{1}{2} \right) = \frac{R^2 (2\pi)^{\frac{3}{2}}}{r_0 \Gamma \left[\frac{1}{4} \right]^2} \\
E_{q\bar{q}}(r_0) &= \frac{r_0}{\pi\alpha'} (K(-1) - E(-1)) = -\frac{r_0}{2\pi\alpha'} \frac{(2\pi)^{\frac{3}{2}}}{\Gamma \left[\frac{1}{4} \right]^2} \\
V_{\text{string}}(L) &= -\frac{(2\pi)^2 R^2 / \alpha'}{\Gamma \left[\frac{1}{4} \right]^4} \frac{1}{L} \sim -\frac{\sqrt{\lambda}}{L} \\
ds^2 &= R^2 [-\cosh^2 \rho dt^2 + d\rho^2 + \sinh^2 \rho d\Omega_3^2] \\
\dot{\rho}^2 + U(\rho) &= 0 \\
\Phi(\rho_0) &= 2 \int_{r_0}^{\infty} \frac{\sinh 2\rho_0}{\sinh \rho \sqrt{\sinh^2 2\rho - \sinh^2 2\rho_0}} d\rho \\
E(\rho_0) &= \frac{R}{2\pi\alpha'} \int_{\rho_0}^{\infty} \frac{\sinh^2 2\rho}{\sinh \rho \sqrt{\sinh^2 2\rho - \sinh^2 2\rho_0}} d\rho \\
E_{q\bar{q}}(\rho_0) &= \frac{R}{2\pi\alpha'} \left[\int_{\rho_0}^{\infty} \left(\frac{2\cosh \rho}{\sqrt{1 - \frac{\sinh^2 2\rho_0}{\sinh^2 2\rho}}} - 2\cosh \rho \right) d\rho - 2\sinh \rho_0 \right] \\
ds^2 &= \frac{r^2}{R^2} \left[-\left(1 - \frac{\mu^4}{r^4} \right) dt^2 + dx_i dx_i \right] + \frac{R^2}{r^2} \frac{1}{1 - \frac{\mu^4}{r^4}} dr^2 + R^2 d\Omega_5^2 \\
ds^2 &= R^2 \left[-\left(\rho^2 - \frac{1}{\rho^2} \right) d\vec{t}^2 + \rho^2 dy_i dy_i + \frac{1}{\rho^2 - \frac{1}{\rho^2}} d\rho^2 + d\Omega_5^2 \right]
\end{aligned}$$

$$\begin{aligned}
\bar{L}(\rho_0) &= \frac{(2\pi)^{\frac{3}{2}} \sqrt{\rho_0^4 - 1}}{\Gamma\left[\frac{1}{4}\right]^2 \rho_0^3} {}_2F_1\left(\frac{3}{4}, \frac{1}{2}, \frac{5}{4}; \frac{1}{\rho_0^4}\right) \\
\bar{E}_{q\bar{q}}(\rho_0) &= \frac{R^2}{\pi\alpha'} \left[1 - \frac{(2\pi)^{\frac{3}{2}}}{2\Gamma\left[\frac{1}{4}\right]^2} \rho_0 F_1\left(-\frac{1}{2}, -\frac{1}{4}, \frac{1}{4}; \frac{1}{\rho_0^4}\right) \right] \\
ds^2 &= \alpha' N e^\phi \left[-dt^2 + dx_i dx_i + dr^2 + e^{2h} (d\theta^2 + \sin^2 \theta d\varphi^2) + \frac{1}{4} (w^i - A^i)^2 \right] \\
w^1 + i w^2 &= e^{-i\psi} (d\tilde{\theta} + i \sin \tilde{\theta} d\tilde{\varphi}), w^3 = d\psi + \cos \tilde{\theta} d\tilde{\varphi} \\
A^1 &= -a(r) d\theta, A^2 = a(r) \sin \theta d\varphi, A^3 = -\cos \theta d\varphi \\
a(r) &= \frac{2r}{\sinh 2r} \\
e^{2h} &= r \coth 2r - \frac{r^2}{\sinh^2 2r} - \frac{1}{4} \\
e^{2\phi} &= e^{2\phi_0} \frac{\sinh 2r}{2e^h} \\
\bar{L}(r_0) &= 2 \int_{r_0}^{\infty} \frac{e^{\phi(r_0)}}{\sqrt{e^{2\phi(r)} - e^{2\phi(r_0)}}} dr \\
\bar{E}_{q\bar{q}}(r_0) &= \frac{N}{\pi} \left[\int_{r_0}^{\infty} \frac{e^{2\phi(r)}}{\sqrt{e^{2\phi(r)} - e^{2\phi(r_0)}}} dr - \int_0^{\infty} e^{\phi(r)} dr \right] \\
\bar{E}_{q\bar{q}}(r_0) &= \frac{N}{\pi} \left[\int_{r_0}^{\infty} \left(\frac{e^{2\phi(r)} + e^{2\phi(r_0)} - e^{2\phi(r_0)}}{\sqrt{e^{2\phi(r)} - e^{2\phi(r_0)}}} - e^{\phi(r)} \right) dr - \int_0^{r_0} e^{\phi(r)} dr \right] \\
&= \frac{N}{\pi} \left[e^{\phi(r_0)} \frac{\bar{L}(r_0)}{2} + \int_{r_0}^{\infty} dr \left(\sqrt{e^{2\phi(r)} - e^{2\phi(r_0)}} - e^{\phi(r)} \right) - \int_0^{r_0} e^{\phi(r)} dr \right] \\
V_{\text{string}}(L) &\approx \frac{e^{\phi_0}}{2\pi\alpha'} L \Rightarrow T_{\text{string}} = \frac{e^{\phi_0}}{2\pi\alpha'} \\
ds^2 &= g_s \alpha' M \left[h^{-\frac{1}{2}}(r) (-dt^2 + dx_i dx_i) + h^{\frac{1}{2}}(r) ds_6^2 \right] \\
ds_6^2 &= \frac{1}{2} K(r) \left[\frac{(dr^2 + (g^5)^2)}{3K^3(r)} + \cosh^2 \frac{r}{2} ((g^3)^2 + (g^4)^2) + \sinh^2 \frac{r}{2} ((g^1)^2 + (g^2)^2) \right] \\
K(r) &= \frac{[\sinh(2r) - 2r]^{\frac{1}{3}}}{2^{\frac{1}{3}} \sinh r} \\
g^1 &= \frac{e^1 - e^3}{\sqrt{2}}, g^2 = \frac{e^2 - e^4}{\sqrt{2}}, g^3 = \frac{e^1 + e^3}{\sqrt{2}}, g^4 = \frac{e^2 + e^4}{\sqrt{2}}, g^5 = e^5 \\
e^1 &= -\sin \theta_1 d\phi_1, e^2 = d\theta_1, e^3 = -\sin \psi d\theta_2 + \cos \psi \sin \theta_2 d\phi_2 \\
e^4 &= \cos \psi d\theta_2 + \sin \psi \sin \theta_2 d\phi_2, \\
e^5 &= d\psi + \cos \theta_1 d\phi_1 + \cos \theta_2 d\phi_2 \\
h(r) &= 2^{\frac{2}{3}} \int_r^{\infty} dx \frac{x \coth x - 1}{\sinh^2 x} (\sinh 2x - 2x)^{\frac{1}{3}} \\
\bar{L}(r_0) &= 2 \int_{r_0}^{\infty} \frac{dr}{\sqrt{6} K(r)} \frac{h(r)}{\sqrt{h(r_0) - h(r)}} \\
\bar{E}_{q\bar{q}}(r_0) &= \frac{g_s M}{\pi} \left[\int_{r_0}^{\infty} \frac{dr}{\sqrt{6} K(r)} \frac{\sqrt{h(r_0)}}{\sqrt{h(r_0) - h(r)}} - \int_0^{r_0} \frac{dr}{\sqrt{6} K(r)} \right] \\
T_{\text{string}} &= \frac{1}{2\pi\alpha'} \frac{\ell_{cf}^2}{g_s \alpha' M \sqrt{h_0}}
\end{aligned}$$

$$\begin{aligned}
ds^2 &= g_s \alpha' N e^{4f(r)} \left[-dt^2 + dx_i dx_i + dr^2 + e^{2h(r)} (d\theta^2 + \sin^2 \theta d\varphi^2) \right. \\
&\quad \left. + \frac{e^{2g(r)}}{4} ((w_1 + a(r)d\theta)^2 + (w_2 - a(r)\sin \theta d\varphi)^2) + \frac{e^{2k(r)}}{4} (w_3 + \cos \theta d\varphi)^2 \right] \\
\partial_\rho a &= \frac{-2}{-1 + 2\rho \coth 2\rho} \left[e^{2k} \frac{(\operatorname{acosh} 2\rho - 1)^2}{\sinh 2\rho} + a(2\rho - a \sinh 2\rho) \right] \\
\partial_\rho k &= \frac{2(1 + a^2 - 2a \cosh 2\rho)^{-1}}{-1 + 2\rho \coth 2\rho} \left[e^{2k} a \sinh 2\rho (\operatorname{acosh} 2\rho - 1) + \left(2\rho - 4a\rho \cosh 2\rho + \frac{a^2}{2} \sinh 4\rho \right) \right] \\
\partial_\rho f &= -\frac{1}{4 \sinh^2 2\rho} \left[\frac{(1 - a \cosh 2\rho)^2 (-4\rho + \sinh 4\rho)}{(1 + a^2 - 2a \cosh 2\rho)(-1 + 2\rho \coth 2\rho)} \right] \\
e^{2g} &= \frac{b \cosh 2\rho - 1}{a \cosh 2\rho - 1}, e^{2h} = \frac{e^{2g}}{4} (2a \cosh 2\rho - 1 - a^2), \text{ con } b(\rho) = \frac{2\rho}{\sinh 2\rho} \\
a(\rho) &= 1 + \mu \rho^2 + \dots, e^{2k(\rho)} = \frac{4}{6 + 3\mu} - \frac{20 + 36\mu + 9\mu^2}{15(2 + \mu)} \rho^2 + \dots \\
e^{2g(\rho)} &= \frac{4}{6 + 3\mu} + \dots, e^{2h(\rho)} = \frac{4\rho^2}{6 + 3\mu} + \dots, e^{2f(\rho)} = 1 + \frac{(2 + \mu)^2}{8} \rho^2 + \dots \\
\bar{L}(\rho_0) &= 2 \int_{\rho_0}^{\rho_\infty} \frac{e^{4f(\rho_0)}}{\sqrt{e^{8f(\rho)} - e^{8f(\rho_0)}}} e^{k(\rho)} d\rho \\
\bar{E}_{q\bar{q}}(\rho_0) &= \frac{Ng_s}{\pi} \left[\int_{\rho_0}^{\rho_\infty} \frac{e^{8f(\rho)}}{\sqrt{e^{8f(\rho)} - e^{8f(\rho_0)}}} e^{k(\rho)} d\rho - \int_0^{\rho_\infty} e^{4f(\rho)} e^{k(\rho)} d\rho \right] \\
T_{\text{string}} &= \frac{g_s}{2\pi\alpha'} \\
\delta x_1(r) + 2(r - r_0) \frac{d\delta x_1(r)}{dr} &= 0, \quad r \rightarrow r_0 \\
\sqrt{r - r_0} \delta x_1(r) &= 1, \quad r \rightarrow r_0 \\
\delta x_1(r) &= 0, r \rightarrow \infty \\
\delta x_1^{(0)}(r) &= C \int_r^\infty d\bar{r} \frac{g(\bar{r}) f(\bar{r})}{(f^2(\bar{r}) - f^2(r_0))^{\frac{3}{2}}} + C' \\
\delta x_1^{(0)}(r) &= - \int_r^\infty d\bar{r} \frac{g(\bar{r})}{f'(\bar{r})} \frac{d}{d\bar{r}} \left(\frac{1}{\sqrt{f^2(\bar{r}) - f^2(r_0)}} \right) \\
&= \frac{g(r)}{f'(r) \sqrt{f^2(r) - f^2(r_0)}} + \frac{L'(r)}{2f'(r)}. \\
\delta x_1^{(0)}(r) &= \frac{g(r_0)}{\sqrt{2}(f'(r_0))^{\frac{3}{2}}} \frac{1}{\sqrt{r - r_0}} + \frac{L'(r_0)}{2f'(r_0)} + \mathcal{O}(\sqrt{r - r_0}) \\
\left[\frac{d}{dr} \left(\frac{(r^4 - r_0^4)^{\frac{3}{2}}}{r^2} \frac{d}{dr} \right) + \omega^2 R^4 \frac{\sqrt{r^4 - r_0^4}}{r^2} \right] \delta x_1(r) &= 0 \quad 0 < r_0 \leq r < \infty \\
\left[\frac{d}{d\rho} \left(\frac{(\rho^4 - 1)^{\frac{3}{2}}}{\rho^2} \frac{d}{d\rho} \right) + \frac{\omega^2 R^4 \sqrt{\rho^4 - 1}}{r_0^2 \rho^2} \right] \delta x_1(\rho) &= 0 \\
\left[\frac{d}{d\rho} \left(\rho^4 \frac{d}{d\rho} \right) + \frac{\omega^2 R^4}{r_0^2} \right] \delta x_1(\rho) &\approx 0, \rho \gg 1 \\
\delta x_1(\rho) &\approx \alpha_0 + \frac{\alpha_1}{\rho^3}, \rho \gg 1
\end{aligned}$$

$$\begin{aligned}
& \left[\frac{d}{d\rho} \left(\frac{(\rho^4 - \rho_0^4)^{\frac{3}{2}}}{\sqrt{\rho^4 - 1}} \frac{d}{d\rho} \right) + \frac{\omega^2 R^4 \rho^4 \sqrt{\rho^4 - \rho_0^4}}{\mu^2 (\rho^4 - 1)^{\frac{3}{2}}} \right] \delta x_1(\rho) = 0, 1 < \rho_0 \leq \rho < \infty \\
& \left[\frac{d}{dr} \left(\frac{(e^{2\phi(r)} - e^{2\phi(r_0)})^{\frac{3}{2}}}{e^{2\phi(r)}} \frac{d}{dr} \right) + \bar{\omega}^2 \sqrt{e^{2\phi(r)} - e^{2\phi(r_0)}} \right] \delta x_1(r) = 0, 0 < r_0 \leq r < \infty \\
& \left[\frac{d}{dr} \left(e^r r^{-\frac{1}{4}} \frac{d}{dr} \right) + \bar{\omega}^2 e^r r^{-\frac{1}{4}} \right] \delta x_1(r) = 0, r \gg 1 \\
& \left[\frac{d^2}{dr^2} + \left(1 - \frac{1}{4r} \right) \frac{d}{dr} + \bar{\omega}^2 \right] \delta x_1(r) = 0, r \gg 1 \\
& \delta x_1(r) \simeq e^{-\frac{1}{2}r} (\beta_0 e^{r\alpha} + \beta_1 e^{-r\alpha}), r \gg 1 \\
& \left[\frac{d}{dr} \left(\frac{K(r)}{h(r)} \left(1 - \frac{h(r)}{h(r_0)} \right)^{\frac{3}{2}} \frac{d}{dr} \right) + \bar{\omega}^2 \frac{1}{6K(r)} \sqrt{1 - \frac{h(r)}{h(r_0)}} \right] \delta x_1(r) = 0, 0 < r_0 \leq r < \infty \\
& \left[\frac{d}{dr} \left(\frac{e^r}{r} \frac{d}{dr} \right) + \bar{\omega}^2 \frac{e^{\frac{r}{3}}}{2^{\frac{3}{2}}} \right] \delta x_1(r) = 0, r \gg 1 \\
& \left[\frac{d^2}{dr^2} + \left(1 - \frac{1}{r} \right) \frac{d}{dr} + \bar{\omega}^2 \frac{r e^{\frac{2}{3}r}}{2^{\frac{4}{3}}} \right] \delta x_1(r) = 0, r \gg 1 \\
& \delta x_1(r) \simeq \alpha_0 + \alpha_1 e^{-r}, r \gg 1 \\
& \left[\frac{d}{d\rho} \left(\frac{(e^{8f(\rho)} - e^{8f(\rho_0)})^{\frac{3}{2}}}{e^{8f(\rho)+k(\rho)}} \frac{d}{d\rho} \right) + \bar{\omega}^2 e^{k(\rho)} \sqrt{e^{8f(\rho)} - e^{8f(\rho_0)}} \right] \delta x_1(\rho) = 0, 0 < \rho_0 \leq \rho < \infty \\
& \left[e^{-k(\rho)} \frac{d}{d\rho} \left(e^{-k(\rho)} \frac{d}{d\rho} \right) + \bar{\omega}^2 \frac{e^{8f_\infty}}{e^{8f_\infty} - e^{8f(\rho_0)}} \right] \delta x_1(\rho) = 0, \rho \gg 1 \\
& \left[\frac{d^2}{dr^2} + \bar{\omega}^2 \right] \delta x_1(r) = 0, r \gg 1 \\
& V(\rho) = 2 \frac{\rho^4 - 1}{\rho^2}, \rho \in [1, \infty) \\
& y(\rho) = y_0 - \frac{1}{4} \text{B} \left(\frac{1}{\rho^4}; \frac{1}{4}, \frac{1}{2} \right) \\
& V(\rho, \rho_0) = 2 \frac{\rho^8(\rho^4 - \rho_0^4) - \rho_0^4(4\rho^4 - 1) - 3\rho^4}{\rho^6(\rho^4 - 1)}, 1 < \rho_0 \leq \rho < \infty \\
& V(r, r_0) = \frac{e^{-2\phi(r)}}{4} \left((e^{2\phi(r)} - 3e^{2\phi(r_0)})\phi^2(r) + 2(e^{2\phi(r)} + e^{2\phi(r_0)})\phi''(r) \right), 0 < r_0 \leq r < \infty \\
& \delta x_1 = \frac{e^{\frac{\phi(r)}{2}}}{(e^{2\phi(r)} - e^{2\phi(r_0)})^{\frac{1}{2}}} \Psi \simeq e^{-\frac{r}{2}} \Psi, r \rightarrow \infty \\
& V(r, r_0) = -\frac{3K(r)}{8h^3(r)h(r_0)} [4h(r)(h(r) + h(r_0))h'(r)k'(r) \\
& \quad - k(r)(3h(r) + 7h(r_0))h'^2(r) + 4h(r)(h(r) + h(r_0))h''(r)] \\
& V(\rho, \rho_0) = \frac{2}{e^{8f(\rho)+2k(\rho)}} \left(2(e^{8f(\rho)} - e^{8f(\rho_0)})f'^2(\rho) + (e^{8f(\rho)} + e^{8f(\rho_0)})(f''(\rho) - k'(\rho)f'(\rho)) \right) \\
& \mathcal{M}_4 = [t, x, r(x), \theta = \tilde{\theta}, \varphi = 2\pi - \tilde{\varphi}, \psi = \pi] \\
& ds_{ind}^2 = \alpha' N e^\phi \left[-dt^2 + (1 + \dot{r}^2)dx^2 + \left(e^{2h} + \frac{1}{4}(1 - a)^2 \right) (d\theta^2 + \sin^2 \theta d\varphi^2) \right]
\end{aligned}$$

$$\begin{aligned}
V_{S^2}(r) &\equiv \frac{1}{4}(1-a(r))^2 + e^{2h(r)} = r \tanh r \\
S_{DBI} &= -T_{D3} \int d^4 \sigma e^{-\phi} \sqrt{g_{ind}} \\
S_{eff} &= 4\pi T_{D3} \mathcal{T}(\alpha' N)^2 \int e^{\phi} r \tanh r \sqrt{1+r^2} dx \\
S_B &= -T_{Dp} \int d^{p+1} \sigma e^{-\phi} \sqrt{-\det(g+2\pi\alpha' F)} + T_{Dp} \int P[C_{p+1}] \\
S_{Dp}^{(F)} &= \frac{T_{Dp}}{2} \int d^{p+1} \sigma \sqrt{-\text{Det}(M)} \bar{\Theta} (1-\Gamma_{Dp}) (\tilde{M}^{-1})^{\alpha\beta} \Gamma_{\beta} D_{\alpha} \Theta \\
D_{\alpha} &= (\partial_{\alpha} x^m) \left(\nabla_m + \frac{1}{16 \cdot 5!} F_{npqrt} \Gamma^{npqrt} (i\sigma_2) \Gamma_m \right), \nabla_m = \partial_m + \frac{1}{4} \omega_m^{\frac{np}{m}} \Gamma_{np} \\
\Gamma_{Dp} &= \frac{\sqrt{-\text{Det}(g)}}{\sqrt{-\text{Det}(g+F)}} \Gamma_{Dp}^{(0)} \otimes (\sigma_3)^{\frac{p+1}{2}} (-i\sigma_2) \sum_q \Gamma^{\alpha_1 \dots \alpha_{2q}} F_{\alpha_1 \alpha_2} \dots F_{\alpha_{2q-1} \alpha_{2q}} \otimes \frac{\sigma_3^q}{2^q q!} \\
\Gamma_{Dp}^{(0)} &= \frac{\epsilon^{\alpha_1 \dots \alpha_{p+1}}}{(p+1)! \sqrt{-\text{Det}(g)}} \Gamma_{\alpha_1 \dots \alpha_{p+1}} \\
ds^2 &= L^2 \left(\frac{\cosh^2(u)}{r^2} (-dt^2 + dr^2) + du^2 + \sinh^2(u) (d\theta^2 + \sin^2 \theta d\phi^2) \right) + L^2 (d\alpha_1^2 + \sin^2 \alpha_1 d\Omega_4^2) \\
F_5 &= 4L^4 \frac{\sinh^2(u) \cosh^2(u)}{r^2} dt \wedge dr \wedge du \wedge d\theta \wedge d\phi \equiv dC_4 \\
C_4 &= 4L^4 \frac{f(u)}{r^2} dt \wedge dr \wedge d\theta \wedge d\phi \\
F &= \frac{q}{r^2} dt \wedge dr + k \sin \theta d\theta \wedge d\phi \\
S_B &= -T_{D3} \int d^4 \sigma \sqrt{-\det(g+F)} + T_{D3} \int P[C_4] \\
q &= \coth(u_{eq}) \sqrt{L^4 \sinh^2(u_{eq}) - k^2} \\
u &= u_{eq} + \delta u, \theta^i = \theta_0^i + \delta \theta^i, F = \frac{q}{r^2} dt \wedge dr + k \sin \theta d\theta \wedge d\phi + f \\
S_B^{(2)} &= \frac{T_{D3}}{4} \int d^4 \sigma \left[\frac{1}{r^2} \left(\frac{1}{L^4 \sinh^4(u_{eq}) + k^2} \left[2kr^2 \sqrt{\sinh^2(u_{eq}) L^4 - k^2} (f_{t\phi} f_{r\theta} - f_{t\theta} f_{r\phi} \right. \right. \right. \\
&\quad \left. \left. + f_{tr} f_{\theta\phi} \right) + L^4 \cosh(u_{eq}) \csc(\theta) \sinh^3(u_{eq}) \left((-f_{\theta\phi}^2 + r^2(f_{t\phi}^2 - f_{r\phi}^2)) \right. \right. \\
&\quad \left. \left. + r^2 \sin^2 \theta (f_{t\theta}^2 - f_{r\theta}^2 + r^2 f_{tr}^2) \right) \right] \\
&\quad \left. - \frac{L^4 \cosh(u_{eq}) \sinh(u_{eq})}{r^2} \left(\csc(\theta) \partial_{\phi} \delta u^2 + \sin(\theta) ((\partial_r \delta u^2 - \partial_t \delta u^2) r^2 + \partial_{\theta} \delta u^2) \right) \right] \\
\hat{g}_{\alpha\beta} &= g_{\alpha\beta} - F_{\alpha\gamma} g^{\gamma\delta} F_{\delta\beta} = \frac{\sinh^4(u_{eq}) L^4 + k^2}{L^2 \sinh^2(u_{eq})} \begin{pmatrix} -\frac{1}{r^2} & 0 & 0 & 0 \\ 0 & \frac{1}{r^2} & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & \sin^2 \theta \end{pmatrix}
\end{aligned}$$

$$S_B^{(2)} = \frac{T_{D3}}{2} \left(\frac{L^4 \cosh(u_{eq}) \sinh^3(u_{eq})}{L^4 \sinh^4(u_{eq}) + k^2} \right) \int d^4 \sigma \left[\sqrt{-\hat{g}} \left(L^2 \hat{g}^{\alpha\beta} \partial_\alpha \delta u \partial_\beta \delta u + \frac{1}{2} \hat{g}^{\alpha\beta} \hat{g}^{\gamma\delta} f_{\alpha\gamma} f_{\beta\delta} \right) \right. \\ \left. + \frac{k \sqrt{\sinh^2(u_{eq}) L^4 - k^2}}{L^4 \cosh(u_{eq}) \sinh^3(u_{eq})} (f_{t\phi} f_{r\theta} - f_{t\theta} f_{r\phi} + f_{tr} f_{\theta\phi}) \right] \\ = \frac{T_{D3}}{2} \left(\frac{L^4 \cosh(u_{eq}) \sinh^3(u_{eq})}{L^4 \sinh^4(u_{eq}) + k^2} \right) \int d^4 \sigma \sqrt{-\hat{g}} \left(L^2 \hat{g}^{\alpha\beta} \partial_\alpha \delta u \partial_\beta \delta u + \frac{1}{2} \hat{g}^{\alpha\beta} \hat{g}^{\gamma\delta} f_{\alpha\gamma} f_{\beta\delta} \right. \\ \left. + \frac{kq}{2L^4 \sinh^2(2u_{eq})} \hat{\epsilon}^{\alpha\beta\gamma\delta} f_{\alpha\beta} f_{\gamma\delta} \right),$$

$$S_F^{(2)} = \frac{T_{D3}}{2} \left(\frac{L^4 \cosh(u_{eq}) \sinh^3(u_{eq})}{L^4 \sinh^4(u_{eq}) + k^2} \right) \int d^4 \sigma \sqrt{\hat{g}} \bar{\Theta} \hat{\Gamma}^\alpha \hat{\nabla}_\alpha \Theta$$

$$\langle W \rangle \simeq e^{-A}$$

$$\mathcal{S}_{\mathcal{A}} = -\text{Tr}_{\mathcal{A}}(\rho_{\mathcal{A}} \ln \rho_{\mathcal{A}})$$

$$\mathcal{S}_{\mathcal{A}} = \frac{2\pi \text{Area}(\gamma_{\mathcal{A}})}{\kappa^2}$$

$$\kappa_{(4)}^2 \mathcal{L} = R - 2\Lambda - \frac{1}{4} \text{Tr}(F_{\mu\nu} F^{\mu\nu})$$

$$F_{\mu\nu}^a = \partial_\mu A_\nu^a - \partial_\nu A_\mu^a + g_{YM} \epsilon^{abc} A_\mu^b A_\nu^c$$

$$G_{\mu\nu} = R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{3}{R^2} g_{\mu\nu} + \frac{1}{2} \text{Tr}[F_{\mu\gamma} F_\nu^\gamma] - \frac{g_{\mu\nu}}{8} \text{Tr}[F_{\gamma\rho} F^{\gamma\rho}]$$

$$D_\mu F^{\mu\nu} = 0$$

$$ds^2 = -M(r)\sigma(r)^2 dt^2 + \frac{1}{M(r)} dr^2 + r^2 h(r)^2 dx^2 + r^2 h(r)^{-2} dy^2$$

$$A=\phi(r)\tau^3dt+\omega(r)\tau^1dx$$

$$\omega(r)=0$$

$$h(r)=1$$

$$\sigma(r)=1$$

$$\phi(r)=\mu\left(1-\frac{r_h}{r}\right)$$

$$M(r)=r^2+\frac{\mu^2r_h^2}{r^2}-\left(\frac{\mu^2}{8}+r_h^2\right)\frac{r_h}{r}$$

$$M=M_1(r-r_h)+M_2(r-r_h)^2+\cdots$$

$$h=h_0+h_2(r-r_h)^2+\cdots$$

$$\sigma=\sigma_0+\sigma_1(r-r_h)+\sigma_2(r-r_h)^2+\cdots$$

$$\omega=\omega_0+\omega_2(r-r_h)^2+\omega_3(r-r_h)^3+\cdots$$

$$\phi=\phi_1(r-r_h)+\phi_2(r-r_h)^2+\cdots$$

$$M'=\frac{3r}{\hat{R}^2}-\frac{1}{8\sigma^2}\left(\frac{g_{YM}^2\phi^2\omega^2}{rh^2M}+r\phi'^2\right)-M\left(\frac{1}{r}+\frac{rh'^2}{h^2}+\frac{\omega'^2}{8rh^2}\right)$$

$$\sigma'=\frac{\sigma}{h^2}\left(rh'^2+\frac{\omega'^2}{8r}\right)+\frac{g_{YM}^2\phi^2\omega^2}{8rM^2h^2\sigma};$$

$$h''=\frac{1}{8r^2h}\left(-\omega'^2+\frac{g_{YM}^2\phi^2\omega^2}{M^2\sigma^2}\right)-h'\left(\frac{2}{r}-\frac{h'}{h}+\frac{M'}{M}+\frac{\sigma'}{\sigma}\right);$$

$$\omega''=-\frac{g_{YM}^2\phi^2\omega}{M^2\sigma^2}+\omega'\left(\frac{2h'}{h}-\frac{M'}{M}-\frac{\sigma'}{\sigma}\right);$$

$$\phi''=\frac{g_{YM}^2\phi\omega^2}{r^2h^2M}-\phi'\left(\frac{2}{r}-\frac{\sigma'}{\sigma}\right)$$

$$\begin{aligned}
& \sigma \rightarrow \lambda \sigma, \phi \rightarrow \lambda \phi \\
& \omega \rightarrow \lambda \omega, h \rightarrow \lambda h \\
& M \rightarrow \lambda^{-2} M, \sigma \rightarrow \lambda \sigma, g_{YM} \rightarrow \lambda^{-1} g_{YM}, \hat{R} \rightarrow \lambda \hat{R} \\
& M \rightarrow \lambda^2 M, r \rightarrow \lambda r, \phi \rightarrow \lambda \phi, \omega \rightarrow \lambda \omega \\
& T = \frac{M_1 \sigma_0}{2\pi} = \frac{1}{16\pi} (24\sigma_0^2 - \phi_1^2) r_h \\
& S = \frac{2\pi}{\kappa_{(4)}^2} A_h = \frac{2\pi^2 V T^2}{\kappa_{(4)}^2} \frac{12^2}{(24\sigma_0^2 - \phi_1^2)^2} \\
& \tilde{S}_{bulk} = - \int dx dy dr \sqrt{-g} \mathcal{L} \\
& G_{yy} = \frac{r^2}{2h^2} (\kappa_{(4)}^2 \mathcal{L} - R) \\
& G_\mu^\mu = -R = G_r^r + G_t^t + G_x^x + \frac{1}{2} (\kappa_{(4)}^2 \mathcal{L} - R) \\
& \mathcal{L} = \frac{2}{r^2 \sigma \kappa_{(4)}^2} \left[\frac{r^3 M \sigma}{h} \left(\frac{h}{r} \right)' \right]' \\
& \tilde{S}_{bulk} = - \int dx dy dr \sqrt{-g} \mathcal{L} = - \frac{2V}{\kappa_{(4)}^2} \left[\frac{r^3 M \sigma}{h} \left(\frac{h}{r} \right)' \right]_{r=r_\infty} \\
& \tilde{S}_{GH} = - \frac{1}{\kappa_{(4)}^2} \int dx dy \sqrt{-g_\infty} \nabla_\mu n^\mu = - \frac{V}{\kappa_{(4)}^2} r^2 \sigma \left[\frac{M'}{2} + M \left(\frac{\sigma'}{\sigma} + \frac{2}{r} \right) \right]_{r=r_\infty} \\
& \tilde{S}_{ct} = \frac{1}{\kappa_{(4)}^2} \int dx dy \sqrt{-g_\infty} = \frac{V}{\kappa_{(4)}^2} [r^2 \sqrt{M} \sigma]_{r=r_\infty}
\end{aligned}$$

$$\begin{aligned}
\Omega &= \lim_{r_\infty \rightarrow \infty} \tilde{S}_{\text{on-shell}} \\
&= \lim_{r_\infty \rightarrow \infty} (\tilde{S}_{bulk} + \tilde{S}_{GH} + \tilde{S}_{ct})
\end{aligned}$$

$$ds^2 = -M(r)dt^2 + r^2 h(r)^2 (dx^2 + dy^2) + \frac{dr^2}{M(r)}$$

$$\begin{aligned}
& A = \phi(r) \tau^3 dt + \omega(r) (\tau^1 dx + \tau^2 dy) \\
h'' &= -\frac{h}{2} \left[\frac{1}{r^2} - \frac{3}{\hat{R}^2 M} + \frac{M'}{rM} + \frac{\phi'^2}{8M} + \frac{\omega'^2}{4r^2 h^2} \right] - \frac{h'}{2} \left[\frac{6}{r} + \frac{h'}{h} + \frac{M'}{M} \right] - \frac{g_{YM}^2 \omega^2}{8r^2 h M} \left[\frac{\phi^2}{M} + \frac{\omega^2}{2r^2 h^2} \right] \\
M'' &= \frac{3}{\hat{R}^2} + \frac{M}{r} \left[-\frac{M'}{M} + \frac{1}{r} + \frac{\omega'^2}{4r h^2} \right] - \frac{h'}{h} \left[M' - \frac{h'}{h} - \frac{2}{r} \right] + \frac{3}{8} \phi'^2 + \frac{g_{YM}^2 \omega^2}{4r^2 h^2} \left[\frac{\phi^2}{M} + \frac{3\omega^2}{2r^2 h^2} \right] \\
\omega'' &= \frac{g_{YM}^2 \omega}{M} \left[\frac{\omega^2}{r^2 h^2} - \frac{\phi^2}{M} \right] - \frac{M' \omega'}{M} \\
\phi'' &= \frac{2g_{YM}^2 \phi \omega^2}{r^2 h^2 M} - 2\phi' \left[\frac{1}{r} + \frac{h'}{h} \right] \\
0 &= -\frac{3}{\hat{R}^2} + \frac{M}{r^2} \left[1 - \frac{\omega'^2}{4h^2} + \frac{M'}{M} r \right] + \frac{h'}{h} \left[M \left(\frac{2}{r} + \frac{h'}{h} \right) + M' \right] + \frac{1}{8} \phi'^2
\end{aligned}$$

$$\omega \rightarrow \lambda \omega, h \rightarrow \lambda h$$

$$M \rightarrow \lambda^{-2} M, \phi \rightarrow \frac{\phi}{\lambda}, \hat{R} \rightarrow \lambda \hat{R}, g_{YM} \rightarrow \frac{g_{YM}}{\lambda}$$

$$M \rightarrow \lambda^2 M, h \rightarrow \frac{h}{\lambda}, \phi \rightarrow \lambda \phi, r \rightarrow \lambda r$$

$$M = M_1(r - r_h) + M_2(r - r_h)^2 + \dots$$

$$h = h_0 + h_1(r - r_h) + h_2(r - r_h)^2 + \dots$$

$$\omega = \omega_0 + \omega_1(r - r_h) + \omega_2(r - r_h)^2 + \dots$$

$$\phi = \phi_1(r - r_h) + \phi_2(r - r_h)^2 + \dots$$



$$\begin{aligned}
M &= r^2 + 2h_1^b r + (h_1^b)^2 + \frac{M_1^b}{r} + \frac{-8h_1^b M_1^b + \rho^2 + 2(\omega_1^b)/3}{8r^2} + \dots \\
h &= 1 + \frac{h_1^b}{r} - \frac{(\omega_1^b)^2}{48r^4} + \dots \\
\omega &= \frac{\omega_1^b}{r} - \frac{h_1^b \omega_1^b}{r^2} + \dots \\
\phi &= \mu + \frac{\rho}{r} - \frac{\rho h_1^b}{r^2} + \dots \\
T &= \frac{M_1}{2\pi} \\
S &= \frac{2\pi}{\kappa_{(4)}^2} A_h = \frac{2\pi}{\kappa_{(4)}^2} r_h^2 h_0^2 \\
\Omega &= \frac{VM_1^b}{\kappa_{(4)}^2} \\
\mathcal{S}_{\mathcal{A}} &= \frac{2\pi}{\kappa_{(d+1)}^2} \int_{\gamma_{\mathcal{A}}} d^{(d-1)} \sigma \sqrt{g_{\text{ind}}^{(d-1)}} \\
ds_{d+1}^2 &= -g_{tt}(r)dt^2 + g_{x_i x_i}(r)dx_i^2 + g_{rr}(r)dr^2, i = 1 \dots d-1 \\
\mathcal{S}_{\mathcal{A}} &= \frac{2\pi\Lambda}{\kappa_{(d+1)}^2} \int d\zeta \sqrt{g_{x_2 x_2}(r) \dots g_{x_{d-1} x_{d-1}}(r)} \sqrt{g_{rr}(r)r'^2 + g_{x_1 x_1}(r)x_1'^2} \\
f^2(r) &= g_{\chi\chi}(r)g_{x_1 x_1}(r), \eta^2(r) = g_{\chi\chi}(r)g_{rr}(r) \\
\mathcal{S}_{\mathcal{A}} &= \frac{2\pi\Lambda}{\kappa_{(d+1)}^2} \int d\zeta \sqrt{\eta^2(r)r'^2 + f^2(r)x_1'^2} \\
x_1'(\zeta) &= \pm \frac{f(r_0)\eta(r)}{f(r)} \frac{r'(\zeta)}{\sqrt{f^2(r) - f^2(r_0)}} \\
L &= 2 \int_{r_0}^{\infty} dr \frac{dx_1}{dr} = 2 \int_{r_0}^{\infty} dr \frac{\eta(r)}{f(r)} \frac{f(r_0)}{\sqrt{f^2(r) - f^2(r_0)}} \\
\mathcal{S}_{\mathcal{A}}(r_0) &= 2 \frac{2\pi\Lambda}{\kappa_{(d+1)}^2} \int_{r_0}^{\infty} dr \frac{f(r)\eta(r)}{\sqrt{f(r)^2 - f(r_0)^2}} \\
\mathcal{S}_{\mathcal{A}disc} &= 2 \frac{2\pi\Lambda}{\kappa_{(d+1)}^2} \int_{r_{\min}}^{\infty} dr \eta(r) \\
\Delta\mathcal{S}_{\mathcal{A}} &= \frac{4\pi\Lambda}{\kappa_{(d+1)}^2} \left(\int_{r_0}^{\infty} dr \frac{f(r)\eta(r)}{\sqrt{f(r)^2 - f(r_0)^2}} - \int_{r_{\min}}^{\infty} dr \eta(r) \right) \\
f_p^2(r) &= g_{yy}g_{xx} = r^4, \eta_p^2(r) = g_{yy}g_{rr} = \frac{r^2}{h^2 N} \\
\Delta\mathcal{S}_{\mathcal{A}} &= \frac{4\pi\Lambda}{\kappa_{(4)}^2} \left(\int_{r_0}^{\infty} dr \frac{r^3}{h\sqrt{N}\sqrt{r^4 - r_0^4}} - \int_{r_{\min}}^{\infty} dr \frac{r}{h\sqrt{N}} \right) \\
\mathcal{S}_{\mathcal{A}}(r_0) &= \frac{4\pi\Lambda}{\kappa_{(4)}^2} \int_{r_0}^{\mathcal{R}} dr \frac{r^3}{h\sqrt{N}\sqrt{r^4 - r_0^4}} = \mathcal{S}_{\mathcal{A}} + \frac{4\pi\Lambda}{\kappa_{(d+1)}^2} \mathcal{R} \\
f_{p+ip}^2(r) &= g_{yy}g_{xx} = r^4 h^4, \eta_{p+ip}^2(r) = g_{yy}g_{rr} = \frac{r^2 h^2}{M} \\
\Delta\mathcal{S}_{\mathcal{A}} &= \frac{4\pi\Lambda}{\kappa_{(4)}^2} \left(\int_{r_0}^{\infty} dr \frac{r^3 h^3}{\sqrt{M}\sqrt{r^4 h^4 - r_0^4 h(r_0)^4}} - \int_{r_{\min}}^{\infty} dr \frac{r h}{\sqrt{M}} \right) \\
\mathcal{S}_{\mathcal{A}}(r_0) &= \frac{4\pi\Lambda}{\kappa_{(4)}^2} \int_{r_0}^{\mathcal{R}} dr \frac{r^3 h^3}{\sqrt{M}\sqrt{r^4 h^4 - r_0^4 h(r_0)^4}} = \mathcal{S}_{\mathcal{A}} + \frac{4\pi\Lambda}{\kappa_{(4)}^2} \mathcal{R} \\
\left[-\frac{d}{dr} \left(P(r, r_0) \frac{d}{dr} \right) + U(r, r_0) \right] \Phi(r) &= \omega^2 Q(r, r_0) \Phi(r), r_0 \leq r < \infty
\end{aligned}$$

$$y = \int_{r_0}^r \sqrt{\frac{Q}{P}} dr, \Phi(r) = (PQ)^{-\frac{1}{4}} \Psi(y)$$

$$\left[-\frac{d^2}{dy^2} + V\right] \Psi = \omega^2 \Psi, 0 \leq y \leq y_0$$

$$V = \frac{U}{Q} + \left[(PQ)^{-\frac{1}{4}} \frac{d^2}{dy^2}\right] (PQ)^{\frac{1}{4}}$$

$$= \frac{U}{Q} + \left[\frac{P^{\frac{1}{4}}}{Q^{\frac{3}{4}}} \frac{d}{dr} \left(\sqrt{\frac{P}{Q}} \frac{d}{dr} \right)\right] (PQ)^{\frac{1}{4}}$$

$$\delta x|_{r=\infty} = 0 \Rightarrow \Psi|_{y=y_0} = 0$$

$$\text{Even solutions : } \frac{d\Psi}{dy}\Big|_{y=0} = 0$$

$$\text{Odd solutions : } \Psi|_{y=0} = 0$$

$$ds^2 = \frac{R^2}{z^2} (-dt^2 + dx_i dx_i + dz^2) + R^2 d\Omega_5^2$$

$$\left(\frac{dz_{\text{cl}}}{dx}\right)^2 = \frac{z_0^4 - (z_{\text{cl}})^4}{(z_{\text{cl}})^4}$$

$$x_{\text{cl}}(z) = \pm z_0 \left[\frac{(2\pi)^{\frac{3}{2}}}{2\Gamma\left[\frac{1}{4}\right]^2} - \frac{1}{4} \operatorname{B}\left(\frac{z^4}{z_0^4}; \frac{3}{4}, \frac{1}{2}\right) \right]$$

$$x\text{-gauge : } \left[\partial_t^2 - \frac{z_{\text{cl}}^4(x)}{z_0^4} \partial_x^2 \right] \delta x_m(t,x) = 0$$

$$r\text{-gauge : } \left[\partial_t^2 - \left(1 - \frac{z^4}{z_0^4} \right) \partial_z^2 + \frac{2}{z} \partial_z \right] \delta x_m(t,z) = 0 \quad m = 2,3$$

$$\left[(1-\tilde{z}^4) \partial_{\tilde{z}}^2 - \frac{2}{\tilde{z}} + \xi^2 \right] f(\tilde{z}) = 0, 0 \leq \tilde{z} \leq 1$$

$$f(\tilde{z}) = \sqrt{1 + \xi^2 \tilde{z}^2} F(q)$$

$$q(\tilde{z}) = \pm 2 \int_{\tilde{z}}^1 \frac{t^2}{(1 + (\xi t)^2) \sqrt{1 - t^4}} dt$$

$$\frac{d^2 F}{dq} + \frac{1}{4} \xi^2 (\xi^4 - 1) F = 0, q \in [-q_*, q_*]$$

$$\omega_n z_0 \sqrt{\omega_n^4 z_0^4 - 1} \int_0^1 \frac{t^2 dt}{(1 + w_n^2 z_0^2) \sqrt{1 - t^4}} = \frac{n\pi}{2}, n = 1, 2, \dots$$

$$\tilde{M}^{\alpha\beta}$$

$$= \frac{L^2 \sinh^2(u_{eq})}{k^2 + L^4 \sinh^4(u_{eq})} \begin{pmatrix} -r^2 & \frac{2r^2 \sqrt{L^4 \sinh^2(u_{eq}) - k^2}}{L^2 \sinh(2u_{eq})} \tilde{\Gamma} & 0 & 0 \\ -\frac{2r^2 \sqrt{L^4 \sinh^2(u_{eq}) - k^2}}{L^2 \sinh(2u_{eq})} \tilde{\Gamma} & r^2 & 0 & 0 \\ 0 & 0 & 1 & -\frac{k}{L^2 \sin \theta \sinh^2(u_{eq})} \tilde{\Gamma} \\ 0 & 0 & \frac{k}{L^2 \sin \theta \sinh^2(u_{eq})} \tilde{\Gamma} & \frac{1}{\sin^2 \theta} \end{pmatrix}$$

$$\Gamma_t = \frac{L \cosh(u_k)}{r} \Gamma_{\underline{0}}, \Gamma_r = \frac{L \cosh(u_k)}{r} \Gamma_{\underline{1}}$$

$$\Gamma_{\theta} = L \sinh(u_k) \Gamma_{\underline{2}}, \Gamma_{\phi} = L \sinh(u_k) \sin \theta \Gamma_{\underline{3}}$$



$$\begin{aligned}
\tilde{M}^{\alpha\beta}\Gamma_\beta D_\alpha &= \frac{L\sinh(u_{eq})}{\sqrt{k^2 + L^4\sinh^4(u_{eq})}} \left[-re^{2R_e\tilde{\Gamma}}\Gamma_{\underline{0}}D_t + re^{2R_e\tilde{\Gamma}}\Gamma_{\underline{1}}D_r + e^{2R_m\tilde{\Gamma}}\Gamma_{\underline{2}}D_\theta + \frac{1}{\sin\theta}e^{2R_m\tilde{\Gamma}}\Gamma_{\underline{3}}D_\phi \right] \\
R_e &= -\frac{1}{2}\sinh^{-1}\left(\frac{\sqrt{L^4\sinh^2(u_{eq}) - k^2}}{\sqrt{L^4\sinh^4(u_{eq}) + k^2}}\right)\Gamma_{01}, R_m = \frac{1}{2}\arcsin\left(\frac{k}{\sqrt{L^4\sinh^4(u_{eq}) + k^2}}\right)\Gamma_{23} \\
\hat{e}^{\underline{0}} &= \frac{\sqrt{\sinh^4(u_{eq})L^4 + k^2}}{Lr\sinh(u_{eq})}dt, \hat{e}^{\underline{1}} = \frac{\sqrt{\sinh^4(u_{eq})L^4 + k^2}}{Lr\sinh(u_{eq})}dr \\
\hat{e}^{\underline{2}} &= \frac{\sqrt{\sinh^4(u_{eq})L^4 + k^2}}{L\sinh(u_{eq})}d\theta, \hat{e}^{\underline{3}} = \frac{\sqrt{\sinh^4(u_{eq})L^4 + k^2}}{L\sinh(u_{eq})}\sin\theta d\phi \\
\tilde{M}^{\alpha\beta}\Gamma_\beta D_\alpha &= e^{\mathcal{R}\tilde{\Gamma}}[\hat{\Gamma}^\alpha e^{\mathcal{R}\tilde{\Gamma}}D_\alpha e^{-\mathcal{R}\tilde{\Gamma}}]e^{\mathcal{R}\tilde{\Gamma}} \\
\nabla_\alpha d\sigma^\alpha &= d + \frac{1}{4}\omega_{\mu\nu}\Gamma_{\underline{\mu\nu}} + \frac{1}{4}\omega^{ij}\Gamma_{\underline{ij}} + \frac{1}{2}\sinh(u_k)e^\mu_{\underline{\mu}}\Gamma_{\underline{\mu 4}} + \frac{1}{2}\cosh(u_k)e^{-\Gamma_{\underline{i}}} \\
e^{\mathcal{R}\tilde{\Gamma}}\nabla_\alpha d\sigma^\alpha e^{-\mathcal{R}\tilde{\Gamma}} &= \hat{\nabla}_\alpha d\sigma^\alpha + \frac{1}{2}\sinh(u_k)e^\mu_{\underline{\mu}}\Gamma_{\underline{\mu 4}}e^{-2R_e\tilde{\Gamma}} + \frac{1}{2}\cosh(u_k)e^{i-\Gamma_{\underline{i 4}}}e^{-2R_m\tilde{\Gamma}}, \\
\tilde{M}^{\alpha\beta}\Gamma_\beta \nabla_\alpha &= e^{\mathcal{R}\tilde{\Gamma}}\left[\hat{\Gamma}^\alpha \hat{\nabla}_\alpha + \frac{L\sinh(u_{eq})}{\sqrt{L^4\sinh^4(u_{eq}) + k^2}}\Gamma_4(\sinh(u_{eq})e^{-2R_e\tilde{\Gamma}} + \cosh(u_{eq})e^{-2R_m\tilde{\Gamma}})\right]e^{\mathcal{R}\tilde{\Gamma}} \\
F_5\Gamma_\alpha d\sigma^\alpha \otimes (i\sigma_2) &= -4\Gamma_{\underline{01234}} \otimes (i\sigma_2) \left(\cosh(u_{eq})e^\mu_{\underline{\mu}}\Gamma_{\underline{\mu}} + \sinh(u_{eq})e^i_{\underline{i}}\right)(1 + \Gamma^{11}) \\
\tilde{M}^{\alpha\beta}\Gamma_\beta F_5\Gamma_\alpha \otimes (i\sigma_2) &= -16e^{\mathcal{R}\tilde{\Gamma}}\Gamma_{\underline{01234}} \\
&\otimes (i\sigma_2) \frac{L\sinh(u_{eq})}{\sqrt{L^4\sinh^4(u_{eq}) + k^2}}(\cosh(u_{eq})e^{-2R_m\tilde{\Gamma}} + \sinh(u_{eq})e^{-2R_e\tilde{\Gamma}})e^{\mathcal{R}\tilde{\Gamma}} \\
\tilde{M}^{\alpha\beta}\Gamma_\beta D_\alpha &= e^{\mathcal{R}\tilde{\Gamma}}\left[\hat{\Gamma}^\alpha \hat{\nabla}_\alpha \right. \\
&\quad \left. + (1 - \Gamma_{D3}^{(0)})\frac{L\sinh(u_{eq})}{\sqrt{L^4\sinh^4(u_{eq}) + k^2}}\Gamma_{\underline{4}}(\cosh(u_{eq})e^{-2R_m\tilde{\Gamma}} + \sinh(u_{eq})e^{-2R_e\tilde{\Gamma}})\right]e^{\mathcal{R}\tilde{\Gamma}} \\
\Gamma_{D3}^{(0)} &= \Gamma_{\underline{0123}} \otimes (i\sigma_2) \\
\Gamma_{D3} &= -\frac{\epsilon^{\alpha_1\alpha_2\alpha_3\alpha_4}\Gamma_{\alpha_1\alpha_2\alpha_3\alpha_4}}{(p+1)!\sqrt{\det(g+F)}} \otimes (i\sigma_2) \times \sum_q \Gamma^{\alpha_1\cdots\alpha_{2q}}F_{\alpha_1\alpha_2}\cdots F_{\alpha_{2q-1}\alpha_{2q}} \otimes \frac{(\sigma_3)^q}{q!2^q} \\
\Gamma_{D3} &= -\Gamma_{D3}^{(0)}\frac{L^4\cosh(u_{eq})\sinh^3(u_{eq})}{L^4\sinh^4(u_{eq})+k^2}\left[1 + \left(\frac{k}{L^2\sinh^2(u_{eq})}\Gamma_{\underline{23}} - \frac{\sqrt{L^4\sinh^2(u_{eq})-k^2}}{L^2\cosh(u_{eq})\sinh(u_{eq})}\Gamma_{\underline{01}}\right) \otimes \sigma_3\right] \\
\bar{\Theta}\Gamma_{D3} &= -\bar{\Theta}e^{\mathcal{R}\tilde{\Gamma}}\Gamma_{D3}^{(0)}e^{-\mathcal{R}\tilde{\Gamma}}
\end{aligned}$$

$$\mathcal{L}_F = \bar{\Theta} e^{\mathcal{R}\tilde{\Gamma}} \left(1 + \Gamma_{D3}^{(0)} \right) \left[\hat{\Gamma}^\alpha \hat{\nabla}_\alpha + \left(1 - \Gamma_{D3}^{(0)} \right) \frac{L \sinh(u_{eq})}{\sqrt{L^4 \sinh^4(u_{eq}) + k^2}} \Gamma_\pm (\cosh(u_{eq}) e^{-2R_m \tilde{\Gamma}} + \sinh(u_{eq}) e^{-2R_e \tilde{\Gamma}}) \right] e^{\mathcal{R}\tilde{\Gamma}} \Theta.$$

$$\mathcal{L}_F = \bar{\Theta}' \left(1 + \Gamma_{D3}^{(0)} \right) \hat{\Gamma}^\alpha \hat{\nabla}_\alpha \Theta'.$$

$$\tilde{\Gamma}\Theta' = \Theta'$$

$$\bar{\Theta}(1-\Gamma_{D3})\tilde{M}^{\alpha\beta}\Gamma_\beta D_\alpha\Theta=\bar{\Theta}\hat{\Gamma}^\alpha\hat{\nabla}_\alpha\Theta$$

$$S_F^{(2)} = \frac{T_{D3}}{2} \left(\frac{L^4 \cosh(u_{eq}) \sinh^3(u_{eq})}{L^4 \sinh^4(u_{eq}) + k^2} \right) \int d^4 \sigma \sqrt{\hat{g}} \bar{\Theta} \hat{\Gamma}^\alpha \hat{\nabla}_\alpha \Theta$$

$$ds^2 = \frac{g_{\mu\nu}(x^\rho,z)dx^\mu dx^\nu + dz^2}{z^2} \equiv g_{\alpha\beta}^{5D} dx^\alpha dx^\beta$$

$$R_{\alpha\beta}-\frac{1}{2}g_{\alpha\beta}^{5D}R-6g_{\alpha\beta}^{5D}=0$$

$$g_{\mu\nu}(x^\rho,z)=\eta_{\mu\nu}+g_{\mu\nu}^{(4)}(x^\rho)z^4+\mathcal{O}(z^6)\left\langle T_{\mu\nu}(x^\rho)\right\rangle=\frac{N_c^2}{2\pi^2}g_{\mu\nu}^{(4)}(x^\rho)$$

$$\varepsilon(\tau)\equiv N_c^2\cdot\frac{3}{8}\pi^2\cdot T_{eff}^4$$

$$\Delta p_L \equiv 1 - \frac{p_L}{\varepsilon/3} \sim 0.7$$

$$\langle T_{--}(x^-)\rangle \propto \mu \delta(x^-) \text{ o m\'as generalmente } \langle T_{--}(x^-)\rangle \propto f(x^-)$$

$$ds^2 = \frac{-dx^-dx^+ + f(x^-)z^4dx^{-2} + dx_\perp^2}{z^2} + \frac{dz^2}{z^2}$$

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Apéndice I.

1. Supergravedad cuántica, supermembranas, agujeros negros cuánticos y supersimetrías en campos cuánticos relativistas. Formalización matemática.

$$\begin{aligned}
& \Gamma^{012}\varepsilon = -\varepsilon, \Gamma^{013456}\varepsilon = \varepsilon \\
& \delta\psi_\mu \equiv \nabla_\mu \epsilon + \frac{1}{288} \left(\Gamma_\mu^{\nu\rho\lambda\sigma} - 8\delta_\mu^\nu \Gamma^{\rho\lambda\sigma} \right) F_{\nu\rho\lambda\sigma} \epsilon \\
& \Gamma^{012345678910} = \mathbb{1} \\
& \Gamma^{0178910}\varepsilon = -\varepsilon \\
& ds_{11}^2 = e^{2A_0} [-dt^2 + dy^2 + e^{-3A_0} (-\partial_z w)^{-\frac{1}{2}} d\vec{u} \cdot d\vec{u} + e^{-3A_0} (-\partial_z w)^{\frac{1}{2}} d\vec{v} \cdot d\vec{v} \\
& \quad + (-\partial_z w) (dz + (\partial_z w)^{-1} (\vec{\nabla}_{\vec{u}} w) \cdot d\vec{u})^2] \\
& e^0 = e^{A_0} dt, e^1 = e^{A_0} dy, e^2 = e^{A_0} (-\partial_z w)^{\frac{1}{2}} (dz + (\partial_z w)^{-1} (\vec{\nabla}_{\vec{u}} w) \cdot d\vec{u}) \\
& e^{i+2} = e^{-\frac{1}{2}A_0} (-\partial_z w)^{-\frac{1}{4}} du_i, e^{i+6} = e^{-\frac{1}{2}A_0} (-\partial_z w)^{\frac{1}{4}} dv_i, i = 1, 2, 3, 4 \\
& C^{(3)} = -e^0 \wedge e^1 \wedge e^2 + \frac{1}{3!} \epsilon_{ijk\ell} ((\partial_z w)^{-1} (\partial_{u_\ell} w) du^i \wedge du^j \wedge du^k - (\partial_{v_\ell} w) dv^i \wedge dv^j \wedge dv^k) \\
& \mathcal{L}_u \equiv \nabla_{\vec{u}} \cdot \nabla_{\vec{u}}, \mathcal{L}_v \equiv \nabla_{\vec{v}} \cdot \nabla_{\vec{v}} \\
& \mathcal{L}_v G_0 = (\mathcal{L}_u G_0) (\partial_z \partial_z G_0) - (\nabla_{\vec{u}} \partial_z G_0) \cdot (\nabla_{\vec{u}} \partial_z G_0) \\
& w = \partial_z G_0, e^{-3A_0} (-\partial_z w)^{\frac{1}{2}} = \mathcal{L}_v G_0 \\
& e^{-3A_0} (\partial_z w)^{-\frac{1}{2}} - (\partial_z w)^{-1} (\nabla_{\vec{u}} w) \cdot (\nabla_{\vec{u}} w) = -\mathcal{L}_u G_0 \\
& ds_{11}^2 = e^{2A_0} [-dt^2 + dy^2 + (-\partial_z w) (dz + (\partial_z w)^{-1} (\partial_u w) du)^2 \\
& \quad + e^{-3A_0} (-\partial_z w)^{-\frac{1}{2}} (du^2 + u^2 d\Omega_3^2) + e^{-3A_0} (-\partial_z w)^{\frac{1}{2}} (dv^2 + v^2 d\Omega_3'^2)] \\
& e^0 = e^{A_0} dt, e^1 = e^{A_0} dy, e^2 = e^{A_0} (-\partial_z w)^{\frac{1}{2}} (dz + (\partial_z w)^{-1} (\partial_u w) du) \\
& e^3 = e^{-\frac{1}{2}A_0} (-\partial_z w)^{-\frac{1}{4}} du, e^4 = e^{-\frac{1}{2}A_0} (-\partial_z w)^{\frac{1}{4}} dv \\
& e^{i+4} = e^{-\frac{1}{2}A_0} (-\partial_z w)^{-\frac{1}{4}} \sigma_i, e^{i+7} = e^{-\frac{1}{2}A_0} (-\partial_z w)^{\frac{1}{4}} \tilde{\sigma}_i, i = 1, 2, 3 \\
& C^{(3)} = -e^0 \wedge e^1 \wedge e^2 + (\partial_z w)^{-1} (u^3 \partial_u w) \text{Vol}(S^3) + (v^3 \partial_v w) \text{Vol}(S'^3) \\
& G_0 = -\frac{1}{2} z^2 \hat{g}_2(u, v) + z \hat{g}_1(u, v) + \hat{g}_0(u, v) \\
& \mathcal{L}_{\vec{v}} \hat{g}_2 + \hat{g}_2 \mathcal{L}_{\vec{u}} \hat{g}_2 - 2 (\vec{\nabla}_{\vec{u}} \hat{g}_2)^2 = 0, \\
& \mathcal{L}_{\vec{v}} \hat{g}_1 + \hat{g}_2 \mathcal{L}_{\vec{u}} \hat{g}_1 - 2 (\vec{\nabla}_{\vec{u}} \hat{g}_1) \cdot (\vec{\nabla}_{\vec{u}} \hat{g}_2) = 0, \\
& \mathcal{L}_{\vec{v}} \hat{g}_0 + \hat{g}_2 \mathcal{L}_{\vec{v}} \hat{g}_0 - (\vec{\nabla}_{\vec{u}} \hat{g}_1)^2 = 0. \\
& \mathcal{L}_{\vec{v}} \hat{g}_2 - \hat{g}_2^3 \mathcal{L}_{\vec{u}} (\hat{g}_2^{-1}) = 0 \\
& \hat{g}_2 = \frac{h_2(\vec{v})}{h_1(\vec{u})} \\
& G_0 = -\frac{1}{2} z^2 \frac{h_2(\vec{v})}{h_1(\vec{u})} + \hat{g}_0(u, v) \\
& \frac{1}{h_1(\vec{u})} \mathcal{L}_{\vec{u}} \hat{g}_0 + \frac{1}{h_2(\vec{v})} \mathcal{L}_{\vec{v}} \hat{g}_0 = 0 \\
& w = \partial_z G_0 = -z \frac{h_2(\vec{v})}{h_1(\vec{u})} \\
& e^{-A_0} e^2 = (-\partial_z w)^{\frac{1}{2}} (dz + (\partial_z w)^{-1} (\vec{\nabla}_{\vec{u}} w) \cdot d\vec{u}) = \left(\frac{h_2(\vec{v})}{h_1(\vec{u})} \right)^{\frac{1}{2}} \left[dz - \frac{z}{h_1(\vec{u})} (\vec{\nabla}_{\vec{u}} h_1(\vec{u})) \cdot d\vec{u} \right] \\
& = (h_1(\vec{u}) h_2(\vec{v}))^{\frac{1}{2}} \left[\frac{dz}{h_1(\vec{u})} - \frac{z}{(h_1(\vec{u}))^2} (\vec{\nabla}_{\vec{u}} h_1(\vec{u})) \cdot d\vec{u} \right] = (h_1(\vec{u}) h_2(\vec{v}))^{\frac{1}{2}} d\hat{z}
\end{aligned}$$



$$\begin{aligned}
\hat{z} &\equiv \frac{z}{h_1(\vec{u})} \\
\mathcal{L}_{\vec{u}}\hat{g}_0 &= -h_0(\vec{u}, \vec{v})h_1(\vec{u}), \mathcal{L}_{\vec{v}}\hat{g}_0 = h_0(\vec{u}, \vec{v})h_2(\vec{v}) \\
\hat{g}_0 &= f_2(\vec{v})h_1(\vec{u}) - f_1(\vec{u})h_2(\vec{v}), \text{ donde } \mathcal{L}_{\vec{u}}f_1 = h_1^2, \mathcal{L}_{\vec{v}}f_2 = h_2^2 \\
G_0 &= -\frac{1}{2}z^2 \frac{h_2(\vec{v})}{h_1(\vec{u})} + f_2(\vec{v})h_1(\vec{u}) - f_1(\vec{u})h_2(\vec{v}) \\
w &= -z \frac{h_2(\vec{v})}{h_1(\vec{u})}, e^{-3A_0} \left(\frac{h_2(\vec{v})}{h_1(\vec{u})} \right)^{\frac{1}{2}} = h_1(\vec{u})h_2^2(\vec{v}) \Rightarrow e^{-2A_0} = h_1(\vec{u})h_2(\vec{v}) \\
ds_{11}^2 &= (h_1(\vec{u})h_2(\vec{v}))^{-1}(-dt^2 + dy^2) + d\hat{z}^2 + h_1(\vec{u})d\vec{u} \cdot d\vec{u} + h_2(\vec{v})d\vec{v} \cdot d\vec{v} \\
ds_{11}^2 &= e^{2A}(\hat{f}_1^2 ds_{AdS_3}^2 + \hat{f}_2^2 ds_{S^3}^2 + \hat{f}_3^2 ds_{S^3}^2 + h_{ij}d\sigma^i d\sigma^j) \\
C^{(3)} &= b_1 \hat{e}^{012} + b_2 \hat{e}^{345} + b_3 \hat{e}^{678} \\
h_{ij}d\sigma^i d\sigma^j &= \frac{\partial_w h \partial_{\bar{w}} h}{h^2} |dw|^2 \\
\partial_w \partial_{\bar{w}} h &= 0 \\
w = \xi + i\rho &\Rightarrow \partial_w = \frac{1}{2}(\partial_\xi - i\partial_\rho), \partial_{\bar{w}} = \frac{1}{2}(\partial_\xi + i\partial_\rho) \\
\partial_{\bar{w}}(-\tilde{h} + ih) &= 0 \\
-\tilde{h} + ih &= \beta w = \beta(\xi + i\rho) \\
h_{ij}d\sigma^i d\sigma^j &= \frac{d\xi^2 + d\rho^2}{4\rho^2} \\
\partial_w G &= \frac{1}{2}(G + \bar{G})\partial_w \log(h) \\
\partial_\xi g_1 + \partial_\rho g_2 &= 0, \partial_\xi g_2 - \partial_\rho g_1 = -\frac{1}{\rho}g_1 \\
\partial_w \Phi = \bar{G}\partial_w h &\Leftrightarrow \partial_\xi \Phi = -\beta g_2, \partial_\rho \Phi = \beta g_1 \\
\left(\partial_\xi^2 + \partial_\rho^2 - \frac{1}{\rho}\partial_\rho \right) \Phi & \\
\partial_\xi \tilde{\Phi} = -\frac{\beta}{\rho}g_1 = -\frac{1}{\rho}\partial_\rho \Phi, \partial_\rho \tilde{\Phi} = -\frac{\beta}{\rho}g_2 = \frac{1}{\rho}\partial_\xi \Phi & \\
\partial_\xi^2 \tilde{\Phi} + \frac{1}{\rho}\partial_\rho(\rho\partial_\rho \tilde{\Phi}) & \\
W_\pm \equiv |G \pm i|^2 + \gamma^{\pm 1}(G\bar{G} - 1) & \\
\gamma(G\bar{G} - 1) \geq 0 & \\
\gamma > 0, |G| \geq 1 & \\
c_1 = \gamma^{1/2} + \gamma^{-1/2} > 0, c_2 = -\gamma^{1/2} < 0, c_3 = -\gamma^{-1/2} < 0, \sigma = +1 & \\
\hat{f}_1^{-2} = \gamma^{-1}(\gamma + 1)^2(G\bar{G} - 1), \hat{f}_2^{-2} = W_+, \hat{f}_3^{-2} = W_- & \\
e^{6A} = h^2(G\bar{G} - 1)W_+W_- = \gamma(\gamma + 1)^{-2}h^2\hat{f}_1^{-2}\hat{f}_2^{-2}\hat{f}_3^{-2} & \\
b_1 = \frac{v_1}{c_1^3} \left[\frac{h(G + \bar{G})}{(G\bar{G} - 1)} + \gamma^{-1}(\gamma + 1)^2\Phi - (\gamma - \gamma^{-1})\tilde{h} \right] & \\
b_2 = \frac{v_2}{c_2^3} \left[-\frac{h(G + \bar{G})}{W_+} + (\Phi - \tilde{h}) \right], b_3 = \frac{v_3}{c_3^3} \left[\frac{h(G + \bar{G})}{W_-} - (\Phi + \tilde{h}) \right] & \\
ds_{AdS_3}^2 = \frac{d\mu^2}{\mu^2} + \mu^2(-dt^2 + dy^2) & \\
\mu \rightarrow \lambda\mu, (t, y) \rightarrow \lambda^{-1}(t, y) & \\
(u, v) \rightarrow \sqrt{\lambda}(u, v), z \rightarrow \lambda^{-1}z & \\
e^{A_0} \rightarrow \lambda e^{A_0}, w \rightarrow \lambda^{-1}w & \\
u = \sqrt{\mu}m_1(\rho, \xi), \quad v = \sqrt{\mu}m_2(\rho, \xi), z = \mu^{-1}m_3(\rho, \xi), & \\
w = \mu^{-1}m_4(\rho, \xi), \quad e^{A_0} = \mu m_5(\rho, \xi), & \\
e^{2A}\hat{f}_1^2\mu^2 = e^{2A_0}, e^{2A}\hat{f}_2^2 = e^{-A_0}(-\partial_z w)^{-\frac{1}{2}}u^2, e^{2A}\hat{f}_3^2 = e^{-A_0}(-\partial_z w)^{\frac{1}{2}}v^2 &
\end{aligned}$$

$$\begin{aligned}
e^{2A} \left(\hat{f}_1^2 \frac{d\mu^2}{\mu^2} + \frac{d\xi^2 + d\rho^2}{4\rho^2} \right) &= e^{-A_0} \left((-\partial_z w)^{-\frac{1}{2}} du^2 + (-\partial_z w)^{\frac{1}{2}} dv^2 \right) \\
&\quad + e^{2A_0} (-\partial_z w) (dz + (-\partial_z w)^{-1} (\partial_u w) du)^2 \\
\frac{\gamma}{(1+\gamma)^2} \frac{1}{(G\bar{G}-1)} \frac{d\mu^2}{\mu^2} + \frac{d\xi^2 + d\rho^2}{4\rho^2} &= \frac{1}{W_+} \frac{du^2}{u^2} + \frac{1}{W_-} \frac{dv^2}{v^2} + \frac{1}{\beta^2 \rho^2 (G\bar{G}-1)} \frac{W_+}{W_-} \left(u^2 dz + (\partial_z w)^{-1} (u^3 \partial_u w) \frac{du}{u} \right)^2 \\
u^2 v^2 &= \frac{\beta^2 \gamma}{(\gamma+1)^2} \mu^2 \rho^2, (-\partial_z w) \frac{v^2}{u^2} = \frac{W_+}{W_-}, e^{A_0} = \frac{\beta \sqrt{\gamma} \mu \rho}{(\gamma+1)} e^{-2A} (W_+ W_-)^{\frac{1}{2}} \\
u &= \sqrt{a \mu \rho} e^{\alpha(\rho, \xi)}, v = \sqrt{a \mu \rho} e^{-\alpha(\rho, \xi)}, z = \mu^{-1} e^{-2\alpha(\rho, \xi)} p(\rho, \xi) \\
a &\equiv \frac{\beta \sqrt{\gamma}}{(\gamma+1)} \\
(\partial_z w)^{-1} (u^3 \partial_u w) &= b_2, (v^3 \partial_v w) = b_3 \\
\frac{\gamma}{(1+\gamma)^2} \frac{1}{(G\bar{G}-1)} \frac{d\mu^2}{\mu^2} + \frac{d\xi^2 + d\rho^2}{4\rho^2} &= \frac{1}{W_+} \frac{du^2}{u^2} + \frac{1}{W_-} \frac{dv^2}{v^2} + \frac{1}{\beta^2 \rho^2 (G\bar{G}-1)} \frac{W_+}{W_-} \left(u^2 dz + b_2 \frac{du}{u} \right)^2 \\
\partial_\xi \alpha &= -\frac{\varepsilon_1}{2\rho} g_1, \partial_\rho \alpha = \frac{1}{2\rho} g_2, b_2 = 2a\rho p + \frac{\varepsilon_2 \beta \rho g_1}{g_1^2 + g_2^2 + g_2} \\
\partial_\xi p &= -\frac{\varepsilon_1 \varepsilon_2 \beta}{2a\rho} (g_2 - 1), \partial_\rho p = -\frac{1}{\rho} \left(p + \frac{\varepsilon_2 \beta}{2a} g_1 \right) \\
\alpha &= -\frac{1}{2\beta} \tilde{\Phi} \\
p &= -\frac{\varepsilon_2}{2a\rho} (\Phi + \beta \xi) \\
b_2 &= \varepsilon_2 \left(\frac{\beta \rho g_1}{g_1^2 + g_2^2 + g_2} - (\Phi + \beta \xi) \right) = \varepsilon_2 \left(\frac{h(G + \bar{G})}{W_+} - (\Phi - \tilde{h}) \right) \\
\gamma = 1, u &= \left(\frac{1}{2} \beta \mu \rho \right)^{\frac{1}{2}} e^{-\frac{1}{2\beta} \tilde{\Phi}}, v = \left(\frac{1}{2} \beta \mu \rho \right)^{\frac{1}{2}} e^{+\frac{1}{2\beta} \tilde{\Phi}}, z = -\frac{\varepsilon_2}{\beta \rho \mu} e^{\frac{1}{\beta} \tilde{\Phi}} (\Phi + \beta \xi) \\
\partial_z w &= -\frac{g_1^2 + g_2^2 + g_2}{g_1^2 + g_2^2 - g_2} e^{-\frac{2}{\beta} \tilde{\Phi}}, (\partial_z w)^{-1} (u^3 \partial_u w) = b_2, (v^3 \partial_v w) = b_3 \\
dw &= (\partial_z w) dz + (\partial_u w) du + (\partial_v w) dv = d \left[\frac{\varepsilon_2}{\beta \rho \mu} e^{-\frac{1}{\beta} \tilde{\Phi}} (\Phi - \beta \xi) \right] \\
w &= \frac{\varepsilon_2}{\beta \rho \mu} e^{-\frac{1}{\beta} \tilde{\Phi}} (\Phi - \beta \xi) \\
b_3 &= \varepsilon_2 \left(\frac{\beta \rho g_1}{g_1^2 + g_2^2 - g_2} - (\Phi - \beta \xi) \right) = \varepsilon_2 \left(\frac{h(G + \bar{G})}{W_-} - (\Phi + \tilde{h}) \right) \\
u^2 z &= -\frac{1}{2} \varepsilon_2 (\Phi + \beta \xi), v^2 w = \frac{1}{2} \varepsilon_2 (\Phi - \beta \xi) \\
\Phi &\rightarrow -\Phi, \tilde{\Phi} \rightarrow -\tilde{\Phi} \Rightarrow u \leftrightarrow v, z \leftrightarrow w \\
\omega &\equiv e^{3A_0} (-\partial_z w)^{\frac{1}{2}} (dz + (\partial_z w)^{-1} (\partial_u w) du) \\
\omega &= \frac{W_+ \mu^2}{4(G\bar{G}-1)} \left(u^2 dz + (\partial_z w)^{-1} (u^3 \partial_u w) \frac{du}{u} \right) \\
&= \frac{\varepsilon_2}{4} \left[\left(\frac{\beta \rho g_1}{g_1^2 + g_2^2 - 1} + 2\Phi \right) \mu d\mu - d(\mu^2 \Phi) \right] \\
&= \frac{\varepsilon_2}{8} \left[\left(\frac{h(G + \bar{G})}{(G\bar{G}-1)} + 4\Phi \right) \mu d\mu - d(2\mu^2 \Phi) \right] = \frac{\varepsilon_2}{v_1} b_1 \mu d\mu - \frac{\varepsilon_2}{4} d(\mu^2 \Phi)
\end{aligned}$$



$$\begin{aligned}
b_1 &= \frac{v_1}{4} \left(\frac{\beta \rho g_1}{g_1^2 + g_2^2 - 1} + 2\Phi \right) \\
C_{tyz}^{(3)} &= -e^0 \wedge e^1 \wedge e^2 = -dt \wedge dy \wedge \omega = -\frac{\varepsilon_2}{v_1} b_1 \mu dt \wedge dy \wedge d\mu + \frac{\varepsilon_2}{4} d(\mu^2 \Phi dt \wedge dy) \\
C_{tyz}^{(3)} &= b_1 \mu dt \wedge dy \wedge d\mu \\
e^0 &= \frac{\mu e^A}{2\sqrt{G\bar{G}-1}} dt, e^1 = \frac{\mu e^A}{2\sqrt{G\bar{G}-1}} dy \\
e^2 &= \frac{\varepsilon_2 e^A}{\rho \sqrt{(G\bar{G}-1)W_+W_-}} \left(\rho g_1 \frac{d\mu}{\mu} + (G\bar{G}-1)(g_2 d\xi - g_1 d\rho) \right), \\
e^3 &= \frac{e^A}{2\sqrt{W_+}} \left(\frac{d\mu}{\mu} + \frac{d\rho}{\rho} + \frac{1}{\rho} (g_1 d\xi + g_2 d\rho) \right), \\
e^4 &= \frac{e^A}{2\sqrt{W_-}} \left(\frac{d\mu}{\mu} + \frac{d\rho}{\rho} - \frac{1}{\rho} (g_1 d\xi + g_2 d\rho) \right), \\
e^{i+4} &= \frac{e^A}{2\sqrt{W_+}} \sigma_i, e^{i+7} = \frac{e^A}{2\sqrt{W_-}} \tilde{\sigma}_i, i = 1, 2, 3 \\
\partial_\xi b_1 &= \varepsilon_2 \partial_\xi \left[-\frac{\beta \rho g_1}{4(G\bar{G}-1)} \right] + \frac{1}{2} \varepsilon_2 \beta g_2, \quad \partial_\rho b_1 = \varepsilon_2 \partial_\rho \left[-\frac{\beta \rho g_1}{4(G\bar{G}-1)} \right] - \frac{1}{2} \varepsilon_2 \beta g_1, \\
\partial_\xi b_2 &= \varepsilon_2 \partial_\xi \left[-\frac{2\beta \rho g_1}{W_+} + \beta \xi \right] - \varepsilon_2 \beta g_2, \quad \partial_\rho b_2 = \varepsilon_2 \partial_\rho \left[-\frac{2\beta \rho g_1}{W_+} + \varepsilon_2 \beta \xi \right] + \varepsilon_2 \beta g_1, \\
\partial_\xi b_3 &= \varepsilon_2 \partial_\xi \left[-\frac{2\beta \rho g_1}{W_-} - \beta \xi \right] - \varepsilon_2 \beta g_2, \quad \partial_\rho b_3 = \varepsilon_2 \partial_\rho \left[-\frac{2\beta \rho g_1}{W_-} - \beta \xi \right] + \varepsilon_2 \beta g_1. \\
b_1 &= -\varepsilon_2 \left(\frac{\beta \rho g_1}{4(G\bar{G}-1)} + \frac{1}{2} \Phi \right), b_2 = -\varepsilon_2 \left(\frac{2\beta \rho g_1}{W_+} - (\Phi + \beta \xi) \right), b_3 = -\varepsilon_2 \left(\frac{2\beta \rho g_1}{W_-} - (\Phi - \beta \xi) \right) \\
\Gamma^{012} \varepsilon &= \eta_1 \varepsilon, \Gamma^{013567} \varepsilon = \eta_2 \varepsilon, \Gamma^{0148910} \varepsilon = -\eta_1 \eta_2 \varepsilon \\
b_1 &= \varepsilon_2 \eta_1 \left(\frac{\beta \rho g_1}{4(G\bar{G}-1)} + \frac{1}{2} \Phi \right), b_2 = \varepsilon_2 \eta_1 \eta_2 \left(\frac{2\beta \rho g_1}{W_+} - (\Phi + \beta \xi) \right) \\
b_3 &= -\varepsilon_2 \eta_2 \left(\frac{2\beta \rho g_1}{W_-} - (\Phi - \beta \xi) \right) \\
\eta_1 &= -1, \eta_2 = +1, \varepsilon_2 = +1 \\
ds_{IIB}^2 &= \sqrt{h_{11}} \left[-e^{3A} dt^2 + e^{3A} h_{ab} dr^a dr^b + \frac{e^{-3A}}{\text{deth}} dw_2^2 + d\mathbf{y}_6^2 \right] \\
e^{2\phi} &= \frac{h_{11}^2}{\text{deth}}, C_0 = -\frac{h_{12}}{h_{11}}, B_2 = e^{3A} h_{1a} dt \wedge dr^a, C_2 = e^{3A} h_{2a} dt \wedge dr^a \\
h_{ab} &= \frac{1}{2} \partial_a \partial_b K(r^1, r^2, \mathbf{y}) \\
\Delta_y K + 2e^{-3A} &= 0 \\
r^1 \rightarrow z, r^2 \rightarrow u_1, w_2 \rightarrow u_2, y_1 \rightarrow u_3, y_2 \rightarrow y, y_{3,4,5,6} \rightarrow v_{3,4,5,6} \\
ds^2 &= \frac{1}{\sqrt{\text{deth}}} (-dt^2 + dy^2) + \frac{\sqrt{\text{deth}}}{h_{11}} (du_2^2 + du_3^2) + \sqrt{\text{deth}} (e^{3A} h_{ab} dr^a dr^b + ds_{\mathbb{R}^4}^2) \\
e^{2\phi} &= \frac{\sqrt{\text{deth}}}{h_{11}}, B_2 = \frac{h_{12}}{h_{11}} du_2 \wedge du_3 \\
C_3 &= e^{3A} h_{1a} dt \wedge dr^a \wedge dy - \frac{v^3}{2} \partial_v \partial_z K d\Omega'_3 \\
C_5 &= \frac{1}{h_{11}} dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy + \frac{v^3}{2} \left(\frac{h_{12}}{h_{11}} \partial_v \partial_z K - \partial_v \partial_{u_1} K \right) du_2 \wedge du_3 \wedge d\Omega'_3
\end{aligned}$$



$$\begin{aligned}
ds_{11}^2 &= e^{\frac{-2\phi}{3}} ds_{10}^2 + e^{\frac{4\phi}{3}} (dx + C_1)^2 \\
C_3' &= C_3 + B_2 \wedge dx \\
ds_{11}^2 &= \frac{h_{11}^{1/3}}{(\text{deth})^{2/3}} (-dt^2 + dy^2) + \frac{(\text{deth})^{1/3}}{h_{11}^{2/3}} (du_2^2 + du_3^2 + du_4^2) \\
&\quad + (\text{deth})^{1/3} h_{11}^{1/3} (e^{3A} h_{ab} dr^a dr^b + ds_{\mathbb{R}^4}^2) \\
C_3 &= \frac{h_{11}}{\text{deth}} dt \wedge dz \wedge dy + \frac{h_{12}}{\text{deth}} dt \wedge du_1 \wedge dy - \frac{h_{12}}{h_{11}} du_2 \wedge du_3 \wedge du_4 - \frac{v^3}{2} \partial_v \partial_z K d\Omega_3' \\
ds_{11}^2 &= \frac{h_{11}^{1/3}}{(\text{deth})^{2/3}} (-dt^2 + dy^2) + \frac{(\text{deth})^{1/3}}{h_{11}^{2/3}} (du_1^2 + du_2^2 + du_3^2 + du_4^2) \\
&\quad + \frac{h_{11}^{4/3}}{(\text{deth})^{2/3}} \left(dz + \frac{h_{12}}{h_{11}} du_1 \right)^2 + (\text{deth})^{1/3} h_{11}^{1/3} (dv^2 + v^2 d\Omega_3'^2) \\
ds_{11}^2 &= e^{2A_0} (-dt^2 + dy^2) + e^{-A_0} (-\partial_z w)^{-\frac{1}{2}} (du_1^2 + du_2^2 + du_3^2 + du_4^2) \\
&\quad + e^{2A_0} (-\partial_z w) \left(dz + (\partial_z w)^{-1} (\partial_{u_1} w) du_1 \right)^2 + e^{-A_0} (-\partial_z w)^{\frac{1}{2}} (dv^2 + v^2 d\Omega_3'^2) \\
C^{(3)} &= -e^{3A_0} (-\partial_z w)^{\frac{1}{2}} dt \wedge dy \wedge dz + e^{3A_0} (-\partial_z w)^{-\frac{1}{2}} (\partial_{x_1} w) dt \wedge dy \wedge dx_1 \\
&\quad + (-\partial_z w)^{-1} (\partial_{u_1} w) du_2 \wedge du_3 \wedge du_4 + (v^3 \partial_v w) d\Omega_3' \\
e^{2A_0} &= \frac{h_{11}^{1/3}}{(\text{deth})^{2/3}}, h_{11} = -\partial_z w, h_{12} = -\partial_{u_1} w \\
t = \eta_0, y = \eta_1, z = \eta_2, \vec{u}, \vec{v} &\text{ constante} \\
C_{tyz}^{(3)} &= -e^0 \wedge e^1 \wedge e^2 \\
\rho &= k_1 \mu^{-1}, \tilde{\Phi}(\xi, \rho) = k_2 \\
t = \eta_0, y = \eta_1, \mu = e^{\eta_2}, \xi = \sigma_1(\eta_2), \rho = \sigma_2(\eta_2) \\
d\hat{s}_3^2 &= e^{2A} \left[\hat{f}_1^2 \left(d\eta_2^2 + e^{2\eta_2} (-d\eta_0^2 + d\eta_1^2) \right) + \frac{(\sigma_1')^2 + (\sigma_2')^2}{4\sigma_2^2} d\eta_2^2 \right] \\
\mathcal{L}_{DBI} &= e^{3A} \hat{f}_1^2 e^{2\eta_2} \left(\hat{f}_1^2 + \frac{(\sigma_1')^2 + (\sigma_2')^2}{4\sigma_2^2} \right)^{\frac{1}{2}} \\
&= h \hat{f}_1^2 e^{2\eta_2} \left[(G\bar{G} - 1) W_+ W_- \left(\hat{f}_1^2 + \frac{(\sigma_1')^2 + (\sigma_2')^2}{4\sigma_2^2} \right) \right]^{\frac{1}{2}} \\
&\quad \frac{(\sigma_1')^2 + (\sigma_2')^2}{\sigma_2^2} = \frac{g_1^2 + g_2^2}{g_1^2} \\
\left[(G\bar{G} - 1) W_+ W_- \left(\hat{f}_1^2 + \frac{(\sigma_1')^2 + (\sigma_2')^2}{4\sigma_2^2} \right) \right] &= \left(\frac{W_+ W_-}{4g_1} \right)^2 \\
\mathcal{L}_{DBI} &= e^{2\eta_2} \frac{\beta \sigma_2 ((g_1^2 + g_2^2)^2 - g_2^2)}{4g_1 (g_1^2 + g_2^2 - 1)} \\
\hat{C}^{(3)} &= b_1 e^{2\eta_2} d\eta_0 \wedge d\eta_1 \wedge d\eta_2 \\
\tilde{C}^{(3)} &= e^{2\eta_2} (b_1 + 2\Lambda + (\partial_\xi \Lambda) \sigma_1' + (\partial_\rho \Lambda) \sigma_2') d\eta_0 \wedge d\eta_1 \wedge d\eta_2 \\
\check{C}^{(3)} &= e^{2\eta_2} \frac{v_1}{c_1^3} \left[\frac{h(G + \bar{G})}{(G\bar{G} - 1)} - 2(\partial_\xi \Phi) \sigma_1' - 2(\partial_\rho \Phi) \sigma_2' \right] d\eta_0 \wedge d\eta_1 \wedge d\eta_2 \\
&= v_1 e^{2\eta_2} \frac{\beta \sigma_2}{4} \left[\frac{g_1}{(g_1^2 + g_2^2 - 1)} + g_2 \frac{\sigma_1'}{\sigma_2} - g_1 \frac{\sigma_2'}{\sigma_2} \right] d\eta_0 \wedge d\eta_1 \wedge d\eta_2 \\
\sigma_2 &= k_1 e^{-\eta_2}, \frac{\sigma_1'}{\sigma_2} = \frac{g_2}{g_1} \\
g_1 \sigma_1' + g_2 \sigma_2' &= 0 \Leftrightarrow \partial_\xi \tilde{\Phi} \sigma_1' + \partial_\rho \tilde{\Phi} \sigma_2' = 0 \\
\eta_0 = t, \eta_1 = -u_1, \eta_2 = u_2, \eta_3 = u_3, \eta_4 = y
\end{aligned}$$

$$\begin{aligned}
d\tilde{s}_5^2 &= \frac{1}{\sqrt{\text{deth}}}(-dt^2 + dy^2) + \frac{h_{22}}{\sqrt{\text{deth}}}du_1^2 + \frac{\sqrt{\text{deth}}}{h_{11}}(du_2^2 + du_3^2) \\
\tilde{B}_2 &= \frac{h_{12}}{h_{11}}du_2 \wedge du_3 \\
\tilde{C}_3 &= -\frac{h_{12}}{\text{deth}}dt \wedge du_1 \wedge dy \\
\tilde{C}_5 &= -\frac{1}{h_{11}}dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy \\
S_{DBI} &= -T_4 \int d^5\sigma e^{-\phi} \sqrt{-\det(\tilde{G}_{\alpha\beta} + F_{\alpha\beta} + \tilde{B}_{\alpha\beta})} = -T_4 \int d^5\sigma \frac{h_{22}}{\text{deth}} \\
S_{WZ} &= -T_4 \int e^{\tilde{B}_2 + \tilde{F}_2} \wedge \oplus_n \tilde{C}_n = T_4 \int d^5\sigma \frac{h_{22}}{\text{deth}} \\
h &= -iw + i\bar{w}, G = \pm \left[i + \sum_{a=1}^{n+1} \frac{\zeta_a \text{Im}(w)}{(\bar{w} - \xi_a)|w - \xi_a|} \right] \\
g_1 &= \pm \sum_{a=1}^{n+1} \frac{\zeta_a \rho (\xi - \xi_a)}{((\xi - \xi_a)^2 + \rho^2)^{\frac{3}{2}}}, g_2 = \pm \left[1 + \sum_{a=1}^{n+1} \frac{\zeta_a \rho^2}{((\xi - \xi_a)^2 + \rho^2)^{\frac{3}{2}}} \right] \\
\tilde{\Phi} &= \pm 2 \left[-\log \rho + \sum_{a=1}^{n+1} \frac{\zeta_a}{\sqrt{(\xi - \xi_a)^2 + \rho^2}} \right], \Phi = \mp 2 \left[\xi + \sum_{a=1}^{n+1} \frac{\zeta_a (\xi - \xi_a)}{\sqrt{(\xi - \xi_a)^2 + \rho^2}} \right] \\
ds_{11} &= -e^{2A_0} dt^2 + e^{2A_1} (dy - P dt)^2 + e^{2A_2} du^2 + e^{2A_3} dv^2 + u^2 e^{2A_4} d\Omega_3^2 + v^2 e^{2A_5} d\Omega_3'^2 \\
&\quad + e^{2A_6} (dz + B_1 du)^2 \\
e^0 &= e^{A_0} dt, e^1 = e^{A_1} (dy - P dt), e^2 = e^{A_6} (dz + B_1 du), \\
e^3 &= e^{A_2} du, e^4 = e^{A_3} dv, e^{i+4} = u e^{A_4} \sigma_i, e^{i+7} = v e^{A_5} \tilde{\sigma}_i, i = 1, 2, 3 \\
\Gamma^{01} \varepsilon &= -\varepsilon, \Gamma^{012} \varepsilon = -\varepsilon, \Gamma^{013456} \varepsilon = \varepsilon \\
\Gamma^{0178910} \varepsilon &= -\varepsilon \\
\delta\psi_\mu &\equiv \nabla_\mu \epsilon + \frac{1}{288} \left(\Gamma_\mu^{\nu\rho\lambda\sigma} - 8\delta_\mu^\nu \Gamma^{\rho\lambda\sigma} \right) F_{\nu\rho\lambda\sigma} \epsilon \\
P &\equiv 1 - e^{A_0 - A_1} \\
\hat{A}_0 &\equiv \frac{1}{2}(A_0 + A_1), \hat{A}_1 \equiv \frac{1}{2}(A_0 - A_1) \\
ds_{11} &= e^{2\hat{A}_0} [-e^{2\hat{A}_1} dt^2 + e^{-2\hat{A}_1} (dy + (e^{2\hat{A}_1} - 1) dt)^2 + (-\partial_z w) (dz + (\partial_z w)^{-1} (\partial_u w) du)^2 \\
&\quad + e^{-3\hat{A}_0} (-\partial_z w)^{-\frac{1}{2}} (du^2 + u^2 d\Omega_3^2) + e^{-3\hat{A}_0} (-\partial_z w)^{\frac{1}{2}} (dv^2 + v^2 d\Omega'^2)] \\
e^0 &= e^{\hat{A}_0 + \hat{A}_1} dt, e^1 = e^{\hat{A}_0 - \hat{A}_1} (dy + (e^{2\hat{A}_1} - 1) dt) \\
e^2 &= e^{\hat{A}_0} (-\partial_z w)^{\frac{1}{2}} (dz + (\partial_z w)^{-1} (\partial_u w) du) \\
e^3 &= e^{-\frac{1}{2}\hat{A}_0} (-\partial_z w)^{-\frac{1}{4}} du, e^4 = e^{-\frac{1}{2}\hat{A}_0} (-\partial_z w)^{\frac{1}{4}} dv, \\
e^{i+4} &= \frac{1}{2} u e^{-\frac{1}{2}\hat{A}_0} (-\partial_z w)^{-\frac{1}{4}} \sigma_i, e^{i+7} = \frac{1}{2} v e^{-\frac{1}{2}\hat{A}_0} (-\partial_z w)^{\frac{1}{4}} \tilde{\sigma}_i, i = 1, 2, 3 \\
C^{(3)} &= -e^0 \wedge e^1 \wedge e^2 + (\partial_z w)^{-1} (u^3 \partial_u w) \text{Vol}(S^3) + (v^3 \partial_v w) \text{Vol}(S'^3) \\
F_1 &\equiv (-\partial_z w)^{\frac{1}{2}} e^{-3\hat{A}_0}, F_2 \equiv (-\partial_z w)^{-\frac{1}{2}} e^{-3\hat{A}_0} + (-\partial_z w)^{-1} (\partial_u w)^2 \\
\mathcal{L}(H) &= e^{2\hat{A}_0} (-\partial_z w)^{-\frac{1}{2}} \left[(-\partial_z w)^{\frac{1}{3}} \partial_u (u^3 \partial_u H) + \frac{1}{v^3} \partial_v (v^3 \partial_v H) + 2(\partial_u w) \partial_u \partial_z H \right. \\
&\quad \left. + \left((-\partial_z w)^{-\frac{1}{2}} e^{-3\hat{A}_0} + (-\partial_z w)^{-1} (\partial_u w)^2 \right) \partial_z^2 H \right] \\
\mathcal{L}(e^{-2\hat{A}_1}) &= 0
\end{aligned}$$

$$\begin{aligned}
ds_{11}^2 &= e^{2A} \left[\hat{f}_1^2 \left(\frac{d\mu^2}{\mu^2} + \mu^2 \left(-e^{2\hat{A}_1} dt^2 + e^{-2\hat{A}_1} (dy + (e^{2\hat{A}_1} - 1) dt)^2 \right) \right) \right. \\
&\quad \left. + \hat{f}_2^2 ds_{S^3}^2 + \hat{f}_3^2 ds_{S'^3}^2 + \frac{d\xi^2 + d\rho^2}{4\rho^2} \right] \\
\mathcal{L}(H) &= 4e^{-A} \left[(G\bar{G} - 1) \frac{1}{\mu} \partial_\mu (\mu^3 \partial_\mu H) + \frac{1}{\rho} \partial_\rho (\rho^3 \partial_\rho H) + \rho^2 \partial_\xi^2 H \right] \\
&\quad \frac{1}{\rho} \partial_\rho (\rho^3 \partial_\rho K) + \rho^2 \partial_\xi^2 K + p(p+2)(G\bar{G} - 1)K \\
K &= \frac{c_1}{u^2} + \frac{c_2}{v^2} = \frac{2}{\beta\rho} \left(c_1 e^{\frac{1}{\beta}\Phi} + c_2 e^{-\frac{1}{\beta}\Phi} \right) \\
&\quad (u^2 + v^2)^{-3} \sim \mu^{-3} \rho^{-3} \\
&\quad \frac{1}{\rho^3} \partial_\rho (\rho^3 \partial_\rho K) + \partial_\xi^2 K + \frac{p(p+2)}{\rho^2} (G\bar{G} - 1)K \\
\mathcal{L}_4(K) &\equiv \frac{1}{\rho^3} \partial_\rho (\rho^3 \partial_\rho K) + \partial_\xi^2 K \\
ds_5^2 &\equiv d\rho^2 + \rho^2 d\Omega_3^2 + d\xi^2 \\
&\quad \mathcal{L}_4 \left(\frac{1}{(\rho^2 + \xi^2)^{\frac{3}{2}}} \right) \\
&\quad \frac{p(p+2)}{\rho^2} (G\bar{G} - 1) \sim \frac{c_0}{(\rho^2 + \xi^2)^{\frac{3}{2}}} \\
K &= \frac{Q}{(\rho^2 + \xi^2)^{\frac{3}{2}}} \left(1 - \frac{c_0}{4} \frac{1}{\sqrt{\rho^2 + \xi^2}} + \dots \right) \\
e^{-2\hat{A}_1} &= V_0 + V_1 \mu^{-2} \\
V_0 &= 1 + \sum_{a=1}^m \frac{k_a}{\left((\xi - \xi_a)^2 + \rho^2 \right)^{\frac{3}{2}}}, V_1 = q_0 + \sum_{a=1}^{m'} \frac{q_a}{\left((\xi - \xi_a)^2 + \rho^2 \right)^{\frac{3}{2}}} \\
e^{-2\hat{A}_1} &= 1 + \alpha + Q\mu^{-2} \\
&\quad \frac{d\mu^2}{\mu^2} + \mu^2 \left(-e^{2\hat{A}_1} dt^2 + e^{-2\hat{A}_1} (dy + (e^{2\hat{A}_1} - 1) dt)^2 \right) \\
&\quad = \frac{d\mu^2}{\mu^2} - \frac{\mu^4}{Q + \mu^2} dt^2 + (Q + \mu^2) \left(dy - \frac{Q}{Q + \mu^2} dt \right)^2 \\
&\quad = \frac{d\mu^2}{\mu^2} + \mu^2 (-dt^2 + dy^2) + Q(dy - dt)^2 \\
ds_{11}^2 &= e^{2\alpha_0} (-dt^2 + dy^2) + e^{2\alpha_1} d\Omega_3^2 + e^{2\alpha_2} d\Omega_3'^2 + g_{ij} dz^i dz^j \\
ds_3^2 &= g_{ij} dz^i dz^j = e^{2\alpha_3} dz^2 + e^{2\alpha_4} du^2 + e^{2\alpha_5} dv^2 \\
e^0 &= e^{\alpha_0} dt, \quad e^1 = e^{\alpha_0} dy, \quad e^2 = e^{\alpha_1} dz, \quad e^3 = e^{\alpha_2} du, \quad e^4 = e^{\alpha_3} dv, \\
e^{i+4} &= e^{\alpha_4} \sigma_i, \quad e^{i+7} = e^{\alpha_5} \tilde{\sigma}_i, \quad i = 1, 2, 3 \\
&\quad \Gamma^{012} \varepsilon = -\varepsilon, \Gamma^{013567} \varepsilon = \varepsilon, \Gamma^{0148910} \varepsilon = -\varepsilon \\
e^0 &= e^{A_0} dt, \quad e^1 = e^{A_0} dy, \quad e^2 = e^{A_1} (dz + B_1 du + B_2 dv) \\
e^3 &= e^{A_2} du, \quad e^4 = e^{A_3} dv, \quad e^{i+4} = u e^{A_4} \sigma_i, \quad e^{i+7} = v e^{A_5} \tilde{\sigma}_i, \quad i = 1, 2, 3 \\
&\quad B_2 \equiv 0 \\
e^0 &= e^{A_0} dt, \quad e^1 = e^{A_0} dy, \quad e^2 = e^{A_1} (dz + B_1 du), \\
e^3 &= e^{A_2} du, \quad e^4 = e^{A_3} dv, \quad e^{i+4} = u e^{A_4} \sigma_i, \quad e^{i+7} = v e^{A_5} \tilde{\sigma}_i, \quad i = 1, 2, 3 \\
ds_{11}^2 &= e^{2A_0} (-dt^2 + dy^2) + e^{2A_2} du^2 + e^{2A_3} dv^2 + u^2 e^{2A_4} d\Omega_3^2 + v^2 e^{2A_5} d\Omega_3'^2 \\
&\quad + e^{2A_1} (dz + B_1 du)^2
\end{aligned}$$

$$\begin{aligned}
F^{(4)} &= e^0 \wedge e^1 \wedge (b_1 e^2 \wedge e^3 + b_2 e^2 \wedge e^4 + b_3 e^3 \wedge e^4) \\
&\quad + (b_4 e^2 + b_5 e^3 + b_6 e^4) \wedge e^5 \wedge e^6 \wedge e^7 + (b_7 e^2 + b_8 e^3 + b_9 e^4) \wedge e^8 \wedge e^9 \wedge e^{10} \\
&\quad \epsilon = e^{\frac{1}{2}A_0} \epsilon_0 \\
&\quad \partial_u(A_5 - A_3) = \partial_z(A_5 - A_3) = 0, \partial_v(A_4 - A_2) = \partial_z(A_4 - A_2) \\
&\quad A_4 = A_2, A_5 = A_3 \\
ds_{11}^2 &= e^{2A_0}(-dt^2 + dy^2) + e^{2A_2}(du^2 + u^2 d\Omega_3^2) + e^{2A_3}(dv^2 + v^2 d\Omega'^2) \\
&\quad + e^{2A_1}(dz + B_1 du)^2 \\
&\quad A_3 = -(A_0 + A_2) \\
&\quad A_1 = -2A_2 \\
ds_{11}^2 &= e^{2A_0}[(-dt^2 + dy^2) + e^{2(A_2 - A_0)}(du^2 + u^2 d\Omega_3^2) + e^{-2(A_2 + 2A_0)}(dv^2 + v^2 d\Omega_{33}'^2) \\
&\quad + e^{-2(A_0 + 2A_2)}(dz + B_1 du)^2] \\
&\quad \partial_z(B_1 e^{-2(A_0 + 2A_2)}) = \partial_u(e^{-2(A_0 + 2A_2)}) \\
B_1 e^{-2(A_0 + 2A_2)} &= -\partial_u w, e^{-2(A_0 + 2A_2)} = -\partial_z w \\
B_1 &= (\partial_z w)^{-1} \partial_u w, e^{-2(A_0 + 2A_2)} = -\partial_z w \\
F_1 &\equiv (-\partial_z w)^{\frac{1}{2}} e^{-3A_0}, F_2 \equiv (-\partial_z w)^{-\frac{1}{2}} e^{-3A_0} + (-\partial_z w)^{-1} (\partial_u w)^2 \\
H_1 &\equiv \mathcal{L}_v w - \partial_z F_1, H_2 \equiv \mathcal{L}_u w + \partial_z F_2 \\
\partial_z H_1 &= \partial_u H_1 = \partial_z H_2 = \partial_v H_2 = 0 \\
w &= \partial_z G_0, F_1 = \mathcal{L}_v G_0, F_2 = -\mathcal{L}_u G_0 \\
\mathcal{L}_v G_0 &= (\partial_z^2 G_0)(\mathcal{L}_u G_0) - (\partial_u \partial_z G_0)^2 \\
&\quad \left(\frac{\partial F}{\partial \eta} \right)_{\zeta, \xi} \\
dw &= \left(\frac{\partial w}{\partial z} \right)_{\vec{u}, \vec{v}} dz + \left(\frac{\partial w}{\partial u_i} \right)_{z, \vec{v}} du_i + \left(\frac{\partial w}{\partial v_i} \right)_{z, \vec{u}} dv_i \\
\left(\frac{\partial z}{\partial u_i} \right)_{w, \vec{v}} &= - \left(\left(\frac{\partial w}{\partial z} \right)_{\vec{u}, \vec{v}} \right)^{-1} \left(\frac{\partial w}{\partial u_i} \right)_{z, \vec{v}}, \left(\frac{\partial z}{\partial v_i} \right)_{w, \vec{u}} = - \left(\left(\frac{\partial w}{\partial z} \right)_{\vec{u}, \vec{v}} \right)^{-1} \left(\frac{\partial w}{\partial v_i} \right)_{z, \vec{u}} \\
&\quad \left(\frac{\partial z}{\partial w} \right)_{\vec{u}, \vec{v}} = \left(\left(\frac{\partial w}{\partial z} \right)_{\vec{u}, \vec{v}} \right)^{-1} \\
e^2 &= (-\partial_z w)^{\frac{1}{2}} (dz + (\partial_z w)^{-1} (\vec{\nabla}_{\vec{u}} w) \cdot d\vec{u}) \\
&= -(-\partial_z w)^{-\frac{1}{2}} ((\partial_z w) dz + (\vec{\nabla}_{\vec{u}} w) \cdot d\vec{u}) = -(-\partial_z w)^{-\frac{1}{2}} (dw - (\vec{\nabla}_{\vec{v}} w) \cdot d\vec{v}) \\
&= -((- \partial_w z)_{\vec{u}, \vec{v}})^{\frac{1}{2}} \left(dw + \left(\left(\frac{\partial z}{\partial w} \right)_{\vec{v}, \vec{u}} \right)^{-1} \left(\frac{\partial z}{\partial v_i} \right)_{w, \vec{v}} dv_i \right) \\
&= -(-\partial_w z)^{\frac{1}{2}} (dw + (\partial_w z)^{-1} (\vec{\nabla}_{\vec{v}} z) \cdot d\vec{v}) \\
C^{(3)} &= -e^0 \wedge e^1 \wedge e^2 + \frac{1}{3!} \epsilon_{ijk\ell} (-(\partial_{u_\ell} z) du^i \wedge du^j \wedge du^k + (\partial_w z)^{-1} (\partial_{v_\ell} z) dv^i \wedge dv^j \wedge dv^k) \\
ds^2 &= G_{xx} (dx + A_\mu dx^\mu)^2 + \hat{g}_{\mu\nu} dx^\mu dx^\nu, \\
B_2 &= B_{\mu x} dx^\mu \wedge (dx + A_\mu dx^\mu) + \hat{B}_2, \\
C_p &= C_{(p-1)x} \wedge (dx + A_\mu dx^\mu) + \hat{C}_p, \\
d\tilde{s}^2 &= G_{xx}^{-1} (dx + B_{\mu x} dx^\mu)^2 + \hat{g}_{\mu\nu} dx^\mu dx^\nu, \\
e^{2\tilde{\phi}} &= G_{xx}^{-1} e^{2\phi} \\
\tilde{B}_2 &= A_\mu dx^\mu \wedge dx + \hat{B}_2 \\
\tilde{C}_p &= \hat{C}_{p-1} \wedge (dx + B_{\mu x} dx^\mu) + C_{(p)x}.
\end{aligned}$$

$$\begin{aligned}
\tilde{g}_{\mu\nu} &= \sqrt{C_0^2 + e^{-2\phi}} g_{\mu\nu}, e^{-\tilde{\phi}} = \frac{e^{-\phi}}{C_0^2 + e^{-2\phi}}, \tilde{C}_0 = -\frac{C_0}{C_0^2 + e^{-2\phi}}, \\
\tilde{B}_2 &= -C_2, \tilde{C}_2 = B_2, \tilde{C}_4 = C_4 + B_2 \wedge C_2 \\
ds^2 &= \frac{\sqrt{\text{deth}}}{h_{11}} (du_2^2 + du_3^2) + \frac{1}{\sqrt{\text{deth}}} (-dt^2 + dy^2) + \sqrt{\text{deth}} (e^{3A} h_{ab} dr^a dr^b + ds_{\mathbb{R}^4}^2) \\
e^{2\phi} &= \frac{\sqrt{\text{deth}}}{h_{11}}, B_2 = \frac{h_{12}}{h_{11}} du_2 \wedge du_3 \\
C_3 &= e^{3A} h_{1a} dt \wedge dr^a \wedge dy, C_5 = \frac{1}{h_{11}} dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy \\
F_p &= dC_{p-1} \quad \text{para } p < 3, \\
F_p &= dC_{p-1} + H_3 \wedge C_{p-3} \quad \text{para } p \geq 3, \\
F_6 &= \star F_4, F_8 = \star F_2. \\
v_3 &= v \cos \phi_1 \\
v_4 &= v \sin \phi_1 \cos \phi_2 \\
v_5 &= v \sin \phi_1 \sin \phi_2 \cos \phi_3 \\
v_6 &= v \sin \phi_1 \sin \phi_2 \sin \phi_3 \\
ds_{\mathbb{R}^4}^2 &= dv^2 + v^2 (d\phi_1^2 + \sin^2 \phi_1 (d\phi_2^2 + \sin^2 \phi_2 d\phi_3^2)) \\
dC_5^e &= -\frac{1}{h_{11}^2} (\partial_z h_{11} dz + \partial_v h_{11} dv) \wedge dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy \\
H_3 \wedge C_3^e &= \left[\partial_z \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{12} - \partial_{u_1} \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{11} \right] dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dz \wedge dy \\
&\quad - \partial_v \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{11} dz \wedge dt \wedge dv \wedge du_2 \wedge du_3 \wedge dy \\
&\quad - \partial_v \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{12} du_1 \wedge dt \wedge dv \wedge du_2 \wedge du_3 \wedge dy \\
F_6^e &= f_1 dt \wedge dz \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy \\
&\quad + f_2 dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy \wedge dv \\
&\quad - f_3 dt \wedge dz \wedge du_2 \wedge du_3 \wedge dy \wedge dv \\
f_1 &= \frac{1}{h_{11}^2} \partial_z h_{11} - \partial_z \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{12} + \partial_{u_1} \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{11}, \\
f_2 &= \frac{1}{h_{11}^2} \partial_v h_{11} - \partial_v \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{12}, \\
f_3 &= \partial_v \left(\frac{h_{12}}{h_{11}} \right) e^{3A} h_{11}. \\
F_4^m &= -v^3 h_{11} (f_2 h_{11} + f_3 h_{12}) dz \wedge d\Omega'_3 - v^3 h_{11} (f_2 h_{12} + f_3 h_{22}) du_1 \wedge d\Omega'_3 \\
&\quad + r^3 f_1 h_{11} \text{deth} dr \wedge d\Omega'_3 \\
F_4^m &= -(v^3 \partial_v h_{11} dz + v^3 \partial_v h_{12} du_1) \wedge d\Omega'_3 \\
&\quad + v^3 (h_{22} \partial_z h_{11} - h_{12} \partial_z h_{12} - h_{12} \partial_{u_1} h_{11} + h_{11} \partial_{u_1} h_{12}) dv \wedge d\Omega'_3 \\
&\quad - \frac{1}{2} v^3 \partial_v \partial_z^2 K dz - \frac{1}{2} v^3 \partial_v \partial_z \partial_{u_1} K du_1 = -\frac{1}{2} d(v^3 \partial_v \partial_z K) + \frac{1}{2} \partial_v (v^3 \partial_v \partial_z K) dv \\
&= -\frac{1}{2} d(v^3 \partial_v \partial_z K) + \frac{v^3}{2} \partial_z \left(\frac{1}{v^3} \partial_v (v^3 \partial_v K) \right) dv = -\frac{1}{2} d(v^3 \partial_v \partial_z K) + \frac{v^3}{2} \partial_z \Delta_y K dv \\
&\quad \frac{v^3}{4} (\partial_{u_1}^2 K \partial_z^3 K - 2 \partial_z \partial_{u_1} K \partial_z^2 \partial_{u_1} K + \partial_z^2 K \partial_z \partial_{u_1}^2 K) dv = v^3 \partial_z (\text{deth}) \\
F_4^m &= -\frac{1}{2} d(v^3 \partial_v \partial_z K) \wedge d\Omega'_3 + \frac{v^3}{2} \partial_z (\Delta_y K + 2 \text{deth}) dv \wedge d\Omega'_3
\end{aligned}$$



$$\begin{aligned}
C_3^m &= -\frac{1}{2}v^3\partial_v\partial_zKd\Omega'_3 \\
dC_3^e &= -[\partial_{u_1}(e^{3A}h_{11}) - \partial_z(e^{3A}h_{12})]dt \wedge du_1 \wedge dz \wedge dy \\
dC_3^e &= -[\partial_{u_1}(e^{3A}h_{11}) - \partial_z(e^{3A}h_{12})]dt \wedge du_1 \wedge dz \wedge dy \\
&\quad -[\partial_v(e^{3A}h_{11})dz + \partial_v(e^{3A}h_{12})du_1] \wedge dt \wedge dv \wedge dy \\
F_6^m &= \frac{v^3\text{deth}}{h_{11}}[h_{12}\partial_v(e^{3A}h_{11}) - h_{11}\partial_v(e^{3A}h_{12})]dz \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
&\quad + \frac{v^3\text{deth}}{h_{11}}[h_{22}\partial_v(e^{3A}h_{11}) - h_{12}\partial_v(e^{3A}h_{12})]du_1 \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
&\quad - \frac{v^3(\text{deth})^2}{h_{11}}[\partial_{u_1}(e^{3A}h_{11}) - \partial_z(e^{3A}h_{12})]du_2 \wedge du_3 \wedge dv \wedge d\Omega'_3, \\
F_6^m &= v^3\left(\frac{h_{12}}{h_{11}}\partial_vh_{11} - \partial_vh_{12}\right)dz \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
&\quad + v^3\left(\frac{h_{12}}{h_{11}}\partial_vh_{12} - \partial_vh_{22}\right)du_1 \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
&\quad + v^3\left(\partial_{u_1}(\text{deth}) - \frac{h_{12}}{h_{11}}\partial_z(\text{deth})\right)dv \wedge dx_2 \wedge du_3 \wedge d\Omega'_3 \\
\frac{1}{2}\frac{h_{12}}{h_{11}}d_\Sigma(v^3\partial_v\partial_zK) \wedge du_2 \wedge du_3 \wedge d\Omega'_3 &- \frac{1}{2}d_\Sigma(v^3\partial_v\partial_1K) \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
F_6^m - H \wedge C_3^m &= \left[\frac{h_{12}}{h_{11}}d_\Sigma\left(\frac{v^3}{2}\partial_v\partial_zK\right) + d\left(\frac{h_{12}}{h_{11}}\frac{v^3}{2}\partial_v\partial_zK\right)\right] \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
&\quad - d_\Sigma\left(\frac{v^3}{2}\partial_v\partial_{u_1}K\right) \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
&\quad + v^3\left(\partial_{u_1}(\text{deth}) - \frac{h_{12}}{h_{11}}\partial_z(\text{deth})\right)dv \wedge dx_2 \wedge dx_3 \wedge d\Omega'_3 \\
d\left(\frac{v^3}{2}\frac{h_{12}}{h_{11}}\partial_v\partial_zK\right) \wedge du_2 \wedge du_3 \wedge d\Omega'_3 &- \frac{h_{12}}{h_{11}}\partial_v\left(\frac{v^3}{2}\partial_v\partial_zK\right)dv \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
-d\left(\frac{v^3}{2}\partial_v\partial_{u_1}K\right) \wedge du_2 \wedge du_3 \wedge d\Omega'_3 &+ \partial_v\left(\frac{v^3}{2}\partial_v\partial_{u_1}K\right)dv \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
\frac{v^3}{2}\partial_{u_1}\left[2\text{deth} + \frac{1}{v^3}\partial_v(v^3\partial_vK)\right] &- \frac{v^3}{2}\frac{h_{12}}{h_{11}}\partial_z\left[2\text{deth} + \frac{1}{v^3}\partial_v(v^3\partial_vK)\right] \\
F_6^m - H \wedge C_3^m &= d\left(\frac{v^3}{2}\frac{h_{12}}{h_{11}}\partial_v\partial_zK - \frac{v^3}{2}\partial_v\partial_{u_1}K\right) \wedge du_2 \wedge du_3 \wedge d\Omega'_3 \\
C_5^m &= \frac{v^3}{2}\left(\frac{h_{12}}{h_{11}}\partial_v\partial_zK - \partial_v\partial_{u_1}K\right)du_2 \wedge du_3 \wedge d\Omega'_3 \\
ds^2 &= \frac{1}{\sqrt{\text{deth}}}(-dt^2 + dy^2) + \frac{\sqrt{\text{deth}}}{h_{11}}(du_2^2 + du_3^2) + \sqrt{\text{deth}}(e^{3A}h_{ab}dr^a dr^b + ds_{\mathbb{R}^4}^2) \\
e^{2\phi} &= \frac{\sqrt{\text{deth}}}{h_{11}}, B_2 = \frac{h_{12}}{h_{11}}du_2 \wedge du_3 \\
C_3 &= e^{3A}h_{1a}dt \wedge dr^a \wedge dy - \frac{v^3}{2}\partial_v\partial_zKd\Omega'_3 \\
C_5 &= \frac{1}{h_{11}}dt \wedge du_1 \wedge du_2 \wedge du_3 \wedge dy + \frac{v^3}{2}\left(\frac{h_{12}}{h_{11}}\partial_v\partial_zK - \partial_v\partial_{u_1}K\right)du_2 \wedge du_3 \wedge d\Omega'_3 \\
ds^2 &= Z^{-1/2}(-dt^2 + dx_1^2 + dx_2^2) + Z^{1/2}(dx_3^2 + \dots + dx_9^2) \\
e^\Phi &= Z^{1/4} \\
C_3 &= Z^{-1}dt \wedge dx_1 \wedge dx_2
\end{aligned}$$



$$\begin{aligned}
Z &= 1 + \frac{Q}{r^2}, r^2 \equiv x_6^2 + \dots + x_9^2 \\
x_2 &= x'_2 c + x'_3 s \\
x_3 &= -x'_2 s + x'_3 c \\
W &\equiv c^2 Z + s^2 \\
ds^2 &= Z^{-1/2}(-dt^2 + dx_1^2) + Z^{1/2}(dx_4^2 + \dots + dx_9^2) \\
&\quad + Z^{-1/2}W(dx_3 - cs(Z-1)W^{-1}dx_2)^2 + Z^{1/2}W^{-1}dx_2^2, \\
e^{2\Phi} &= Z^{1/2}, \\
C_3 &= Z^{-1}sdt \wedge dx_1 \wedge (dx_3 - cs(Z-1)W^{-1}dx_2) \\
&\quad + W^{-1}cdt \wedge dx_1 \wedge dx_2. \\
ds^2 &= Z^{-1/2}(-dt^2 + dx_1^2) + Z^{1/2}(dx_4^2 + \dots + dx_9^2) + Z^{1/2}W^{-1}(dx_2^2 + dx_3^2) \\
e^{2\Phi} &= W^{-1}Z, B_2 = -cs(Z-1)W^{-1}dx_2 \wedge dx_3 \\
C_2 &= Z^{-1}sdt \wedge dx_1, C_4 = W^{-1}cdt \wedge dx_1 \wedge dx_2 \wedge dx_3 \\
ds^2 &= Z^{-1/2}(-dt^2 + dx_1^2 + dx_4^2) + Z^{1/2}(dx_5^2 + \dots + dx_9^2) + Z^{1/2}W^{-1}(dx_2^2 + dx_3^2) \\
e^{2\Phi} &= Z^{1/2}W^{-1} \\
B_2 &= -cs(Z-1)W^{-1}dx_2 \wedge dx_3 \\
C_3 &= Z^{-1}sdt \wedge dx_1 \wedge dx_4 \\
C_5 &= W^{-1}cdt \wedge dx_1 \wedge dx_2 \wedge dx_3 \wedge dx_4 \\
h_{11} &= c^2 Z + s^2, h_{22} = s^2 Z + c^2 \\
h_{12} &= cs(Z-1), \text{deth} = Z \\
dB_2 \wedge C_3 &= [-cs\partial_l((Z-1)W^{-1})dx_l \wedge dx_2 \wedge dx_3] \wedge [Z^{-1}sdt \wedge dx_1 \wedge dx_4] \\
&= cs^2 W^{-2}Z^{-1}(\partial_l Z)dt \wedge dx_1 \wedge dx_2 \wedge dx_3 \wedge dx_4 \wedge dx_l \\
dC_5 &= c\partial_l(W^{-1})dx_l \wedge dt \wedge dx_1 \wedge \dots \wedge dx_4 \\
&= c^3 W^{-2}(\partial_l Z)dt \wedge dx_1 \wedge dx_2 \wedge dx_3 \wedge dx_4 \wedge dx_l \\
F_6 &= cW^{-1}Z^{-1}(\partial_l Z)dt \wedge dx_1 \wedge dx_2 \wedge dx_3 \wedge dx_4 \wedge dx_l \\
&= cZ^{-1}(\partial_l Z)e^0 \wedge e^1 \wedge e^2 \wedge e^3 \wedge e^4 \wedge e^l \\
F_4^{(m)} &= c \frac{\partial_l Z}{Z} \frac{\epsilon_{(l-4),abcd}}{4!} e^{4+a} \wedge e^{4+b} \wedge e^{4+c} \wedge e^{4+d} \\
&= c(\partial_l Z) \frac{\epsilon_{(l-4),abcd}}{4!} dx^{4+a} \wedge dx^{4+b} \wedge dx^{4+c} \wedge dx^{4+d} \\
F_4^{(m)} &= -c \frac{x_l}{r} (\partial_r Z) \frac{\epsilon_{(l-5),abc}}{3!} dx^5 \wedge dx^{5+a} \wedge dx^{5+b} \wedge dx^{5+c} \\
C_3^{(m)} &= -c \frac{x_5 x_l}{r} (\partial_r Z) \frac{\epsilon_{(l-5),abc}}{3!} dx^{5+a} \wedge dx^{5+b} \wedge dx^{5+c} \\
&= -cx_5 r^3 (\partial_r Z) d\Omega'_3 \\
ds^2 &= Z^{-1/6} W^{1/3} [Z^{-1/2}(-dt^2 + dx_1^2 + dx_4^2) + Z^{1/2}(dx_5^2 + \dots + dx_9^2)] \\
&\quad + Z^{1/3} W^{-2/3} (dx_2^2 + dx_3^2 + dx_{11}^2) \\
C_3 &= Z^{-1}sdt \wedge dx_1 \wedge dx_4 - cs(Z-1)W^{-1}dx_2 \wedge dx_3 \wedge dx_{11} - cx_5 r^3 (\partial_r Z) d\Omega'_3 \\
x_4 &\rightarrow cu_1 + sz, x_5 \rightarrow -su_1 + cz \\
x_1 &\rightarrow y, x_2 \rightarrow u_2, x_3 \rightarrow u_3, x_{11} \rightarrow -u_4, x_{6,7,8,9} \rightarrow v_{1,2,3,4} \\
ds^2 &= W^{1/3} Z^{-2/3} (-dt^2 + dy^2) + W^{4/3} Z^{-2/3} (dz - cs(Z-1)W^{-1}du_1)^2 \\
&\quad + W^{1/3} Z^{1/3} (dv_1^2 + \dots + dv_4^2) + W^{-2/3} Z^{1/3} (du_1^2 + du_2^2 + du_3^2 + du_4^2) \\
C_3 &= Z^{-1}sdt \wedge dy \wedge (sdz + cdu_1) + cs(Z-1)W^{-1}du_2 \wedge du_3 \wedge du_4 + c(su_1 - cz)v^3 (\partial_v Z) d\Omega'_3 \\
e^{A_0} &= W^{1/6} Z^{-1/3}, (-\partial_z w) = W \\
(\partial_{u_1} w) &= cs(Z-1) \\
(\partial_{v_l} w) &= c(su_1 - cz)v_l \frac{\partial_v Z}{v} \\
\delta C_3 &= -c^2 dt \wedge dy \wedge dz + csdt \wedge dy \wedge du_1 \\
w &\equiv -zW + cs(Z-1)u_1
\end{aligned}$$

$$G_0 = -\frac{1}{2}Z(cz - su_1)^2 - \frac{1}{2}(sz + cu_1)^2 + f(v)$$

Entiéndase que la supergravedad cuántica, para efectos de este trabajo, comporta la simetría entre dos partículas o antipartículas, según sea el caso, de las cuales, una de ellas, es una superpartícula (Véase la definición proporcionada por este autor respecto de las superpartículas en sentido lato), a propósito de la deformación o perforación del espacio cuántico de que se trate, combinando en consecuencia, relatividad general y supersimetría. Entiéndase por supersimetría, para efectos de este trabajo, comporta la interacción de dos partículas o antipartículas, según sea el caso, de las cuales, una de ellas, es una superpartícula (Véase la definición proporcionada por este autor), a propósito de la deformación o perforación del espacio cuántico de que se trate, por acción de las superpartículas. Entiéndase que las supermembranas, para efectos de este trabajo, comporta la existencia de infinitas dimensiones a propósito de la deformación o perforación del espacio cuántico de que se trate, por acción de las superpartículas. Finalmente, entiéndase por superespacio, para efectos de este trabajo, como la existencia de un espacio cuántico relativista, el mismo que posee dimensiones ordinarias y anticonmutativas, a propósito de la deformación o perforación del espacio cuántico de que se trate, por acción de las superpartículas.

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Iosif Bena, Anthony Houppe, Dimitrios Toulikas y Nicholas P. Warner, Maze Topiary in Supergravity, arXiv:2312.02286v2 [hep-th] 21 Jan 2025.



Apéndice J.

1. Agujeros negros cuánticos, supermembranas, superspacios, dimensión temporal y supergravedad cuántica para campos cuánticos relativistas o curvos.

$$S_{EH}[g] = \frac{c^3}{16\pi G_N^{(d)}} \int_{\mathcal{M}} d^d x \sqrt{|g|} R(g) + (-1)^d \frac{c^3}{8\pi G_N^{(d)}} \int_{\partial\mathcal{M}} d^{d-1}\Sigma \mathcal{K} + S_{\text{materia}}$$

$$d^{d-1}\Sigma \equiv n^2 d^{d-1}\Sigma_\rho n^\rho$$

$$d^{d-1}\Sigma_\rho = \frac{1}{(d-1)! \sqrt{|g|}} \epsilon_{\rho\mu_1 \dots \mu_{d-1}} dx^{\mu_1} \wedge \dots dx^{\mu_{d-1}}$$

$$\vec{F} = -\frac{8(d-3)\pi G_N^{(d)} mM}{(d-2)\omega_{d-2}} \frac{\vec{x}_{d-1}}{|\vec{x}_{d-1}|^{d-1}}$$

$$\mathcal{Z} = \int Dg e^{+iS_{EH}/\hbar}$$

$$\frac{S_{EH}}{\hbar} = \frac{2\pi}{\ell_{\text{Planck}}^{d-2}} \int d^d x \dots$$

$$\frac{\ell_{\text{Planck}}^{d-2}}{2\pi} = \frac{16\pi G_N^{(d)} \hbar}{c^3}$$

$$\ell_{\text{Planck}} = \frac{\ell_{\text{Planck}}}{2\pi}.$$

$$\lambda_{\text{Compton}} = \frac{\hbar}{Mc}$$

$$R_s = \left(\frac{16\pi M G_N^{(d)} c^{-2}}{(d-2)\omega_{(d-2)}} \right)^{\frac{1}{d-3}}$$

$$M_{\text{Planck}} = \left(\frac{\hbar^{d-3}}{G_N^{(d)} c^{d-5}} \right)^{\frac{1}{d-2}}$$

$$\frac{c^3}{G_N^{(d)} \hbar} = \left(\frac{M_{\text{Planck}} c}{\hbar} \right)^{d-2}$$

$$M \sim M_{\text{Planck}} \Rightarrow \lambda_{\text{Compton}} \sim R_s \sim \ell_{\text{Planck}}$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 0. \Rightarrow R_{\mu\nu}$$

$$ds^2 = W(r)(dct)^2 - W^{-1}(r)dr^2 - R^2(r)d\Omega_{(2)}^2$$

$$d\Omega_{(2)}^2 = d\theta^2 + \sin^2 \theta d\varphi^2$$

$$ds^2 = W(dct)^2 - W^{-1}dr^2 - r^2 d\Omega_{(2)}^2, W = 1 + \frac{\omega}{r}$$

$$M = -\frac{\omega c^2}{2G_N^{(4)}}, \Rightarrow \omega = -R_s$$

$$M = \frac{1}{8\pi G_N^{(4)}} \int_{S_\infty^2} d^2 S_i (\partial_j g_{ij} - \partial_i g_{jj})$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = \frac{8\pi G_N^{(4)}}{c^4} T_{\text{materia } \mu\nu}$$

$$ds^2 = \left(1 - \frac{R_s}{r}\right) dv^2 - 2dvdr - r^2 d\Omega_{(2)}^2$$



$$\begin{aligned}
v &= ct + r + R_S \log \left| 1 - \frac{R_S}{r} \right| \\
ds^2 &= \frac{4R_S^3 e^{-r/R_S}}{r} [(dcT)^2 - dX^2] - r^2 d\Omega_{(2)}^2 \\
\left(\frac{r}{R_S} - 1 \right) e^{r/R_S} &= X^2 - c^2 T^2 \\
\frac{ct}{R_S} &= \ln \left(\frac{X + cT}{X - cT} \right) = 2 \operatorname{arcth}(cT/X) \\
ds^2 &= W^{\frac{2M}{\omega}-1} W dt^2 - W^{1-\frac{2M}{\omega}} [W^{-1} dr^2 + r^2 d\Omega_{(2)}^2] \\
\varphi &= \varphi_0 + \frac{\Sigma}{\omega} \ln W \\
W &= 1 + \frac{\omega}{r}, \omega = \pm 2\sqrt{M^2 + \Sigma^2} \\
A &= \int_{r=R_S} d\theta d\varphi r^2 = 4\pi R_S^2 \\
\kappa^2 &= -\frac{1}{2} (\nabla^\mu k^\nu) (\nabla_\mu k_\nu) \Big|_{\text{horizonte}} \\
ds^2 &= g_{tt}(r) dt^2 + g_{rr}(r) dr^2 - r^2 d\Omega_{(2)}^2 \\
\kappa &= \frac{1}{2} \frac{\partial_r g_{tt}}{\sqrt{-g_{tt} g_{rr}}}, \\
\kappa &= \frac{1}{4G_N^{(4)} M} \\
dE &= T dS \\
dM &\sim \frac{1}{G_N^{(4)}} \kappa dA \\
dM &= \frac{1}{8\pi G_N^{(4)}} \kappa dA \\
M &= \frac{1}{4\pi G_N^{(4)}} \kappa A \\
T &= \frac{\hbar \kappa}{2\pi c} \\
S &= \frac{1}{4\hbar G_N^{(4)}} \\
S &= \frac{1}{32\pi^2} \frac{A}{\ell_{\text{Planck}}^2} \\
T &= \frac{\hbar c^3}{8\pi G_N^{(4)} M}, S = \frac{4\pi G_N^{(4)} M^2}{\hbar c} \\
dM c^2 &= T dS, M c^2 = 2TS \\
C^{-1} &= \frac{\partial T}{\partial M} = \frac{-\hbar c^3}{8\pi G_N^{(4)} M^2} < 0 \\
S(E) &= \log \rho(E) \\
\rho(M) &\sim e^{M^2}
\end{aligned}$$



$$\begin{aligned}
Z &= \text{Tr} e^{-(H-\mu_i C_i)/T} \\
W &= E - TS - \mu_i C_i \\
e^{-\beta W} &= Z \\
S &= \frac{1}{T} (E - \mu_i C_i) + \log Z \\
Z &= \int Dg e^{-\tilde{S}_{EH}/\hbar} \\
S_{EH}[g] &= \frac{c^3}{16\pi G_N^{(4)}} \int_{\mathcal{M}} d^4x \sqrt{|g|} R + \frac{c^3}{8\pi G_N^{(4)}} \int_{\partial\mathcal{M}} (\mathcal{K} - \mathcal{K}_0) \\
Z &= e^{-\tilde{S}_{EH}(\text{on-shell})} \\
\left(\frac{r}{R_S} - 1\right) e^{r/R_S} &= X^2 - T^2 \\
\left(\frac{r}{R_S} - 1\right) e^{r/R_S} &= X^2 + \mathcal{T}^2 > 0 \\
\frac{X+T}{X-T} &= e^{t/R_S} \\
\frac{X-i\mathcal{T}}{X+\mathcal{T}} &= e^{-2i\text{Arg}(X+i\mathcal{T})} = e^{-it/R_S} \\
n_\mu &= -\frac{\delta_{\mu r}}{\sqrt{-n^2}} = -\sqrt{-g_{rr}} \delta_{\mu r} \\
ds_{(3)}^2 &= h_{\mu\nu} dx^\mu dx^\nu = g_{tt} dt^2 - r^2 d\Omega_{(2)}^2 \Big|_{r=r_0} \\
\nabla_\mu n_\nu &= -\sqrt{-g_{rr}} \{ \delta_{\mu r} \delta_{\nu r} \partial_r \log \sqrt{-g_{rr}} - \Gamma_{\mu\nu}^r \} \\
\mathcal{K} &= h^{\mu\nu} \nabla_\mu n_\nu = \frac{1}{\sqrt{-g_{rr}}} \left\{ \frac{1}{2} \partial_r \log g_{tt} + \frac{2}{r} \right\} \Big|_{r=r_0} \\
\mathcal{K}_0 &= \frac{2}{r} \Big|_{r=r_0} \\
g_{tt} &\sim 1 - \frac{2M}{r}, g_{rr} \sim -\left(1 + \frac{2M}{r}\right) \\
(\mathcal{K} - \mathcal{K}_0)|_{r=r_0} &\sim -\frac{M}{r_0^2} \\
\frac{i}{8\pi} \int_{r_0 \rightarrow \infty} d^3x \sqrt{|h|} (\mathcal{K} - \mathcal{K}_0) &= \lim_{r_0 \rightarrow \infty} \frac{i}{8\pi} \int_0^{-i\beta} dt \int_{S^2} d\Omega^2 r_0^2 \sqrt{g_{tt}(r_0)} (\mathcal{K} - \mathcal{K}_0) \\
&= \lim_{r_0 \rightarrow \infty} \frac{\beta}{2} r_0^2 (\mathcal{K} - \mathcal{K}_0) = -\frac{\beta M}{2} \\
S &= \beta M + \log Z = \frac{\beta M}{2} = 4\pi M^2 \\
S_{EM}[g, A] &= S_{EH}[g] + \frac{1}{c} \int d^d x \sqrt{|g|} \left[-\frac{1}{4} F^2 \right] \\
F_{\mu\nu} &= 2\partial_{[\mu} A_{\nu]} \\
A'_\mu &= A_\mu + \partial_\mu \Lambda \\
R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R - \frac{8\pi G_N^{(4)}}{c^3} T_{\mu\nu} &= 0 \\
\nabla_\mu F^{\mu\nu} &= 0 \\
T_{\mu\nu} &= \frac{-2c}{\sqrt{|g|}} \frac{\delta S_M[A]}{\delta g^{\mu\nu}} = F_{\mu\rho} F_\nu^\rho - \frac{1}{4} g_{\mu\nu} F^2 \\
\nabla_\mu {}^* F^{\mu\nu} &
\end{aligned}$$

$$\begin{aligned}
A &\equiv A_\mu dx^\mu, F = \frac{1}{2} F_{\mu\nu} dx^\mu \wedge dx^\nu \equiv dA \\
d^*F &= 0, dF = 0. (\partial_{[\alpha} F_{\beta\gamma]} = 0) \\
\frac{1}{c^2} \int d^d x \sqrt{|g|} [-A_\mu j^\mu] \\
\nabla_\mu j^\mu &= 0, (d^*j = 0) \\
\partial_\mu j^\mu &= 0 \\
\nabla_\mu F^{\mu\nu} &= \frac{1}{c} j^\nu, \left(d^*F = \frac{1}{c} \star j \right) \\
j^\mu(y) &= qc \int_\gamma dX^\mu \frac{1}{\sqrt{|g|}} \delta^{(4)}(y - X(\xi)) \\
j^\mu(y^0, \vec{y}) &= qc \int dX^0 \frac{dX^\mu}{dX^0} \frac{1}{\sqrt{|g|}} \delta^{(3)}(\vec{y} - \vec{X}) \delta(y^0 - X^0) = qV^\mu \frac{\delta^{(3)}(\vec{y} - \vec{X}(y^0))}{\sqrt{|g|}} \\
j^\mu(y^0, \vec{y}) &= qc \delta^{\mu 0} \frac{\delta^{(3)}(\vec{y})}{\sqrt{|g|}} \\
-\frac{q}{c} \int_{\gamma(\xi)} A_\mu \dot{x}^\mu d\xi &= -\frac{q}{c} \int_\gamma A \\
S_{M,q}[X^\mu(\xi)] &= -Mc \int d\xi \sqrt{g_{\mu\nu}(X) \dot{X}^\mu \dot{X}^\nu} - \frac{q}{c} \int A_\mu \dot{X}^\mu \\
0 &= \int_V d^*j \\
\int_V d^*j &= \int_{x^0=x_2^0}^* j - \int_{x^0=x_1^0}^* j = 0 \\
q &= \frac{1}{c} \int_{x^0=\text{constant}}^* j, \\
q &= \int_{S_\infty^2}^* F \\
S_{EM}[g, A] &= \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{|g|} \left[R - \frac{1}{4} F^2 \right] \\
q &= \frac{1}{16\pi G_N^{(d)}} \int_{S_\infty^{d-2}}^* F \\
E_r = F_{0r} &\sim \frac{4G_N^{(4)} q}{r^2} \\
F_{tr} &\sim \pm \frac{1}{R^2(r)} \\
ds^2 &= f(r) dt^2 - f^{-1}(r) dr^2 - r^2 d\Omega_{(2)}^2 \\
F_{tr} &= -\frac{4G_N^{(4)} q}{r^2} \\
f(r) &= r^{-2} (r - r_+) (r - r_-) \\
r_\pm &= G_N^{(4)} M \pm r_0, r_0 = G_N^{(4)} (M^2 - 4q^2)^{1/2} \\
A_\mu &= \delta_{\mu t} \frac{-4G_N^{(4)} q}{r} \\
A &= 4\pi r_+^2
\end{aligned}$$

$$A_{\text{extreme}} = 4\pi r_+^2 = 4\pi \left(G_N^{(4)} M \right)^2$$

$$F_{12} = -G_N^{(4)} \frac{M_1 M_2}{r_{12}^2} + 4G_N^{(4)} \frac{q_1 q_2}{r_{12}^2}$$

$$ds^2 = H^{-2} dt^2 - H^2 d\vec{x}_3^2$$

$$A_\mu = -2\text{sign}(q)(H^{-1} - 1)\delta_{\mu t},$$

$$H = 1 + \frac{G_N^{(4)} M}{|\vec{x}_3|}$$

$$\partial_{\underline{i}} \partial_{\underline{i}} H = 0$$

$$ds^2 = H^{-2} dt^2 - H^2 d\vec{x}_3^2,$$

$$A_\mu = \delta_{\mu t} \alpha (H^{-1} - 1), \alpha = \pm 2$$

$$\partial_{\underline{i}} \partial_{\underline{i}} H = 0.$$

$$H(\vec{x}_3) = 1 + \sum_{i=1}^N \frac{2G_N^{(4)} |q_i|}{|\vec{x}_3 - \vec{x}_{3,i}|}$$

$$ds^2 = \frac{\rho^2}{R_{AdS}^2} dt^2 - R_{AdS}^2 \frac{d\rho^2}{\rho^2} - R_{AdS}^2 d\Omega_{(2)}^2$$

$$A_t = -\frac{2\rho}{R_{AdS}}, F_{\rho t} = -\frac{2}{R_{AdS}}$$

$$ds^2 = H^{-2} W dt^2 - H^2 [W^{-1} d\rho^2 + \rho^2 d\Omega_{(2)}^2]$$

$$A_\mu = \delta_{\mu t} \alpha (H^{-1} - 1)$$

$$H = 1 + \frac{h}{\rho}, W = 1 + \frac{\omega}{\rho}, \omega = h \left[1 - \left(\frac{\alpha}{2} \right)^2 \right]$$

$$\alpha = -\frac{4G_N^{(4)} q}{r_\pm}, h = r_\pm, \omega = \pm 2r_0$$

$$S[g, A^I] = \frac{1}{16\pi G_N^{(4)}} \int d^4x \sqrt{|g|} \left[R - \frac{1}{4} \sum_{I=1}^{I=N} (F^I)^2 \right]$$

$$ds^2 = H^{-2} W dt^2 - H^2 [W^{-1} d\rho^2 + \rho^2 d\Omega_{(2)}^2]$$

$$A_\mu^i = \delta_{\mu t} \alpha^i (H^{-1} - 1)$$

$$H = 1 + \frac{h}{\rho}, W = 1 + \frac{\omega}{\rho}, \omega = h \left[1 - \sum_{i=1}^{i=N} \left(\frac{\alpha^i}{2} \right)^2 \right]$$

$$\alpha^i = -\frac{4G_N^{(4)} q^i}{r_\pm}, h = r_\pm, \omega = \pm 2r_0$$

$$r_\pm = G_N^{(4)} M \pm r_0, r_0 = G_N^{(4)} \sqrt{M^2 - 4 \sum_{i=1}^{i=N} q_i^2}$$

$$dM = \frac{1}{8\pi G_N^{(4)}} \kappa dA + \Phi dq$$

$$\kappa = \frac{1}{G_N^{(4)}} \frac{\sqrt{M^2 - 4q^2}}{(M + \sqrt{M^2 - 4q^2})^2}$$



$$T = \frac{1}{2\pi G_N^{(4)}} \frac{\sqrt{M^2 - 4q^2}}{(M + \sqrt{M^2 - 4q^2})^2}, S = \pi G_N^{(4)} (M + \sqrt{M^2 - 4q^2})^2$$

$$\vec{\tilde{E}} = \vec{B}, \vec{\tilde{B}} = -\vec{E}$$

$$\tilde{F} = aF + b^{\star}F, \Rightarrow^{\star} \tilde{F} = -bF + a^{\star}F, a^2 + b^2 \neq 0$$

$$\vec{F} \equiv \begin{pmatrix} F \\ \star F \end{pmatrix}, \star \vec{F} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \vec{F}$$

$$\nabla_\mu \vec{F}^{\mu\nu} = 0$$

$$\vec{\tilde{F}} = M\vec{F}, M = \begin{pmatrix} a & b \\ -b & a \end{pmatrix}$$

$$\int_{S^2_\infty} \star \vec{F} = \begin{pmatrix} 16\pi G_N^{(4)} q \\ p \end{pmatrix} \equiv 16\pi G_N^{(4)} \vec{q}, \vec{q} = \begin{pmatrix} q \\ p/16\pi G_N^{(4)} \end{pmatrix}$$

$$\tilde{A}_\mu(x) = -\int_0^1 d\lambda \lambda x^\nu \frac{\epsilon_{\mu\nu}^{\rho\sigma}}{\sqrt{|g|}} \partial_\rho A_\sigma(\lambda x)$$

$$G_{\mu\nu} - \vec{F}_\mu^{T\rho} \vec{F}_{\nu\rho}$$

$$\vec{F} \equiv \begin{pmatrix} e^{-2F} \\ \star F \end{pmatrix}, G_{\mu\nu} + (\vec{F}_\mu^\rho)^T \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \vec{F}_{\nu\rho}$$

$$M = \begin{pmatrix} a & 0 \\ 0 & 1/a \end{pmatrix}, \quad e' = a^{-1}e$$

$$M = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}, \quad e' = \frac{1}{e}$$

$$\begin{cases} \tilde{F} = \cos \xi F + \sin \xi \star F \\ \star \tilde{F} = -\sin \xi F + \cos \xi \star F \end{cases}$$

$$\begin{cases} \tilde{F}_{tr} = \frac{-4G_N^{(4)} \cos \xi q}{r^2} \\ \tilde{F}_{\theta\varphi} = 4G_N^{(4)} \sin \xi q \sin \theta \end{cases}$$

$$\tilde{q} = \cos \xi q, \tilde{p} = -16G_N^{(4)} \sin \xi q, \Rightarrow \tilde{q}^2 + \left(\frac{\tilde{p}}{16\pi G_N^{(4)}} \right) = q^2$$

$$ds^2 = f(r) dt^2 - f^{-1}(r) dr^2 - r^2 d\Omega_{(2)}^2$$

$$F_{tr} = -\frac{4G_N^{(4)} q}{r^2}, F_{\theta\varphi} = -\frac{p}{4\pi} \sin \theta,$$

$$f(r) = r^{-2}(r-r_+)(r-r_-),$$

$$r_{\pm} = G_N^{(4)} M \pm r_0, r_0 = G_N^{(4)} \left\{ M^2 - 4 \left[q^2 + \left(\frac{p}{16\pi G_N^{(4)}} \right)^2 \right] \right\}^{1/2}$$

$$[Q^\alpha, M_{ab}] = \Gamma_s(M_{ab})^\alpha{}_\beta Q^\beta$$

$$\{Q^\alpha, Q^\beta\} = i(\gamma^a \mathcal{C}^{-1})^{\alpha\beta} P_a$$

$$[M_{ab}, M_{cd}] = -M_{eb} \Gamma_v(M_{cd})^e{}_a - M_{ae} \Gamma_v(M_{cd})^e{}_b$$

$$[P_a, M_{bc}] = -P_e \Gamma_v(M_{bc})^e{}_a,$$

$$[Q^\alpha, M_{ab}] = \Gamma_s(M_{ab})^\alpha{}_\beta Q^\beta,$$

$$\{Q^\alpha, Q^\beta\} = i(\gamma^a \mathcal{C}^{-1})^{\alpha\beta} P_a.$$

$$A_\mu = e^a{}_\mu P_a + \frac{1}{2} \omega_\mu{}^{ab} M_{ab} + \bar{\psi}_{\mu\alpha} Q^\alpha$$

$$\begin{aligned}
\Lambda &= \sigma^a P_a + \frac{1}{2} \sigma^{ab} M_{ab} + \bar{\epsilon}_\alpha Q^\alpha \\
\delta A_\mu &= \partial_\mu \Lambda + \Lambda, A_\mu \equiv \mathbb{D}_\mu \Lambda \\
S[e^a{}_\mu, \omega_\mu{}^{ab}, \psi_\mu] &= \int d^4x e [R(e, \omega) + 2e^{-1} \epsilon^{\mu\nu\rho\sigma} \bar{\psi}_\mu \gamma_5 \gamma_\nu \nabla_\rho \psi_\sigma] \\
R(e, \omega) &= e_a{}^\mu e_b{}^\nu R_{\mu\nu}{}^{ab}(\omega) \\
\omega_{abc} &= -\Omega_{abc} + \Omega_{bca} - \Omega_{cab}, \\
\Omega_{\mu\nu}{}^a &= \Omega_{\mu\nu}{}^a(e) + \frac{1}{2} T_{\mu\nu}{}^a, \\
\Omega_{\mu\nu}{}^a(e) &= \partial_{[\mu} e^a{}_{\nu]}, T_{\mu\nu}{}^a = i \bar{\psi}_\mu \gamma^a \psi_\nu. \\
\left\{ \begin{aligned} \delta_\xi x^\mu &= \xi^\mu \\ \delta_\xi e^a{}_\mu &= -\xi^\nu \partial_\nu e^a{}_\mu - \partial_\mu \xi^\nu e^a{}_\nu \\ \delta_\xi \psi_\mu &= -\xi^\nu \partial_\nu \psi_\mu - \partial_\mu \xi^\nu \psi_\nu \end{aligned} \right. \\
\left\{ \begin{aligned} \delta_\sigma e^a{}_\mu &= \sigma^a{}_b e^b{}_\mu \\ \delta_\sigma \psi_\mu &= \frac{1}{2} \sigma^{ab} \gamma_{ab} \psi_\mu \\ \delta_\epsilon e^a{}_\mu &= -i \bar{\epsilon} \gamma \psi_\mu \\ \delta_\epsilon \psi_\mu &= \nabla_\mu \epsilon \end{aligned} \right. \\
[M_{ab}, M_{cd}] &= -M_{eb} \Gamma_\nu (M_{cd})^e{}_a - M_{ae} \Gamma_\nu (M_{cd})^e{}_b, \\
[P_a, M_{bc}] &= -P_e \Gamma_\nu (M_{bc})^e{}_a, \\
[Q^{ai}, M_{ab}] &= \Gamma_s (M_{ab})^\alpha{}_\beta Q^{\beta i}, \\
\{Q^{ai}, Q^{\beta j}\} &= i \delta^{ij} (\gamma^a \mathcal{C}^{-1})^{\alpha\beta} P_a - i (\mathcal{C}^{-1})^{\alpha\beta} Q^{ij} - \gamma_5 (\mathcal{C}^{-1})^{\alpha\beta} P^{ij} \\
A_\mu &= e^a{}_\mu P_a + \frac{1}{2} \omega_\mu{}^{ab} M_{ab} + \frac{1}{2} A^{ij}{}_\mu Q^{ij} + \bar{\psi}^i{}_\mu Q^{i\alpha} \\
&\quad \frac{1}{n!} (\gamma^{a_1 \dots a_n} \mathcal{C}^{-1})^{\alpha\beta} Z_{a_1 \dots a_n}^{ij} \\
[Z_{c_1 \dots c_n}^{kl}, M_{ab}] &= -n \Gamma_\nu (M_{ab})^e{}_{[c_1} Z_{|e|c_2 \dots c_n]}^{kl} \\
&\quad \mathcal{C}^{-1}, \gamma_5 \mathcal{C}^{-1}, \gamma_5 \gamma_a \mathcal{C}^{-1}, \gamma_{abc} \mathcal{C}^{-1}, \gamma_{abcd} \mathcal{C}^{-1} \\
&\quad \gamma_a \mathcal{C}^{-1}, \gamma_{ab} \mathcal{C}^{-1}, \gamma_5 \gamma_{ab} \mathcal{C}^{-1}, \gamma_5 \gamma_{abc} \mathcal{C}^{-1} \\
\{Q^{ai}, Q^{\beta j}\} &= i \delta^{ij} (\gamma^a \mathcal{C}^{-1})^{\alpha\beta} P_a + i (\mathcal{C}^{-1})^{\alpha\beta} Z^{[ij]} + \gamma_5 (\mathcal{C}^{-1})^{\alpha\beta} \tilde{Z}^{[ij]} \\
&\quad + (\gamma^a \mathcal{C}^{-1})^{\alpha\beta} Z_a^{(ij)} + i (\gamma_5 \gamma^a \mathcal{C}^{-1})^{\alpha\beta} Z_a^{[ij]} \\
&\quad + i (\gamma^{ab} \mathcal{C}^{-1})^{\alpha\beta} Z_{ab}^{(ij)} + (\gamma_5 \gamma^{ab} \mathcal{C}^{-1})^{\alpha\beta} \tilde{Z}_{ab}^{(ij)} \\
\delta g_{\mu\nu} &= -2 \nabla_{(\mu} k_{\nu)} \\
\delta_\epsilon B &\sim \epsilon F = 0, \\
\delta_\epsilon F &\sim \left\{ \partial \epsilon + B \epsilon \right\} = 0 \\
\delta_\epsilon \psi_\mu &= \nabla_\mu \epsilon = 0 \\
(1 - \gamma^0 \gamma^1) \epsilon &= 0 \\
k^\mu &\sim \bar{\epsilon} \gamma^\mu \epsilon \\
\delta_\epsilon |s\rangle &\sim \bar{\epsilon}_\alpha Q^{i\alpha} |s\rangle = 0 \\
\bar{\epsilon} \mathfrak{M}_\epsilon &= 0, \\
\mathfrak{M} &\equiv i \delta^{ij} \gamma^a P_a + i Z^{[ij]} + \gamma_5 \tilde{Z}^{[ij]} + \gamma^a Z_a^{(ij)} + i \gamma_5 \gamma^a Z_a^{[ij]} + i \gamma^{ab} Z_{ab}^{(ij)} + \gamma_5 \gamma^{ab} \tilde{Z}_{ab}^{(ij)} \\
\mathfrak{M} &= i p \gamma^0 (1 \pm \gamma^0 \gamma^1) \\
\mathfrak{M} &= i \gamma^0 M \left(\delta^{ij} + \frac{Q}{M} \gamma^0 \epsilon^{ij} \right) \\
(\delta^{ij} \pm \gamma^0 \epsilon^{ij}) \epsilon^j &= 0
\end{aligned}$$

$$\begin{aligned}
M &= |Z| \\
M &= |Z_1| = |Z_2| \\
M &= |Z_1| \neq |Z_2|. \\
M &= |Z_1| = |Z_2| = |Z_3| = |Z_4|, \\
M &= |Z_1| = |Z_2| \neq |Z_{3,4}|, \\
M &= |Z_1| \neq |Z_{2,3,4}|. \\
&(\delta^{ij} + \gamma^0 \alpha^{ij}) \\
\mathfrak{M} &= i\gamma^0 M \left(\delta^{ij} + \frac{Z^{(p)}}{M} \gamma^0 \gamma^1 \dots \gamma^p \alpha^{ij} \right) \\
M &\geq |Z_i|, i = 1, \dots, [N/2] \\
\left\{ e^a{}_\mu, \psi_\mu = \begin{pmatrix} \psi_\mu^1 \\ \psi_\mu^2 \end{pmatrix}, A_\mu \right\} \\
S &= \int d^4x e \{ R(e, \omega) + 2e^{-1} \epsilon^{\mu\nu\rho\sigma} \bar{\psi}_\mu \gamma_5 \gamma_\nu \nabla_\rho \psi_\sigma - \mathcal{F}^2 + \mathcal{J}_{(m)}{}^{\mu\nu} (\mathcal{J}_{(e)\mu\nu} + \mathcal{J}_{(m)\mu\nu}) \} \\
&\begin{cases} \mathcal{F}_{\mu\nu} = \tilde{F}_{\mu\nu} + \mathcal{J}_{(m)\mu\nu} \\ \tilde{F}_{\mu\nu} = F_{\mu\nu} + \mathcal{J}_{(e)\mu\nu} \\ F_{\mu\nu} = 2\partial_{[\mu} A_{\nu]} \end{cases} \\
&\begin{cases} \mathcal{J}_{(e)\mu\nu} = i\bar{\psi}_\mu \sigma^2 \psi_\nu \\ \mathcal{J}_{(m)\mu\nu} = -\frac{1}{2e} \epsilon^{\mu\nu\rho\sigma} \bar{\psi}_\rho \gamma_5 \sigma^2 \psi_\sigma \end{cases} \\
T_{\mu\nu}{}^a &= i\bar{\psi}_\mu \gamma^a \psi_\nu (\equiv i\bar{\psi}_{j\mu} \gamma^a \psi_\nu^j) \\
&\begin{cases} \delta_\epsilon e^a{}_\mu = -i\bar{\epsilon} \gamma^a \psi_\mu \\ \delta_\epsilon A_\mu = -i\bar{\epsilon} \sigma^2 \psi_\mu \\ \delta_\epsilon \psi_\mu = \tilde{\nabla}_\mu \epsilon \end{cases} \\
&\tilde{\nabla}_\mu = \nabla_\mu + \frac{1}{4} \tilde{F} \gamma_\mu \sigma^2 \\
&\begin{cases} \tilde{F}'_{\mu\nu} = \cos \theta \tilde{F}_{\mu\nu} + \sin \theta \star \tilde{F}_{\mu\nu} \\ \psi'_\mu = e^{\frac{i}{2}\theta\gamma_5} \psi_\mu \end{cases} \\
&\{e^a{}_\mu, A^{(n)}{}_\mu, \phi, a, \psi_\mu^i, \lambda^i\} \\
S &= \int d^4x \sqrt{|g|} \left\{ R + 2(\partial\phi)^2 + \frac{1}{2} e^{4\phi} (\partial a)^2 - e^{-2\phi} \sum_{n=1}^6 F^{(n)} F^{(n)} + a \sum_{n=1}^6 F^{(n)\star} F^{(n)} \right\} \\
&\tau = a + i e^{-2\phi} \\
&\tilde{F}_{\mu\nu}^{(n)} = \partial_\mu \tilde{A}_\nu^{(n)} - \partial_\nu \tilde{A}_\mu^{(n)} \\
\Lambda &= \begin{pmatrix} a & b \\ c & d \end{pmatrix}, ad - bc = 1 \\
&\begin{pmatrix} \tilde{F}^{(n)}{}_{\mu\nu} \\ F^{(n)}{}_{\mu\nu} \end{pmatrix} \rightarrow \Lambda \begin{pmatrix} \tilde{F}^{(n)}{}_{\mu\nu} \\ F^{(n)}{}_{\mu\nu} \end{pmatrix} \\
&\tau \rightarrow \frac{a\tau + b}{c\tau + d} \\
M^2 + \frac{|Z_1 Z_2|^2}{M^2} - |Z_1|^2 - |Z_2|^2 &\geq 0 \\
\partial_{\underline{i}} \partial_{\underline{i}} \mathcal{H}_1 = \partial_{\underline{i}} \partial_{\underline{i}} \mathcal{H}_2 &= 0
\end{aligned}$$



$$\sum_{n=1}^N \left(k^{(n)}\right)^2=0, \sum_{n=1}^N \left|k^{(n)}\right|^2=\frac{1}{2}$$

$$e^{-2U}=2\Im\, \mathfrak{m}(\mathcal{H}_1\overline{\mathcal{H}}_2)$$

$$\partial_{[\underline{i}}\omega_{\underline{j}]}=\epsilon_{ijk}\Re(\mathcal{H}_1\partial_{\underline{k}}\overline{\mathcal{H}}_2-\overline{\mathcal{H}}_2\partial_{\underline{k}}\mathcal{H}_1)$$

$$ds^2=e^{2U}\big(dt^2+\omega_{\underline{i}}dx^{\underline{i}}\big)^2-e^{-2U}d\vec{x}^2$$

$$\lambda=\frac{\mathcal{H}_1}{\mathcal{H}_2}$$

$$A_t^{(n)}=2e^{2U}\Re(k^{(n)}\mathcal{H}_2)$$

$$\tilde{A}_t^{(n)}=-2e^{2U}\Re(k^{(n)}\mathcal{H}_1)$$

$$\mathcal{H}_1=i\mathcal{H}_2=\frac{1}{\sqrt{2}}V^{-1}$$

$$\Upsilon=-2\frac{\sum_n\overline{\Gamma^{(n)}}^2}{M}$$

$$M^2+|\Upsilon|^2-4\sum_n|\Gamma^{(n)}|^2=0$$

$$\frac{1}{2}|Z_{1,2}|^2=\sum_n|\Gamma^{(n)}|^2\pm\left[\left(\sum_n|\Gamma^{(n)}|^2\right)^2-\left|\sum_n\Gamma^{(n)2}\right|^2\right]^{\frac{1}{2}}$$

$$A=4\pi(|M|^2-|\Upsilon|^2)=4\pi||Z_1|^2-|Z_2|^2\mid$$

$$A=8\pi\sqrt{\det\left[\begin{pmatrix}\vec{p}^t\\ \vec{q}^t\end{pmatrix}(\vec{p}\vec{q})\right]}$$

$$\begin{pmatrix}\vec{p} \\ \vec{q}\end{pmatrix}'=R\otimes S\begin{pmatrix}\vec{p} \\ \vec{q}\end{pmatrix}$$

$$\left\{\hat{\hat{Q}}^{\hat{\hat{\alpha}}},\hat{\hat{Q}}^{\hat{\hat{\beta}}}\right\}~=~i\left(\hat{\Gamma}^{\hat{a}}\hat{\hat{C}}^{-1}\right)^{\hat{\hat{\alpha}}\hat{\hat{\beta}}}\hat{\hat{P}}_{\hat{\hat{a}}}+\frac{1}{2}\left(\hat{\Gamma}^{\hat{a}_1\hat{a}_2}\hat{\hat{C}}^{-1}\right)^{\hat{\hat{\alpha}}\hat{\hat{\beta}}}\hat{\hat{\mathcal{Z}}}^{(2)}_{\hat{\hat{a}}_1\hat{\hat{a}}_2}+\frac{i}{5!}\left(\hat{\Gamma}^{\hat{a}_1\cdots\hat{a}_5}\hat{\hat{C}}^{-1}\right)^{\hat{\hat{\alpha}}\hat{\hat{\beta}}}\hat{\hat{\mathcal{Z}}}^{(5)}_{\hat{\hat{a}}_1\cdots\hat{\hat{a}}_5}$$

$$+\frac{1}{6!}\Big(\hat{\Gamma}^{\hat{a}_1\cdots\hat{a}_6}\hat{\hat{C}}^{-1}\Big)^{\hat{\hat{\alpha}}\hat{\hat{\beta}}}\mathcal{Z}^{(6)}_{\hat{\hat{a}}_1\cdots\hat{\hat{a}}_6}+\frac{i}{9!}\Big(\hat{\Gamma}^{\hat{a}_1\cdots\hat{a}_9}\hat{\hat{C}}^{-1}\Big)^{\hat{\hat{\alpha}}\hat{\hat{\beta}}}\mathcal{Z}^{(9)}_{a_1\cdots a_9}$$

$$+\frac{1}{10!}\Big(\hat{\Gamma}^{\hat{a}_1\cdots\hat{a}_{10}}\hat{\hat{C}}^{-1}\Big)^{\hat{\hat{\alpha}}\hat{\hat{\beta}}}\hat{\hat{\mathcal{Z}}}^{(10)}_{\hat{\hat{a}}_1,\cdots,\hat{\hat{a}}_{10}}.$$

$$\{\hat{\hat{e}}_{\hat{\hat{\mu}}}^{\hat{\hat{a}}},\hat{\hat{C}}_{\hat{\hat{\mu}}\hat{\hat{\rho}}\hat{\hat{L}}},\hat{\hat{\psi}}_{\hat{\hat{\mu}}}\}$$

$$\hat{S}=\frac{1}{16\pi G_N}\int\,d^{11}\hat{x}\sqrt{|\hat{g}|}\left[\hat{R}-\frac{1}{2\cdot 4!}\hat{G}^2-\frac{1}{(144)^2}\frac{1}{\sqrt{|\hat{g}}}}\hat{\hat{e}}^{\hat{G}}\hat{G}^{\hat{G}}\hat{\hat{C}}^{\hat{C}}\right]$$

$$\hat{\hat{G}}=4\partial\hat{\hat{C}}\\ \delta_{\hat{\hat{\chi}}}\hat{\hat{C}}=3\partial\hat{\hat{\chi}}$$



$$\left\{ \begin{array}{l} \delta_{\hat{\epsilon}} \hat{e}_{\hat{\mu}}^{\hat{a}} = -\frac{i}{2} \hat{\bar{\epsilon}} \hat{\Gamma}^{\hat{a}} \hat{\psi}_{\hat{\mu}}, \\ \delta_{\hat{\epsilon}} \hat{\psi}_{\hat{\mu}} = 2 \nabla_{\hat{\mu}} \hat{\epsilon} + \frac{i}{144} \left(\hat{\Gamma}^{\hat{\alpha} \hat{\beta} \hat{\gamma} \hat{\delta}}_{\hat{\mu}} - 8 \hat{\Gamma}^{\hat{\beta} \hat{\gamma} \hat{\delta}} \hat{\eta}_{\hat{\mu}}^{\hat{\alpha}} \right) \hat{\epsilon} \hat{G}_{\hat{\alpha} \hat{\beta} \hat{\gamma} \hat{\delta}}, \\ \delta_{\hat{\epsilon}} \hat{G}_{\hat{\mu} \hat{\nu} \hat{\rho}} = \frac{3}{2} \hat{\bar{\epsilon}} \hat{\Gamma}_{[\hat{\mu} \hat{\nu}} \hat{\psi}_{\hat{\rho}]} . \end{array} \right.$$

$$\begin{aligned} & \partial \left({}^* \hat{G} + \frac{35}{2} \hat{C} \hat{G} \right) \\ & {}^* \hat{G} = 7(\partial \hat{\hat{C}} - 10 \hat{C} \partial \hat{C}) \equiv \hat{\hat{G}} \\ & \delta_{\hat{\chi}} \hat{\hat{C}} = 6 \partial \hat{\chi} \\ & \delta_{\hat{\chi}} \hat{\hat{C}} = -30 \partial \hat{\chi} \hat{C} \\ & \hat{P}_{\hat{a}} = (\hat{P}_{\hat{a}}, \hat{Z}^{(0)}), \hat{Z}_{\hat{a} \hat{b}}^{(2)} = (\hat{Z}_{\hat{a} \hat{b}}^{(2)}, \hat{Z}_{\hat{a}}^{(1)}), \hat{Z}_{\hat{a}_1 \dots \hat{a}_5}^{(5)} = (\hat{Z}_{\hat{a}_1 \dots \hat{a}_5}^{(5)}, \hat{Z}_{\hat{a}_1 \dots \hat{a}_4}^{(4)}) \\ & \hat{Z}_{\hat{a}_1 \dots \hat{a}_6}^{(6)} = (\hat{Z}_{\hat{a}_1 \dots \hat{a}_6}^{(6)}, \cdot), \hat{Z}_{\hat{a}_1 \dots \hat{a}_9}^{(9)} = (\cdot, \hat{Z}_{\hat{a}_1 \dots \hat{a}_8}^{(8)}) \\ & \{ \hat{Q}^{\hat{\alpha}}, \hat{Q}^{\hat{\beta}} \} = i(\hat{\Gamma}^{\hat{a}} \hat{C}^{-1})^{\hat{\alpha} \hat{\beta}} \hat{P}_{\hat{a}} + \sum_{n=0,1,4,8} \frac{c_n}{n!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_n} \hat{\Gamma}_{11} \hat{C}^{-1}) \hat{\alpha} \hat{\beta} \hat{Z}_{\hat{a}_1 \dots \hat{a}_n}^{(n)} \\ & + \sum_{n=2,5,6} \frac{c_n}{n!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_n} \hat{\Gamma}_{11} \hat{C}^{-1}) \hat{\alpha} \hat{\beta} \hat{Z}_{\hat{a}_1 \dots \hat{a}_n}^{(n)} \\ & \{ \hat{g}_{\hat{\mu} \hat{\nu}}, \hat{B}_{\hat{\mu} \hat{\nu}}, \hat{\phi}, \hat{C}^{(3)}_{\hat{\mu} \hat{\nu} \hat{\rho}}, \hat{C}^{(1)}_{\hat{\mu}}, \}. \\ & \frac{c_5}{5!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_5} \hat{\Gamma}_{11} \hat{C}^{-1})^{\hat{\alpha} \hat{\beta}} \hat{Z}_{\hat{a}_1 \dots \hat{a}_5}^{(5)} + \frac{c_9}{9!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_9} \hat{C}^{-1}) \hat{\alpha} \hat{\beta} \hat{Z}_{\hat{a}_1 \dots \hat{a}_9}^{(9)} \\ & \hat{g}_{\hat{\mu} \hat{\nu}} = e^{-\frac{2}{3} \hat{\phi}} \hat{g}_{\hat{\mu} \hat{\nu}} - e^{\frac{4}{3} \hat{\phi}} \hat{C}^{(1)} \hat{\mu}^{(1)} \hat{\nu}, \quad \hat{C}_{\hat{\mu} \hat{\nu} \hat{\rho}} = \hat{C}^{(3)} \hat{\mu} \hat{\nu} \hat{\rho} \\ & \hat{g}_{\hat{\mu} \hat{z}} = -e^{\frac{4}{3} \hat{\phi}} \hat{C}^{(1)} \hat{\mu}, \quad \hat{C}_{\hat{\mu} \hat{\nu} \hat{z}} = \hat{B}_{\hat{\mu} \hat{\nu}}, \\ & \hat{g}_{\hat{z} \hat{z}} = -e^{\frac{4}{3} \hat{\phi}}. \\ & \left(\hat{e}_{\hat{\mu}}^{\hat{a} \hat{u}} = \begin{pmatrix} e^{-\frac{1}{3} \hat{\phi}} \hat{e}_{\hat{\mu}}^{\hat{a}} & e^{\frac{2}{3} \hat{\phi}} \hat{C}^{(1)} \hat{\mu} \\ 0 & e^{\frac{2}{3} \hat{\phi}} \end{pmatrix}, \right. \\ & \left. (\hat{e}_{\hat{a}}^{\hat{\mu}}) = \begin{pmatrix} e^{\frac{1}{3} \hat{\phi}} \hat{e}^{\hat{\mu}}_{\hat{a}} & -e^{\frac{1}{3} \hat{\phi}} \hat{C}^{(1)} \hat{a} \\ 0 & e^{-\frac{2}{3} \hat{\phi}} \end{pmatrix}. \right. \\ & \hat{S} = \frac{2\pi t_{\text{Planck}}^{(11)}}{16\pi G_N^{(1)}} \int d^{10} \hat{x} \sqrt{|\hat{g}|} \left\{ e^{-2\hat{\phi}} \left[\hat{R} - 4(\partial \hat{\phi})^2 + \frac{1}{2 \cdot 3!} \hat{H}^2 \right] \right. \\ & \left. - \left[\frac{1}{4} (\hat{G}^{(2)})^2 + \frac{1}{2 \cdot 4!} (\hat{G}^{(4)})^2 \right] - \frac{1}{144} \frac{1}{\sqrt{|\hat{g}|}} \hat{\epsilon} \partial \hat{C}^{(3)} \partial \hat{C}^{(3)} \hat{B} \right\} \\ & \hat{g}_{\hat{\mu} \hat{\nu}} \rightarrow e^{\frac{2}{3} \hat{\phi}_0} \hat{\eta}_{\hat{\mu} \hat{\nu}} \\ & \hat{g}_{\hat{\mu} \hat{\nu}} \rightarrow e^{\frac{2}{3} \hat{\phi}_0} \hat{g}_{\hat{\mu} \hat{\nu}}, \quad \hat{C}^{(1)} \hat{\mu} \rightarrow e^{\frac{1}{3} \hat{\phi}_0} \hat{C}^{(1)} \hat{\mu}, \\ & \hat{B}_{\hat{\mu} \hat{\nu}} \rightarrow e^{\frac{2}{3} \hat{\phi}_0} \hat{B}_{\hat{\mu} \hat{\nu}}, \quad \hat{C}^{(3)} \hat{\mu} \hat{\nu} \hat{\rho} \rightarrow e^{\hat{\phi}_0} \hat{C}^{(3)}_{\hat{\mu} \hat{\nu} \hat{\rho}} \end{aligned}$$

$$\hat{S} = \frac{g_A^2}{16\pi G_{NA}^{100}} \int d^{10}\hat{x} \sqrt{|\hat{g}|} \left\{ e^{-2\hat{\phi}} \left[\hat{R} - 4(\partial\hat{\phi})^2 + \frac{1}{2 \cdot 3!} \hat{H}^2 \right] - \left[\frac{1}{4} (\hat{G}^{(2)})^2 + \frac{1}{2 \cdot 4!} (\hat{G}^{(4)})^2 \right] - \frac{1}{144} \frac{1}{\sqrt{|\hat{g}|}} \hat{\epsilon} \partial \hat{C}^{(3)} \partial \hat{C}^{(3)} \hat{B} \right\}$$

$$\frac{2\pi\ell_{\rm Planck}^{(11)}e^{\frac{8}{3}\hat{\phi}_0}}{16\pi G_N^{(11)}}=\frac{g_A^2}{16\pi G_{NA}^{(10)}}$$

$$G_N^{(10)}=\frac{G_N^{(11)}}{2\pi\ell_{\rm Planck}^{(11)}g_A^{2/3}}$$

$$R_{11}=\frac{1}{2\pi}\lim_{r\rightarrow\infty}\int\sqrt{|\hat{g}_{\underline{z}\underline{z}}|}dz=\ell_{\rm Planck}^{(11)}e^{\frac{2}{3}\hat{\phi}_0}=\ell_{\rm Planck}^{(11)}g_A^{2/3}$$

$$G_{NA}^{(10)}=\frac{G_N^{(11)}}{2\pi R_{11}}=\frac{G_N^{(11)}}{V_{11}^8}$$

$$G_{NA}^{(10)}=\frac{\left(\ell_{\rm Planck}^{(11)}\right)^8}{32\pi^2g_A^{2/3}}$$

$$G_{NA}^{(10)}=8\pi^6g_A^2\ell_s^8$$

$$\ell_{\rm Planck}^{(11)}=2\pi\ell_sg_A^{1/3}$$

$$R_{11}=\ell_sg_A$$

$$\hat{G}^{(10-k)}=(-1)^{[k/2]}\star\hat{G}^{(k)}$$

$$\hat{G}=d\hat{C}-\hat{H}\wedge\hat{C}$$

$$d\hat{G}-\hat{H}\wedge\hat{G}=0,d^\star\hat{G}+\hat{H}\wedge^\star\hat{G}$$

$$\hat{H}^{(7)}=e^{-2\hat{\phi}\star}\hat{H}$$

$$dH=0,d\big(e^{2\phi\star}\hat{H}^{(7)}\big)$$

$$d\big(e^{-2\phi\star}\hat{H}\big)+\frac{1}{2}\star\hat{G}\wedge\hat{G}=0,d\hat{H}^{(7)}+\frac{1}{2}\star\hat{G}\wedge\hat{G}$$

$$\hat{H}^{(7)}=d\hat{B}^{(6)}-\frac{1}{2}\sum_{n=1}^{n=4}\star\hat{G}^{(2n+2)}\wedge\hat{C}^{(2n-1)}$$

$$\left\{\begin{array}{l} \hat{\hat{\epsilon}}=e^{-\frac{1}{6}(\hat{\phi}^{-}-\hat{\phi}_0)}\hat{\epsilon} \\ \hat{\hat{\psi}}_{\hat{a}}=e^{\frac{1}{6}(\hat{\phi}-\hat{\phi}_0)}\left(2\hat{\psi}_{\hat{a}}-\frac{1}{3}\hat{\Gamma}_{\hat{a}}\hat{\lambda}\right) \\ \hat{\hat{\psi}}_{\hat{z}}=\frac{2i}{3}e^{\frac{1}{6}(\hat{\phi}-\hat{\phi}_0)}\hat{\Gamma}_{11}\hat{\lambda} \end{array}\right.$$

$$\begin{aligned}
\delta_{\hat{\epsilon}} \hat{\epsilon}_{\hat{\mu}}^{\hat{a}} &= -i \bar{\hat{\epsilon}} \hat{\Gamma}^{\hat{a}} \hat{\psi}_{\hat{\mu}}, \\
\delta_{\hat{\epsilon}} \hat{\psi}_{\hat{\mu}} &= \left\{ \partial_{\hat{\mu}} - \frac{1}{4} \left(\psi_{\hat{\mu}} + \frac{1}{2} \Gamma_{11} \hat{H}_{\mu} \right) \right\} \hat{\epsilon} + \frac{i}{8} e^{\hat{\epsilon}} \sum_{n=1,2} \frac{1}{(2n)!} \hat{\epsilon}^{(2n)} \hat{\Gamma}_{\hat{\mu}} (-\hat{\Gamma}_{11})^n \hat{\epsilon}, \\
\delta_{\hat{\epsilon}} \hat{\beta}_{\hat{\mu}\hat{\nu}} &= -2i \bar{\hat{\epsilon}} \hat{\Gamma}_{[\hat{\mu}} \hat{\Gamma}_{11} \hat{\psi}_{\hat{\nu}]}, \\
\delta_{\hat{\epsilon}} \hat{C}^{(1)}_{\hat{\mu}} &= -e^{\hat{\phi}} \hat{\Gamma}_{11} \left(\hat{\psi}_{\hat{\mu}} - \frac{1}{2} \hat{\Gamma}_{\hat{\mu}} \hat{\lambda} \right), \\
\delta_{\hat{\epsilon}} \hat{C}^{(3)}_{\hat{\mu}\hat{\nu}\hat{\rho}} &= 3e^{\hat{\phi}} \bar{\hat{\epsilon}} \hat{\Gamma}_{\hat{\mu}\hat{\nu}} \left(\hat{\psi}_{\hat{\rho}} - \frac{1}{3!} \hat{\Gamma}_{\hat{\rho}} \hat{\lambda} \right) + 3\hat{C}^{(1)}_{[\hat{\mu}} \delta_{\hat{\epsilon}} \hat{B}_{\hat{\nu}\hat{\rho}]}, \\
\delta_{\hat{\epsilon}} \hat{\lambda} &= \left(\partial \partial \phi \hat{\phi} + \frac{1}{12} \hat{\Gamma}_{11} \hat{H} \right) \hat{\epsilon} + \frac{i}{4} e^{\hat{\phi}} \sum_{n=1,2} \frac{5-2n}{(2n)!} \hat{\zeta}^{(2n)} (-\hat{\Gamma}_{11})^n \hat{\epsilon}, \\
\delta_{\hat{\epsilon}} \hat{\phi} &= -\frac{i}{2} \hat{\epsilon} \hat{\epsilon} \\
\left\{ \hat{Q}^{i\hat{\alpha}}, \hat{Q}^{j\hat{\beta}} \right\} &= i \delta^{ij} (\hat{\Gamma}^{\hat{a}} \hat{C}^{-1})^{\hat{\alpha}\hat{\beta}} \hat{P}_{\hat{a}} + (\hat{\Gamma}^{\hat{a}} \hat{C}^{-1})^{\hat{\alpha}\hat{\beta}} \hat{Z}_{\hat{a}}^{(1)(ij)} + \frac{i}{3!} (\hat{\Gamma}^{\hat{a}_1 \hat{a}_2 \hat{a}_3} \hat{C}^{-1})^{\hat{\alpha}\hat{\beta}} \hat{Z}_{\hat{a}_1 \hat{a}_2 \hat{a}_3}^{(3)[ij]} \\
&\quad + \frac{i}{5!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_5} \hat{C}^{-1})^{\hat{\alpha}\hat{\beta}} \hat{Z}_{\hat{a}_1 \dots \hat{a}_5}^{(5)(ij)} + \frac{i}{7!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_7} \hat{C}^{-1})^{\hat{\alpha}\hat{\beta}} \hat{Z}_{\hat{a}_1 \dots \hat{a}_7}^{(7)[ij]} \\
&\quad + \frac{i}{9!} (\hat{\Gamma}^{\hat{a}_1 \dots \hat{a}_9} \hat{C}^{-1})^{\hat{\alpha}\hat{\beta}} \hat{Z}_{\hat{a}_1 \dots \hat{a}_9}^{(9)(ij)}. \\
&\quad \hat{Z}_{\hat{a}}^{(1)(ij)} = \hat{Z}_{\hat{a}}^{(1)0} \delta^{ij} \hat{Z}_{\hat{a}}^{(1)1} \sigma^1 + \hat{Z}_{\hat{a}}^{(1)3} \sigma^3 \\
&\quad \hat{Z}_{\hat{a}_1 \hat{a}_2 \hat{a}_3}^{(3)} = \hat{Z}_{\hat{a}_1 \hat{a}_2 \hat{a}_3}^{(i)} i \sigma^2 \\
&\quad \{ \hat{J}_{\hat{\mu}\hat{\nu}}, \hat{B}_{\hat{\mu}\hat{\nu}}, \hat{\phi} \} \\
&\quad \{ C^{(0)}, C^{(2)} \hat{\mu}\hat{\nu}, C^{(4)} \hat{\mu}\hat{\nu}\hat{\rho}\hat{\sigma} \} \\
&\quad \left\{ \begin{aligned} \hat{\mathcal{H}} &= 3\partial\mathcal{B} \\ \hat{G}^{(1)} &= \partial\hat{C}^{(0)} \\ \hat{G}^{(3)} &= 3(\partial\hat{C}^{(2)} - \partial\hat{B}\hat{C}^{(0)}) \\ \hat{G}^{(5)} &= 5(\partial\hat{C}^{(4)} - 6\partial\hat{B}\hat{C}^{(2)}) \\ \hat{G}^{(5)} &= +^* \hat{G}^{(5)} \\ (\hat{G}^{(5)})^2 &= (^* \hat{G}^{(5)})^2 = -(\hat{G}^{(5)})^2 \Rightarrow 0 \end{aligned} \right. \\
S_{\text{NSD}} &= \frac{g_B^2}{16\pi G_{NB}^{(10)}} \int d^{10} \hat{x} \sqrt{|\hat{J}|} \left\{ e^{-2\hat{\phi}} \left[\hat{R}(\hat{J}) - 4(\partial\hat{\phi})^2 + \frac{1}{2 \cdot 5!} \hat{\mathcal{H}}^2 \right] \right. \\
&\quad + \frac{1}{2} (\hat{G}^{(0)})^2 + \frac{1}{2 \cdot 3!} (\hat{G}^{(3)})^2 + \frac{1}{4 \cdot 3!} (\hat{G}^{(5)})^2 \\
&\quad \left. - \frac{1}{192} \frac{1}{\sqrt{|\hat{J}|}} \in \partial\hat{C}^{(4)} \partial\hat{C}^{(2)} \hat{B} \right\} \\
g_B &= e^{\hat{\phi}_0}
\end{aligned}$$

$$\begin{aligned}
\delta_{\hat{\varepsilon}} \hat{e}_{\hat{\mu}}^{\hat{a}} &= -i \bar{\hat{\varepsilon}} \hat{\Gamma}^{\hat{a}} \hat{\zeta}_{\hat{\mu}}, \\
\delta_{\hat{\varepsilon}} \hat{\zeta}_{\hat{\mu}} &= \nabla_{\hat{\mu}} \hat{\varepsilon} - \frac{1}{8} \hat{\mathcal{H}}_{\hat{\mu}} \sigma_3 \hat{\varepsilon} + \frac{1}{8} e^{\hat{\phi}} \sum_{n=1,2,3} \frac{1}{(2n-1)!} \hat{l}^{(2n-1)} \hat{\Gamma}_{\hat{\mu}} \mathcal{P}_n \hat{\varepsilon}, \\
\delta_{\hat{\varepsilon}} \hat{\mathcal{B}}_{\hat{\mu}\hat{\nu}} &= -2i \bar{\hat{\varepsilon}} \sigma^3 \hat{\Gamma}_{[\hat{\mu}} \hat{\zeta}_{\hat{\nu}]}, \\
\delta_{\hat{\varepsilon}} \hat{\mathcal{C}}^{(2n-2)}_{\hat{\mu}_1 \dots \hat{\mu}_{2n-2}} &= i(2n-2) e^{-\hat{\phi}} \bar{\hat{\varepsilon}} \mathcal{P}_n \hat{\Gamma}_{[\hat{\mu}_1 \dots \hat{\mu}_{2n-3}} \left(\hat{\zeta}_{\hat{\mu}_{2n-2}] } - \frac{1}{2(2n-2)} \hat{\Gamma}_{\hat{\mu}_{2n-2}] \hat{\chi}} \right) \\
&\quad + \frac{1}{2} (2n-2)(2n-3) \hat{\mathcal{C}}^{(2n-4)}_{[\hat{\mu}_1 \dots \hat{\mu}_{2n-4}} \delta_{\hat{\varepsilon}} \hat{\mathcal{B}}_{\hat{\mu}_{2n-3} \hat{\mu}_{2n-4}}, \\
\delta_{\hat{\varepsilon}} \hat{\chi} &= \left(\partial \hat{\phi} - \frac{1}{12} \hat{\mathcal{H}} \sigma^3 \right) \hat{\varepsilon} + \frac{1}{2} e^{\hat{\phi}} \sum_{n=1,2,3} \frac{(n-3)}{(2n-1)!} \hat{\zeta}^{(2n-1)} \mathcal{P}_n \varepsilon, \\
\delta_{\hat{\varepsilon}} \hat{\phi} &= -\frac{i}{2} \bar{\hat{\varepsilon}} \hat{\chi}
\end{aligned}$$

$$\begin{aligned}
\mathcal{P}_n &= \begin{cases} \sigma^1, & n \\ i\sigma^2, & n \end{cases} \\
\hat{J}_{E\mu\nu} &= e^{-\varphi/2} J_{\mu\nu} \\
\begin{pmatrix} \hat{\hat{B}} \\ \hat{B} \end{pmatrix} &= \begin{pmatrix} \hat{\mathcal{C}}^{(2)} \\ \hat{B} \end{pmatrix} \\
\hat{D} &= \hat{\mathcal{C}}^{(4)} - 3\hat{B}\hat{\mathcal{C}}^{(2)} \\
\begin{cases} \hat{\mathcal{H}} &= 3\partial\hat{\mathcal{B}} \\ \hat{F} &= \hat{G}^{(5)} = +^*\hat{F} \\ &= 5\left(\partial\hat{D} - \hat{\hat{B}}^T\eta\hat{\mathcal{H}}\right) \end{cases} \\
\eta &= i\sigma^2 = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} = -\eta^{-1} = -\eta^T \\
\Lambda\eta\Lambda^T &= \eta, \Rightarrow \eta\Lambda\eta^T = (\Lambda^{-1})^T, \Lambda \in SL(2, \mathbb{R}) \\
\hat{\mathcal{M}} &= e^{\hat{\phi}} \begin{pmatrix} |\hat{t}|^2 & \hat{\mathcal{C}}^{(0)} \\ \hat{\mathcal{C}}^{(0)} & 1 \end{pmatrix}, \hat{\mathcal{M}}^{-1} = e^{\hat{\phi}} \begin{pmatrix} 1 & -\hat{\mathcal{C}}^{(0)} \\ -\hat{\mathcal{C}}^{(0)} & |\hat{t}|^2 \end{pmatrix} \\
\hat{t} &= \hat{\mathcal{C}}^{(0)} + ie^{-\hat{\phi}} \\
\hat{\mathcal{M}}' &= \Lambda\hat{\mathcal{M}}\Lambda^T, \\
\hat{\mathcal{B}}' &= \Lambda\hat{\mathcal{B}} \\
\hat{t}' &= \frac{a\hat{t} + b}{c\hat{t} + d}
\end{aligned}$$

$$\begin{aligned}
\hat{S}_{\text{NSD}} &= \frac{g_B^2}{16\pi G_N^{100}} \int d^{10} \hat{x} \sqrt{|\hat{J}_E|} \left\{ \hat{R}(\hat{J}_E) + \frac{1}{4} \text{Tr}(\partial \hat{\mathcal{M}} \hat{\mathcal{M}}^{-1})^2 \right. \\
&\quad \left. + \frac{1}{2 \cdot 3!} \hat{\mathcal{H}}^T \hat{\mathcal{M}}^{-1} \hat{\mathcal{H}} + \frac{1}{4 \cdot 5!} \hat{F}^2 - \frac{1}{2^7 \cdot 3^3} \frac{1}{\sqrt{|\hat{J}_E|}} \in \hat{D} \hat{\mathcal{H}}^T \eta \hat{\mathcal{H}} \right\} \\
g_B' &= 1/g_B \\
\hat{j}' &= |c\hat{\lambda} + d|\hat{j} \\
\hat{j}' &= e^{-\hat{\phi}} \hat{j} \\
R' &= R/g_B
\end{aligned}$$

$$\begin{aligned}
\{Q^{i\alpha}, Q^{j\beta}\} = & i(\Gamma^a \mathcal{C}^{-1})^{\alpha\beta} \left(\delta^{ij} P_a + \sigma^{1ij} Z_a^{(1)1} + \sigma^{3ij} Z_a^{(1)3} \right) \\
& + (\mathcal{C}^{-1})^{\alpha\beta} \left(\delta^{ij} Z^{(0)0} + \sigma^{1ij} Z^{(0)1} + \sigma^{3ij} Z^{(0)3} \right) \\
& + \frac{i}{2!} (\Gamma^{a_1 a_2} \mathcal{C}^{-1})^{\alpha\beta} \sigma^{2ij} Z_{a_1 a_2}^{(2)} + \frac{1}{3!} (\Gamma^{a_1 a_2 a_3} \mathcal{C}^{-1})^{\alpha\beta} \sigma^{2ij} Z_{a_1 a_2 a_3}^{(3)} \\
& + \frac{1}{4!} (\Gamma^{a_1 \cdot a_4} \mathcal{C}^{-1})^{\alpha\beta} \left(\sigma^{1ij} Z_{a_1 \cdot a_4}^{(4)1} + \sigma^{3ij} Z_{a_1 \cdot a_4}^{(4)3} \right) \\
& + \frac{i}{5!} (\Gamma^{a_1 \cdot a_5} \mathcal{C}^{-1})^{\alpha\beta} \left(\delta^{ij} Z_{a_1 \cdot a_5}^{(5)0} + \sigma^{1ij} Z_{a_1 \cdot a_5}^{(5)1} + \sigma^{3ij} Z_{a_1 \cdot a_5}^{(5)3} \right) \\
& + \frac{i}{6!} (\Gamma^{a_1 \cdots a_6} \mathcal{C}^{-1})^{\alpha\beta} \sigma^{2ij} \left(Z_{a_1 \cdots a_6}^{(6)} + Z_{a_1 \cdots a_6}^{(6)!} \right) \\
& + \frac{1}{7!} (\Gamma^{a_1 \cdots a_7} \mathcal{C}^{-1})^{\alpha\beta} \sigma^{2ij} \left(Z_{a_1 \cdots a_7}^{(7)} + Z_{a_1 \cdots a_7}^{(7)!} \right) \\
& + \frac{1}{8!} (\Gamma^{a_1 \cdot a_8} \mathcal{C}^{-1})^{\alpha\beta} \left(\sigma^{1ij} Z_{a_1 \cdot a_8}^{(8)1} + \sigma^{3ij} Z_{a_1 \cdot a_8}^{(8)3} \right) \\
\hat{g}_{\mu\nu} = & g_{\mu\nu} - k^2 A^{(1)}{}_{\mu} A^{(1)}{}_{\nu}, g_{\mu\nu} = \hat{g}_{\mu\nu} - \hat{g}_{\mu\bar{x}} \hat{g}_{\nu\bar{x}} / \hat{g}_{\bar{x}\bar{x}}, \\
\hat{B}_{\mu\nu} = & B_{\mu\nu} + A^{(1)}{}_{[\mu} A^{(2)}{}_{\nu]}, B_{\mu\nu} = \hat{B}_{\mu\nu} + \hat{g}_{[\mu|\bar{x}} \hat{B}_{\nu] \bar{x}} / \hat{g}_{\bar{x}\bar{x}}, \\
\hat{\phi} = & \phi + \frac{1}{2} \log k, \phi = \hat{\phi} - \frac{1}{4} \log |\hat{g}_{\bar{x}\bar{x}}|, \\
\hat{g}_{\mu\bar{x}} = & -k^2 A^{(1)}{}_{\mu}, A^{(1)}{}_{\mu} = \hat{g}_{\mu\bar{x}} / \hat{g}_{\bar{x}\bar{x}}, \\
\hat{B}_{\mu\bar{x}} = & -A^{(2)}{}_{\mu}, A^{(2)}{}_{\mu} = -\hat{B}_{\mu\bar{x}}, \\
\hat{g}_{\bar{x}\bar{x}} = & -k^2, k = |\hat{g}_{\bar{x}\bar{x}}|^{1/2} \\
\hat{C}^{(2n-1)}{}_{\mu_1 \cdots \mu_{2n-1}} = & C^{(2n-1)}{}_{\mu_1 \cdots \mu_{2n-1}} + (2n-1) A^{(1)}{}_{[\mu_1} C^{(2n-2)}{}_{\mu_2 \cdots \mu_{2n-1}]}, \\
\hat{C}^{(2n+1)}{}_{\mu_1 \cdots \mu_{2n}\bar{x}} = & C^{(2n)}{}_{\mu_1 \cdots \mu_{2n}}, \\
C^{(2n-1)}{}_{\mu_1 \cdots \mu_{2n-1}} = & \hat{C}^{(2n-1)}{}_{\mu_1 \cdots \mu_{2n-1}} - (2n-1) \hat{g}_{[\mu_1|\bar{x}} \hat{C}^{(2n-1)}{}_{\mu_2 \cdots \mu_{2n-1}]\bar{x}} / \hat{g}_{\bar{x}\bar{x}}, \\
C^{(2n)}{}_{\mu_1 \cdots \mu_{2n}} = & \hat{C}^{(2n+1)}{}_{\mu_1 \cdots \mu_{2n}\bar{x}} \\
\hat{j}_{\mu\nu} = & g_{\mu\nu} - k^{-2} A^{(2)}{}_{\mu} A^{(2)}{}_{\nu}, g_{\mu\nu} = \hat{j}_{\mu\nu} - \hat{j}_{\mu\bar{y}} \hat{j}_{\nu\bar{y}} / \hat{j}_{\bar{y}\bar{y}}, \\
\hat{B}_{\mu\nu} = & B_{\mu\nu} + A^{(1)}{}_{[\mu} A^{(2)}{}_{\nu]}, B_{\mu\nu} = \hat{B}_{\mu\nu} + \hat{j}_{[\mu|\bar{y}} \hat{B}_{\nu]\bar{y}} / \hat{j}_{\bar{y}\bar{y}}, \\
\hat{\phi} = & \phi - \frac{1}{2} \log k, \phi = \hat{\phi} - \frac{1}{4} \log |\hat{j}_{\bar{y}\bar{y}}|, \\
\hat{j}_{\mu\bar{y}} = & -k^{-2} A^{(2)}{}_{\mu}, A^{(1)}{}_{\mu} = \hat{B}_{\mu\bar{y}}, \\
\hat{B}_{\mu\bar{y}} = & A^{(1)}{}_{\mu}, A^{(2)}{}_{\mu} = \hat{j}_{\mu\bar{y}} / \hat{j}_{\bar{y}\bar{y}}, \\
\hat{j}_{\bar{y}\bar{y}} = & -k^{-2}, k = |\hat{j}_{\bar{y}\bar{y}}|^{-1/2} \\
\hat{C}^{(2n)}{}_{\mu_1 \cdots \mu_{2n}} = & C^{(2n)}{}_{\mu_1 \cdots \mu_{2n}} - (2n) A^{(2)}{}_{[\mu_1} C^{(2n-1)}{}_{\mu_2 \cdots \mu_{2n}]}, \\
\hat{C}^{(2n)}{}_{\mu_1 \cdots \mu_{2n-1}\bar{y}} = & -C^{(2n-1)}{}_{\mu_1 \cdots \mu_{2n-1}}, \\
C^{(2n)}{}_{\mu_1 \cdots \mu_{2n}} = & \hat{C}^{(2n)}{}_{\mu_1 \cdots \mu_{2n}} + (2n) \hat{j}_{[\mu_1|\bar{y}} \hat{C}^{(2n)}{}_{\mu_2 \cdots \mu_{2n}]\bar{y}} / \hat{j}_{\bar{y}\bar{y}}, \\
C^{(2n-1)}{}_{\mu_1 \cdots \mu_{2n-1}} = & -\hat{C}^{(2n)}{}_{\mu_1 \cdots \mu_{2n-1}\bar{y}}
\end{aligned}$$

$$\begin{aligned}
\hat{j}_{\mu\nu} &= \hat{g}_{\mu\nu} - (\hat{g}_{\mu\bar{x}}\hat{g}_{\nu\bar{x}} - \hat{B}_{\mu\bar{x}}\hat{B}_{\nu\bar{x}})/\hat{g}_{\bar{x}\bar{x}}, \hat{j}_{\mu\bar{y}} = \hat{B}_{\mu\bar{x}}/\hat{g}_{\bar{x}\bar{x}}, \\
\hat{B}_{\mu\nu} &= \hat{B}_{\mu\nu} + 2\hat{g}_{[\mu|\bar{x}}\hat{B}_{\nu]\bar{x}}/\hat{g}_{\bar{x}\bar{x}}, \hat{B}_{\mu\bar{y}} = \hat{g}_{\mu\bar{x}}/\hat{g}_{\bar{x}\bar{x}}, \\
\hat{\phi} &= \hat{\phi} - \frac{1}{2}\log|\hat{g}_{\bar{x}\bar{x}}|, \hat{j}_{\bar{y}\bar{y}} = 1/\hat{g}_{\bar{x}\bar{x}}, \\
\hat{C}_{\mu_1\ldots\mu_{2n}}^{(2n)} &= \hat{C}_{\mu_1\ldots\mu_{2n}\bar{x}}^{(2n+1)} + 2n\hat{B}_{[\mu_1|\bar{x}}\hat{C}_{\mu_2\ldots\mu_{2n}]}^{(2n-1)} \\
&\quad - 2n(2n-1)\hat{B}_{[\mu_1|\bar{x}}\hat{g}_{\mu_2|\bar{x}}\hat{C}_{\mu_3\ldots\mu_{2n}]\bar{x}}^{(2n-1)}/\hat{g}_{\bar{x}\bar{x}}, \\
\hat{C}_{\mu_1\ldots\mu_{2n-1}\bar{y}}^{(2n)} &= -\hat{C}_{\mu_1\ldots\mu_{2n-1}}^{(2n-1)} \\
&\quad + (2n-1)\hat{g}_{[\mu_1|\bar{x}}\hat{C}_{\mu_2\ldots\mu_{2n-1}]\bar{x}}^{(2n-1)}/\hat{g}_{\bar{x}\bar{x}} \\
\hat{g}_{\mu\nu} &= \hat{j}_{\mu\nu} - (\hat{j}_{\mu\bar{y}}\hat{j}_{\nu\bar{y}} - \hat{B}_{\mu\bar{y}}\hat{B}_{\nu\bar{y}})/\hat{j}_{\bar{y}\bar{y}}, \hat{g}_{\mu\bar{x}} = \hat{B}_{\mu\bar{y}}/\hat{j}_{\bar{y}\bar{y}}, \\
\hat{B}_{\mu\nu} &= \hat{B}_{\mu\nu} + 2\hat{j}_{[\mu|\bar{y}}\hat{B}_{\nu]\bar{y}}/\hat{j}_{\bar{y}\bar{y}}, \hat{B}_{\mu\bar{x}} = \hat{j}_{\mu\bar{y}}/\hat{j}_{\bar{y}\bar{y}}, \\
\hat{\phi} &= \hat{\phi} - \frac{1}{2}\log|\hat{j}_{\bar{y}\bar{y}}|, \hat{g}_{\bar{x}\bar{x}} = 1/\hat{j}_{\bar{y}\bar{y}}, \\
\hat{C}_{\mu_1\ldots\mu_{2n+1}}^{(2n+1)} &= -\hat{C}_{\mu_1\ldots\mu_{2n+1}\bar{y}}^{(2n+2)} + (2n+1)\hat{B}_{[\mu_1|\bar{y}}\hat{C}_{\mu_2\ldots\mu_{2n+1}]}^{(2n)} \\
&\quad - 2n(2n+1)\hat{B}_{[\mu_1|\bar{y}}\hat{j}_{\mu_2|\bar{y}}\hat{C}_{\mu_3\ldots\mu_{2n+1}]\bar{y}}^{(2n)}/\hat{j}_{\bar{y}\bar{y}}, \\
\hat{C}_{\mu_1\ldots\mu_{2n}\bar{x}}^{(2n+1)} &= \hat{C}_{\mu_1\ldots\mu_{2n}}^{(2n)} \\
&\quad + 2n\hat{j}_{[\mu_1|\bar{y}}\hat{C}_{\mu_2\ldots\mu_{2n}]\bar{y}}^{(2n)}/\hat{j}_{\bar{y}\bar{y}} \\
&\quad \hat{j}_{\bar{y}\bar{y}} = 1/\hat{g}_{\bar{x}\bar{x}}, \hat{g}_{\bar{x}\bar{x}} = 1/\hat{j}_{\bar{y}\bar{y}} \\
&\quad \hat{g}_{\bar{x}\bar{x}} \rightarrow (R_A/\ell_s)^2, \hat{j}_{\bar{y}\bar{y}} \rightarrow (R_B/\ell_s)^2 \\
&\quad R_{A,B} = \ell_s^2/R_{B,A} \\
&\quad g_{A,B} = g_{B,A}/R_{B,A} \\
&\quad g \ll 1, \ell_s \ll 1 \\
&\quad \ell_s/R_c \ll 1 \\
S_{NG}^{(p)}[X^\mu(\xi)] &= -T_{(p)} \int d^{p+1}\xi \sqrt{|g_{ij}|} \\
&\quad g_{ij} = g_{\mu\nu}(X)\partial_i X^\mu \partial_j X^\nu \\
S_p^{(p)}[X^\mu, \gamma_{ij}] &= -\frac{T_{(p)}}{2} \int d^{p+1}\xi \sqrt{|\gamma|} [\gamma^{ij} \partial_i X^\mu \partial_j X^\nu g_{\mu\nu} + (1-p)] \\
&\quad \gamma_{ij} = g_{ij} \\
&\quad \gamma_{ij} = \Omega(\xi) g_{ij} \\
S_{NG}^{(p)}[X^\mu(\xi)] &= -\frac{T_{(p)}}{K_0^\alpha} \int d^{p+1}\xi K(X)^\alpha \sqrt{|g_{ij}|} - \frac{\mu}{K_0^\alpha(p+1)!} \int d^{p+1}\xi \epsilon^{i_1\ldots i_{p+1}} A_{(p+1)i_1\ldots i_{p+1}} \\
&\quad A_{(p+1)i_1\ldots i_{p+1}} = A_{(p+1)\mu_1\ldots\mu_{p+1}}(X) \partial_{i_1} X^{\mu_1} \ldots \partial_{i_{p+1}} X^{\mu_{p+1}} \\
&\quad \delta A_{(p+1)} = (p+1) \partial \Lambda_{(p)} \\
S_{NG}^{(p)}[X^\mu(\xi)] &= -\frac{T_{(p)}}{K_0^\alpha} \int d^{p+1}\xi K^\alpha(X) \sqrt{|g_{ij} + F_{ij}|} + \ldots, F_{ij} = 2\partial_{[i} V_{j]} \\
S_{NG}^{(p)}[X^\mu(\xi)] &= -\frac{T_{(p)}}{K_0^\alpha} \int d^{p+1}\xi K^\alpha(X) \sqrt{|g_{ij}|} \left\{ 1 - \frac{1}{2} \mathcal{F}^2 + \ldots \right\} \\
S &= \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{|g|} \left[R + 2(\partial \log K)^2 + \frac{(-1)^{p+1}}{2 \cdot (p+2)!} K^\beta F_{(p+2)}^2 \right] \\
&\quad F_{(p+2)} = (p+2) \partial A_{(p+1)}
\end{aligned}$$

$$\begin{aligned}
S &= -T \int d^2\xi \sqrt{|\hat{g}_{ij}|} - \frac{T}{2} \int d^2\xi \epsilon^{ij} \hat{B}_{ij} \\
S^{(p)} &= -T_{(p)} e^{2\phi_0} \int d^{p+1}\xi e^{-2\phi} \sqrt{|g_{ij}|} + \dots \\
S &= -T_{S5} e^{2\hat{\phi}_0} \int d^6\xi e^{-2\hat{\phi}} \sqrt{|\hat{g}_{ij}|} - \frac{T_{S5} e^{2\hat{\phi}_0}}{6!} \int d^6\xi \epsilon^{i_1 \dots i_6} \hat{B}_{i_1 \dots i_6}^{(6)} \\
S &= -T_{Dp} e^{\hat{\phi}_0} \int d^{p+1}\xi e^{-\hat{\phi}} \sqrt{|\hat{g}_{ij} + 2\pi\alpha' \mathcal{F}_{ij}|} - \frac{T_{SDp} e^{\hat{\phi}_0}}{6!} \int d^6\xi \epsilon^{i_1 \dots i_{p+1}} \hat{C}_{i_1 \dots i_{p+1}}^{(p+1)} \\
\mathcal{F}_{ij} &= F_{ij} + \frac{1}{2\pi\alpha'} \hat{B}_{ij}, F_{ij} = 2\partial_{[i} V_{j]} \\
R_{A,B} &= \ell_s^2 / R_{B,A} \\
g_{A,B} &= g_B \ell_s / R_{B,A} \\
g' &= 1/g \\
R'_i &= R_i / \sqrt{g} \\
M' &= g^{1/2} M \\
M_{F1} &= \frac{R_9}{\ell_s^2} \\
M_{D1} = M'_{F1} &= g^{1/2} M_{F1} = g^{1/2} \frac{R_9}{\ell_s^2} = \frac{R'_9}{g' \ell_s^2} \\
M_{D0} = M'_{D1} &= \frac{R_9}{g \ell_s^2} = \frac{\ell_s^2 / R_9}{g' \ell_s / R_9 \ell_s^2} = \frac{1}{g' \ell_s} \\
M_{D0} = M'_{D1} &= \frac{R_9}{g \ell_s^2} = \frac{R_9}{g' \ell_s / R_9 \ell_s^2} = \frac{R_8 R_9}{g' \ell_s^3} \\
M_{Dp} &= \frac{R_{10-p} \dots R_9}{g \ell_s^{p+1}} \\
M_{S5} = g^{1/2} M'_{D5} &= g^{1/2} \frac{R_5 \dots R_9}{g \ell_s^6} = g'^{-1/2} \frac{R'_5 / g'^{1/2} \dots R'_9 / g^{1/2}}{g'^{-1} \ell_s^6} = \frac{R_5 \dots R_9}{g^2 \ell_s^6}
\end{aligned}$$

Objeto Supermasivo	Masa	Masa	Objeto Masivo
F1m	R_9^{-1}		
Do	$g_A^{-1} \ell_s^{-1}$	R_{10}^{-1}	WM(+, - ¹⁰)
F1w	$R_9 \ell_s^{-2}$	$R_{10} R_9 \left(\ell_{\text{Planck}}^{(11)} \right)^{-3}$	M2(+, - ⁸ , + ²)
D2	$R_9 R_8 g_A^{-1} \ell_s^{-3}$	$R_9 R_8 \left(\ell_{\text{Planck}}^{(11)} \right)^{-3}$	M2(+, - ⁷ , + ² , -)
D4	$R_9 \dots R_6 g_A^{-1} \ell_s^{-5}$	$R_{10} R_9 \dots R_5 \left(\ell_{\text{Planck}}^{(11)} \right)^{-6}$	M5(+, - ⁵ , + ⁵)
S5A	$R_9 \dots R_5 g_A^{-2} \ell_s^{-6}$	$R_9 \dots R_5 \left(\ell_{\text{Planck}}^{(11)} \right)^{-6}$	M5(+, - ⁴ , + ⁵ , -)
D6	$R_9 \dots R_4 g_A^{-1} \ell_s^{-7}$	$R_{10}^2 R_9 \dots R_4 \left(\ell_{\text{Planck}}^{(11)} \right)^{-9}$	KK7M(+, - ³ , + ⁶ , - [*])

KK6A	$R_9^2 R_8 \dots R_4 g_A^{-2} \ell_s^{-8}$	$R_{10} R_9^2 \dots R_4 \left(\ell_{\text{Planck}}^{(11)} \right)^{-9}$	KK7M(+, ³ , + ⁵ , +*, +)
D8	$R_9 \dots R_2 g_A^{-1} \ell_s^{-9}$	$R_{10}^3 R_9 \dots R_4 \left(\ell_{\text{Planck}}^{(11)} \right)^{-12}$	KK9M(+, -, + ⁸ , +*)
KK8A	$R_9^3 R_8 \dots R_2 g_A^{-3} \ell_s^{-11}$	$R_{10} R_9^3 R_8 \dots R_2 \left(\ell_{\text{Planck}}^{(11)} \right)^{-12}$	KK9M(+, -, + ⁷ , +*, +)
KK9A	$R_9^3 R_8 \dots R_1 g_A^{-4} \ell_s^{-12}$	$R_{10} R_9^3 R_8 \dots R_1 \left(\ell_{\text{Planck}}^{(11)} \right)^{-12}$	KK9M(+, + ⁸ , +*, -)

Objeto Supermasivo	Masa	Objeto Masivo	Masa
F1m	R_9^{-1}	KK6A	$R_9^2 R_8 \dots R_4 g_B^{-2} \ell_s^{-8}$
F1w	$R_9 \ell_s^{-2}$	D7	$R_9 \dots R_3 g_B^{-1} \ell_s^{-8}$
D1	$R_9 g_B^{-1} \ell_s^{-2}$	Q7	$R_9 \dots R_3 g_B^{-3} \ell_s^{-8}$
D3	$R_9 \dots R_7 g_B^{-1} \ell_s^{-4}$	D9	$R_9 \dots R_1 g_B^{-1} \ell_s^{-10}$
D5	$R_9 \dots R_5 g_B^{-1} \ell_s^{-6}$	Q9	$R_9 \dots R_1 g_B^{-4} \ell_s^{-10}$
S5B	$R_9 \dots R_5 g_B^{-2} \ell_s^{-6}$		

$$S = -T_{M2} \int d^3 \xi \sqrt{|\hat{g}_{ij}|} - \frac{T_{M2}}{3!} \int d^3 \xi \epsilon^{i_1 \dots i_3} \hat{C}_{i_1 \dots i_3}$$

$$\ell_s = \ell_{\text{Planck}}^{(11)} / R_{10}^{1/2}$$

$$g_A = R_{10}^{3/2} / \ell_{\text{Planck}}^{(11)}$$

$$M_{M2} = M_{F1A} = \frac{R_9}{\ell_s^2} = \frac{R_9 R_{10}}{\left(\ell_{\text{Planck}}^{(11)} \right)^3}$$

$$M_{M2} = \frac{R_8 R_9}{\left(\ell_{\text{Planck}}^{(11)} \right)^3} = \frac{R_8 R_9}{g_A \ell_s^3}$$

$$M_{M5} = M_{D4} = \frac{R_6 \dots R_9}{g_A \ell_s^5} = \frac{R_6 \dots R_{10}}{\left(\ell_{\text{Planck}}^{(11)} \right)^6}$$

$$M_{M5} = \frac{R_5 \dots R_9}{\left(\ell_{\text{Planck}}^{(11)} \right)^6} = \frac{R_5 \dots R_9}{g_A^2 \ell_s^6}$$

$$M_{D0} = \frac{1}{g_A \ell_s} = \frac{1}{R_{10}}$$

Objeto Súper	Masa
WM	o
M2	$R_{10} R_9 \left(\ell_{\text{Planck}}^{(11)} \right)^{-3}$



$$\begin{array}{ll}
\text{M5} & R_{10} \dots R_6 \left(\ell_{\text{Planck}}^{(11)} \right)^{-6} \\
\text{KK7M} & R_{10}^2 R_9 \dots R_4 \left(\ell_{\text{Planck}}^{(11)} \right)^{-9} \\
\text{KK9M} & R_{10}^3 R_9 \dots R_4 \left(\ell_{\text{Planck}}^{(11)} \right)^{-12}
\end{array}$$

$$S = \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{|g|} \left[R + 2(\partial\varphi)^2 + \frac{(-1)^{p+1}}{2 \cdot (p+2)!} e^{-2a\varphi} F_{(p+2)}^2 \right]$$

$$F_{(p+2)} = dA_{(p+1)}, F_{(p+2)\mu_1 \dots \mu_{p+2}} = (p+2) \partial_{[\mu_1} A_{(p+1)\mu_2 \dots \mu_{p+2}]}$$

$$ds^2 = f[W dt^2 - d\vec{y}_p^2] - g^{-1}[W^{-1} d\rho^2 - \rho^2 d\Omega_{(\tilde{p}+2)}^2]$$

$$A_{t\underline{t}y^1 \dots \underline{y}^p} = \alpha(H^{-1} - 1)$$

$$\tilde{p} \equiv d - p - 4$$

$$H = 1 + \frac{h}{\rho^{\tilde{p}+1}}, W = 1 + \frac{\omega}{\rho^{\tilde{p}+1}}$$

$$ds^2 = H^{\frac{2x-2}{\tilde{p}+1}} [W dt^2 - d\vec{y}_p^2] - H^{\frac{-(2x-2)}{\tilde{p}+1}} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(\tilde{p}+2)}^2]$$

$$e^{-2a\varphi} = e^{-2a\varphi_0} H^{2x}, A_{t\underline{t}y^1 \dots \underline{y}^p} = e^{a\varphi_0} \alpha(H^{-1} - 1),$$

$$H = 1 + \frac{h}{\rho^{\tilde{p}+1}}, W = 1 + \frac{\omega}{\rho^{\tilde{p}+1}},$$

$$\omega = h \left[1 - \frac{a^2}{4x} \alpha^2 \right],$$

$$x = \frac{\frac{a^2}{2} c}{1 + \frac{a^2}{2} c}, c = \frac{(p+1) + (\tilde{p}+1)}{(p+1)(\tilde{p}+1)}$$

$$ds^2 = H^{\frac{2x-2}{\tilde{p}+1}} (dt^2 - d\vec{y}_p^2) - H^{\frac{-(2x-2)}{\tilde{p}+1}} d\vec{x}_{(\tilde{p}+3)}^2,$$

$$e^{-2a\varphi} = e^{-2a\varphi_0} H^{2x}, A_{t\underline{t}g^1 \dots \underline{y}^p} = e^{a\varphi_0} \alpha(H^{-1} - 1)$$

$$\partial_{\underline{m}} \partial_{\underline{m}} H = 0,$$

$$x = \frac{\frac{a^2}{2}}{1 + \frac{a^2}{2}}, c = \frac{(p+1) + (\tilde{p}+1)}{(p+1)(\tilde{p}+1)}, \alpha^2 = \frac{4x}{a^2}$$

$$H = 1 + \frac{h}{|\vec{x}_{(\tilde{p}+3)}|^{\tilde{p}+1}}$$

$$H = 1 + \sum_{I=1}^N \frac{h_I}{|\vec{x}_{(\tilde{p}+3)} - \vec{x}_{(\tilde{p}+3)I}|^{\tilde{p}+1}}$$

$$ds^2 = W dt^2 - d\vec{y}_p^2 - W^{-1} d\rho^2 + \rho^2 d\Omega_{(\tilde{p}+2)}^2$$

$$W = 1 + \frac{\omega}{\rho^{\tilde{p}+1}}$$

$$S = S_a + S_p$$

$$\begin{aligned}
S_p[X^\mu, \gamma_{ij}] &= -\frac{T}{2} \int d^{p+1} \xi \sqrt{|\gamma|} [e^{-2b\varphi} \gamma^{ij} \partial_i X^\mu \partial_j X^\nu g_{\mu\nu} - (p-1)] \\
&\quad - \frac{\mu}{(p+1)!} \int d^{p+1} \xi A_{(p+1)\mu_1 \dots \mu_{p+1}} \partial_{i_1} X^{\mu_1} \dots \partial_{i_{p+1}} X^{\mu_{p+1}} \\
&\quad Y^i(\xi) = \xi^i \\
&\quad X^m(\xi) = 0 \\
&\quad a = -(p+1)b \\
&\quad \mu = T/\alpha \\
&\quad H = \epsilon + \frac{h}{|\vec{x}_{(\tilde{p}+3)}|^{\tilde{p}+1}} \\
&\quad h = \frac{16\pi G_N^{(d)} T}{(\tilde{p}+1)\alpha^2 \omega_{(\tilde{p}+2)}} \\
d\hat{s}_E^2 &= H_{M2}^{-2/3} [W dt^2 - d\vec{y}_2^2] - H_{M2}^{1/3} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(7)}^2] \\
\hat{C}_{t\underline{y}^1 \underline{y}^2} &= \alpha(H_{M2}^{-1} - 1), \\
H_{M2} &= 1 + \frac{h_{M2}}{\rho^6}, W = 1 + \frac{\omega}{\rho^6}, \\
\omega &= h_{M2} [1 - \alpha^2] \\
d\hat{s}^2 &= H_{M2}^{-2/3} [dt^2 - d\vec{y}_2^2] - H_{M2}^{1/3} d\vec{x}_8^2, \\
\hat{C}_{ty_1 y_2} &= \pm(H_{M2}^{-1} - 1), \\
H_{M2} &= 1 + \frac{h_{M2}}{|\vec{x}_8|^6} \\
h_{M2} &= \frac{16\pi G_N^{(11)} T_{M2}}{6\omega_{(7)}} \\
T_{M2} &= \frac{M_{M2}}{(2\pi)^2 R_9 R_{10}} = \frac{1}{(2\pi)^2 \left(\ell_{\text{Planck}}^{(11)}\right)^3} = \frac{2\pi}{\left(\ell_{\text{Planck}}^{(11)}\right)^3} \\
h_{M2} &= \frac{\left(\ell_{\text{Planck}}^{(11)}\right)^6}{6\omega_{(7)}} \\
d\hat{s}^2 &= H_{M5}^{-1/3} [W dt^2 - d\vec{y}_5^2] - H_{M5}^{2/3} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(4)}^2] \\
\tilde{\hat{C}}_{t\underline{t} y^1 \dots \underline{y}^5} &= \alpha(H_{M5}^{-1} - 1), \\
H_{M5} &= 1 + \frac{h_{M5}}{\rho^3}, W = 1 + \frac{\omega}{\rho^3}, \\
\omega &= h_{M5} [1 - \alpha^2] \\
d\hat{s}^2 &= H_{M5}^{-1/3} [dt^2 - d\vec{y}_5^2] - H_{M5}^{2/3} d\vec{x}_5^2, \\
\tilde{\hat{C}}_{t\underline{t}^1 \dots \underline{y}^5} &= \pm(H_{M5}^{-1} - 1), \\
H_{M5} &= 1 + \frac{h_{M5}}{|\vec{x}_5|^3} \\
h_{M5} &= \frac{\left(\ell_{\text{Planck}}^{(11)}\right)^3}{3\omega_{(4)}}
\end{aligned}$$

$$S = \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{|g_s|} \left\{ e^{-2\phi} \left[R_s - 4(\partial\phi)^2 + \frac{(-1)^{p_1+1}}{2 \cdot (p_1+2)!} F_{(p_1+2)}^2 \right] + \frac{(-1)^{p_2+1}}{2 \cdot (p_2+2)!} F_{(p_2+2)}^2 \right\}$$

$$S = \frac{1}{16\pi G_N^{(d)}} \int d^d x \sqrt{|g|} \left[R + \frac{4}{(d-2)} (\partial\phi)^2 + \frac{(-1)^{p_1+1}}{2 \cdot (p_1+2)!} e^{-4\frac{(p_1+1)}{(d-2)}\phi} F_{(p_1+2)}^2 + \frac{(-1)^{p_2+1}}{2 \cdot (p_2+2)!} e^{2\frac{(\tilde{p}_2-p_1)}{(d-2)}\phi} F_{(p_2+2)}^2 \right]$$

$$\phi = \sqrt{\frac{(d-2)}{2}} \varphi$$

$$a_1 = \frac{2(p_1+1)}{\sqrt{2(d-2)}}, \quad (\text{NS} - \text{NS})$$

$$a_2 = \frac{-(\tilde{p}_2 - p_2)}{\sqrt{2(d-2)}}. \quad (\text{RR})$$

$$a_3 = -\frac{2(\tilde{p}_1+1)}{\sqrt{2(d-2)}}$$

$$d\tilde{s}_E^2 = H_{F1}^{-3/4} [W dt^2 - dy^2] - H_{F1}^{1/4} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(7)}^2],$$

$$d\hat{s}_S^2 = H_{F1}^{-1} [W dt^2 - dy^2] - [W^{-1} d\rho^2 + \rho^2 d\Omega_{(7)}^2],$$

$$e^{-2(\hat{\phi}-\hat{\phi}_0)} = H_{F1},$$

$$\hat{B}_{t\underline{y}} = \alpha(H_{F1}^{-1} - 1),$$

$$H_{F1} = 1 + \frac{h_{F1}}{\rho^6}, W = 1 + \frac{\omega}{\rho^6},$$

$$\omega = h_{F1}[1 - \alpha^2]$$

$$d\tilde{s}_E^2 = H_{F1}^{-3/4} [dt^2 - dy^2] - H_{F1}^{1/4} d\vec{x}_8^2$$

$$d\hat{s}_S^2 = H_{F1}^{-1} [dt^2 - dy^2] - d\vec{x}_8^2,$$

$$e^{-2(\hat{\phi}-\hat{\phi}_0)} = H_{F1},$$

$$\hat{B}_{t\underline{y}} = \pm(H_{F1}^{-1} - 1),$$

$$H_{F1} = 1 + \frac{h_{F1}}{|\vec{x}_8|^6}$$

$$h_{F1} = \frac{2^5 \pi^6 \ell_s^6 g^2}{3\omega_{(7)}}$$

$$d\tilde{s}_E^2 = H_{S5}^{-1/4} [W dt^2 - d\vec{y}_5^2] - H_{S5}^{3/4} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(3)}^2]$$

$$d\hat{s}_S^2 = [W dt^2 - d\vec{y}_5^2] - H_{S5} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(3)}^2]$$

$$e^{-2(\hat{\phi}-\hat{\phi}_0)} = H_{S5}^{-1},$$

$$\hat{B}^{(6)}_{t\underline{ty}^1 \dots \underline{y}^5} = \alpha e^{-2\hat{\phi}_0} (H_{S5}^{-1} - 1)$$

$$H_{S5} = 1 + \frac{h_{ps}}{\rho^2}, W = 1 + \frac{\omega}{\rho^2},$$

$$\omega = h_{S5}[1 - \alpha^2]$$

$$\begin{aligned}
d\tilde{s}_E^2 &= H_{S5}^{-1/4} [dt^2 - d\vec{y}_5^2] - H_{S5}^{3/4} d\vec{x}_4^2, \\
d\hat{s}_S^2 &= [dt^2 - d\vec{y}_5^2] - H_{S5} d\vec{x}_4^2, \\
e^{-2(\hat{\phi}-\hat{\phi}_0)} &= H_{S5}^{-1}, \\
\tilde{B}^{(6)}{}_{\underline{t}\underline{t}\underline{y}^1\dots\underline{y}^5} &= \pm e^{-2\hat{\phi}_0} (H_{S5}^{-1} - 1), \\
H_{S5} &= 1 + \frac{h_{S5}}{|\vec{x}_4|^2} \\
h_{S5} &= \ell_s^2 \\
d\tilde{s}_E^2 &= H_{Dp}^{-\frac{(7-p)}{8}} [W dt^2 - d\vec{y}_p^2] - H_{Dp}^{\frac{(p+1)}{8}} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(8-p)}^2] \\
d\hat{s}_S^2 &= H_{Dp}^{-1/2} [W dt^2 - d\vec{y}_p^2] - H_{Dp}^{1/2} [W^{-1} d\rho^2 + \rho^2 d\Omega_{(8-p)}^2], \\
e^{-2(\hat{\phi}-\hat{\phi}_0)} &= H_{Dp}^{\frac{(p-3)}{2}}, \\
\hat{C}^{(p+1)}{}_{\underline{t}\underline{y}^1\dots\underline{y}^p} &= \alpha e^{-\hat{\phi}_0} (H_{Dp}^{-1} - 1), \\
H_{Dp} &= 1 + \frac{h_{Dp}}{\rho^{7-p}}, W = 1 + \frac{\omega}{\rho^{7-p}}, \\
\omega &= h_{Dp} [1 - \alpha^2] \\
H &\sim h \log \rho, W \sim \omega \log \rho \\
H &\sim h\rho, W \sim \omega\rho \\
d\tilde{s}_E^2 &= H_{Dp}^{\frac{p-7}{8}} [dt^2 - d\vec{y}_p^2] - H_{Dp}^{\frac{p+1}{8}} d\vec{x}_{9-p}^2, \\
d\hat{s}_S^2 &= H_{Dp}^{-1/2} [dt^2 - d\vec{y}_p^2] - H_{Dp}^{1/2} d\vec{x}_{9-p}^2, \\
e^{-2(\hat{\phi}-\hat{\phi}_0)} &= H_{Dp}^{\frac{(p-3)}{2}}, \\
\hat{C}^{(p+1)}{}_{\underline{t}\underline{y}^1\dots\underline{y}^p} &= \pm e^{-\hat{\phi}_0} (H_{Dp}^{-1} - 1), \\
H_{Dp} &= 1 + \frac{h_{Dp}}{|\vec{x}_{9-p}|^{7-p}} \\
H_{D7} &= 1 + h_{D7} \log |\vec{x}_2| \\
H_{D8} &= 1 + h_{D8} |x| \\
h_{Dp} &= \frac{(2\pi\ell_s)^{(7-p)} g}{(7-p)\omega_{(8-p)}} \\
h_{D7} &=, h_{D8} = \frac{g}{4\pi\ell_s} \\
\left\{ \begin{aligned} d\hat{s}^2 &= H_{M2}^{-2/3} \left[dt^2 - dy^2 - e^{\frac{4}{3}\hat{\phi}_0} dz^2 \right] - H_{M2}^{1/3} d\vec{x}_8^2 \\ \hat{C}_{\underline{t}\underline{y}\underline{z}} &= \pm e^{\frac{2}{3}\hat{\phi}_0} (H_{M2}^{-1} - 1) \\ H_{M2} &= 1 + \frac{h_{M2}}{|\vec{x}_8|^6} \end{aligned} \right. \\
h_{M2} &= \frac{(\ell_{\text{Planck}}^{(11)})^6}{6\omega_{(7)}} = \frac{(2\pi\ell_s g^{1/3})^6}{6\omega_{(7)}} = \frac{(2\pi\ell_s)^6 g^2}{6\omega_{(7)}} = h_{D2}
\end{aligned}$$



$$\begin{aligned}
H_{M2} &= 1 + h_{M2} \sum_{n=-\infty}^{n=+\infty} \frac{1}{(|\vec{x}_7|^2 + (z + 2\pi n R_{11})^2)^3} \\
u_n &= \frac{(z - 2\pi n R_{11})}{|\vec{x}_7|}, u_n \in \left[\frac{2\pi n R_{11}}{|\vec{x}_7|}, \frac{2\pi(n+1)R_{11}}{|\vec{x}_7|} \right] \\
H_{M2} &= 1 + \frac{h_{M2}}{|\vec{x}_7|^6} \sum_{n=-\infty}^{n=+\infty} \frac{1}{(1 + u_n^2)^3} \sim 1 + \frac{h_{M2}}{|\vec{x}_7|^6} \frac{1}{2\pi R_{11}/|\vec{x}_7|} \int_{-\infty}^{+\infty} \frac{du}{(1 + u^2)^3} \\
&= 1 + \frac{h_{M2} \omega_{(5)}}{2\pi R_{11} \omega_{(4)}} \frac{1}{|\vec{x}_7|^5} \\
\frac{h_{M2} \omega_{(5)}}{2\pi R_{11} \omega_{(4)}} &= h_{D2} \\
H_p &= 1 + h_p \sum_{n=-\infty}^{n=+\infty} \frac{1}{(|\vec{x}_{n+1}|^2 + (z + 2\pi n R)^2)^{n/2}} \sim 1 + \frac{h_p \omega_{(n-1)}}{2\pi R \omega_{(n-2)}} \frac{1}{|\vec{x}_{n+1}|^{n-1}} \\
h'_p &= \frac{h_p \omega_{(n-1)}}{2\pi R \omega_{(n-2)}} \\
h'_p &= \frac{h_p \omega_{(n-1)}}{V^m \omega_{(n-m-1)}}, V^m = (2\pi)^m R_1 \dots R_m \\
\frac{h_{Dp} \omega_{(6-p)}}{2\pi R \omega_{(5-p)}} &= \frac{(2\pi \ell_s)^{7-p} g}{2\pi R} \frac{\omega_{(6-p)}}{(7-p) \omega_{(8-p)} \omega_{(6-p)}} \\
\frac{\omega_{(n-1)}}{n \omega_{(n+1)} \omega_{(n-2)}} &= \frac{1}{(n-1) \omega_{(n)}} \\
\begin{cases} \hat{e}_{\underline{i}^2}{}^j = H_{M2}^{-1/3} \delta_i^j \\ \hat{e}_{\underline{m}}^n = H_{M2}^{1/6} \delta_m^n \end{cases} \\
\begin{cases} \hat{\omega}_{\underline{m}}{}^{nl} = -\frac{1}{3} H_{M2}^{-1} \partial_{\underline{q}} H_{M2} \eta_m^{[n} \eta^{p]q} \\ \hat{\omega}_{\underline{2}}{}^{mj} = \frac{2}{3} H_{M2}^{-3/2} \partial_{\underline{q}} H_{M2} \eta_i^{[m} \eta^{j]q} \\ \hat{G}_{\underline{m} \underline{n} y^1 y^2} = \mp H_{M2}^{-2} \partial_{\underline{m}} H_{M2} \end{cases} \\
\begin{cases} \delta_{\hat{\hat{e}}} \hat{\hat{\psi}}_{\underline{i}} = \frac{1}{3} H_{M2}^{-3/2} \partial_{\underline{n}} H_{M2} \hat{\hat{\Gamma}}_{(i)}^n \left(1 \mp \frac{i}{2} \epsilon(i) j k \hat{\hat{\Gamma}}^{(i)jk} \right) \hat{\hat{e}} = 0, \\ \delta_{\hat{\hat{e}}} \hat{\hat{\psi}}_{\underline{m}} \end{cases} \\
= 2 \partial_{\underline{m}} \hat{\hat{e}} - \frac{1}{6} H_{M2}^{-1} \partial_{\underline{n}} H_{M2} \left[\hat{\hat{\Gamma}}^{mn} \mp i \left(\hat{\hat{\Gamma}}^{mn} + 2 \delta^{mn} \right) \hat{\hat{\Gamma}}^{012} \right] \hat{\hat{e}} \\
\left(1 \mp i \hat{\hat{\Gamma}}^{012} \right) \hat{\hat{e}} = 0 \\
\delta_{\hat{\hat{e}}} \hat{\hat{\psi}}_{\underline{m}} = 2 \left(\partial_{\underline{m}} + \frac{1}{6} H_{M2}^{-1} \partial_{\underline{m}} H_{M2} \right) \hat{\hat{e}} \\
\hat{\hat{e}} = H_{M2}^{-1/6} \hat{\hat{e}}_0, \left(1 \mp i \hat{\hat{\Gamma}}^{012} \right) \hat{\hat{e}}_0 \\
\begin{cases} \hat{\hat{e}}_{\underline{i}}^j = H_{M5}^{-1/6} \delta_i^j, \\ \hat{\hat{e}}_{\underline{m}}^n = H_{M5}^{1/3} \delta_m^n \end{cases}
\end{aligned}$$

$$\begin{aligned}
& \begin{cases} \hat{\omega}_{\underline{m}}^{nl} = -\frac{2}{3} H_{M5}^{-1} \partial_{\underline{q}} H_{M2} \eta_m^{[n} \eta^{p]q} \\ \hat{\omega}_{\underline{m}}^{mj} = \frac{1}{3} H_{M5}^{-3/2} \partial_{\underline{q}} H_{M2} \eta_i^{[m} \eta^{j]q} \\ \hat{G}_{\underline{m}_1 \dots \underline{m}_4} = \pm \epsilon_{m_1 \dots m_5} \partial_{\underline{m}_5} H_{M5} \\ \hat{\epsilon} = H_{M5}^{-1/12} \hat{\epsilon}_0, (1 \mp \hat{\Gamma}^{012345}) \hat{\epsilon}_0 \end{cases} \\
& \begin{cases} \delta_{\hat{\epsilon}} \hat{\psi}_{\hat{\mu}} = \left\{ \partial_{\hat{\mu}} - \frac{1}{4} \hat{\psi}_{\hat{\mu}} + \frac{i}{8} e^{\hat{\phi}} \frac{1}{(p+2)!} \hat{\epsilon}^{(p+2)} \hat{\Gamma}_{\hat{\mu}} (-\hat{\Gamma}_{11})^{\frac{p+2}{2}} \right\} \hat{\epsilon} \\ \delta_{\hat{\epsilon}} \hat{\lambda} = \left\{ \partial \hat{\phi} - \frac{i}{4} e^{\hat{\phi}} \frac{(p-3)}{(p+2)!} \hat{\zeta}^{(p+2)} (-\hat{\Gamma}_{11})^{\frac{p+2}{2}} \right\} \hat{\epsilon} \\ \delta_{\hat{\epsilon}} \hat{\zeta}_{\hat{\mu}} = \left\{ \partial_{\hat{\mu}} - \frac{1}{4} \hat{\psi}_{\hat{\mu}} + \frac{1}{8} e^{\hat{\phi}} \frac{1}{(p+2)!} \hat{\epsilon}^{(p+2)} \hat{\Gamma}_{\hat{\mu}} \mathcal{P}_{\frac{p+3}{2}} \right\} \hat{\epsilon} \\ \delta_{\hat{\epsilon}} \hat{\chi} = \left\{ \partial \hat{\phi} + \frac{1}{4} e^{\hat{\phi}} \left(\frac{(p-3)}{(p+2)!} \hat{\epsilon}^{(p+2)} \mathcal{P}_{\frac{p+3}{2}} \right) \right\} \hat{\epsilon} \end{cases} \\
& \mathcal{P}_n = \begin{cases} \sigma^1, n \text{ par} \\ i\sigma^2, n \text{ impar} \end{cases} \\
& \begin{cases} \hat{\psi}_{\underline{i}} = -\frac{1}{2} H_{Dp}^{-3/2} \partial_{\underline{n}} H_{Dp} \eta_{ij} \Gamma^{nj} \\ \hat{\psi}_{\underline{m}} = \frac{1}{2} H_{Dp}^{-1} \partial_{\underline{n}} H_{Dp} \eta_{mq} \hat{\Gamma}^{nq} \\ \hat{G}^{(p+2)} = \mp e^{-\hat{\phi}_0} H_{Dp}^{\frac{p}{4}-2} \partial_{\underline{m}} H_{Dp} \hat{\Gamma}^m \hat{\Gamma}^{01 \dots p} \end{cases} \\
\text{IIA: } \hat{\epsilon} = H_{Dp}^{-1/8} \hat{\epsilon}_0, \quad \left[1 \mp i \hat{\Gamma}^{01 \dots p} (-\hat{\Gamma}_{11})^{\frac{p+2}{2}} \right] \hat{\epsilon}_0 = 0 \\
\text{IIB: } \hat{\epsilon} = H_{Dp}^{-1/8} \hat{\epsilon}_0, \quad \left(1 \pm i \hat{\Gamma}^{01 \dots p} \mathcal{P}_{\frac{p+3}{2}} \right) \hat{\epsilon}_0 = 0 \\
& \begin{cases} \delta_{\hat{\epsilon}} \hat{\psi}_{\hat{\mu}} = \left\{ \partial_{\hat{\mu}} - \frac{1}{4} \left(\hat{\psi}_{\hat{\mu}} + \frac{1}{2} \hat{H}_{\hat{\mu}} \mathcal{O} \right) \right\} \hat{\epsilon} \\ \delta_{\hat{\epsilon}} \hat{\lambda} = \left\{ \partial \hat{\phi} - \frac{1}{12} \hat{H} \mathcal{O} \right\} \epsilon \end{cases} \\
& \psi_{\underline{i}} = -H_{F1}^{-3/2} \partial_{\underline{m}} H_{F1} \Gamma_i^m \\
& H_{\underline{i}} = \mp 2 H_{F1}^{-1} \partial_{\underline{m}} H_{F1} \Gamma^{01} \\
& H_{\underline{m}} = \pm \epsilon_{ij} H_{F1}^{-3/2} \partial_{\underline{m}} H_{F1} \Gamma^{mj} \\
& \hat{\epsilon} = H_{F1}^{1/4} \hat{\epsilon}_0, (1 \pm \hat{\Gamma}^{01} \mathcal{O}) \hat{\epsilon}_0 \\
& \begin{cases} \delta_{\hat{\epsilon}} \hat{\psi}_{\hat{\mu}} = \left\{ \partial_{\hat{\mu}} - \frac{1}{4} \left(\hat{\omega}_{\hat{\mu}} + \frac{1}{7!} e^{2\hat{\phi}} \hat{H}^{(7)} \hat{a}_1 \dots \hat{a}_7 \right. \right. \\ \quad \left. \left. \hat{\Gamma}_{\hat{\mu} \hat{a}_1 \dots \hat{a}_7} \mathcal{O} \right) \right\} \hat{\epsilon} \\ \delta_{\hat{\epsilon}} \hat{\lambda} = \left\{ \partial \hat{\phi} + \frac{1}{2} \hat{H}^{(7)} \mathcal{O} \right\} \hat{\epsilon} \\ \hat{\epsilon} = \hat{\epsilon}_0, (1 \pm \hat{\Gamma}^{0 \dots 5} \mathcal{O}) \hat{\epsilon}_0 \end{cases}
\end{aligned}$$

$$\left\{ \begin{array}{l} d\tilde{s}_E^2 = H_{F1}^{-6/7} dt^2 - H_{F1}^{1/7} d\vec{x}_8^2 \\ d s_s^2 = H_{F1}^{-1} dt^2 - d\vec{x}_8^2 \\ A_t = \pm(H_{F1}^{-1} - 1) \\ e^{-2(\phi-\phi_0)} = H_{F1}^{1/2} \\ K/K_0 = H_{F1}^{-1/2} \end{array} \right. \left\{ \begin{array}{l} d\tilde{s}_E^2 = H_{S5}^{-2/3} dt^2 - H_{S5}^{1/3} d\vec{x}_4^2 \\ d s_s^2 = dt^2 - H_{S5} d\vec{x}_4^2, \\ A_t = \pm(H_{S5}^{-1} - 1), \\ e^{-2(\phi-\phi_0)} = H_{S5}^{-1}, \\ K/K_0 = 1 \end{array} \right. \left\{ \begin{array}{l} d\tilde{s}_E^2 = H_{Dp}^{\frac{7-p}{8-p}} dt^2 - H_{Dp}^{\frac{1}{8-p}} d\vec{x}_{9-p}^2 \\ d s_s^2 = H_{Dp}^{-1/2} dt^2 - H_{Dp}^{1/2} d\vec{x}_4^2 \\ A_t = \pm(H_{Dp}^{-1} - 1) \\ e^{-2(\phi-\phi_0)} = H_{Dp}^{\frac{p-6}{4}} \\ K/K_0 = H_{Dp}^{-p/4} \end{array} \right.$$

5 – pluridimensión ||+| + + + + - - -

$$P_p \epsilon = (1 \pm \Gamma^{01 \cdots p} \mathcal{O}_p) \epsilon$$

$$[P_p, P_{p'}] = 0$$

$$\begin{array}{l} D_p \perp D_{(p+4)}(p) \\ q_{F1} \sim \int_{S^7} e^{-2\varphi^* \hat{\mathcal{H}}} \\ de^{-2\varphi^* \hat{\mathcal{H}}} = 0 \\ q_{F1} \sim \int_{S^7} (e^{-2\varphi^* \hat{\mathcal{H}}} - {}^* \hat{G}^{(3)} \hat{C}^{(0)} - \hat{G}^{(5)} \hat{C}^{(2)}) \end{array}$$

$$\begin{array}{l} \int_{S^5} \hat{G}^{(5)} \int_{S^2} \hat{C}^{(2)} \\ q_{F1} \sim \int_{S^2} dV^{(1)} \\ F1 \perp D_p(0) \\ D_p \perp D_{p+2}(p-1), p \geq 1 \\ D_p \perp D_{p+4}(p) \\ F1 \perp S5 \text{ B}(0) \\ D_p \perp D_p(p-2), p \geq 2 \\ D_p \perp S5(p-1), p \geq 1 \end{array}$$

$$\begin{array}{l} F1||S5, \; F1 \perp D_p(0), \\ S5 \perp S5(1), S5 \perp S5(1), S5 \perp D_p(p-1)(p>1), \\ D_p \perp D_{p'}(m)p+p'=4+2m, \\ W||F1, \; W||S5, \; W||D_p, \\ KK \perp D_p(p-2). \end{array}$$



$$\begin{aligned}
& M2 \perp M2(0), M2 \perp M5(1), M5 \perp M5(1), M5 \perp M5(3), \\
& W \parallel M2, W \parallel M5, \\
& KK \parallel M2, KK \perp M2(0), KK \parallel M5, KK \perp M5(1), KK \perp M5(3), \\
& W \parallel KK, W \perp KK(2), W \perp KK(4). \\
& d\hat{s}_s^2 = H_{Dp}^{-1/2} H_{F1}^{-1} dt^2 - H_{Dp}^{+1/2} H_{F1}^{-1} dy^2 - H_{Dp}^{-1/2} d\vec{z}_p^2 - H_{Dp}^{+1/2} d\vec{x}_{8-p}^2, \\
& e^{-2(\hat{\phi} - \hat{\phi}_0)} = H_{Dp}^{\frac{(p-3)}{2}} H_{F1}, \\
& \hat{C}^{(p+1)} t_{t\perp} \dots \dots \underline{\underline{z}}^p = \pm e^{-\hat{\phi}_0} (H_{Dp}^{-1} - 1), \\
& \hat{B}_{t\perp} = \pm (H_{F1}^{-1} - 1), \\
& H_{Dp,F1} = 1 + \frac{h_{Dp,F1}}{|\vec{x}_{8-p}|^{6-p}} \\
& \left\{ \begin{aligned} ds^2 &= H^\alpha [W dt^2 - d\vec{y}_{p-1}^2 - dz^2] - H^\beta [W^{-1} d\rho^2 + \rho^2 d\Omega^2] \\ e^{-2(\phi - \phi_0)} &= H^\gamma \\ A_{(p+1)t y^1 \dots y^{p-1} \underline{\underline{z}}} &= \alpha (H^{-1} - 1) \\ W &= 1 + \frac{\omega}{\rho^n}, H = 1 + \frac{h}{\rho^n} \\ \begin{pmatrix} t \\ z \end{pmatrix} &\rightarrow \begin{pmatrix} \cosh \gamma & \sinh \gamma \\ \sinh \gamma & \cosh \gamma \end{pmatrix} \begin{pmatrix} t \\ z \end{pmatrix} \\ W dt^2 - dz^2 &\rightarrow dt^2 - dz^2 + \cosh^2 \gamma (W - 1) (dt + \tanh^2 \gamma dz)^2 \\ H_W^{-1} dt^2 - H_W [dz - (H_W^{-1} - 1) dt]^2, & H_W = 1 + \frac{h_W}{\rho^n} \\ ds^2 &= H^\alpha \{ H_W^{-1} dt^2 - H_W [dz - (H_W^{-1} - 1) dt]^2 - d\vec{y}_{p-1}^2 \} - H^\beta d\vec{x}^2 \\ ds^2 &= H_W^{-1} dt^2 - H_W [dz + \alpha (H_W^{-1} - 1) dt]^2 - d\vec{x}_{d-2}^2, \\ H_W &= 1 + \frac{h_W}{|\vec{x}_{d-2}|^{d-4}}, \alpha = \pm 1 \\ H_W &= 1 + \frac{h_W}{|\vec{x}_{d-2}|^{d-4}} \delta(u_\alpha) \cdot u_\alpha = \frac{1}{\sqrt{2}} (t - \alpha z) \\ h_W &= -\alpha \frac{\sqrt{2} |p^z| 8\pi G_N^{(d)}}{(d-4) \omega_{(d-3)}} \\ \delta(u_\alpha) &\sim -\alpha \frac{\sqrt{2}}{2\pi R_z} \\ h_W &= \frac{|N| 8G_N^{(d)}}{R_z^2 (d-4) \omega_{(d-3)}} \end{aligned} \right. \\
& \left\{ \begin{aligned} d\hat{s}_s^2 &= H_{D1}^{-1/2} H_{D5}^{-1/2} \{ H_W^{-1} dt^2 - H_W [dy^1 + \alpha_W (H_W^{-1} - 1) dt]^2 \} \\ &\quad - H_{D1}^{1/2} H_{D5}^{-1/2} d\vec{y}_4^2 - H_{D1}^{1/2} H_{D5}^{1/2} d\vec{x}_4^2, \\ e^{-2(\hat{\phi} - \hat{\phi}_0)} & \\ \hat{C}_{D5}/H_{D1} & \\ t \hat{t} \underline{\underline{y}}^{(2)} &= \alpha_{D1} (H_{D1}^{-1} - 1) \\ \hat{C}_{tt^1 \dots y^5}^{(6)} &= \alpha_{D5} (H_{D5}^{-1} - 1) \end{aligned} \right.
\end{aligned}$$

$$\begin{aligned}
H_i &= 1 + \frac{r_i^2}{|\vec{x}_4|^2}, i = D1, D5, W \\
r_{D5}^2 &= N_{D5} h_{D5} = N_{D5} \ell_s^2 g \\
r_{D1}^2 &= N_{D1} h_{D1} \frac{\omega_{(5)}}{V^4 \omega_{(1)}} = \frac{N_{D1} \ell_s^6 g}{V}, V \equiv R_2 \dots R_5 \\
r_W^2 &= h_W \frac{\omega_{(5)}}{V^4 \omega_{(1)}} = \frac{N_W \ell_s^8 g^2}{R^2 V} \\
d\tilde{s}_E^2 &= (H_{D1} H_{D5} H_W)^{-2/3} dt^2 - (H_{D1} H_{D5} H_W)^{1/3} d\vec{x}_4^2 \\
ds_s^2 &= (H_{D1} H_{D5})^{-1/2} H_W^{-1} dt^2 - (H_{D1} H_{D5})^{1/2} d\vec{x}_4^2 \\
A^{(D1, D5, W)t} &= \alpha_{D1, D5, W} (H_{D1, D5, W}^{-1} - 1) \\
K_V/K_{V0} &= H_{D1}/H_{D5}, e^{-2(\phi - \phi_0)} = K_R/K_{R0} = (H_{D1} H_{D5})^{-1/4} H_W^{1/2} \\
A &= \omega_{(3)} \left(\lim_{|\vec{x}_4| \rightarrow 0} |\vec{x}_4|^6 H_{D1} H_{D5} H_W \right)^{1/2} = 2\pi^2 (r_{D1} r_{D5} r_W)^{1/2} \\
&= 2\pi^2 \sqrt{N_{D1} N_{D5} N_W} \frac{\ell_s^8 g^2}{RV} \\
S &= \frac{A}{4G_N^{(5)}}, G_N^{(5)} = \frac{G_N^{(10)}}{(2\pi)^5 RV} = \frac{\pi \ell_s^8 g^2}{4 RV} \\
S &= 2\pi \sqrt{N_{D1} N_{D5} N_W} \\
M &= \frac{N_{D1} R}{g \ell_s} + \frac{N_{D5} RV}{g \ell_s^6} + \frac{N_W}{R} \\
ds^2 &= H^{-2} dt^2 - H d\vec{x}_4^2, H = H_{D1} = H_{D5} = H_W \\
\rho(E) &\sim e^{\sqrt{\pi(c-24E_0)EL/3}} \\
\rho(E) &= e^{2\pi\sqrt{N_{D1} N_{D5} N_W}} \\
e_a{}^\mu e_b{}^\nu g_{\mu\nu} &= \eta_{ab}, e_\mu{}^a e_\nu{}^b \eta_{ab} = g_{\mu\nu} \\
\nabla_\mu \xi^\nu &= \partial_\mu \xi^\nu + \Gamma_{\mu\rho}{}^\nu \xi^\rho, \\
\mathcal{D}_\mu \xi^a &= \partial_\mu \xi^a + \omega_{\mu b}{}^a \xi^b, \\
\nabla_\mu \psi &= \partial_\mu \psi - \frac{1}{4} \omega_\mu{}^{ab} \Gamma_{ab} \psi \\
[\nabla_\mu, \nabla_\nu] \xi^\rho &= R_{\mu\nu\sigma}{}^\rho(\Gamma) \xi^\sigma + T_{\mu\nu}{}^\sigma \nabla_\sigma \xi^\rho \\
[\mathcal{D}_\mu, \mathcal{D}_\nu] \xi^a &= R_{\mu\nu b}{}^a(\omega) \xi^b \\
R_{\mu\nu\rho}^\sigma(\Gamma) &= 2\partial_{[\mu} \Gamma_{\nu]\rho}^\sigma + 2\Gamma_{[\mu|\lambda}{}^\sigma \Gamma_{\nu]\rho}^\lambda \\
R_{\mu\nu a}{}^b(\omega) &= 2\partial_{[\mu} \omega_{\nu]a}^b - 2\omega_{[\mu|a}{}^c \omega_{|\nu]c}^b \\
\nabla_\mu e_a^\mu &= 0 \\
\omega_{\mu a}{}^b &= \Gamma_{\mu a}{}^b + e_a{}^\nu \partial_\mu e_\nu^b \\
R_{\mu\nu\rho}{}^\sigma(\Gamma) &= e_\rho{}^a e^\sigma{}_b R_{\mu\nu a}{}^b(\omega) \\
\nabla_\mu g_{\rho\sigma} &= 0 \\
\Gamma_{\mu\nu}^\rho &= \left\{ \begin{smallmatrix} \rho \\ \mu\nu \end{smallmatrix} \right\} + K_{\mu\nu}^\rho = \Gamma_{\mu\nu}^\rho(g) + K_{\mu\nu}^\rho \\
\left\{ \begin{smallmatrix} \rho \\ \mu\nu \end{smallmatrix} \right\} &= \frac{1}{2} g^{\rho\sigma} \{ \partial_\mu g_{\nu\sigma} + \partial_\nu g_{\mu\sigma} - \partial_\sigma g_{\mu\nu} \} \\
K_{\mu\nu}^\rho &= \frac{1}{2} g^{\rho\sigma} \{ T_{\mu\sigma\nu} + T_{\nu\sigma\mu} - T_{\mu\nu\sigma} \} \\
\omega_{abc} &= \omega_{abc}(e) + K_{abc}, \omega_{abc}(e) = -\Omega_{abc} + \Omega_{bca} - \Omega_{cab}, \Omega_{ab}^c = e_a{}^\mu e_b{}^\nu \partial_{[\mu} e^c{}_{\nu]}
\end{aligned}$$

$$\begin{aligned}
\{\hat{\hat{a}}, \hat{\hat{b}}\} &= +2\hat{\eta}^{\hat{a}\hat{b}} \\
\hat{\hat{\Gamma}}_{\hat{1}\hat{0}} &= i\hat{\hat{\Gamma}}^{\hat{0}} \dots \hat{\hat{\Gamma}}^{\hat{b}} \equiv -i\hat{\hat{\Gamma}}_{\hat{1}\hat{1}} \\
\hat{\hat{\Gamma}}^{\hat{a}\star} &= -\hat{\hat{\Gamma}}^{\hat{a}} \\
\hat{\hat{\Gamma}}^{\hat{0}\dagger} &= +\hat{\hat{\Gamma}}^{\hat{0}} \\
\hat{\hat{\Gamma}}^{\hat{i}\dagger} &= -\hat{\hat{\Gamma}}^{\hat{i}}, \hat{i} = 1, \dots, 10. \\
\hat{\hat{\Gamma}}^{\hat{0}T} &= -\hat{\hat{\Gamma}}^{\hat{0}} \\
\hat{\hat{\Gamma}}^{\hat{i}T} &= +\hat{\hat{\Gamma}}^{\hat{i}}, \hat{i} = 1, \dots, 10. \\
\hat{\hat{\Gamma}}_{\hat{0}}^{\hat{a}} \hat{\hat{\Gamma}}^{\hat{a}\hat{0}} &= \hat{\hat{\Gamma}}^{\hat{a}\dagger} \\
\hat{\hat{\mathcal{D}}} &= i\hat{\hat{\Gamma}}^0 \\
\hat{\hat{\mathcal{D}}} \hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} \hat{\hat{\mathcal{D}}}^{-1} &= (-1)^{[n/2]} \left(\hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} \right)^\dagger \\
\hat{\hat{\mathcal{C}}} &= \hat{\hat{\mathcal{D}}} = i\hat{\hat{\Gamma}}^0 \\
\hat{\hat{\mathcal{C}}}^T &= \hat{\hat{\mathcal{C}}}^\dagger = \hat{\hat{\mathcal{C}}}^{-1} = -\hat{\hat{\mathcal{C}}} \\
\hat{\hat{\mathcal{C}}} \hat{\hat{\Gamma}}^{\hat{a}} \hat{\hat{\mathcal{C}}}^{-1} &= -\hat{\hat{\Gamma}}^{\hat{a}\hat{2}} \\
\hat{\hat{\mathcal{C}}} \hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} \hat{\hat{\mathcal{C}}}^{-1} &= (-1)^{n+[n/2]} \left(\hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} \right)^T \\
\bar{\hat{\lambda}} &= \hat{\lambda}^\dagger \hat{\mathcal{D}} \\
\hat{\lambda}^c &= \hat{\lambda}^T \hat{\mathcal{C}} \\
\bar{\hat{\lambda}} &= \hat{\lambda}^c \\
\bar{\hat{\epsilon}}^{\hat{a}_1 \dots \hat{a}_n} \hat{\hat{\psi}} &= (-1)^{n+[n/2]} \hat{\hat{N}}^{\hat{a}_1 \dots \hat{a}_n} \hat{\hat{\epsilon}} \\
\left(\bar{\hat{\epsilon}} \hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} \hat{\hat{\psi}} \right)^\dagger &= (-1)^{[n/2]} \bar{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} \hat{\hat{\epsilon}} \\
\hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} &= i \frac{(-1)^{[n/2]+1}}{(11-n)!} \hat{\hat{\epsilon}}^{\hat{a}_1 \dots \hat{a}_n \hat{b}_1 \dots \hat{b}_{11-n}} \hat{\hat{\Gamma}}_{\hat{b}_1 \dots \hat{b}_{11-n}} \\
\begin{cases} \hat{\hat{\Gamma}}^{\hat{a}} = \hat{\hat{\Gamma}}^{\hat{a}}, \hat{a} = 0, \dots, 9 \\ \hat{\hat{\Gamma}}^{10} = +i\hat{\hat{\Gamma}}^0 \dots \hat{\hat{\Gamma}}^9 \end{cases} \\
\hat{\hat{\Gamma}}_{11} &= -\hat{\hat{\Gamma}}^0 \dots \hat{\hat{\Gamma}}^9 = i\hat{\hat{\Gamma}}^{10} \\
\hat{\hat{\Gamma}}_{11} \hat{\hat{\psi}}^{(\pm)} &= \pm \hat{\hat{\psi}}^{(\pm)} \\
\hat{\hat{\Gamma}}_{11} &= \mathbb{I}_{16 \times 16} \otimes \sigma^3 = \begin{pmatrix} \mathbb{I}_{16 \times 16} & 0 \\ 0 & -\mathbb{I}_{16 \times 16} \end{pmatrix} \\
\hat{\hat{\psi}} &= \begin{pmatrix} \hat{\hat{\psi}}^{(+)} \\ \hat{\hat{\psi}}^{(-)} \end{pmatrix} \\
\hat{\hat{\Gamma}}_{11} \hat{\hat{\Gamma}}^{\hat{a}_1 \dots \hat{a}_n} &= \frac{(-1)^{[(10-n)/2]+1}}{(10-n)!} \hat{\hat{\epsilon}}^{\hat{a}_1 \dots \hat{a}_n \hat{b}_1 \dots \hat{b}_{10-n}} \hat{\hat{\Gamma}}_{\hat{b}_1 \dots \hat{b}_{10-n}} \\
\begin{cases} \hat{\hat{\Gamma}}^a = \Gamma^a \otimes \sigma^2, a = 0, \dots, 8 \\ \hat{\hat{\Gamma}}^9 = \mathbb{I}_{16 \times 16} \otimes i\sigma^1 \end{cases} \\
\Gamma^8 &= \Gamma^0 \dots \Gamma^7 \\
\Gamma_{(8)9} &= i\Gamma^8 = i\Gamma^0 \dots \Gamma^7 \\
\gamma_5 &= -i\gamma^0 \gamma^1 \gamma^2 \gamma^3 = \frac{i}{4!} \epsilon_{abcd} \gamma^{abcd} \\
\gamma^{a_1 \dots a_n} &= \frac{(-1)^{[n/2]} i}{(4-n)!} \epsilon^{a_1 \dots a_n b_1 \dots b_{4-n}} \gamma_{b_1 \dots b_{4-n}} \gamma_5
\end{aligned}$$

$$\begin{aligned}
& n^\mu n_\mu = \varepsilon, \begin{cases} \varepsilon = +1, & \Sigma \text{ espacio} \\ \varepsilon = -1, & \Sigma \text{ tiempo} \end{cases} \\
& h_{\mu\nu} = g_{\mu\nu} - \varepsilon n_\mu n_\nu \\
& \mathcal{K}_{\mu\nu} \equiv h_\mu{}^\alpha h_\nu{}^\beta \nabla_{(\alpha} n_{\beta)} \\
& \mathcal{K}_{\mu\nu} = \frac{1}{2} \mathcal{E}_n h_{\mu\nu} \\
& \mathcal{K} = h^{\mu\nu} \mathcal{K}_{\mu\nu} = h^{\mu\nu} \nabla_\mu n_\nu \\
& \left\{ \begin{aligned} x^1 &= \rho_{n-1} \sin \varphi \\ x^2 &= \rho_{n-1} \cos \varphi \\ x^3 &= \rho_{n-2} \cos \theta_1 \\ &\vdots \\ x^k &= \rho_{n-k+1} \cos \theta_{k-2}, 3 \leq k \leq n+1 \end{aligned} \right. \\
& \left\{ \begin{aligned} \rho_l &= [(x^1)^2 + \dots + (x^{n+1-l})^2]^{1/2} = r \prod_{m=1}^l \sin \theta_{n-m}, \\ \rho_0 &= r = [(x^1)^2 + \dots + (x^{n+1})^2]^{1/2} \end{aligned} \right. \\
& d\Omega^n \equiv d\varphi \prod_{i=1}^{n-1} \sin^i \theta_i d\theta_i \\
& d\Omega^n = \frac{1}{n! r^{n+1}} \epsilon_{\mu_1 \dots \mu_{n+1}} x^{\mu_{n+1}} dx^{\mu_1} \dots dx^{\mu_n} \\
& \begin{cases} d^{n+1}x = r^n dr d\Omega^n, \\ r^n d\Omega^n = d^n y \sqrt{|g|} \end{cases} \\
& \omega_{(n)} = \int_{S^n} d\Omega^n = \frac{2\pi^{\frac{n+1}{2}}}{\Gamma\left(\frac{n+1}{2}\right)} \\
& \Gamma(x+1) = x\Gamma(x), \Gamma(0) = 1, \Gamma(1/2) = \pi^{1/2} \\
& d\vec{x}^2 = d\rho_0^2 + \rho_0^2 d\theta_{n-1}^2 + \dots + \rho_{n-2}^2 d\theta_1^2 + \rho_{n-1}^2 d\varphi^2 \\
& = dr^2 + r^2 \{ d\theta_{n-1}^2 + \sin^2 \theta_{n-1} [d\theta_{n-2}^2 + \sin^2 \theta_{n-2} (d\theta_{n-3}^2 + \sin^2 \theta_{n-3} (\dots \\
& \quad \dots \sin^2 \theta_2 (d\theta_1^2 + \sin^2 \theta_1 d\varphi^2) \dots)] \} \\
& = dr^2 + r^2 d\Omega_{(n)}^2 \\
& \mathbf{P} = P^A \mathbf{G}_A \\
& [\mathbf{P}, \mathbf{Q}] = \mathbf{P}\mathbf{Q} - (-1)^{pq} \mathbf{Q}\mathbf{P} \\
& [\mathbf{P}, \mathbf{Q}] = P^A Q^B [\mathbf{G}_A, \mathbf{G}_B], \text{ si } P^A \text{ o } Q^B \\
& [\mathbf{P}, \mathbf{Q}] = P^A Q^B \{\mathbf{G}_A, \mathbf{G}_B\}, \text{ si } P^A \text{ y } Q^B \\
& [\mathbf{G}_A, \mathbf{G}_B] = \mathbf{G}_A \mathbf{G}_B - \mathbf{G}_B \mathbf{G}_A \\
& \{\mathbf{G}_A, \mathbf{G}_B\} = \mathbf{G}_A \mathbf{G}_B + \mathbf{G}_B \mathbf{G}_A \\
& \mathbf{D}\mathbf{Z} = d\mathbf{Z} + [\mathbf{A}, \mathbf{Z}] \\
& \mathbf{A} \rightarrow \mathbf{A}' = g(\mathbf{A} - g^{-1} d\mathbf{g})g^{-1} \\
& \delta \mathbf{A} = -\mathbf{D}\lambda \\
& \langle \dots \rangle_r: \underbrace{\mathfrak{g} \times \dots \times \mathfrak{g}}_r \rightarrow \mathbb{C} \\
& \langle \dots \mathbf{P}\mathbf{Q} \dots \rangle_r = (-1)^{pq} \langle \dots \mathbf{Q}\mathbf{P} \dots \rangle_r \\
& \langle (g\mathbf{Z}_1 g^{-1}) \dots (g\mathbf{Z}_r g^{-1}) \rangle_r = \langle \mathbf{Z}_1 \dots \mathbf{Z}_r \rangle_r, \\
& \langle [\lambda, \mathbf{Z}_1] \mathbf{Z}_2 \dots \mathbf{Z}_r \rangle_r + \dots + \langle \mathbf{Z}_1 \dots \mathbf{Z}_{r-1} [\lambda, \mathbf{Z}_r] \rangle_r \\
& \langle [\mathbf{A}, \mathbf{Z}_1] \mathbf{Z}_2 \dots \mathbf{Z}_r \rangle_r + (-1)^{p_1} \langle \mathbf{Z}_1 [\mathbf{A}, \mathbf{Z}_2] \mathbf{Z}_3 \dots \mathbf{Z}_r \rangle_r + \\
& \quad \dots + (-1)^{p_1 + \dots + p_{r-1}} \langle \mathbf{Z}_1 \dots \mathbf{Z}_{r-1} [\mathbf{A}, \mathbf{Z}_r] \rangle_r \\
& \langle \mathbf{D}(\mathbf{Z}_1 \dots \mathbf{Z}_r) \rangle_r = d\langle \mathbf{Z}_1 \dots \mathbf{Z}_r \rangle_r
\end{aligned}$$

$$\begin{aligned}
\langle D(\mathbf{Z}_1 \cdots \mathbf{Z}_r) \rangle_r &= \langle (D\mathbf{Z}_1)\mathbf{Z}_2 \cdots \mathbf{Z}_r \rangle_r + (-1)^{p_1} \langle \mathbf{Z}_1 (D\mathbf{Z}_2)\mathbf{Z}_3 \cdots \mathbf{Z}_r \rangle_r + \\
&\quad + \cdots + (-1)^{p_1 + \cdots + p_{r-1}} \langle \mathbf{Z}_1 \cdots \mathbf{Z}_{r-1} (D\mathbf{Z}_r) \rangle_r \\
[\mathbf{P}_a, \mathbf{P}_b] &= 0 \\
[J_{ab}, \mathbf{P}_c] &= \eta_{cb} \mathbf{P}_a - \eta_{ca} \mathbf{P}_b \\
[J_{ab}, J_{cd}] &= \eta_{cb} J_{ad} - \eta_{ca} J_{bd} + \eta_{db} J_{ca} - \eta_{da} J_{cb} \\
\langle J_{ab} \mathbf{P}_c \rangle &= \varepsilon_{abc} \\
\langle [A, J_{ab}] \mathbf{P}_c \rangle + \langle J_{ab} [A, \mathbf{P}_c] \rangle &= 0 \\
\mathbf{A} &= e^a \mathbf{P}_a + \frac{1}{2} \omega^{ab} J_{ab} \\
\omega_a^e \varepsilon_{ebc} + \omega_b^e \varepsilon_{aec} + \omega_c^e \varepsilon_{abe} &= 0 \\
D_\omega \varepsilon_{abc} &= 0 \\
\delta_{abcd}^{efgh} &= 0 \\
0 &= \omega^d{}_e \delta_{abcd}^{efgh} \\
&= \omega^d{}_e (\delta_a^e \delta_{bcd}^{fgh} - \delta_b^e \delta_{acd}^{fgh} + \delta_c^e \delta_{abd}^{fgh} - \delta_d^e \delta_{abc}^{fgh}) \\
&= \omega^d{}_a \delta_{bbd}^{fgh} - \omega^d{}_b \delta_{acd}^{fgh} + \omega^d{}_c \delta_{abd}^{fgh} \\
&= \omega^e{}_a \delta_{ebc}^{fgh} + \omega^e{}_b \delta_{aec}^{fgh} + \omega^e{}_c \delta_{abe}^{fgh} \\
\omega_a^e \varepsilon_{ebc} + \omega_b^e \varepsilon_{aec} + \omega_c^e \varepsilon_{abe} &= 0 \\
Q_{\text{CS}}^{(2n+1)} &\equiv (n+1) \int_0^1 dt \langle \mathbf{A}(t d\mathbf{A} + t^2 \mathbf{A}^2)^n \rangle \\
Q_{\text{CS}}^{(3)} &= \left\langle \mathbf{A} d\mathbf{A} + \frac{2}{3} \mathbf{A}^3 \right\rangle \\
Q_{\text{CS}}^{(5)} &= \left\langle \mathbf{A} (d\mathbf{A})^2 + \frac{3}{2} \mathbf{A}^3 d\mathbf{A} + \frac{3}{5} \mathbf{A}^5 \right\rangle \\
dQ_{\text{CS}}^{(2n+1)} &= \langle \mathbf{F}^{n+1} \rangle \\
d\delta_{\text{gauge}} Q_{\text{CS}}^{(2n+1)} &= 0 \\
Q_{\text{CS}}^{(2n+1)}(\mathbf{A}') &= Q_{\text{CS}}^{(2n+1)}(\mathbf{A}) + (-1)^{n+1} \frac{n!(n+1)!}{(2n+1)!} \langle (g^{-1} dg)^{2n+1} \rangle + d\Omega_{\text{fin}}^{(2n)} \\
\delta_{\text{gauge}} Q_{\text{CS}}^{(2n+1)} &= d\Omega^{(2n)} \\
L_{\text{YM}} &= -\frac{1}{4} \langle \mathbf{F} \wedge \star \mathbf{F} \rangle \\
L_{\text{CS}}^{(2n+1)} &= (n+1)k \int_0^1 dt \langle \mathbf{A}(t d\mathbf{A} + t^2 \mathbf{A}^2)^n \rangle \\
[\mathbf{P}_a, \mathbf{P}_b] &= J_{ab} \\
[J_{ab}, \mathbf{P}_c] &= \eta_{cb} \mathbf{P}_a - \eta_{ca} \mathbf{P}_b \\
[J_{ab}, J_{cd}] &= \eta_{cb} J_{ad} - \eta_{ca} J_{bd} + \eta_{db} J_{ca} - \eta_{da} J_{cb} \\
\mathbf{A} &= \frac{1}{\ell} e^a \mathbf{P}_a + \frac{1}{2} \omega^{ab} J_{ab} \\
\boldsymbol{\lambda} &= \frac{1}{\ell} \lambda^a \mathbf{P}_a + \frac{1}{2} \lambda^{ab} J_{ab} \\
\delta e^a &= \lambda^a{}_b e^b - D_\omega \lambda^a \\
\delta \omega^{ab} &= -D_\omega \lambda^{ab} + \frac{1}{\ell^2} (\lambda^a e^b - \lambda^b e^a) \\
\mathbf{F} &= \frac{1}{\ell} T^a \mathbf{P}_a + \frac{1}{2} \left(R^{ab} + \frac{1}{\ell^2} e^a e^b \right) J_{ab} \\
T^a &= D_\omega e^a \\
R^{ab} &= d\omega^{ab} + \omega^a{}_c \omega^{cb}
\end{aligned}$$



$$\begin{aligned}
\langle J_{a_1 a_2} \cdots J_{a_{2n-1} a_{2n}} P_{a_{2n+1}} \rangle &= \frac{2^n}{n+1} \varepsilon_{a_1 \cdots a_{2n+1}} \\
\mathbf{e} &= \frac{1}{\ell} e^a \mathbf{P}_a \\
\boldsymbol{\omega} &= \frac{1}{2} \omega^{ab} J_{ab} \\
\mathbf{T} &= d\mathbf{e} + [\boldsymbol{\omega}, \mathbf{e}] \\
\mathbf{R} &= d\boldsymbol{\omega} + \boldsymbol{\omega}^2 \\
\mathbf{T} &= \frac{1}{\ell} T^a \mathbf{P}_a \\
\mathbf{R} &= \frac{1}{2} R^{ab} J_{ab} \\
L_{\text{CS}}^{(2n+1)} &= (n+1)k \int_0^1 dt \langle (\mathbf{R} + t^2 \mathbf{e}^2)^n \mathbf{e} \rangle + dB_{\text{CS}}^{(2n)} \\
L_{\text{CS}}^{(2n+1)} &= \frac{k}{\ell} \varepsilon_{a_1 \cdots a_{2n+1}} \int_0^1 dt \left(R^{a_1 a_2} + \frac{t^2}{\ell^2} e^{a_1} e^{a_2} \right) \times \cdots \times \\
&\times \left(R^{a_{2n-1} a_{2n}} + \frac{t^2}{\ell^2} e^{a_{2n-1}} e^{a_{2n}} \right) e^{a_{2n+1}} + dB_{\text{CS}}^{(2n)} \\
\mathcal{Q}_{A \leftarrow \bar{A}}^{(2n+1)} &\equiv (n+1) \int_0^1 dt \langle \boldsymbol{\Theta} \mathbf{F}_t^n \rangle \\
\boldsymbol{\Theta} &\equiv \mathbf{A} - \bar{\mathbf{A}} \\
\mathbf{A}_t &\equiv \bar{\mathbf{A}} + t\boldsymbol{\Theta} \\
\mathbf{F}_t &\equiv d\mathbf{A}_t + \mathbf{A}_t^2 \\
\langle \mathbf{F}^{n+1} \rangle - \langle \bar{\mathbf{F}}^{n+1} \rangle &= \int_0^1 dt \frac{d}{dt} \langle \mathbf{F}_t^{n+1} \rangle \\
\langle \mathbf{F}^{n+1} \rangle - \langle \bar{\mathbf{F}}^{n+1} \rangle &= (n+1) \int_0^1 dt \left\langle \mathbf{F}_t^n \frac{d}{dt} \mathbf{F}_t \right\rangle \\
\frac{d}{dt} \mathbf{F}_t &= D_t \boldsymbol{\Theta} \\
\langle \mathbf{F}^{n+1} \rangle - \langle \bar{\mathbf{F}}^{n+1} \rangle &= (n+1) \int_0^1 dt \langle \mathbf{F}_t^n D_t \boldsymbol{\Theta} \rangle \\
\langle \mathbf{F}^{n+1} \rangle - \langle \bar{\mathbf{F}}^{n+1} \rangle &= (n+1) \int_0^1 dt \langle D_t (\mathbf{F}_t^n \boldsymbol{\Theta}) \rangle \\
\langle \mathbf{F}^{n+1} \rangle - \langle \bar{\mathbf{F}}^{n+1} \rangle &= (n+1) d \int_0^1 dt \langle \mathbf{F}_t^n \boldsymbol{\Theta} \rangle \\
\langle \mathbf{F}^{n+1} \rangle - \langle \bar{\mathbf{F}}^{n+1} \rangle &= d\mathcal{Q}_{A \leftarrow \bar{A}}^{(2n+1)} \\
\langle \mathbf{F}^{n+1} \rangle &= d\mathcal{Q}_{A \leftarrow 0}^{(2n+1)} \\
S_{\text{T}}[\mathbf{A}, \bar{\mathbf{A}}] &= k \int_M \mathcal{Q}_{A \leftarrow \bar{A}}^{(2n+1)} \\
\delta_{\text{dif}} \mathbf{A} &= -\mathcal{E}_{\xi} \mathbf{A} \\
\delta_{\text{dif}} \bar{\mathbf{A}} &= -\mathcal{E}_{\xi} \bar{\mathbf{A}} \\
\delta_{\text{gauge}} \mathbf{A} &= -D\lambda \\
\delta_{\text{gauge}} \bar{\mathbf{A}} &= -\bar{D}\lambda \\
S_{\text{T}}^{(2n+1)}[\bar{\mathbf{A}}, \mathbf{A}] &= -S_{\text{T}}^{(2n+1)}[\mathbf{A}, \bar{\mathbf{A}}]
\end{aligned}$$

$$\begin{aligned}
\Theta &\rightarrow -\Theta \\
\mathbf{A}_t &\rightarrow \mathbf{A}_{1-t} \\
\mathbf{F}_t &\rightarrow \mathbf{F}_{1-t} \\
\int_0^1 f(t)dt &= \int_0^1 f(1-t)dt \\
\delta S_T^{(2n+1)} &= (n+1)k \int_M \left(\langle \delta \mathbf{A} \mathbf{F}^n \rangle - \langle \delta \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) + \int_{\partial M} \Xi \\
\Xi &\equiv n(n+1)k \int_0^1 dt \langle \delta \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
\mathcal{Q}_{A \leftarrow \overline{A}}^{(2n+1)} &= (n+1) \int_0^1 dt \langle \Theta \mathbf{F}_t^n \rangle \\
\delta \Theta &= \delta \mathbf{A} - \delta \overline{\mathbf{A}} \\
\delta \mathbf{A}_t &= \delta \overline{\mathbf{A}} + t \delta \Theta \\
\delta \mathbf{F}_t &= D_t \delta \mathbf{A}_t \\
\delta \mathcal{Q}_{A \leftarrow \overline{A}}^{(2n+1)} &= (n+1) \int_0^1 dt \langle \delta \Theta \mathbf{F}_t^n \rangle + n(n+1) \int_0^1 dt \langle \Theta D_t \delta \mathbf{A}_t \mathbf{F}_t^{n-1} \rangle \\
\langle \Theta D_t \delta \mathbf{A}_t \mathbf{F}_t^{n-1} \rangle &= \langle D_t \Theta \delta \mathbf{A}_t \mathbf{F}_t^{n-1} \rangle + d \langle \delta \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
\frac{d}{dt} \mathbf{F}_t &= D_t \Theta \\
\frac{d}{dt} \delta \mathbf{A}_t &= \delta \Theta \\
n \langle D_t \Theta \delta \mathbf{A}_t \mathbf{F}_t^{n-1} \rangle &= \frac{d}{dt} \langle \delta \mathbf{A}_t \mathbf{F}_t^n \rangle - \langle \delta \Theta \mathbf{F}_t^n \rangle \\
n \langle \Theta D_t \delta \mathbf{A}_t \mathbf{F}_t^{n-1} \rangle &= \frac{d}{dt} \langle \delta \mathbf{A}_t \mathbf{F}_t^n \rangle - \langle \delta \Theta \mathbf{F}_t^n \rangle + n d \langle \delta \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
\delta \mathcal{Q}_{A \leftarrow \overline{A}}^{(2n+1)} &= (n+1) \int_0^1 dt \frac{d}{dt} \langle \delta \mathbf{A}_t \mathbf{F}_t^n \rangle + n(n+1) d \int_0^1 dt \langle \delta \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
\delta \mathcal{Q}_{A \leftarrow \overline{A}}^{(2n+1)} &= (n+1) \left(\langle \delta \mathbf{A} \mathbf{F}^n \rangle - \langle \delta \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) + n(n+1) d \int_0^1 dt \langle \delta \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
\langle \mathbf{F}^n \mathbf{G}_A \rangle &= 0 \\
\langle \overline{\mathbf{F}}^n \mathbf{G}_A \rangle &= 0 \\
\int_0^1 dt \langle \delta \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \Big|_{\partial M} & \\
d \star J &= 0 \\
\star J_{\text{gauge}} &= n(n+1)k d \int_0^1 dt \langle \lambda \Theta \mathbf{F}_t^{n-1} \rangle \\
\star J_{\text{dif}} &= n(n+1)k d \int_0^1 dt \langle l_\xi \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
Q_{\text{gauge}}(\lambda) &= n(n+1)k \int_{\partial \Sigma} \int_0^1 dt \langle \lambda \Theta \mathbf{F}_t^{n-1} \rangle \\
Q_{\text{dif}}(\xi) &= n(n+1)k \int_{\partial \Sigma} \int_0^1 dt \langle l_\xi \mathbf{A}_t \Theta \mathbf{F}_t^{n-1} \rangle \\
\delta_\lambda Q_{\text{gauge}}(\eta) &= -Q_{\text{gauge}}([\lambda, \eta]) \\
\delta_\lambda Q_{\text{dif}}(\xi) &= -Q_{\text{gauge}}(\mathcal{L}_\xi \lambda)
\end{aligned}$$



$$\begin{aligned}
\{Q_\eta, Q_\lambda\} &= Q_{[\eta, \lambda]} \\
C(S_T^{(2n+1)}) &= -S_T^{(2n+1)} \\
PT(S_T^{(2n+1)}) &= -S_T^{(2n+1)} \\
CPT(S_T^{(2n+1)}) &= S_T^{(2n+1)} \\
\delta L_T^{(2n+1)} &= (n+1)k \left(\langle \delta \mathbf{A} \mathbf{F}^n \rangle - \langle \delta \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) + d\Xi \\
\Xi &= n(n+1)k \int_0^1 dt \langle \delta \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\delta_{\text{gauge}} \mathbf{A} &= -D\lambda \\
\delta_{\text{gauge}} \overline{\mathbf{A}} &= -\overline{D}\lambda \\
\delta_{\text{gauge}} L_T^{(2n+1)} &= -(n+1)k d \left(\langle \lambda \mathbf{F}^n \rangle - \langle \lambda \overline{\mathbf{F}}^n \rangle \right) + d\Xi_{\text{gauge}} \\
\star J'_{\text{gauge}} &\equiv (n+1)k \left(\langle \lambda \mathbf{F}^n \rangle - \langle \lambda \overline{\mathbf{F}}^n \rangle \right) - \Xi_{\text{gauge}} \\
\Xi_{\text{gauge}} &= -n(n+1)k d \int_0^1 dt \langle \lambda \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + (n+1)k \left(\langle \lambda \mathbf{F}^n \rangle - \langle \lambda \overline{\mathbf{F}}^n \rangle \right) \\
\star J'_{\text{gauge}} &= n(n+1)k d \int_0^1 dt \langle \lambda \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\delta_{\text{dif}} \mathbf{A} &= -\mathcal{E}_\xi \mathbf{A} \\
\delta_{\text{dif}} \overline{\mathbf{A}} &= -\mathcal{E}_\xi \overline{\mathbf{A}} \\
\delta_{\text{dif}} L_T^{(2n+1)} &= -(n+1)k \left(\langle \mathcal{E}_\xi \mathbf{A} \mathbf{F}^n \rangle - \langle \mathcal{E}_\xi \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) + d\Xi_{\text{dif}} \\
\mathcal{E}_\xi \mathbf{A} &= I_\xi \mathbf{F} + DI_\xi \mathbf{A} \\
\langle \mathcal{E}_\xi \mathbf{A} \mathbf{F}^n \rangle &= \langle I_\xi \mathbf{F} \mathbf{F}^n \rangle + \langle DI_\xi \mathbf{A} \mathbf{F}^n \rangle \\
\langle \mathcal{E}_\xi \mathbf{A} \mathbf{F}^n \rangle &= d \langle I_\xi \mathbf{A} \mathbf{F}^n \rangle \\
\delta_{\text{dif}} L_T^{(2n+1)} &= -(n+1)k d \left(\langle I_\xi \mathbf{A} \mathbf{F}^n \rangle - \langle I_\xi \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) + d\Xi_{\text{dif}} \\
\delta_{\text{dif}} L_T^{(2n+1)} &= -\mathcal{E}_\xi L_T^{(2n+1)} \\
&= -dI_\xi L_T^{(2n+1)}, \\
\star J'_{\text{dif}} &\equiv (n+1)k \left(\langle I_\xi \mathbf{A} \mathbf{F}^n \rangle - \langle I_\xi \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) - \Xi_{\text{dif}} - I_\xi L_T^{(2n+1)} \\
\Xi_{\text{dif}} + I_\xi L_T^{(2n+1)} &= -n(n+1)k d \int_0^1 dt \langle I_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + \\
&\quad + (n+1)k \left(\langle I_\xi \mathbf{A} \mathbf{F}^n \rangle - \langle I_\xi \overline{\mathbf{A}} \overline{\mathbf{F}}^n \rangle \right) \\
\star J'_{\text{dif}} &= n(n+1)k d \int_0^1 dt \langle I_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
S_T^{(2n+1)}[\mathbf{A}, \overline{\mathbf{A}}] &= S_{\text{CS}}^{(2n+1)}[\mathbf{A}] - S_{\text{CS}}^{(2n+1)}[\overline{\mathbf{A}}] + \int_{\partial M} \mathcal{B}^{(2n)} \\
-S_{\text{CS}}^{(2n+1)}[\overline{\mathbf{A}}] &= -\int_M L_{\text{CS}}^{(2n+1)}(\overline{\mathbf{A}}) = \int_{-M} L_{\text{CS}}^{(2n+1)}(\overline{\mathbf{A}}) \\
L_T^{(2n+1)}(\mathbf{A}, \overline{\mathbf{A}}) &= (n+1)k \int_0^1 dt \langle \boldsymbol{\Theta} \mathbf{F}_t^n \rangle
\end{aligned}$$

$$\begin{aligned} dQ_{A \leftarrow \bar{A}}^{(2n+1)} + dQ_{\tilde{A} \leftarrow \tilde{A}}^{(2n+1)} + dQ_{\tilde{A} \leftarrow A}^{(2n+1)} &= 0 \\ dQ_{A \leftarrow \bar{A}}^{(2n+1)} + dQ_{\bar{A} \leftarrow \tilde{A}}^{(2n+1)} + dQ_{\tilde{A} \leftarrow A}^{(2n+1)} &= \langle F^{n+1} \rangle - \langle \bar{F}^{n+1} \rangle + \langle \bar{F}^{n+1} \rangle + \\ &\quad - \langle \tilde{F}^{n+1} \rangle + \langle \tilde{F}^{n+1} \rangle - \langle F^{n+1} \rangle \\ &= 0. \end{aligned}$$

$$\begin{aligned} Q_{A \leftarrow \bar{A}}^{(2n+1)} + Q_{\bar{A} \leftarrow \tilde{A}}^{(2n+1)} + Q_{\tilde{A} \leftarrow A}^{(2n+1)} &= dQ_{A \leftarrow \tilde{A} \leftarrow \bar{A}}^{(2n)} \\ Q_{A \leftarrow \bar{A}}^{(2n+1)} &= Q_{A \leftarrow \tilde{A}}^{(2n+1)} + Q_{\tilde{A} \leftarrow \bar{A}}^{(2n+1)} + dQ_{A \leftarrow \tilde{A} \leftarrow \bar{A}}^{(2n)} \\ t^i &\geq 0, i = 0, \dots, r+1 \end{aligned}$$

$$\sum_{i=0}^{r+1} t^i = 1$$

$$A_t = \sum_{i=0}^{r+1} t^i A_i$$

$$\begin{aligned} F_t &= dA_t + A_t^2 \\ T_{r+1} &= (A_0 A_1 \cdots A_{r+1}) \\ d: \Omega^a(M) \times \Omega^b(T_{r+1}) &\rightarrow \Omega^{a+1}(M) \times \Omega^b(T_{r+1}) \\ d_t: \Omega^a(M) \times \Omega^b(T_{r+1}) &\rightarrow \Omega^a(M) \times \Omega^{b+1}(T_{r+1}) \\ l_t: \Omega^a(M) \times \Omega^b(T_{r+1}) &\rightarrow \Omega^{a-1}(M) \times \Omega^{b+1}(T_{r+1}) \\ l_t A_t &= 0 \end{aligned}$$

$$l_t F_t = d_t A_t$$

$$d^2 = 0$$

$$d_t^2 = 0$$

$$[l_t, d] = d_t$$

$$[l_t, d_t] = 0$$

$$\{d, d_t\} = 0$$

$$\int_{\partial T_{r+1}} \frac{l_t^p}{p!} \pi = \int_{T_{r+1}} \frac{l_t^{p+1}}{(p+1)!} d\pi + (-1)^{p+q} d \int_{T_{r+1}} \frac{l_t^{p+1}}{(p+1)!} \pi$$

$$\pi = \sum_p \alpha_p \left\langle A_t^{a_p} F_t^{b_p} (d_t A_t)^{c_p} (d_t F_t)^{d_p} \right\rangle$$

$$a_p + 2b_p + c_p + 2d_p = m$$

$$c_p + d_p = q$$

$$(p+1)d_t l_t^p \pi = l_t^{p+1} d\pi - d l_t^{p+1} \pi$$

$$[l_t^{p+1}, d] = (p+1)d_t l_t^p$$

$$[l_t^2, d] = l_t [l_t, d] + [l_t, d] l_t$$

$$= l_t d_t + d_t l_t$$

$$= 2 d_t l_t.$$

$$[l_t^{k+2}, d] = l_t [l_t^{k+1}, d] + [l_t, d] l_t^{k+1}$$

$$= (k+1) l_t d_t l_t^k + d_t l_t^{k+1}$$

$$= (k+2) d_t l_t^{k+1}$$

$$(p+1) \int_{\partial T_{r+1}} l_t^p \pi = \int_{T_{r+1}} l_t^{p+1} d\pi - \int_{T_{r+1}} d l_t^{p+1}$$

$$d \int_{T_s} \alpha = (-1)^s \int_{T_s} d\alpha$$



$$\begin{aligned}
(p+1) \int_{\partial T_{r+1}} l_t^p \pi &= \int_{T_{r+1}} l_t^{p+1} d\pi + (-1)^{p+q} d \int_{T_{r+1}} l_t^{p+1} \pi \\
\pi &= \langle \mathbf{F}_t^{n+1} \rangle \\
\int_{\partial T_{p+1}} \frac{l_t^p}{p!} \langle \mathbf{F}_t^{n+1} \rangle &= (-1)^p d \int_{T_{p+1}} \frac{l_t^{p+1}}{(p+1)!} \langle \mathbf{F}_t^{n+1} \rangle \\
\int_{\partial T_1} \langle \mathbf{F}_t^{n+1} \rangle &= d \int_{T_1} l_t \langle \mathbf{F}_t^{n+1} \rangle \\
\mathbf{A}_t &= t^0 \mathbf{A}_0 + t^1 \mathbf{A}_1 \\
\partial T_1 &= (\mathbf{A}_1) - (\mathbf{A}_0) \\
\int_{\partial T_1} \langle \mathbf{F}_t^{n+1} \rangle &= \langle \mathbf{F}_1^{n+1} \rangle - \langle \mathbf{F}_0^{n+1} \rangle \\
l_t \langle \mathbf{F}_t^{n+1} \rangle &= (n+1) \langle (l_t \mathbf{F}_t) \mathbf{F}_t^n \rangle \\
l_t \mathbf{F}_t &= d_t \mathbf{A}_t \\
&= dt^0 \mathbf{A}_0 + dt^1 \mathbf{A}_1 \\
&= dt^1 (\mathbf{A}_1 - \mathbf{A}_0). \\
\langle \mathbf{F}_1^{n+1} \rangle - \langle \mathbf{F}_0^{n+1} \rangle &= (n+1) d \int_{T_1} dt^1 \langle (\mathbf{A}_1 - \mathbf{A}_0) \mathbf{F}_t^n \rangle \\
\langle \mathbf{F}_1^{n+1} \rangle - \langle \mathbf{F}_0^{n+1} \rangle &= d \mathcal{Q}_{A_1 \leftarrow A_0}^{(2n+1)} \\
\mathcal{Q}_{A_1 \leftarrow A_0}^{(2n+1)} &= \int_{(A_0 A_1)} l_t \langle \mathbf{F}_t^{n+1} \rangle \\
&= (n+1) \int_0^1 dt^1 \langle (\mathbf{A}_1 - \mathbf{A}_0) \mathbf{F}_t^n \rangle. \\
\int_{\partial T_2} l_t \langle \mathbf{F}_t^{n+1} \rangle &= -d \int_{T_2} \frac{l_t^2}{2} \langle \mathbf{F}_t^{n+1} \rangle \\
\mathbf{A}_t &= t^0 \mathbf{A}_0 + t^1 \mathbf{A}_1 + t^2 \mathbf{A}_2 \\
\partial T_2 &= (\mathbf{A}_1 \mathbf{A}_2) - (\mathbf{A}_0 \mathbf{A}_2) + (\mathbf{A}_0 \mathbf{A}_1) \\
\int_{\partial T_2} l_t \langle \mathbf{F}_t^{n+1} \rangle &= \int_{(A_1 A_2)} l_t \langle \mathbf{F}_t^{n+1} \rangle - \int_{(A_0 A_2)} l_t \langle \mathbf{F}_t^{n+1} \rangle + \int_{(A_0 A_1)} l_t \langle \mathbf{F}_t^{n+1} \rangle \\
\int_{\partial T_2} l_t \langle \mathbf{F}_t^{n+1} \rangle &= \mathcal{Q}_{A_2 \leftarrow A_1}^{(2n+1)} - \mathcal{Q}_{A_2 \leftarrow A_0}^{(2n+1)} + \mathcal{Q}_{A_1 \leftarrow A_0}^{(2n+1)} \\
l_t^2 \langle \mathbf{F}_t^{n+1} \rangle &= n(n+1) \langle (d_t \mathbf{A}_t)^2 \mathbf{F}_t^{n-1} \rangle \\
\int_{T_2} \frac{l_t^2}{2} \langle \mathbf{F}_t^{n+1} \rangle &= \mathcal{Q}_{A_2 \leftarrow A_1 \leftarrow A_0}^{(2n)} \\
\mathcal{Q}_{A_2 \leftarrow A_1 \leftarrow A_0}^{(2n)} &\equiv n(n+1) \int_0^1 dt \int_0^t ds \langle (\mathbf{A}_2 - \mathbf{A}_1)(\mathbf{A}_1 - \mathbf{A}_0) \mathbf{F}_t^{n-1} \rangle \\
\mathbf{A}_t &= \mathbf{A}_0 + s(\mathbf{A}_2 - \mathbf{A}_1) + t(\mathbf{A}_1 - \mathbf{A}_0) \\
\mathcal{Q}_{A_2 \leftarrow A_1}^{(2n+1)} - \mathcal{Q}_{A_2 \leftarrow A_0}^{(2n+1)} + \mathcal{Q}_{A_1 \leftarrow A_0}^{(2n+1)} &= -d \mathcal{Q}_{A_2 \leftarrow A_1 \leftarrow A_0}^{(2n)} \\
\mathcal{Q}_{A_2 \leftarrow A_0}^{(2n+1)} &= \mathcal{Q}_{A_2 \leftarrow A_1}^{(2n+1)} + \mathcal{Q}_{A_1 \leftarrow A_0}^{(2n+1)} + d \mathcal{Q}_{A_2 \leftarrow A_1 \leftarrow A_0}^{(2n)} \\
L_T^{(2n+1)}(\mathbf{A}, \bar{\mathbf{A}}) &= k \mathcal{Q}_{A \leftarrow \bar{A}}^{(2n+1)} \\
\mathcal{Q}_{A \leftarrow \bar{A}}^{(2n+1)} &= \mathcal{Q}_{A \leftarrow \tilde{A}}^{(2n+1)} + \mathcal{Q}_{\tilde{A} \leftarrow \bar{A}}^{(2n+1)} + d \mathcal{Q}_{A \leftarrow \tilde{A} \leftarrow \bar{A}}^{(2n)} \\
\mathcal{Q}_{A \leftarrow \tilde{A} \leftarrow \bar{A}}^{(2n)} &\equiv n(n+1) \int_0^1 dt \int_0^t ds \langle (\mathbf{A} - \tilde{\mathbf{A}})(\tilde{\mathbf{A}} - \bar{\mathbf{A}}) \mathbf{F}_{st}^{n-1} \rangle \\
\mathbf{A}_{st} &= \bar{\mathbf{A}} + s(\mathbf{A} - \tilde{\mathbf{A}}) + t(\tilde{\mathbf{A}} - \bar{\mathbf{A}})
\end{aligned}$$

$$\begin{aligned}
A &= a_0 + a_1, \\
\bar{A} &= \bar{a}_0 + \bar{a}_1 \\
Q_{a_0+a_1 \leftarrow \bar{a}_0+\bar{a}_1}^{(2+1)} &= Q_{a_0+a_1 \leftarrow a_0}^{(2+1)} + Q_{a_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n+1)} + dQ_{a_0+a_1 \leftarrow a_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n)} \\
Q_{a_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n+1)} &= Q_{a_0 \leftarrow \bar{a}_0}^{(2n+1)} + Q_{\bar{a}_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n+1)} + dQ_{a_0 \leftarrow \bar{a}_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n)} \\
Q_{a_0+a_1 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n+1)} &= Q_{a_0+a_1 \leftarrow a_0}^{(2n+1)} + Q_{\bar{a}_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n+1)} + Q_{a_0 \leftarrow \bar{a}_0}^{(2n+1)} + \\
&\quad + dQ_{a_0 \leftarrow \bar{a}_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n)} + dQ_{a_0+a_1 \leftarrow a_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n)} \\
Q_{a_0+a_1 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n+1)} &= Q_{a_0+a_1 \leftarrow a_0}^{(2n+1)} - Q_{\bar{a}_0+\bar{a}_1 \leftarrow \bar{a}_0}^{(2n+1)} + Q_{a_0 \leftarrow \bar{a}_0}^{(2n+1)} + \\
&\quad + dQ_{a_0 \leftarrow \bar{a}_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n)} + dQ_{a_0+a_1 \leftarrow a_0 \leftarrow \bar{a}_0+\bar{a}_1}^{(2n)} \\
Q_{A \leftarrow \bar{A}}^{(2n+1)} &= Q_{A \leftarrow \bar{A}}^{(2n+1)} + Q_{\bar{A} \leftarrow \bar{A}}^{(2n+1)} + dQ_{A \leftarrow \bar{A} \leftarrow \bar{A}}^{(2n)} \\
Q_{CS}^{(2n+1)}(A) &= Q_{A \leftarrow 0}^{(2n+1)} \\
Q_{A \leftarrow \bar{A}}^{(2n+1)} &= Q_{CS}^{(2n+1)}(A) - Q_{CS}^{(2n+1)}(\bar{A}) + d\mathcal{B}^{(2n)} \\
\mathcal{B}^{(2n)} &= -n(n+1) \int_0^1 dt \int_0^t ds \langle A \bar{A} F_{st}^{n-1} \rangle \\
A_{st} &= sA + (1-t)\bar{A} \\
F_{st} &= \bar{F} + \bar{D}(sA - t\bar{A}) + (sA - t\bar{A})^2 \\
[P_a, P_b] &= J_{ab} \\
[J_{ab}, P_c] &= \eta_{cb} P_a - \eta_{ca} P_b \\
[J_{ab}, J_{cd}] &= \eta_{cb} J_{ad} - \eta_{ca} J_{bd} + \eta_{db} J_{ca} - \eta_{da} J_{cb} \\
\langle J_{a_1 a_2} \cdots J_{a_{2n-1} a_{2n}} P_{a_{2n+1}} \rangle &= \frac{2^n}{n+1} \varepsilon_{a_1 \cdots a_{2n+1}} \\
L_G^{(2n+1)} &= k Q_{A \leftarrow \bar{A}}^{(2n+1)} \\
\bar{A} &= \bar{\omega} \\
A &= e + \omega \\
\bar{F} &= \bar{R} \\
F &= R + e^2 + T \\
R &= d\omega + \omega^2 \\
T &= de + [\omega, e] \\
\tilde{A} &= \omega \\
Q_{e+\omega \leftarrow \bar{\omega}}^{(2n+1)} &= Q_{e+\omega \leftarrow \omega}^{(2n+1)} + Q_{\omega \leftarrow \bar{\omega}}^{(2n+1)} + dQ_{e+\omega \leftarrow \omega \leftarrow \bar{\omega}}^{(2n)} \\
Q_{e+\omega \leftarrow \omega}^{(2n+1)} &= (n+1) \int_0^1 dt \langle e F_t^n \rangle \\
A_t &= \omega + te \\
F_t &= R + t^2 e^2 + tT \\
Q_{e+\omega \leftarrow \omega}^{(2n+1)} &= (n+1) \int_0^1 dt \langle e (R + t^2 e^2)^n \rangle \\
Q_{\omega \leftarrow \bar{\omega}}^{(2n+1)} &= 0 \\
Q_{e+\omega \leftarrow \omega \leftarrow \bar{\omega}}^{(2n)} &= n(n+1) \int_0^1 dt \int_0^t ds \langle e \theta F_{st}^{n-1} \rangle \\
\theta &\equiv \omega - \bar{\omega} \\
A_{st} &= \bar{\omega} + se + t\theta \\
F_{st} &= \bar{R} + D_{\bar{\omega}}(se + t\theta) + s^2 e^2 + st[e, \theta] + t^2 \theta^2
\end{aligned}$$

$$\begin{aligned}
\mathcal{Q}_{\mathbf{e}+\boldsymbol{\omega}\leftarrow\boldsymbol{\omega}\leftarrow\bar{\boldsymbol{\omega}}}^{(2n)} &= n(n+1)\int_0^1 dt\int_0^t ds\left\langle\mathbf{e}\boldsymbol{\theta}(\bar{\mathbf{R}}+t\mathbf{D}_{\bar{\boldsymbol{\omega}}}\boldsymbol{\theta}+s^2\mathbf{e}^2+t^2\boldsymbol{\theta}^2)^{n-1}\right\rangle \\
L_{\mathbf{G}}^{(2n+1)} &= (n+1)k\int_0^1 dt\langle\mathbf{e}(\mathbf{R}+t^2\mathbf{e}^2)^n\rangle+ \\
&+n(n+1)k\int_0^1 dt\int_0^t ds\left\langle\mathbf{e}\boldsymbol{\theta}(\bar{\mathbf{R}}+t\mathbf{D}_{\bar{\boldsymbol{\omega}}}\boldsymbol{\theta}+s^2\mathbf{e}^2+t^2\boldsymbol{\theta}^2)^{n-1}\right\rangle \\
&\langle\mathbf{J}_{ab}(\mathbf{R}+\mathbf{e}^2)^{n-1}\mathbf{T}\rangle=0 \\
&\langle\mathbf{P}_a(\mathbf{R}+\mathbf{e}^2)^n\rangle=0 \\
&\mathcal{R}_{abc}\equiv\langle\mathbf{F}^{n-1}\mathbf{J}_{ab}\mathbf{P}_c\rangle \\
&\mathcal{R}_{abc}T^c=0 \\
&\mathcal{R}_{abc}\left(R^{ab}+\frac{1}{\ell^2}e^ae^b\right)=0 \\
\mathcal{R}_{abc} &= \frac{2}{n+1}\varepsilon_{abca_1\cdots a_{2n-2}}\left(R^{a_1a_2}+\frac{1}{\ell^2}e^{a_1a_2}\right)\cdots \\
&\cdots\left(R^{a_{2n-3}a_{2n-2}}+\frac{1}{\ell^2}e^{a_{2n-3}a_{2n-2}}\right) \\
&\int_0^1 dt\langle(\delta\bar{\boldsymbol{\omega}}+t\delta\boldsymbol{\theta}+t\delta\mathbf{e})(\boldsymbol{\theta}+\mathbf{e})\mathbf{F}_t^{n-1}\rangle\Big|_{\partial M} \\
&\mathbf{A}_t=\bar{\boldsymbol{\omega}}+t(\mathbf{e}+\boldsymbol{\theta}) \\
&\mathbf{F}_t=\bar{\mathbf{R}}+t\mathbf{D}_{\bar{\boldsymbol{\omega}}}(\mathbf{e}+\boldsymbol{\theta})+t^2(\mathbf{e}^2+[\mathbf{e},\boldsymbol{\theta}]+\boldsymbol{\theta}^2) \\
&\delta\bar{\boldsymbol{\omega}}|_{\partial M}=0 \\
&\int_0^1 dt\left\langle t(\delta\boldsymbol{\theta}\mathbf{e}-\boldsymbol{\theta}\delta\mathbf{e})(\bar{\mathbf{R}}+t^2\mathbf{e}^2+t^2\boldsymbol{\theta}^2)^{n-1}\right\rangle\Big|_{\partial M} \\
&\delta\theta^{[ab}e^{c]}=\theta^{[ab}\delta e^{c]} \\
&\bar{\boldsymbol{\omega}}\rightarrow\bar{\boldsymbol{\omega}}+\bar{\mathbf{e}}_{\mathbf{g}} \\
&\mathcal{Q}_{\boldsymbol{\omega}\leftarrow\bar{\boldsymbol{\omega}}}^{(2n+1)}=0\rightarrow\mathcal{Q}_{\boldsymbol{\omega}+\mathbf{e}_g\leftarrow\bar{\boldsymbol{\omega}}+\bar{\mathbf{e}}_g}^{(2n+1)}\neq 0 \\
&\lambda_{\alpha_1}\cdots\lambda_{\alpha_n}=\lambda_{\gamma(\alpha_1,\dots,\alpha_n)} \\
&K_{\alpha_1\cdots\alpha_n}^\rho=\begin{cases} 1, & \text{cuando } \rho=\gamma(\alpha_1,\dots,\alpha_n) \\ 0, & \infty \end{cases} \\
&K_{\alpha_1\cdots\alpha_n}^\rho=K_{\alpha_1\cdots\alpha_{n-1}}^\sigma K_{\sigma\alpha_n}^\rho=K_{\alpha_1\sigma}^\rho K_{\alpha_2\cdots\alpha_n}^\sigma \\
&(\lambda_\alpha)_\mu{}^\nu=K_{\mu\alpha}{}^\nu \\
&(\lambda_\alpha)_\mu{}^\sigma(\lambda_\beta)_\sigma{}^\nu=K_{\alpha\beta}{}^\sigma(\lambda_\sigma)_\mu{}^\nu=(\lambda_{\gamma(\alpha,\beta)})_\mu{}^\nu \\
&S_p\times S_q=\left\{\lambda_\gamma\mid\lambda_\gamma=\lambda_{\alpha_p}\lambda_{\alpha_q},\text{ con }\lambda_{\alpha_p}\in S_p,\lambda_{\alpha_q}\in S_q\right\} \\
&0_S\lambda_\alpha=\lambda_\alpha 0_S=0_S \\
&[\mathbf{T}_{a_0},\mathbf{T}_{b_0}]=C_{a_0b_0}{}^{c_0}\mathbf{T}_{c_0}+C_{a_0b_0}{}^{c_1}\mathbf{T}_{c_1} \\
&[\mathbf{T}_{a_0},\mathbf{T}_{b_1}]=C_{a_0b_1}{}^c\mathbf{T}_{c_1} \\
&[\mathbf{T}_{a_1},\mathbf{T}_{b_1}]=C_{a_1b_1}{}^c\mathbf{T}_{c_0}+C_{a_1b_1}{}^{c_1}\mathbf{T}_{c_1} \\
&C_{(A,\alpha)(B,\beta)}{}^{(C,\gamma)}=K_{\alpha\beta}{}^\gamma C_{AB}{}^C \\
&[\mathbf{T}_{(A,\alpha)},\mathbf{T}_{(B,\beta)}]\equiv\lambda_\alpha\lambda_\beta[\mathbf{T}_A,\mathbf{T}_B] \\
&=\lambda_{\gamma(\alpha,\beta)}C_{AB}{}^CT_C \\
&=C_{AB}^C\mathbf{T}_{(C,\gamma(\alpha,\beta))} \\
&K_{\alpha\beta}^\rho=\begin{cases} 1, & \text{cuando } \rho=\gamma(\alpha,\beta) \\ 0, & \infty \end{cases}
\end{aligned}$$

$$\begin{aligned}
[T_{(A,\alpha)}, T_{(B,\beta)}] &= K_{\alpha\beta}{}^\rho C_{AB}{}^C T_{(C,\rho)} \\
C_{(A,\alpha)(B,\beta)}{}^{(C,\gamma)} &= K_{\alpha\beta}{}^\gamma C_{AB}{}^C \\
C_{(A,\alpha)(B,\beta)}{}^{(C,\gamma)} &= -(-1)^{q(A)q(B)} C_{(B,\beta)(A,\alpha)}{}^{(C,\gamma)} \\
K_{i,N+1}^j &= K_{N+1,i}^j = 0 \\
K_{i,N+1}{}^{N+1} &= K_{N+1,i}^{N+1} = 1 \\
K_{N+1,N+1}^j &= 0 \\
K_{N+1,N+1}{}^{N+1} &= 1 \\
[T_{(A,i)}, T_{(B,j)}] &= K_{ij}{}^k C_{AB}{}^C T_{(C,k)} + K_{ij}{}^{N+1} C_{AB}{}^C T_{(C,N+1)} \\
[T_{(A,N+1)}, T_{(B,j)}] &= C_{AB}{}^C T_{(C,N+1)} \\
[T_{(A,N+1)}, T_{(B,N+1)}] &= C_{AB}{}^C T_{(C,N+1)} \\
[T_{(A,i)}, T_{(B,j)}] &= C_{(A,i)(B,j)}{}^{(C,k)} T_{(C,k)} \\
C_{(A,i)(B,j)}{}^{(C,k)} &= K_{ij}{}^k C_{AB}{}^C \\
T_{(A,N+1)} &= 0_S T_A = \mathbf{0} \\
C_{(A,i)(B,j)}{}^{(C,k)} &= \begin{cases} 0, & \text{cuando } i+j \neq k \\ C_{AB}^C, & \text{cuando } i+j = k \end{cases} \\
S_E^{(N)} &= \{\lambda_\alpha, \alpha = 0, \dots, N, N+1\} \\
\lambda_\alpha \lambda_\beta &= \begin{cases} \lambda_{\alpha+\beta}, & \text{cuando } \alpha+\beta \leq N \\ \lambda_{N+1}, & \text{cuando } \alpha+\beta \geq N+1 \end{cases} \\
K_{\alpha\beta}^\gamma &= \begin{cases} \delta_{\alpha+\beta}^\gamma, & \text{cuando } \alpha+\beta \leq N \\ \delta_{N+1}^\gamma, & \text{cuando } \alpha+\beta \geq N+1 \end{cases} \\
C_{(A,\alpha)(B,\beta)}{}^{(C,\gamma)} &= \begin{cases} \delta_{\alpha+\beta}^\gamma C_{AB}{}^C, & \text{cuando } \alpha+\beta \leq N \\ \delta_{N+1}^\gamma C_{AB}^C, & \text{cuando } \alpha+\beta \geq N+1 \end{cases} \\
C_{(A,i)(B,j)}{}^{(C,k)} &= \delta_{i+j}^k C_{AB}{}^C \\
\lambda^\alpha \lambda^\beta &= \lambda^{\alpha+\beta} \\
\lambda^\alpha &= 0 \text{ cuando } \alpha > N \\
[V_p, V_q] &\subset \bigoplus_{r \in i(p,q)} V_r \\
S_p \times S_q &\subset \bigcap_{r \in i(p,q)} S_r \\
W_p &\equiv S_p \otimes V_p \\
\mathfrak{G}_R &\equiv \bigoplus_{p \in I} W_p \\
[W_p, W_q] &\subset (S_p \times S_q) \otimes [V_p, V_q] \\
&\subset \bigcap_{s \in i(p,q)} S_s \otimes \bigoplus_{r \in i(p,q)} V_r \\
&\subset \bigoplus_{r \in i(p,q)} \left(\bigcap_{s \in i(p,q)} S_s \right) \otimes V_r. \\
&\bigcap_{s \in i(p,q)} S_s \subset S_r
\end{aligned}$$

$$\begin{aligned}
[W_p, W_q] &\subset \bigoplus_{r \in i(p,q)} S_r \otimes V_r \\
[W_p, W_q] &\subset \bigoplus_{r \in i(p,q)} W_r \\
C_{(a_p, \alpha_p)(b_q, \beta_q)}^{(c_r, \gamma_r)} &= K_{\alpha_p \beta_q}^{\gamma_r} C_{a_p b_q}^{c_r} \\
[V_0, V_0] &\subset V_0, \\
[V_0, V_1] &\subset V_1, \\
[V_1, V_1] &\subset V_0. \\
S_0 &= \left\{ \lambda_{2m}, \text{ con } m = 0, \dots, \left\lfloor \frac{N}{2} \right\rfloor \right\} \cup \{\lambda_{N+1}\}, \\
S_1 &= \left\{ \lambda_{2m+1}, \text{ con } m = 0, \dots, \left\lfloor \frac{N-1}{2} \right\rfloor \right\} \cup \{\lambda_{N+1}\} \\
S_0 \times S_0 &\subset S_0, \\
S_0 \times S_1 &\subset S_1, \\
S_1 \times S_1 &\subset S_0 \\
\mathfrak{G}_R &= W_0 \oplus W_1 \\
W_0 &= S_0 \otimes V_0 \\
W_1 &= S_1 \otimes V_1 \\
C_{(a_p, \alpha_p)(b_q, \beta_q)}^{(c_r, \gamma_r)} &= \begin{cases} \delta_{\alpha_p + \beta_q}^{\gamma_r} C_{a_p b_q}^{c_r}, & \text{cuando } \alpha_p + \beta_q \leq N \\ \delta_{N+1}^{\gamma_r} C_{a_p b_q}^{c_r}, & \text{cuando } \alpha_p + \beta_q \geq N+1 \end{cases} \\
C_{(a_p, \alpha_p)(b_q, \beta_q)}^{(c_r, \gamma_r)} &= \begin{cases} \delta_{\alpha_p + \beta_q}^{\gamma_r} C_{a_p b_q}^{c_r}, & \text{cuando } \alpha_p + \beta_q \leq N \\ \delta_{N+1}^{\gamma_r} C_{a_p b_q}^{c_r}, & \text{cuando } \alpha_p + \beta_q \geq N+1 \end{cases} \\
N_0 &= 2 \left\lfloor \frac{N}{2} \right\rfloor \\
N_1 &= 2 \left\lfloor \frac{N-1}{2} \right\rfloor + 1 \\
S_0 &= \{\lambda_0, \lambda_2, \lambda_4\}, \\
S_1 &= \{\lambda_1, \lambda_3, \lambda_4\} \\
[V_0, V_0] &\subset V_0, \\
[V_0, V_1] &\subset V_1, \\
[V_0, V_2] &\subset V_2, \\
[V_1, V_1] &\subset V_0 \oplus V_2, \\
[V_1, V_2] &\subset V_1, \\
[V_2, V_2] &\subset V_0 \oplus V_2 \\
S_p &= \left\{ \lambda_{2m+p}, \text{ con } m = 0, \dots, \left\lfloor \frac{N-p}{2} \right\rfloor \right\} \cup \{\lambda_{N+1}\}, p = 0, 1, 2 \\
S_0 \times S_0 &\subset S_0, \\
S_0 \times S_1 &\subset S_1, \\
S_0 \times S_2 &\subset S_2, \\
S_1 \times S_1 &\subset S_0 \cap S_2, \\
S_1 \times S_2 &\subset S_1, \\
S_2 \times S_2 &\subset S_0 \cap S_2 \\
C_{(a_p, \alpha_p)(b_q, \beta_q)}^{(c_r, \gamma_r)} &= \begin{cases} \delta_{\alpha_p + \beta_q}^{\gamma_q} C_{a_p b_q}^{c_r}, & \text{cuando } \alpha_p + \beta_q \leq N \\ \delta_{N+1}^{\gamma_r} C_{a_p b_q}^{c_r}, & \text{cuando } \alpha_p + \beta_q \geq N+1 \end{cases}
\end{aligned}$$

$$C_{(a_p, \alpha_p)(b_q, \beta_q)}^{(c_r, \gamma_r)} = \begin{cases} \delta_{\alpha_p + \beta_q}^{\gamma_r} C_{a_p b_q b_r}^{c_r}, & \text{cuando } \alpha_p + \beta_q \leq N \\ \delta_{N+1}^{\gamma_r} C_{a_p b_q}, & \text{cuando } \alpha_p + \beta_q \geq N + 1 \end{cases}$$

$$N_p = 2 \left\lfloor \frac{N-p}{2} \right\rfloor + p, p = 0, 1, 2$$

$$[S_{\kappa\lambda}, S_{\mu\nu}] = -i(\delta_{\mu\lambda} S_{\kappa\nu} - \delta_{\mu\kappa} S_{\lambda\nu} + \delta_{\nu\lambda} S_{\mu\kappa} - \delta_{\nu\kappa} S_{\mu\lambda})$$

	a	b
a	a	b
b	b	a

$$[G_i, G_j] = i\varepsilon_{ijk} G_k,$$

$$J_i = a G_i,$$

$$K_i = b G_i,$$

$$[J_i, J_j] = i\varepsilon_{ijk} J_k,$$

$$[J_i, K_j] = i\varepsilon_{ijk} K_k,$$

$$[K_i, K_j] = i\varepsilon_{ijk} J_k$$

$$M_{ij} = \varepsilon_{ijk} J_k,$$

$$M_{i4} = K_i,$$

$$[M_{\kappa\lambda}, M_{\mu\nu}] = -i(\delta_{\mu\lambda} M_{\kappa\nu} - \delta_{\mu\kappa} M_{\lambda\nu} + \delta_{\nu\lambda} M_{\mu\kappa} - \delta_{\nu\kappa} M_{\mu\lambda})$$

$$[P_a, P_b] = J_{ab},$$

$$[J^{ab}, P_c] = \delta_{ec}^{ab} P^e,$$

$$[P_a, Z_{b_1 \cdots b_5}] = -\frac{1}{5!} \varepsilon_{ab_1 \cdots b_5 c_1 \cdots c_5} Z^{c_1 \cdots c_5},$$

$$[J^{ab}, J_{cd}] = \delta_{ecd}^{abf} J^e{}_f,$$

$$[J^{ab}, Z_{c_1 \cdots c_5}] = \frac{1}{4!} \delta_{dc_1 \cdots c_5}^{abe_1 \cdots e_4} Z^d{}_{e_1 \cdots e_4},$$

$$[Z^{a_1 \cdots a_5}, Z_{b_1 \cdots b_5}] = \eta^{[a_1 \cdots a_5][c_1 \cdots c_5]} \varepsilon_{c_1 \cdots c_5 b_1 \cdots b_5 e} P^e + \delta_{db_1 \cdots b_5}^{a_1 \cdots a_5 e} J^d{}_e$$

$$-\frac{1}{3!3!5!} \varepsilon_{c_1 \cdots c_{11}} \delta_{d_1 d_2 d_3 b_1 \cdots b_5}^{a_1 \cdots a_5 c_4 c_5 c_6} [c_1 c_2 c_3][d_1 d_2 d_3] Z^{c_7 \cdots c_{11}},$$

$$[P_a, Q] = -\frac{1}{2} \Gamma_a Q,$$

$$[J_{ab}, Q] = -\frac{1}{2} \Gamma_{ab} Q,$$

$$[Z_{abcde}, Q] = -\frac{1}{2} \Gamma_{abcde} Q,$$

$$\{Q, \overline{Q}\} = \frac{1}{8} \left(\Gamma^a P_a - \frac{1}{2} \Gamma^{ab} J_{ab} + \frac{1}{5!} \Gamma^{abcde} Z_{abcde} \right)$$

$$\text{Superspacios de } \mathfrak{G}_R \quad \text{Generadores}$$

$$J_{ab} = \lambda_0 J_{ab}^{((0.5)}$$

$$S_0 \otimes V_0 \quad Z_{ab} = \lambda_2 J_{ab}^{(0.sp)}$$

$$\mathbf{0} = \lambda_3 J_{ab}^{(osp)}$$



$$\begin{array}{ll}
S_1 \otimes V_1 & \begin{aligned} Q &= \lambda_1 Q^{(\text{osp})} \\ \mathbf{0} &= \lambda_3 Q^{(\text{osp})} \\ P_a &= \lambda_2 P_a^{(\text{osp})} \end{aligned} \\
S_2 \otimes V_2 & \begin{aligned} Z_{abcde} &= \lambda_2 Z_{abcde}^{(\text{osp})} \\ \mathbf{0} &= \lambda_3 P_a^{(\text{osp})} \\ 0 &= \lambda_3 Z_{abcde}^{(\text{osp})} \end{aligned}
\end{array}$$

$$\begin{aligned}
[J^{ab}, J_{cd}] &= \delta_{ecd}^{abf} J_f^e, \\
[J^{ab}, P_c] &= \delta_{ec}^{ab} P^e, \\
[J^{ab}, Z_{cd}] &= \delta_{ecd}^{aff} Z_f^e, \\
[J^{ab}, Z_{c_1 \dots c_5}] &= \frac{1}{4!} a_{dc_1 \dots c_5}^{abe_4} Z_{e_1 \dots e_4}^d, \\
[J_{ab}, Q] &= -\frac{1}{2} \Gamma_{ab} Q, \\
[P_a, P_b] &= \mathbf{0}, \\
[P_a, Z_{bc}] &= \mathbf{0}, \\
[P_a, Z_{b_1 \dots b_5}] &= \mathbf{0}, \\
[Z_{ab}, Z_{cd}] &= \mathbf{0}, \\
[Z_{ab}, Z_{c_1 \dots c_5}] &= \mathbf{0}, \\
[Z_{a_1 \dots a_5}, Z_{b_1 \dots b_5}] &= \mathbf{0}, \\
[P_a, Q] &= \mathbf{0}, \\
[Z_{ab}, Q] &= \mathbf{0}, \\
[Z_{abcde}, Q] &= \mathbf{0}, \\
\{Q, \bar{Q}\} &= \frac{1}{8} \left(\Gamma^a P_a - \frac{1}{2} \Gamma^{ab} Z_{ab} + \frac{1}{5!} \Gamma^{abcde} Z_{abcde} \right) \\
[P_a, Q] &= -\frac{1}{2} \Gamma_a Q', \\
[Z_{ab}, Q] &= -\frac{1}{2} \Gamma_{ab} Q', \\
[Z_{abcde}, Q] &= -\frac{1}{2} \Gamma_{abcde} Q'
\end{aligned}$$

Superspacios de $\mathfrak{G}_R$ Generadores

$$\begin{array}{ll}
S_0 \otimes V_0 & \begin{aligned} J_{ab} &= \lambda_0 J_{ab}^{(\text{osp})} \\ Z_{ab} &= \lambda_2 J_{ab}^{(\text{osp})} \\ \mathbf{0} &= \lambda_4 J_{ab}^{(\text{osp})} \end{aligned} \\
S_1 \otimes V_1 & \begin{aligned} Q &= \lambda_1 Q^{(\text{osp})} \\ Q' &= \lambda_3 Q^{(\text{osp})} \end{aligned}
\end{array}$$



$$\begin{aligned}
S_2 \otimes V_2 \quad & \mathbf{0} = \lambda_4 \mathbf{Q}^{((s,p))} \\
& \mathbf{P}_a = \lambda_2 \mathbf{P}_a^{(osp)} \\
& \mathbf{Z}_{abcde} = \lambda_2 \mathbf{Z}_{abcde}^{(osp)} \\
& \mathbf{0} = \lambda_4 \mathbf{P}_a^{(osp)} \\
& \mathbf{0} = \lambda_4 \mathbf{Z}_{abcde}^{(osp)}
\end{aligned}$$

$$\begin{aligned}
& \lambda_\alpha \lambda_\beta = \lambda_{(\alpha+\beta) \bmod 4} \\
S_0 &= \{\lambda_0, \lambda_2\}, \\
S_1 &= \{\lambda_1, \lambda_3\}, \\
S_2 &= \{\lambda_0, \lambda_2\}. \\
& S_E^{(2)} \qquad \qquad \mathbb{Z}_4
\end{aligned}$$

0	1	2	3	0	1	2	3
1	2	3	3	1	2	3	0
2	3	3	3	2	3	0	1
3	3	3	3	3	0	1	2

Subespacios de $\mathfrak{G}_R$ Generadores

$$\begin{aligned}
S_0 \otimes V_0 \quad & J_{ab} = \lambda_0 J_{ab}^{((osp))} \\
& Z_{ab} = \lambda_2 J_{ab}^{(osp)} \\
S_1 \otimes V_1 \quad & Q = \lambda_1 Q^{((o.5p))} \\
& Q' = \lambda_3 Q^{((asp))} \\
& P'_a = \lambda_0 P_a^{(osp)} \\
S_2 \otimes V_2 \quad & abccde = \lambda_0 Z_{abcde}^{(osp)} \\
& P_a = \lambda_2 P_a^{(osp)} \\
& a_bcc ee = \lambda_2 Z_{abcde}^{(osp)}
\end{aligned}$$

$$\begin{aligned}
\{Q, Q\} &\sim P + Z_2 + Z_5, \\
\{Q', Q'\} &\sim P + Z_2 + Z_5, \\
\{Q, Q'\} &\sim P' + J + Z'_5 \\
[P, P] &\sim J \\
[P', P'] &\sim J \\
[P, P'] &\sim Z_2
\end{aligned}$$



$$\begin{aligned}
& [Z_2, Z_2] \sim J, \\
& [Z_2, Z_5] \sim Z'_5, \\
& [Z_2, Z'_5] \sim Z_5, \\
& [Z_5, Z_5] \sim P' + J + Z'_5, \\
& [Z_5, Z'_5] \sim P + Z_2 + Z_5, \\
& [Z'_5, Z'_5] \sim P' + J + Z'_5 \\
& \quad [P, Q] \sim Q', \\
& \quad [Z_2, Q] \sim Q', \\
& \quad [Z_5, Q] \sim Q' \\
& \quad [P, Q'] \sim Q, \\
& \quad [Z_2, Q'] \sim Q, \\
& \quad [Z_5, Q'] \sim Q, \\
& \quad [P', Q] \sim Q, \\
& \quad [J, Q] \sim Q, \\
& \quad [Z'_5, Q] \sim Q, \\
& \quad [P', Q'] \sim Q' \\
& \quad [J, Q'] \sim Q' \\
& \quad [Z'_5, Q'] \sim Q' \\
& |T_{(A_1, \alpha_1)} \cdots T_{(A_n, \alpha_n)}| \equiv \alpha_\gamma K_{\alpha_1 \cdots \alpha_n}{}^\gamma |T_{A_1} \cdots T_{A_n}| \\
& \quad \sum_{p=1}^n X_{A_0 \cdots A_n}^{(p)} = 0 \\
& X_{A_0 \cdots A_n}^{(p)} = (-1)^{q(A_0)(q(A_1) + \cdots + q(A_{p-1}))} C_{A_0 A_p}^B \times \\
& \quad \times |T_{A_1} \cdots T_{A_{p-1}} T_B T_{A_{p+1}} \cdots T_{A_n}|, \\
& X_{(A_0, \alpha_0) \cdots (A_n, \alpha_n)}^{(p)} = (-1)^{q(A_0, \alpha_0)(q(A_1, \alpha_1) + \cdots + q(A_{p-1}, \alpha_{p-1}))} \times \\
& \quad \times C_{(A_0, \alpha_0)(A_p, \alpha_p)}^{(B, \beta)} |T_{(A_1, \alpha_1)} \cdots T_{(A_{p-1}, \alpha_{p-1})} \\
& \quad \times T_{(B, \beta)} T_{(A_{p+1}, \alpha_{p+1})} \cdots T_{(A_n, \alpha_n)}|, \\
& X_{(A_0, \alpha_0) \cdots (A_n, \alpha_n)}^{(p)} = \alpha_\gamma K_{\alpha_0 \cdots \alpha_n}{}^\gamma X_{A_0 \cdots A_n}^{(p)} \\
& \quad \sum_{p=1}^n X_{(A_0, \alpha_0) \cdots (A_n, \alpha_n)}^{(p)} = 0 \\
& |T_{(A_1, \alpha_1)} \cdots T_{(A_n, \alpha_n)}| = \sum_{m=0}^M \alpha_\gamma^{\beta_1 \cdots \beta_m} K_{\beta_1 \cdots \beta_m \alpha_1 \cdots \alpha_n}{}^\gamma |T_{A_1} \cdots T_{A_n}| \\
& \quad \text{STr}(T_{(A_1, \alpha_1)} \cdots T_{(A_n, \alpha_n)}) = K_{\gamma \alpha_1 \cdots \alpha_n}{}^\gamma \text{Str}(T_{A_1} \cdots T_{A_n}) \\
& \quad |T_{(A_{p_1}, \alpha_{p_1})} \cdots T_{(A_{p_n}, \alpha_{p_n})}| = \alpha_\gamma K_{\alpha_{p_1} \cdots \alpha_{p_n}}{}^\gamma |T_{A_{p_1}} \cdots T_{A_{p_n}}| \\
& \quad |T_{(A_1, i_1)} \cdots T_{(A_n, i_n)}| \equiv \alpha_j K_{i_1 \cdots i_n}{}^j |T_{A_1} \cdots T_{A_n}| \\
& Y_{(A_0, i_0) \cdots (A_n, i_n)}^{(p)} = (-1)^{q(A_0, i_0)(q(A_1, i_1) + \cdots + q(A_{p-1}, i_{p-1}))} \times \\
& \quad \times C_{(A_0, i_0)(A_p, i_p)}^{(B, j)} |T_{(A_1, i_1)} \cdots T_{(A_{p-1}, i_{p-1})} \\
& \quad \times T_{(B, j)} T_{(A_{p+1}, i_{p+1})} \cdots T_{(A_n, i_n)}|. \\
& Y_{(A_0, i_0) \cdots (A_n, i_n)}^{(p)} = \alpha_k K_{i_0 i_p}{}^j K_{i_1 \cdots i_{p-1} j i_{p+1} \cdots i_n}{}^k X_{A_0 \cdots A_n}^{(p)}
\end{aligned}$$

$$\begin{aligned}
K_{i_0 i_p}{}^j K_{i_1 \dots i_{p-1} j i_{p+1} \dots i_n}{}^k &= K_{i_0 i_p}{}^\gamma K_{i_1 \dots i_{p-1} \gamma i_{p+1} \dots i_n}{}^k + \\
&\quad - K_{i_0 i_p}{}^{N+1} K_{i_1 \dots i_{p-1} (N+1) i_{p+1} \dots i_n}{}^k \\
&= K_{i_0 \dots i_n}{}^k - K_{i_0 i_p}{}^{N+1} K_{i_1 \dots i_{p-1} (N+1) i_{p+1} \dots i_n}{}^k. \\
K_{i_0 i_p}{}^j K_{i_1 \dots i_{p-1} j i_{p+1} \dots i_n}{}^k &= K_{i_0 \dots i_n}{}^k \\
Y_{(A_0, i_0) \dots (A_n, i_n)}^{(p)} &= \alpha_k K_{i_0 \dots i_n}{}^k X_{A_0 \dots A_n}^{(p)} \\
\sum_{p=1}^n Y_{(A_0, i_0) \dots (A_n, i_n)}^{(p)} &= 0 \\
|T_{(a_{p_1}, i_{p_1})} \dots T_{(a_{p_n}, i_{p_n})}| &= \alpha_j K_{i_{p_1} \dots i_{p_n}}{}^j |T_{a_{p_1}} \dots T_{a_{p_n}}| \\
\text{STr}(T_{(A_1, i_1)} \dots T_{(A_n, i_n)}) &= K_{j_1 i_1}{}^2 K_{j_2 i_2}{}^{j_3} \dots K_{j_{n-1} i_{n-1}}{}^{j_n} \times \\
&\quad \times K_{j_n i_n}{}^1 \text{Str}(T_{A_1} \dots T_{A_n}). \\
K_{j_1 i_1 \dots i_n}{}^{j_1} &= K_{j_1 i_1}{}^2 K_{j_2 i_2}{}^3 \dots K_{j_{n-1} i_{n-1}}{}^n K_{j_n i_n}{}^{j_1} \\
&= K_{j_1 i_1}{}^2 K_{j_2 i_2}{}^2 \dots K_{j_{n-1} i_{n-1}}{}^n K_{j_n i_n}{}^1, \\
\text{STr}(T_{(A_1, i_1)} \dots T_{(A_n, i_n)}) &= K_{j_1 i_1 \dots i_n}{}^{j_1} \text{Str}(T_{A_1} \dots T_{A_n}) \\
\text{STr}(T_{(A_1, i_1)} \dots T_{(A_n, i_n)}) &= K_{i_1 \dots i_n}{}^0 \text{Str}(T_{A_1} \dots T_{A_n}) \\
\text{STr}(T_{(A_1, 0)} \dots T_{(A_n, 0)}) &= \text{Str}(T_{A_1} \dots T_{A_n}) \\
\{Q, \bar{Q}\} &= 2\gamma^a P_a \\
\{Q, \bar{Q}\} &= \frac{1}{8} \left(\Gamma^a P_a - \frac{1}{2} \Gamma^{ab} Z_{ab} + \frac{1}{5!} \Gamma^{abcde} Z_{abcde} \right) \\
[J^{ab}, Z_{cd}] &= \delta_{ecd}^{abf} Z_f^e, \\
[J^{ab}, Z_{c_1 \dots c_5}] &= \frac{1}{4!} \delta_{dc_1 \dots c_5}^{abe_1 \dots e_4} Z_{e_1 \dots e_4}^d \\
\lambda_\alpha \lambda_\beta &= \begin{cases} \lambda_{\alpha+\beta}, & \text{cuando } \alpha + \beta \leq 2 \\ \lambda_3, & \text{cuando } \alpha + \beta \geq 3 \end{cases} \\
\langle T_{(A_1, i_1)} \dots T_{(A_n, i_n)} \rangle &= \alpha_j K_{i_1 \dots i_n}{}^j \langle T_{A_1} \dots T_{A_n} \rangle \\
K_{i_1 \dots i_n}{}^j &= \delta_{i_1 + \dots + i_n}^j \\
\langle J_{a_1 b_1} \dots J_{a_6 b_6} \rangle_M &= \alpha_0 \langle J_{a_1 b_1} \dots J_{a_6 b_6} \rangle_{osp} \\
\langle J_{a_1 b_1} \dots J_{a_5 b_5} P_c \rangle_M &= \alpha_2 \langle J_{a_1 b_1} \dots J_{a_5 b_5} P_c \rangle_{osp}, \\
\langle J_{a_1 b_1} \dots J_{a_5 b_5} Z_{cd} \rangle_M &= \alpha_2 \langle J_{a_1 b_1} \dots J_{a_5 b_5} J_{cd} \rangle_{osp}, \\
\langle J_{a_1 b_1} \dots J_{a_5 b_5} Z_{c_1 \dots c_5} \rangle_M &= \alpha_2 \langle J_{a_1 b_1} \dots J_{a_5 b_5} Z_{c_1 \dots c_5} \rangle_{osp}, \\
\langle QJ_{a_1 b_1} \dots J_{a_4 b_4} \bar{Q} \rangle_M &= \alpha_2 \langle QJ_{a_1 b_1} \dots J_{a_4 b_4} \bar{Q} \rangle_{osp} \\
G &= \begin{bmatrix} C_{\alpha\beta} & 0 \\ 0 & 1 \end{bmatrix}
\end{aligned}$$

$$\begin{aligned}
\mathbf{P}_a &= \begin{bmatrix} \frac{1}{2}(\Gamma_a)^\alpha & 0 \\ 0 & 0 \end{bmatrix}, \\
J_{ab} &= \begin{bmatrix} \frac{1}{2}(\Gamma_{ab})^\alpha & 0 \\ 0 & 0 \end{bmatrix}, \\
\mathbf{Z}_{abcde} &= \begin{bmatrix} \frac{1}{2}(\Gamma_{abcde})^\alpha & 0 \\ 0 & 0 \end{bmatrix} \\
\mathbf{Q}^\gamma &= \begin{bmatrix} 0 & C^{\gamma\alpha} \\ \delta_\beta^\gamma & 0 \end{bmatrix} \\
A &= \begin{bmatrix} A^\alpha{}_\beta & A^\alpha \\ A_\beta & 0 \end{bmatrix} \\
\{A_1\} &= A_1 \\
\{A_1 \cdots A_n\} &= \frac{1}{n} \sum_{p=1}^n A_p \{A_1 \cdots \hat{A}_p \cdots A_n\} \\
\langle A_1 \cdots A_6 \rangle &= \text{STr}\{A_1 \cdots A_6\} \\
\langle A_1 \cdots A_6 \rangle &= \text{Tr}\{A_1 \cdots A_6\} \\
\langle A_1 \cdots A_6 \rangle &= \text{Tr}(\{A_1 \cdots A_5\}A_6) \\
\langle \chi \zeta A_1 \cdots A_4 \rangle &= -\frac{2}{5} \bar{\chi} \{A_1 \cdots A_4\} \zeta \\
\{A_1 A_2\} &= 2\text{Tr}(A_1 A_2) \mathbb{1} + A_1 A_2 \Gamma_{[4]}, \\
\{A_1 A_2 A_3\} &= \sum_{\langle ijk \rangle} \left(\text{Tr}(A_i A_j) A_k - \frac{4}{3} [A_i A_j A_k] \right) \Gamma_{[2]} + A_1 A_2 A_3 \Gamma_{[6]}, \\
\{A_1 \cdots A_4\} &= \sum_{\langle ijkl \rangle} \left(\frac{1}{2} \text{Tr}(A_i A_j) \text{Tr}(A_k A_l) - \frac{2}{3} \text{Tr}(A_i A_j A_k A_l) \right) \mathbb{1} + \\
&+ \sum_{\langle ijkl \rangle} \left(\frac{1}{2} \text{Tr}(A_i A_j) A_k A_l - \frac{4}{3} A_i [A_j A_k A_l] \right) \Gamma_{[4]} + \\
&+ (A_1 \cdots A_4) \Gamma_{[8]}, \\
\{A_1 \cdots A_5\} &= \sum_{\langle ijklm \rangle} \left(\frac{1}{2} \text{Tr}(A_i A_j) \text{Tr}(A_k A_l) A_m - \frac{2}{3} \text{Tr}(A_i A_j A_k A_l) A_m + \right. \\
&- \frac{4}{3} \text{Tr}(A_i A_j) [A_k A_l A_m] + \frac{32}{15} [A_i A_j A_k A_l A_m] \left. \right) \Gamma_{[2]} + \\
&+ \sum_{\langle ijklm \rangle} \left(\frac{1}{6} \text{Tr}(A_i A_j) A_k A_l A_m - \frac{2}{3} A_i A_j [A_k A_l A_m] \right) \Gamma_{[6]} + \\
&+ (A_1 \cdots A_5) \Gamma_{[10]} \\
\text{Tr}(A_{i_1} \cdots A_{i_n}) &= (A_{i_1})_{c_1}^{c_1} (A_{i_1})_{c_3}^{c_2} \cdots (A_{i_n})^{c_n} \\
[A_{i_1} \cdots A_{i_n}]^{ab} &= (A_{i_1})^a_{c_1} (A_{i_2})_{c_1}^{c_1} \cdots (A_{i_n})^{c_{n-1}b}_{c_2}
\end{aligned}$$



$$\begin{aligned}
\langle J^5 P \rangle_{\text{osp}} &= \frac{1}{2} \varepsilon_{a_1 \dots a_{11}} L_1^{a_1 a_2} \dots L_5^{a_9 a_{10}} B_1^{a_{11}}, \\
\langle J^6 \rangle_{\text{osp}} &= \frac{1}{3} \sum_{\langle i_1 \dots i_6 \rangle} \left[\frac{1}{4} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \right. \\
&\quad \left. - \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \frac{16}{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4} L_{i_5} L_{i_6}) \right], \\
\langle J^5 Z \rangle_{\text{osp}} &= \frac{1}{3} \varepsilon_{a_1 \dots a_{11}} \sum_{\langle i_1 \dots i_5 \rangle} \left[-\frac{5}{4} L_{i_1}^{a_1 a_2} \dots L_{i_4}^{a_7 a_8} (L_{i_5})_{bc} B_5^{bca_9 a_{10} a_{11}} + \right. \\
&\quad + 10 L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} L_{i_3}^{a_5 a_6} (L_{i_4})^{a_7} (L_{i_5})^{a_8} {}_c B_5^{ba_9 a_{10} a_{11}} + \\
&\quad + \frac{1}{4} L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} L_{i_3}^{a_5 a_6} B_5^{a_7 \dots a_{11}} \text{Tr}(L_{i_4} L_{i_5}) + \\
&\quad \left. - L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} [L_{i_3} L_{i_4} L_{i_5}]_5^{a_5 a_6} B_5^{a_7 \dots a_{11}} \right], \\
\langle QJ^4 \bar{Q} \rangle_{\text{osp}} &= -\frac{1}{240} \varepsilon_{a_1 \dots a_8 abc} L_1^{a_1 a_2} \dots L_4^{a_7 a_8} (\bar{\chi} \Gamma^{abc} \zeta) + \\
&\quad + \frac{1}{60} \sum_{\langle i_1 \dots i_4 \rangle} \left[\frac{3}{4} \text{Tr}(L_{i_1} L_{i_2}) L_{i_3}^{a_1 a_2} L_{i_4}^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \right. \\
&\quad \left. - 2 L_{i_1 a_2}^{a_1} [L_{i_2} L_{i_3} L_{i_4}]^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \right. \\
&\quad \left. + \frac{3}{4} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \bar{\chi} \zeta - \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \bar{\chi} \zeta \right] \\
\langle J^5 P \rangle_{\text{osp}} &= L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_1^c \langle J_{a_1 b_1} \dots J_{a_5 b_5} P_c \rangle_{\text{osp}} \\
\langle J^6 \rangle_{\text{osp}} &= L_1^{a_1 b_1} \dots L_6^{a_6 b_6} \langle J_{a_1 b_1} \dots J_{a_6 b_6} \rangle_{\text{osp}} \\
\langle J^5 Z \rangle_{\text{osp}} &= L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_5^{c_1 \dots c_5} \langle J_{a_1 b_1} \dots J_{a_5 b_5} Z_{c_1 \dots c_5} \rangle_{\text{osp}} \\
\langle QJ^4 \bar{Q} \rangle_{\text{osp}} &= L_1^{a_1 b_1} \dots L_4^{a_4 b_4} \bar{\chi}_\alpha \zeta^\beta \langle Q^\alpha J_{a_1 b_1} \dots J_{a_4 b_4} \bar{Q}_\beta \rangle_{\text{osp}} \\
\text{Tr}(A_{i_1} \dots A_{i_n}) &= (A_{i_1})^{c_1} {}_{c_2} (A_{i_1})^{c_2} \dots (A_{i_n})^{c_n} {}_{c_1}, \\
[A_{i_1} \dots A_{i_n}]^{ab} &= (A_{i_1})^a {}_{c_1} (A_{i_2})^{c_1} {}_{c_2} \dots (A_{i_n})^{c_{n-1} b}. \\
\langle J^5 P \rangle_{\text{osp}} &= \frac{1}{2} \varepsilon_{a_1 \dots a_{11}} L_1^{a_1 a_2} \dots L_5^{a_9 a_{10}} B_1^{a_{11}}, \\
\langle J^6 \rangle_{\text{osp}} &= \frac{1}{3} \sum_{\langle i_1 \dots i_6 \rangle} \left[\frac{1}{4} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \right. \\
&\quad \left. - \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \frac{16}{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4} L_{i_5} L_{i_6}) \right]
\end{aligned}$$

$$\begin{aligned}\langle J^5 Z \rangle_{\text{osp}} = & \frac{1}{3} \varepsilon_{a_1 \dots a_{11}} \sum_{\langle i_1 \dots i_5 \rangle} \left[-\frac{5}{4} L_{i_1}^{a_1 a_2} \dots L_{i_4}^{a_7 a_8} (L_{i_5})_{bc} B_5^{bca_9 a_{10} a_{11}} + \right. \\ & + 10 L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} L_{i_3}^{a_5 a_6} (L_{i_4})^{a_7}{}_b (L_{i_5})^{a_8}{}_c B_5^{bca_9 a_{10} a_{11}} + \\ & + \frac{1}{4} L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} L_{i_3}^{a_5 a_6} B_5^{a_7 \dots a_{11}} \text{Tr}(L_{i_4} L_{i_5}) + \\ & \left. - L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} [L_{i_3} L_{i_4} L_{i_5}]^{a_5 a_6} B_5^{a_7 \dots a_{11}} \right],\end{aligned}$$

$$\begin{aligned}\langle QJ^4 \bar{Q} \rangle_{\text{osp}} = & -\frac{1}{240} \varepsilon_{a_1 \dots a_8 abc} L_1^{a_1 a_2} \dots L_4^{a_7 a_8} (\bar{\chi} \Gamma^{abc} \zeta) + \\ & + \frac{1}{60} \sum_{\langle i_1 \dots i_4 \rangle} \left[\frac{3}{4} \text{Tr}(L_{i_1} L_{i_2}) L_{i_3}^{a_1 a_2} L_{i_4}^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \right. \\ & - 2 L_{i_1}^{a_1 a_2} [L_{i_2} L_{i_3} L_{i_4}]^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \\ & \left. + \frac{3}{4} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \bar{\chi} \zeta - \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \bar{\chi} \zeta \right],\end{aligned}$$

$$\langle J^5 P \rangle_{\text{osp}} = L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_1^c \langle J_{a_1 b_1} \dots J_{a_5 b_5} P_c \rangle_{\text{osp}},$$

$$\langle J^6 \rangle_{\text{osp}} = L_1^{a_1 b_1} \dots L_6^{a_6 b_6} \langle J_{a_1 b_1} \dots J_{a_6 b_6} \rangle_{\text{osp}},$$

$$\langle J^5 Z \rangle_{\text{osp}} = L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_5^{c_1 \dots c_5} \langle J_{a_1 b_1} \dots J_{a_5 b_5} Z_{c_1 \dots c_5} \rangle_{\text{osp}},$$

$$\langle QJ^4 \bar{Q} \rangle_{\text{osp}} = L_1^{a_1 b_1} \dots L_4^{a_4 b_4} \bar{\chi}_\alpha \zeta^\beta \left\langle Q^\alpha J_{a_1 b_1} \dots J_{a_4 b_4} \bar{Q}_\beta \right\rangle_{\text{osp}}.$$

$$\begin{aligned}\langle J^6 \rangle_{\text{M}} = & \frac{1}{3} \alpha_0 \sum_{\langle i_1 \dots i_6 \rangle} \left[\frac{1}{4} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \right. \\ & \left. - \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \frac{16}{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4} L_{i_5} L_{i_6}) \right]\end{aligned}$$

$$\langle J^5 P \rangle_{\text{M}} = \frac{1}{2} \alpha_2 \varepsilon_{a_1 \dots a_{11}} L_1^{a_1 a_2} \dots L_5^{a_9 a_{10}} B_1^{a_{11}}$$

$$\langle J^5 Z_2 \rangle_{\text{M}} =$$

$$\begin{aligned}\alpha_2 \sum_{\langle i_1 \dots i_5 \rangle} \left[\frac{1}{2} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} B_2) + \right. \\ \left. - \frac{4}{3} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4} L_{i_5} B_2) + \right. \\ \left. - \frac{2}{3} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} B_2) + \right. \\ \left. + \frac{32}{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4} L_{i_5} B_2) \right]\end{aligned}$$

$$\begin{aligned}
\langle J^5 Z_5 \rangle_M &= \frac{1}{3} \alpha_2 \varepsilon_{a_1 \dots a_{11}} \sum_{\langle i_1 \dots i_5 \rangle} \left[-\frac{5}{4} L_{i_1}^{a_1 a_2} \dots L_{i_4}^{a_7 a_8} (L_{i_5})_{bc} B_5^{bca_9 a_{10} a_{11}} + \right. \\
&\quad + 10 L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} L_{i_3}^{a_5 a_6} (L_{i_4})^{a_7}{}_b (L_{i_5})^{a_8}{}_c B_5^{bca_9 a_{10} a_{11}} + \\
&\quad + \frac{1}{4} \text{Tr}(L_{i_1} L_{i_2}) L_{i_3}^{a_1 a_2} L_{i_4}^{a_3 a_4} L_{i_5}^{a_5 a_6} B_5^{a_7 \dots a_{11}} + \\
&\quad \left. - L_{i_1}^{a_1 a_2} L_{i_2}^{a_3 a_4} [L_{i_3} L_{i_4} L_{i_5}]^{a_5 a_6} B_5^{a_7 \dots a_{11}} \right], \\
\langle QJ^4 \bar{Q} \rangle_M &= -\frac{\alpha_2}{240} \varepsilon_{a_1 \dots a_8 abc} L_1^{a_1 a_2} \dots L_4^{a_7 a_8} (\bar{\chi} \Gamma^{abc} \zeta) + \\
&\quad + \frac{\alpha_2}{60} \sum_{\langle i_1 \dots i_4 \rangle} \left[\frac{3}{4} \text{Tr}(L_{i_1} L_{i_2}) L_{i_3}^{a_1 a_2} L_{i_4}^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \right. \\
&\quad - 2 L_{i_1}^{a_1 a_2} [L_{i_2} L_{i_3} L_{i_4}]^{a_3}{}_{a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \\
&\quad \left. + \frac{3}{4} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \bar{\chi} \zeta - \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \bar{\chi} \zeta \right], \\
\langle J^6 \rangle_M &= L_1^{a_1 b_1} \dots L_6^{a_6 b_6} \langle J_{a_1 b_1} \dots J_{a_6 b_6} \rangle_M, \\
\langle J^5 P \rangle_M &= L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_1^c \langle J_{a_1 b_1} \dots J_{a_5 b_5} P_c \rangle_M, \\
\langle J^5 Z_2 \rangle_M &= L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_2^{cd} \langle J_{a_1 b_1} \dots J_{a_5 b_5} Z_{cd} \rangle_M, \\
\langle J^5 Z_5 \rangle_M &= L_1^{a_1 b_1} \dots L_5^{a_5 b_5} B_5^{c_1 \dots c_5} \langle J_{a_1 b_1} \dots J_{a_5} c_{c_1 c_5} \rangle_M, \\
\langle QJ^4 \bar{Q} \rangle_M &= L_1^{a_1 b_1} \dots L_4^{a_4 b_4} \bar{\chi}_\alpha \zeta^\beta \langle Q^\alpha J_{a_1 b_1} \dots J_{a_4 b_4} Q_\beta \rangle_M. \\
\langle \dots \rangle'_M &= \langle \dots \rangle_{6=6} + \beta_{4+2} \langle \dots \rangle_{6=4+2} + \beta_{2+2+2} \langle \dots \rangle_{6=2+2+2} \\
\langle J^6 \rangle'_M &= \frac{1}{3} \alpha_0 \sum_{\langle i_1 \dots i_6 \rangle} \left[\frac{1}{4} \gamma_5 \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \right. \\
&\quad \left. - \kappa_{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} L_{i_6}) + \frac{16}{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4} L_{i_5} L_{i_6}) \right] \\
\langle J^5 P \rangle'_M &= \frac{1}{2} \alpha_2 \varepsilon_{a_1 \dots a_{11}} L_1^{a_1 a_2} \dots L_5^{a_9 a_{10}} B_1^{a_{11}} \\
\langle J^5 Z_2 \rangle'_M &= \alpha_2 \sum_{\langle i_1 \dots i_5 \rangle} \left[\frac{1}{2} \gamma_5 \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} B_2) + \right. \\
&\quad - \frac{4}{3} \kappa_{15} \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4} L_{i_5} B_2) + \\
&\quad - \frac{2}{3} \kappa_{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \text{Tr}(L_{i_5} B_2) + \\
&\quad \left. + \frac{32}{15} \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4} L_{i_5} B_2) \right],
\end{aligned}$$

$$\begin{aligned}\langle J^5 \mathbf{Z}_5 \rangle'_M = & \frac{1}{3} \alpha_2 \varepsilon_{a_1 \dots a_{11}} \sum_{\langle i_1 \dots i_5 \rangle} \left[-\frac{5}{4} L_{i_1}^{a_1 a_2} \dots L_{i_4}^{a_4 a_8} (L_{i_5})_{bc} b_5^{bca_9 a_{10} a_{11}} + \right. \\ & + 10 L_{i_1}^{a_1 a_2} L_{i_2 a_4}^{a_3} L_{i_3}^{a_3 a_6} (L_{i_4})^{a_7}{}_b (L_{i_5})^{a_8}{}_c B_5^{bca_9 a_{10} a_{11}} + \\ & + \frac{1}{4} \kappa_{15} \text{Tr}(L_{i_1} L_{i_2}) L_{i_3}^{a_1 a_2} L_{i_4 a_4 a_4}^{a_5 a_6} B_5^{a_7 \dots a_{11}} + \\ & \left. - L_{i_1 a_2 a_2}^{a_3 a_4} [L_{i_3} L_{i_4} L_{i_5} a_5 a_6] B_5^{a_7 \dots a_{11}} \right],\end{aligned}$$

$$\begin{aligned}\langle QJ^4 \overline{Q} \rangle'_M = & -\frac{\alpha_2}{240} \varepsilon_{a_1 \dots a_8 abc} L_1^{a_1 a_2} \dots L_4^{a_7 a_8} (\bar{\chi} \Gamma^{abc} \zeta) + \\ & + \frac{\alpha_2}{60} \sum_{\langle i_1 \dots i_4 \rangle} \left\{ \frac{3}{4} \kappa_9 \text{Tr}(L_{i_1} L_{i_2}) L_{i_3}^{a_1 a_2} L_{i_4}^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \right. \\ & - 2 L_{i_1}^{a_2 a_2} [L_{i_2} L_{i_3} L_{i_4}]_3^{a_3 a_4} (\bar{\chi} \Gamma_{a_1 \dots a_4} \zeta) + \\ & + \frac{3}{4} (5\gamma_9 - 4) \text{Tr}(L_{i_1} L_{i_2}) \text{Tr}(L_{i_3} L_{i_4}) \bar{\chi} \zeta + \\ & \left. - \kappa_3 \text{Tr}(L_{i_1} L_{i_2} L_{i_3} L_{i_4}) \bar{\chi} \zeta \right\},\end{aligned}$$

$$\kappa_n = 1 + \frac{1}{n} \beta_{4+2} \text{Tr}(\mathbb{1}),$$

$$\gamma_n = \kappa_n + \frac{1}{15} \beta_{2+2+2} [\text{Tr}(\mathbb{1})]^2$$

$$\beta_{4+2} = \frac{1}{\text{Tr}(\mathbb{1})} n(\kappa_n - 1),$$

$$\beta_{2+2+2} = \frac{15}{[\text{Tr}(\mathbb{1})]^2} (\gamma_n - \kappa_n)$$

$$\kappa_m = 1 + \frac{n}{m} (\kappa_n - 1)$$

$$\gamma_m = \gamma_n + \left(\frac{n}{m} - 1\right) (\kappa_n - 1)$$

$$\beta_{4+2} = 0 \Leftrightarrow \kappa_n = 1$$

$$\beta_{2+2+2} = 0 \Leftrightarrow \gamma_n = \kappa_n.$$

$$\boldsymbol{A} = \boldsymbol{\omega} + \boldsymbol{e} + \boldsymbol{b}_2 + \boldsymbol{b}_5 + \overline{\boldsymbol{\psi}}$$

$$\overline{\boldsymbol{A}} = \overline{\boldsymbol{\omega}} + \overline{\boldsymbol{e}} + \overline{\boldsymbol{b}}_2 + \overline{\boldsymbol{b}}_5 + \bar{\chi}$$

$$\boldsymbol{\omega} = \frac{1}{2} \omega^{ab} \boldsymbol{J}_{ab},$$

$$\boldsymbol{e} = \frac{1}{\ell} e^a \boldsymbol{P}_a,$$

$$\boldsymbol{b}_2 = \frac{1}{2} b_2^{ab} \boldsymbol{Z}_{ab},$$

$$\boldsymbol{b}_5 = \frac{1}{5!} b_5^{abcde} \boldsymbol{Z}_{abcde},$$

$$\overline{\boldsymbol{\psi}} = \bar{\psi}_\alpha Q^\alpha$$

$$\boldsymbol{F} = \boldsymbol{R} + \boldsymbol{F}_P + \boldsymbol{F}_2 + \boldsymbol{F}_5 + \text{D}_\omega \overline{\boldsymbol{\psi}}$$

$$\begin{aligned}
R &= \frac{1}{2} R^{ab} J_{ab}, \\
F_P &= \left(\frac{1}{\ell} T^a + \frac{1}{16} \bar{\psi} \Gamma^a \psi \right) P_a, \\
F_2 &= \frac{1}{2} \left(D_\omega b^{ab} - \frac{1}{16} \bar{\psi} \Gamma^{ab} \psi \right) Z_{ab}, \\
F_5 &= \frac{1}{5!} \left(D_\omega b^{a_1 \dots a_5} + \frac{1}{16} \bar{\psi} \Gamma^{a_1 \dots a_5} \psi \right) Z_{a_1 \dots a_5},
\end{aligned}$$

$$\begin{aligned}
D_\omega \bar{\psi} &= D_\omega \bar{\psi} Q \\
D_\omega \psi &= d\psi + \frac{1}{4} \omega^{ab} \Gamma_{ab} \psi, \\
D_\omega \bar{\psi} &= d\bar{\psi} - \frac{1}{4} \omega^{ab} \bar{\psi} \Gamma_{ab}
\end{aligned}$$

$$\delta_{\text{gauge}} A = -D\lambda,$$

$$\delta_{\text{gauge}} \bar{A} = -\bar{D}\lambda$$

$$\lambda = \frac{1}{2} \lambda^{ab} J_{ab} + \frac{1}{\ell} \kappa^a P_a + \frac{1}{2} \kappa^{ab} Z_{ab} + \frac{1}{5!} \kappa^{abcde} Z_{abcde} + \bar{\epsilon} Q$$

- $\lambda = (1/\ell) \kappa^a P_a,$

$$\begin{aligned}
\delta e^a &= -D_\omega \kappa^a, \\
\delta b_2^{ab} &= 0, \\
\delta b_5^{abcde} &= 0, \\
\delta \omega^{ab} &= 0, \\
\delta \psi &= 0
\end{aligned}$$

- $Z_2: \lambda = (1/2) \kappa^{ab} Z_{ab},$

$$\begin{aligned}
\delta e^a &= 0, \\
\delta b_2^{ab} &= -D_\omega \kappa^{ab}, \\
\delta b_5^{abcde} &= 0, \\
\delta \omega^{ab} &= 0, \\
\delta \psi &= 0.
\end{aligned}$$

- $Z_5: \lambda = (1/5!) \kappa^{abcde} Z_{abcde},$

$$\begin{aligned}
\delta e^a &= 0, \\
\delta b_2^{ab} &= 0, \\
\delta b_5^{abcde} &= -D_\omega \kappa^{abcde}, \\
\delta \omega^{ab} &= 0, \\
\delta \psi &= 0.
\end{aligned}$$

- $\lambda = (1/2) \lambda^{ab} J_{ab},$

$$\begin{aligned}
\delta e^a &= \lambda^a{}_b e^b, \\
\delta b_2^{ab} &= -2\lambda^{[a}{}_c b_2^{b]c}, \\
\delta b_5^{abcde} &= 5\lambda^{[a}{}_f b_5^{bcde]}, \\
\delta \omega^{ab} &= -D_\omega \lambda^{ab}, \\
\delta \psi &= \frac{1}{4} \lambda^{ab} \Gamma_{ab} \psi.
\end{aligned}$$

- $\lambda = \bar{\epsilon} Q,$



$$\begin{aligned}
\delta e^a &= \frac{\ell}{8} \bar{\varepsilon} \Gamma^a \psi, \\
\delta b_2^{ab} &= -\frac{1}{8} \bar{\varepsilon} \Gamma^{ab} \psi, \\
\delta b_5^{abcde} &= \frac{1}{8} \bar{c} \Gamma^{abcde} \psi, \\
\delta \omega^{ab} &= 0, \\
\delta \psi &= -D_\omega \varepsilon \\
L_M^{(11)}(A, \bar{A}) &= Q_{A \leftarrow \bar{A}}^{(11)} \\
L_M^{(11)}(A, \bar{A}) &= Q_{A \leftarrow \bar{\omega}}^{(11)} + Q_{\bar{\omega} \leftarrow \bar{A}}^{(11)} + dQ_{A \leftarrow \bar{\omega} \leftarrow \bar{A}}^{(10)} \\
Q_{A \leftarrow \bar{\omega}}^{(11)} &= Q_{A \leftarrow \omega}^{(11)} + Q_{\omega \leftarrow \bar{\omega}}^{(11)} + dQ_{A \leftarrow \omega \leftarrow \bar{\omega}}^{(10)} \\
L_M^{(11)}(A, \bar{A}) &= Q_{A \leftarrow \omega}^{(11)} - Q_{\bar{A} \leftarrow \bar{\omega}}^{(11)} + Q_{\omega \leftarrow \bar{\omega}}^{(11)} + d\mathcal{B}_M^{(10)} \\
\mathcal{B}_M^{(10)} &= Q_{A \leftarrow \omega \leftarrow \bar{\omega}}^{(10)} + Q_{A \leftarrow \bar{\omega} \leftarrow \bar{A}}^{(10)} \\
A_0 &= \omega \\
A_1 &= \omega + e \\
A_2 &= \omega + e + b_2 \\
A_3 &= \omega + e + b_2 + b_5 \\
A_4 &= \omega + e + b_2 + b_5 + \bar{\psi} \\
Q_{A_4 \leftarrow A_0}^{(11)} &= Q_{A_4 \leftarrow A_3}^{(11)} + Q_{A_3 \leftarrow A_0}^{(11)} + dQ_{A_4 \leftarrow A_3 \leftarrow A_0}^{(10)} \\
Q_{A_3 \leftarrow A_0}^{(11)} &= Q_{A_3 \leftarrow A_2}^{(11)} + Q_{A_2 \leftarrow A_0}^{(11)} + dQ_{A_3 \leftarrow A_2 \leftarrow A_0}^{(10)} \\
Q_{A_2 \leftarrow A_0}^{(11)} &= Q_{A_2 \leftarrow A_1}^{(11)} + Q_{A_1 \leftarrow A_0}^{(11)} + dQ_{A_2 \leftarrow A_1 \leftarrow A_0}^{(10)} \\
Q_{A_4 \leftarrow A_0}^{(11)} &= 6 \left[H_a e^a + \frac{1}{2} H_{ab} b_2^{ab} + \frac{1}{5!} H_{abcde} b_5^{abcde} - \frac{5}{2} \bar{\psi} \mathcal{R} D_\omega \psi \right] \\
H_a &\equiv \langle R^5 P_a \rangle_M, \\
H_{ab} &\equiv \langle R^5 Z_{ab} \rangle_M, \\
H_{abcce} &\equiv \langle R^5 Z_{abcde} \rangle_M, \\
\mathcal{R}^\alpha{}_\beta &\equiv \left\langle Q^\alpha R^4 \bar{Q}_\beta \right\rangle_M, \\
H_a &= \frac{\alpha_2}{64} R_a^{(5)}, \\
H_{ab} &= \alpha_2 \left[\frac{5}{2} \left(R^4 - \frac{3}{4} R^2 R^2 \right) R_{ab} + 5 R^2 R_{ab}^3 - 8 R_{ab}^5 \right] \\
H_{abcde} &= -\frac{5}{16} \alpha_2 \left[5 R_{[ab} R_{cde]}^{(4)} - 40 R_{[a}^f R_{b}^g R_{cde]}^{(3)} + \right. \\
&\quad \left. - R^2 R_{abcde}^{(3)} + 4 R_{abcdefg}^{(2)} (R^3)^{fg} \right], \\
\mathcal{R} &= -\frac{\alpha_2}{40} \left\{ \left(R^4 - \frac{3}{4} R^2 R^2 \right) \mathbb{1} + \frac{1}{96} R_{abc}^{(4)} \Gamma^{abc} + \right. \\
&\quad \left. - \frac{3}{4} \left[R^2 R^{ab} - \frac{8}{3} (R^3)^{ab} \right] R^{cd} \Gamma_{abcd} \right\} \\
R^n &= R_{a_1}^{a_1} \dots R_{a_n}^{a_n} \\
R_{ab}^n &= R_{a c_1}^{c_1} R_{c_1 c_2}^{c_2} \dots R_{c_{n-1} b}^{c_{n-1}} \\
R_{a_1 \dots a_{d-2n}}^{(n)} &= \varepsilon_{a_1 \dots a_{d-2n} b_1 \dots b_{2n}} R^{b_1 b_2} \dots R^{b_{2n-1} b_{2n}}
\end{aligned}$$



$$\begin{aligned}
Q_{\omega \leftarrow \bar{\omega}}^{(11)} &= 3 \int_0^1 dt \theta^{ab} L_{ab}(t) \\
L_{ab}(t) &= \langle \mathbf{R}_t^5 \mathbf{J}_{ab} \rangle_{\mathbf{M}} \\
\mathbf{R}_t &= \frac{1}{2} R_t^{ab} \mathbf{J}_{ab}, \\
R_t^{ab} &= \bar{R}^{ab} + t D_{\bar{\omega}} \theta^{ab} + t^2 \theta^a{}_c \theta^{cb} \\
L_{ab}(t) &= \alpha_0 \left[\frac{5}{2} \left(R_t^4 - \frac{3}{4} R_t^2 R_t^2 \right) (R_t)_{ab} + 5 R_t^2 (R_t^3)_{ab} - 8 (R_t^5)_{ab} \right] \\
Q_{\omega \leftarrow \bar{\omega}}^{(11)} &= Q_{\omega \leftarrow 0}^{(11)} - Q_{\bar{\omega} \leftarrow 0}^{(11)} + dQ_{\omega \leftarrow A_3 \leftarrow \bar{\omega}}^{(10)} \\
\langle \mathbf{F}^5 \mathbf{G}_A \rangle_{\mathbf{M}} &= 0 \\
H_a &= 0, \\
H_{ab} &= 0, \\
H_{abcde} &= 0, \\
\mathcal{R} D_{\omega} \psi &= 0 \\
L_{ab} - 10 (D_{\omega} \bar{\psi}) \mathcal{Z}_{ab} (D_{\omega} \psi) + 5 H_{abc} \left(T^c + \frac{1}{16} \bar{\psi} \Gamma^c \psi \right) + \\
+ \frac{5}{2} H_{abcd} \left(D_{\omega} b^{cd} - \frac{1}{16} \bar{\psi} \Gamma^{cd} \psi \right) + \frac{1}{24} H_{abc_1 \dots c_5} \left(D_{\omega} b^{c_1 \dots c_5} + \frac{1}{16} \bar{\psi} \Gamma^{c_1 \dots c_5} \psi \right) &= 0 \\
L_{ab} &\equiv \langle \mathbf{R}^5 \mathbf{J}_{ab} \rangle_{\mathbf{M}}, \\
(\mathcal{Z}_{ab})^{\alpha} &\equiv \left\langle \mathbf{Q}^{\alpha} \mathbf{R}^3 \mathbf{J}_{ab} \bar{\mathbf{Q}}_{\beta} \right\rangle_{\mathbf{M}}, \\
H_{abc} &\equiv \langle \mathbf{R}^4 \mathbf{J}_{ab} \mathbf{P}_c \rangle_{\mathbf{M}}, \\
H_{abcd} &\equiv \langle \mathbf{R}^4 \mathbf{J}_{ab} \mathbf{Z}_{cd} \rangle_{\mathbf{M}}, \\
H_{abcdefg} &\equiv \langle \mathbf{R}^4 \mathbf{J}_{ab} \mathbf{Z}_{cdefg} \rangle_{\mathbf{M}}, \\
H_c &= \frac{1}{2} R^{ab} H_{abc}, \\
H_{cd} &= \frac{1}{2} R^{ab} H_{abcd}, \\
H_{cdefg} &= \frac{1}{2} R^{ab} H_{abcdefg}, \\
\mathcal{R}_{\beta}^{\alpha} &= \frac{1}{2} R^{ab} (\mathcal{Z}_{ab})^{\alpha}{}_{\beta} \\
L_{ab} &= \alpha_0 \left[\frac{5}{2} \left(R^4 - \frac{3}{4} R^2 R^2 \right) R_{ab} + 5 R^2 R_{ab}^3 - 8 R_{ab}^5 \right] \\
\mathcal{Z}_{ab} &= \frac{\alpha_2}{40} \left\{ 2 \left(R_{ab}^3 - \frac{3}{4} R^2 R_{ab} \right) \mathbb{1} - \frac{1}{48} R_{abcde}^{(3)} \Gamma^{cde} + \right. \\
&\quad \left. - \frac{3}{4} \left(R_{ab} R^{cd} - \frac{1}{2} R^2 \delta_{ab}^{cd} \right) R^{ef} \Gamma_{cdef} + \right. \\
&\quad \left. - \left[\delta_{ab}^{cg} R_{gh} R^{hd} R^{ef} - R_a^c R_b^d R^{ef} + \frac{1}{2} \delta_{ab}^{ef} (R^3)^{cd} \right] \Gamma_{cdef} \right\} \\
H_{abc} &= \frac{\alpha_2}{32} R_{abc}^{(4)}
\end{aligned}$$

$$\begin{aligned}
H_{abcd} = & \alpha_2 \delta_{ab}^{ef} g_{cd}^{gh} \left[\frac{3}{4} R^2 R_{ef} R_{gh} - R_{ef}^3 R_{gh} - R_{ef} R_{gh}^3 + \right. \\
& - \frac{4}{5} (R_{eh} R_{fg}^3 + R_{eh}^3 R_{fg} - R_{eh}^2 R_{fg}^2) + \frac{1}{2} R^2 R_{eh} R_{fg} + \\
& \left. + \frac{1}{8} \eta_{[ef][gh]} \left(R^4 - \frac{3}{4} R^2 R^2 \right) - \eta_{fg} \left(R^2 R_{eh}^2 - \frac{8}{5} R_{eh}^4 \right) \right], \\
H_{abc_1 \dots c_5} = & \frac{\alpha_2}{80} \delta_{c_1 \dots c_5}^{\text{clefg}} \left(-\frac{5}{3} R_{abccde}^{(3)} R_{fg} + 10 R_{abcdepq}^{(2)} R^p{}_f R^q{}_g + \right. \\
& - \frac{1}{6} R_{ab} R_{cdefg}^{(3)} + \frac{1}{4} R^2 R_{abcdefg}^{(2)} - \frac{2}{3} R_{abcdefgpq}^{(1)} (R^3)^{pq} + \\
& + \frac{1}{3} R^p{}_a R^q{}_b R_{clefgpq}^{(2)} - \frac{1}{3} R^q{}_a R_{bcdefgp}^{(2)} R^p{}_q + \frac{1}{3} R^q{}_b R_{acclefgp}^{(2)} R^p{}_q + \\
& \left. - \frac{10}{3} \eta_{ga} R_{bcdep}^{(3)} R^p{}_f + \frac{10}{3} \eta_{gb} R_{acdep}^{(3)} R^p{}_f - \frac{5}{24} \eta_{[ab][cd]} R_{efg}^{(4)} \right).
\end{aligned}$$

$$\langle \mathbf{F}^n \mathbf{G}_A \rangle = 0$$

$$\begin{aligned}
ds^2 = & e^{-2\xi|z|} \left(dz^2 + \tilde{g}_{\alpha\beta}^{(d)} dx^\alpha dx^\beta \right) + \gamma_{\kappa\lambda}^{(10-d)} dy^\kappa dy^\lambda \\
R^{ab} = & \tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b + 2\xi \theta(z) (\tilde{e}^a \kappa^b - \tilde{e}^b \kappa^a) - \kappa^a \kappa^b, \\
R^{aZ} = & -2e^{\xi|z|} \xi \delta(z) E^Z \tilde{e}^a - 2\xi \theta(z) \tilde{T}^a + D_{\tilde{\omega}} \kappa^a,
\end{aligned}$$

$$T^a = \kappa^a E^Z + e^{-\xi|z|} \tilde{T}^a,$$

$$T^Z = -e^{-\xi|z|} \kappa^a \tilde{e}_a$$

$$\xi e^{\xi|z|} \delta(z) E^Z \varepsilon_{abcd} (\tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b) \tilde{e}^c = \mathcal{T}_d$$

$$\varepsilon_{abcd} (\tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b) (\tilde{R}^{cd} - \xi^2 \tilde{e}^c \tilde{e}^d) = \mathcal{T}$$

$$\xi e^{\xi|z|} \delta(z) E^Z \varepsilon_{abcd} (\tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b) \tilde{e}^c = \mathcal{T}_d,$$

$$\varepsilon_{abcd} (\tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b) (\tilde{R}^{cd} - \xi^2 \tilde{e}^c \tilde{e}^d) = \mathcal{T},$$

$$\begin{aligned}
\mathcal{T}_d = & 2E^Z e^{\xi|z|} \xi \delta(z) \varepsilon_{abcd} \left(\frac{1}{2} \kappa^a \kappa^b - \xi \theta(z) (\tilde{e}^a \kappa^b - \tilde{e}^b \kappa^a) \right) \tilde{e}^c + \\
& + \varepsilon_{abcd} [\tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b + 2\xi \theta(z) (\tilde{e}^a \kappa^b - \tilde{e}^b \kappa^a) - \kappa^a \kappa^b] \times \\
& \times \left(\frac{1}{2} D_{\tilde{\omega}} \kappa^c - \xi \theta(z) \tilde{T}^c \right),
\end{aligned}$$

$$\begin{aligned}
\mathcal{T} = & -4\varepsilon_{abcd} \left(\tilde{R}^{ab} - \xi^2 \tilde{e}^a \tilde{e}^b + \xi \theta(z) (\tilde{e}^a \kappa^b - \tilde{e}^b \kappa^a) - \frac{1}{2} \kappa^a \kappa^b \right) \times \\
& \times \left(\xi \theta(z) (\tilde{e}^c \kappa^d - \tilde{e}^d \kappa^c) - \frac{1}{2} \kappa^c \kappa^d \right)
\end{aligned}$$

$$\delta L = E(\phi) \delta \phi + \mathrm{d} \Xi(\phi, \delta \phi)$$

$$\delta_{\mathrm{gauge}} L = \mathrm{d} \Omega$$

$$\star J_{\mathrm{gauge}} = \Omega - \Xi_{\mathrm{gauge}}$$

$$\delta_{\mathrm{dif}} L = -\mathcal{E}_{\xi} L$$

$$= -(\mathrm{d} \mathcal{I}_{\xi} + \mathcal{I}_{\xi} \mathrm{d}) L$$

$$= -\mathrm{d} \mathcal{I}_{\xi} L,$$

$$\star J_{\mathrm{dif}} = -\Xi_{\mathrm{dif}} - \mathcal{I}_{\xi} L$$

$$\Xi = n(n+1)k \int_0^1 dt \langle \delta \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle$$

$$\star J_{\mathrm{gauge}} = \Omega - \Xi_{\mathrm{gauge}}$$

$$\Omega = 0$$



$$\begin{aligned}
\delta_{\text{gauge}} \mathbf{A} &= -\mathbf{D}\boldsymbol{\lambda}, \\
\delta_{\text{gauge}} \bar{\mathbf{A}} &= -\bar{\mathbf{D}}\bar{\boldsymbol{\lambda}} \\
\delta_{\text{gauge}} \mathbf{A}_t &= -\mathbf{D}_t \boldsymbol{\lambda} \\
\Xi_{\text{gauge}} &= -n(n+1)k \int_0^1 dt \langle \mathbf{D}_t \boldsymbol{\lambda} \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\Xi_{\text{gauge}} &= -n(n+1)k \int_0^1 dt \langle \boldsymbol{\lambda} \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + n(n+1)k \int_0^1 dt \langle \boldsymbol{\lambda} \mathbf{D}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\Xi_{\text{gauge}} &= -n(n+1)k \int_0^1 dt \langle \boldsymbol{\lambda} \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + (n+1)k \int_0^1 dt \frac{d}{dt} \langle \boldsymbol{\lambda} \mathbf{F}_t^n \rangle \\
\Xi_{\text{gauge}} &= -n(n+1)k \int_0^1 dt \langle \boldsymbol{\lambda} \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + (n+1)k \left(\langle \boldsymbol{\lambda} \mathbf{F}^n \rangle - \langle \boldsymbol{\lambda} \bar{\mathbf{F}}^n \rangle \right) \\
\star J_{\text{gauge}} &= n(n+1)k \int_0^1 dt \langle \boldsymbol{\lambda} \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\star J_{\text{dif}} &= -\Xi_{\text{dif}} - \mathbf{I}_\xi L_{\text{T}}^{(2n+1)} \\
\delta_{\text{dif}} \mathbf{A} &= -\mathbf{E}_\xi \mathbf{A}, \\
\delta_{\text{dif}} \bar{\mathbf{A}} &= -\mathbf{E}_\xi \bar{\mathbf{A}} \\
\delta_{\text{dif}} \mathbf{A}_t &= -\mathbf{E}_\xi \mathbf{A}_t \\
\Xi_{\text{dif}} &= -n(n+1)k \int_0^1 dt \langle \mathbf{E}_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\mathbf{E}_\xi \mathbf{A}_t &= \mathbf{I}_\xi \mathbf{F}_t + \mathbf{D}_t \mathbf{I}_\xi \mathbf{A}_t \\
\Xi_{\text{dif}} &= -n(n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \mathbf{F}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle - n(n+1)k \int_0^1 dt \langle \mathbf{D}_t \mathbf{I}_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle \\
\Xi_{\text{dif}} &= -\mathbf{I}_\xi L_{\text{T}}^{(2n+1)} + (n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \boldsymbol{\Theta} \mathbf{F}_t^n \rangle + \\
&\quad -n(n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + \\
&\quad +n(n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \mathbf{A}_t \mathbf{D}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle. \\
\frac{d}{dt} \mathbf{F}_t &= \mathbf{D}_t \boldsymbol{\Theta} \\
\frac{d}{dt} \mathbf{I}_\xi \mathbf{A}_t &= \mathbf{I}_\xi \boldsymbol{\Theta} \\
\Xi_{\text{dif}} &= -\mathbf{I}_\xi L_{\text{T}}^{(2n+1)} + (n+1)k \int_0^1 dt \frac{d}{dt} \langle \mathbf{I}_\xi \mathbf{A}_t \mathbf{F}_t^n \rangle + \\
&\quad -n(n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle, \\
\Xi_{\text{dif}} + \mathbf{I}_\xi L_{\text{T}}^{(2n+1)} &= -n(n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle + \\
&\quad + (n+1)k \left(\langle \mathbf{I}_\xi \mathbf{A} \mathbf{F}^n \rangle - \langle \mathbf{I}_\xi \bar{\mathbf{A}} \bar{\mathbf{F}}^n \rangle \right). \\
\star J_{\text{dif}} &= n(n+1)k \int_0^1 dt \langle \mathbf{I}_\xi \mathbf{A}_t \boldsymbol{\Theta} \mathbf{F}_t^{n-1} \rangle
\end{aligned}$$

$$\begin{aligned}
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} &\equiv \det \begin{bmatrix} \delta_{b_1}^{a_1} & \delta_{b_2}^{a_1} & \dots & \delta_{b_n}^{a_1} \\ \delta_{b_1}^{a_2} & \delta_{b_2}^{a_2} & \dots & \delta_{b_n}^{a_2} \\ \vdots & \vdots & \ddots & \vdots \\ \delta_{b_1}^{a_n} & \delta_{b_2}^{a_n} & \dots & \delta_{b_n}^{a_n} \end{bmatrix} \\
\delta_{b_1 \dots b_r a_{r+1} \dots a_n}^{a_1 \dots a_r a_{r+1} \dots a_n} &= \frac{(d-r)!}{(d-n)!} \delta_{b_1 \dots b_r}^{a_1 \dots a_r} \\
\delta_{a_1 \dots a_n}^{a_1 \dots a_n} &= \frac{d!}{(d-n)!} \\
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} A_{b_1 \dots b_n} &= n! A_{a_1 \dots a_n}, \\
\delta_{b_1 \dots b_n}^{a_n} A_{a_1 \dots a_n} &= n! A_{b_1 \dots b_n} \\
\varepsilon_{a_1 \dots a_d} &= \delta_{a_1 \dots a_d}^{1 \dots d}, \\
\varepsilon^{a_1 \dots a_d} &= \delta_{1 \dots d}^{a_1 \dots a_d} \\
\varepsilon^{a_1 \dots a_d} \varepsilon_{b_1 \dots b_d} &= \delta_{b_1 \dots b_d}^{a_1 \dots a_d} \\
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} &= \sum_{p=1}^n (-1)^{p+1} \delta_{b_p}^{a_1} \mathcal{V}_{b_1 \dots b_p \dots b_n}^{a_2 \dots a_n} \\
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} &= \sum_{p=1}^{n-1} \sum_{q=p+1}^n (-1)^{p+q+1} \delta_{b_p b_q}^{a_1 a_2} \delta_{b_1 \dots \hat{b}_p \dots \hat{b}_q \dots b_n}^{a_2 \dots a_n} \\
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} &= (-1)^{r(r+1)/2} \sum_{p_1=1}^{n-r+1} \sum_{p_2=p_1+1}^{n-r+2} \dots \sum_{p_{r-1}=p_{r-2}+1}^{n-1} \\
&\quad \sum_{p_r=p_{r-1}+1}^n (-1)^{p_1+\dots+p_r} \delta_{b_{p_1} \dots b_{p_r}}^{a_{r+1} \dots a_n} \delta_{b_1 \dots b_{p_1} \dots b_{p_r} \dots b_n}^{a_1 \dots a_r} \\
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} &= (-1)^{r(r+1)/2} \sum_{p_1=1}^{n-r+1} \sum_{p_2=p_1+1}^{n-r+2} \dots \sum_{p_{r-1}=p_{r-2}+1}^{n-1} \\
&\quad \sum_{p_r=p_{r-1}+1}^n (-1)^{p_1+\dots+p_r} \delta_{b_1 \dots b_r}^{a_{p_1} \dots a_{p_r}} \delta_{b_{r+1} \dots b_n}^{a_1 \dots \hat{a}_{p_1} \dots \hat{a}_{p_r} \dots a_n} \\
\delta_{b_1 \dots b_n}^{a_1 \dots a_n} &= \binom{n}{p} \delta_{b_1 \dots b_p}^{a_1 \dots a_p} \delta_{b_{p+1} \dots b_n}^{a_{p+1} \dots a_n} = \binom{n}{p} \delta_{[b_1 \dots b_p]}^{a_1 \dots a_p} \delta_{b_{p+1} \dots b_n}^{a_{p+1} \dots a_n} \\
\eta_{[a_1 \dots a_p][b_1 \dots b_p]} &\equiv \delta_{a_1 \dots a_p}^{c_1 \dots c_p} (\eta_{b_1 c_1} \dots \eta_{b_p c_p}), \\
\eta_{[a_1 \dots a_p][b_1 \dots b_p]} &\equiv \delta_{c_1 \dots c_p}^{a_1 \dots a_p} (\eta^{b_1 c_1} \dots \eta^{b_p c_p}) \\
\eta_{[b_1 \dots b_p][a_1 \dots a_p]} &= \eta_{[a_1 \dots a_p][b_1 \dots b_p]}, \\
\eta^{b_1 \dots b_p}[a_1 \dots a_p] &= \eta^{a_1 \dots a_p}[b_1 \dots b_p], \\
\eta_{[a_1 \dots a_p][b_1 \dots b_p]} A^{b_1 \dots b_p} &= p! A_{a_1 \dots a_p}, \\
\eta_{[a_1 \dots a_p][b_1 \dots b_p]} A_{b_1 \dots b_p} &= p! A^{a_1 \dots a_p}, \\
\eta_{[a_1 \dots a_p][c_1 \dots c_p]} \eta_{[c_1 \dots c_p][b_1 \dots b_p]} &= p! \delta_{b_1 \dots b_p \dots a_p}^{a_1} \\
\{\Gamma_a, \Gamma_b\} &= 2\eta_{ab} \\
\Gamma_{a_1 \dots a_p} &\equiv \Gamma_{[a_1} \dots \Gamma_{a_p]} \\
&= \frac{1}{p!} \delta_{a_1 \dots a_p}^{b_1 \dots b_p} \Gamma_{b_1} \dots \Gamma_{b_p}
\end{aligned}$$



$$\begin{aligned}
\Gamma_* &\equiv \Gamma_0 \cdots \Gamma_{d-1} \\
&= \frac{1}{d!} \varepsilon^{a_1 \cdots a_d} \Gamma_{a_1} \cdots \Gamma_{a_d} \\
&= \frac{1}{d!} \varepsilon^{a_1 \cdots a_d} \Gamma_{a_1 \cdots a_d} \\
\Gamma_*^2 &= (-1)^{(d-2)(d+1)/2} \\
\Gamma_* \Gamma_a &= (-1)^{d+1} \Gamma_a \Gamma_*, \\
\Gamma_* \Gamma_{a_1 \cdots a_p} &= (-1)^{p(d+1)} \Gamma_{a_1 \cdots a_p} \Gamma_*, \\
\Gamma_{a_1 \cdots a_{d-k}} &= \frac{1}{k!} (-1)^{k(k-1)/2} \varepsilon_{a_1 \cdots a_d} \Gamma^{a_{d-k+1} \cdots a_d} \Gamma_* \\
\Gamma_a^T &= \xi C \Gamma_a C^{-1} \\
C^T &= \lambda C
\end{aligned}$$

$d \bmod 8$	λ	ξ	S	A
0	+1	+1	1,4	2,3
0	+1	-1	3,4	1,2
1	+1	+1	1,4	2,3
2	+1	+1	1,4	2,3
2	-1	-1	1,2	3,4
3	-1	-1	1,2	3,4
4	-1	+1	2,3	1,4
4	-1	-1	1,2	3,4
5	-1	+1	2,3	1,4
6	-1	+1	2,3	1,4
6	+1	-1	3,4	1,2
7	+1	-1	3,4	1,2

$$\begin{aligned}
\left(C \Gamma_{a_1 \cdots a_p} \right)^T &= (-1)^{p(p-1)/2} \lambda \xi^p \left(C \Gamma_{a_1 \cdots a_p} \right) \\
\text{Tr} \left(\Gamma^{a_1 \cdots a_p} \Gamma_{b_1 \cdots b_q} \right) &= (-1)^{p(p-1)/2} m \delta_{b_1 \cdots b_q}^{a_1 \cdots a_p} \\
M &= \sum_{p \geq 0} \frac{1}{p!} k_{a_1 \cdots a_p} \Gamma^{a_1 \cdots a_p} \\
k_{a_1 \cdots a_p} &= \frac{1}{m} (-1)^{p(p-1)/2} \text{Tr} \left(M \Gamma_{a_1 \cdots a_p} \right) \\
\Gamma_{a_1 \cdots a_i} \Gamma^{b_1 \cdots b_j} &= \sum_{k=|i-j|}^{i+j} \frac{i! j!}{s! t! u!} \delta_{[a_i}^{[b_1} \cdots \delta_{a_{t+1}}^{b_s} \Gamma_{a_1 \cdots a_t]}^{b_{s+1} \cdots b_j]}
\end{aligned}$$



$$\begin{aligned}
s &= \frac{1}{2}(i+j-k) \\
t &= \frac{1}{2}(i-j+k) \\
u &= \frac{1}{2}(-i+j+k) \\
\Gamma^{a_1 \cdots a_i} \Gamma_{b_1 \cdots b_j} &= \sum_{s=0}^{\min(i,j)} \frac{1}{t! u!} (-1)^{s(s-1)/2} \delta_{d_1 \cdots d_t b_1 \cdots b_j}^{a_1 \cdots a_i e_1 \cdots e_u} \Gamma_{e_1 \cdots e_u}^{d_1 \cdots d_t} \\
&\quad t = i-s, \\
&\quad u = j-s \\
[\Gamma^{a_1 \cdots a_i}, \Gamma_{b_1 \cdots b_j}] &= \sum_{s=0}^{\min(i,j)} \frac{1}{t! u!} (-1)^{s(s-1)/2} \times \\
&\quad \times [1 - (-1)^{ij-s^2}] \delta_{d_1 \cdots d_t b_1 \cdots b_j}^{a_1 \cdots a_i e_1 \cdots e_u} \Gamma_{e_1 \cdots e_u}^{d_1 \cdots d_t}, \\
\{\Gamma^{a_1 \cdots a_i}, \Gamma_{b_1 \cdots b_j}\} &= \sum_{s=0}^{\min(i,j)} \frac{1}{t! u!} (-1)^{s(s-1)/2} \times \\
&\quad \times [1 + (-1)^{ij-s^2}] \delta_{d_1 \cdots d_t b_1 \cdots b_j}^{a_1 \cdots a_i e_1 \cdots e_u} \Gamma_{e_1 \cdots e_u}^{d_1 \cdots d_t}. \\
\Gamma_{a_1 \cdots a_i} \Gamma_{b_1 \cdots b_j} &= \sum_{s=0}^{\min(i,j)} D_{a_1 \cdots a_i b_1 \cdots b_j}(s) \\
D_{a_1 \cdots a_i b_1 \cdots b_j}(s) &= (-1)^{s(i-s)+s(s-1)/2} \sum_{p_1=1}^{1+j-s} \cdots \sum_{p_s=p_{s-1}+1}^j \\
&\quad \sum_{q_1=1}^{1+i-s} \cdots \sum_{q_s=q_{s-1}+1}^i (-1)^{p_1+\cdots+p_s+q_1+\cdots+q_s} \\
&\quad \eta_{[a_{q_1} \cdots a_{q_s}][b_{p_1} \cdots b_{p_s}]} \Gamma_{a_1 \cdots \hat{a}_{q_1} \cdots \hat{a}_{q_s} \cdots a_i b_1 \cdots \hat{b}_{p_1} \cdots \hat{b}_{p_s} \cdots b_j}. \\
A_i B_j &= \sum_{s=0}^{\min(i,j)} \binom{i}{s} \binom{j}{s} (-1)^{s(s-1)/2} \eta_{[b_1 \cdots b_s][c_1 \cdots c_s]} \\
&\quad A^{a_1 \cdots a_{i-s} b_1 \cdots b_s} B^{c_1 \cdots c_s a_{i-s+1} \cdots a_{i+j-2s}} \Gamma_{a_1 \cdots a_{i+j-2s}}, \\
A_i B_j &= \sum_{s=0}^{\min(i,j)} \binom{i}{s} \binom{j}{s} (-1)^{s(s-1)/2} \eta_{[b_1 \cdots b_s][c_1 \cdots c_s]} \\
&\quad A^{a_1 \cdots a_{i-s} b_1 \cdots b_s} B^{c_1 \cdots c_s a_{i-s+1} \cdots a_{i+j-2s}} \Gamma_{a_1 \cdots a_{i+j-2s}}, \\
A_i &= A^{a_1 \cdots a_i} \Gamma_{a_1 \cdots a_i}, \\
B_j &= B^{b_1 \cdots b_j} \Gamma_{b_1 \cdots b_j} \\
\psi' &= \exp\left(\frac{1}{4} \lambda^{ab} \Gamma_{ab}\right) \psi \\
\delta \psi &= \frac{1}{4} \lambda^{ab} \Gamma_{ab} \psi \\
\bar{\psi}_\alpha &= \psi^\beta C_{\beta\alpha} \\
\psi^\alpha &= \bar{\psi}_\beta C^{\beta\alpha} \\
C^{\alpha\gamma} C_{\gamma\beta} &= C_{\beta\gamma} C^{\gamma\alpha} = \delta_\beta^\alpha
\end{aligned}$$

$$\begin{aligned}\bar{\chi}\bar{\zeta} &= \bar{\zeta}\chi \\ \bar{\chi}S\bar{\zeta} &= -\bar{\zeta}S\chi \\ \bar{\chi}A\bar{\zeta} &= \bar{\zeta}A\chi\end{aligned}$$

$$AdS_2\times S^2\times S^4\times \Sigma$$

$$(a;q)_0:=1, (a;q)_n:=\prod_{k=0}^{n-1}(1-aq^k), (q)_n:=\prod_{k=1}^n(1-q^k),$$

$$(a;q)_\infty:=\prod_{k=0}^\infty(1-aq^k), (q)_\infty:=\prod_{k=1}^\infty(1-q^k)$$

$$\begin{aligned}&\langle W_{\mathcal{R}_1}\cdots W_{\mathcal{R}_k}\rangle^G(t;q)\\&= \frac{1}{|\mathrm{Weyl}(G)|}\frac{(q)_\infty^{2\mathrm{rank}(G)}}{\left(q^{\frac{1}{2}}t^{\pm 2};q\right)_\infty^{\mathrm{rank}(G)}}\overline{\!\!\!\!\!\prod\!\!\!\!\!}\prod_{\alpha\in\mathrm{root}(G)}ds\frac{(s^\alpha;q)_\infty(qs^\alpha;q)_\infty}{\left(q^{\frac{1}{2}}t^2s^\alpha;q\right)_\infty\left(q^{\frac{1}{2}}t^{-2}s^\alpha;q\right)_\infty}\prod_{i=1}^k\chi_{\mathcal{R}_i}\end{aligned}$$

$$\mathcal{I}^G(t;q):=\mathrm{Tr}(-1)^Fq^{J+\frac{H+C}{4}}t^{H-C}$$

$$\langle \mathcal{W}_{\mathcal{R}_1}\cdots \mathcal{W}_{\mathcal{R}_k}\rangle^G(t;q)=\frac{\langle W_{\mathcal{R}_1}\cdots W_{\mathcal{R}_k}\rangle^G(t;q)}{\mathcal{I}^G(t;q)}$$

$$\langle W_{\mathcal{R}_1}\cdots W_{\mathcal{R}_k}\rangle_{\frac{1}{2}\mathrm{BPS}}^G(\mathfrak{q})$$

$$=\frac{1}{|\mathrm{Weyl}(G)|}\frac{1}{(1-\mathfrak{q}^2)^{\mathrm{rank}(G)}}\overline{\!\!\!\!\!\prod\!\!\!\!\!}\prod_{\alpha\in\mathrm{root}(G)}ds\frac{(1-s^\alpha)}{(1-\mathfrak{q}^2s^\alpha)}\prod_{i=1}^k\chi_{\mathcal{R}_i}^{\mathfrak{g}}$$

$$\frac{1}{N!}\overline{\!\!\!\!\!\prod\!\!\!\!\!}\prod_{i=1}^N\frac{ds_i}{2\pi i s_i}\frac{\prod_{i\neq j}1-\frac{s_i}{s_j}}{\prod_{i,j}1-\mathfrak{t}\frac{s_i}{s_j}}P_\mu(s;\mathfrak{t})P_\lambda(s^{-1};\mathfrak{t})=\frac{\delta_{\mu\lambda}}{(\mathfrak{t};\mathfrak{t})_{N-l(\mu)}\prod_{j\geq 1}(\mathfrak{t};\mathfrak{t})_{m_j(\mu)}}$$

$$\frac{1}{|\mathrm{Weyl}(G)|}\overline{\!\!\!\!\!\prod\!\!\!\!\!}\prod_{\alpha\in\mathrm{root}(G)}ds(1-s^\alpha)\prod_{i=1}^k\chi_{\mathcal{R}_i}^{\mathfrak{g}}$$

$$\langle T_B T_B \rangle^G(t;q)=\sum_{v\in \mathrm{Rep}(B)}\frac{1}{|\mathrm{Weyl}(B)|}\frac{(q)_\infty^{2\mathrm{rank}(G)}}{\left(q^{\frac{1}{2}}t^{\pm 2};q\right)_\infty^{\mathrm{rank}(G)}}\oint\prod_{\alpha\in\mathrm{root}(G)}ds$$

$$\times \frac{\left(q^{\frac{|\alpha(B)|}{2}}s^\alpha;q\right)_\infty\left(q^{1+\frac{|\alpha(B)|}{2}}s^\alpha;q\right)_\infty}{\left(q^{\frac{1+|\alpha(B)|}{2}}t^2s^\alpha;q\right)_\infty\left(q^{\frac{1+|\alpha(B)|}{2}}t^{-2}s^\alpha;q\right)_\infty}Z_{\mathrm{bubb}}^{(B,v)}(t,s;q).$$

$$(z_e,z_m)\in\mathbb{Z}_2\times\mathbb{Z}_2$$

$$(z_e,z_m)=(0,0), (z_e,z_m)=(1,0)$$

$$(z_e,z_m)=(0,0), (z_e,z_m)=(0,1)$$

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$$\chi_{\mathrm{sp}}^{\mathrm{sp}(2N+1)}=\prod_{i=1}^N\left(s_i^{\frac{1}{2}}+s_i^{-\frac{1}{2}}\right)$$

$$\begin{aligned}
\chi_{\square}^{\text{so}(2N+1)} &= 1 + \sum_{i=1}^N (s_i + s_i^{-1}). \\
\chi_{\lambda}^{\text{so}(2N+1)} &= \frac{\det(s_j^{\lambda_i+N-i+1/2} - s_j^{-(\lambda_i+N-i+1/2)})}{\det(s_j^{N-i+1/2} - s_j^{-(N-i+1/2)})} \\
\langle W_{\mathcal{R}_1} \cdots W_{\mathcal{R}_k} \rangle^{SO(2N+1)} &= \int d\mu^{SO(2N+1)} \exp \left[\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) \frac{\bar{P}_n(s)^2 - \bar{P}_{2n}(s)}{2} \right] \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{so}(2N+1)}(s) \\
d\mu^{SO(2N+1)} &= \frac{1}{2^N N!} \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} (1 - s_i)(1 - s_i^{-1}) \\
&\quad \times \prod_{1 \leq i < j \leq N} (1 - s_i s_j)(1 - s_i^{-1} s_j^{-1})(1 - s_i s_j^{-1})(1 - s_i^{-1} s_j) \\
f_n(q, t) &= \frac{q^{n/2}(t^{2n} + t^{-2n}) - 2q^n}{1 - q^n} \\
P_m(s) &:= \sum_{i=1}^N (s_i^m + s_i^{-m}) \\
\bar{P}_m(s) &:= 1 + P_m(s) = 1 + \sum_{i=1}^N (s_i^m + s_i^{-m}) \\
\bar{M}_n(s) &= \frac{\bar{P}_n(s)^2 - \bar{P}_{2n}(s)}{2} = P_n(s) + \frac{P_n(s)^2 - P_{2n}(s)}{2} \\
\exp \left(\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) \bar{M}_n(s) \right) &= \sum_{\lambda} \frac{1}{z_{\lambda}} f_{\lambda}(q, t) \bar{M}_{\lambda}(s) \\
z_{\lambda} &= \prod_{i=1}^{\infty} i^{m_i} m_i!, \quad f_{\lambda}(q, t) = \prod_{i=1}^{\ell(\lambda)} f_{\lambda_i}(q, t), \quad \bar{M}_{\lambda}(s) = \prod_{i=1}^{\ell(\lambda)} \bar{M}_{\lambda_i}(s) \\
\langle W_{\mathcal{R}_1} \cdots W_{\mathcal{R}_k} \rangle^{SO(2N+1)} &= \sum_{\lambda} \frac{1}{z_{\lambda}} f_{\lambda}(q, t) \int d\mu^{SO(2N+1)} \bar{M}_{\lambda}(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{so}(2N+1)}(s) \\
&\quad \int d\mu^{SO(2N+1)} \bar{M}_{\lambda}(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{so}(2N+1)}(s) \\
\int d\mu^{SO(2N+1)} \bar{P}_{\mu}(s) &= \sum_{\nu \in R_{2N+1}(|\mu|)} \chi_{\nu}^S(\mu) + \sum_{\nu \in W_{2N+1}(|\mu|)} \chi_{\nu}^S(\mu) \\
R_n(p) &= \{ \lambda \vdash p \mid \ell(\lambda) \leq n \text{ and } \forall \lambda_i \text{ is even} \} \\
W_n(p) &= \{ \lambda \vdash p \mid \ell(\lambda) = n \text{ and } \forall \lambda_i \text{ is odd} \} \\
s_{\lambda} &= \sum_{\mu \vdash \lambda} \frac{\chi_{\lambda}^S(\mu)}{z_{\mu}} p_{\mu} \\
\bar{M}_{\lambda}(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{so}(2N+1)}(s) &= \sum_{\mu} a_{\lambda, \mathcal{R}}^{\mu} \bar{P}_{\mu}(s)
\end{aligned}$$

$$\begin{aligned}
& \int d\mu^{SO(2N+1)} \bar{M}_\lambda(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{so}(2N+1)}(s) = \sum_{\mu} a_{\lambda, \mathcal{R}}^{\mu} \int d\mu^{SO(2N+1)} \bar{P}_{\mu}(s) \\
& = \sum_{\mu} a_{\lambda, \mathcal{R}}^{\mu} \left(\sum_{v \in R_{2N+1}(|\mu|)} \chi_v^S(\mu) + \sum_{v \in W_{2N+1}(|\mu|)} \chi_v^S(\mu) \right) \\
& \quad \left(\chi_{\text{sp}}^{\text{so}(2N+1)} \right)^2 = \prod_{i=1}^N (1 + s_i)(1 + s_i^{-1}) \\
& \quad d\mu^{SO(2N+1)} \left(\chi_{\text{sp}}^{\text{so}(2N+1)} \right)^2 = d\mu^{USp(2N)} \\
& d\mu^{USp(2N)} = \frac{1}{2^N N!} \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} (1 - s_i^2)(1 - s_i^{-2}) \\
& \quad \times \prod_{1 \leq i < j \leq N} (1 - s_i s_j)(1 - s_i^{-1} s_j^{-1})(1 - s_i s_j^{-1})(1 - s_i^{-1} s_j) \\
& \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(2N+1)} = \int d\mu^{USp(2N)} \exp \left[\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) \bar{M}_n(s) \right] \\
& = \int d\mu^{USp(2N)} \exp \left[\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) \left(P_n(s) + \frac{P_n(s)^2 - P_{2n}(s)}{2} \right) \right] \\
& j^{SO(3)}(t; q) = j^{USp(2)}(t; q) \\
& = - \frac{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}}{(q; q)_{\infty}^2} \sum_{\substack{p_1, p_2 \in \mathbb{Z} \\ p_1 < p_2}} \frac{\left(q^{\frac{1}{2}} t^{-2} \right)^{p_1 + p_2 - 2}}{\left(1 - q^{p_1 - \frac{1}{2}} t^2 \right) \left(1 - q^{p_2 - \frac{1}{2}} t^2 \right)}. \\
& \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(3)}(t; q) \\
& = \frac{1}{2} \frac{(q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}} \oint \frac{ds}{2\pi i s} \frac{(s^{\pm}; q)_{\infty} (qs^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s^{\pm}; q \right)_{\infty}} \left(s^{\frac{1}{2}} + s^{-\frac{1}{2}} \right)^2 \\
& \langle T_{\left(\frac{1}{2}\right)} T_{\left(\frac{1}{2}\right)} \rangle^{USp(2)/\mathbb{Z}_2}(t; q) \\
& = \frac{(q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}} \oint \frac{ds}{2\pi i s} \frac{\left(q^{\frac{1}{2}} s^{\pm 2}; q \right)_{\infty} \left(q^{\frac{3}{2}} s^{\pm 2}; q \right)_{\infty}}{(qt^2 s^{\pm 2}; q)_{\infty} (qt^{-2} s^{\pm 2}; q)_{\infty}} \\
& \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(3)}(t; q) = \left\langle T_{\left(\frac{1}{2}\right)} T_{\left(\frac{1}{2}\right)} \right\rangle^{USp(2)/\mathbb{Z}_2}(t; q) \\
& \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(3)}(t; q) = \langle T_{\left(\frac{1}{2}\right)} T_{\left(\frac{1}{2}\right)} \rangle^{USp(2)/\mathbb{Z}_2}(t; q) \\
& = \langle W_{\square} W_{\square} \rangle^{SU(2)}(t; q) = \frac{\left(q^{\frac{1}{2}} t^{22}; q \right)_{\infty}}{(qt^{\pm 4}; q)_{\infty}} \sum_{m \in \mathbb{Z} \setminus \{0, n\}} \frac{t^{2m} - t^{-2m}}{t^2 - t^{-2}} \frac{q^{\frac{m-1}{2}}}{1 - q^m}.
\end{aligned}$$

$$\begin{aligned}
\langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \left\langle T\left(\frac{1}{2}\right) T\left(\frac{1}{2}\right) \right\rangle_{\frac{1}{2}\text{BPS}}^{\text{US}(2)/\mathbb{Z}_2}(q) \\
&= \langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{\text{SU}(2)}(q) = \frac{1+q^2}{1-q^4} = \frac{1}{1-q^2}. \\
\langle \underbrace{W_{\text{sp}} W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(t; q) &= \langle \underbrace{W_{\square} W_{\square} \cdots W_{\square}}_{2k} \rangle^{\text{SU}(2)}(t; q) \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \frac{2+3q^2+q^4}{1-q^4} \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \frac{5+9q^2+5q^4+q^6}{1-q^4} \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \frac{14+28q^2+20q^4+7q^6+q^8}{1-q^4} \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \frac{42+90q^2+75q^4+35q^6+9q^8+q^{10}}{1-q^4} \\
\langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \mathcal{J}_{\frac{1}{2}\text{BPS}}^{\text{SO}(3)}(q) \sum_{i=0}^k a_{\text{ksp}}^{\text{so}(3)}(i) q^{2i} \\
&= \frac{1}{1-q^4} \sum_{i=0}^k a_{\text{ksp}}^{\text{so}(3)}(i) q^{2i} \\
a_{\text{ksp}}^{\text{so}(3)}(i) &= (2i+1) \frac{(2k)!}{(k-i)!(k+i+1)!} \\
&= C_{k+i+1, 2i+1} \\
C_{n,m} &= \frac{m}{n} \binom{2n-m-1}{n-1} \\
C_k &= \frac{1}{k+1} \binom{2k}{k} \\
&= \prod_{1 \leq i \leq j \leq k-1} \frac{i+j+2}{i+j} \\
\frac{1-\sqrt{1-4x}}{2x} &= \sum_{k=0}^{\infty} a_{\text{ksp}}^{\text{so}(3)}(0) x^k \\
\frac{1}{x^{i+1}} \left(\frac{1-\sqrt{1-4x}}{2} \right)^{2i+1} &= \sum_{k=0}^{\infty} a_{\text{ksp}}^{\text{so}(3)}(i) x^k \\
\sum_{k=0}^{\infty} x^k \langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(3)}(q) &= \frac{1}{1-q^4} \cdot \frac{2(1-\sqrt{1-4x})}{4x - (1-\sqrt{1-4x})^2 q^2} \\
\langle W_{\square} W_{\square} \rangle^{\text{SO}(3)}(t; q) &= \frac{1}{2} \frac{(q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}} \oint \frac{ds}{2\pi i s} \frac{(s^{\pm}; q)_{\infty} (qs^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s^{\pm}; q \right)_{\infty}} (1+s+s^{-1})^2
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle^{SO(3)}(t; q) &= \langle W_{\square\square} W_{\square\square} \rangle^{SU(2)}(t; q) \\
&= \frac{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}}{(qt^{\pm 4}; q)_{\infty}} \left[\frac{3}{2} \sum_{m \in \mathbb{Z} \setminus \{0\}} \left(\frac{t^{2m} - t^{-2m}}{t - t^{-2}} \frac{q^{\frac{m-1}{2}}}{1 - q^m} \right) - \frac{2}{1 - q} - \frac{q^{\frac{1}{2}}(t^2 + t^{-2})}{1 - q^2} \right] \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \langle W_{\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SU(2)}(q) \\
&= \frac{1 + q^2 + q^4}{1 - q^4} \\
&= \frac{1 - q^6}{(1 - q^2)(1 - q^4)} \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle^{SO(3)}(t; q) &= \sum_{i=0}^k \binom{k}{i} (-1)^i \langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2(k-i)} \rangle^{SO(3)}(t; q) \\
\langle W_{\square} \rangle^{SO(3)}(t; q) &= \langle W_{\text{sp}} W_{\text{sp}} \rangle^{SO(3)}(t; q) - J^{SO(3)}(t; q) \\
&= - \frac{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}}{(q; q)_{\infty}^2} \sum_{\substack{p_1, p_2 \in \mathbb{Z} \\ p_1 + 1 < p_2}} \frac{\left(q^{\frac{1}{2}} t^{-2}\right)^{p_1 + p_2 - 2}}{\left(1 - q^{p_1 - \frac{1}{2}} t^2\right) \left(1 - q^{p_2 - \frac{1}{2}} t^2\right)} \\
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - J_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \frac{q^2}{1 - q^4}. \\
\langle W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - 3 \langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&\quad + 3 \langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - J_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \frac{1 + 3q^2 + 2q^4 + q^6}{1 - q^4}. \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \frac{3 + 6q^2 + 6q^4 + 3q^6 + q^8}{1 - q^4} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \frac{6 + 15q^2 + 15q^4 + 10q^6 + 4q^8 + q^{10}}{1 - q^4} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \frac{15 + 36q^2 + 40q^4 + 29q^6 + 15q^8 + 5q^{10} + q^{12}}{1 - q^4} \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \frac{\sum_{i=0}^k a_k^{SO(3)}(i) q^{2i}}{1 - q^4} \\
a_k^{SO(3)}(i) &= c_k^{(i)} - c_k^{(i+1)} \\
(1 + x + x^2)^n &= \sum_{i=-k}^k c_k^{(i)} x^{k+i}
\end{aligned}$$

$$\begin{aligned}
R_n &= \sum_{i=0}^n (-1)^{n-i} \binom{n}{i} C_i \\
\sum_{n=1}^{\infty} R_n x^n &= \frac{1}{2x} \left(1 - \frac{\sqrt{1-3x}}{\sqrt{1+x}} \right) \\
\langle W_{(k)} W_{(k)} \rangle^{SO(3)}(t; q) &= \frac{\left(q^{\frac{1}{2}t^{\pm 2}}; q \right)_{\infty}}{(qt^{\pm 4}; q)_{\infty}} \left[\frac{2k+1}{2} \sum_{m \in \mathbb{Z} \setminus \{0\}} \left(\frac{t^{2m} - t^{-2m}}{t - t^{-2}} \frac{q^{\frac{m-1}{2}}}{1 - q^m} \right) \right. \\
&\quad \left. - \sum_{m=1}^{2k} (2k - m + 1) \left(\frac{t^{2m} - t^{-2m}}{t - t^{-2}} \frac{q^{\frac{m-1}{2}}}{1 - q^m} \right) \right] \\
\langle W_{(k)} W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \frac{1 + q^2 + \dots + q^{4k}}{1 - q^4} \\
&= \frac{1 - q^{4k+2}}{(1 - q^2)(1 - q^4)} \\
\langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \frac{1}{(1 - q^2)(1 - q^4)} \\
\langle \underbrace{W_{\square} \dots W_{\square}}_k \rangle^{SO(3)}(t; q) &= \sum_{k_1+k_2+k_3=k} \binom{k}{k_1, k_2, k_3} (-3)^{k_2} \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_{4k_1+2k_2} \rangle^{SO(3)}(t; q) \\
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - 3 \langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) + \mathcal{I}_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \frac{q^4}{1 - q^4}, \\
\langle W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_{12} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - 9 \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_{10} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&\quad + 30 \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_8 \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - 45 \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_6 \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&\quad + 30 \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_4 \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - 9 \langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) + \mathcal{I}_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \frac{1 + 3q^2 + 5q^4 + 4q^6 + 3q^8 + 2q^{10} + q^{12}}{1 - q^4}. \\
\langle \underbrace{W_{\square} \dots W_{\square}}_k \rangle^{SO(3)}(t; q) &= \sum_{k_1+k_2+k_3+k_4=k} \binom{k}{k_1, k_2, k_3, k_4} (-1)^{k_2+k_4} 5^{k_2} 6^{k_3} \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_{6k_2+4k_1+2k_2} \rangle^{SO(3)}(t; q). \\
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) &= \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_6 \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - 5 \langle \underbrace{W_{\text{sp}} \dots W_{\text{sp}}}_4 \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) + 6 \langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - \mathcal{I}_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \frac{q^6}{1 - q^4},
\end{aligned}$$

$$\begin{aligned}
& \langle W_{\square}^{W_{\square\square}W_{\square\square}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{18}(q) - 15 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{16}(q) + 93 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{14}(q) \\
&- 308 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{12}(q) + 588 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{10}(q) - 651 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{8}(q) \\
&+ 399 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{6}(q) - 123 \underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}}_{4}(q) + 18 \langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) - \mathcal{I}_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
&= \frac{1 + 3q^2 + 5q^4 + 7q^6 + 6q^8 + 5q^{10} + 4q^{12} + 3q^{14} + 2q^{16} + q^{18}}{1 - q^4}.
\end{aligned}$$

$$\chi_{(k)}^{\text{so}(3)} = \sum_{n=0}^k (-1)^n \binom{2k-n}{n} \chi_{\text{sp}}^{\text{so}(3)2k-2n}$$

$$\langle W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) = \frac{q^{2k}}{1 - q^4}$$

$$\langle W_{(k)} W_{(l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) = \frac{q^{2(l-k)}(1 - q^{4k+2})}{(1 - q^2)(1 - q^4)}$$

$$\langle W_{(k)} W_{(k)} W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) = \frac{1 + q^2 - 3q^{2k+2} + q^{6k+4}}{(1 - q^2)^2(1 - q^4)}$$

$$\langle W_{(\infty)} W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) = \frac{1}{(1 - q^2)^3}$$

$$\langle W_{(k)} W_{(k)} W_{(k)} W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) = \frac{2k + 1 - 3q^2 - (2k + 1)q^4 + 4q^{4k+4} - q^{8k+6}}{(1 - q^2)^3(1 - q^4)}$$

$$\begin{aligned}
& \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(5)}(t; q) \\
&= \frac{1}{8} \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm}; q\right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \frac{(s_i^{\pm}; q)_{\infty} (q s_i^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm}; q\right)_{\infty}} \\
&\times \frac{(s_1^{\pm} s_2^{\mp}; q)_{\infty} (s_1^{\pm} s_2^{\pm}; q)_{\infty} (q s_1^{\pm} s_2^{\mp}; q)_{\infty} (q s_1^{\pm} s_2^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\pm}; q\right)_{\infty}} \prod_{i=1}^2 \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}}\right)^2 \\
&\left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{USp(4)/\mathbb{Z}_2}(t; q)
\end{aligned}$$

$$\begin{aligned}
&= \frac{1}{2} \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \frac{\left(q^{\frac{1}{2}} s_i^{\pm 2}; q\right)_{\infty} \left(q^{\frac{3}{2}} s_i^{\pm 2}; q\right)_{\infty}}{\left(q t^2 s_i^{\pm 2}; q\right)_{\infty} \left(q t^{-2} s_i^{\pm 2}; q\right)_{\infty}} \\
&\times \frac{(s_1^{\pm} s_2^{\mp}; q)_{\infty} \left(q^{\frac{1}{2}} s_1^{\pm} s_2^{\pm}; q\right)_{\infty} (q s_1^{\pm} s_2^{\mp}; q)_{\infty} \left(q^{\frac{3}{2}} s_1^{\pm} s_2^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q t^2 s_1^{\pm} s_2^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q t^{-2} s_1^{\pm} s_2^{\pm}; q\right)_{\infty}} \\
&\langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(5)}(t; q) = \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{USp(4)/\mathbb{Z}_2}(t; q)
\end{aligned}$$

$$\begin{aligned}
\langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(5)}(q) &= \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle_{\frac{1}{2}\text{BPS}}^{USp(4)/\mathbb{Z}_2}(q) \\
&= \frac{1 + q^2 + q^4 + q^6}{(1 - q^4)(1 - q^8)} \\
&= \frac{1}{(1 - q^2)(1 - q^4)} \\
\langle \underbrace{V_{\text{sp}} \cdots W_{\text{sp}}}_4 \rangle_{\frac{1}{2}\text{BPS}}^{SSin(5)}(q) &= \frac{3 + 6q^2 + 8q^4 + 9q^6 + 6q^8 + 3q^{10} + q^{12}}{(1 - q^4)(1 - q^8)} \\
\langle \underbrace{S_{\text{sp}} \cdots W_{\text{sp}}}_6 \rangle_{\frac{1}{2}\text{BPS}}^{SSin(5)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)} (14 + 40q^2 + 66q^4 + 85q^6 \\
&\quad + 81q^8 + 59q^{10} + 34q^{12} + 15q^{14} + 5q^{16} + q^{18}) \\
\langle \underbrace{S_{\text{sp}} \cdots W_{\text{sp}}}_8 \rangle_{\frac{1}{2}\text{BPS}}^{Sin(5)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)} (84 + 300q^2 + 581q^4 + 840q^6 + 945q^8 + 842q^{10} \\
&\quad + 616q^{12} + 378q^{14} + 195q^{16} + 83q^{18} + 28q^{20} + 7q^{22} + q^{24}) \\
\langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{Sp(5)}(q) &= \frac{\sum_{i=0}^{3k} a_{k\text{ sp}}^{\text{so}(5)}(i) q^{2i}}{(1 - q^4)(1 - q^8)} \\
a_{k\text{ sp}}^{a_{\text{sp}}(0)} &= C_k C_{k+2} - C_{k+1}^2 \\
&= \frac{24(2k+1)!(2k-1)!}{(k-1)!k!(k+2)!(k+3)!} \\
&= \prod_{1 \leq i \leq j \leq k-1} \frac{i+j+4}{i+j}, \\
{}_3F_2\left(1, \frac{1}{2}, \frac{3}{2}; 3, 4; 16x\right) &= \sum_{k=0}^{\infty} a_{k\text{ sp}}^{\text{so}(5)}(0) x^k \\
{}_pF_q(a_1, \dots, a_p; b_1, \dots, b_q; z) &= \sum_{k=0}^{\infty} \frac{(a_1)_k (a_2)_k \cdots (a_p)_k}{(b_1)_k (b_2)_k \cdots (b_q)_k} \frac{z^k}{k!} \\
a_{k\text{ sp}}^{\text{so}(5)}(1) &= \frac{60(2k)!(2k+2)!}{(k-1)!k!(k+3)!(k+4)!} \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{\sum_{i=0}^{2k} a_k^{\text{son}(5)}(i) q^{2i}}{(1 - q^4)(1 - q^8)}
\end{aligned}$$



$$\begin{aligned}
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{q^4}{(1-q^4)(1-q^8)}, \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1+q^2+q^4+q^6+q^8}{(1-q^4)(1-q^8)} \\
&= \frac{1-q^{10}}{(1-q^2)(1-q^4)(1-q^8)}, \\
\langle W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{q^2+3q^4+3q^6+3q^8+2q^{10}+q^{12}}{(1-q^4)(1-q^8)} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1-q^4)(1-q^8)} (3+3q^2+9q^4+15q^6+12q^8 \\
&\quad +12q^{10}+6q^{12}+q^{16}) \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1-q^4)(1-q^8)} (1+10q^2+24q^4+36q^6+44q^8 \\
&\quad +41q^{10}+31q^{12}+19q^{14}+10q^{16}+4q^{18}+q^{20}) \\
a_{k\square}^{so(5)}(0) &= \sum_{i=0}^{\lfloor \frac{k}{2} \rfloor} C_i C_{i+1} \binom{k}{2i} - \sum_{i=0}^{\lfloor \frac{k+1}{2} \rfloor} C_i^2 \binom{k}{2i-1} \\
&= -k {}_3F_2 \left(\frac{3}{2}, \frac{1}{2} - \frac{k}{2}; 3, 3; 16 \right) + {}_3F_2 \left(\frac{3}{2}, \frac{1}{2} - \frac{k}{2}; 2, 3; 16 \right) \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{\sum_{i=0}^{3k} a_k^{\mathfrak{N}^{(5)}}(i) q^{2i}}{(1-q^4)(1-q^8)} \\
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{q^2+q^6}{(1-q^4)(1-q^8)} \\
&= \frac{q^2}{(1-q^4)^2}, \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1+q^2+3q^4+2q^6+3q^8+q^{10}+q^{12}}{(1-q^4)(1-q^8)} \\
&= \frac{(1-q^6)(1-q^8)}{(1-q^2)(1-q^4)^3}, \\
\langle W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1-q^4)(1-q^8)} (1+6q^2+9q^4+16q^6+15q^8 \\
&\quad +15q^{10}+9q^{12}+6q^{14}+2q^{16}+q^{18}) \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1-q^4)(1-q^8)} (6+22q^2+54q^4+82q^6+15q^8 \\
&\quad +15q^{10}+9q^{12}+6q^{14}+2q^{16}+q^{18}) \\
\langle \underbrace{W_{(2)} \cdots W_{(2)}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{\sum_{i=0}^{4k} a_k^{\text{son}(5)}(i) q^{2i}}{(1-q^4)(1-q^8)}
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{q^4 + q^8}{(1 - q^4)(1 - q^8)} \\
&= \frac{q^4}{(1 - q^4)^2}, \\
\langle W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1 + q^2 + 2q^4 + 2q^6 + 3q^8 + 2q^{10} + 3q^{12} + q^{14} + q^{16}}{(1 - q^4)(1 - q^8)} \\
\langle W_{\square\square} W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)} \\
&\quad \times (1 + 3q^2 + 9q^4 + 13q^6 + 20q^8 + 21q^{10} \\
&\quad + 22q^{12} + 18q^{14} + 15q^{16} + 9q^{18} + 6q^{20} + 2q^{22} + q^{24}) \\
\langle W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{q^{4k-4} + q^{4k}}{(1 - q^4)(1 - q^8)} \\
&= \frac{q^{4k-4}}{(1 - q^4)^2}. \\
\langle W_{(k)} W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{\sum_{i=0}^{8k} a_2^{\text{so}(5)}(i) q^{2i}}{(1 - q^4)(1 - q^8)} \\
\langle W_{\square\square\square} W_{\square\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)} (1 + q^2 + 2q^4 + 3q^6 + 4q^8 + 4q^{10} \\
&\quad + 6q^{12} + 5q^{14} + 5q^{16} + 4q^{18} + 3q^{20} + q^{22} + q^{24}) \\
\langle W_{\square\square\square} W_{\square\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)} (1 + q^2 + 2q^4 + 3q^6 + 5q^8 + 5q^{10} \\
&\quad + 8q^{12} + 8q^{14} + 10q^{16} + 9q^{18} + 10q^{20} + 7q^{22} + 7q^{24} \\
&\quad + 4q^{26} + 3q^{28} + q^{30} + q^{32}) \\
\langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= 1 + q^2 + 3q^4 + 4q^6 + 9q^8 + 11q^{10} + 21q^{12} + 26q^{14} + 44q^{16} + 54q^{18} + 84q^{20} + \dots \\
\langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) &= \frac{1 - q^{24}}{(1 - q^2)(1 - q^4)^2(1 - q^6)(1 - q^8)^2(1 - q^{12})} \\
\langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(7)}(t; q) &= \frac{1}{48} \frac{(q)_\infty^6}{\left(q^{\frac{1}{2}} t t^\pm; q\right)_\infty^3} \int \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \frac{(s_i^\pm; q)_\infty (q s_i^\pm; q)_\infty}{\left(q^{\frac{1}{2}} t^2 s_i^\pm; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm; q\right)_\infty} \\
&\quad \times \prod_{i < j} \frac{(s_i^\pm s_j^\mp; q)_\infty (s_i^\pm s_j^\pm; q)_\infty (q s_i^\pm s_j^\mp; q)_\infty (q s_i^\pm s_j^\pm; q)_\infty}{\left(q^{\frac{1}{2}} t^2 s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1}{2}} t^2 s_i^\pm s_j^\pm; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm s_j^\pm; q\right)_\infty} \prod_{i=1}^3 \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}}\right)^2.
\end{aligned}$$

$$\begin{aligned}
& \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{USp(6)/\mathbb{Z}_2} (t; q) \\
&= \frac{1}{6} \frac{(q)_\infty^6}{\left(q^{\frac{1}{2}} \pm^{\pm 2}; q\right)_\infty^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \frac{\left(q^{\frac{1}{2}} s_i^{\pm 2}; q\right)_\infty \left(q^{\frac{3}{2}} s_i^{\pm 2}; q\right)_\infty}{\left(q t^2 s_i^{\pm 2}; q\right)_\infty \left(q t^{-2} s_i^{\pm 2}; q\right)_\infty} \\
&\times \prod_{i < j} \frac{\left(s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1}{2}} s_i^\pm s_j^\pm; q\right)_\infty \left(q s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{3}{2}} s_i^\pm s_j^\pm; q\right)_\infty}{\left(q^{\frac{1}{2}} t^2 s_i^\pm s_j^\mp; q\right)_\infty \left(q t^2 s_i^\pm s_j^\pm; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm s_j^\mp; q\right)_\infty \left(q t^{-2} s_i^\pm s_j^\pm; q\right)_\infty} \\
&\langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(7)}(q) = \langle T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)/\mathbb{Z}_2}(q) \\
&= \frac{1 + q^2 + q^4 + 2q^6 + q^8 + q^{10} + q^{12}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&= \frac{1}{(1 - q^2)(1 - q^4)(1 - q^6)} \cdot \frac{\sum_{i=0}^{6k} a_{k\text{sp}}^{\text{so}(7)}(i) q^{2i}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPin}(7)}^{\text{SPS}}(q) = \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(7)}(q) = \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} (4 + 9q^2 + 15q^4 \\
&\quad + 25q^6 + 29q^8 + 32q^{10} + 33q^{12} + 26q^{14} + 20q^{16} \\
&\quad + 13q^{18} + 6q^{20} + 3q^{22} + q^{24}), \\
&\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(7)}(q) = \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} (30 + 105q^2 + 235q^4 \\
&\quad + 435q^6 + 650q^8 + 855q^{10} + 1010q^{12} + 1055q^{14} \\
&\quad + 1006q^{16} + 865q^{18} + 665q^{20} + 470q^{22} + 299q^{24} \\
&\quad + 170q^{26} + 89q^{28} + 40q^{30} + 15q^{32} + 5q^{34} + q^{36}). \\
&a_{k\text{sp}}^{\text{so}(7)}(0) = \prod_{1 \leq i \leq j \leq k-1} \frac{i+j+6}{i+j} \\
&{}_4F_3\left(1, \frac{1}{2}, \frac{5}{2}, \frac{3}{2}; 4, 5, 6; 64x\right) = \sum_{k=0}^{\infty} a_{k\text{sp}}^{\text{so}(7)}(0) x^k \\
&\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{\sum_{i=0}^{3k} a_k^{\text{son}(7)}(i) q^{2i}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{q^6}{(1 - q^4)(1 - q^8)(1 - q^{12})}, \\
&\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{1 + q^2 + q^4 + q^6 + q^8 + q^{10} + q^{12}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&= \frac{1 - q^{14}}{(1 - q^2)(1 - q^4)(1 - q^8)(1 - q^{12})}, \\
&\langle W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{q^4 + 3q^6 + 3q^8 + 3q^{10} + 3q^{12} + 3q^{14} + 2q^{16} + q^{18}}{(1 - q^4)(1 - q^8)(1 - q^{12})}
\end{aligned}$$

$$\begin{aligned}
& \left\langle \underbrace{W_{\square} \cdots W_{\square}}_k \right\rangle_{\frac{1}{2}\text{SPS}}^{SO(7)}(q) = \frac{\sum_{i=0}^{5k} a_{k\square}^{so(7)}(i) q^{2i}}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \left\langle \underbrace{W_{\square} \cdots W_{\square}}_k \right\rangle_{\frac{1}{2}\text{SPS}}^{SO(7)}(q) = \frac{\sum_{i=0}^{6k} a_{k\square}^{so(7)}(i) q^{2i}}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \left\langle W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{q^2 + q^6 + q^{10}}{(1-q^4)(1-q^8)(1-q^{12})}, \\
& \left\langle W_{\square} W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{1}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \quad \times (1 + q^2 + 3q^4 + 2q^6 + 5q^8 + 3q^{10} + 5q^{12} + 2q^{14} + 3q^{16} + q^{18} + q^{20}). \\
& \left\langle W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{q^4 + q^8 + q^{12}}{(1-q^4)(1-q^8)(1-q^{12})}, \\
& \left\langle W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{1}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \quad \times (1 + q^2 + 3q^4 + 4q^6 + 7q^8 + 6q^{10} + 9q^{12} + 6q^{14} + 7q^{16} + 4q^{18} + 3q^{20} + q^{22} + q^{24}). \\
& \left\langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{\sum_{i=0}^{3lk} a_{k(l)}^{50}(7)(i) q^{2i}}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \left\langle W_{\square\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{q^4 + q^8 + q^{12}}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \left\langle W_{\square\square} W_{\square\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) = \frac{1}{(1-q^4)(1-q^8)(1-q^{12})} \\
& \quad \times (1 + q^2 + 2q^4 + 2q^6 + 4q^8 + 3q^{10} + 5q^{12} + 3q^{14} + 5q^{16} + 2q^{18} + 3q^{20} + q^{22} + q^{24}) \\
& \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(2N+1)}(t; q) \\
& = \frac{1}{2^N N!} \frac{(q)_{\infty}^{2N}}{\left(\frac{1}{q^2} t^{\pm}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \frac{(s_i^{\pm}; q)_{\infty} (qs_i^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm}; q\right)_{\infty}} \\
& \quad \times \prod_{i < j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (qs_i^{\pm} s_j^{\mp}; q)_{\infty} (qs_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \prod_{i=1}^N \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}}\right)^2.
\end{aligned}$$



$$\begin{aligned}
& \langle T_{\left(\frac{1}{2}\right)^N} T_{\left(\frac{1}{2}\right)^N} \rangle^{USp(2N)/\mathbb{Z}_2}(t; q) \\
&= \frac{1}{N!} \frac{(q)_\infty^{2N}}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_\infty^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \frac{\left(q^{\frac{1}{2}} s_i^{\pm 2}; q\right)_\infty \left(q^{\frac{3}{2}} s_i^{\pm 2}; q\right)_\infty^{\pm 2}}{(qt^{-2} s_i^{\pm 2}; q)_\infty} \\
&\times \prod_{i < j} \frac{(s_i^\pm s_j^\mp; q)_\infty \left(q^{\frac{1}{2}} s_i^\pm s_j^\pm; q\right)_\infty (qs_i^\pm s_j^\mp; q)_\infty \left(q^{\frac{3}{2}} s_i^\pm s_j^\pm; q\right)_\infty}{(t^2 s_i^\pm s_j^\mp; q)_\infty (qt^2 s_i^\pm s_j^\pm; q)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm s_j^\mp; q\right)_\infty (qt^{-2} s_i^\pm s_j^\pm; q)_\infty}. \\
&\langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N+1)}(q) = \langle T_{\left(\frac{1}{2}\right)^N} T_{\left(\frac{1}{2}\right)^N} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)/\mathbb{Z}_2}(q) \\
&= \prod_{n=1}^N \frac{1}{(1 - q^{2n})}. \\
&j_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N+1)}(q) = \prod_{n=1}^N \frac{1}{1 - q^{4n}} \\
&\langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N+1)}(q) = \prod_{n=1}^N (1 + q^{2n}) \\
&\langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N+1)}(q) = \frac{\sum_{i=0}^{\frac{N(N+1)k}{2}} a_{k\text{ sp}}^{\text{so}(2N+1)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
&a_{k\text{ sp}}^{\text{so}(2N+1)}(0) = \det(C_{2N-i-j+k}) \\
&= \prod_{1 \leq i \leq j \leq k-1} \frac{i+j+2N}{i+j}, \\
&\langle \underbrace{W_\square \cdots W_\square}_k \rangle_{\frac{1}{2}\text{BPS}}^{S(2N+1)}(q) = \frac{\sum_{i=0}^{Nk} a_{k\square}^{\text{so}(2N+1)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
&\langle W_\square \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) = \frac{q^{2N}}{\prod_{n=1}^N (1 - q^{4n})}, \\
&\langle W_\square W_\square \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) = \frac{1 + q^2 + q^4 + \cdots + q^{4N}}{\prod_{n=1}^N (1 - q^{4n})} \\
&= \frac{1 - q^{4N+2}}{(1 - q^2) \prod_{n=1}^N (1 - q^{4n})} \\
&\langle \mathcal{W}_\square \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) = q^{2N}, \\
&\langle \mathcal{W}_\square \mathcal{W}_\square \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) = \frac{1 - q^{4N+2}}{1 - q^2} \\
&\langle \mathcal{W}_\square \mathcal{W}_\square \rangle_{\frac{1}{2}\text{BPS}, c}^{SO(2N+1)}(q) = \langle \mathcal{W}_\square \mathcal{W}_\square \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) - \langle \mathcal{W}_\square \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q)^2 \\
&= \frac{1 - q^{4N}}{1 - q^2}.
\end{aligned}$$

$$\begin{aligned}
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS } SO(2N+1)}(q) &= \frac{\sum_{i=0}^{(2N-1)k} a_{k \square \square}^{\text{so}(2N+1)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) &= \frac{q^2 + q^6 + \cdots + q^{4N-2}}{\prod_{n=1}^N (1 - q^{4n})} \\
&= \frac{q^2 (1 - q^{4N})}{(1 - q^4) \prod_{n=1}^N (1 - q^{4n})}. \\
\langle \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) &= \frac{q^2 (1 - q^{4N})}{1 - q^4} \\
\langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) &= \frac{\sum_{i=0}^{Nlk} a_{k(l)}^{\text{so}(2N+1)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
(z_e, z_m) &\in \mathbb{Z}_2 \times \mathbb{Z}_2 \\
(z_e, z_m) &= (0,0), (z_e, z_m) = (1,0) \\
(z_e, z_m) &= (0,0), (z_e, z_m) = (0,1) \\
(z_e, z_m) &= (0,0), (z_e, z_m) = (1,1) \\
\chi_{\square}^{\text{usp}(2N)} &= \sum_{i=1}^N (s_i + s_i^{-1}) \\
\chi_{\square}^{\mu\text{sp}(2N)} &= \frac{\det(s_j^{\lambda_i + N - i + 1} - s_j^{-\lambda_i - N + i - 1})}{\det(s_j^{N - i + 1} - s_j^{-N + i - 1})} \\
\langle W_{\mathcal{R}_1} \cdots W_{\mathcal{R}_k} \rangle^{USp(2N)} &= \int d\mu^{USp(2N)} \exp \left(\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) L_n(s) \right) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{usp}(2N)}(s), \\
L_n(s) &= \frac{P_n(s)^2 + P_{2n}(s)}{2} \\
\exp \left(\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) L_n(s) \right) &= \sum_{\lambda} \frac{1}{z_{\lambda}} f_{\lambda}(q, t) L_{\lambda}(s) \\
\langle W_{\mathcal{R}_1} \cdots W_{\mathcal{R}_k} \rangle^{USp(2N)} &= \sum_{\lambda} \frac{1}{z_{\lambda}} f_{\lambda}(q, t) \int d\mu^{USp(2N)} L_{\lambda}(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{usp}(2N)}(s) \\
L_{\lambda}(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{usp}(2N)}(s) &= \sum_{\mu} b_{\lambda, \mathcal{R}}^{\mu} P_{\mu}(s) \\
\int d\mu^{USp(2N)} P_{\mu}(s) &= \sum_{v \in R_{2N}^c(|\mu|)} \chi_v^S(\mu), \\
R_n^c(p) &= \{ \lambda \vdash p \mid \ell(\lambda) \leq n \text{ and } \forall \lambda'_i \text{ is even} \} \\
\int d\mu^{USp(2N)} L_{\lambda}(s) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\mu\text{sp}(2N)}(s) &= \sum_{\mu} b_{\lambda, \mathcal{R}}^{\mu} \int d\mu^{USp(2N)} P_{\mu}(s) \\
&= \sum_{\mu} b_{\lambda, \mathcal{R}}^{\mu} \sum_{v \in R_{2N}^c(|\mu|)} \chi_v^S(\mu). \\
\langle W_{\mathcal{R}_1} \cdots W_{\mathcal{R}_k} \rangle^{USp(2N)} &= \sum_{\lambda} \frac{1}{z_{\lambda}} f_{\lambda}(q, t) \sum_{\mu} b_{\lambda, \mathcal{R}}^{\mu} \sum_{v \in R_{2N}^c(|\mu|)} \chi_v^S(\mu)
\end{aligned}$$

$$\begin{aligned}
& \langle W_{\square} W_{\square} \rangle^{USp(2)}(t; q) \\
&= \frac{1}{2} \frac{(q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \oint \frac{ds}{2\pi i s} \frac{(s^{\pm 2}; q)_{\infty} (qs^{\pm 2}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s^{\pm 2}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s^{\pm 2}; q\right)_{\infty}} (s + s^{-1})^2 \\
& \quad \langle T_{(1)} T_{(1)} \rangle^{SO(3)}(t; q) \\
&= \frac{(q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \oint \frac{ds}{2\pi i s} \frac{\left(q^{\frac{1}{2}} s^{\pm}; q\right)_{\infty} \left(q^{\frac{3}{2}} s^{\pm}; q\right)_{\infty}}{(qt^2 s^{\pm}; q)_{\infty} (qt^{-2} s^{\pm}; q)_{\infty}} \\
& \quad \langle W_{\square} W_{\square} \rangle^{USp(2)}(t; q) \\
&= \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(3)}(t; q) \\
&= \langle T_{(1)} T_{(1)} \rangle^{SO(3)}(t; q) = \left\langle T_{\left(\frac{1}{2}\right)} T_{\left(\frac{1}{2}\right)} \right\rangle^{USp(2)/\mathbb{Z}_2}(t; q) \\
& \quad \langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2)}(q) = \langle T_{(1)} T_{(1)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(3)}(q) \\
& \quad = \frac{1 + q^2}{1 - q^4} \\
& \quad = \frac{1}{1 - q^2}
\end{aligned}$$

$$\begin{aligned}
& \langle W_{\square\square} W_{\square\square} \rangle^{USp(2)}(t; q) \\
&= \frac{1}{2} \frac{(q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \oint \frac{ds}{2\pi i s} \frac{(s^{\pm 2}; q)_{\infty} (qs^{\pm 2}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s^{\pm 2}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s^{\pm 2}; q\right)_{\infty}} (1 + s^2 + s^{-2})^2 \\
& \quad \langle W_{(2k)} \rangle^{USp(2)}(t; q) = \langle W_{(k)} \rangle^{SO(3)}(t; q). \\
& \quad \langle W_{(2k)} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2)}(q) = \frac{q^{2k}}{(1 - q^4)} \\
& \quad \langle W_{(k)} W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2)}(q) = \frac{1 - q^{2k+2}}{(1 - q^2)(1 - q^4)} \\
& \quad \langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2)}(q) = \frac{1}{(1 - q^2)(1 - q^4)}
\end{aligned}$$

$$\begin{aligned}
& \langle W_{\square} W_{\square} \rangle^{USp(4)}(t; q) \\
&= \frac{1}{8} \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \frac{(s_i^{\pm 2}; q)_{\infty} (qs_i^{\pm 2}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm 2}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm 2}; q\right)_{\infty}} \\
& \quad \times \frac{(s_1^{\pm} s_2^{\mp}; q)_{\infty} (s_1^{\pm} s_2^{\pm}; q)_{\infty} (qs_1^{\pm} s_2^{\mp}; q)_{\infty} (qs_1^{\pm} s_2^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\pm}; q\right)_{\infty}} \left[\sum_{i=1}^2 (s_i + s_i^{-1}) \right]^2,
\end{aligned}$$

$$\begin{aligned}
& \langle T_{(1,0)} T_{(1,0)} \rangle^{SO(5)}(t; q) \\
&= \frac{1}{2} \frac{(q)_\infty^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_\infty^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \frac{\left(q^{\frac{1}{2}} s_1^\pm; q\right)_\infty \left(s_2^\pm; q\right)_\infty \left(q^{\frac{3}{2}} s_1^\pm; q\right)_\infty \left(q^{\frac{1}{2}} t^2 s_2^\pm; q s_\infty^\pm; q\right)_\infty \left(q t^{-2} s_1^\pm; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_2^\pm; q\right)_\infty}{\left(q^{\frac{1}{2}} s_1^\pm s_2^\mp; q\right)_\infty \left(q^{\frac{1}{2}} s_1^\pm s_2^\pm; q\right)_\infty \left(q^{\frac{3}{2}} s_1^\pm s_2^\mp; q\right)_\infty \left(q^{\frac{3}{2}} s_1^\pm s_2^\pm; q\right)_\infty} \\
&\times \frac{\left(q t^2 s_1^\pm s_2^\mp; q\right)_\infty \left(q t^2 s_1^\pm s_2^\pm; q\right)_\infty \left(q t^{-2} s_1^\pm s_2^\mp; q\right)_\infty \left(q t^{-2} s_1^\pm s_2^\pm; q\right)_\infty}{\left(q t^2 s_1^\pm s_2^\mp; q\right)_\infty \left(q t^2 s_1^\pm s_2^\pm; q\right)_\infty \left(q t^{-2} s_1^\pm s_2^\mp; q\right)_\infty \left(q t^{-2} s_1^\pm s_2^\pm; q\right)_\infty}.
\end{aligned}$$

$$\begin{aligned}
& \langle W_\square W_\square \rangle^{USp(4)}(t; q) = \langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(5)}(t; q) \\
&= \langle T_{(1,0)} T_{(1,0)} \rangle^{SO(5)}(t; q) = \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{USp(4)/\mathbb{Z}_2}(t; q).
\end{aligned}$$

$$\begin{aligned}
& \langle W_\square W_\square \rangle_{\frac{1}{2}\text{BPS}}^{USp(4)}(q) = \langle T_{(1,0)} T_{(1,0)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) \\
&= \frac{1}{(1 - q^2)(1 - q^4)}, \\
& \langle W_\square \rangle^{USp(4)}(t; q) = \langle W_\square \rangle^{SO(5)}(t; q) \\
& \langle \underbrace{W_\square \cdots W_\square}_k \rangle^{USp(4)}(t; q) = \langle \underbrace{W_\square \cdots W_\square}_k \rangle^{SO(5)}(t; q) \\
& \langle \underbrace{W_{\square\square} \cdots W_{\square\square}}_k \rangle^{USp(4)}(t; q) = \langle \underbrace{W_{\square\square} \cdots W_{\square\square}}_k \rangle^{SO(5)}(t; q) \\
& \langle \underbrace{W_{(2l)} \cdots W_{(2l)}}_k \rangle^{USp(4)}(t; q) = \langle \underbrace{W_{(l^2)} \cdots W_{(l^2)}}_k \rangle^{SO(5)}(t; q) \\
& \langle \underbrace{W_{(2l)} \cdots W_{(2l)}}_k \rangle_{\frac{1}{2}\text{BPS}}^{USp(4)}(q) = \frac{\sum_{i=0}^{3lk} a_k^{\text{usp}(4)}(i) q^{2i}}{(1 - q^4)(1 - q^8)} \\
& \langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{USp(4)}(q) = \langle W_{(\infty^2)} W_{(\infty^2)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) \\
&= 1 + q^2 + 4q^4 + 5q^6 + 13q^8 + 16q^{10} + 33q^{12} + 41q^{14} + 73q^{16} + 90q^{18} + 145q^{20} + \cdots. \\
& \langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{USp(4)}(q) = \langle W_{(\infty^2)} W_{(\infty^2)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(5)}(q) \\
&= \frac{1 - q^{16}}{(1 - q^2)(1 - q^4)^3(1 - q^6)(1 - q^8)^2}. \\
& \langle \underbrace{W_{(l^2)} \cdots W_{(l^2)}}_k \rangle^{USp(4)}(t; q) = \langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \rangle^{SO(5)}(t; q)
\end{aligned}$$

$$\begin{aligned}
& \langle W_\square W_\square \rangle^{USp(6)}(t; q) \\
&= \frac{1}{48} \frac{(q)_\infty^6}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_\infty^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \frac{(s_i^\pm; q)_\infty (q s_i^{\pm 2}; q)_\infty}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm 2}; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm 2}; q\right)_\infty} \\
&\times \prod_{i < j} \frac{(s_i^\pm s_j^\mp; q)_\infty (s_i^\pm s_j^\pm; q)_\infty (q s_i^\pm s_j^\mp; q)_\infty (q s_i^\pm s_j^\pm; q)_\infty}{\left(q^{\frac{1}{2}} t^2 s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1}{2}} t^2 s_i^\pm s_j^\pm; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1}{2}} t^{-2} s_i^\pm s_j^\pm; q\right)_\infty} \left[\sum_{i=1}^3 (s_i + s_i^{-1}) \right]^2.
\end{aligned}$$

$$\begin{aligned}
& \langle T_{(1,0,0)} T_{(1,0,0)} \rangle^{SO(7)}(t; q) \\
&= \frac{1}{8} \frac{(q)_\infty^6}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_\infty^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \frac{\left(q^{\frac{1}{2}\delta_{i,1}} s_i^\pm; q\right)_\infty \left(q^{1+\frac{1}{2}\delta_{i,1}} s_i^\pm; q\right)_\infty}{\left(q^{\frac{1+\delta_{i,1}}{2}} t^2 s_i^\pm; q\right)_\infty \left(q^{\frac{1+\delta_{i,1}}{2}} t^{-2} s_i^\pm; q\right)_\infty} \\
&\times \prod_{i < j} \frac{\left(q^{\frac{1}{2}\delta_{i+j,1}} s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1}{2}\delta_{i+j,1}} s_i^\pm s_j^\mp; q\right)_\infty \left(q^{1+\frac{1}{2}\delta_{i+j,1}} s_i^\pm s_j^\mp; q\right)_\infty \left(q^{1+\frac{1}{2}\delta_{i+j,1}} s_i^\pm s_j^\mp; q\right)_\infty}{\left(q^{\frac{1+\delta_{i+j,1}}{2}} t^2 s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1+\delta_{i+j,1}}{2}} t^2 s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1+\delta_{i+j,1}}{2}} t^{-2} s_i^\pm s_j^\mp; q\right)_\infty \left(q^{\frac{1+\delta_{i+j,1}}{2}} t^{-2} s_i^\pm s_j^\mp; q\right)_\infty}.
\end{aligned}$$

$$\begin{aligned}
\langle W_\square W_\square \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) &= \langle T_{(1,0,0)} T_{(1,0,0)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(7)}(q) \\
&= \frac{1 + q^2 + q^4 + q^6 + q^8 + q^{10}}{(1 - q^4)(1 - q^8(1 - q^{12}))} \\
&= \frac{1}{(1 - q^2)(1 - q^4)(1 - q^8)}.
\end{aligned}$$

$$\langle \underbrace{W_\square \cdots W_\square}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) = \frac{\sum_{i=0}^{5k} a_{k\square}^{\text{usp}(6)}(i) q^{2i}}{(1 - q^4)(1 - q^8)(1 - q^{12})}$$

$$\begin{aligned}
\langle W_\square W_\square W_\square W_\square \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} (3 + 6q^2 + 9q^4 + 12q^6 + 14q^8 \\
&\quad + 15q^{10} + 12q^{12} + 9q^{14} + 6q^{16} + 3q^{18} + q^{20})
\end{aligned}$$

$$\det \begin{pmatrix} F_0 & F_1 & F_2 \\ F_1 & F_2 & F_3 \\ F_2 & F_3 & F_4 \end{pmatrix} = \sum_{k=0}^{\infty} \frac{a_{k\square}^{\text{usp}(6)}(0)}{(2k)!} x^{2k}$$

$$F_m(x) := \sum_{j=0}^m \binom{m}{j} (I_{2j-m}(2x) - I_{2j-m+2}(2x))$$

$$I_k(2x) := \sum_{n=0}^{\infty} \frac{x^{2n+k}}{n!(n+k)!}$$

$$\left\langle W_\square \right\rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) = \frac{q^4 + q^8}{(1 - q^4)(1 - q^8)(1 - q^{12})},$$

$$\left\langle W_\square W_\square \right\rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) = \frac{1 + q^2 + 2q^4 + 2q^6 + 3q^8 + 2q^{10} + 3q^{12} + q^{14} + q^{16}}{(1 - q^4)(1 - q^8)(1 - q^{12})}.$$

$$\left\langle \begin{matrix} W_\square & W_\square \\ \square & \square \\ \square & \square \end{matrix} \right\rangle_{\frac{1}{2}\text{BPS}}^{USS(6)}(q) = \frac{1 + q^2 + q^4 + 2q^6 + 2q^8 + 2q^{10} + 2q^{12} + q^{14} + q^{16} + q^{18}}{(1 - q^4)(1 - q^8)(1 - q^{12})}$$

$$\langle \underbrace{W_{\square\square} \cdots W_{\square\square}}_k \rangle^{USp(6)}(t; q) = \langle \underbrace{W_\square \cdots W_\square}_k \rangle^{SO(7)}(t; q)$$

$$\begin{aligned}
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) &= \frac{q^2 + q^6 + q^{10}}{(1 - q^4)(1 - q^8)(1 - q^{12})}, \\
\langle W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\quad \times (1 + q^2 + 3q^4 + 2q^6 + 5q^8 + 3q^{10} \\
&\quad + 5q^{12} + 2q^{14} + 3q^{16} + q^{18} + q^{20}) \\
\langle \underbrace{W_{(2l)} \cdots W_{(2l)}}_k \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) &= \frac{\sum_{i=0}^{5lk} a_{k(2l)}^{\text{usp}(6)}(i) q^{2i}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
\langle W_{\square\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(6)}(q) &= \frac{q^4 + q^8 + 2q^{12} + q^{16} + q^{20}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
\langle W_{\square} W_{\square} \rangle^{USp(2N)}(t; q) &= \frac{1}{2^N N!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \frac{(s_i^{\pm 2}; q)_{\infty} (q s_i^{\pm 2}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm 2}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm 2}; q\right)_{\infty}} \\
&\quad \times \prod_{i < j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \left[\sum_{i=1}^N (s_i + s_i^{-1}) \right]^2. \\
\langle T_{(1,0^{N-1})} T_{(1,0^{N-1})} \rangle^{SO(2N+1)}(t; q) &= \frac{1}{2^{N-1} (N-1)!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \frac{\left(q^{\frac{1}{2} \delta_{i,1}} s_i^{\pm}; q\right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i,1}} s_i^{\pm}; q\right)_{\infty}}{\left(\frac{1+\delta_{i,1}}{2} t^2 s_i^{\pm}; q\right)_{\infty} \left(q^{\frac{1+\delta_{i,1}}{2}} t^{-2} s_i^{\pm}; q\right)_{\infty}} \\
&\quad \times \prod_{i < j} \frac{\left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1+\delta_{i+j,1}}{2}} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1+\delta_{i+j,1}}{2}} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1+\delta_{i+j,1}}{2}} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1+\delta_{i+j,1}}{2}} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}. \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \left\langle T_{(1,0^{N-1})} T_{(1,0^{N-1})} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(2N+1)}(q) \\
&= \frac{1}{(1 - q^2) \prod_{n=1}^{N-1} (1 - q^{4n})}. \\
\mathcal{J}_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \prod_{n=1}^N \frac{1}{1 - q^{4n}} \\
\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \frac{1 - q^{4N}}{1 - q^2} \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \frac{\sum_{i=0}^{(2N-1)k} a_{k\square}^{\text{usp}(2N)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
\det(F_{i+j-2}(x)) &= \sum_{k=0}^{\infty} \frac{a_{k\square}^{\text{usp}(2N)}(0)}{(2k)!} x^{2k}
\end{aligned}$$

$$\begin{aligned}
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \frac{\sum_{i=0}^{2(N-1)k} a_{k\square}^{\text{usp}(2N)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
\langle \underbrace{W_{\square\square} \cdots W_{\square\square}}_k \rangle^{USp(2N)}(t; q) &= \langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle^{SO(2N+1)}(t; q) \\
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \frac{q^2 + q^6 + \cdots + q^{4N-2}}{\prod_{n=1}^N (1 - q^{4n})} \\
&= \frac{q^2(1 - q^{4N})}{(1 - q^4) \prod_{n=1}^N (1 - q^{4n})} \\
\langle \underbrace{W_{(2l)} \cdots W_{(2l)}}_k \rangle_{\frac{1}{2}\text{BPS}}^{USp(2N)}(q) &= \frac{\sum_{i=0}^{(2N-1)lk} a_{k(2l)}^{\text{usp}(2N)}(i) q^{2i}}{\prod_{n=1}^N (1 - q^{4n})} \\
(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &\in (\mathbb{Z}_2 \times \mathbb{Z}_2) \times (\mathbb{Z}_2 \times \mathbb{Z}_2) \\
\text{Spin}(2N), \\
SO(2N) &= \text{Spin}(2N)/\mathbb{Z}_2^V, \\
Ss(2N) &= \text{Spin}(2N)/\mathbb{Z}_2^S, \\
Sc(2N) &= \text{Spin}(2N)/\mathbb{Z}_2^C, \\
SO(2N)/\mathbb{Z}_2 &= \text{Spin}(2N)/(\mathbb{Z}_2^S \times \mathbb{Z}_2^C).
\end{aligned}$$

$$\begin{aligned}
Ss(2N)_+ : (z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (0,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (1,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (0,0; 0,1), \\
Ss(2N)_- : (z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (0,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (1,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (0,1; 0,1). \\
Ss(2N)_+ : (z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (0,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (1,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (0,0; 1,0), \\
Ss(2N)_- : (z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (0,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (1,0; 0,0), \\
&(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) = (0,1; 1,0). \\
(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (0,0; 0,0) \\
(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (n_{SS}, n_{SC}; 1,0), \\
(z_{e,S}, z_{e,C}; z_{m,S}, z_{m,C}) &= (n_{CS}, n_{CC}; 0,1), \\
(z_e, z_m) &\in \mathbb{Z}_4 \times \mathbb{Z}_4 \\
\text{Spin}(2N), \\
SO(2N) &= \text{Spin}(2N)/\mathbb{Z}_2, \\
SO(2N)/\mathbb{Z}_2 &= \text{Spin}(2N)/\mathbb{Z}_4. \\
(z_e, z_m) &= (0,0), (z_e, z_m) = (n, 1) \\
\chi_{\text{sp}}^{\text{so}(2N)} &= \frac{1}{2} \left[\prod_{i=1}^N \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}} \right) + \prod_{i=1}^N \left(s_i^{\frac{1}{2}} - s_i^{-\frac{1}{2}} \right) \right]
\end{aligned}$$



$$\begin{aligned}
\chi_{\text{sp}}^{\text{so}(2N)} &= \frac{1}{2} \left[\prod_{i=1}^N \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}} \right) - \prod_{i=1}^N \left(s_i^{\frac{1}{2}} - s_i^{-\frac{1}{2}} \right) \right] \\
\chi_{\square}^{\text{so}(2N)} &= \sum_{i=1}^N (s_i + s_i^{-1}) \\
\chi_{\lambda}^{\text{so}(2N)} &= \frac{\det(s_j^{\lambda_i+N-i} + s_j^{-\lambda_i-N+i}) + \det(s_j^{\lambda_i+N-i} - s_j^{-\lambda_i-N+i})}{\det(s_j^{N-i} + s_j^{-N+i})} \\
\langle W_{\mathcal{R}_1} \cdots W_{\mathcal{R}_k} \rangle^{SO(2N)} &= \int d\mu^{SO(2N)} \exp \left(\sum_{n=1}^{\infty} \frac{1}{n} f_n(q, t) M_n(s) \right) \prod_{i=1}^k \chi_{\mathcal{R}_i}^{\text{so}(2N)}(s), \\
d\mu^{SO(2N)} &= \frac{1}{2^{N-1} N!} \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \prod_{1 \leq i < j \leq N} (1 - s_i s_j) (1 - s_i^{-1} s_j^{-1}) (1 - s_i s_j^{-1}) (1 - s_i^{-1} s_j) \\
M_n(s) &= \frac{P_n(s)^2 + P_{2n}(s)}{2} \\
\mathcal{J}^{SO(4)}(t; q) &= \mathcal{J}^{SU(2)}(t; q) \times \mathcal{J}^{SU(2)}(t; q) \\
&= \frac{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^2}{(q; q)_{\infty}^4} \left(\sum_{\substack{p_1, p_2 \in \mathbb{Z} \\ p_1 < p_2}} \frac{\left(q^{\frac{1}{2}} t^{-2} \right)^{p_1 + p_2 - 2}}{\left(1 - q^{p_1 - \frac{1}{2}} t^2 \right) \left(1 - q^{p_2 - \frac{1}{2}} t^2 \right)} \right)^2. \\
\langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(4)}(t; q) &= \frac{1}{4} \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \\
&\times \frac{(s_1^{\pm} s_2^{\mp}; q)_{\infty} (s_1^{\pm} s_2^{\pm}; q)_{\infty} (q s_1^{\pm} s_2^{\mp}; q)_{\infty} (q s_1^{\pm} s_2^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\pm}; q \right)_{\infty}} \\
&\times \left(s_1^{\frac{1}{2}} s_2^{\frac{1}{2}} + s_1^{-\frac{1}{2}} s_2^{-\frac{1}{2}} \right)^2. \\
\left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{SO(4)/\mathbb{Z}_2}(t; q) &= \frac{1}{2} \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \\
&\times \frac{\left(q^{\frac{1}{2}} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} (s_1^{\pm} s_2^{\pm}; q)_{\infty} \left(q^{\frac{3}{2}} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} (q s_1^{\pm} s_2^{\pm}; q)_{\infty}}{(q t^2 s_1^{\pm} s_2^{\mp}; q)_{\infty} \left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\pm}; q \right)_{\infty} (q t^{-2} s_1^{\pm} s_2^{\mp}; q)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\pm}; q \right)_{\infty}}. \\
\langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(4)}(t; q) &= \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{SO(4)/\mathbb{Z}_2}(t; q) \\
&= \mathcal{J}^{SU(2)}(t; q) \langle W_{\square} W_{\square} \rangle^{SU(2)}(t; q),
\end{aligned}$$

$$\begin{aligned}
\langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) &= \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}\right)} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(4)/\mathbb{Z}_2}(q) \\
&= \frac{1+q^2}{(1-q^4)^2} \\
&= \frac{1}{(1-q^2)(1-q^4)}. \\
\underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{2k}}_{2k}^{\text{Spin}(4)}(t; q) &= \mathcal{J}^{SU(2)}(t; q) \underbrace{\langle W_{\square} \cdots W_{\square} \rangle_{2k}}_{2k}^{SU(2)}(t; q) \\
\underbrace{\langle W_{\text{sp}} \cdots W_{\text{sp}} \rangle_{2k}}_{2k}^{\text{Spin}(4)}(q) &= \mathcal{J}_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) \sum_{i=0}^k a_{k\text{sp}}^{\text{so}(4)}(i) q^{2i} \\
&= \frac{1}{(1-q^4)^2} \sum_{i=0}^k a_{k\text{sp}}^{\text{so}(4)}(i) q^{2i} \\
a_{k\text{sp}}^{\text{so}(4)}(i) &= \frac{(2i+1)(2k)!}{(k-i)!(k+i+1)!} \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) &= \frac{2+3q^2+q^4}{(1-q^4)^2}, \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) &= \frac{5+9q^2+5q^4+q^6}{(1-q^4)^2}, \\
\langle W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) &= \frac{14+28q^2+20q^4+7q^6+q^8}{(1-q^4)^2}. \\
\langle W_{\text{sp}}^{2k} W_{\text{sp}}^{2m} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) &= \left(\mathcal{J}_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) \right)^{-1} \langle W_{\text{sp}}^{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) \langle W_{\text{sp}}^{2m} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) \\
&= (1-q^4)^2 \langle W_{\text{sp}}^{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) \langle W_{\text{sp}}^{2m} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q). \\
\langle W_{\text{sp}}^{2k} W_{\text{sp}}^{2m} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(4)}(q) &= \frac{1}{(1-q^4)^2} \left(\sum_{i=0}^k a_{k\text{sp}}^{\text{so}(4)}(i) q^{2i} \right) \left(\sum_{j=0}^m a_{m\text{sp}}^{\text{so}(4)}(j) q^{2j} \right) \\
\langle W_{\square} W_{\square} \rangle^{SO(4)}(t; q) &= \frac{1}{4} \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \\
&\times \frac{(s_1^{\pm} s_2^{\mp}; q)_{\infty} (s_1^{\pm} s_2^{\pm}; q)_{\infty} (q s_1^{\pm} s_2^{\mp}; q)_{\infty} (q s_1^{\pm} s_2^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\pm}; q \right)_{\infty}} \\
&\times (s_1 + s_2 + s_1^{-1} + s_2^{-1})^2. \\
\langle T_{(1,0)} T_{(1,0)} \rangle^{SO(4)}(t; q) &= \frac{(q)_{\infty}^4}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^2} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \\
&\times \frac{\left(q^{\frac{1}{2}} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} s_1^{\pm} s_2^{\pm}; q \right)_{\infty} \left(q^{\frac{3}{2}} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q^{\frac{3}{2}} s_1^{\pm} s_2^{\pm}; q \right)_{\infty}}{\left(q t^2 s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q t^2 s_1^{\pm} s_2^{\pm}; q \right)_{\infty} \left(q t^{-2} s_1^{\pm} s_2^{\mp}; q \right)_{\infty} \left(q t^{-2} s_1^{\pm} s_2^{\pm}; q \right)_{\infty}}.
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle^{SO(4)}(t; q) &= \langle W_{\square} W_{\square} \rangle^{SU(2)}(t; q)^2 \\
\left\langle W_{\square} W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{1}{(1 - q^2)^2} \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \rangle^{SO(4)}(t; q) &= \langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle^{SU(2)}(t; q)^2 \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{\sum_{i=0}^{2k} a_k^{so(4)}(i) q^{2i}}{(1 - q^4)^2} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{4 + 12q^2 + 13q^4 + 6q^6 + q^8}{(1 - q^4)^2} \\
\left\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{25 + 90q^2 + 131q^4 + 100q^6 + 43q^8 + 10q^{10} + q^{12}}{(1 - q^4)^2} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{1}{(1 - q^4)^2} (196 + 784q^2 + 1344q^4 + 1316q^6 \\
&\quad + 820q^8 + 336q^{10} + 89q^{12} + 14q^{14} + q^{16}). \\
a_{k\square}^{so(4)}(0) &= C_k^2 \\
\langle W_{\square} W_{\square} \rangle^{SO(4)^-}(t; q) &= \frac{1}{2} \frac{(q)_{\infty}^2 (-q; q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty} \left(-q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \oint \frac{ds}{2\pi i s} \\
&\times \frac{(s^{\pm}; q)_{\infty} (-s; q)_{\infty} (qs^{\pm}; q)_{\infty} (-qs^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s^{\pm}; q\right)_{\infty} \left(-q^{\frac{1}{2}} t^2 s^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s^{\pm}; q\right)_{\infty} \left(-q^{\frac{1}{2}} t^{-2} s^{\pm}; q\right)_{\infty}} (s + s^{-1})^2. \\
\langle T_{(1)} T_{(1)} \rangle^{SO(4)^-}(t; q) &= \frac{(q)_{\infty}^2 (-q; q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty} \left(-q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \oint \frac{ds}{2\pi i s} \\
&\times \frac{\left(q^{\frac{1}{2}} s^{\pm}; q\right)_{\infty} \left(-q^{\frac{1}{2}} s; q\right)_{\infty} \left(q^{\frac{3}{2}} s^{\pm}; q\right)_{\infty} \left(-q^{\frac{3}{2}} s^{\pm}; q\right)_{\infty}}{(qt^2 s^{\pm}; q)_{\infty} (-qt^2 s^{\pm}; q)_{\infty} (qt^{-2} s^{\pm}; q)_{\infty} (-qt^{-2} s^{\pm}; q)_{\infty}}. \\
\langle W_{\square} W_{\square} \rangle^{SO(4)^-}(t; q) &= \langle T_{(1)} T_{(1)} \rangle^{SO(4)^-}(t; q) \\
&= \langle W_{\square} W_{\square} \rangle^{SU(2)}(t^2; q^2). \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)^-}(q) &= \langle T_{(1)} T_{(1)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)^-}(q) \\
&= \frac{1}{1 - q^4}. \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)^-}(q) &= \frac{\sum_{i=0}^k a_{k\square}^{so(4)^-}(i) q^{4i}}{1 - q^8} \\
a_{k\square}^{so(4)^-}(i) &= (2i + 1) \frac{(2k)!}{(k - i)! (k + i + 1)!}, \\
\langle W_{\lambda_1} \cdots W_{\lambda_k} \rangle^{O(4)^+}(t; q) &= \frac{1}{2} \left[\langle W_{\lambda_1} \cdots W_{\lambda_k} \rangle^{SO(4)}(t; q) + \langle W_{\lambda_1} \cdots W_{\lambda_k} \rangle^{SO(4)^-}(t; q) \right]. \\
\langle W_{\square} W_{\square} \rangle^{O(4)^+}(t; q) &= \langle T_{(1)} T_{(1)} \rangle^{O(4)^+}(t; q)
\end{aligned}$$

$$\begin{aligned}
\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)^+}(q) &= \frac{\sum_{i=0}^{2k+1} a_{k\square}^{o(4)}(i) q^{2i}}{(1-q^4)(1-q^8)} \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)^+}(q) &= \frac{1+q^2+q^4+q^6}{(1-q^4)(1-q^8)} \\
&= \frac{1}{(1-q^2)(1-q^4)}, \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)^+}(q) &= \frac{3+6q^2+9q^4+9q^6+6q^8+3q^{10}}{(1-q^4)(1-q^8)}, \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)^+}(q) &= \frac{1}{(1-q^4)(1-q^8)} (15+45q^2+80q^4+95q^6+85q^8 \\
&\quad +55q^{10}+20q^{12}+5q^{14}). \\
a_{k\square}^{v(4)}(0) &= \frac{1}{2}(C_k^2 + C_k) \\
\square &= (1,1), \bar{\square} = (1,-1) \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle^{SO(4)}(t; q) &= \langle \underbrace{W_{\bar{\square}} \cdots W_{\bar{\square}}}_k \rangle^{SO(4)}(t; q) \\
&= \mathcal{I}^{SU(2)}(t; q) \langle \underbrace{W_{\square\square} \cdots W_{\square\square}}_k \rangle^{SU(2)}(t; q). \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{\sum_{i=0}^k a_{k\square}^{so(4)}(i) q^{2i}}{(1-q^4)^2} \\
\langle W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{q^2}{(1-q^4)^2}, \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{1+q^2+q^4}{(1-q^4)^2}, \\
\langle W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{1+3q^2+2q^4+q^6}{(1-q^4)^2} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{3+6q^2+6q^4+3q^6+q^8}{(1-q^4)^2} \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{6+15q^2+15q^4+10q^6+4q^8+q^{10}}{(1-q^4)^2}, \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{15+36q^2+40q^4+29q^6+15q^8+5q^{10}+q^{12}}{(1-q^4)^2} \\
\left\langle \left(W_{\square} \right)^k \left(W_{\bar{\square}} \right)^m \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= (1-q^4)^2 \left\langle \left(W_{\square} \right)^k \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) \left\langle \left(W_{\bar{\square}} \right)^m \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q)
\end{aligned}$$

$$\begin{aligned}
\langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \rangle^{SO(4)}(t; q) &= \langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \rangle^{SU(2)}(t; q)^2 \\
\langle \underbrace{W_{\square\square} \cdots W_{\square\square}}_k \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{\sum_{i=0}^{2k} a_{k\square\square}^{so(4)-}(i) q^{2i}}{(1 - q^4)^2} \\
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{q^4}{(1 - q^4)^2} \\
\langle W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{(1 + q^2 + q^4)^2}{(1 - q^4)^2} \\
\langle W_{\square\square} W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{(1 + 3q^2 + 2q^4 + q^6)^2}{(1 - q^4)^2} \\
a_{k\square\square}^{so(4)-}(0) &= R_k^2 \\
a_{k\square\square}^{so(4)-}(1) &= 2R_k R_{k+1} \\
\langle (W_{\square\square})^k \rangle^{SO(4)} &= \left\langle \left(W_{\square} \right)^k \left(W_{\square} \right)^k \right\rangle^{SO(4)} \\
\langle W_{(2k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{q^{4k}}{(1 - q^4)^2} \\
\langle W_{(k)} W_{(k)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{(1 - q^{2k+2})^2}{(1 - q^2)^2 (1 - q^4)^2} \\
\langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \frac{1}{(1 - q^2)^2 (1 - q^4)^2} \\
\langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \rangle^{SO(4)-}(t; q) &= \langle \underbrace{W_{(l)} \cdots W_{(l)}}_k \rangle^{SU(2)}(t^2; q^2) \\
\langle W_{(2l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)-}(q) &= \frac{q^{4l}}{1 - q^8}, \\
\langle W_{(l)} W_{(l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)-}(q) &= \frac{1 - q^{4l+4}}{(1 - q^4)(1 - q^8)}. \\
\langle W_{(2l)} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)+}(q) &= \frac{q^{4l}}{(1 - q^4)(1 - q^8)}, \\
\langle W_{(l)} W_{(l)} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)+}(q) &= \frac{1 - q^2 + q^4 - q^{2l+2} - q^{2l+6} + q^{4l+6}}{(1 - q^2)^2 (1 - q^4)(1 - q^8)} \\
\langle W_{(\infty)} W_{(\infty)} \rangle_{\frac{1}{2}\text{BPS}}^{O(4)+}(q) &= \frac{1 - q^2 + q^4}{(1 - q^2)^2 (1 - q^4)(1 - q^8)} \\
\langle \underbrace{W_{(l,l)} \cdots W_{(l,l)}}_k \rangle^{SO(4)}(t; q) &= \langle \underbrace{W_{(l,-l)} \cdots W_{(l,-l)}}_k \rangle^{SO(4)}(t; q) \\
&= j^{SU(2)}(t; q) \langle \underbrace{W_{(2l)} \cdots W_{(2l)}}_k \rangle^{SU(2)}(t; q). \\
\langle W_{(l,l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \langle W_{(l,-l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) = \frac{q^{2l}}{(1 - q^4)^2}, \\
\langle W_{(l,l)} W_{(l,l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) &= \langle W_{(l,-l)} W_{(l,-l)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(4)}(q) = \frac{1 - q^{4l+2}}{(1 - q^2)(1 - q^4)^2}. \\
\langle W_{(l,l)}^k W_{(l,-l)}^k \rangle^{SO(4)} &= \langle W_{(2l)}^k \rangle^{SO(4)}
\end{aligned}$$

$$\begin{aligned}
j^{SO(6)}(t; q) &= j^{SU(4)}(t; q) \\
&= -\frac{\left(q^{\frac{1}{2}}t^{\pm 2}; q\right)_{\infty}}{(q; q)_{\infty}^2} \sum_{\substack{p_1, p_2, p_3, p_4 \in \mathbb{Z} \\ p_1 < p_2 < p_3 < p_4}} \frac{\left(q^{\frac{1}{2}}t^{-2}\right)^{p_1+p_2+p_3+p_4-8}}{(1-q^{p_1-1}t^4)(1-q^{p_2-1}t^4)(1-q^{p_3-1}t^4)(1-q^{p_4-1}t^4)} \\
\langle W_{\text{sp}} W_{\overline{\text{sp}}} \rangle^{\text{Spin}(6)}(t; q) &= \frac{1}{24} \frac{(q)_{\infty}^6}{\left(q^{\frac{1}{2}}t^{\pm 2}; q\right)_{\infty}^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \\
&\times \prod_{i \neq j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \\
&\times (1 + s_1 s_2 + s_1 s_3 + s_2 s_3)(1 + s_1^{-1} s_2^{-1} + s_1^{-1} s_3^{-1} + s_2^{-1} s_3^{-1}) \\
\left\langle T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} \right\rangle^{SO(6)/\mathbb{Z}_2}(t; q) &= \frac{1}{6} \frac{(q)_{\infty}^6}{\left(q^{\frac{1}{2}}t^{\pm 2}; q\right)_{\infty}^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \\
&\times \prod_{i < j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} \left(q^{\frac{1}{2}}s_i^{\pm} s_j^{\pm}; q\right)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} \left(q^{\frac{3}{2}}s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} (q t^2 s_i^{\pm} s_j^{\pm}; q)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} (q t^{-2} s_i^{\pm} s_j^{\pm}; q)_{\infty}} \\
\langle W_{\text{sp}} W_{\overline{\text{sp}}} \rangle^{\text{Spin}(6)}(t; q) &= \langle T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} \rangle^{SO(6)/\mathbb{Z}_2}(t; q) \\
&= \langle W_{\square} W_{\overline{\square}} \rangle^{SU(4)}(t; q). \\
\langle W_{\text{sp}} W_{\overline{\text{sp}}} \rangle^{\text{Spin}(6)}(t; q) &= \left\langle T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} T_{\left(\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\right)} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(6)/\mathbb{Z}_2}(q) = \frac{1}{(1-q^2)(1-q^4)(1-q^6)} \\
\underbrace{\langle W_{\text{sp}} W_{\overline{\text{sp}}} \cdots W_{\text{sp}} W_{\overline{\text{sp}}} \rangle_{2k}^{\text{Spin}(6)}}(t; q) &= \underbrace{\langle W_{\square} W_{\overline{\square}} \cdots W_{\square} W_{\overline{\square}} \rangle_{2k}^{SU(4)}}(t; q) \\
\underbrace{\langle W_{\text{sp}} W_{\overline{\text{sp}}} \cdots W_{\text{sp}} W_{\overline{\text{sp}}} \rangle_{2k}^{\text{Spin}(6)}}_{\frac{1}{2}\text{BPS}}(q) &= \frac{\sum_{i=0}^{3k} a_{k \text{ sp}}^{\text{so}(6)}(i) q^{2i}}{(1-q^4)(1-q^6)(1-q^8)} \\
\left\langle W_{\text{sp}} W_{\overline{\text{sp}}} W_{\text{sp}} W_{\overline{\text{sp}}} \right\rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(6)}(q) &= \frac{2 + 4q^2 + 6q^4 + 7q^6 + 5q^8 + 3q^{10} + q^{12}}{(1-q^4)(1-q^6)(1-q^8)} \\
\left\langle W_{\text{sp}} W_{\overline{\text{sp}}} W_{\text{sp}} W_{\overline{\text{sp}}} W_{\text{sp}} W_{\overline{\text{sp}}} \right\rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(6)}(q) &= \frac{1}{(1-q^4)(1-q^6)(1-q^8)} (6 + 18q^2 + 35q^4 + 50q^6 \\
&\quad + 53q^8 + 45q^{10} + 29q^{12} + 14q^{14} + 5q^{16} + q^{18}).
\end{aligned}$$

$$\det \begin{pmatrix} I_0(2x) & I_1(2x) & I_2(2x) & I_3(2x) \\ I_1(2x) & I_0(2x) & I_1(2x) & I_2(2x) \\ I_2(2x) & I_1(2x) & I_0(2x) & I_1(2x) \\ I_3(2x) & I_2(2x) & I_1(2x) & I_0(2x) \end{pmatrix} = \sum_{k=0}^{\infty} \frac{a_{k \text{ sp}}^{\text{so}(6)}(0)}{(k!)^2} x^{2k}$$

$$\begin{aligned} \langle W_{\square} W_{\square} \rangle^{SO(6)}(t; q) &= \frac{1}{24} \frac{(q)_{\infty}^6}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \\ &\times \prod_{i \neq j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \\ &\times (s_1 + s_2 + s_3 + s_1^{-1} + s_2^{-1} + s_3^{-1})^2. \end{aligned}$$

$$\begin{aligned} \langle T_{(1,0,0)} T_{(1,0,0)} \rangle^{SO(6)}(t; q) &= \frac{1}{4} \frac{(q)_{\infty}^6}{\left(q^{\frac{1}{2}} t^2; q\right)_{\infty}^3} \oint \prod_{i=1}^3 \frac{ds_i}{2\pi i s_i} \\ &\times \prod_{i < j} \frac{\left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \\ &\times \frac{\left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i+j,1,1})} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i+j,1,1})} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}. \end{aligned}$$

$$\begin{aligned} \langle W_{\square} W_{\square} \rangle^{SO(6)}(t; q) &= \langle T_{(1,0,0)} T_{(1,0,0)} \rangle^{SO(6)}(t; q) \\ &= \left\langle W_{\square} \cdots W_{\square} \right\rangle^{SU(4)}(t; q). \end{aligned}$$

$$\begin{aligned} \langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \langle T_{(1,0,0)} T_{(1,0,0)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) \\ &= \frac{1 + q^2 + 2q^4 + q^6 + q^8}{(1 - q^4)(1 - q^6)(1 - q^8)} \\ &= \frac{1}{(1 - q^2)(1 - q^4)^2}. \end{aligned}$$

$$\underbrace{\langle W_{\square} \cdots W_{\square} \rangle^{SO(6)}(t; q)}_{2k} = \underbrace{\langle W_{\square} \cdots W_{\square} \rangle^{SU(4)}(t; q)}_{2k}$$

$$\underbrace{\langle W_{\square} \cdots W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}}_{2k}(q) = \frac{\sum_{i=0}^{4k} a_{k, \square}^{\text{son}(6)}(i) q^{2i}}{(1 - q^4)(1 - q^6)(1 - q^8)}$$

$$\begin{aligned} \langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^6)(1 - q^8)} (3 + 7q^2 + 15q^4 + 18q^6 + 20q^8 \\ &\quad + 14q^{10} + 9q^{12} + 3q^{14} + q^{16}), \end{aligned}$$

$$\begin{aligned} \langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^6)(1 - q^8)} (16 + 60q^2 + 149q^4 + 249q^6 \\ &\quad + 334q^8 + 347q^{10} + 301q^{12} + 206q^{14} + 119q^{16} \\ &\quad + 53q^{18} + 20q^{20} + 5q^{22} + q^{24}). \end{aligned}$$



$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle^{SO(6)^-}(t; q) &= \frac{1}{8} \frac{(q)_{\infty}^4 (-q; q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^2 \left(-q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \\
&\times \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \frac{(s_i^{\pm}; q)_{\infty} (-s_i^{\pm}; q)_{\infty} (qs_i^{\pm}; q)_{\infty} (-qs_i^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm}; q\right)_{\infty} \left(-q^{\frac{1}{2}} t^2 s_i^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm}; q\right)_{\infty} \left(-q^{\frac{1}{2}} t^{-2} s_i^{\pm}; q\right)_{\infty}} \\
&\times \frac{(s_1^{\pm} s_2^{\mp}; q)_{\infty} (s_1^{\pm} s_2^{\pm}; q)_{\infty} (qs_1^{\pm} s_2^{\mp}; q)_{\infty} (qs_1^{\pm} s_2^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_1^{\pm} s_2^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_1^{\pm} s_2^{\pm}; q\right)_{\infty}} \\
&\times (s_1 + s_2 + s_1^{-1} + s_2^{-1})^2.
\end{aligned}$$

$$\begin{aligned}
\langle T_{(1,0)} T_{(1,0)} \rangle^{SO(6)^-}(t; q) &= \frac{1}{2} \frac{(q)_{\infty}^4 (-q; q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^2 \left(-q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}} \oint \prod_{i=1}^2 \frac{ds_i}{2\pi i s_i} \\
&\times \frac{\left(q^{\frac{1}{2}} s_1^{\pm}; q\right)_{\infty} (s_2^{\pm}; q)_{\infty} \left(-q^{\frac{1}{2}} s_1^{\pm}; q\right)_{\infty} (-s_2^{\pm}; q)_{\infty}}{(qt^2 s_1^{\pm}; q)_{\infty} \left(q^{\frac{1}{2}} t^2 s_2^{\pm}; q\right)_{\infty} (-qt^2 s_1^{\pm}; q)_{\infty} \left(-q^{\frac{1}{2}} t^2 s_2^{\pm}; q\right)_{\infty}} \\
&\times \frac{\left(q^{\frac{3}{2}} s_1^{\pm}; q\right)_{\infty} (qs_2^{\pm}; q)_{\infty} \left(-q^{\frac{3}{2}} s_1^{\pm}; q\right)_{\infty} (-qs_2^{\pm}; q)_{\infty}}{(qt^{-2} s_1^{\pm}; q)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_2^{\pm}; q\right)_{\infty} (-qt^{-2} s_1^{\pm}; q)_{\infty} \left(-q^{\frac{1}{2}} t^{-2} s_2^{\pm}; q\right)_{\infty}} \\
&\times \frac{\left(q^{\frac{1}{2}} s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} s_1^{\pm} s_2^{\pm}; q\right)_{\infty} \left(q^{\frac{3}{2}} s_1^{\pm} s_2^{\mp}; q\right)_{\infty} \left(q^{\frac{3}{2}} s_1^{\pm} s_2^{\pm}; q\right)_{\infty}}{(qt^2 s_1^{\pm} s_2^{\mp}; q)_{\infty} (qt^2 s_1^{\pm} s_2^{\pm}; q)_{\infty} (qt^{-2} s_1^{\pm} s_2^{\mp}; q)_{\infty} (qt^{-2} s_1^{\pm} s_2^{\pm}; q)_{\infty}},
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \langle T_{(1,0)} T_{(1,0)} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) \\
&= \frac{1 + q^2 + q^6 + q^8}{(1 + q^6)(1 - q^4)(1 - q^8)} \\
&= \frac{1}{(1 - q^2)(1 - q^8)}. \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{\sum_{i=0}^{4k} a_k^{SO(6)^-}(i) q^{2i}}{(1 + q^6)(1 - q^4)(1 - q^8)^8}, \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{3 + 5q^2 + 3q^4 + 6q^6 + 8q^8 + 4q^{10} + 3q^{12} + 3q^{14} + q^{16}}{(1 + q^6)(1 - q^4)(1 - q^8)}, \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{1}{(1 + q^6)(1 - q^4)(1 - q^8)} (14 + 30q^2 + 31q^4 + 49q^6 \\
&\quad + 66q^8 + 55q^{10} + 49q^{12} + 46q^{14} + 29q^{16} + 15q^{18} \\
&\quad + 10q^{20} + 5q^{22} + q^{24}) \\
a_k^{SO(6)^-}(0) &= C_k C_{k+2} - C_{k+1}^2 \\
\langle W_{\lambda_1} \cdots W_{\lambda_k} \rangle^{O(6)^+}(t; q) &= \frac{1}{2} \left[\langle W_{\lambda_1} \cdots W_{\lambda_k} \rangle^{SO(6)^+}(t; q) + \langle W_{\lambda_1} \cdots W_{\lambda_k} \rangle^{SO(6)^-}(t; q) \right]
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(6)^+}(q) &= \frac{1 + q^2 + q^4 + q^6 + q^8 + q^{10}}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&= \frac{1}{(1 - q^2)(1 - q^4)(1 - q^8)}, \\
\langle W_{\square} W_{\square} W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(6)^+}(q) &= \frac{1}{(1 - q^2)(1 - q^4)(1 - q^8)} (3 + 6q^2 + 9q^4 + 12q^6 \\
&\quad + 15q^8 + 15q^{10} + 12q^{12} + 9q^{14} + 6q^{16} + 3q^{18}). \\
\langle \underbrace{W_{\square} \cdots W_{\square}}_k \rangle^{SO(6)}(t; q) &= \langle \underbrace{W_{\square} W_{\square} W_{\square} W_{\square}}_k \rangle^{SU(4)}(t; q), \\
\left\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \right\rangle^{SO(6)}(t; q) &= \left\langle \underbrace{W_{\square} \cdots W_{\square}}_{2k} \right\rangle^{SU(4)}(t; q) \\
\left\langle W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \frac{q^2}{(1 - q^2)(1 - q^4)(1 - q^8)}, \\
\left\langle W_{\square} W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\quad \times (1 + 2q^2 + 4q^4 + 6q^6 + 7q^8 + 7q^{10} \\
&\quad + 6q^{12} + 5q^{14} + 3q^{16} + q^{18}), \\
\left\langle \underbrace{W_{\square} W_{\square}}_{\frac{1}{2}\text{BPS}} \right\rangle^{SO(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\quad \times (1 + q^2 + 2q^4 + 3q^6 + 3q^8 + 3q^{10} \\
&\quad + 3q^{12} + 2q^{14} + q^{16} + q^{18}). \\
\left\langle W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{q^2}{(1 + q^2)(1 - q^4)(1 - q^8)}, \\
\left\langle W_{\square} W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{1 + q^2 + 2q^4 + q^8}{(1 + q^2)(1 - q^4)(1 - q^8)}, \\
\left\langle W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{O(6)^+}(q) &= \frac{q^2}{(1 - q^4)^2(1 - q^8)}, \\
\left\langle W_{\square} W_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{O(6)^+}(q) &= \frac{1 + q^2 + 2q^4 + q^6 + 2q^8}{(1 - q^4)^2(1 - q^8)}. \\
\langle \underbrace{W_{(l)} \cdots W_{(l)}}_{2k} \rangle^{SO(6)}(t; q) &= \langle \underbrace{W_{(l^2)} \cdots W_{(\overline{l^2})}}_k \rangle^{SU(4)}(t; q)
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \frac{q^4 + q^8}{(1 + q^4)(1 - q^6)(1 - q^8)}, \\
\langle W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)}(q) &= \frac{1}{(1 - q^4)(1 - q^8)(1 - q^{12})} \\
&\quad \times (1 + q^2 + 3q^4 + 4q^6 + 6q^8 + 6q^{10} + 7q^{12} \\
&\quad + 6q^{14} + 4q^{16} + 4q^{18} + q^{20} + q^{22}), \\
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{q^4 + q^8}{(1 + q^6)(1 - q^4)(1 - q^8)}, \\
\langle W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(6)^-}(q) &= \frac{1 + 2q^8 - 2q^{10} + q^{12} - q^{14} - q^{18}}{(1 + q^6)(1 - q^2)(1 - q^4)(1 - q^8)}. \\
\\
\langle W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(6)^+}(q) &= \frac{q^4 + q^8}{(1 + q^6)(1 - q^4)(1 - q^8)}, \\
\langle W_{\square\square} W_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(6)^+}(q) &= \frac{1 + q^2 + 2q^4 + 2q^6 + 4q^8 + 3q^{10} + 4q^{12} + 2q^{14} + 2q^{16} + q^{18}}{(1 - q^4)(1 - q^8)(1 - q^{12})}. \\
\\
\langle W_{\text{sp}} W_{\text{sp}} \rangle^{\text{Spin}(2N)}(t; q) &= \frac{1}{2^{N-1}N!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}}t^{\pm 2}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \\
&\quad \times \prod_{i \neq j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \\
&\quad \times \frac{1}{4} \left[\prod_{i=1}^N \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}} \right) + \prod_{i=1}^N \left(s_i^{\frac{1}{2}} - s_i^{-\frac{1}{2}} \right) \right]^2. \\
\\
\langle W_{\text{sp}} W_{\overline{\text{sp}}} \rangle^{\text{Spin}(2N)}(t; q) &= \frac{1}{2^{N-1}N!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}}t t^{\pm 2}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \\
&\quad \times \prod_{i \neq j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \\
&\quad \times \frac{1}{4} \left[\prod_{i=1}^N \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}} \right) + \prod_{i=1}^N \left(s_i^{\frac{1}{2}} - s_i^{-\frac{1}{2}} \right) \right] \left[\prod_{i=1}^N \left(s_i^{\frac{1}{2}} + s_i^{-\frac{1}{2}} \right) - \prod_{i=1}^N \left(s_i^{\frac{1}{2}} - s_i^{-\frac{1}{2}} \right) \right] \\
\\
\langle T_{(\frac{1}{2}^N)} T_{(\frac{1}{2}^N)} \rangle^{SO(2N)/\mathbb{Z}_2}(t; q) &= \frac{1}{N!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}}t^{\pm 2}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \\
&\quad \times \prod_{i < j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} \left(q^{\frac{1}{2}}s_i^{\pm} s_j^{\pm}; q\right)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} \left(q^{\frac{3}{2}}s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}}t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} (q t^2 s_i^{\pm} s_j^{\pm}; q)_{\infty} \left(q^{\frac{1}{2}}t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} (q t^{-2} s_i^{\pm} s_j^{\pm}; q)_{\infty}}.
\end{aligned}$$

$$\begin{aligned}
\langle W_{\text{sp}} W_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N)}(q) &= \langle T\left(\frac{1}{2}^N\right) T\left(\frac{1}{2}^N\right) \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N)/\mathbb{Z}_2}(q) \\
&= \prod_{n=1}^N \frac{1}{1 - q^{2n}} \\
\langle W_{\text{sp}} W_{\overline{\text{sp}}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N)}(q) &= \left\langle T\left(\frac{1}{2}^N\right) T\left(\frac{1}{2}^N\right) \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(2N)/\mathbb{Z}_2}(q) \\
&= \prod_{n=1}^N \frac{1}{1 - q^{2n}} \\
j_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N)}(q) &= \frac{1}{1 - q^{2N}} \prod_{n=1}^{N-1} \frac{1}{1 - q^{4n}} \\
\langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N)}(q) &= \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\overline{\text{sp}}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N)}(q) \\
&= \prod_{n=1}^{N-1} (1 + q^{2n}). \\
\langle \underbrace{W_{\text{sp}} \cdots W_{\text{sp}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N=4n)}(q) &= \langle \underbrace{W_{\text{sp}} \cdots W_{\overline{\text{sp}}}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2N=4n+2)}(q) \\
&= \frac{\sum_{i=0}^{\frac{N(N+1)k}{2}} a_{k\text{sp}}^{\text{sp}(2N)}(i) q^{2i}}{1 - q^{2N} \prod_{n=1}^{N-1} (1 - q^{4n})}, \\
\langle W_{\square} W_{\square} \rangle^{SO(2N)}(t; q) &= \frac{1}{2^{N-1} N!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^N} \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \\
&\times \prod_{i \neq j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \left[\sum_{i=1}^N (s_i + s_i^{-1}) \right]^2. \\
\left\langle T_{(1,0^{N-1})} T_{(1,0^{N-1})} \right\rangle^{SO(2N)}(t; q) &= \frac{1}{2^{N-2} (N-1)!} \frac{(q)_{\infty}^{2N}}{\left(q^{\frac{1}{2}} t^{\pm 2}; q\right)_{\infty}^N} \\
&\times \oint \prod_{i=1}^N \frac{ds_i}{2\pi i s_i} \prod_{i < j} \frac{\left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^2 s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^2 s_i^{\pm} s_j^{\pm}; q\right)_{\infty}} \\
&\times \frac{\left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^{-2} s_i^{\pm} s_j^{\mp}; q\right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^{-2} s_i^{\pm} s_j^{\pm}; q\right)_{\infty}}.
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N)}(q) &= \frac{1 + q^2 + \dots + q^{2N-4} + 2q^{2N-2} + q^{2N} + \dots + q^{4N-4}}{(1 - q^{2N}) \prod_{n=1}^{N-1} (1 - q^{4n})} \\
&= \frac{1}{(1 - q^2)(1 - q^{2(N-1)}) \prod_{n=1}^{N-2} (1 - q^{4n})}. \\
\langle \underbrace{W_{\square} \dots W_{\square}}_{2k} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N)}(q) &= \frac{\sum_{i=0}^{(2N-2)k} a_k^{so(2N)}(i) q^{2i}}{(1 - q^{2N}) \prod_{n=1}^{N-1} (1 - q^{4n})}. \\
\langle W_{\square} W_{\square} \rangle^{SO(2N)^-}(t; q) &= \frac{1}{2^{N-1}(N-1)!} \frac{(q)_{\infty}^{2N-2} (-q; q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^{N-1} \left(-q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}} \\
&\times \oint \prod_{i=1}^{N-1} \frac{ds_i}{2\pi i s_i} \frac{(s_i^{\pm}; q)_{\infty} (-s_i^{\pm}; q)_{\infty} (q s_i^{\pm}; q)_{\infty} (-q s_i^{\pm}; q)_{\infty}}{t^2 s_i^{\pm}; q)_{\infty} \left(-q^{\frac{1}{2}} t^2 s_i^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm}; q \right)_{\infty} \left(-q^{\frac{1}{2}} t^{-2} s_i^{\pm}; q \right)_{\infty}} \\
&\times \prod_{i < j} \frac{(s_i^{\pm} s_j^{\mp}; q)_{\infty} (s_i^{\pm} s_j^{\pm}; q)_{\infty} (q s_i^{\pm} s_j^{\mp}; q)_{\infty} (q s_i^{\pm} s_j^{\pm}; q)_{\infty}}{\left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^2 s_i^{\pm} s_j^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}} t^{-2} s_i^{\pm} s_j^{\pm}; q \right)_{\infty}} \left[\sum_{i=1}^{N-1} (s_i + s_i^{-1}) \right]^2. \\
\langle T_{(1,0^{N-2})} T_{(1,0^{N-2})} \rangle^{SO(2N)^-}(t; q) &= \frac{1}{2^{N-2}(N-2)!} \frac{(q)_{\infty}^{2N-2} (-q; q)_{\infty}^2}{\left(q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}^{N-1} \left(-q^{\frac{1}{2}} t^{\pm 2}; q \right)_{\infty}} \oint \prod_{i=1}^{N-1} \frac{ds_i}{2\pi i s_i} \\
&\times \frac{\left(q^{\frac{1}{2} \delta_{i,1}} s_i^{\pm}; q \right)_{\infty} \left(-q^{\frac{1}{2} \delta_{i,1}} s_i^{\pm}; q \right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i,1}} s_i^{\pm}; q \right)_{\infty} \left(-q^{1+\frac{1}{2} \delta_{i,1}} s_i^{\pm}; q \right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i,1})} t^2 s_i^{\pm}; q \right)_{\infty} \left(-q^{\frac{1}{2}(1+\delta_{i,1})} t^2 s_i^{\pm}; q \right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i,1})} t^{-2} s_i^{\pm}; q \right)_{\infty} \left(-q^{\frac{1}{2}(1+\delta_{i,1})} t^{-2} s_i^{\pm}; q \right)_{\infty}} \\
&\times \prod_{i < j} \frac{\left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q \right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^2 s_i^{\pm} s_j^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^2 s_i^{\pm} s_j^{\pm}; q \right)_{\infty}} \\
&\times \frac{\left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\mp}; q \right)_{\infty} \left(q^{1+\frac{1}{2} \delta_{i+j,1}} s_i^{\pm} s_j^{\pm}; q \right)_{\infty}}{\left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^{-2} s_i^{\pm} s_j^{\mp}; q \right)_{\infty} \left(q^{\frac{1}{2}(1+\delta_{i+j,1})} t^{-2} s_i^{\pm} s_j^{\pm}; q \right)_{\infty}}. \\
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2N)^-}(q) &= \left\langle T_{(1,0^{N-2})} T_{(1,0^{N-2})} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(2N)^-}(q) \\
&= \frac{1 + q^2 + \dots + q^{2N-4}}{\prod_{n=1}^{N-1} (1 - q^{4n})} \\
&= \frac{1 - q^{2(N-1)}}{1 - q^2} \prod_{n=1}^{N-1} \frac{1}{1 - q^{4n}}
\end{aligned}$$

$$\begin{aligned}
\langle W_{\square} W_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2N)^+}(q) &= \frac{1 + q^2 + \dots + q^{4N-2}}{\prod_{n=1}^N (1 - q^{4n})} \\
&= \frac{1}{1 - q^2} \prod_{n=1}^{N-1} \frac{1}{1 - q^{4n}}. \\
j_{\frac{1}{2}\text{BPS}}^{O(2N)^+}(q) &= \prod_{n=1}^N \frac{1}{1 - q^{4n}} \\
\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2N)^+}(q) &= \frac{1 - q^{4N}}{1 - q^2} \\
j^{SO(2\infty+1)}(t; q) &= j^{USp(2\infty)}(t; q) = j^{SO(2\infty)}(t; q) = j^{O(2\infty)^+}(t; q) \\
&= \prod_{n,m,l=0}^{\infty} \frac{\left(1 - q^{n+m+l+\frac{3}{2}} t^{-4m+4l\pm 2}\right)^2}{(1 - q^{n+m+l+1} t^{-4m+4l\pm 4})(1 - q^{n+m+1} t^{-4m+4l})(1 - q^{n+m+l+3} t^{-4m+4l})}. \\
i^{AdS_5 \times \mathbb{RP}^5}(t; q) &= \frac{q^{\frac{1}{2}}(t^2 + t^{-2}) - q - q^2}{(1 - qt^4)(1 - qt^{-4})} - \frac{q^{\frac{1}{2}}(t^2 + t^{-2})}{\left(1 + q^{\frac{1}{2}}t^2\right)\left(1 + q^{\frac{1}{2}}t^{-2}\right)(1 - q)} \\
i^X(t; q) &:= \text{Tr}(-1)^F q^{\frac{h+j}{2}} t^{2(q_2 - q_3)} \\
i_{\frac{1}{2}\text{BPS}}^X(q) &:= \text{Tr}(-1)^F q^{2(q_2 - q_3)} \\
\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS},c}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2\infty)}(q) \\
&= \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS},c}^{SO(2\infty)}(q) = \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2\infty)^+}(q) = \frac{1}{1 - q^2}. \\
\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle^{SO(2\infty+1)}(t; q) &= \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle^{USp(2\infty)}(t; q) \\
&= \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle^{SO(2\infty)}(t; q) = \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle^{O(2\infty)^+}(t; q) = \frac{1 - q}{\left(1 - q^{\frac{1}{2}}t^2\right)\left(1 - q^{\frac{1}{2}}t^{-2}\right)}. \\
i^{\text{string}}(t; q) &= -q + q^{\frac{1}{2}}t^2 + q^{\frac{1}{2}}t^{-2} \\
\left\langle \mathcal{W}_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2\infty)}(q) = \langle \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2\infty)^+}(q) \\
&= \frac{q^2}{(1 - q^4)} \\
&= q^2 + q^6 + q^{10} + q^{14} + q^{18} + \dots. \\
\left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2\infty)}(q) = \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle_{\frac{1}{2}\text{BPS}}^{O(2\infty)^+}(q) \\
&= \frac{1 + q^2 + q^4}{(1 - q^4)^2} \\
&= 1 + q^2 + 3q^4 + 2q^6 + 5q^8 + 3q^{10} + 7q^{12} + 4q^{14} + 9q^{16} + \dots.
\end{aligned}$$

$$\begin{aligned} \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle_{\frac{1}{2}\text{BPS},c}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS},c}^{USp(2\infty)}(q) \langle \mathcal{W}_{\square} \mathcal{W}_{\square} \rangle_{\frac{1}{2}\text{BPS},c}^{O(2\infty)^+} \\ &= \frac{1}{(1-q^2)(1-q^4)} = 1 + q^2 2q^4 + 2q^6 + 3q^8 + 3q^{10} + 4q^{12} + 4q^{14} + 5q^{16} + 5q^{18} + \dots \end{aligned}$$

$$\begin{aligned} \langle \mathcal{W}_{\lambda} \mathcal{W}_{\lambda} \rangle_{\frac{1}{2}\text{BPS},c}^G(q) &:= \langle \mathcal{W}_{\lambda} \mathcal{W}_{\lambda} \rangle_{\frac{1}{2}\text{BPS}}^G(q) - \langle \mathcal{W}_{\lambda} \rangle_{\frac{1}{2}\text{BPS}}^G(q)^2. \\ \left\langle \mathcal{W}_{\square} \right\rangle^{SO(2\infty+1)}(t;q) &= \langle \mathcal{W}_{\square\square} \rangle^{USp(2\infty)}(t;q) = \left\langle \mathcal{W}_{\square} \right\rangle^{O(2\infty)^+}(t;q) \\ &= \frac{q^{\frac{1}{2}}(t^2 + t^{-2}) - q - q^2}{(1-qt^4)(1-qt^{-4})} \end{aligned}$$

$$\begin{aligned} \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle^{SO(2\infty+1)}(t;q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle^{USp(2\infty)}(t;q) = \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle^{O(2\infty)^+}(t;q) \\ &= \frac{1}{(1-qt^4)(1-qt^{-4})} \left(1 + (t^2 + t^{-2})q^{\frac{1}{2}} + (3 + t^4 + t^{-4})q - 3(t^2 + t^{-2})q^{\frac{3}{2}} \right. \\ &\quad \left. - (t^2 + t^{-2})q^2 - 3(t^2 + t^{-2})q^{\frac{5}{2}} + (3 + t^4 + t^{-4})q^3 + (t^2 + t^{-2})q^{\frac{7}{2}} + q^4 \right). \end{aligned}$$

$$\begin{aligned} \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle_c^{SO(2\infty+1)}(t;q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_c^{USp(2\infty)}(t;q) = \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle_c^{O(2\infty)^+}(t;q) \\ &= \frac{(1-q) \left(1 + q - q^{\frac{3}{2}}(t^2 + t^{-2}) \right)}{\left(1 - q^{\frac{1}{2}}t^2 \right) \left(1 - q^{\frac{1}{2}}t^{-2} \right) (1-qt^4)(1-qt^{-4})}. \end{aligned}$$

$$\begin{aligned} \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(2\infty+1)}(q) &= \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Sin}(4\infty)}(q) = \left\langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \right\rangle_{\frac{1}{2}\text{BPS}}^{S^{\text{Sin}(4\infty+2)}}(q) \\ &= \prod_{n=1}^{\infty} \frac{1}{1 - q^{4n-2}}. \end{aligned}$$

$$\langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle_{\frac{1}{2}\text{BPS}}^{\text{Spin}(\infty)}(q) = \sum_{n \geq 0} d_{\{\text{sp}, \text{sp}\}}^{(H)}(n) q^{2n}$$

$$i_{\frac{1}{2}\text{BPS}}^{\text{fat string}}(q) = \frac{q^2}{1-q^4} = q^2 + q^6 + q^{10} + \dots$$

$$S_{\text{D5}} = T_5 \int d^6\sigma \sqrt{\det(g + 2\pi\alpha' F)} - iT_5 \int 2\pi\alpha' F \wedge C_{(4)}$$

$$S_{\text{D5}AdS_2 \times \mathbb{RP}^4} = T_5 \int d^6\sigma \sqrt{\det g} = T_5 \text{vol}(AdS_2) \text{vol}(\mathbb{RP}^4)$$

$$ds_{AdS_2}^2 = \frac{1}{r^2}(-dt^2 + dr^2), ds_4^2 = g_{ij}d\sigma^i d\sigma^j, i, j = 1, 2, 3, 4$$

$$S = T_5 \int d^6\sigma \sqrt{g^{(4)}} \frac{1}{2} \frac{1}{r^2} [r^2 (\partial_t \phi)^2 - r^2 (\partial_r \phi)^2 + (\nabla_i \phi \nabla^i \phi - 4\phi^2)].$$

$$\phi(t,r,\Theta) = \sum_w \phi_w(t,r) Y^w(\Theta).$$



$$\begin{aligned}
S &= T_5 \sum_w \frac{2}{3} \pi^2 \int d^2 \sigma \frac{1}{r^2} (r^2 (\partial_t \phi_w)^2 - r^2 (\partial_r \phi_w)^2 - w(w+1) \phi_w^2) \\
h &= \frac{1}{2} + \sqrt{\frac{1}{4} + m^2} \\
\langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle^{\text{Spin}(2\infty+1)}(t; q) &= \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle^{\text{Spin}(4\infty)}(t; q) = \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle^{\text{Spin}(4\infty+2)}(t; q) \\
&= \prod_{n=0}^{\infty} \prod_{m=0}^{\infty} \frac{(1 - q^{1+n+m} t^{4n-4m})(1 - q^{2+n+m} t^{4n-4m})}{\left(1 - q^{\frac{1}{2}+n+m} t^{2+4n-4m}\right) \left(1 - q^{\frac{1}{2}+n+m} t^{-2+4n+4m}\right)} \\
\langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle^{\text{Spin}(2\infty+1)}(q) &= \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle^{\text{Spin}(4\infty)}(q) = \langle \mathcal{W}_{\text{sp}} \mathcal{W}_{\text{sp}} \rangle^{\text{Spin}(4\infty+2)}(q) \\
&= \prod_{n=1}^{\infty} \frac{(1 - q^n)^{2n-1}}{\left(1 - q^{n-\frac{1}{2}}\right)^{2n}} \\
&= 1 + 2q^{1/2} + 2q^2 + 6q^{3/2} + 7q^2 + 10q^{5/2} + 21q^3 + 22q^{7/2} + \dots \\
i^{\text{fat string}}(t; q) &= \frac{q^{\frac{1}{2}}(t^2 + t^{-2}) - q - q^2}{(1 - qt^4)(1 - qt^{-4})} \\
\langle \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2\infty)}(q) = \langle \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2\infty)^+}(q) \\
&= \frac{q^4}{(1 - q^4)} \\
\langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2\infty)}(q) = \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2\infty)^+}(q) \\
&= \frac{1 + q^2 + q^8}{(1 - q^4)^2} \\
\langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS},c}^{SO(2\infty+1)}(q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{USp(2\infty)}(q) = \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_{\frac{1}{2}\text{BPS}}^{O(2\infty)^+}(q) \\
&= \frac{1}{(1 - q^2)(1 - q^4)} \\
\langle \mathcal{W}_{\square\square} \rangle^{SO(2\infty+1)}(t; q) &= \langle \mathcal{W}_{\square\square} \rangle^{USp(2\infty)}(t; q) = \langle \mathcal{W}_{\square\square} \rangle^{O(2\infty)^+}(t; q) \\
&= \frac{q(1 + t^4 + t^{-4}) - q^{\frac{3}{2}}(t^2 + t^{-2}) - q^2}{(1 - qt^4)(1 - qt^{-4})} \\
\langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle^{SO(2\infty+1)}(t; q) &= \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle^{USp(2\infty)}(t; q) = \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle^{O(2\infty)^+}(t; q) \\
&= \frac{1}{(1 - qt^4)(1 - qt^{-4})} \left(1 + (t^2 + t^{-2})q^{\frac{1}{2}} + q - (t^2 + t^{-2})q^{\frac{3}{2}} + (t^8 + t^4 + t^{-4} + t^{-8})q^2 \right. \\
&\quad \left. - (2t^6 + 5t^2 + 5t^{-2} + 2t^{-6})q^{\frac{5}{2}} + q^3 + 3(t^2 + t^{-2})q^{\frac{7}{2}} + q^4 \right).
\end{aligned}$$

$$\langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_c^{SO(2\infty+1)}(t; q) = \left\langle \mathcal{W}_{\square} \mathcal{W}_{\square} \right\rangle_c^{USp(2\infty)}(t; q) = \langle \mathcal{W}_{\square\square} \mathcal{W}_{\square\square} \rangle_c^{O(2\infty)^+}(t; q)$$

$$= \frac{(1-q) \left(1+q - q^{\frac{3}{2}}(t^2+t^{-2}) \right)}{\left(1 - q^{\frac{1}{2}}t^2 \right) \left(1 - q^{\frac{1}{2}}t^{-2} \right) (1-qt^4)(1-qt^{-4})}.$$

$$\chi_{\lambda}^{\text{usp}(2N)} = \det \left(E_{\lambda'_i-i+j} - E_{\lambda'_i-i-j} \right)_{1 \leq i, j \leq l(\lambda')}$$

$$\chi_{\lambda}^{\text{so}(2N+1)} = \det(\bar{H}_{\lambda_i-i+j} - \bar{H}_{\lambda_i-i-j})_{1 \leq i, j \leq l(\lambda)}$$

$$E_k = e_k(s_1, \ldots, s_N, s_1^{-1}, \ldots, s_N^{-1})$$

$$\bar{H}_k = h_k(s_1, \ldots, s_N, s_1^{-1}, \ldots, s_N^{-1}, 1)$$

$$P_k = p_k(s_1, \ldots, s_N, s_1^{-1}, \ldots, s_N^{-1})$$

$$\bar{P}_k = p_k(s_1, \ldots, s_N, s_1^{-1}, \ldots, s_N^{-1}, 1)$$

$$p_k(x_1, \ldots, x_n) = \sum_{i=1}^n x_i^k$$

$$\chi_{\square}^{\text{usp}(2N)} = P_1,$$

$$\chi_{\square}^{\text{so}(2N+1)} = \bar{P}_1,$$

$$\chi_{\square\square}^{\text{usp}(2N)} = \frac{P_2}{2} + \frac{P_1^2}{2},$$

$$\chi_{\square\square}^{\text{so}(2N+1)} = \frac{\bar{P}_2}{2} + \frac{\bar{P}_1^2}{2} - 1,$$

$$\chi_{\square}^{\text{usp}(2N)} = -\frac{P_2}{2} + \frac{P_1^2}{2} - 1,$$

$$\chi_{\square}^{\text{so}(2N+1)} = -\frac{\bar{P}_2}{2} + \frac{\bar{P}_1^2}{2},$$

$$\chi_{\square\square\square}^{\text{usp}(2N)} = \frac{P_3}{3} + \frac{P_2P_1}{2} + \frac{P_1^3}{6},$$

$$\chi_{\square\square\square}^{\text{so}(2N+1)} = \frac{\bar{P}_3}{3} + \frac{\bar{P}_2\bar{P}_1}{2} + \frac{\bar{P}_1^3}{6} - \bar{P}_1,$$

$$\chi_{\square\square}^{\text{usp}(2N)} = -\frac{P_3}{3} + \frac{P_1^3}{3} - P_1,$$

$$\chi_{\square\square}^{\text{so}(2N+1)} = -\frac{\bar{P}_3}{3} + \frac{\bar{P}_1^3}{3} - \bar{P}_1,$$

$$\chi_{\square\square\square}^{\text{usp}(2N)} = \frac{P_3}{3} - \frac{P_2P_1}{2} + \frac{P_1^3}{6} - P_1, \quad \chi_{\square\square\square}^{\text{so}(2N+1)} = \frac{\bar{P}_3}{3} - \frac{\bar{P}_2\bar{P}_1}{2} + \frac{\bar{P}_1^3}{6}.$$

$$T_p\mathcal{M}^{\mathbb{C}} = T_p\mathcal{M}^{+} \oplus T_p\mathcal{M}^{-},$$

$$T_p\mathcal{M}^{\pm} = \{Z \in T_p\mathcal{M}^{\mathbb{C}}/\mathcal{I}_pZ = \pm iZ\}$$

$$\mathcal{I}_p(\mathcal{P}^{\pm}Z) = \mathcal{I}_pZ^{\pm} = \pm i(\mathcal{P}^{\pm}Z) = \pm iZ^{\pm}$$

$$\mathcal{N}(u,v)\equiv [\mathcal{J}u,\mathcal{J}v]-\mathcal{J}[u,\mathcal{J}v]-\mathcal{J}[\mathcal{J}u,v]-[u,v]$$

$$\phi(p)\equiv z^\mu\equiv x^\mu+iy^\mu, \psi(p)\equiv w^\mu\equiv u^\mu+iv^\mu, 1\leq \mu, v\leq n$$

$$\frac{\partial u^\nu}{\partial x^\mu} = \frac{\partial v^\nu}{\partial y^\mu}, \frac{\partial u^\nu}{\partial y^\mu} = -\frac{\partial v^\nu}{\partial x^\mu}.$$

$$\mathcal{J}_p\frac{\partial}{\partial x^\mu}=\frac{\partial}{\partial y^\mu}, \mathcal{J}_p\frac{\partial}{\partial y^\mu}=-\frac{\partial}{\partial x^\mu}, \text{ entonces}$$

$$\mathcal{J}_p\frac{\partial}{\partial u^\mu}=\frac{\partial}{\partial v^\mu}, \mathcal{J}_p\frac{\partial}{\partial v^\mu}=-\frac{\partial}{\partial u^\mu}$$

$$(\mathcal{J}_p)=\begin{pmatrix} 0 & \mathbb{I} \\ -\mathbb{I} & 0 \end{pmatrix}, \forall p\in \mathcal{M}$$

$$(\mathcal{J}_p)=\begin{pmatrix} i\mathbb{I} & 0 \\ 0 & -i\mathbb{I} \end{pmatrix}, \forall p\in \mathcal{M}$$

$$\alpha=\sum_{p+q=k}\alpha^{(p,q)}, \alpha^{(p,q)}=\frac{1}{p!q!}\alpha_{\mu_1\ldots\mu_p\bar{\nu}_1\ldots\bar{\nu}_q}dz^{\mu_1}\wedge\ldots\wedge dz^{\mu_p}\wedge d\bar{z}^{\bar{\nu}_1}\ldots\wedge d\bar{z}^{\bar{\nu}_q}$$

$$\begin{aligned}
d &= \partial + \bar{\partial}, \\
\partial \alpha^{(p,q)} &= \alpha^{(p+1,q)}, \bar{\partial} \alpha^{(p,q)} = \alpha^{(p,q+1)}, \\
\alpha^{(p+1,q)} &= \frac{1}{p! q!} \partial_\mu \alpha_{\mu_1 \dots \mu_p \bar{\nu}_1 \dots \bar{\nu}_q}^{(p,q)} dz^\mu \wedge dz^{\mu_1} \wedge \dots \wedge dz^{\mu_p} \wedge d\bar{z}^{\bar{\nu}_1} \wedge \dots \wedge d\bar{z}^{\bar{\nu}_q}, \\
\alpha^{(p,q+1)} &= \frac{1}{p! q!} \partial_{\bar{\nu}} \alpha_{\mu_1 \dots \mu_p \bar{\nu}_1 \dots \bar{\nu}_q}^{(p,qz^{\mu_1})} \wedge \dots \wedge dz^{\mu_p} \wedge d\bar{z}^{\bar{\nu}} \wedge d\bar{z}^{\bar{\nu}_1} \wedge \dots \wedge d\bar{z}^{\bar{\nu}_q}, \\
\partial^2 &= \partial \bar{\partial} + \bar{\partial} \partial = \bar{\partial}^2 \\
\mathcal{G}_p(\mathcal{J}_p u, \mathcal{J}_p v) &= \mathcal{G}_p(u, v) \quad \forall p \in \mathcal{M}, u, v \in T_p \mathcal{M} \\
\mathcal{G}_p(u, v) &\equiv g_p(u, v) + g_p(\mathcal{J}_p u, \mathcal{J}_p v) \quad \forall p \in \mathcal{M}, u, v \in T_p \mathcal{M} \\
\mathcal{G}_p(\mathcal{J}_p u, u) &= \mathcal{G}_p(\mathcal{J}_p^2 u, \mathcal{J}_p u) = -\mathcal{G}_p(u, \mathcal{J}_p u) = -\mathcal{G}_p(\mathcal{J}_p u, u) \\
\mathcal{G}_{\mu\nu} &\equiv \mathcal{G}\left(\frac{\partial}{\partial z^\mu}, \frac{\partial}{\partial z^\nu}\right) = \mathcal{G}\left(\mathcal{J}_p \frac{\partial}{\partial z^\mu}, \mathcal{J}_p \frac{\partial}{\partial z^\nu}\right) = -\mathcal{G}\left(\frac{\partial}{\partial z^\mu}, \frac{\partial}{\partial z^\nu}\right) \\
\mathcal{G}_{\bar{\mu}\bar{\nu}} &\equiv \mathcal{G}\left(\frac{\partial}{\partial \bar{z}^\mu}, \frac{\partial}{\partial \bar{z}^\nu}\right) = \mathcal{G}\left(\mathcal{J}_p \frac{\partial}{\partial \bar{z}^\mu}, \mathcal{J}_p \frac{\partial}{\partial \bar{z}^\nu}\right) = -\mathcal{G}\left(\frac{\partial}{\partial \bar{z}^\mu}, \frac{\partial}{\partial \bar{z}^\nu}\right) \\
\mathcal{G}_{\mu\bar{\nu}} &\equiv \mathcal{G}\left(\frac{\partial}{\partial z^\mu}, \frac{\partial}{\partial \bar{z}^\nu}\right) = \mathcal{G}\left(\mathcal{J}_p \frac{\partial}{\partial z^\mu}, \mathcal{J}_p \frac{\partial}{\partial \bar{z}^\nu}\right) = \mathcal{G}\left(\frac{\partial}{\partial \bar{z}^\nu}, \frac{\partial}{\partial z^\mu}\right) = \mathcal{G}_{\bar{\nu}\mu} \\
\mathcal{G} &= \mathcal{G}_{\mu\bar{\nu}} dz^\mu \otimes d\bar{z}^{\bar{\nu}} + \mathcal{G}_{\bar{\mu}\nu} d\bar{z}^{\bar{\mu}} \otimes dz^\nu \\
\omega_p(u, v) &\equiv \mathcal{G}_p(\mathcal{J}_p u, v) \quad \forall u, v \in T_p \mathcal{M}. \\
\omega_p(u, v) &= \mathcal{G}_p(\mathcal{J}_p^2 u, \mathcal{J}_p v) = \mathcal{G}_p(-u, \mathcal{J}_p v) = -\mathcal{G}_p(\mathcal{J}_p v, u) = -\omega(u, v). \\
\omega_p(\mathcal{J}_p u, \mathcal{J}_p v) &= \mathcal{G}_p(\mathcal{J}_p^2 u, \mathcal{J}_p v) = \mathcal{G}_p(-u, \mathcal{J}_p v) = \omega(u, v) \\
\omega &= i\mathcal{G}_{\mu\bar{\nu}} dz^\mu \wedge d\bar{z}^{\bar{\nu}} \\
[\omega] &\in H_{\bar{\partial}}^{(1,1)}(\mathcal{M}; \mathbb{C}) \\
\omega &= i\partial_\mu \partial_{\bar{\nu}} \mathcal{K}(z, \bar{z}) dz^\mu \wedge d\bar{z}^{\bar{\nu}} \\
\mathcal{K}' &= \mathcal{K} + f + f', \\
\omega(\mathcal{K}') &= i\partial \bar{\partial}(\mathcal{K} + f + f') = i\partial \bar{\partial} \mathcal{K} = \omega(\mathcal{K}). \\
\Gamma_{\mu\nu}^\rho &= \mathcal{G}^{\rho\bar{\rho}} \partial_\mu \mathcal{G}_{\bar{\rho}\nu}, \Gamma_{\bar{\mu}\bar{\nu}}^{\bar{\rho}} = \mathcal{G}^{\bar{\rho}\rho} \partial_{\bar{\mu}} \mathcal{G}_{\rho\bar{\nu}} \\
R_{\mu\bar{\nu}} &= \frac{1}{2} \partial_\mu \partial_{\bar{\nu}} (\log \det \mathcal{G}). \\
\mathfrak{R} &\equiv iR_{\mu\bar{\nu}} dz^\mu \wedge d\bar{z}^{\bar{\nu}} \\
\Psi_{(i)} &= e^{-(qf_{(i,j)} + \bar{q}\bar{f}_{(i,j)})} \Psi_{(j)}, \\
\mathcal{K}_{(i)} &= \mathcal{K}_{(j)} + f_{(i,j)} + \bar{f}_{(i,j)}. \\
\mathcal{Q} &\equiv (2i)^{-1} (dz^\mu \partial_\mu \mathcal{K} - d\bar{z}^{\bar{\nu}} \partial_{\bar{\nu}} \mathcal{K}) \\
\mathcal{Q}_{(i)} &= \mathcal{Q}_{(j)} - \frac{i}{2} \partial f_{(i,j)}, \\
\mathfrak{D}_\mu &\equiv \nabla_\mu + iq\mathcal{Q}_\mu, \mathfrak{D}_{\bar{\nu}} \equiv \nabla_{\bar{\nu}} - i\bar{q}\mathcal{Q}_{\bar{\nu}} \\
D_\mu &\equiv \partial_\mu + iq\mathcal{Q}_\mu, D_{\bar{\nu}} \equiv \partial_{\bar{\nu}} - i\bar{q}\mathcal{Q}_{\bar{\nu}} \\
\Omega &\equiv \begin{pmatrix} \mathcal{X}^\Lambda \\ \mathcal{F}_\Sigma \end{pmatrix} \rightarrow \begin{cases} \langle \Omega | \bar{\Omega} \rangle & \equiv \bar{\mathcal{X}}^\Lambda \mathcal{F}_\Lambda - \mathcal{X}^\Lambda \bar{\mathcal{F}}_\Lambda = -ie^{-\mathcal{K}} \\ \partial_{\bar{\nu}} \Omega & = 0 \\ \langle \partial_\mu \Omega | \Omega \rangle & = 0 \end{cases} \\
\mathcal{V} &\equiv \begin{pmatrix} \mathcal{L}^\Lambda \\ \mathcal{M}_\Sigma \end{pmatrix} \rightarrow \begin{cases} \langle \mathcal{V} | \bar{\mathcal{V}} \rangle & \equiv \bar{\mathcal{L}}^\Lambda \mathcal{M}_\Lambda - \mathcal{L}^\Lambda \bar{\mathcal{M}}_\Lambda = -i \\ \mathfrak{D}_{\bar{\nu}} \mathcal{V} & = \left(\partial_{\bar{\nu}} + \frac{1}{2} \partial_{\bar{\mathcal{L}}} \mathcal{K} \right) \mathcal{V} = 0, \\ \langle \mathfrak{D}_\mu \mathcal{V} | \mathcal{V} \rangle & = 0 \end{cases}
\end{aligned}$$

$$\begin{aligned}
u_\mu &\equiv \mathfrak{D}_\mu \mathcal{V} = \begin{pmatrix} f^\Lambda{}_\mu \\ h_{\Sigma\mu} \end{pmatrix}, \bar{u}_{\bar{\nu}} = \overline{u_\nu}, \\
\mathfrak{D}_{\bar{\nu}} u_\mu &= \mathcal{G}_{\mu\bar{\nu}} \mathcal{V} \langle u_\mu | \bar{u}_{\bar{\nu}} \rangle = i \mathcal{G}_{\mu\bar{\nu}} \\
\langle u_\mu | \bar{\mathcal{V}} \rangle &= 0, \langle u_\mu | \mathcal{V} \rangle = 0 \\
\langle \mathfrak{D}_\mu u_\nu | \mathcal{V} \rangle &= \langle u_\nu | u_\mu \rangle = 0. \\
\mathcal{A} &= i \langle \mathcal{A} | \bar{\mathcal{V}} \rangle \mathcal{V} - i \langle \mathcal{A} | \mathcal{V} \rangle \bar{\mathcal{V}} + i \langle \mathcal{A} | u_\mu \rangle \mathcal{G}^{\mu\bar{\nu}} \bar{u}_{\bar{\nu}} - i \langle \mathcal{A} | \bar{u}_{\bar{\nu}} \rangle \mathcal{G}^{\mu\bar{\nu}} u_\mu \\
\mathcal{C}_{\mu\nu\rho} &\equiv \langle \mathfrak{D}_\mu u_\nu | u_\rho \rangle \rightarrow \mathfrak{D}_\mu u_\nu = i \mathcal{C}_{\mu\nu\rho} \mathcal{G}^{\rho\bar{\epsilon}} \bar{u}_{\bar{\epsilon}} \\
\mathfrak{D}_{\bar{\mu}} \mathcal{C}_{\nu\rho\epsilon} &= 0, \mathfrak{D}_{[\mu} \mathcal{C}_{\nu]\rho\epsilon} = 0 \\
\mathcal{M}_\Lambda &= \mathcal{N}_{\Lambda\Sigma} \mathcal{L}^\Sigma, h_{\Lambda\mu} = \bar{\mathcal{N}}_{\Lambda\Sigma} f^\Sigma{}_\mu. \\
\mathcal{L}^\Lambda \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} \bar{\mathcal{L}}^\Sigma &= -\frac{1}{2}, \\
\mathcal{L}^\Lambda \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} f^\Sigma{}_\mu &= \mathcal{L}^\Lambda \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} \bar{f}^\Sigma{}_{\bar{\nu}} = 0 \\
f^\Lambda{}_\mu \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} \bar{f}^\Sigma{}_{\bar{\nu}} &= -\frac{1}{2} \mathcal{G}_{\mu\bar{\nu}}. \\
(\partial_\mu \mathcal{N}_{\Lambda\Sigma}) \mathcal{L}^\Sigma &= -2i \mathfrak{Im}(\mathcal{N})_{\Lambda\Sigma} f^\Sigma{}_\mu \\
\partial_\mu \bar{\mathcal{N}}_{\Lambda\Sigma} f^\Sigma{}_\nu &= -2 \mathcal{C}_{\mu\nu\rho} \mathcal{G}^{\rho\bar{\rho}} \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} \bar{f}^\Sigma{}_{\bar{\rho}} \\
\mathcal{C}_{\mu\nu\rho} &= f^\Lambda{}_\mu f^\Sigma{}_\nu \partial_\rho \bar{\mathcal{N}}_{\Lambda\Sigma} \\
\mathcal{L}^\Sigma \partial_{\bar{\nu}} \mathcal{N}_{\Lambda\Sigma} &= 0, \\
\partial_{\bar{\nu}} \bar{\mathcal{N}}_{\Lambda\Sigma} f^\Sigma{}_\mu &= 2i \mathcal{G}_{\mu\bar{\nu}} \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} \mathcal{L}^\Sigma. \\
U^{\Lambda\Sigma} &\equiv f^\Lambda{}_\mu \mathcal{G}^{\mu\bar{\nu}} \bar{f}^\Sigma{}_{\bar{\nu}} = -\frac{1}{2} \mathfrak{Im}(\mathcal{N})^{-1|\Lambda\Sigma} - \bar{\mathcal{L}}^\Lambda \mathcal{L}^\Sigma \\
\mathcal{T}_\Lambda &\equiv 2i \mathcal{L}_\Lambda = 2i \mathcal{L}^\Sigma \mathfrak{Im} \mathcal{N}_{\Sigma\Lambda}, \\
\mathcal{T}^\mu{}_\Lambda &\equiv -\bar{f}_\Lambda{}^\mu = -\mathcal{G}^{\mu\bar{\nu}} \bar{f}^\Sigma{}_{\bar{\nu}} \mathfrak{Im} \mathcal{N}_{\Sigma\Lambda}. \\
\partial_\mu \mathcal{N}_{\Lambda\Sigma} &= 4 \mathcal{T}_{\mu(\Lambda} \mathcal{T}_{\Sigma)}, \\
\partial_{\bar{\nu}} \mathcal{N}_{\Lambda\Sigma} &= 4 \bar{\mathcal{C}}_{\bar{\nu}\bar{\rho}\bar{\epsilon}} \mathcal{T}^{\bar{\nu}}{}_{(\Lambda} \mathcal{T}^{\bar{\rho}}{}_{\Sigma)}. \\
e^{-\mathcal{K}} &= -2 \mathfrak{Im} \mathcal{N}_{\Lambda\Sigma} \mathcal{X}^\Lambda \bar{\mathcal{X}}^\Sigma \\
\partial_\mu \mathcal{X}^\Lambda [2 \mathcal{F}_\Lambda - \partial_\Lambda (\mathcal{X}^\Sigma \mathcal{F}_\Sigma)] &= 0. \\
\mathcal{F}_\Lambda &= \partial_\Lambda \mathcal{F}(\mathcal{X}) \\
\mathcal{N}_{\Lambda\Sigma} &= \bar{\mathcal{F}}_{\Lambda\Sigma} + 2i \frac{\mathfrak{Im} \mathcal{F}_{\Lambda\Lambda'} \mathcal{X}^{\Lambda'} \mathfrak{Im} \mathcal{F}_{\Sigma\Sigma'} \mathcal{X}^{\Sigma'}}{\mathcal{X}^\Omega \mathfrak{Im} \mathcal{F}_{\Omega\Omega'} \mathcal{X}^{\Omega'}}. \\
\mathcal{C}_{\mu\nu\rho} &= e^{\mathcal{K}} \partial_\mu \mathcal{X}^\Lambda \partial_\nu \mathcal{X}^\Sigma \partial_\rho \mathcal{X}^\Omega \mathcal{F}_{\Lambda\Sigma\Omega}, \\
[P_\mu, P_\nu] &= 0 \\
[P_\mu, J_{\nu\rho}] &= (\eta_{\mu\nu} P_\rho - \eta_{\mu\rho} P_\nu), \\
[J_{\mu\nu}, J_{\rho\gamma}] &= -(\eta_{\mu\rho} J_{\nu\gamma} + \eta_{\nu\gamma} J_{\mu\rho} - \eta_{\mu\gamma} J_{\nu\rho} - \eta_{\nu\rho} J_{\mu\gamma}), \\
[T_r, T_s] &= f_{rs}^t T_t, \\
[P_\mu, T_s] &= [J_{\mu\nu}, T_s] = 0 \\
[Q_\alpha^L, J_{\mu\nu}] &= (\sigma_{\mu\nu})_\alpha{}^\beta Q_\beta^L, \\
[Q_\alpha^L, P_\mu] &= [\bar{Q}_\alpha^L, P_\mu] = 0 \\
\{Q_\alpha^L, \bar{Q}_{\beta M}\} &= 2(\sigma^\mu)_{\alpha\dot{\beta}} P_\mu \delta_M^L, \\
\{Q_\alpha^L, Q_\beta^M\} &= \epsilon_{\alpha\beta} Z^{LM}, \\
[Q_\alpha^L, T_r] &= S_r{}^L{}_M Q_\alpha^M \neq 0,
\end{aligned}$$

$$\begin{aligned}
[\delta_{\epsilon_1}, \delta_{\epsilon_2}] &= (\bar{\epsilon}_1 \gamma^\mu \epsilon_2) \partial_\mu + \dots \\
\delta_\epsilon B &\sim \bar{\epsilon} F, \\
\delta_\epsilon F &\sim B \epsilon. \\
\delta_\epsilon B &\sim \bar{\epsilon} F, \\
\delta_\epsilon F &\sim \partial \epsilon + B \epsilon. \\
S_{EH}[\mathbf{e}, \omega] &= \int \star \mathbf{R}(\omega) \wedge \mathbf{e} \wedge \mathbf{e}. \\
S &= \int d^4 x \sqrt{|g|} \{ R + \mathcal{G}_{ij}(\phi) \partial_\mu \phi^i \partial^\mu \phi^j + 2 \Im m \mathcal{N}_{\Lambda\Sigma}(\phi) F^\Lambda{}_{\mu\nu} F^{\Sigma\mu\nu} \\
&\quad - 2 \Re e \mathcal{N}_{\Lambda\Sigma}(\phi) F^\Lambda{}_{\mu\nu} \star F^{\Sigma\mu\nu} \}, \\
\mathcal{E}_{\mu\nu} &= G_{\mu\nu} + \mathcal{G}_{ij} \left[\partial_\mu \phi^i \partial_\nu \phi^j - \frac{1}{2} \partial_\rho \phi^i \partial^\rho \phi^j \right] + 8 \Im m \mathcal{N}_{\Lambda\Sigma} F_\mu^{\Lambda+\rho} F_{\nu\rho}^{\Sigma-} = 0, \\
\mathcal{E}_i &= \nabla_\mu (\mathcal{G}_{ij} \partial^\mu \phi^j) - \frac{1}{2} \partial_i \mathcal{G}_{jk} \partial_\rho \phi^j \partial^\rho \phi^k + \partial_i [\tilde{F}_\Lambda^{\nu\mu*} F_{\mu\nu}^\Lambda] = 0, \\
\mathcal{E}_\Lambda^\mu &= \nabla_\nu \star \tilde{F}_\Lambda^{\nu\mu} = 0, \\
\tilde{F}_{\Lambda\mu\nu} &\equiv -\frac{1}{4\sqrt{|g|}} \frac{\delta S}{\delta \star F^{\Lambda\mu\nu}} = \Re e \mathcal{N}_{\Lambda\Sigma} F_{\mu\nu}^\Sigma + \Im m \mathcal{N}_{\Lambda\Sigma}^* F_{\mu\nu}^\Sigma. \\
\mathcal{B}^{\Lambda\mu} &\equiv \nabla_\nu \star F^{\Lambda\nu\mu} = 0. \\
\mathcal{E}_\mu^M &\equiv \begin{pmatrix} \mathcal{B}_\mu^\Lambda \\ \mathcal{E}_{\Lambda\mu} \end{pmatrix} \\
\mathcal{E}_\mu^M = 0 &\rightarrow m^M{}_N \mathcal{E}_\mu^N = 0, m^M{}_N \in \text{GL}(2n_v + 2, \mathbb{R}). \\
F_\mu^M &\equiv \begin{pmatrix} F^\Lambda \\ \tilde{F}_\Lambda \end{pmatrix}, F'^M = m^M{}_N F^N. \\
\tilde{F}'_{\Lambda\mu\nu} &\equiv -\frac{1}{4\sqrt{|g|}} \frac{\delta S'}{\delta \star F'^{\Lambda\mu\nu}}. \\
\mathfrak{i}: \text{Diff}(\mathcal{M}_{\text{escalar}}) &\rightarrow \text{GL}(2n_v + 2, \mathbb{R}) \\
\{\phi, F^M, \mathcal{N}_{\Sigma\Lambda}(\phi)\} &\xrightarrow{\xi} \{\xi(\phi), (\mathfrak{i}(\xi))^M{}_N F^N, \mathcal{N}'_{\Sigma\Lambda}(\xi(\phi))\}. \\
\mathfrak{i}: \text{Diff}(\mathcal{M}_{\text{escalar}}) &\rightarrow \text{Sp}(2n_v + 2, \mathbb{R}) \\
\mathcal{N}'(\xi(\phi)) &= (A\mathcal{N}(\phi) + B)(C\mathcal{N}(\phi) + D)^{-1} \\
m &\equiv \begin{pmatrix} D & C \\ B & A \end{pmatrix} \in \text{Sp}(2n_v + 2, \mathbb{R}) \\
\mathfrak{i}: \text{Isometrías}(\mathcal{M}_{\text{escalar}}, \mathcal{G}_{ij}) &\rightarrow \text{Sp}(2n_v + 2, \mathbb{R}). \\
S &= \int d^4 x \sqrt{|g|} \{ R + h_{uv}(q) \partial_\mu q^u \partial^\mu q^v + \mathcal{G}_{i\bar{j}}(z, \bar{z}) \partial_\mu z^i \partial^\mu \bar{z}^{\bar{j}} \\
&\quad + 2 \Im_{\Lambda\Sigma}(z, \bar{z}) F^\Lambda{}_{\mu\nu} F^{\Sigma\mu\nu} - 2 \Re_{\Lambda\Sigma}(z, \bar{z}) F^\Lambda{}_{\mu\nu} \star F^{\Sigma\mu\nu} \} \\
\mathcal{F} &= -\frac{1}{3!} \kappa_{ijk}^0 z^i z^j z^k + \frac{ic}{2} + \frac{i}{(2\pi)^3} \sum_{\{d_i\}} n_{\{d_i\}} Li_3 \left(e^{2\pi i d_i z^i} \right) \\
\mathcal{F}_P &= -\frac{1}{3!} \kappa_{ijk}^0 z^i z^j z^k + \frac{ic}{2}, \\
\mathcal{F}_{NP} &= \frac{i}{(2\pi)^3} \sum_{\{d_i\}} n_{\{d_i\}} Li_3 \left(e^{2\pi i d_i z^i} \right). \\
F(\mathcal{X}) &= -\frac{1}{3!} \kappa_{ijk}^0 \frac{\mathcal{X}^i \mathcal{X}^j \mathcal{X}^k}{\mathcal{X}^0} + \frac{ic(\mathcal{X}^0)^2}{2} + \frac{i(\mathcal{X}^0)^2}{(2\pi)^3} \sum_{\{d_i\}} n_{\{d_i\}} Li_3 \left(e^{2\pi i d_i \frac{\mathcal{X}^i}{\mathcal{X}^0}} \right)
\end{aligned}$$

$$\begin{aligned}
z^i &= \frac{\mathcal{X}^i}{\mathcal{X}^0} \\
\mathcal{B} &\equiv \mathcal{M} - I^-(\mathcal{I}^+), \\
\nabla^\mu \xi^2 &= -2\kappa \xi^\mu. \\
\delta M &= \frac{1}{8\pi} \kappa \delta A + \Omega \delta J + \Phi \delta Q \\
\delta A &\geq 0. \\
T &= \frac{\kappa}{2\pi}. \\
S_{\text{bh}} &= \frac{A}{4}. \\
\delta_\epsilon B &\sim \bar{\epsilon} F = 0, \\
\delta_\epsilon F &\sim \partial \epsilon + B \epsilon = 0, \\
\mathbf{g} &= \left(1 - \frac{2M}{r}\right) dt \otimes dt - \left(1 - \frac{2M}{r}\right)^{-1} dr \otimes dr - r^2 (d\theta \otimes d\theta + \sin^2 \theta d\phi \otimes d\phi) \\
\delta_\epsilon \Psi_\mu|_{\text{Schw.}} &= 0 \\
\mathbf{g} &= \left(1 - \frac{2M}{r} + \frac{q^2}{r^2}\right) dt \otimes dt - \left(1 - \frac{2M}{r} + \frac{q^2}{r^2}\right)^{-1} dr \otimes dr - r^2 (d\theta \otimes d\theta + \sin^2 \theta d\phi \otimes d\phi) \\
T &= \frac{\kappa}{2\pi} = \frac{r_+ - r_-}{4\pi r_+^2} = \frac{\sqrt{M^2 - q^2}}{2\pi^2 r_+^2} = 0. \\
\delta_\epsilon \Psi_\mu|_{\text{RN extr.}} &= 0 \\
\mathbf{g} &= e^{2U(\tau)} dt \otimes dt - e^{-2U(\tau)} \gamma_{mn} dx^m \otimes dx^n, \\
\gamma_{mn} dx^m \otimes dx^n &= \frac{r_0^2}{\sinh^2 r_0 \tau} \left[\frac{r_0^2}{\sinh^2 r_0 \tau} d\tau \otimes d\tau + h_{\mathcal{S}^2} \right] \\
h_{\mathcal{S}^2} &= d\theta \otimes d\theta + \sin^2 \theta d\phi \otimes d\phi, \\
\mathbf{g} &= e^{2\tilde{U}(\tau)} dt \otimes dt - e^{-2U(\tau)} [\delta_{ab} dx^a \otimes dx^b], \\
\lim_{\tau \rightarrow -\infty} e^{-2U} &= \frac{A}{4\pi} \lim_{\tau \rightarrow -\infty} \tau^2, \quad \lim_{\tau \rightarrow -\infty} \tau \frac{d\phi^i}{d\tau} = 0, \quad i = 1, \dots, n_v \\
\lim_{\tau \rightarrow -\infty} \phi^i &= \phi_h^i, \quad \mathcal{G}^{ij}(\phi_h) \partial_j V_{\text{bh}}(\phi_h) = 0 \\
\partial_j V_{\text{bh}}(\phi_h) &= 0, \\
S &= \pi V_{\text{bh}}(\phi_h(Q)) = 0, \\
I_{\text{FGK}}[U, z^i] &= \int d\tau \{ (\dot{U})^2 + \mathcal{G}_{ij} \dot{z}^i \dot{\bar{z}}^j - e^{2U} V_{\text{bh}}(z, \bar{z}, Q) \}, \\
(\dot{U})^2 + \mathcal{G}_{ij} \dot{z}^i \dot{\bar{z}}^j + e^{2U} V_{\text{bh}}(z, \bar{z}, Q) &= r_0^2. \\
V_{\text{bh}}(z, \bar{z}, Q) &\equiv \frac{1}{2} \mathcal{M}_{MN}(\mathcal{N}) \mathcal{Q}^M \mathcal{Q}^N, \\
(Q^M) &= \begin{pmatrix} p^\Lambda \\ q_\Lambda \end{pmatrix}, \\
(\mathcal{M}_{MN}(\mathcal{N})) &\equiv \begin{pmatrix} I + RI^{-1}R & -RI^{-1} \\ -I^{-1}R & I^{-1} \end{pmatrix} \\
X &\equiv \frac{1}{\sqrt{2}} e^{U+i\alpha} \\
\mathcal{F}_\Lambda &\equiv \frac{\partial \mathcal{F}}{\partial \mathcal{X}^\Lambda} \text{ y } \mathcal{F}_{\Lambda\Sigma} \equiv \frac{\partial^2 \mathcal{F}}{\partial \mathcal{X}^\Lambda \partial \mathcal{X}^\Sigma}, \text{ se tiene } \mathcal{F}_\Lambda = \mathcal{F}_{\Lambda\Sigma} \mathcal{X}^\Sigma \\
(\mathcal{V}^M) &= \begin{pmatrix} \mathcal{L}^\Lambda \\ \mathcal{M}_\Lambda \end{pmatrix} = e^{\mathcal{K}/2} \begin{pmatrix} \mathcal{X}^\Lambda \\ \mathcal{F}_\Lambda \end{pmatrix}
\end{aligned}$$

$$\begin{aligned}
\frac{\mathcal{M}^\Lambda}{X} &= \mathcal{F}_{\Lambda\Sigma} \frac{\mathcal{L}^\Lambda}{X}. \\
\mathcal{R}^M &\equiv \Re(\mathcal{V}^M/X), \mathcal{J}^M \equiv \Im(\mathcal{V}^M/X), \\
\mathcal{R}^M &= -\mathcal{M}_{MN}(\mathcal{F})\mathcal{J}^M \\
d\mathcal{R}^M &= -\mathcal{M}_{MN}(\mathcal{F})d\mathcal{J}^M \\
\frac{\partial \mathcal{J}^M}{\partial \mathcal{R}_N} &= \frac{\partial \mathcal{J}^N}{\partial \mathcal{R}_M} = -\frac{\partial \mathcal{R}^M}{\partial \mathcal{J}_N} = -\frac{\partial \mathcal{R}^N}{\partial \mathcal{J}_M} = -\mathcal{M}^{MN}(\mathcal{F}). \\
H^M &\equiv \mathcal{J}^M(X, z, \bar{X}, \bar{z}) \\
z^i &= \frac{\mathcal{V}^i/X}{\mathcal{V}^0/X} = \frac{\tilde{H}^i(H) + iH^i}{\tilde{H}^0(H) + iH^0}, e^{-2U} = \frac{1}{2|X|^2} = \tilde{H}_M(H)H^M. \\
\dot{\alpha} &= 2|X|^2 \dot{H}^M H_M - \left[\frac{1}{2i} \dot{z}^i \partial_i \mathcal{K} + c.c. \right]. \\
\mathbf{W}(H) &\equiv \tilde{H}_M(H)H^M = e^{-2U}, \\
\tilde{H}_M &= \frac{1}{2} \frac{\partial W}{\partial H^M} \equiv \frac{1}{2} \partial_M W, H^M = \frac{1}{2} \frac{\partial W}{\partial \tilde{H}^M}. \\
-I_{\text{H-FGK}}[H] &= \int d\tau \left\{ \frac{1}{2} g_{MN} \dot{H}^M \dot{H}^N - V \right\} \\
r_0^2 &= \frac{1}{2} g_{MN} \dot{H}^M \dot{H}^N + V \\
g_{MN} &\equiv \partial_M \partial_N \log W - 2 \frac{H_M H_N}{W} \\
V(H) &\equiv \left\{ -\frac{1}{4} g_{MN} + \frac{H_M H_N}{2W^2} \right\} Q^M Q^N \\
V_{\text{bh}} &= -WV. \\
g_{MN} \ddot{H}^N + [PQ, M] \dot{H}^P \dot{H}^Q + \partial_M V &= 0 \\
[PQ, M] &\equiv \partial_{(P} g_{Q)M} - \frac{1}{2} \partial_M g_{PQ} \\
\tilde{H}_M (\ddot{H}^M - r_0^2 H^M) + \frac{(\dot{H}^M H_M)^2}{W} &= 0, \\
\dot{H}^M H_M &= 0, \\
\tilde{H}_M (\ddot{H}^M - r_0^2 H^M) &= 0, \\
H^M &= A^M - \frac{B^M}{\sqrt{2}} \tau, \\
H^M &= A^M \cosh(r_0 \tau) + B^M \sinh(r_0 \tau) \\
H^M &= A^M - \frac{Q^M}{\sqrt{2}} \tau, \\
H^0 = H_0 = H_i = 0, p^0 = p_0 = q_i = 0, \\
H^i &= a^i - \frac{p^i}{\sqrt{2}} \tau, r_0 = 0, \\
H^M &= H^M(a, b), \\
\{H^P = 0, Q^P = 0\} &\Rightarrow \mathcal{E}_P = 0, \\
\begin{pmatrix} iH^i \\ \tilde{H}_i \end{pmatrix} &= \frac{e^{\mathcal{K}/2}}{X} \begin{pmatrix} \mathcal{X}^i \\ \frac{\partial F(\mathcal{X})}{\partial \mathcal{X}^i} \end{pmatrix}, \begin{pmatrix} \tilde{H}^0 \\ 0 \end{pmatrix} = \frac{e^{\mathcal{K}/2}}{X} \begin{pmatrix} \mathcal{X}^0 \\ \frac{\partial F(\mathcal{X})}{\partial \mathcal{X}^0} \end{pmatrix}, \\
e^{-2U} &= \tilde{H}_i H^i, z^i = i \frac{H^i}{\tilde{H}^0}, \\
\frac{\partial F(H)}{\partial \tilde{H}^0} &= 0
\end{aligned}$$

$$\begin{aligned}
F(H) &= \frac{i}{3!} \kappa_{ijk}^0 \frac{H^i H^j H^k}{\tilde{H}^0} + \frac{ic(\tilde{H}^0)^2}{2} + \frac{i(\tilde{H}^0)^2}{(2\pi)^3} \sum_{\{d_i\}} n_{\{d_i\}} Li_3 \left(e^{-2\pi d_i \frac{H^i}{\tilde{H}^0}} \right) \\
\tilde{H}_i &= -i \frac{\partial F(H)}{\partial H^i}, \\
&-\frac{1}{3!} \kappa_{ijk}^0 \frac{H^i H^j H^k}{(\tilde{H}^0)^3} + c + \frac{1}{4\pi^3} \sum_{\{d_i\}} n_{\{d_i\}} \left[Li_3 \left(e^{-2\pi d_i \frac{H^i}{\tilde{H}^0}} \right) + Li_2 \left(e^{-2\pi d_i \frac{H^i}{\tilde{H}^0}} \right) \left[\frac{\pi d_i H^i}{\tilde{H}^0} \right] \right] \\
&\lim_{|w| \rightarrow 0} Li_s(w) = w, \forall s \in \mathbb{N} \\
&-\frac{1}{3!} \kappa_{ijk}^0 \frac{H^i H^j H^k}{(\tilde{H}^0)^3} + c + \frac{1}{4\pi^3} \sum_{\{d_i\}} n_{\{d_i\}} \left[e^{-2\pi d_i \frac{H^i}{\tilde{H}^0}} + e^{-2\pi d_i \frac{H^i}{\tilde{H}^0}} \left[\frac{\pi d_i H^i}{\tilde{H}^0} \right] \right] = 0, \Im m z^i \gg 1 \\
&-\frac{1}{3!} \kappa_{ijk}^0 \frac{H^i H^j H^k}{(\tilde{H}^0)^3} + c + \frac{1}{4\pi^3} \sum_{\{d_i\}} n_{\{d_i\}} e^{-2\pi d_i \frac{H^i}{\tilde{H}^0}} \left[\frac{\pi d_i H^i}{\tilde{H}^0} \right] \\
&-\frac{1}{3!} \kappa_{ijk}^0 \frac{H^i H^j H^k}{(\tilde{H}^0)^3} + c + \frac{\hat{n}}{4\pi^3} e^{-2\pi \hat{d}_i \frac{H^i}{\tilde{H}^0}} \left[\frac{\pi \hat{d}_i H^i}{\tilde{H}^0} \right] \\
e^{-2U} &= \mathbf{W}(H) = \alpha |\kappa_{ijk}^0 H^i H^j H^k|^{2/3} \\
V_{\text{bh}} &= \frac{W(H)}{4} \partial_{ij} \log W(H) Q^i Q^j, \\
z^i &= i(3!c)^{1/3} \frac{H^i}{(\kappa_{ijk}^0 H^i H^j H^k)^{1/3}}, \\
\kappa_{ijk}^0 \Im m z^i \Im m z^j \Im m z^k &> \frac{3c}{2}. \\
c &> \frac{c}{4} \\
e^{-\mathcal{K}} &= 6c \\
c > 0 &\Rightarrow h^{11} > h^{21} \\
X^{3,1} &\Rightarrow \kappa_{111}^0 = 48, \kappa_{222}^0 = \kappa_{333}^0 = 8 \\
Y^{3,1} &\Rightarrow \kappa_{122}^0 = 6, \kappa_{222}^0 = 18, \kappa_{333}^0 = 8 \\
\mathbf{W}(H) &= \alpha |48(H^1)^3 + 8[(H^2)^3 + (H^3)^3]|^{2/3}, \\
\mathbf{W}(H) &= \alpha |18(H^2)^2[H^1 + H^2] + 8(H^3)^3|^{2/3}. \\
H &= a \cosh(r_0 \tau) + \frac{b}{r_0} \sinh(r_0 \tau), b = s_b \sqrt{r_0^2 a^2 + \frac{p^2}{2}}, \\
z^1 &= i(3!c)^{1/3} \lambda^{-1/3} = s_{2,3} z^{2,3} \\
\lambda &= [48 + 8(s_2 + s_3)] \text{ para } X^{3,1} \\
\lambda &= [18 + 18s_2 + 8s_3] \text{ para } Y^{3,1}. \\
ds^2 &= \left[\frac{1}{2} (3!c)^{1/3} \left[a \cosh(r_0 \tau) + \frac{b}{r_0} \sinh(r_0 \tau) \right]^2 \right]^{-1} dt^2 \\
&-\frac{1}{2} (3!c)^{1/3} \left[a \cosh(r_0 \tau) + \frac{b}{r_0} \sinh(r_0 \tau) \right]^2 \left[\frac{r_0^4}{\sinh^4 r_0 \tau} d\tau^2 + \frac{r_0^2}{\sinh^2 r_0 \tau} d\Omega_{(2)}^2 \right] \\
2\pi i d_i z^i &\sim -\frac{1}{3} \sum_{i=1}^3 d_i, d_i \geq 1
\end{aligned}$$

$$s_2 = s_3 = -1, \text{ para } X^{3,1},$$

$$s_2 = -s_3 = 1, \text{ para } Y^{3,1}.$$

$$\mathcal{G}_{ij^*} = \partial_i \partial_{j^*} \mathcal{K}$$

$$a = -s_b \frac{\Im m z^1}{\sqrt{3c}}.$$

$$M = r_0 \sqrt{1 + \frac{3cp^2}{2r_0^2(\Im m z^1)^2}},$$

$$S_{\pm} = r_0^2 \pi \left(\sqrt{1 + \frac{3cp^2}{2r_0^2(\Im m z^1)^2}} \pm 1 \right)^2.$$

$$S_+ S_- = \frac{\pi^2 \alpha^2}{4} p^4 \lambda^{4/3}.$$

$$e^{-2U} = \mathbf{W}(H) = \alpha |H^1 H^2 H^3|^{2/3},$$

$$z^i = i c^{1/3} \frac{H^i}{(H^1 H^2 H^3)^{1/3}}.$$

$$H^i = a^i \cosh (r_0 \tau) + \frac{b^i}{r_0} \sinh (r_0 \tau), b^i = s_b^i \sqrt{r_0^2 (a^i)^2 + \frac{(p^i)^2}{2}}.$$

$$a^i = -s_b^i \frac{\Im m z_{\infty}^i}{\sqrt{3c}}.$$

$$M = \frac{r_0}{3} \sum_i \sqrt{1 + \frac{3c(p^i)^2}{2r_0^2(\Im m z_{\infty}^i)^2}}$$

$$S_{\pm} = r_0^2 \pi \prod_i \left(\sqrt{1 + \frac{3c(p^i)^2}{2r_0^2(\Im m z_{\infty}^i)^2}} \pm 1 \right)^{2/3}$$

$$S_+ S_- = \frac{\pi^2 \alpha^2}{4} \prod_i (p^i)^{4/3}$$

$$-\frac{1}{3!}\kappa_{ijk}^0\frac{H^iH^jH^k}{(\tilde{H}^0)^3}+\frac{\hat{n}}{4\pi^3}e^{-2\pi\hat{a}_l\frac{H^l}{\tilde{H}^0}}\left[\frac{\pi\hat{d}_nH^n}{\tilde{H}^0}\right]$$

$$\tilde{H}^0 = \frac{\pi \hat{d}_l H^l}{W_a \left(s_a \sqrt{\frac{3 \hat{n} (\hat{d}_n H^n)^3}{2 \kappa_{ijk}^0 H^i H^j H^k}} \right)},$$

$$\tilde{H}_i = \frac{1}{2} \kappa_{ijk}^0 \frac{H^j H^k}{\pi \hat{d}_l H^l} W_a \left(s_a \sqrt{\frac{3 \hat{n} (\hat{d}_m H^m)^3}{2 \kappa_{pqr}^0 H^p H^q H^r}} \right).$$

$$e^{-2U} = \mathbf{W}(H) = \frac{\kappa_{ijk}^0 H^i H^j H^k}{2\pi \hat{d}_m H^m} W_a \left(s_a \sqrt{\frac{3 \hat{n} (\hat{d}_l H^l)^3}{2 \kappa_{pqr}^0 H^p H^q H^r}} \right)$$

$$z^i = i \frac{H^i}{\pi \hat{d}_m H^m} W_a \left(s_a \sqrt{\frac{3 \hat{n} (\hat{d}_l H^l)^3}{2 \kappa_{pqr}^0 H^p H^q H^r}} \right).$$

$$\begin{aligned}
s_0 &\equiv \text{sign} \left[\kappa_{ijk}^0 \frac{H^i H^j H^k}{\hat{d}_m H^m} \right], \\
s_{-1} &\equiv -1 \\
H^i &= a^i - \frac{p^i}{\sqrt{2}} \tau, r_0 = 0. \\
S &= \frac{1}{2} \kappa_{ijk}^0 \frac{p^i p^j p^k}{\hat{d}_m p^m} W_a(s_a \beta), \\
\beta &= \sqrt{\frac{3 \hat{n} (\hat{d}_l p^l)^3}{2 \kappa_{pqr}^0 p^p p^q p^r}}, \\
M = \dot{U}(0) &= \frac{1}{2\sqrt{2}} \left[\frac{3 \kappa_{ijk}^0 p^i a^j a^k}{\kappa_{pqr}^0 a^p a^q a^r} \left[1 - \frac{1}{1 + W_a(s_a \alpha)} \right] - \frac{d_l p^l}{d_n a^n} \left[1 - \frac{3}{2(1 + W_a(s_a \alpha))} \right] \right], \\
\alpha &= \sqrt{\frac{3 \hat{n} (d_l a^l)^3}{2 \kappa_{pqr}^0 a^p a^q a^r}}. \\
W_a(x) e^{-2W_a(x)} &\gg e^{-2W_a(x)}. \\
e^{-2U} &= \frac{\kappa_{ijk}^0 H^i H^j H^k}{2\pi \hat{d}_m H^m} W_0 \left(\sqrt{\frac{3 \hat{n} (\hat{d}_l H^l)^3}{2 \kappa_{pqr}^0 H^p H^q H^r}} \right), \\
z^i &= i \frac{H^i}{\pi \hat{d}_m H^m} W_0 \left(\sqrt{\frac{3 \hat{n} (\hat{d}_l H^l)^3}{2 \kappa_{pqr}^0 H^p H^q H^r}} \right). \\
\frac{\kappa_{ijk}^0 H^i H^j H^k}{2\pi \hat{d}_n H^n} &> 0 \quad \forall \tau \in (-\infty, 0], \\
\frac{\kappa_{ijk}^0 a^i a^j a^k}{2\pi \hat{d}_m a^m} W_0(\alpha) &= 1, \\
e^{-2U} &\xrightarrow{\tau \rightarrow -\infty} \frac{\kappa_{ijk}^0 p^i p^j p^k}{8\pi \hat{d}_m p^m} W_0(\beta) \tau^2. \\
L &= \bigoplus_{k=0}^N L_k, \\
u_j \circ u_k &\in L_{(j+k) \bmod (N+1)}. \\
u_0 \circ v_0 &\in L_0, \forall u_0, v_0 \in L_0, \\
u_0 \circ u_1 &\in L_1, \forall u_0 \in L_0, v_1 \in L_1, \\
u_1 \circ v_1 &\in L_0, \forall u_1, v_1 \in L_1, \\
\bullet \text{ Supersimetrizaci3n: } &\forall x_i \in L_i, \forall x_j \in L_j, j = 0, 1, x_i \circ x_j = -(-1)^{ij} x_j \circ x_i. \\
\bullet \text{ Identidades de Jacobi: } &\forall x_k \in L_k, \forall x_l \in L_l, \forall x_m \in L_m, k, l, m = 0, 1, \quad x_k \circ (x_l \circ x_m) (-1)^{km} + x_l \circ (x_m \circ x_k) (-1)^{lk} + x_m \circ (x_k \circ x_l) (-1)^{ml} = 0. \\
&x_\mu \circ x_\nu \equiv x_\mu x_\nu - (-1)^{g_\mu g_\nu} x_\nu x_\mu = c_{\mu\nu}^\omega x_\omega \\
Li_w(z) &= \sum_{j=1}^{\infty} \frac{z^j}{j^w}, z, w \in \mathbb{C}. \\
Li_{w-1}(z) &= z \frac{\partial Li_w(z)}{\partial z}.
\end{aligned}$$

$$\begin{aligned}
Li_1(z) &= -\log(1-z) \\
Li_0(z) &= \frac{z}{1-z}, Li_{-n}(z) = \left(z \frac{\partial}{\partial z}\right)^n \frac{z}{1-z}. \\
Li_w(z) &= \int_0^z \frac{Li_{w-1}(s)}{s} ds \\
z &= W(z)e^{W(z)}, \forall z \in \mathbb{C}. \\
\frac{dW(z)}{dz} &= \frac{W(z)}{z(1+W(z))}, \forall z \notin \{0, -1/e\}, \left. \frac{dW(z)}{dz} \right|_{z=0} \\
\lim_{x \rightarrow -1/e} \frac{dW_0(x)}{dx} &= \infty, \lim_{x \rightarrow -1/e} \frac{dW_{-1}(x)}{dx} = -\infty.
\end{aligned}$$

$$\begin{aligned}
&X \times S \\
\mu^{1,1} &\in \Omega^{0,1}(X, \mathcal{T}_X) \otimes C^\infty(S) \\
\mu^{1,1} &= \mu_j^i(z, \bar{z}, t) d\bar{z}_i \frac{\partial}{\partial z_j} \\
\gamma^{1,0} &\in \Omega^{1,0}(X) \otimes C^\infty(S), \gamma^{1,2} \in \Omega^{1,2}(X) \otimes C^\infty(S) \\
\bar{\partial} \mu^{1,1} + \frac{1}{2} [\mu^{1,1}, \mu^{1,1}] + \Omega^{-1} \vee (\partial \gamma^{1,0} \wedge \partial \gamma^{1,2}) &= 0 \\
\Omega^{0,\cdot}(X, V) &= (\Omega^{0,j}(X, V)[-j], \bar{\partial}) \\
\partial_\Omega(\mu) \wedge \Omega &= L_\mu(\Omega) \\
\partial_\Omega: \Omega^{0,\cdot}(X, \mathcal{T}_X) &\rightarrow \Omega^{0,\cdot}(X) \\
\Omega^{0,\cdot}(X, \mathcal{T}_X) &\xrightarrow{\partial_\Omega} \Omega^{0,\cdot}(X) \\
[\mu, \mu'] &= L_\mu \mu' \\
[\mu, \nu] &= L_\mu \nu \\
(\text{Sym}(\mathcal{L}^\vee[-1]), \delta_{\mathcal{L}}) & \\
\Phi^*: C^\cdot(\mathcal{L}') &\rightarrow C^\cdot(\mathcal{L}) \\
\Psi_\infty: \nu &\mapsto 1 - e^{-\nu}, \mu \mapsto e^{-\nu} \mu \\
[\mu]_1 &= \bar{\partial} \mu + \partial_\Omega \mu \\
[\mu_1, \mu_2]_2 &= \partial_\Omega(\mu_1 \wedge \mu_2) \\
[\nu, \mu_1, \mu_2]_3 &= \partial_\Omega(\nu \mu_1 \wedge \mu_2) \\
[\nu_1, \dots, \nu_{k-2}, \mu_1, \mu_2]_k &= \partial_\Omega(\nu_1 \cdots \nu_{k-2} \mu_1 \wedge \mu_2) \\
[\nu_1, \dots, \nu_{k-3}, \mu_1, \mu_2, \gamma]_k &= \nu_1 \cdots \nu_{k-3}(\mu \wedge \mu') \vee \partial \gamma. \\
[\nu_1, \dots, \nu_{k-2}, \mu, \gamma]_k &= \nu_1 \cdots \nu_{k-2} \mu \vee \partial \gamma. \\
[x \otimes a]_1 &= [x]_1^{\mathcal{L}} \otimes a + (-1)^{|x|} x \otimes d_{\mathcal{A}} a \\
[x_1 \otimes a_1, \dots, x_k \otimes a_k]_k &= [x_1, \dots, x_k]_k^{\mathcal{L}} \otimes (a_1 \cdots a_k), k \geq 2. \\
F_A &= [A]_1 + \frac{1}{2} [A, A]_2 + \frac{1}{3!} [A, A, A]_3 + \cdots \\
(A, B) &\in \mathcal{L}[1] \oplus \mathcal{L}^![-2] \\
\begin{matrix} -n & & -n+1 & -1 & & 0 \\ \Omega^0(X; S) & \xrightarrow{\partial} & \Omega^1(X; S) & & \text{PV}^1(X; S) & \xrightarrow{\partial_\Omega} & \text{PV}^0(X; S). \end{matrix} \\
\Omega^i(X; S) &= \Omega^{i,\cdot,\cdot}(X; S) \\
&= \bigoplus_{j,k} \text{PV}^{i,j}(X) \otimes \Omega^k(S)[-j-k] \\
n &= \dim_{\mathbb{C}}(X) + \dim_{\mathbb{R}}(S) - 1 \\
&\int_{X \times S}^{\Omega} \mu \vee \gamma + \int_{X \times S}^{\Omega} \nu \beta \\
S_{BF,0} &= \int^{\Omega} \left[\beta \wedge (\bar{\partial} + d_S) \nu + \gamma \wedge (\bar{\partial} + d_S) \mu + \beta \wedge \partial_\Omega \mu + \frac{1}{2} [\mu, \mu] \vee \gamma + [\mu, \nu] \beta \right] \left[\beta \wedge (\bar{\partial} + d_S) \nu + \gamma \wedge (\bar{\partial} + d_S) \mu + \beta \wedge \partial_\Omega \mu + \frac{1}{2} [\mu, \mu] \vee \gamma + [\mu, \nu] \beta \right].
\end{aligned}$$



$$S_{BF,\infty} = \int^{\Omega} \left[\beta \wedge (\bar{\partial} + d_S) \nu + \gamma \wedge (\bar{\partial} + d_S) \mu + \beta \wedge \partial_{\Omega} \mu + \frac{1}{2} \frac{1}{1-\nu} \mu^2 \vee \partial \gamma \right] \left[\beta \wedge (\bar{\partial} + d_S) \nu + \gamma \wedge (\bar{\partial} + d_S) \mu + \beta \wedge \partial_{\Omega} \mu + \frac{1}{2} \frac{1}{1-\nu} \mu^2 \vee \partial \gamma \right].$$

$$\mu \mapsto e^{-\nu} \mu, \nu \mapsto 1 - e^{-\nu}, \beta \mapsto (\beta - \mu \vee \gamma) e^{\nu}, \gamma \mapsto e^{\nu} \gamma$$

$$\Omega^0(X;S)_{\beta} \stackrel{\partial}{\rightarrow} \Omega^1(X;S)_{\gamma} \qquad \qquad \qquad \text{PV}^1(X;S)_{\mu} \stackrel{\partial_{\Omega}}{\rightarrow} \text{PV}^0(X;S)_{\nu}.$$

$$\begin{aligned} S_{BF} + gJ \\ \{S_{BF} + gJ, S_{BF} + gJ\} = 0 \\ \{S_{BF}, J\} = \{J, J\} = 0 \end{aligned}$$

$$J=\frac{1}{6}\gamma\wedge\partial\gamma\wedge\partial\gamma$$

$$\gamma^{1,i;j}\in\Omega^{1,i}(X)\otimes\Omega^j(S)$$

$$\deg(J)=6$$

$$\{\beta\wedge\partial_{\Omega}\mu,J\}=\frac{1}{2}\partial\beta\wedge\partial\gamma\wedge\partial\gamma=0$$

$$\left\{\frac{1}{2}\frac{1}{1-\nu}\partial\gamma\vee\mu^2,\frac{1}{6}\gamma\wedge\partial\gamma\wedge\partial\gamma\right\}=\frac{1}{2}(\mu\vee\partial\gamma)\wedge\partial\gamma\wedge\partial\gamma$$

$$\mathrm{Sym}^3(\boxplus)\cong\boxplus\oplus\boxplus$$

$$\wedge^3(\exists)\cong\boxplus\oplus\boxplus$$

$$\gamma\mapsto\sqrt{g}\gamma,\beta\mapsto\sqrt{g}\beta$$

$$\frac{1}{\sqrt{g}}(S_{BF,\infty}+J)$$

$$\tilde{J}=\frac{1}{6}e^{\nu}\gamma\wedge\partial(e^{\nu}\gamma)\wedge\partial(e^{\nu}\gamma)$$

$$\bar{\partial}\nu+d_S\nu+\partial_{\Omega}\mu=0$$

$$\bar{\partial}\mu+d_S\mu+\frac{1}{2}\frac{1}{1-\nu}\partial_{\Omega}(\mu^2)+\frac{1}{2}(\partial\gamma\wedge\partial\gamma)\vee(g\Omega^{-1})=0$$

$$(\bar{\partial}+d_S)\gamma+\partial\beta+\frac{1}{1-\nu}(\mu\vee\partial\gamma)=0$$

$$(\bar{\partial}+d_S)\beta+\frac{1}{2}\frac{1}{(1-\nu)^2}\mu^2\vee\partial\gamma=0$$

- $\mu=\sum_{i,j}\mu^{i;j}\mu^{i;j}\in\text{PV}^{1,i}(X)\otimes\Omega^j(\mathbb{R}), i=0,\dots,5, j=0,1$
- $\nu=\sum_{i,j}\nu^{i;j}\nu^{i;j}\in\text{PV}^{0,i}(X,\text{ T}_X)\otimes\Omega^j(\mathbb{R}), i=0,\dots,5, j=0,1$
- $\gamma=\sum_{i,j}\gamma^{i;j}\gamma^{i;j}\in\Omega^{1,i}(X)\otimes\Omega^j(\mathbb{R}), i=0,\dots,5, j=0,1$
- $\beta=\sum_{i,j}\beta^{i;j}\beta^{i;j}\in\Omega^{0,i}(X)\otimes\Omega^j(\mathbb{R}), i=0,\dots,5, j=0,1$

$$\bar{\partial}\mu^{1;0}+\frac{1}{2}[\mu^{1;0},\mu^{1;0}]+\left(\frac{1}{2}\partial\gamma^{1;0}\wedge\partial\gamma^{1;0}+\partial\gamma^{2;0}\wedge\partial\gamma^{0;0}\right)\vee(g\Omega^{-1})=0$$

$$\alpha\stackrel{\text{def}}{=} \partial\gamma^{0;0}$$

$$\Omega_X^{2,\,\text{hol}}\stackrel{\wedge\alpha}{\rightarrow}\Omega_X^{4,\,\text{hol}}\cong\Omega\mathcal{T}_X^{\text{hol}}$$

$$\bar{\partial}\xi+\frac{1}{2}[\xi,\xi]=\alpha\vee\rho$$

$$\mu^{1;0}=\mu_i^j(z,\bar{z},t)\mathrm{d}\bar{z}_j\partial_{z_i}$$

$$\mu^{0;1}=\mu_i^t(z,\bar{z},t)\mathrm{d}t\partial_{z_i}$$

$$\mu^{0;0}=\mu_i(z,\bar{z},t)\partial_{z_i}$$

$$\beta^{3;0}=\beta^{ijk}(z,\bar{z},t)\mathrm{d}\bar{z}_i\,\mathrm{d}\bar{z}_j\,\mathrm{d}\bar{z}_k,\beta^{2;1}=\beta_t^{ij}(z,\bar{z},t)\mathrm{d}\bar{z}_i\,\mathrm{d}\bar{z}_j\,\mathrm{d}t$$

$$\gamma^{2;0}=\gamma^{ijk}(z,\bar{z},t)\mathrm{d}z_i\,\mathrm{d}\bar{z}_j\,\mathrm{d}\bar{z}_k,\gamma^{1;1}=\gamma_t^{ij}(z,\bar{z},t)\mathrm{d}z_i\,\mathrm{d}\bar{z}_j\,\mathrm{d}t.$$



$$\begin{aligned}
&\mu^i \partial_{z_i} \in \text{Vect}(\mathbb{C}^5) \cong \mathcal{O}(\mathbb{C}^5) \partial_{z_i}, \nu \in \mathcal{O}(\mathbb{C}^5) \\
&\beta \in \mathcal{O}(\mathbb{C}^5), \gamma^i \, dz_i \in \Omega^1(\mathbb{C}^5) \cong \mathcal{O}(\mathbb{C}^5) dz_i \\
&\mu_{(m_j)}^i: \mu^i \mapsto \partial_{z_1}^{m_1} \partial_{z_2}^{m_2} \partial_{z_3}^{m_3} \partial_{z_4}^{m_4} \partial_{z_5}^{m_5} \mu^i \\
&\nu_{(m_j)}: \nu \mapsto \partial_{z_1}^{m_1} \partial_{z_2}^{m_2} \partial_{z_3}^{m_3} \partial_{z_4}^{m_4} \partial_{z_5}^{m_5} \nu \\
&\gamma_{(m_j)}^i: \gamma^i \mapsto \partial_{z_1}^{m_1} \partial_{z_2}^{m_2} \partial_{z_3}^{m_3} \partial_{z_4}^{m_4} \partial_{z_5}^{m_5} \gamma^i \\
&\beta_{(m_j)}: \beta \mapsto \partial_{z_1}^{m_1} \partial_{z_2}^{m_2} \partial_{z_3}^{m_3} \partial_{z_4}^{m_4} \partial_{z_5}^{m_5} \beta \\
&i(q_1, \dots, q_5) = \frac{\sum_{i=1}^5 q_i}{\prod_{i=1}^5 (1 - q_i)} + \frac{\sum_{i=1}^5 q_i^{-1}}{\prod_{i=1}^5 (1 - q_i^{-1})} \\
&\quad \frac{q_1^{m_1+1} \dots q_i^{m_i} \dots q_5^{m_5+1}}{q_1^{m_1} \dots q_i^{m_i+1} \dots q_5^{m_5}} \\
&\sum_{i=1}^5 \left(\sum_{(m_i) \in \mathbb{Z}_{\geq 0}^5} q_1^{m_1} \dots q_i^{m_i+1} \dots q_5^{m_5} - \sum_{(m_i) \in \mathbb{Z}_{\geq 0}^5} q_1^{m_1+1} \dots q_i^{m_i} \dots q_5^{m_5+1} \right) \\
&\quad - \frac{\sum_{i=1}^5 q_1 \dots \hat{q}_i \dots q_5}{\prod_{i=1}^5 (1 - q_i)} + \frac{\sum_{i=1}^5 q_i}{\prod_{i=1}^5 (1 - q_i)} \\
&\prod_{i=1}^5 \prod_{(m_i) \in \mathbb{Z}_{\geq 0}^5} \frac{1 - q_1^{m_1+1} \dots q_i^{m_i} \dots q_5^{m_5+1}}{1 - q_1^{m_1} \dots q_i^{m_i+1} \dots q_5^{m_5}} \\
&E(5,10)_+ = \text{Vect}_0(\mathbb{C}^5) \\
&E(5,10)_- = \Omega_{cl}^2(\mathbb{C}^5) \\
&[\alpha, \alpha'] = \Omega^{-1} \vee (\alpha \wedge \alpha') \\
&\delta = \{S_{BF, \infty} + J, -\} \\
&\delta^{(1)} = \bar{\partial} + d_{\mathbb{R}} + \partial_{\Omega}|_{\mu \rightarrow \nu} + \partial|_{\beta \rightarrow \gamma} \\
&\text{Vect}(\mathbb{C}^5) \xrightarrow{\partial_{\Omega}} \mathcal{O}(\mathbb{C}^5) \\
&\mathcal{O}(\mathbb{C}^5) \xrightarrow{\partial} \Omega^1(\mathbb{C}^5) \\
&\mu = \mu \otimes 1 \in \Pi \text{Vect}_0(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}) \\
&[\gamma] = [\gamma] \otimes 1 \in (\Omega^1(\mathbb{C}^5)/d(\mathcal{O}(\mathbb{C}^5))) \otimes \Omega^0(\mathbb{R}) \\
&(\mu, [\gamma], b) \in \text{Vect}_0(\mathbb{C}^5) \oplus \Pi \Omega^1(\mathbb{C}^5)/\partial \mathcal{O}(\mathbb{C}^5) \oplus \mathbb{C} \\
&[[\gamma], [\gamma']] = \Omega^{-1} \vee (\partial \gamma \wedge \partial \gamma') \in \text{Vect}_0(\mathbb{C}^5) \\
&\partial: \Omega^1(\mathbb{C}^5)/d\left(\begin{smallmatrix} \mathbb{C}^5 \\ q \end{smallmatrix}\right) \xrightarrow{\cong} \Omega_{cl}^2(\mathbb{C}^5) \\
&_K \leftrightarrow (\Pi \mathcal{E}, \delta^{(1)}) \rightleftharpoons (E(5,10) \oplus \mathbb{C}_b, 0), \\
&\quad \tilde{K} \partial \gamma + \partial K \gamma = \gamma \\
&\quad \gamma - \tilde{K} \partial \gamma = \partial K \gamma \\
&\quad [\mu, \mu', [\gamma]]_3 = \varphi(\mu, \mu', [\gamma]) \\
&\varphi: E(5,10) \times E(5,10) \times E(5,10) \rightarrow \mathbb{C}_b \\
&\quad \varphi(\mu, \mu', \alpha) = \langle \mu \wedge \mu', \alpha \rangle|_{z=0} \\
&\mathbb{C}^{\text{even}}(\mathcal{L}) = \mathbb{C}^{2 \cdot, +}(\mathcal{L}) \oplus \mathbb{C}^{2 \cdot + 1, -}(\mathcal{L}) \\
&\mathbb{C}^{\text{odd}}(\mathcal{L}) = \mathbb{C}^{2 \cdot, -}(\mathcal{L}) \oplus \mathbb{C}^{2 \cdot + 1, +}(\mathcal{L}) \\
&\quad t_{11d} = V \oplus \Pi S \\
&\text{Sym}^2(S) \cong V \oplus \wedge^2 V \oplus \wedge^5 V
\end{aligned}$$

$$\begin{aligned}
[\psi, \psi'] &= \Gamma_{\wedge^1}(\psi, \psi') \\
\mathfrak{so}_{11d} &= \mathfrak{so}(11, \mathbb{C}) \ltimes \mathfrak{t}_{11d} \\
c_{M2} &\in \mathbb{C}^{2,+}(\mathfrak{so}_{11d}; \Omega^1(\mathbb{R}^{11})[2]) \\
c_{M2}(\psi, \psi') &= \Gamma_{\wedge^2}(\psi, \psi') \in \Omega^2(\mathbb{R}^{11}) \\
V &= L \oplus L^\vee \oplus \mathbb{C}_t, S = \wedge^1 L \\
&\mathfrak{sl}(L) \oplus \wedge^2 L \oplus \wedge^2 L^\vee \oplus L \oplus L^\vee \oplus \mathbb{C} \\
S = \wedge^1(L) &= \mathbb{C} \oplus L \oplus \wedge^2 L \oplus \wedge^3 L \oplus \wedge^4 L \oplus \wedge^5 L \\
\text{Stab}(Q) &= \mathfrak{sl}(L) \oplus \wedge^2 L^\vee \oplus L^\vee \subset \mathfrak{so}(11, \mathbb{C}) \\
&L \oplus \text{Stab}(Q) \oplus \Pi(\wedge^2 L^\vee) \oplus \mathbb{C} \\
[Q, -] &: \mathfrak{so}(11, \mathbb{C}) \rightarrow S \\
[\psi, \psi']_2 &= \psi \wedge \psi' \in \wedge^4 L^\vee \cong L_v \\
[v, v', \psi]_3 &= 4 \langle v \wedge v', \psi \rangle \in \mathbb{C}_b \\
[z_i \wedge z_j, z_k \wedge z_l]_2 &= \epsilon_{ijklm} \partial_{z_m} \\
\left[\partial_{z_i}, \partial_{z_j}, z_k \wedge z_\ell \right]_3 &= 4(\delta_k^i \delta_\ell^j - \delta_\ell^i \delta_k^j) \\
H^*(\text{m2brane}^Q) \oplus (L^\vee \xrightarrow{\mathbb{B}} \Pi L^\vee) \\
[v, \lambda] &= \langle v, \lambda \rangle \in \mathbb{C}_b \\
[v, \psi] &= \langle v, \psi \rangle \in \Pi L_{\bar{\lambda}} \\
&\mathfrak{g} \rightsquigarrow \text{m2brane}^Q \\
H: \Omega^2(\mathbb{R}^{11}) &\rightarrow \Omega^1(\mathbb{R}^{11}) \\
\tilde{\psi} &= \psi - H\Gamma_{\wedge^2}(Q, \psi) \in \Pi S \oplus \Pi \Omega^1 \\
[v, \tilde{\psi}] &= -L_v(H\Gamma_{\wedge^2}(Q, \psi)) = -\langle v, \Gamma_{\wedge^2}(Q, \psi) \rangle - d\langle v, H\Gamma_{\wedge^2}(Q, \psi) \rangle \\
v \otimes \psi &\mapsto \langle v, H\Gamma_{\wedge^2}(Q, \psi) \rangle \in L_\lambda \\
K \cup (\mathfrak{g}, \delta) &\xrightarrow[i]{q} (H^*(\text{m2brane}^Q), 0) \\
H^*(\text{m2brane}^Q) &\rightsquigarrow \mathcal{L}(\mathbb{C}^5 \times \mathbb{R}) \\
L^\vee &\mapsto 0 \\
\wedge^2 L_1^\vee &\mapsto 0 \\
z_i \wedge z_j \in \wedge^2 L_2^\vee &\mapsto \frac{1}{2}(z_i dz_j - z_j dz_i) \in \Omega^{1,0}(\mathbb{C}^5) \widehat{\otimes} \Omega^0(\mathbb{R}) \\
A_{ij} \in \mathfrak{sl}(5) &\mapsto \sum_{ij} A_{ij} z_i \partial_{z_j} \in \text{PV}^{1,0}(\mathbb{C}^5) \widehat{\otimes} \Omega^0(\mathbb{R}) \\
\partial_{z_j} \in L &\mapsto \partial_{z_i} \in \text{PV}^{1,0}(\mathbb{C}^5) \widehat{\otimes} \Omega^0(\mathbb{R}^5) \\
1 \in \mathbb{C}_b &\mapsto 1 \in \Omega^{0,0}(\mathbb{C}^5) \widehat{\otimes} \Omega^0(\mathbb{R}). \\
[z_i \wedge z_j, z_k \wedge z_l] &= \epsilon_{ijklm} \partial_{z_m} \\
[\partial_{z_i}, z_j dz_k - z_k dz_j] &= \delta_j^i dz_k - \delta_k^i dz_j \\
\Phi^{(2)}(\partial_{z_i}, z_j \wedge z_k) &= \frac{1}{2}(\delta_j^i z_k - \delta_k^i z_j). \\
[\Phi^{(1)}(\partial_{z_i}), \Phi^{(1)}(z_j \wedge z_k)] &= \partial \Phi^{(2)}(\partial_{z_i}, z_j \wedge z_k) \\
\Phi^{(1)}[\partial_{z_i}, \partial_{z_j}, z_k \wedge z_l]_3 &= [\Phi^{(1)}(\partial_{z_i}), \Phi^{(1)}(\partial_{z_j}), \Phi^{(1)}(z_k \wedge z_l)]_3 \\
&+ [\partial_{z_i}, \Phi^{(2)}(\partial_{z_j}, z_k \wedge z_l)] + [\partial_{z_j}, \Phi^{(2)}(\partial_{z_i}, z_k \wedge z_l)] \\
H^*(\text{m2brane}^Q) &\rightarrow E(\overline{5}, 10) \\
L \oplus \text{Stab}(Q) \oplus \Pi(\wedge^3 L) \oplus \mathbb{C}_b &\rightarrow E(5, 10) \oplus \mathbb{C}_{b'}
\end{aligned}$$

$$\begin{aligned}
& L_1^\vee \mapsto 0 \\
& \wedge^2 L_1^\vee \mapsto 0 \\
& z_i \wedge z_j \in \wedge^2 L_2 \mapsto dz_i \wedge dz_j \in \Omega_{cl}^2(\mathbb{C}^5) \\
& A_{ij} \in \mathfrak{sl}(5) \mapsto \sum_{ij} A_{ij} z_i \partial_{z_j} \in \text{Vect}_0(\mathbb{C}^5) \\
& \partial_{z_i} \in L \mapsto \partial_{z_i} \in \text{Vect}_0(\mathbb{C}^5) \\
& b \in \mathbb{C}_b \mapsto b \in \mathbb{C}_{b'}. \\
& \mu_2(\psi, \psi', v, v') = \langle v \wedge v', \Gamma(\psi, \psi') \rangle \\
& A \in \Pi \Omega^{0,\cdot}(\mathbb{C}^2) \widehat{\otimes} \Omega^\cdot(\mathbb{R}^7), \\
& \bar{\partial} A + d_{\mathbb{R}^7} A + \partial_{z_1} A \wedge \partial_{z_2} A = 0 \\
& \mathbb{C}^2 \times \mathbb{R}^7 \\
& Q + Q_{nm} \\
& Q_{nm} \in \wedge^2(L^\vee) \\
& \mathbb{C}^5 \times \mathbb{R} = \mathbb{C}_{z_i}^2 \times \mathbb{C}_{w_a}^3 \times \mathbb{R} \\
& Q_{nm} = dz_1 \wedge dz_2 \\
& \prod_{(n_1, n_2) \in \mathbb{Z}_{\geq 0}^2} \frac{1}{1 - q^{-n_1 + n_2}} \\
& A_{(n_1, n_2)}: A \mapsto \partial_{w_1}^{n_1} \partial_{w_2}^{n_2} A(0) \\
& \sum_{(n_1, n_2) \in \mathbb{Z}_{\geq 0}^2} q^{-n_1 + n_2} \\
& \prod_{(n_1, n_2) \in \mathbb{Z}_{\geq 0}^2} \frac{1}{1 - q^{-n_1 + n_2}} \\
& q_1 q_2 = 1, q_3 q_4 q_5 = 1 \\
& i(q) = \frac{1}{(1 - q)(1 - q^{-1})} \\
& \gamma_{nm} = \frac{1}{2} (z_1 dz_2 - z_2 dz_1) \in \Omega^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}) \\
& dz_1 \wedge dz_2 \in \Omega_{cl}^2(\mathbb{C}^5) \\
& [f_i \partial_{z_i}, dz_1 \wedge dz_2] = \partial f_i \wedge dz_j - \partial f_j \wedge dz_i \\
& [g_a \partial_{w_a}, dz_1 \wedge dz_2] = 0 \\
& [h^{ab} dw_a \wedge dw_b, dz_1 \wedge dz_2] = \epsilon_{abc} h^{ab} \partial_{w_c}. \\
& \partial_{z_1} f + \partial_{z_2} g = 0 \\
& H(E(5,10), [dz_1 \wedge dz_2, -]) \simeq \text{Vect}_0(\mathbb{C}^2) \\
& 0 \rightarrow \mathbb{C} \rightarrow \mathcal{O}(\mathbb{C}^2) \rightarrow \text{Vect}_0(\mathbb{C}^2) \rightarrow 0 \\
& H(\widehat{E(5,10)}, [dz_1 \wedge dz_2, -]) \simeq \mathcal{O}(\mathbb{C}^2) \\
& \varphi(\mu, \mu', \alpha) = \langle \mu \wedge \mu', \alpha \rangle|_{z=0} \\
& (f_i \partial_{z_i}, g_j \partial_{z_j}) \mapsto (f_1 g_2 - f_2 g_1)(z_1 = z_2 = 0) \\
& Z \times M \\
& \alpha \in \Pi \Omega^{0,\cdot}(Z) \widehat{\otimes} \Omega^\cdot(M) \\
& \{\alpha^I(z, \bar{z}) d\bar{z}_I, \alpha^J(z, \bar{z}) d\bar{z}_J\}_{pb} = (\partial_{z_1} \alpha^I \partial_{z_2} \alpha^J \pm \partial_{z_2} \alpha^I \partial_{z_1} \alpha^J) d\bar{z}_I \wedge d\bar{z}_J \\
& \frac{1}{2} \int_{Z \times M} (\alpha \wedge d\alpha) \wedge \omega_Z^{2,0} + \frac{1}{6} \int_{Z \times M} \alpha \wedge \{\alpha, \alpha\}_{pb} \wedge \omega_Z^{2,0}
\end{aligned}$$



$$\begin{aligned}
& \mu_z \in \mathrm{PV}^{1,\cdot}(\mathbb{C}_z^2) \otimes \mathrm{PV}^{0,\cdot}(\mathbb{C}_w^3) \otimes \Omega^\cdot(\mathbb{R}) \\
& \mu_w \in \mathrm{PV}^{0,\cdot}(\mathbb{C}_z^2) \otimes \mathrm{PV}^{1,\cdot}(\mathbb{C}_w^3) \otimes \Omega^\cdot(\mathbb{R}) \\
& \frac{1}{2} \int_{\mathbb{C}^2 \times \mathbb{C}^3 \times \mathbb{R}} \frac{1}{1-v} (\partial \gamma \vee \mu^2) \wedge (\mathrm{d}^2 z \wedge \mathrm{d}^3 w) \\
& \quad \frac{1}{6} \int_{\mathbb{C}^2 \times \mathbb{C}^3 \times \mathbb{R}} \gamma \partial \gamma \partial \gamma \\
& \int \frac{1}{1-v} \left(\frac{1}{2} \partial^w \gamma_w \vee \mu_w^2 + \partial^z \gamma_w \vee \mu_w \mu_z + \partial^w \gamma_z \vee \mu_w \mu_z + \frac{1}{2} \partial^z \gamma_z \vee \mu_z^2 \right) \wedge (\mathrm{d}^2 z \wedge \mathrm{d}^3 w) \\
& \quad + \frac{1}{2} \int \frac{1}{1-v} (\mathrm{d}^2 z \vee \mu_z^2) \wedge (\mathrm{d}^2 z \wedge \mathrm{d}^3 w) \\
& \frac{1}{6} \int (\gamma_w \partial^z \gamma_w \partial^z \gamma_w + \gamma_w \partial^w \gamma_w \partial^z \gamma_z + \gamma_w \partial^w \gamma_z \partial^w \gamma_z) + \frac{1}{2} \int (\gamma_w \partial^w \gamma_w) \wedge \mathrm{d}^2 z \\
& \quad \frac{1}{2} \int (\mathrm{d}^2 z \vee \mu_z^2) \wedge (\mathrm{d}^2 z \wedge \mathrm{d}^3 w) + \frac{1}{2} \int (\gamma_w \wedge \partial^w \gamma_w) \wedge \mathrm{d}^2 z \\
& \quad \Omega_Z^{1,\cdot} \widehat{\otimes} \Omega_W^{0,\cdot} \xrightarrow{\omega_Z^{2,0} \otimes 1} \mathrm{PV}_Z^{1,\cdot} \widehat{\otimes} \mathrm{PV}_W^{0,\cdot} \\
& \quad \Omega_Z^{0,\cdot} \widehat{\otimes} \Omega_W^{1,\cdot} \xrightarrow{1 \otimes \partial^w} \Omega_Z^{0,\cdot} \widehat{\otimes} \Omega_W^{2,\cdot} \xrightarrow{1 \otimes \Omega_W} \mathrm{PV}_Z^{0,\cdot} \widehat{\otimes} \mathrm{PV}_W^{1,\cdot} \\
& \quad \Omega_Z^{0,\cdot} \widehat{\otimes} \Omega_W^{0,\cdot} \hat{\otimes} \Omega_L = \bigoplus_{k=0}^3 \Omega_Z^{0,\cdot} \hat{\otimes} \Omega_W^{k,\cdot} \hat{\otimes} \Omega_L \\
& \mu_z = (1 - \tilde{\alpha}^3) (\partial_{z_1} \wedge \partial_{z_2}) \vee \partial^z \alpha^0, \mu_w = (\partial_{w_1} \wedge \partial_{w_2} \wedge \partial_{w_3}) \vee \alpha^2, v = \tilde{\alpha}^3 \\
& \quad \beta = \alpha^0, \gamma_w = \alpha^1, \gamma_z = 0 \\
& \quad \int \sum_{k=0}^3 \alpha^k (\bar{\partial} + \mathrm{d}_{\mathbb{R}}) \alpha^{3-k} \\
& \quad \int \alpha^0 \partial^w \alpha^2 - \int \alpha^0 \partial^z \alpha^0 \partial^z \alpha^3 \\
& \quad \int \frac{1}{2} \alpha^1 \partial^w \alpha^1 \\
& \frac{1}{2} \int \frac{1}{1 - \tilde{\alpha}^3} \partial^w \alpha^1 (\tilde{\alpha}^2)^2 \mathrm{d}^2 z + \int \alpha^2 \partial^z \alpha^0 \partial^z \alpha^1 + \frac{1}{2} \int (1 - \alpha^3) \partial^z \alpha^0 \partial^z \alpha^0 \\
& \quad \frac{1}{6} \int \alpha^1 \partial^z \alpha^1 \partial^z \alpha^1 \\
& \quad S_{pCS}(\alpha) + \int \frac{1}{2} \frac{1}{1 - \tilde{\alpha}^3} \partial^w \alpha^1 (\tilde{\alpha}^2)^2 \mathrm{d}^2 z \\
& \quad \frac{1}{6} \int \frac{1}{1 - \tilde{\alpha}^3} \alpha^2 (\tilde{\alpha}^2)^2 \\
& \frac{1}{2} \int \frac{1}{1 - \tilde{\alpha}^3} \partial^w \alpha^1 (\tilde{\alpha}^2)^2 + \frac{1}{6} \int \frac{1}{1 - \tilde{\alpha}^3} \partial^w (\alpha^2) \alpha^2 (\tilde{\alpha}^2)^2 \\
& \quad \mathrm{PV}^{i,j}(X) = \Omega^{0,j}(X, \wedge^i \mathrm{T}_X) \\
& \quad \partial_\Omega: \mathrm{PV}^{i,\cdot}(X) \rightarrow \mathrm{PV}^{i-1,\cdot}(X) \\
& \quad (\mathrm{PV}^\cdot(X)[[u]][2], \bar{\partial} + u \partial_\Omega) \\
& \quad (\partial_\Omega \otimes 1) \delta_{\Delta \subset X \times X} \in [\mathrm{PV}^\cdot(X)]^{\hat{\otimes} 2} \\
& I_{BCOV}(\Sigma) = \mathrm{Tr}_X \langle \exp \Sigma \rangle_0 = \sum_{n \geq 0} \mathrm{Tr}_X \langle \Sigma^{\otimes n} \rangle_0 \\
& \langle u^{k_1} \mu_1 \otimes \cdots \otimes u^{k_m} \mu_m \rangle_0 = \left(\int_{\overline{\mathcal{M}}_{0,m}} \psi_1^{k_1} \cdots \psi_m^{k_m} \right) \mu_1 \cdots \mu_m = \binom{m-3}{k_1, \dots, k_m} \mu_1 \cdots \mu_m \\
& \Sigma \mapsto [u(\exp(\Sigma/u) - 1)]_+
\end{aligned}$$

$$\begin{aligned}
& \left(\bigoplus_{i+j \leq d-1} u^i \text{PV}^{j,\cdot}(X)[2], \bar{\partial} + u\partial_\Omega \right) \\
& \text{PV}^{1,\cdot} \xrightarrow{u\partial_\Omega} u\text{PV}^{0,\cdot} \\
& \text{PV}^{2,\cdot} \xrightarrow{u\partial_\Omega} u\text{PV}^{1,\cdot} \xrightarrow{u\partial_\Omega} u^2\text{PV}^{0,\cdot} \\
& \text{PV}^{3,\cdot} \xrightarrow{u\partial_\Omega} u\text{PV}^{2,\cdot} \xrightarrow{u\partial_\Omega} u^2\text{PV}^{1,\cdot} \xrightarrow{u\partial_\Omega} u^3\text{PV}^{0,\cdot} \\
& \alpha = \sum_n \alpha_n u^n \in \mathcal{E}_{mKS}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2) \\
& I_{IIA} = \int_{\mathbb{C}^4 \times \mathbb{R}^2} \alpha_0^3 + \dots \\
& \eta \in \text{PV}^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2), \mu + uv \in \text{PV}^{1,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2) \oplus u\text{PV}^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2) \\
& \Pi \in \text{PV}^{3,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2), \sigma \in \text{PV}^{3,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2) \\
& I_{IIA} = \frac{1}{2} \text{Tr}_{\mathbb{C}^4 \times \mathbb{R}^2} \frac{1}{1-v} \mu^2 \wedge \Pi + \text{Tr}_{\mathbb{C}^4 \times \mathbb{R}^2} \frac{1}{1-v} \eta \wedge \mu \wedge \sigma + \frac{1}{2} \text{Tr}_{\mathbb{C}^4 \times \mathbb{R}^2} \frac{1}{1-v} \eta \wedge \Pi^2 + \dots \\
& \text{PV}^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2)_\eta \\
& u^{-1}\Omega^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2)_\beta \xrightarrow{u\partial} \Omega^{1,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2)_\gamma \\
& \Omega^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^\cdot(\mathbb{R}^2)_\theta \\
& \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \eta \theta + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \mu \vee \gamma + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega v \beta \\
& \tilde{I}_{IIA} = \frac{1}{2} \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \frac{1}{1-v} \mu^2 \vee \partial \gamma + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \frac{1}{1-v} (\eta \wedge \mu) \vee \partial \theta + \frac{1}{2} \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \frac{1}{1-v} \eta \wedge \partial \gamma \wedge \partial \gamma \\
& M \times V \rightarrow M \times V_{\mathbb{R}} \\
& \mathbb{C}^4 \times \mathbb{C} \times \mathbb{R}_t \rightarrow \mathbb{C}^4 \times \mathbb{R}_x \times \mathbb{R}_t \cong \mathbb{C}^4 \times \mathbb{R}^2 \\
& v_{11d} \in \text{PV}^{0,\cdot}(\mathbb{C}^5) \otimes \Omega^\cdot(\mathbb{R}) \\
& v(z_i, x, t) = v_{11d}(z_i, x, y=0, t)|_{d\bar{z}_5=dx} \\
& \beta(z_i, x, t) = \beta_{11d}(z_i, x, y=0, t)|_{d\bar{z}_5=dx} \\
& \mu_{11d} = \mu_{11d}^0 + \theta_{11d} \partial_{z_5} \\
& \mu_{11d}^0 \in \text{PV}^{1,\cdot}(\mathbb{C}^4) \otimes \Omega^{0,\cdot}(\mathbb{C}_{z_5}) \otimes \Omega^\cdot(\mathbb{R}_t) \\
& \theta_{11d} \in \Omega^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^{0,\cdot}(\mathbb{C}_{z_5}) \otimes \Omega^\cdot(\mathbb{R}_t) \\
& \mu(z_i, x, t) = \mu_{11d}^0(z_i, x, y=0, t)|_{d\bar{z}_5=dx} \\
& \theta(z_i, x, t) = \theta_{11d}(z_i, x, y=0, t)|_{d\bar{z}_5=dx} \\
& \gamma_{11d} = \gamma_{11d}^0 + \eta_{11d} dz_5 \\
& \gamma_{11d}^0 \in \Omega^{1,\cdot}(\mathbb{C}^4) \otimes \Omega^{0,\cdot}(\mathbb{C}_{z_5}) \otimes \Omega^\cdot(\mathbb{R}_t) \\
& \eta_{11d} \in \text{PV}^{0,\cdot}(\mathbb{C}^4) \otimes \Omega^{0,\cdot}(\mathbb{C}_{z_5}) \otimes \Omega^\cdot(\mathbb{R}_t) \\
& \gamma(z_i, x, t) = \gamma_{11d}^0(z_i, x, y=0, t)|_{d\bar{z}_5=dx} \\
& \eta(z_i, x, t) = \eta_{11d}(z_i, x, y=0, t)|_{d\bar{z}_5=dx} \\
& \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \frac{1}{1-v} \mu^2 \vee \partial \gamma + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^\Omega \frac{1}{1-v} (\theta \wedge \mu) \vee \partial \eta \\
& \int_{\mathbb{C}^4 \times \mathbb{R}^2} \eta \wedge \partial \gamma \wedge \partial \gamma \\
& \tilde{\theta} = \frac{1}{1-v} \theta, \tilde{\eta} = (1-v)\eta, \tilde{\beta} = \beta + \frac{1}{1-v} \eta \wedge \theta
\end{aligned}$$

$$\begin{array}{c} \int_{\mathbb{C}^4 \times \mathbb{R}^2}^{\Omega_{\mathbb{C}^4}} \frac{1}{1-\nu} \mu^2 \vee \partial \gamma + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^{\Omega_{\mathbb{C}^4}} \frac{1}{1-\nu} \tilde{\eta} \wedge \partial \gamma \wedge \partial \gamma + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^{\Omega_{\mathbb{C}^4}} (\tilde{\theta} \wedge \mu) \vee \partial \left(\frac{1}{1-\nu} \tilde{\eta} \right) \\ + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^{\Omega_{\mathbb{C}^4}} \frac{1}{1-\nu} (\tilde{\eta} \wedge \tilde{\theta}) \partial_{\Omega} \mu \\ - \int_{\mathbb{C}^4 \times \mathbb{R}^2}^{\Omega_{\mathbb{C}^4}} \left(\frac{1}{1-\nu} \tilde{\eta} \right) \partial_{\Omega} (\tilde{\theta} \mu) + \int_{\mathbb{C}^4 \times \mathbb{R}^2}^{\Omega_{\mathbb{C}^4}} \left(\frac{1}{1-\nu} \tilde{\eta} \right) \tilde{\theta} \partial_{\Omega} \mu \\ \hline - \qquad \qquad \qquad + \qquad \qquad \qquad - \qquad \qquad \qquad + \\ \text{PV}^{1,\bullet}(\mathbb{C}^5) \xrightarrow{u\hat{c}_{\Omega}} u\text{PV}^{0,\bullet}(\mathbb{C}^5) \end{array}.$$

$$\text{PV}^{3,\bullet}(\mathbb{C}^5) \xrightarrow{u\hat{c}_{\Omega}} u\text{PV}^{2,\bullet}(\mathbb{C}^5) \xrightarrow{u\hat{c}_{\Omega}} u^2\text{PV}^{1,\bullet}(\mathbb{C}^5) \xrightarrow{u\hat{c}_{\Omega}} u^3\text{PV}^{0,\bullet}(\mathbb{C}^5)$$

$$\begin{array}{c} I_{\text{typeI}} = \text{Tr}_{\mathbb{C}^5} \frac{1}{1-\nu} \mu^2 \vee \sigma + \dots \\ \hline - \qquad \qquad \qquad + \qquad \qquad \qquad - \\ \text{PV}^{1,\bullet}(\mathbb{C}^5)_{\mu} \xrightarrow{\hat{c}_{\Omega}} \text{PV}^{0,\bullet}(\mathbb{C}^5)_{\nu} \\ \Omega^{0,\bullet}(\mathbb{C}^5)_{\tilde{\beta}} \xrightarrow{\partial} \Omega^{1,\bullet}(\mathbb{C}^5)_{\tilde{\gamma}}. \end{array}$$

$$\begin{array}{c} \tilde{I}_{\text{typeI}} = \frac{1}{2} \int_{\mathbb{C}^5}^{\Omega} \frac{1}{1-\nu} \mu^2 \vee \partial \tilde{\gamma} \\ \hline - \qquad \qquad \qquad + \\ \text{PV}^{1,\bullet}(\mathbb{C}^5)_{\mu} \xrightarrow{\hat{c}_{\Omega}} \text{PV}^{0,\bullet}(\mathbb{C}^5)_{\nu} \\ \Omega^{0,\bullet}(\mathbb{C}^5)_{\beta} \xrightarrow{\partial} \Omega^{1,\bullet}(\mathbb{C}^5)_{\gamma}. \end{array}$$

$$\begin{array}{c} \mathcal{M}_{t=0}, \mathcal{M}_{t=1} \subset \mathcal{E}_{\partial} \\ \mathcal{M}_{t=0} \stackrel{\cong}{\times} \mathcal{E}_{\partial} \mathcal{M}_{t=1}. \\ \mathcal{M}_{t=0}: \gamma|_{t=0} = \beta|_{t=0} = 0 \\ \mathcal{M}_{t=1}: \gamma|_{t=1} = \beta|_{t=1} = 0 \\ \mathbb{C}^5 \times [0,1] \rightarrow \mathbb{C}^5 \end{array}$$

$$\begin{array}{c} - \qquad \qquad \qquad + \\ \text{PV}^{1,\bullet}(\mathbb{C}^5)_{\mu} \xrightarrow{\hat{c}_{\Omega}} \text{PV}^{0,\bullet}(\mathbb{C}^5)_{\nu}. \\ \hline - \qquad \qquad \qquad + \qquad \qquad \qquad - \\ \text{PV}^{1,\bullet}(\mathbb{C}^5)_{\mu} \xrightarrow{\hat{c}_{\Omega}} \text{PV}^{0,\bullet}(\mathbb{C}^5)_{\nu} \\ \Omega^{0,\bullet}(\mathbb{C}^5)_{\beta} \xrightarrow{\partial} \Omega^{1,\bullet}(\mathbb{C}^5)_{\gamma} \xrightarrow{\mathbf{1}} \Omega^{0,\bullet}(\mathbb{C}^5)_{\tilde{\beta}} \xrightarrow{\partial} \Omega^{1,\bullet}(\mathbb{C}^5)_{\tilde{\gamma}}. \end{array}$$

$$\begin{array}{c} \mathcal{M}_{t=0} \hookrightarrow \tilde{\mathcal{M}}_{t=0} \rightarrow \mathcal{E}_{\partial} \\ \hline - \qquad \qquad \qquad + \qquad \qquad \qquad - \\ \text{PV}^{1,\bullet}(\mathbb{C}^5)_{\mu} \xrightarrow{\hat{c}_{\Omega}} \text{PV}^{0,\bullet}(\mathbb{C}^5)_{\nu} \\ \Omega^{0,\bullet}(\mathbb{C}^5)_{\tilde{\beta}} \xrightarrow{\partial} \Omega^{1,\bullet}(\mathbb{C}^5)_{\tilde{\gamma}} \end{array}$$

$$\begin{array}{c} \mathbb{C}^5 \times S^1 \rightarrow \mathbb{C}^5 \\ \mu + \epsilon \mu' \in \text{PPV}^{1,\bullet}(\mathbb{C}^5)[\epsilon] \end{array}$$



$$\mathrm{PV}^{1,\bullet}(\mathbb{C}^5)_\mu \xrightarrow{\partial} \mathrm{PV}^{0,\bullet}(\mathbb{C}^5)_\nu$$

$$\epsilon\Omega^{0,\bullet}(\mathbb{C}^5)_{\beta'} \xrightarrow{\partial_\Omega} \epsilon\Omega^{1,\bullet}(\mathbb{C}^5)_{\gamma'}.$$

$$\epsilon\mathrm{PV}^{1,\bullet}(\mathbb{C}^5)_{\mu'} \xrightarrow{\partial_\Omega} \epsilon\mathrm{PV}^{0,\bullet}(\mathbb{C}^5)_{\nu'}$$

$$\Omega^{0,\bullet}(\mathbb{C}^5)_\beta \xrightarrow{\partial_\Omega} \Omega^{1,\bullet}(\mathbb{C}^5)_\gamma.$$

$$\int_{\mathbb{C}^5}^\Omega (\beta' \wedge \bar{\partial} \nu + \beta \wedge \bar{\partial} \nu' + \gamma' \wedge \bar{\partial} \mu + \gamma \wedge \bar{\partial} \mu' + \beta' \wedge \partial_\Omega \mu + \beta \wedge \partial_\Omega \mu') \\ + \int_{\mathbb{C}^5}^\Omega \left(\frac{1}{2} \frac{1}{1-\nu} \mu^2 \vee \partial \gamma' + \frac{1}{1-\nu} (\mu \wedge \mu') \vee \partial \gamma' + \frac{1}{2} \frac{\nu'}{(1-\nu)^2} \mu^2 \vee \partial \gamma \right)$$

$$+ \frac{1}{2} \int_{\mathbb{C}^5} \gamma' \wedge \partial \gamma \wedge \partial \gamma$$

$$X \times \mathbb{C}^2 \times \mathbb{R} \rightarrow \mathbb{C}^2 \times \mathbb{R}$$

$$\alpha, \eta \in \Pi\Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$A_{\mathrm{grav}}, B_{\mathrm{grav}} \in \Pi\Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\int_{\mathbb{C}^2 \times \mathbb{R}}^\Omega (\eta \bar{\partial} \alpha + B_{\mathrm{grav}} \bar{\partial} A_{\mathrm{grav}} + B \bar{\partial} A + \psi \bar{\partial} \chi)$$

$$+ \int_{\mathbb{C}^2 \times \mathbb{R}}^\Omega \left(\frac{1}{2} \eta \{ \alpha, \alpha \} + B_{\mathrm{grav}} \{ \alpha, A_{\mathrm{grav}} \} + B \{ \alpha, A \} + \psi \{ \alpha, \chi \} \right)$$

$$+ \frac{1}{6} \int_{\mathbb{C}^2 \times \mathbb{R}} B_{\mathrm{grav}} \partial B_{\mathrm{grav}} \partial B_{\mathrm{grav}}$$

$$\mathrm{PV}^{0,\cdot}(X \times \mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}) \simeq H^{\cdot}(X, \mathcal{O}) \otimes \mathrm{PV}^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$= \mathrm{PV}^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}) \oplus \Pi \bar{\Omega}_X \mathrm{PV}^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$v_{11d} = v + \bar{\Omega}_X \tilde{v}$$

$$\Pi\mathrm{PV}^{1,\cdot}(X \times \mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}) \simeq \Pi H^{\cdot}(X, \mathcal{O}) \otimes \mathrm{PV}^{1,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\oplus \Pi H^{\cdot}(X, T_X) \otimes \mathrm{PV}^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$= \Pi\mathrm{PV}^{1,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}) \oplus \bar{\Omega}_X \mathrm{PV}^{1,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\oplus H^1(X, T_X) \otimes \mathrm{PV}^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}) \oplus \Pi H^2(X, T_X) \otimes \mathrm{PV}^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\mu_{11d} = \mu + \bar{\Omega}_X \tilde{\mu}$$

$$+ e^i \chi_i + f^a A_a + (\Omega_X^{-1} \vee \omega^2) A_{\mathrm{grav}}$$

$$H^2(X, \Omega_X^2)_\perp \subset H^2(X, \Omega_X^2) \cong H^2(X, T_X)$$

$$\beta_{11d} = \beta + \bar{\Omega}_X \tilde{\beta}$$

$$\gamma_{11d} = \gamma + \bar{\Omega}_X \tilde{\gamma} + e_i \psi^i + f_a B^a + \omega \wedge B_{\mathrm{grav}}$$

$$\alpha, \chi \in \Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\eta = (\mathrm{d}^2 z)^{-1} \vee \partial \tilde{\gamma}, \psi = (\mathrm{d}^2 z)^{-1} \vee \partial \gamma \in \Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\alpha, A_{\mathrm{grav}} \in \Pi\Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}), \eta, B_{\mathrm{grav}}$$

$$\in \Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R})$$

$$\chi, \chi_i \in \Omega^0 \cdot (\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}), \quad \psi, \psi^i \in \Omega^0(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}), \quad i = 1, \dots, h^{2,1}$$

$$A_a \in \Pi\Omega^{0,\cdot}(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}), \quad B^a \in \Pi\Omega^0(\mathbb{C}^2) \otimes \Omega^{\cdot}(\mathbb{R}), a = 1, \dots, h^{1,1} - 1.$$

$$\int_{\mathbb{C}^2 \times \mathbb{R}}^\Omega \left(\frac{1}{2} \partial \alpha \wedge \partial \alpha \wedge \eta + \partial A_{\mathrm{grav}} \wedge \partial A_{\mathrm{grav}} \wedge B_{\mathrm{grav}} \right)$$

$$+ \int_{\mathbb{C}^2 \times \mathbb{R}}^\Omega (\partial \alpha \wedge \partial \chi \wedge \psi + \partial \alpha \wedge \partial \chi_i \wedge \psi^i + \partial \alpha \wedge \partial A_a \wedge B^a)$$

$$\frac{1}{6} \int_{\mathbb{C}^2 \times \mathbb{R}} B_{\mathrm{grav}} \partial B_{\mathrm{grav}} \partial B_{\mathrm{grav}}$$

$$H^{\cdot}(X, \Omega^{\cdot}) = \oplus_{i,j} H^i(X, \Omega_X^j)$$



$$\begin{aligned}
& H^*(X, \Omega^*) \otimes \mathcal{O}(\mathbb{C}^2) \\
& [[\omega] \otimes f, [\omega] \otimes g] = [\omega^2] \otimes \{f, g\} \in H^{2,2}(X, \Omega^2) \otimes \mathcal{O}(\mathbb{C}^2) \\
& \mathbb{C}_w^4 \times \mathbb{C}_z \times \mathbb{R} \\
& I_{M2}(\gamma) = N \int_{\mathbb{C}_z} \gamma + \dots \\
& I_{D2}(\gamma) = N \int_{\mathbb{R} \times \mathbb{C}_z} \gamma + \dots \\
& \bar{\partial} \mu + \frac{1}{2} [\mu, \mu] + \partial \gamma \partial \gamma = N \Omega^{-1} \delta_{w=0} \\
& \partial_\Omega \mu = 0 \\
& F_{M2} = \frac{6}{(2\pi i)^4} \frac{\sum_{a=1}^4 \bar{w}_a \, d\bar{w}_1 \cdots \widehat{d}_a \cdots d\bar{w}_4}{\|w\|^8} \partial_z \\
& \bar{\partial}(NF_{M2}) + \frac{1}{2} [NF_{M2}, NF_{M2}] = N \Omega^{-1} \delta_{w=0} \\
& \partial_\Omega(NF_{M2}) = 0 \\
& [F_{M2}, F_{M2}] = 0 \\
& \frac{\partial}{\partial z}, z \frac{\partial}{\partial z} - \frac{1}{4} \sum_{a=1}^4 w_a \frac{\partial}{\partial w_a}, z \left(z \frac{\partial}{\partial z} - \frac{1}{2} \sum_{a=1}^4 w_a \frac{\partial}{\partial w_a} \right) \in \text{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}) \\
& \sum_{a,b=1}^4 B_{ab} w_a \frac{\partial}{\partial w_b} \in \text{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), (B_{ab}) \in \mathfrak{sl}(4) \\
& (\wedge^2 \mathbb{C}^4)_{+1} \oplus (\wedge^2 \mathbb{C}^4)_{-1} \\
& \frac{1}{2} (w_a \, dw_b - w_b \, dw_a) \in \Omega^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), a, b = 1, 2, 3, 4 \\
& \frac{1}{2} z (w_a \, dw_b - w_b \, dw_a) \in \Omega^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), a, b = 1, 2, 3, 4 \\
& i_{M2}: \mathfrak{osp}(6|1) \hookrightarrow E(5,10) \\
& dw_a \wedge dw_b, a, b = 1, 2, 3, 4 \\
& z \, dw_a \wedge dw_b + \frac{1}{2} \, dz \wedge (w_a \, dw_b - w_b \, dw_a), a, b = 1, 2, 3, 4 \\
& (\mathbb{C}^5 \times \mathbb{R}) \setminus \{w = 0\} \cong (\mathbb{C}_w^4 \setminus 0) \times \mathbb{C}_z \times \mathbb{R} \\
& [F, -]: \text{PV}^{i,\cdot}(\mathbb{C}_w^4 \setminus 0) \otimes \text{PV}^{j,\cdot}(\mathbb{C}_z) \otimes \Omega^*(\mathbb{R}) \rightarrow \text{PV}^{i,+3}(\mathbb{C}_w^4 \setminus 0) \otimes \text{PV}^{j,\cdot}(\mathbb{C}_z) \otimes \Omega^*(\mathbb{R}) \\
& [F, -]: \Omega^{i,\cdot}(\mathbb{C}_w^4 \setminus 0) \otimes \Omega^{j,\cdot}(\mathbb{C}_z) \otimes \Omega^*(\mathbb{R}) \rightarrow \Omega^{i,+3}(\mathbb{C}_w^4 \setminus 0) \otimes \Omega^{j,\cdot}(\mathbb{C}_z) \otimes \Omega^*(\mathbb{R}) \\
& \delta^{(1)} = \bar{\partial} + d_{\mathbb{R}} + \partial_\Omega|_{\mu \rightarrow \nu} + \partial|_{\beta \rightarrow \gamma} \\
& + \qquad \qquad \qquad - \\
& \hline
& H^\bullet(\mathbb{C}^4 \setminus 0, T) \otimes H^\bullet(\mathbb{C}, \mathcal{O}) \qquad H^\bullet(\mathbb{C}^4 \setminus 0, \mathcal{O}) \otimes H^\bullet(\mathbb{C}, \mathcal{O}) \\
& H^\bullet(\mathbb{C}^4 \setminus 0, \mathcal{O}) \otimes H^\bullet(\mathbb{C}, T) \\
& H^\bullet(\mathbb{C}^4 \setminus 0, \mathcal{O}) \otimes H^\bullet(\mathbb{C}, \mathcal{O}) \qquad H^\bullet(\mathbb{C}^4 \setminus 0, \mathcal{O}) \otimes H^\bullet(\mathbb{C}, \Omega^1) \\
& \qquad \qquad \qquad H^\bullet(\mathbb{C}^4 \setminus 0, \Omega^1) \otimes H^\bullet(\mathbb{C}, \mathcal{O}) \\
& \mathbb{C}[z] \hookrightarrow H^*(\mathbb{C}, \mathcal{F}) \\
& \mathbb{C}[w_1, \dots, w_4] \hookrightarrow H^0(\mathbb{C}^4 \setminus 0, \mathcal{O}) \\
& \mathbb{C}[w_1, \dots, w_4] \{\partial_{w_i}\} \hookrightarrow H^0(\mathbb{C}^4 \setminus 0, T) \\
& \mathbb{C}[w_1, \dots, w_4] \{dw_i\} \hookrightarrow H^0(\mathbb{C}^4 \setminus 0, \Omega^1)
\end{aligned}$$



$$\begin{aligned}
& (w_1 \cdots w_4)^{-1} \mathbb{C}[w_1^{-1}, \dots, w_4^{-1}] \hookrightarrow H^3(\mathbb{C}^4 \setminus 0, \mathcal{O}) \\
& (w_1 \cdots w_4)^{-1} \mathbb{C}[w_1^{-1}, \dots, w_4^{-1}] \{\partial_{w_i}\} \hookrightarrow H^3(\mathbb{C}^4 \setminus 0, \mathcal{T}) \\
& (w_1 \cdots w_4)^{-1} \mathbb{C}[w_1^{-1}, \dots, w_4^{-1}] \{dw_i\} \hookrightarrow H^3(\mathbb{C}^4 \setminus 0, \Omega^1) \\
& H(\mathcal{L}(\mathbb{C}^5 \times \mathbb{R} \setminus \{w = 0\}), \bar{\partial})
\end{aligned}$$

$$\begin{array}{ccc}
- & & + \\
\hline
H^3(\mathbb{C}^4 \setminus 0, \mathcal{O})[z] \{\partial_{w_i}\} & \xrightarrow{\hat{\partial}_\Omega} & H^3(\mathbb{C}^4 \setminus 0, \mathcal{O})[z] \\
H^3(\mathbb{C}^4 \setminus 0, \mathcal{O})[z] \partial_z & \xrightarrow{\hat{\partial}_\Omega} & \\
H^3(\mathbb{C}^4 \setminus 0, \mathcal{O})[z] & \xrightarrow{\partial} & H^3(\mathbb{C}^4 \setminus 0, \mathcal{O})[z] dz \\
& \xrightarrow{\partial} & H^3(\mathbb{C}^4 \setminus 0, \Omega^1)[z] \{dw_i\}.
\end{array}$$

$$\begin{aligned}
[F] &= (w_1 \cdots w_4)^{-1} \partial_z \in H^3(\mathbb{C}^4 \setminus 0, \mathcal{O})[z] \partial_z \\
[[F], z(w_a dw_b - w_b dw_a)] &= (w_1 \cdots w_4)^{-1} (w_a dw_b - w_b dw_a) = 0 \\
\{w_1 = w_2 = t = 0\} &\subset \mathbb{C}_z^3 \times \mathbb{C}_w^2 \times \mathbb{R}
\end{aligned}$$

$$I_{M5} = N \int_{\mathbb{C}_z^3} \partial_\Omega^{-1} \mu \vee \Omega + \cdots$$

$$N \int_{\mathbb{C}^2 \times \mathbb{R}} \partial_\Omega^{-1} \mu \vee \Omega_{\mathbb{C}^4} + \cdots$$

$$\begin{aligned}
\bar{\partial} \partial \gamma + \partial_\Omega \left(\frac{1}{1-\nu} \mu \right) \wedge \partial \gamma &= N \delta_{w_1=w_2=t=0} \\
(\bar{\partial} + d_{\mathbb{R}}) \mu + \partial \gamma \partial \gamma &= 0
\end{aligned}$$

$$F_{M5} = \frac{1}{(2\pi i)^3} \frac{\bar{w}_1 d\bar{w}_2 \wedge dt - \bar{w}_2 d\bar{w}_1 \wedge dt + t d\bar{w}_1 \wedge d\bar{w}_2}{(\|w\|^2 + t^2)^{5/2}} \wedge dw_1 \wedge dw_2$$

$$\bar{\partial}(NF_{M5}) + d_{\mathbb{R}}(NF_{M5}) = N \delta_{w_1=w_2=t=0}$$

$$(NF_{M5}) \wedge (NF_{M5}) = 0.$$

$$\bar{\partial} F + d_{\mathbb{R}} F = N \delta_{w_1=w_2=t=0}$$

$$(\mathbb{C}^2 \times \mathbb{R}) \setminus 0 \simeq S^4 \times \mathbb{R}$$

$$\oint NF = N$$

$$\begin{aligned}
& \overset{S^4}{\mathbb{C}^5 \times \mathbb{R}} = \mathbb{C}_z^3 \times \mathbb{C}_w^2 \times \mathbb{R}_t \\
& \mathfrak{sl}(4) \oplus \mathfrak{sl}(2)
\end{aligned}$$

$$\frac{\partial}{\partial z_i} \in \mathrm{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), i = 1, 2, 3$$

$$A_{ij} z_i \frac{\partial}{\partial z_j} \in \mathrm{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), (A_{ij}) \in \mathfrak{sl}(3)$$

$$\sum_{i=1}^3 z_i \frac{\partial}{\partial z_i} - \frac{3}{2} \sum_{a=1}^2 w_a \frac{\partial}{\partial w_a} \in \mathrm{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R})$$

$$z_j \left(\sum_{i=1}^3 z_i \frac{\partial}{\partial z_i} - 2 \sum_{a=1}^2 w_a \frac{\partial}{\partial w_a} \right) \in \mathrm{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R})$$

$$w_1 \frac{\partial}{\partial w_2}, w_2 \frac{\partial}{\partial w_1}, \frac{1}{2} \left(w_1 \frac{\partial}{\partial w_1} - w_2 \frac{\partial}{\partial w_2} \right) \in \mathrm{PV}^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R})$$

$$L \otimes R \oplus \wedge^2 L \otimes R \cong \mathbb{C}^3 \otimes \mathbb{C}^2 \oplus \wedge^2 \mathbb{C}^3 \otimes \mathbb{C}$$

$$z_i dw_a \in \Omega^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), a = 1, 2, i = 1, 2, 3$$



$$\begin{aligned}
& \frac{1}{2} w_a (z_i \, dz_j - z_j \, dz_i) \in \Omega^{1,0}(\mathbb{C}^5) \otimes \Omega^0(\mathbb{R}), a = 1,2,k = 1,2,3 \\
& i_{M5} \colon \mathfrak{osp}(6 \mid 1) \hookrightarrow E. \\
& dz_i \wedge dw_a, i = 1,2,3, a = 1,2. \\
& w_a \, dz_i \wedge \, dz_j + \frac{1}{2} \, dw_a \wedge (z_i \, dz_j - z_j \, dz_i), i,j = 1,2,3, a = 1,2 \\
& \Omega^{0,\cdot}(\mathbb{C}^5) \otimes \Omega^{\cdot}(\mathbb{R}) \\
& \Omega^{0,\cdot}(\mathbb{C}^5) \otimes \Omega^{\cdot}(\mathbb{R}) \cong \Omega^{\cdot}(\mathbb{C}^5 \times \mathbb{R}) / (dz_1, \dots, dz_5) \\
& \mathbb{C}^5 \times \mathbb{R} \setminus \mathbb{C}^3 \cong \mathbb{C}_z^3 \times (\mathbb{C}_w^2 \times \mathbb{R} \setminus 0) \\
& \Omega^{\cdot}(\mathbb{C}^5 \times \mathbb{R} \setminus \mathbb{C}^3) / (dz_1, dz_2, dz_3, dw_1, dw_2) \\
& \Omega^{0,\cdot}(\mathbb{C}_z^3) \otimes (\Omega^{\cdot}(\mathbb{C}_w^2 \times \mathbb{R} \setminus 0) / (dw_1, dw_2)) \\
& \mathbb{C}[w_1, w_2] \hookrightarrow H^0(\Omega^{\cdot}(\mathbb{C}_w^2 \times \mathbb{R} \setminus 0) / (dw_1, dw_2)) \\
& w_1^{-1} w_2^{-1} \mathbb{C}[w_1, w_2] \hookrightarrow H^2(\Omega^{\cdot}(\mathbb{C}_w^2 \times \mathbb{R} \setminus 0) / (dw_1, dw_2)) \\
& \frac{\bar{w}_1 \, d\bar{w}_2 \wedge \, dt - \bar{w}_2 \, d\bar{w}_1 \wedge \, dt + t \, d\bar{w}_1 \wedge \, d\bar{w}_2}{(\|w\|^2 + t^2)^{5/2}} \\
& \delta^{(1)} = \bar{\partial} + d_{\mathbb{R}} + \partial_{\Omega}|_{\mu \rightarrow \nu} + \partial|_{\beta \rightarrow \gamma}. \\
& \mathbb{C}^5 \times \mathbb{R} \setminus \{w_1 = w_2 = t = 0\} \cong \mathbb{C}_z^3 \times (\mathbb{C}_w^2 \times \mathbb{R} \setminus 0) \\
& + \hspace{10em} - \\
& \frac{w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ \partial_{w_i} \} \xrightarrow{\hat{e}_{\Omega}} w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \hspace{1em} w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ \partial_{z_i} \} \xrightarrow{\hat{e}_{\Omega}} w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ dw_i \} }{w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ \partial_{z_i} \} \xrightarrow{\partial} w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ dz_i \} \xrightarrow{\partial} w_1^{-1} w_2^{-1} \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ dw_i \}.} \\
& [[F], f^i(z, w) dz_i] = \epsilon_{ijk} w_1^{-1} w_2^{-1} \partial_{z_j} f^i(z, w) \partial_{z_k} \\
& \mathbb{C}[w_1^{-1}, w_2^{-1}][z_1, z_2, z_3] \{ \partial_{z_i} \} \subset H^0(\mathbb{C}^3, \mathbf{T}) \otimes H^2(\Omega^{\cdot}(\mathbb{C}^2 \times \mathbb{R} \setminus 0) / (dw_1, dw_2)) \\
& [[F], w_a (z_i \, dz_j - z_j \, dz_i)] = 2 \epsilon_{ijk} (w_1^{-1} w_2^{-1}) \cdot w_a \partial_{z_k} = 0 \\
& \partial_t A_i = \Delta A_i + \left[A_j, 2 \partial_j A_i - \partial_i A_j + [A_j, A_i] \right] + (C_A^{\varepsilon} A)_i + \xi_i^{\varepsilon} \\
& \lim_{\varepsilon \downarrow 0} |C_A^{\varepsilon} - \bar{C}_A^{\varepsilon}| = 0 \\
& |\mathbf{E} W_{\ell}[\mathcal{F}_s(A^a)] - \mathbf{E} W_{\ell}[\mathcal{F}_s(A^b)]| \gtrsim t^{\frac{10}{9}} \\
& \chi^{\varepsilon}(t,x) = \varepsilon^{-5} \chi(\varepsilon^{-2} t, \varepsilon^{-1} x) \\
& E = \mathfrak{g}^3 \oplus \mathbf{V} \\
& \mathfrak{g} \ni v \mapsto \mathrm{Ad}_g v \in \mathfrak{g} \text{ and } \mathfrak{g}^3 \ni (v_1, v_2, v_3) \mapsto (\mathrm{Ad}_g v_1, \mathrm{Ad}_g v_2, \mathrm{Ad}_g v_3) \in \mathfrak{g}^3 \\
& C_{\mathrm{YM}}^{\varepsilon} \in L_G(\mathfrak{g}), C_{\mathrm{Higgs}}^{\varepsilon} \in L_G(\mathbf{V}) \\
& C_A^{\varepsilon} = C_{\mathrm{YM}}^{\varepsilon} + \dot{C}_A \in L(\mathfrak{g}^3) \text{ and } C_{\Phi}^{\varepsilon} = C_{\mathrm{Higgs}}^{\varepsilon} + \dot{C}_{\Phi} \in L(\mathbf{V}). \\
& \partial_t A_i = \Delta A_i + \left[A_j, 2 \partial_j A_i - \partial_i A_j + [A_j, A_i] \right] \\
& \hspace{10em} - \mathbf{B}((\partial_i \Phi + A_i \Phi) \otimes \Phi) + (C_A^{\varepsilon} A)_i + \xi_i^{\varepsilon} \\
& \partial_t \Phi = \Delta \Phi + 2 A_j \partial_j \Phi + A_j^2 \Phi - |\Phi|^2 \Phi + C_{\Phi}^{\varepsilon} \Phi + \xi_{\mathrm{H}}^{\varepsilon} \\
& (A(0), \Phi(0)) = (a, \varphi) \in \mathcal{C}^{\infty}, \\
& \langle \mathbf{B}(u \otimes v), h \rangle_{\mathfrak{g}} = \langle u, h v \rangle_{\mathbf{V}}. \\
& \partial_t X = \Delta X + X \partial X + X^3 + C^{\varepsilon} X + \chi^{\varepsilon} * \xi \\
& X = (A, \Phi) \colon [0, T] \times \mathbf{T}^3 \rightarrow E \\
& C = \{C^{\varepsilon}\}_{\varepsilon \in (0,1)} = \{C_A^{\varepsilon}, C_{\Phi}^{\varepsilon}\}_{\varepsilon \in (0,1)} \\
& \mathfrak{G}^{\varrho} \stackrel{\mathrm{def}}{=} \mathcal{C}^{\varrho}(\mathbf{T}^3, G)
\end{aligned}$$

$$\begin{aligned}
g \cdot A &\stackrel{\text{def}}{=} \text{Ad}_g(A) - (\text{d}g)g^{-1}, \text{ and } g \cdot \Phi \stackrel{\text{def}}{=} g\Phi \\
g \cdot \text{SYM}\mathcal{H}(C, (a, \varphi)) &\stackrel{\text{law}}{=} \text{SYM}\mathcal{H}(C, g(0) \cdot (a, \varphi)) \\
g^{-1}(\partial_t g) &= \partial_j(g^{-1}\partial_j g) + [A_j, g^{-1}\partial_j g] \\
\check{C} &\in L_G(\mathfrak{g}) \\
[\text{SYM}\mathcal{H}(C, (a, \varphi))] &\stackrel{\text{law}}{=} [\text{SYM}\mathcal{H}(C, g(0) \cdot (a, \varphi))] \\
\partial_s A_i &= \Delta A_i + [A_j, 2\partial_j A_i - \partial_i A_j + [A_j, A_i]] \\
A(0) &= a \\
\text{d}y_t &= y_t \text{d}\ell_A, y_0 = 1 \\
\ell_A(t) &= \int_0^t \langle A(\ell_s), \dot{\ell}_s \rangle \text{d}s \\
W_\ell(A) &= \text{Trhol}(A, \ell) \\
\mathring{C}_A &= \check{C} + c \in L_G(\mathfrak{g}^3), C^\varepsilon = (C^\varepsilon_{\text{YM}} + \dot{C}_A, C^\varepsilon_\Phi) \in L_G(\mathfrak{g}^3) \oplus L_G(\mathbf{V}) \\
\text{SYM}\mathcal{H}(C, x) &= (A^{(1)}, \Phi^{(1)}), \text{SYM}\mathcal{H}(C, g \cdot x) = (A^{(2)}, \Phi^{(2)}) \\
|\mathbf{E}W_\ell[\mathcal{F}_s(A_t^{(1)})] - \mathbf{E}W_\ell[\mathcal{F}_s(A_t^{(2)})]| &\geq \sigma t^{1+r} \\
|\mathbf{E}W_\ell[\mathcal{F}_s(A_t)] - \mathbf{E}W_\ell[\mathcal{F}_s(\tilde{A}_t)]| &\geq \sigma t^{1+r} \\
\partial_t \tilde{X} &= \Delta \tilde{X} + \tilde{X}\partial \tilde{X} + \tilde{X}^3 + \chi^\varepsilon * \xi + C^\varepsilon \tilde{X} + (c \text{d}\tilde{g}\tilde{g}^{-1}, 0), \tilde{X}(0) = \tilde{x} \\
\partial_t \tilde{g} &= \Delta \tilde{g} - (\partial_j \tilde{g})\tilde{g}^{-1}(\partial_j \tilde{g}) + [\tilde{A}_j, (\partial_j \tilde{g})\tilde{g}^{-1}]\tilde{g} \\
(\partial_t g)g^{-1} &= \partial_j((\partial_j g)g^{-1}) + [B_j, (\partial_j g)g^{-1}] \\
\partial_t Y &= \Delta Y + Y\partial Y + Y^3 + C^\varepsilon Y + (C^\varepsilon \text{d}g g^{-1}, 0) + \text{Ad}_g(\chi^\varepsilon * \xi) \\
\mathbf{E}W_\ell[\mathcal{F}_s(A_t^{(1)})] &= \mathbf{E}W_\ell[\mathcal{F}_s(\tilde{A}_t)] + O(t^M) \\
|\mathbf{E}W_\ell[\mathcal{F}_s(\tilde{A}_t)] - \mathbf{E}W_\ell[\mathcal{F}_s(A_t^{(2)})]| &\geq \sigma t^{1+r} \\
\mathcal{P}_t \star f &= \int_0^t P_{t-s} f_s \text{d}s \\
|f|_{C^\beta} &= \sup_{s \in (0,1)} s^{-\beta/2} |P_s f|_\infty \\
|f|_{C^\beta(\mathbf{R}^3)} &\stackrel{\text{def}}{=} \max_{|k| \leq [\beta]} |\partial^k f|_\infty + \max_{|k| = [\beta]} |\partial^k f|_{C^{\beta - [\beta]}} < \infty, \\
|f|_{C^\eta} &= \sup_{x \neq y} |x - y|^{-\eta} |f(x) - f(y)| \\
O &= [-1, 2] \times \mathbf{T}^3 \\
|\xi|_{C^\beta(O)} &= \sup_{z \in O} \sup_{\varphi \in \mathcal{B}^r} \sup_{\lambda \in (0,1]} \lambda^{-\beta} |\langle \xi, \varphi_z^\lambda \rangle| \\
\varphi_{(s,y)}^\lambda(t,x) &= \lambda^{-5} \varphi((t-s)\lambda^{-2}, (x-y)\lambda^{-1}). \\
\partial_t X &= \Delta X + X\partial X + X^3 + \chi^\varepsilon * \xi + C^\varepsilon X + (ch, 0), \\
\partial_t h_i &= \Delta h_i - [h_j, \partial_j h_i] + [[A_j, h_j], h_i] + \partial_i[A_j, h_j], \\
\omega \in (-1/2, 0), \kappa &\stackrel{\text{def}}{=} \frac{1}{100}(\omega + 1/2) \wedge \frac{1}{100}(-\omega) \in \left(0, \frac{1}{200}\right) \\
|f_1 f_2|_{\mathcal{D}_{\alpha}^{\gamma, \eta}} &\lesssim |f_1|_{\mathcal{D}_1^{\gamma_1, \eta_1}} |f_2|_{\mathcal{D}_2^{\gamma_2, \eta_2}} \\
|\partial f|_{\mathcal{D}_{\alpha-1}^{\gamma-1, \eta-1}} &\lesssim |f|_{\mathcal{D}_{\alpha}^{\gamma, \eta}} \\
\mathcal{X} &= PX(0) + \Psi + \mathcal{P}\mathbf{1}_+ \{ \mathcal{X}\partial \mathcal{X} + \mathcal{X}^3 + \dot{C}\mathcal{X} + c\mathcal{H} \} \\
&\stackrel{\text{def}}{=} PX(0) + \Psi + \mathcal{P}\mathbf{1}_+ \{ Q^{\text{YM}\mathcal{H}}(\mathcal{X}) + c\mathcal{H} \} \\
\mathcal{H} &= Ph(0) + \mathcal{P}\mathbf{1}_+ (\mathcal{H}\partial \mathcal{H} + \mathcal{X}\mathcal{H}^2) + \mathcal{P}'\mathbf{1}_+ (\mathcal{X}\mathcal{H})
\end{aligned}$$



$$\begin{aligned}
\Psi &= \mathcal{P} \mathbf{1}_+ \xi \mathbf{1}_+ \Xi \in \mathcal{D}_{-\frac{1}{2}\kappa}^{\frac{3}{2}+2\kappa, -\frac{1}{2}\kappa} \\
\mathcal{X} &= \mathcal{Y} + \Psi, \\
\mathcal{Y} &= PX(0) + \mathcal{P} \mathbf{1}_+ \{\mathcal{X}^3 + c\mathcal{H} + \check{C}\mathcal{X}\} + \mathcal{P}^{\Psi\partial\Psi}(\Psi\partial\Psi) \\
&\quad + \mathcal{P} \mathbf{1}_+ (\mathcal{Y}\partial\Psi + \Psi\partial\mathcal{Y} + \mathcal{Y}\partial\mathcal{Y}) \\
&= PX(0) + \mathcal{P} \mathbf{1}_+ \{\tilde{Q}^{\text{YMH}}(\mathcal{Y}) + c\mathcal{H}\} + \mathcal{P}^{\Psi\partial\Psi}(\Psi\partial\Psi), \\
\tilde{Q}^{\text{YMH}}(\mathcal{Y}) &= \mathcal{X}^3 + \dot{C}\mathcal{X} + \mathcal{Y}\partial\Psi + \Psi\partial\mathcal{Y} + \mathcal{Y}\partial\mathcal{Y} \\
\tilde{Q}^{\text{YMH}}: \mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega} &\rightarrow \mathcal{D}_{-\frac{3}{2}\kappa}^{\kappa, \omega - \frac{3}{2}\kappa} \\
\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega} \times \mathcal{D}_{-\frac{3}{2}\kappa}^{\frac{1}{2}+2\kappa, -\frac{3}{2}\kappa} &\rightarrow \mathcal{D}_{-\frac{3}{2}\kappa}^{\kappa, \omega - \frac{3}{2}\kappa}. \\
(\mathcal{Y}, \mathcal{H}) &\in \mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega} \times \mathcal{D}_0^{1+2\kappa, 0} \\
|\mathcal{P} \mathbf{1}_+ f|_{\mathcal{D}_{\gamma+2, \bar{\eta}}} &\lesssim t^{\theta/2} |f|_{\mathcal{D}_{\gamma, \eta}}, |\mathcal{P}' \mathbf{1}_+ f|_{\mathcal{D}_{\gamma+1, \bar{\eta}-1}} \lesssim t^{\theta/2} |f|_{\mathcal{D}_{\gamma, \eta}} \\
|\mathcal{P}^w \mathbf{1}_+ f|_{\mathcal{D}_{\gamma+2, \bar{\eta}}} &\lesssim t^{\theta/2} (|f|_{\mathcal{D}_{\gamma, \eta}} + |w|_{\mathcal{C}^{\eta\wedge\alpha}(O)}) \\
\tau^{-1/q} &\lesssim 2 + \|Z\|_{\frac{3}{2}+2\kappa; O} + |X(0)|_{\mathcal{C}^3} + |h(0)|_{\mathcal{C}^3} \\
&\quad + |\mathcal{P} \star \mathbf{1}_+ \xi|_{\mathcal{C}([-1, 3], \mathcal{C}^{-1/2-\kappa})} + |\mathcal{P} \star (\Psi\partial\Psi)|_{\mathcal{C}([-1, 3], \mathcal{C}^{-2\kappa})} \\
|\mathcal{P}' \mathbf{1}_+ (\mathcal{X}\mathcal{H})|_{\mathcal{D}_0^{1+2\kappa, -\kappa}} &\lesssim t^{1/4} |\mathcal{X}\mathcal{H}|_{\mathcal{D}_{-\frac{1}{2}\kappa}^{\frac{1}{2}+\kappa, -\frac{1}{2}\kappa}} \\
&\lesssim t^{1/4} |\mathcal{X}|_{\mathcal{D}_{-\frac{1}{2}\kappa}^{\frac{3}{2}+2\kappa, -\frac{1}{2}\kappa}} |\mathcal{H}|_{\mathcal{D}_0^{1+2\kappa, 0}} \lesssim t^{1/4} \\
|\mathcal{P} \mathbf{1}_+ (\mathcal{H}\partial\mathcal{H})|_{\mathcal{D}_0^{1+2\kappa, 0}} &\lesssim t, |\mathcal{P} \mathbf{1}_+ (\mathcal{X}\mathcal{H}^2)|_{\mathcal{D}_0^{\frac{3}{2}+\kappa, -\kappa}} \lesssim t^{3/4} \\
|\mathcal{P} \mathbf{1}_+ (\tilde{Q}^{\text{YMH}}(\mathcal{Y}))|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa, \omega}} &\lesssim t^{1/4-\kappa/2} |\tilde{Q}^{\text{YMH}}(\mathcal{Y})|_{\mathcal{D}_{-\frac{3}{2}\kappa}^{\kappa, \omega - \frac{3}{2}\kappa}} \lesssim t^{1/4-\kappa/2} \\
|\mathcal{P}^{\Psi\partial\Psi}(\Psi\partial\Psi)|_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega}} &\lesssim t^{-(\kappa+\omega)/2} \\
|\mathcal{P}\mathcal{H}|_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega}}^{\frac{3}{2}} &\lesssim |\mathcal{P}Ph(0)|_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega}} + \left| \mathcal{P}O_{\mathcal{D}_0^{1+2\kappa, -\kappa}}(t^{1/4}) \right|_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega}} \\
&\lesssim t^{1-\frac{\omega}{2}} |\text{Ph}(0)|_{\mathcal{D}_0^{0+, 0}} + t^{1/4+1-\frac{\kappa}{2}-\frac{\omega}{2}} \lesssim t^{1-\omega/2} \\
B_0 &= PX(0) + \Psi, h_0 = Ph(0) \\
B_n &= \mathcal{P}(\dot{C}B_{n-4}\mathbf{1}_{n\geq 4}) + \sum_{\substack{k_1+k_2=n-1 \\ n}} \mathcal{P}(B_{k_1}\partial B_{k_2}) + \sum_{k_1+k_2+k_3=n-2} \mathcal{P}(B_{k_1}B_{k_2}B_{k_3}) \\
\mathcal{X} &= \sum_{i=0}^n B_i + c\mathbf{1}_{n=5}\mathcal{P}h_0 + r_n, \mathcal{H} = h_0 + q_0 \\
r_0 &= O_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa, \omega}}(t^{-\omega/2-\kappa/2}), q_0 = O_{\mathcal{D}_0^{1+2\kappa, -\kappa}}(t^{1/4}) \\
\eta(0) &= -1/2 - \kappa, \quad \eta(n) = -1/2 + 2\kappa \quad (1 \leq n \leq 5) \\
b(0) &= 0, \quad b(n) = (1/4 - \kappa/2)n - 3\kappa/2 \quad (1 \leq n \leq 5) \\
|B_n|_{\mathcal{D}_{\alpha(n)}^{\frac{3}{2}+2\kappa, \eta(n)}} &\lesssim t^{b(n)} \quad \forall 0 \leq n \leq 5
\end{aligned}$$

$$r_0 = O_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa,\omega}}(t^{-\omega/2-\kappa/2}), r_n = O_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}}(t^{(n+1)/4-\kappa_n}) \forall 1 \leq n \leq 5,$$

$$\begin{aligned} \frac{1}{2}\eta(n) + b(n) &= -\frac{1}{4} - \frac{\kappa}{2} + \left(\frac{1}{4} - \frac{\kappa}{2}\right)n \quad \forall n \geq 0 \\ |\mathcal{P}(B_{k_1}\partial B_{k_2})|_{\mathcal{D}_{\alpha(n)}^{\frac{3}{2}+2\kappa,-\frac{1}{2}+2\kappa}} &\lesssim t^{\frac{1}{2}(\eta(k_1)+\eta(k_2)+\frac{3}{2}-2\kappa)} |B_{k_1}\partial B_{k_2}|_{\mathcal{D}_{\alpha(k_1)+\alpha(k_2)-1}^{\kappa,\eta(k_1)+\eta(k_2)-1}} \\ &\lesssim t^{\frac{1}{2}(\eta(k_1)+\eta(k_2)+\frac{3}{2}-2\kappa)} |B_{k_1}|_{\mathcal{D}_{\alpha(k_1)}^{\frac{3}{2}+2\kappa,\eta(k_1)}} |\partial B_{k_2}|_{\mathcal{D}_{\alpha(k_2)-1}^{\frac{1}{2}+2\kappa(k_2)-1}} \\ &\lesssim t^{\frac{1}{2}(\eta(k_1)+\eta(k_2)+\frac{3}{2}-2\kappa)} \cdot t^{b(k_1)} \cdot t^{b(k_2)} = t^{b(n)}. \\ |\mathcal{P}(B_{k_1}B_{k_2}B_{k_3})|_{\mathcal{D}_{\alpha(n)}^{\frac{3}{2}+2\kappa,-\frac{1}{2}+2\kappa}} &\lesssim t^{\frac{1}{2}(\sum_i \eta(k_i)+\frac{5}{2}-2\kappa)} \prod_{i=1}^3 |B_{k_i}|_{\mathcal{D}_{\alpha(k_i)}^{\frac{3}{2}+2\kappa,\eta(k_i)}} \\ &\lesssim t^{\frac{1}{2}(\sum_i \eta(k_i)+\frac{5}{2}-2\kappa)} \cdot t^{\sum_i b(k_i)} = t^{b(n)}. \\ |\mathcal{P}(B_{n-4})|_{\mathcal{D}_{\alpha(n)}^{\frac{3}{2}+2\kappa,-\frac{1}{2}+2\kappa}} &\lesssim t^{\frac{1}{2}(\eta(n-4)+\frac{5}{2}-2\kappa)} |B_{n-4}|_{\mathcal{D}_{\alpha(n-4)}^{\frac{3}{2}+2\kappa,\eta(n-4)}} \\ &\lesssim t^{(\frac{1}{4}-\frac{\kappa}{2})n+\frac{\kappa}{2}} \leq t^{b(n)} \end{aligned}$$

$$\begin{aligned} \sum_{i=0}^n B_i + c\mathbf{1}_{n=5}\mathcal{P}h_0 + r_n &= B_0 + \mathcal{P}\mathbf{1}_+ \{Q^{\text{YMH}}(\mathcal{X}) + c\mathcal{H}\} \\ r_n = \mathcal{P}\mathbf{1}_+ \left\{ Q^{\text{YMH}} \left(\sum_{i=0}^{n-1} B_i + r_{n-1} \right) + c(h_0 + q_0 - \mathbf{1}_{n=5}h_0) \right\} &- \sum_{i=1}^n B_i \\ r_n = \mathcal{P}\mathbf{1}_+ \left(\sum_{k_1+k_2 \geq n} B_{k_1}\partial B_{k_2} + \sum_{k_1+k_2+k_3 \geq n-1} B_{k_1}B_{k_2}B_{k_3} \right. & \\ \left. + r_{n-1}\partial r_{n-1} + \left(\sum_{i=0}^{n-1} B_i \right) \partial r_{n-1} + r_{n-1} \partial \left(\sum_{i=0}^{n-1} B_i \right) \right. & \\ \left. + r_{n-1}^3 + 3r_{n-1}^2 \left(\sum_{i=0}^{n-1} B_i \right) + 3r_{n-1} \left(\sum_{i=0}^{n-1} B_i \right)^2 + c(h_0 + q_0 - \mathbf{1}_{n=5}h_0) \right. & \\ \left. + \mathring{C} \left(\sum_{i=(n-3) \vee 0}^{n-1} B_i + r_{n-1} \right) \right) \Bigg), & \end{aligned}$$

$$\begin{aligned} |\mathcal{P}(B_{k_1}\partial B_{k_2})|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{1}{2}(\eta(k_1)+\eta(k_2)+1-\omega)} \cdot t^{b(k_1)} \cdot t^{b(k_2)} \\ &\leq t^{-\omega/2-\kappa+(1/4-\kappa/2)n} = t^{(n+1)/4-\kappa_n} \\ |\mathcal{P}(B_{k_1}B_{k_2}B_{k_3})|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{1}{2}(\sum_i \eta(k_i)+2-\omega)} \cdot t^{\sum_i b(k_i)} \\ &\leq t^{\frac{1}{4}-\frac{\omega}{2}-\frac{3}{2}\kappa+(1/4-\kappa/2)(n-1)} = t^{(n+1)/4-\kappa_n} \end{aligned}$$

$$\begin{aligned}
|\mathcal{P}B_i|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{1}{2}(\eta(i)+2-\omega)} |B_i|_{\mathcal{D}_{\alpha(i)}^{\frac{3}{2}+2\kappa,\eta(i)}} \\
&\lesssim t^{\frac{1}{2}(\eta(i)+2-\omega)} \cdot t^{b(i)} = t^{\frac{3}{4}\frac{\omega}{2}-\frac{\kappa}{2}+(\frac{1}{4}-\frac{\kappa}{2})i} \leq t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}(B_0\partial r_{n-1})|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{1/4-\kappa/2} |B_0\partial r_{n-1}|_{\mathcal{D}_{-\frac{3}{2}-2\kappa}^{\kappa,\omega-\frac{3}{2}-\kappa}} \\
&\lesssim t^{1/4-\kappa/2} |B_0|_{\mathcal{D}_{-\frac{1}{2}-\kappa}^{\frac{3}{2}+2\kappa,-\frac{1}{2}-\kappa}} |\partial r_{n-1}|_{\mathcal{D}_{-1-\kappa}^{\frac{1}{2}+2\kappa,\omega-1}} \\
&\lesssim t^{1/4-\kappa/2} t^{n/4-\kappa_{n-1}} = t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}(r_{n-1}\partial B_0)|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}(B_0^2 r_{n-1})|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{1/2-\kappa} |B_0^2 r_{n-1}|_{\mathcal{D}_{-1-3\kappa}^{\frac{1}{2},\omega-1-2\kappa}} \\
&\lesssim t^{1/2-\kappa} |B_0|_{\mathcal{D}_{-\frac{1}{2}-\kappa}^{\frac{3}{2}+2\kappa,-\frac{1}{2}-\kappa}}^2 |r_{n-1}|_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa,\omega}} \\
&\lesssim t^{1/2-\kappa} t^{n/4-\kappa_{n-1}} = t^{\frac{1}{4}\frac{\kappa}{2}} t^{(n+1)/4-\kappa_n} \leq t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}(B_0 r_{n-1}^2)|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{\omega}{2}+\frac{3}{4}-\frac{\kappa}{2}} |B_0 r_{n-1}^2|_{\mathcal{D}^{1,2\omega-\frac{1}{2}-3\kappa}} \\
&\lesssim t^{\frac{\omega}{2}+\frac{3}{4}-\frac{\kappa}{2}} |B_0|_{\mathcal{D}_{-\frac{1}{2}-\kappa}^{\frac{3}{2}+2\kappa,-\frac{1}{2}-\kappa}} |r_{n-1}|_{\mathcal{D}_{-\kappa}^{3+2\kappa,\omega}}^2 \\
&\lesssim t^{\frac{\omega}{2}+\frac{3}{4}-\frac{\kappa}{2}} (t^{n/4-\kappa_{n-1}})^2 \leq t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}(r_{n-1}\partial r_{n-1})|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{(1+\omega)/2} |r_{n-1}\partial r_{n-1}|_{\mathcal{D}_{-1-2\kappa}^{\frac{1}{2}+\kappa,2\omega-1}} \\
&\lesssim t^{(1+\omega)/2} |r_{n-1}|_{\mathcal{D}_{-\kappa}^{\frac{3}{2}+2\kappa,\omega}} \left| \partial r_{n-1} \right|_{\mathcal{D}_{-1-\kappa}^{\frac{1}{2}+2}} \\
&\lesssim t^{(1+\omega)/2} (t^{n/4-\kappa_{n-1}})^2 \leq t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}(r_{n-1}^3)|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{1+\omega} |r_{n-1}^3|_{\mathcal{D}_{-3\kappa}^{\frac{3}{2},3\omega}} \lesssim t^{1+\omega} |r_{n-1}|_{\mathcal{D}_{-\kappa}^3}^3 + 2\kappa, \omega \\
&\lesssim t^{1+\omega} (t^{n/4-\kappa_{n-1}})^3 \leq t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}q_0|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{1}{2}(2-\kappa-\omega)} |q_0|_{\mathcal{D}_0^{\frac{1}{2}+2\kappa,-\kappa}} \lesssim t^{\frac{1}{2}(2-\kappa-\omega)} \cdot t^{\frac{1}{4}} \leq t^{(n+1)/4-\kappa_n} \\
|\mathcal{P}h_0|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{1}{2}(2-\omega)} |h_0|_{\mathcal{D}_0^{\frac{1}{2}+2\kappa,0}} \lesssim t^{(n+1)/4-\kappa_n} \\
\Psi &= \mathcal{R}\Psi = \mathcal{P} \star \mathbf{1}_+ \xi \\
(\mathcal{R}X)(t) &= X(0) + \Psi_t + \mathcal{P}_t \star (\Psi \partial \Psi) + O_{L^\infty} \left(t^{\frac{1}{4}-3\kappa/2} \right). \\
\mathcal{R}B_1 &= \mathcal{P} \star (PX(0)P'X(0) + PX(0)\partial\Psi + P'X(0)\Psi + \Psi\partial\Psi) \\
|\Psi|_{L^\infty} &\lesssim t^{-\frac{1}{4}-\kappa/2}, |\partial\Psi|_{L^\infty} \lesssim t^{-\frac{3}{4}-\kappa/2} \\
|r_1|_{\mathcal{D}_0^{\frac{3}{2}+2\kappa,\omega}} &\lesssim t^{\frac{1}{2}-\kappa_1}, |(\mathcal{R}r_1)(t)|_{L^\infty} \lesssim t^{\frac{1}{2}-\kappa_1+\omega/2} = t^{\frac{1}{4}-3\kappa/2} \\
\mathcal{N}(A): (0, \infty) &\rightarrow \mathcal{C}^\infty(\mathbf{T}^3, \mathfrak{g}^3 \otimes (\mathfrak{g}^3)^3), \mathcal{N}_s(A) \stackrel{\text{def}}{=} P_s A \otimes \nabla P_s A \\
\mathcal{L} &= \mathbf{T}^3 \times \{v \in \mathbf{R}^3: |v| \leq 1/4\}
\end{aligned}$$

$$\begin{aligned}
|f|_{\gamma\text{-gr}} &= \sup_{\ell \in \mathcal{L}} \frac{|\int_{\ell} f|}{|\ell|^{\gamma}}, \\
\int_{\ell} f &= \int_0^1 |v|f(y+tv)dt \in F \\
[[A; B]]_{\gamma, \delta} &\stackrel{\text{def}}{=} \sup_{s \in (0,1)} s^{\delta} |\mathcal{N}_s A - \mathcal{N}_s B|_{\gamma\text{-gr}}, \quad [[A]]_{\gamma, \delta} = [[A; 0]]_{\gamma, \delta} \\
[[A; B]]_{\beta, \delta} &\stackrel{\text{def}}{=} \sup_{s \in (0,1)} s^{\delta} |\mathcal{N}_s A - \mathcal{N}_s B|_{\mathcal{C}^{\beta}} \\
\|A\|_{\alpha, \theta} &\stackrel{\text{def}}{=} \sup_{s \in (0,1)} |P_s A|_{\alpha\text{-gr}; < s^{\theta}} \\
|A|_{\alpha\text{-gr}; < r} &\stackrel{\text{def}}{=} \sup_{\ell \in \mathcal{L}, |\ell| < r} \frac{|A(\ell)|}{|\ell|^{\alpha}} \\
|A|_{\mathcal{C}^{\eta}} &\lesssim \|A\|_{\alpha, \theta} \\
\Sigma(A) &\stackrel{\text{def}}{=} \|A\|_{\alpha, \theta} + [[A]]_{\gamma, \delta} < \infty \\
\Sigma(A, B) &\stackrel{\text{def}}{=} \|A - B\|_{\alpha, \theta} + [[A; B]]_{\gamma, \delta} < \infty \\
\alpha \in (0, 1/2), \theta > 0, \gamma \in (1/2, 1), \delta \in (0, 1) \\
\mu &\stackrel{\text{def}}{=} \gamma - 1 + 2(1 - \delta) \in (-1/2, 0), \text{ and } \eta + \mu > -1 \\
|\mathcal{F}_s A|_{\infty} &\lesssim s^{\frac{\eta}{2}} \Sigma(A), |\partial \mathcal{F}_s A|_{\infty} \lesssim s^{\frac{\eta}{2} - \frac{1}{2}} \Sigma(A) \\
|P_s A|_{\infty} &\lesssim s^{\frac{\eta}{2}} |A|_{\mathcal{C}^{\eta}}, |\partial P_s A|_{\infty} \lesssim s^{\frac{\eta}{2} - \frac{1}{2}} |A|_{\mathcal{C}^{\eta}} \\
\lambda &\stackrel{\text{def}}{=} (1 - \zeta)\eta/2 - \theta(1 - \alpha)\zeta < 0 \\
|P_s A|_{\gamma\text{-gr}} &\lesssim s^{\lambda} \|A\|_{\alpha, \theta} \\
0 < \nu &\leq \min \left\{ \frac{\eta}{2} + \frac{\mu}{2} + \frac{1}{2}, 1 + 3\eta/2, \mu + \frac{1}{2} \right\} \\
|\mathcal{P}_s \star \mathcal{N}A|_{\gamma\text{-gr}} &\lesssim \int_0^s |\mathcal{N}_r A|_{\gamma\text{-gr}} dr \lesssim \int_0^s r^{-\delta} [[A]]_{\gamma, \delta} dr \lesssim s^{1-\delta} [[A]]_{\gamma, \delta} \\
\mathcal{F}_s A &= P_s A + \mathcal{P}_s \star \mathcal{N}A + R_s A \\
|R_s A|_{\infty} &\lesssim s^{\nu} (\Sigma(A) + \Sigma(A)^3) \\
R_s &= \int_0^s P_{s-r} \{ (P_r A + \mathcal{P}_r \star \mathcal{N}A + R_r) \partial (P_r A + \mathcal{P}_r \star \mathcal{N}A + R_r) \\
&\quad + (\mathcal{F}_r A)^3 \} dr - \mathcal{P}_s \star \mathcal{N}A \\
|\mathcal{P}_s \star \mathcal{N}A|_{\infty} &\lesssim \int_0^s |P_{s-r}(\mathcal{N}_r A)|_{\infty} dr \lesssim \int_0^s (s-r)^{\frac{\gamma-1}{2}} |\mathcal{N}_r A|_{\mathcal{C}^{\gamma-1}} ds \\
&\lesssim \int_0^s (s-r)^{\frac{\gamma-1}{2}} |\mathcal{N}_r A|_{\gamma\text{-gr}} ds \\
&\lesssim \int_0^s (s-r)^{\frac{\gamma-1}{2}} r^{-\delta} [[A]]_{\gamma, \delta} ds \asymp s^{\mu/2} [[A]]_{\gamma, \delta}, \\
|\partial \mathcal{P}_s \star \mathcal{N}A|_{\infty} &\lesssim s^{\mu/2 - \frac{1}{2}} [[A]]_{\gamma, \delta} \\
\int_0^s \left\{ r^{\frac{\eta}{2} + \frac{\mu}{2} - \frac{1}{2}} + r^{\mu - \frac{1}{2}} + r^{3\eta/2} \right\} dr &\asymp s^{\frac{\eta}{2} + \frac{\mu}{2} + \frac{1}{2}} + s^{\mu + \frac{1}{2}} + s^{1+3\eta/2} \\
|R|_{\mathcal{B}} &\stackrel{\text{def}}{=} \sup_{s \in (0, T)} \left\{ s^{-\nu} |R_s|_{\infty} + s^{-\nu + \frac{1}{2}} |\partial R_s|_{\infty} \right\} \\
s^{-\nu} |R|_{\infty} &\lesssim \Sigma(A) + \Sigma(A)^3 + s^{\kappa'} |R|_{\mathcal{B}} (\Sigma(A) + \Sigma(A)^2) \\
&\quad + s^{\kappa''} |R|_{\mathcal{B}}^2 (1 + \Sigma(A)) + s^{1+2\nu} |R|_{\mathcal{B}}^3
\end{aligned}$$



$$\begin{aligned}
\kappa' &= \min \left\{ \frac{\eta}{2} + \frac{1}{2}, 1 + \eta \right\} = \frac{\eta}{2} + \frac{1}{2} \\
\kappa'' &= \min \left\{ \frac{1}{2} + \nu, 1 + \frac{\eta}{2} + \nu \right\} = \frac{1}{2} + \nu \\
|R|_{\mathcal{B}} &\lesssim \Sigma(A) + \Sigma(A)^3 + T^\kappa (|R|_{\mathcal{B}} + |R|_{\mathcal{B}}^3) (1 + \Sigma(A)^2) \\
\mathcal{F}_s A &= P_s A + O_{\Omega_{\gamma-\text{gr}}[\mathfrak{g}, \mathfrak{g}]}(s^\nu (\Sigma(A) + \Sigma(A)^3)), P_s A = O_{\Omega_{\gamma-\text{gr}}}(s^\lambda) \dots \\
\mathcal{F}_s \tilde{A} &= \mathcal{F}_s A + P_s r + O_{L^\infty[\mathfrak{g}, \mathfrak{g}]}(s^{\eta/2 + \frac{1}{2}} |r|_\infty) \\
\mathcal{F}_s A + Q_s &= P_s(A + r) + \int_0^s P_{s-u} \{ (\mathcal{F}_u A + Q_u) \partial (\mathcal{F}_u A + Q_u) + (\mathcal{F}_u A + Q_u)^3 \} du \\
Q_s &= P_s r + \int_0^s P_{s-u} (Q_u \partial \mathcal{F}_u A + (\mathcal{F}_u A) \partial Q_u + Q_u \partial Q_u \\
&\quad + (\mathcal{F}_u A)^2 Q_u + (\mathcal{F}_u A) Q_u^2 + Q_u^3) du \\
&\int_0^s \left\{ |Q_u|_\infty u^{\eta/2 - \frac{1}{2}} + u^{\eta/2} |Q_u|_{\mathcal{C}^1} + |Q_u|_\infty |Q_u|_{\mathcal{C}^1} \right. \\
&\quad \left. + u^\eta |Q_u|_\infty + u^{\eta/2} |Q_u|_\infty^2 + |Q_u|_\infty^3 \right\} du \\
&\lesssim |Q|_\infty s^{\eta/2 + \frac{1}{2}} + s^{\eta/2 + \frac{1}{2}} |Q|_{L_{1/2}^\infty \mathcal{C}^1} + s^{1/2} |Q|_\infty |Q|_{L_{1/2}^\infty \mathcal{C}^1} \\
&\quad + s^{\eta+1} |Q|_\infty + s^{\eta/2+1} |Q|_\infty^2 + s |Q|_\infty^3, \\
&\quad |Q|_\infty + |Q|_{L_{1/2}^\infty \mathcal{C}^1} \lesssim |r|_\infty \\
Q_s &= P_s r + O_{L^\infty}(s^{\eta/2 + \frac{1}{2}} |r|_\infty) \\
\tilde{\mathcal{X}} &= B + c\bar{h} + O_{\mathcal{D}_0^{\frac{3}{2}+2\kappa, \omega}}(t^{3/2-\kappa_5}) \\
\tilde{\mathcal{X}} &= B + c\mathcal{P}\text{Ph}(0) + O_{\mathcal{D}_0^{\frac{3}{2}+2\kappa, \omega}}(t^{3/2-\kappa_5}), \\
\tilde{A} &= A + ctP_t h(0) + O_{L^\infty}(t^{5/4-7\kappa/2}). \\
\mathcal{F}_s \tilde{A} &= \mathcal{F}_s A + P_s \left(ctP_t h(0) + O_{L^\infty}(t^{5/4-7\kappa/2}) \right) + O_{L^\infty[\mathfrak{g}, \mathfrak{g}]}(s^{\eta/2 + \frac{1}{2}} t) \\
|f|_{p\text{-var}} &\stackrel{\text{def}}{=} \sup_{P \subset [0,1]} \left(\sum_{[s,t] \in P} |f(t) - f(s)|^p \right)^{1/p} \\
dJ^\gamma(x) &= J^\gamma(x) d\gamma(x), J^\gamma(0) = \text{id} \\
J^{\gamma+\zeta}(1) &= J^\gamma(1) + \int_0^1 d\zeta(x) + \int_0^1 \int_0^x \{ d\zeta(x) d\gamma(y) + d\gamma(x) d\zeta(y) \} \\
&\quad + O\{v(w^2 + w^{L-1}) + w^L + w^{L+1} + v^{L+1} + v^2(1 + w + v + w^{L-3})\} \\
J^\gamma(1) &= \text{id} + I^\gamma + \int_0^1 \dots \int_0^{x_{L-1}} J^\gamma(x_L) d\gamma(x_L) d\gamma(x_{L-1}) \dots d\gamma(x_1) \\
I^{\gamma+\zeta} &= I^\gamma + \int_0^1 d\zeta(x) + \int_0^1 \int_0^x \{ d\zeta(x) d\gamma(y) + d\gamma(x) d\zeta(y) \} \\
&\quad + O\{v^2 + v(w^2 + w^{L-1})\} + O\{v^2(w + v + w^{L-3} + v^{L-3})\} \\
|\ell_a|_{\frac{1}{\gamma}\text{-var}} &\leq |\ell_a|_{\gamma\text{-H\"{o}l}} \leq |a|_{\gamma\text{-gr}}
\end{aligned}$$

$$\begin{aligned}
W_\ell(\mathcal{F}_s \tilde{A}) &= W_\ell(\mathcal{F}_s A) + t \text{Tr} \int_\ell \text{ch}(0) + t \text{Tr} \int_{[0,1]^2} d\ell_{A(0)}(x_1) d\ell_{\text{ch}(0)}(x_2) \\
&+ O\left(t^{5/4-7\kappa/2} + ts + ts^\lambda \|\Psi_t^{\text{YM}}\|_{\alpha,\theta} + ts^\lambda \|\mathcal{P}_t \star (\Psi \partial \Psi)^{\text{YM}}\|_{\alpha,\theta} + |A(0)|_{L^\infty} s^{\eta/2+\frac{1}{2}} t\right. \\
&\quad \left.+ s^\nu t(\Sigma(A) + \Sigma(A)^3) + t(u^2 + u^{L-1}) + s^{\nu L} + u^L + u^{L+1} + t^2 u\right) \\
\gamma(x) &= \int_0^x (\mathcal{F}_s A)_1(y, 0, 0) dy = \ell_{\mathcal{F}_s A}(x) \\
W_\ell(\mathcal{F}_s A) &= \text{Tr} J^\nu(1), W_\ell(\mathcal{F}_s \tilde{A}) = \text{Tr} J^{\nu+\zeta}(1) \\
D_h &\stackrel{\text{def}}{=} ct P_{t+s} h(0) = cth(0) + O_{L^\infty}(t(t+s)), \\
D_{\text{err}} &\stackrel{\text{def}}{=} O_{L^\infty[\mathfrak{g}, \mathfrak{g}]} \left(s^{\eta/2+\frac{1}{2}} t \right) + O_{L^\infty}(t^{5/4-7\kappa/2}).
\end{aligned}$$

$$\begin{aligned}
W_\ell(\mathcal{F}_s \tilde{A}) &= W_\ell(\mathcal{F}_s A) + \text{Tr} \left(\int_0^1 d\zeta(x) \right) \\
&+ \text{Tr} \int_0^1 \int_0^x \{ d\zeta(x) d\gamma(y) + d\gamma(x) d\zeta(y) \} \\
&+ O\{v(w^2 + w^{L-1}) + w^L + w^{L+1} + v^{L+1} + v^2(1 + w + v + w^{L-3})\} \\
&\quad \text{Tr} \left(\int_0^1 d\zeta(x) \right) = t \text{Tr} \int_\ell \text{ch}(0) + O(t^{5/4-7\kappa/2} + ts) \\
\text{Tr} \int_0^1 \int_0^x \{ d\zeta(x) d\gamma(y) + d\gamma(x) d\zeta(y) \} &= \text{Tr} \int_{[0,1]^2} \{ d\zeta(x) d\gamma(y) \} \\
&\quad \mathcal{F}_s A = P_s A + O_{\gamma\text{-gr}}(s^\nu(\Sigma(A) + \Sigma(A)^3)) \\
&\quad A = A(0) + \Psi_t^{\text{YM}} + \mathcal{P}_t \star (\Psi \partial \Psi)^{\text{YM}} + O_{L^\infty}(t^{1/4-3\kappa/2}) \\
\int_0^1 d\gamma(x) &= \int_0^1 d\ell_{A(0)}(x) + O\left(s^\lambda \|\Psi_t^{\text{YM}}\|_{\alpha,\theta} + s^\lambda \|\mathcal{P}_t \star (\Psi \partial \Psi)^{\text{YM}}\|_{\alpha,\theta}\right) \\
&\quad + O\left(t^{1/4-3\kappa/2} + s^\nu(\Sigma(A) + \Sigma(A)^3)\right) \\
\int_0^1 d\zeta(y) &= t \int_0^1 d\ell_{\text{ch}(0)}(y) + O\left(t^{5/4-7\kappa/2} + s^{\eta/2+\frac{1}{2}} t\right) \\
\text{Tr} \int_{[0,1]^2} \{ d\zeta(x) d\gamma(y) \} &= t \text{Tr} \left(\int_{[0,1]^2} d\ell_{A(0)}(x_1) d\ell_{\text{ch}(0)}(x_2) \right) \\
&+ O\left(ts^\lambda \|\Psi_t^{\text{YM}}\|_{\alpha,\theta} + ts^\lambda \|\mathcal{P}_t \star (\Psi \partial \Psi)^{\text{YM}}\|_{\alpha,\theta} + t^{5/4-7\kappa/2}\right) \\
&+ O\left(ts^\nu(\Sigma(A) + \Sigma(A)^3) + s^{\eta/2+\frac{1}{2}} t |A(0)|_{L^\infty}\right) \\
w = |\gamma|_{\frac{1}{\gamma}\text{-var}} &\leq |\mathcal{F}_s A|_{\gamma\text{-gr}} = |\mathcal{P}_s A|_{\gamma\text{-gr}} + O(s^\nu) = O(s^\lambda) \|A\|_{\alpha,\theta} + O(s^\nu) \\
v = |\zeta|_{\frac{1}{\gamma}\text{-var}} &\leq |D_h + D_{\text{err}}|_{\gamma\text{-gr}} \lesssim s^{\eta/2+\frac{1}{2}} t + t = O(t) \\
[A + B]_{\gamma,\delta} &\lesssim [A]_{\gamma,\delta} + |B|_{\mathcal{C}\bar{\eta}} (|A|_{\mathcal{C}\eta} + |B|_{\mathcal{C}\bar{\eta}}) \\
\mathcal{N}_s(A + B) &= \mathcal{N}_s A + \mathcal{N}_s B + P_s A \otimes \nabla P_s B + P_s B \otimes \nabla P_s A \\
Z_{s,t} &= s^\delta \mathcal{N}_s \Psi_t
\end{aligned}$$

$$\begin{aligned}
& \mathbf{E} \left[\left| \sup_{(s,t) \neq (\bar{s}, \bar{t})} \frac{|Z_{s,t} - Z_{\bar{s}, \bar{t}}|_{\gamma\text{-gr}}}{(|t - \bar{t}| + |s - \bar{s}|)^{\bar{\kappa}}} \right|^p \right]^{1/p} < \infty \\
& (\mathbf{E} | Z_{s,t} - Z_{\bar{s}, \bar{t}} \bar{\tau}_{\gamma\text{-gr}}^p)^{1/p} \lesssim (|t - \bar{t}| + |s - \bar{s}|)^{\kappa/2} \\
& \mathbf{E} \left| \sup_{t \in [0,1]} t^{-\bar{\kappa}} [\Psi_t]_{\gamma, \delta} \right|^p < \infty \\
& |C_{s, \bar{s}, t, \bar{t}}(x)| \lesssim (|s - \bar{s}| + |t - \bar{t}|)^{\kappa} |x|^{4\delta-4-2\kappa} \\
& |\nabla C_{s, \bar{s}, t, \bar{t}}| \lesssim (|s - \bar{s}| + |t - \bar{t}|)^{\kappa} |x - y|^{4\delta-5-2\kappa} \\
& |\nabla C_{r,s}(x)| \lesssim |x|^{-2}, |\nabla (C_{r,r} - C_{r,s})(x)| \lesssim |r - s|^{\kappa} |x|^{-2-2\kappa} \\
& d((x, v), (\bar{x}, \bar{v})) \stackrel{\text{def}}{=} |x - \bar{x}| \vee |x + v - (\bar{x} + \bar{v})| \\
& \left(\mathbf{E} \left| \int_{\ell} (Z_{s,t} - Z_{\bar{s}, \bar{t}})^p \right|^{1/p} \lesssim (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell| \right)^{2\delta-1-\kappa} \\
& \left(\mathbf{E} \left| \left(\int_{\ell} - \int_{\bar{\ell}} \right) (Z_{s,t} - Z_{\bar{s}, \bar{t}}) \right|^p \right)^{1/p} \lesssim (|t - \bar{t}| + |s - \bar{s}|)^{\kappa/2} d(\ell, \bar{\ell})^{2\delta-3/2-\kappa} \\
& \mathbf{E} \left| \int_{\ell} (Z_{s,t} - Z_{\bar{s}, \bar{t}}) \right|^2 = |\ell|^2 \int_{[0,1]^2} C((r - \bar{r})v) dr d\bar{r} \\
& \lesssim (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell|^{4\delta-2-2\kappa} \int_{[0,1]^2} |r - \bar{r}|^{4\delta-4-2\kappa} dr d\bar{r} \\
& \lesssim (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell|^{4\delta-2-2\kappa} \\
& \mathbf{E} | (\int_{\ell} - \int_{\bar{\ell}}) (Z_{s,t} - Z_{\bar{s}, \bar{t}}) |^2 = |\ell|^2 \int_{[0,1]^2} \{ C((r - \bar{r})v) - C(rv - \bar{r}\bar{v}) \\
& \quad - C(r\bar{v} - \bar{r}v) + C((r - \bar{r})\bar{v}) \} dr d\bar{r} \\
& = |\ell|^2 \int_{[0,1]^2} \{ 2C((r - \bar{r})v) - 2C(rv - \bar{r}\bar{v}) \} dr d\bar{r}, \\
& (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell|^{4\delta-2-2\kappa} \int_0^h r^{4\delta-4-2\kappa} dr \\
& \asymp (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell|^{4\delta-2-2\kappa} h^{4\delta-3-2\kappa} \\
& = (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell| |v - \bar{v}|^{4\delta-3-2\kappa} \\
& (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell|^{4\delta-3-2\kappa} |v - \bar{v}| \int_h^1 r^{4\delta-5-2\kappa} dr \\
& \asymp (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell|^{4\delta-3-2\kappa} |v - \bar{v}| h^{4\delta-4-2\kappa} \\
& = (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell| |v - \bar{v}|^{4\delta-3-2\kappa} \\
& \mathbf{E} | \left(\int_{\ell} - \int_{\bar{\ell}} \right) (Z_{s,t} - Z_{\bar{s}, \bar{t}}) |^2 \lesssim (|t - \bar{t}| + |s - \bar{s}|)^{\kappa} |\ell| |v - \bar{v}|^{4\delta-3-2\kappa} \\
& \left(\mathbf{E} \left| \int_{\ell} A \right|^p \right)^{1/p} \leq M_p |\ell|^{\alpha} \\
& \left(\mathbf{E} \left| \left(\int_{\ell} - \int_{\bar{\ell}} \right) A \right|^p \right)^{1/p} \leq M_p d(\ell, \bar{\ell})^{\beta} \\
& (\mathbf{E} | A |_{\gamma\text{-gr}}^p)^{1/p} \leq \lambda M_p
\end{aligned}$$

$$\begin{aligned}
& \sup_{\ell \in D} \frac{|\int_{\ell} A|}{|\ell|^{\gamma}} \lesssim \sup_{N \geq 1} \sup_{\substack{a \in D_N \\ |a| \leq K2^{-N/\omega}}} \frac{|\int_a A|}{2^{-\gamma N/\omega}} + \sup_{N \geq 1} \sup_{\substack{a, b \in D_N \\ d(a, b) \leq K2^{-N}}} \frac{|(\int_a - \int_b) A|}{2^{-\gamma N/\omega}}. \\
& \mathbf{E} \left| \sup_{\ell \in D} \frac{|\int_{\ell} A|}{|\ell|^{\gamma}} \right|^p \lesssim M_p^p \sum_{N \geq 1} \{2^{N(6-p(\alpha-\gamma)/\omega)} + 2^{N(12-p(\beta-\gamma/\omega))}\} \\
& \alpha = \frac{1}{2} - \varepsilon, \theta = \varepsilon, \gamma = \frac{1}{2} + \varepsilon, \delta = 1 - \varepsilon, \nu = -\bar{\eta} = \varepsilon/2 \\
& \|f\|_{\alpha, \theta} \lesssim |f|_{\mathcal{C}^{\bar{\eta}}} \\
& \mathcal{P} \star (\Psi \partial \Psi) = \lim_{\varepsilon \downarrow 0} \mathcal{P} \star (\Psi_{\varepsilon} \partial \Psi_{\varepsilon}), \\
& \mathbf{E} \|\Psi_t\|_{\alpha, \theta}^p + \mathbf{E} |\mathcal{P}_t \star (\Psi \partial \Psi)|_{\mathcal{C}^{\bar{\eta}}}^p + \mathbf{E} [\Psi]_{\gamma, \delta}^p = O(t^{p\varepsilon}). \\
& Q_t = \left\{ \|Z\|_{\frac{3}{2}+2\kappa; O} + |\mathcal{P} \star \mathbf{1}_{+\xi}|_{\mathcal{C}([-1,3], \mathcal{C}^{-1/2-\kappa})} + |\mathcal{P} \star (\Psi \partial \Psi)|_{\mathcal{C}([-1,3], \mathcal{C}^{-2\kappa})} \right. \\
& \quad \left. + t^{-\varepsilon} \|\Psi_t\|_{\alpha, \theta} + t^{-\varepsilon} |\mathcal{P}_t \star (\Psi \partial \Psi)|_{\mathcal{C}^{\bar{\eta}}} + t^{-\varepsilon} [\Psi_t]_{\gamma, \delta} < M \right\} \\
& \{ \|Z\|_{\frac{3}{2}+2\kappa; O} + |\mathcal{P} \star \mathbf{1}_{+\xi}|_{\mathcal{C}([-1,3], \mathcal{C}^{-1/2-\kappa})} + |\mathcal{P} \star (\Psi \partial \Psi)|_{\mathcal{C}([-1,3], \mathcal{C}^{-2\kappa})} < M \} \supset Q_t, \\
& \mathbf{E} W_{\ell}(\mathcal{F}_s \tilde{A}) \mathbf{1}_{Q_t} - \mathbf{E} W_{\ell}(\mathcal{F}_s A) \mathbf{1}_{Q_t} = \mathbf{P}[Q_t] \left\{ t \text{Tr} \left(\int_{\ell} \text{ch}(0) \right) \right. \\
& \quad \left. + t \text{Tr} \left(\int_{[0,1]^2} d\ell_{A(0)}(x_1) d\ell_{\text{ch}(0)}(x_2) \right) + O(t^{1+r+\beta/6} + t^{1+3r/2} + t^{1+r+\nu\beta}) \right\} \\
& \|A\|_{\alpha, \theta} \leq |A(0)|_{\infty} + \|\Psi_t\|_{\alpha, \theta} + \|\mathcal{P}_t \star (\Psi \partial \Psi)\|_{\alpha, \theta} + O(t^{1/4-3\kappa/2}) \lesssim t^r \\
& \beta = -\frac{r}{4\lambda} > 0 \\
& u = s^{\lambda} \|A\|_{\alpha, \theta} \lesssim s^{\lambda} t^r = t^{3r/4} \\
& [\Psi_t]_{\gamma, \delta} \lesssim [\Psi_t]_{\gamma, \delta} + |\Psi|_{\mathcal{C}^{\eta}} \{ |A(0)|_{L^{\infty}} + |\mathcal{P}_t(\Psi \partial \Psi)|_{\mathcal{C}^{\bar{\eta}}} + O(t^{1/4-3\kappa/2}) \} \lesssim t^r \\
& t^{5/4-7\kappa/2} + ts \lesssim t^{1+3r/2}, \\
& ts^{\lambda} (\|\Psi_t\|_{\alpha, \theta} + \|\mathcal{P}_t \star (\Psi \partial \Psi)\|_{\alpha, \theta}) \lesssim s^{\lambda} t^{1+\varepsilon} \lesssim t^{1+3r/2}, \\
& |A(0)|_{L^{\infty}} s^{\eta/2 + \frac{1}{2}} t \lesssim t^{r+1} t^{(\frac{\eta}{2} + \frac{1}{2})\beta} \lesssim t^{1+r+\beta/6}, \\
& s^{\nu} t (\Sigma(A) + \Sigma(A)^3) \lesssim t^{1+r+\nu\beta}, \\
& t(u^2 + u^{L-1}) \lesssim t^{1+3r/2}, \\
& |\mathbf{E}\{W_{\ell}(\mathcal{F}_s \tilde{A}) - W_{\ell}(\mathcal{F}_s A)\} \mathbf{1}_{Q_t}| \gtrsim t^{1+r} \\
& A(0) = t^r \text{ch}(0). \\
& \int_{\ell} \text{ch}(0) = \int_{\ell} c_1^{(1)} h_1(0) = c_1^{(1)} \zeta(1). \\
& |\mathbf{E} W_{\ell}(\mathcal{F}_s \tilde{A}) \mathbf{1}_{Q_t} - \mathbf{E} W_{\ell}(\mathcal{F}_s A) \mathbf{1}_{Q_t}| \gtrsim t \\
& |\mathbf{E} W_{\ell}(\mathcal{F}_s \tilde{A}) \mathbf{1}_{Q_t} - \mathbf{E} W_{\ell}(\mathcal{F}_s A) \mathbf{1}_{Q_t}| \gtrsim t^{1+r} \left| \text{Tr} \left(\left\{ c_1^{(1)} \zeta(1) \right\}^2 \right) \right| - o(t^{1+r}) \gtrsim t^{1+r} \\
& u(x, y, z) = \begin{cases} e^{\psi(y)X} & \text{if } y \in \left[0, \frac{1}{4}\right] \\ 1 & \text{if } y \in \left[\frac{1}{4}, \frac{3}{4}\right] \\ e^{\psi(y-1)X} & \text{if } y \in \left[\frac{3}{4}, 1\right] \end{cases} \\
& h_2(x, 0, 0) \stackrel{\text{def}}{=} (\partial_2 u) u^{-1}(x, 0, 0) = X, x \in [0, 1]
\end{aligned}$$

$$\begin{aligned}
\int_{\ell} ch &= \int_0^1 c_1^{(2)} h_2(x, 0, 0) dx = c_1^{(2)} X \\
\delta\psi_a^+ &= D_{\bar{a}}\epsilon^+ + \frac{1}{16} F_{\#} \gamma_{\bar{a}} \epsilon^-, \quad \delta\rho^+ = \gamma^a D_a \epsilon^+, \\
\delta\psi_{\bar{a}}^- &= D_a \epsilon^- + \frac{1}{16} F_{\#}^T \gamma_a \epsilon^+, \quad \delta\rho^- = \gamma^{\bar{a}} D_{\bar{a}} \epsilon^-. \\
R_{a\bar{b}} + \frac{1}{16} \Phi^{-1} \langle F, \Gamma_{a\bar{b}} F \rangle &= 0, S = 0, \Gamma^A D_A F = 0, \\
\gamma^b D_b \psi_{\bar{a}}^+ - D_{\bar{a}} \rho^+ &= \frac{1}{16} \gamma^b F_{\#} \gamma_{\bar{a}} \psi_b^-, \quad \gamma^a D_a \rho^+ - D^{\bar{a}} \psi_{\bar{a}}^+ = -\frac{1}{16} F_{\#} \rho^-, \\
\gamma^{\bar{b}} D_{\bar{b}} \psi_a^- - D_a \rho^- &= \frac{1}{16} \gamma^{\bar{b}} F_{\#}^T \gamma_a \psi_{\bar{b}}^+, \quad \gamma^{\bar{a}} D_{\bar{a}} \rho^- - D^a \psi_a^- = -\frac{1}{16} F_{\#}^T \rho^+. \\
&\{g_{\mu\nu}, B_{\mu\nu}, \phi, A_{\mu_1 \dots \mu_n}^{(n)}, \psi_{\mu}^{\pm}, \lambda^{\pm}\} \\
\psi_{\mu} &= \psi_{\mu}^+ + \psi_{\mu}^- \quad \gamma^{(10)} \psi_{\mu}^{\pm} = \mp \psi_{\mu}^{\pm} \\
\lambda &= \lambda^+ + \lambda^- \quad \gamma^{(10)} \lambda^{\pm} = \pm \lambda^{\pm} \\
\psi_{\mu} &= \begin{pmatrix} \psi_{\mu}^+ \\ \psi_{\mu}^- \end{pmatrix} \quad \gamma^{(10)} \psi_{\mu}^{\pm} = \psi_{\mu}^{\pm} \\
\lambda &= \begin{pmatrix} \lambda^+ \\ \lambda^- \end{pmatrix} \quad \gamma^{(10)} \lambda^{\pm} = -\lambda^{\pm}. \\
\rho^{\pm} &:= \gamma^{\mu} \psi_{\mu}^{\pm} - \lambda^{\pm}, \\
S_B &= \frac{1}{2\kappa^2} \int \sqrt{-g} \left[e^{-2\phi} \left(\mathcal{R} + 4(\partial\phi)^2 - \frac{1}{12} H^2 \right) - \frac{1}{4} \sum_n \frac{1}{n!} \left(F_{(n)}^{(B)} \right)^2 \right], \\
F^{(B)} &= \sum_n F_{(n)}^{(B)} = \sum_n e^B \wedge dA_{(n-1)}, \\
F_{(n)}^{(B)} &= (-)^{[n/2]} * F_{(10-n)}^{(B)}, \\
S_F &= -\frac{1}{2\kappa^2} \int \sqrt{-g} \left[e^{-2\phi} \left(2\bar{\psi}^{+\mu} \gamma^{\nu} \nabla_{\nu} \psi_{\mu}^+ - 4\bar{\psi}^{+\mu} \nabla_{\mu} \rho^+ - 2\bar{\rho}^+ \nabla \rho^+ \right. \right. \\
&\quad \left. \left. - \frac{1}{2} \bar{\psi}^{+\mu} \mathbb{H} \psi_{\mu}^+ - \bar{\psi}_{\mu}^+ H^{\mu\nu\lambda} \gamma_{\nu} \psi_{\lambda}^+ - \frac{1}{2} \rho^+ H^{\mu\nu\lambda} \gamma_{\mu\nu} \psi_{\lambda}^+ + \frac{1}{2} \rho^+ H / \rho^+ \right) \right. \\
&\quad \left. + e^{-2\phi} \left(2\bar{\psi}^{-\mu} \gamma^{\nu} \nabla_{\nu} \psi_{\mu}^- - 4\bar{\psi}^{-\mu} \nabla_{\mu} \rho^- - 2\bar{\rho}^- \nabla \rho^- \right. \right. \\
&\quad \left. \left. + \frac{1}{2} \bar{\psi}^{-\mu} H H \psi_{\mu}^- + \bar{\psi}_{\mu}^- H^{\mu\nu\lambda} \gamma_{\nu} \psi_{\lambda}^- + \frac{1}{2} \rho^- H^{\mu\nu\lambda} \gamma_{\mu\nu} \psi_{\lambda}^- - \frac{1}{2} \rho^- H / \rho^- \right) \right. \\
&\quad \left. - \frac{1}{4} e^{-\phi} \left(\bar{\psi}_{\mu}^+ \gamma^{\nu} F^{(B)} \gamma^{\mu} \psi_{\nu}^- + \rho^+ F^{(B)} \rho^- \right) \right]. \\
\mathcal{R}_{\mu\nu} - \frac{1}{4} H_{\mu\lambda\rho} H_{\nu}{}^{\lambda\rho} + 2\nabla_{\mu} \nabla_{\nu} \phi - \frac{1}{4} e^{2\phi} \sum_n \frac{1}{(n-1)!} F_{\mu\lambda_1 \dots \lambda_{n-1}}^{(B)} F_{\nu}^{(B)\lambda_1 \dots \lambda_{n-1}} &= 0 \\
\nabla^{\mu} (e^{-2\phi} H_{\mu\nu\lambda}) - \frac{1}{2} \sum_n \frac{1}{(n-2)!} F_{\mu\nu\lambda_1 \dots \lambda_{n-2}}^{(B)} F^{(B)\lambda_1 \dots \lambda_{n-2}} &= 0 \\
\nabla^2 \phi - (\nabla \phi)^2 + \frac{1}{4} \mathcal{R} - \frac{1}{48} H^2 &= 0 \\
dF^{(B)} - H \wedge F^{(B)} &= 0
\end{aligned}$$

$$\begin{aligned}
& \gamma^\nu \left[\left(\nabla_\nu \mp \frac{1}{24} H_{\nu\lambda\rho} \gamma^{\lambda\rho} - \partial_\nu \phi \right) \psi_\mu^\pm \pm \frac{1}{2} H_{\nu\mu}{}^\lambda \psi_\lambda^\pm \right] - \left(\nabla_\mu \mp \frac{1}{8} H_{\mu\nu\lambda} \gamma^{\nu\lambda} \right) \rho^\pm \\
&= \frac{1}{16} e^\phi \sum_n (\pm)^{[(n+1)/2]} \gamma^\nu \psi_{(n)}^{(B)} \gamma_\mu \psi_\nu^\mp, \\
& \left(\nabla_\mu \mp \frac{1}{8} H_{\mu\nu\lambda} \gamma^{\nu\lambda} - 2\partial_\mu \phi \right) \psi^{\mu\pm} - \gamma^\mu \left(\nabla_\mu \mp \frac{1}{24} H_{\mu\nu\lambda} \gamma^{\nu\lambda} - \partial_\mu \phi \right) \rho^\pm \\
&= \frac{1}{16} e^\phi \sum_n (\pm)^{[(n+1)/2]} F_{(n)}^{(B)} \rho^\mp, \\
& \epsilon = \epsilon^+ + \epsilon^- \quad \gamma^{(10)} \epsilon^\pm = \mp \epsilon^\pm \\
& \epsilon = \begin{pmatrix} \epsilon^+ \\ \epsilon^- \end{pmatrix} \quad \gamma^{(10)} \epsilon^\pm = \epsilon^\pm \\
& \delta e_\mu^a = \bar{\epsilon}^+ \gamma^a \psi_\mu^+ + \bar{\epsilon}^- \gamma^a \psi_\mu^- \\
& \delta B_{\mu\nu} = 2\bar{\epsilon}^+ \gamma_{[\mu} \psi_{\nu]}^+ - 2\bar{\epsilon}^- \gamma_{[\mu} \psi_{\nu]}^-, \\
& \delta\phi - \frac{1}{4} \delta \log(-g) = -\frac{1}{2} \bar{\epsilon}^+ \rho^+ - \frac{1}{2} \bar{\epsilon}^- \rho^-, \\
& (e^B \wedge \delta A)_{\mu_1 \dots \mu_n}^{(n)} = \frac{1}{2} (e^{-\phi} \bar{\psi}_\nu^+ \gamma_{\mu_1 \dots \mu_n} \gamma^\nu \epsilon^- - e^{-\phi} \bar{\epsilon}^+ \gamma_{\mu_1 \dots \mu_n} \rho^-) \\
& \quad \mp \frac{1}{2} (e^{-\phi} \bar{\epsilon}^+ \gamma^\nu \gamma_{\mu_1 \dots \mu_n} \psi_\nu^- + e^{-\phi} \bar{\rho}^+ \gamma_{\mu_1 \dots \mu_n} \epsilon^-), \\
& \delta \psi_\mu^\pm = \left(\nabla_\mu \mp \frac{1}{8} H_{\mu\nu\lambda} \gamma^{\nu\lambda} \right) \epsilon^\pm + \frac{1}{16} e^\phi \sum_n (\pm)^{[(n+1)/2]} F_{(n)}^{(B)} \gamma_\mu \epsilon^\mp, \\
& \delta \rho^\pm = \gamma^\mu \left(\nabla_\mu \mp \frac{1}{24} H_{\mu\nu\lambda} \gamma^{\nu\lambda} - \partial_\mu \phi \right) \epsilon^\pm. \\
& B_{(i)} = B_{(j)} - d\Lambda_{(ij)} \\
& \Lambda_{(ij)} + \Lambda_{(jk)} + \Lambda_{(ki)} = d\Lambda_{(ijk)}, \\
& A_{(i)} = e^{d\Lambda_{(ij)}} \wedge A_{(j)} + d\hat{\Lambda}_{(ij)}, \\
& B'_{(i)} = B_{(i)} - d\lambda_{(i)}, A'_{(i)} = e^{d\lambda_{(i)}} A_{(i)}, \\
& \Omega_{\text{cl}}^2(M) \rightarrow G_{\text{NS}} \rightarrow \text{Diff}(M), \\
& \delta_{v+\lambda} g = \mathcal{L}_v g, \delta_{v+\lambda} \phi = \mathcal{L}_v \phi, \delta_{v+\lambda} B_{(i)} = \mathcal{L}_v B_{(i)} - d\lambda_{(i)}, \\
& d\lambda_{(i)} = d\lambda_{(j)} - \mathcal{L}_v d\Lambda_{(ij)} \\
& \lambda_{(i)} = \lambda_{(j)} - i_v d\Lambda_{(ij)}, \\
& 0 \rightarrow T^*M \rightarrow E \rightarrow TM \rightarrow 0 \\
& v_{(i)} + \lambda_{(i)} = v_{(j)} + (\lambda_{(j)} - i_{v_{(j)}} d\Lambda_{(ij)}), \\
& \langle V, V \rangle = i_v \lambda, \\
& \tilde{E} = \det T^*M \otimes E \\
& \langle \hat{E}_A, \hat{E}_B \rangle = \Phi^2 \eta_{AB}, \quad \eta = \frac{1}{2} \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}. \\
& \tilde{F} = \{(x, \{\hat{E}_A\}): x \in M, \text{ and } \{\hat{E}_A\} \tilde{E}_x\}. \\
& V^A \mapsto V'^A = M_B^A V^B, \hat{E}_A \mapsto \hat{E}'_A = \hat{E}_B (M^{-1})^B{}_A. \\
& V^M = \begin{cases} v^\mu & \text{for } M = \mu \\ \lambda_\mu & \text{for } M = \mu + d \end{cases} \\
& E_{(p)}^{\otimes n} = (\det T^*M)^p \otimes E \otimes \dots \otimes E. \\
& \{\Gamma_A, \Gamma_B\} = 2\eta_{AB} \\
& V^A \Gamma_A \Psi_{(i)} = i_v \Psi_{(i)} + \lambda_{(i)} \wedge \Psi_{(i)} \\
& \Psi_{(i)} = e^{d\Lambda_{(ij)}} \wedge \Psi_{(j)}
\end{aligned}$$

$$\begin{aligned}
S_{(p)}^{\pm} &= (\det T^* M)^p \otimes S^{\pm}(E) \\
\langle \Psi, \Psi' \rangle &= \sum (-)^{[(n+1)/2]} \Psi^{(d-n)} \wedge \Psi'^{(n)} \in \Gamma((\det T^* M)^{2p}) \\
\hat{E}_A &= \begin{cases} \hat{E}_a = (\det e)(\hat{e}_a + i_{\hat{e}_a} B) & \text{para } A = a \\ E^a = (\det e)e^a & \text{para } A = a + d \end{cases} \\
\langle \hat{E}_A, \hat{E}_B \rangle &= (\det e)^2 \eta_{AB} \\
V^{(B)} &= v^a (\det e) \hat{e}_a + \lambda_a (\det e) e^a \\
&= v_{(i)} + \lambda_{(i)} - i_{v_{(i)}} B_{(i)}, \\
M &= (\det A)^{-1} \begin{pmatrix} 1 & 0 \\ \omega & 1 \end{pmatrix} \begin{pmatrix} A & 0 \\ 0 & (A^{-1})^T \end{pmatrix}, \\
\hat{E}_A &= \begin{cases} \hat{E}_a = e^{-2\phi} (\det e)(\hat{e}_a + i_{\hat{e}_a} B) & \text{para } A = a \\ E^a = e^{-2\phi} (\det e)e^a & \text{para } A = a + d \end{cases} \\
V^{(B,\phi)} &= e^{2\phi} (v_{(i)} + \lambda_{(i)} - i_{v_{(i)}} B_{(i)}). \\
\Psi^{(B)} &= \sum_n \frac{1}{n!} \Psi_{a_1 \dots a_n}^{(B)} e^{a_1} \wedge \dots \wedge e^{a_n} = e^{B_{(i)}} \wedge \Psi_{(i)}, \\
\Psi^{(B,\phi)} &= e^{p\phi} e^{B_{(i)}} \wedge \Psi_{(i)}. \\
L_V W &= \mathcal{L}_v w + \mathcal{L}_v \zeta - i_w d\lambda \\
\mathcal{L}_v w^\mu &= v^\nu \partial_\nu w^\mu - w^\nu \partial_\nu v^\mu + p(\partial_\nu v^\nu) w^\mu, \\
\mathcal{L}_v \zeta_\mu &= v^\nu \partial_\nu \zeta_\mu + (\partial_\mu v^\nu) \zeta_\nu + p(\partial_\nu v^\nu) \zeta^\mu. \\
\partial_M &= \begin{cases} \partial_\mu & \text{para } M = \mu \\ 0 & \text{para } M = \mu + d \end{cases} \\
L_V W^M &= V^N \partial_N W^M + (\partial^M V^N - \partial^N V^M) W_N + p(\partial_N V^N) W^M \\
m \cdot W &= \begin{pmatrix} a & 0 \\ -\omega & -a^T \end{pmatrix} \begin{pmatrix} w \\ \zeta \end{pmatrix} - p \text{tra} \begin{pmatrix} w \\ \zeta \end{pmatrix} \\
L_V \alpha^{M_1 \dots M_n} &= V^N \partial_N \alpha^{M_1 \dots M_n} + (\partial^{M_1} V^N - \partial^N V^{M_1}) \alpha_N^{M_2 \dots M_n} \\
&+ \dots + (\partial^{M_n} V^N - \partial^N V^{M_n}) \alpha^{M_1 \dots M_{n-1}}_N + p(\partial_N V^N) W^M \\
L_V \Psi &= V^N \partial_N \Psi + \frac{1}{4} (\partial_M V_N - \partial_N V_M) \Gamma^{MN} \Psi + p(\partial_M V^M) \Psi \\
[[V, W]] &= \frac{1}{2} (L_V W - L_W V) \\
&= [v, w] + \mathcal{L}_v \zeta - \mathcal{L}_w \lambda - \frac{1}{2} d(i_v \zeta - i_w \lambda), \\
[[U, V]]^M &= U^N \partial_N V^M - V^N \partial_N U^M - \frac{1}{2} (U_N \partial^M V^N - V_N \partial^M U^N). \\
(d\Psi)_{(i)} &= \frac{1}{2} \Gamma^M \partial_M \Psi_{(i)} = d\Psi_{(i)}, \\
D_M W^A &= \partial_M W^A + \tilde{\Omega}_M{}^A{}_B W^B \\
\tilde{\Omega}_M{}^A{}_B &= \Omega_M{}^A{}_B - \Lambda_M \delta^A{}_B, \\
\Omega_M^{AB} &= -\Omega_M{}^{BA}. \\
D_M \alpha^{A_1 \dots A_n} &= \partial_M \alpha^{A_1 \dots A_n} + \Omega_M{}^{A_1}{}_B \alpha^{BA_2 \dots A_n} \\
&+ \dots + \Omega_M{}^{A_n}{}_B \alpha^{A_1 \dots A_{n-1}B} - p \Lambda_M \alpha^{A_1 \dots A_n}. \\
D_M \Psi &= \left(\partial_M + \frac{1}{4} \Omega_M^{AB} \Gamma_{AB} - p \Lambda_M \right) \Psi \\
W &= W^A \hat{E}_A = w^a \hat{E}_a + \zeta_a E^a \\
(D_M^\nabla W^A) \hat{E}_A &= \begin{cases} (\nabla_\mu w^a) \hat{E}_a + (\nabla_\mu \zeta_a) E^a & \text{para } M = \mu \\ 0 & \text{para } M = \mu + d \end{cases}
\end{aligned}$$

$$\begin{aligned}
T(V) \cdot \alpha &= L_V^D \alpha - L_V \alpha \\
T_{ABC} &= -3\tilde{\Omega}_{[ABC]} + \tilde{\Omega}_D{}^D{}_B \eta_{AC} - \Phi^{-2} \langle \hat{E}_A, L_{\Phi^{-1}\hat{E}_B} \hat{E}_C \rangle \\
T &\in \Gamma(\Lambda^3 E \oplus E) \\
T^M{}_{PN} &= (T_1)^M{}_{PN} - (T_2)_P \delta^M{}_N, \\
(T_1)_{MNP} &= -3\tilde{\Omega}_{[MNP]} = -3\Omega_{[MNP]} \\
(T_2)_M &= -\tilde{\Omega}_Q{}_M = \Lambda_M - \Omega_Q{}_M \\
\Gamma^M D_M \Psi &= \Gamma^M \left(\partial_M \Psi + \frac{1}{4} \Omega_{MNP} \Gamma^{NP} \Psi - \frac{1}{2} \Lambda_M \Psi \right) \\
&= \Gamma^M \partial_M \Psi + \frac{1}{4} \Omega_{[MNP]} \Gamma^{MNP} \Psi - \frac{1}{2} (\Lambda_M - \Omega_N{}_M) \Gamma^M \Psi \\
&= 2 \, d\Psi - \frac{1}{12} (T_1)_{[MNP]} \Gamma^{MNP} \Psi - \frac{1}{2} (T_2)_M \Gamma^M \Psi. \\
\Gamma^M D_M \Psi &= 2 \, d\Psi \\
L_{\Phi^{-1}\hat{E}_A}^{\hat{E}_A} \hat{E}_B &= (L_{\Phi^{-1}\hat{E}_A} \Phi) \Phi^{-1} \hat{E}_B + \Phi \left(L_{\Phi^{-1}\hat{E}_A} (\Phi^{-1} \hat{E}_B) \right) \\
L_{\Phi^{-1}\hat{E}_A} \Phi &= \begin{cases} -e^{-2\phi} (\text{dete}) (i_{\hat{e}_a} i_{\hat{e}_b} de^b + 2i_{\hat{e}_a} d\phi) & \text{para } A = a \\ 0 & \text{para } A = a + d \end{cases} \\
L_{\Phi^{-1}\hat{E}_A} \Phi^{-1} \hat{E}_B &= \begin{pmatrix} [\hat{e}_a, \hat{e}_b] + i_{\hat{e}_a} \hat{e}_b B - i_{\hat{e}_a} i_{\hat{e}_b} H & \mathcal{L}_{\hat{e}_a} e^b \\ -\hat{e}_{\hat{e}_b} e^a & 0 \end{pmatrix}_{AB} \\
T_1 &= -4H, T_2 = -4 \, d\phi, \\
R(U, V, W) &= [D_U, D_V]W - D_{[[U, V]]}W \\
&= [D_{fU}, D_{gV}]hW - D_{[[fU, gV]]}hW \\
&= fgh([D_U, D_V]W - D_{[[U, V]]}W) - \frac{1}{2} h \langle U, V \rangle D_{(f \, dg - g \, df)} W, \\
E &= \mathcal{C}_+ \oplus \mathcal{C}_-, \\
\langle \hat{E}_a^+, \hat{E}_b^+ \rangle &= \Phi^2 \eta_{ab}, \\
\langle \hat{E}_{\bar{a}}^-, \hat{E}_{\bar{b}}^- \rangle &= -\Phi^2 \eta_{\bar{a}\bar{b}}, \\
\langle \hat{E}_a^+, \hat{E}_{\bar{a}}^- \rangle &= 0. \\
\hat{E}_A &= \begin{cases} \hat{E}_a^+ & \text{para } A = a \\ \hat{E}_{\bar{a}}^- & \text{para } A = \bar{a} + d \end{cases} \\
\langle \hat{E}_A, \hat{E}_B \rangle &= \Phi^2 \eta_{AB}, \text{ where } \eta_{AB} = \begin{pmatrix} \eta_{ab} & 0 \\ 0 & -\eta_{\bar{a}\bar{b}} \end{pmatrix} \\
\hat{E}^A &= \begin{cases} \hat{E}^{+a} & \text{para } A = a \\ -\hat{E}^{-\bar{a}} & \text{para } A = \bar{a} + d' \end{cases} \\
\hat{E}_a^+ &= e^{-2\phi} \sqrt{-g} (\hat{e}_a^+ + e_a^+ + i_{\hat{e}_a^+} B), \\
\hat{E}_{\bar{a}}^- &= e^{-2\phi} \sqrt{-g} (\hat{e}_{\bar{a}}^- - e_{\bar{a}}^- + i_{\hat{e}_{\bar{a}}^-} B), \\
\Phi &= e^{-2\phi} \sqrt{-g} \\
g &= \eta_{ab} e^{+a} \otimes e^{+b} = \eta_{\bar{a}\bar{b}} e^{-\bar{a}} \otimes e^{-\bar{b}}, \\
g(\hat{e}_a^+, \hat{e}_b^+) &= \eta_{ab}, g(\hat{e}_{\bar{a}}^-, \hat{e}_{\bar{b}}^-) = \eta_{\bar{b}\bar{a}}, \\
G &= \Phi^{-2} (\eta^{ab} \hat{E}_a^+ \otimes \hat{E}_b^+ + \eta^{\bar{a}\bar{b}} \hat{E}_{\bar{a}}^- \otimes \hat{E}_{\bar{b}}^-) \\
G_{MN} &= \frac{1}{2} \begin{pmatrix} g - Bg^{-1}B & -Bg^{-1} \\ g^{-1}B & g^{-1} \end{pmatrix}_{MN} \\
\Gamma^{(+)} &= \frac{1}{d!} \epsilon^{a_1 \dots a_d} \Gamma_{a_1} \dots \Gamma_{a_d}, \quad \Gamma^{(-)} = \frac{1}{d!} \epsilon^{\bar{a}_1 \dots \bar{a}_d} \Gamma_{\bar{a}_1} \dots \Gamma_{\bar{a}_d}. \\
\Gamma_a \cdot \Psi^{(B)} &= i_{\hat{e}_a^+} \Psi^{(B)} + e_a^+ \wedge \Psi^{(B)}, \Gamma_{\bar{a}} \cdot \Psi^{(B)} = i_{\hat{e}_{\bar{a}}^-} \Psi^{(B)} - e_{\bar{a}}^- \wedge \Psi^{(B)}
\end{aligned}$$

$$\begin{aligned}
\Gamma^{(+)}\Psi_{(n)}^{(B)} &= (-)^{[n/2]} * \Psi_{(n)}^{(B)}, & \Gamma^{(-)}\Psi_{(n)}^{(B)} &= (-)^d (-)^{[n+1/2]} * \Psi_{(n)}^{(B)} \\
\Gamma^{(-)}\Gamma^A\Gamma^{(-)-1} &= G_B^A\Gamma^B \\
DG &= 0, D\Phi = 0 \\
W &= w_+^a \hat{E}_a^+ + w_-^{\bar{a}} \hat{E}_{\bar{a}}^-, \\
D_M W^A &= \begin{cases} \partial_M w_+^a + \Omega_M^a{}_{\bar{b}} w_+^{\bar{b}} & \text{para } A = a \\ \partial_M w_-^{\bar{a}} + \Omega_M^{\bar{a}}{}_{\bar{b}} w_-^{\bar{b}} & \text{para } A = \bar{a}' \end{cases} \\
\Omega_{Mab} &= -\Omega_{Mba}, \Omega_{M\bar{a}\bar{b}} = -\Omega_{M\bar{b}\bar{a}}. \\
\nabla_\mu v^\nu &= (\partial_\mu v^a + \omega_\mu^{+a}{}_{\bar{b}} v^{\bar{b}})(\hat{e}_a^+)^\nu = (\partial_\mu v^{\bar{a}} + \omega_\mu^{-\bar{a}}{}_{\bar{b}} v^{\bar{b}})(\hat{e}_{\bar{a}}^-)^\nu. \\
D_M^\nabla W^a &= \begin{cases} \nabla_\mu w_+^a & \text{para } M = \mu \\ 0 & \text{para } M = \mu + d' \end{cases}, D_M^\nabla W^{\bar{a}} = \begin{cases} \nabla_\mu w_-^{\bar{a}} & \text{para } M = \mu \\ 0 & \text{para } M = \mu + d' \end{cases} \\
W &= w_+^a \hat{E}_a^+ + w_-^{\bar{a}} \hat{E}_{\bar{a}}^- = (w_+^a + w_-^a) \hat{E}_a^+ + (w_{+a} - w_{-a}) E^a, \\
T_1 &= -4H, T_2 = -4 d\phi \\
D_M W^A &= D_M^\nabla W^A + \Sigma_M^A{}_{\bar{B}} W^{\bar{B}} \\
\Sigma_{Mab} &= -\Sigma_{Mba}, \Sigma_{M\bar{a}\bar{b}} = -\Sigma_{M\bar{b}\bar{a}}. \\
(T_1)_{ABC} &= -4H_{ABC} - 3\Sigma_{[ABC]}, (T_2)_A = -4 d\phi_A - \Sigma_C{}^C{}_A. \\
dx^\mu &= \frac{1}{2} \Phi^{-1} (\hat{e}_a^{+\mu} \hat{E}^{+a} - \hat{e}_{\bar{a}}^{-\mu} \hat{E}^{-\bar{a}}), \\
d\phi &= \frac{1}{2} \partial_a \phi (\Phi^{-1} \hat{E}^{+a}) - \frac{1}{2} \partial_{\bar{a}} \phi (\Phi^{-1} \hat{E}^{-\bar{a}}). \\
\Lambda^3 T^* M &\hookrightarrow \Lambda^3 E \simeq \Lambda^3 C_+ \oplus (\Lambda^2 C_+ \otimes C_-) \oplus (C_+ \otimes \Lambda^2 C_-) \oplus \Lambda^3 C_-, \\
d\phi_A &= \begin{cases} \frac{1}{2} \partial_a \phi & A = a \\ \frac{1}{2} \partial_{\bar{a}} \phi & A = \bar{a} + d \end{cases}, H_{ABC} = \begin{cases} \frac{1}{8} H_{abc} & (A, B, C) = (a, b, c) \\ \frac{1}{8} H_{ab\bar{c}} & (A, B, C) = (a, b, \bar{c} + d) \\ \frac{1}{8} H_{a\bar{b}\bar{c}} & (A, B, C) = (a, \bar{b} + d, \bar{c} + d) \\ \frac{1}{8} H_{\bar{a}\bar{b}\bar{c}} & (A, B, C) = (\bar{a} + d, \bar{b} + d, \bar{c} + d) \end{cases} \\
\Sigma_{[abc]} &= -\frac{1}{6} H_{abc}, \quad \Sigma_{\bar{a}bc} = -\frac{1}{2} H_{\bar{a}bc}, \quad \Sigma_a{}^a{}_{\bar{b}} = -2\partial_b \phi, \\
\Sigma_{[\bar{a}\bar{b}\bar{c}]} &= +\frac{1}{6} H_{\bar{a}\bar{b}\bar{c}}, \quad \Sigma_{a\bar{b}\bar{c}} = +\frac{1}{2} H_{a\bar{b}\bar{c}}, \quad \Sigma_{\bar{a}}{}^{\bar{a}}{}_{\bar{b}} = -2\partial_{\bar{b}} \phi. \\
D_a w_+^b &= \nabla_a w_+^b - \frac{1}{6} H_a{}^b{}_{\bar{c}} w_+^{\bar{c}} - \frac{2}{9} (\delta_a{}^b \partial_c \phi - \eta_{ac} \partial^b \phi) w_+^c + A_a^{+b}{}_{\bar{c}} w_+^{\bar{c}}, \\
D_{\bar{a}} w_+^b &= \nabla_{\bar{a}} w_+^b - \frac{1}{2} H_{\bar{a}}{}^b{}_{\bar{c}} w_+^{\bar{c}}, \\
D_a w_-^{\bar{b}} &= \nabla_a w_-^{\bar{b}} + \frac{1}{2} H_a{}^{\bar{b}}{}_{\bar{c}} w_-^{\bar{c}}, \\
D_{\bar{a}} w_-^{\bar{b}} &= \nabla_{\bar{a}} w_-^{\bar{b}} + \frac{1}{6} H_{\bar{a}}^{\bar{b}}{}_{\bar{c}} w_-^{\bar{c}} - \frac{2}{9} (\delta_{\bar{a}}{}^{\bar{b}} \partial_{\bar{c}} \phi - \eta_{\bar{a}\bar{c}} \partial^{\bar{b}} \phi) w_-^{\bar{c}} + A_{\bar{a}}^{-\bar{b}}{}_{\bar{c}} w_-^{\bar{c}}, \\
A_{abc}^+ &= -A_{acb}^+, \quad A_{[abc]}^+ = 0, \quad A_a^{+a}{}_{\bar{b}} = 0, \\
A_{\bar{a}\bar{b}\bar{c}}^- &= -A_{\bar{a}\bar{c}\bar{b}}^-, \quad A_{[\bar{a}\bar{b}\bar{c}]}^- = 0, \quad A_{\bar{a}}^{-\bar{a}}{}_{\bar{b}} = 0 \\
D_{\bar{a}} w_+^b &= \nabla_{\bar{a}} w_+^b - \frac{1}{2} H_{\bar{a}}{}^b{}_{\bar{c}} w_+^{\bar{c}}, \\
D_a w_-^{\bar{b}} &= \nabla_a w_-^{\bar{b}} + \frac{1}{2} H_a{}^{\bar{b}}{}_{\bar{c}} w_-^{\bar{c}},
\end{aligned}$$

$$\begin{aligned}
D_a w_+^a &= \nabla_a w_+^a - 2(\partial_a \phi) w_+^a, \\
D_{\bar{a}} w_-^{\bar{a}} &= \nabla_{\bar{a}} w_-^{\bar{a}} - 2(\partial_{\bar{a}} \phi) w_-^{\bar{a}}, \\
D_M \epsilon^+ &= \partial_M \epsilon^+ + \frac{1}{4} \Omega_M^{ab} \gamma_{ab} \epsilon^+, \\
D_M \epsilon^- &= \partial_M \epsilon^- + \frac{1}{4} \Omega_M^{\bar{a}\bar{b}} \gamma_{\bar{a}\bar{b}} \epsilon^-. \\
\\
D_{\bar{a}} \epsilon^+ &= \left(\nabla_{\bar{a}} - \frac{1}{8} H_{\bar{a}bc} \gamma^{bc} \right) \epsilon^+, \\
D_a \epsilon^- &= \left(\nabla_a + \frac{1}{8} H_{a\bar{b}\bar{c}} \gamma^{\bar{b}\bar{c}} \right) \epsilon^-, \\
\gamma^a D_a \epsilon^+ &= \left(\gamma^a \nabla_a - \frac{1}{24} H_{abc} \gamma^{abc} - \gamma^a \partial_a \phi \right) \epsilon^+, \\
\gamma^{\bar{a}} D_{\bar{a}} \epsilon^- &= \left(\gamma^{\bar{a}} \nabla_{\bar{a}} + \frac{1}{24} H_{\bar{a}\bar{b}\bar{c}} \gamma^{\bar{a}\bar{b}\bar{c}} - \gamma^{\bar{a}} \partial_{\bar{a}} \phi \right) \epsilon^-. \\
R_{a\bar{b}} w_+^a &= [D_a, D_{\bar{b}}] w_+^a \\
R_{\bar{a}b} w_-^{\bar{a}} &= [D_{\bar{a}}, D_b] w_-^{\bar{a}} \\
\frac{1}{2} R_{a\bar{b}} \gamma^a \epsilon^+ &= [\gamma^a D_a, D_{\bar{b}}] \epsilon^+, \\
\frac{1}{2} R_{\bar{a}b} \gamma^{\bar{a}} \epsilon^- &= [\gamma^{\bar{a}} D_{\bar{a}}, D_b] \epsilon^-. \\
-\frac{1}{4} S \epsilon^+ &= (\gamma^a D_a \gamma^b D_b - D^{\bar{a}} D_{\bar{a}}) \epsilon^+ \\
-\frac{1}{4} S \epsilon^- &= (\gamma^{\bar{a}} D_{\bar{a}} \gamma^{\bar{b}} D_{\bar{b}} - D^a D_a) \epsilon^- \\
R_{ab} &= \mathcal{R}_{ab} - \frac{1}{4} H_{acd} H_b^{cd} + 2 \nabla_a \nabla_b \phi + \frac{1}{2} e^{2\phi} \nabla^c (e^{-2\phi} H_{cab}), \\
S &= \mathcal{R} + 4 \nabla^2 \phi - 4 (\partial \phi)^2 - \frac{1}{12} H^2 \\
ds_{10}^2 &= ds^2(\mathbb{R}^{9-d,1}) + ds_d^2 \\
\{g, B, \phi\} &\in \frac{O(10,10)}{O(9,1) \times O(1,9)} \times \mathbb{R}^+ \\
\delta_V G &= L_V G, \delta_V \Phi = L_V \Phi \\
\psi_{\bar{a}}^+ &\in \Gamma(C_- \otimes S^+(C_+)), \psi_{\bar{a}}^- \in \Gamma(C_+ \otimes S^+(C_-)), \\
\rho^+ &\in \Gamma(S^+(C_+)), \rho^- \in \Gamma(S^+(C_-)) \\
\epsilon^+ &\in \Gamma(S^+(C_+)), \epsilon^- \in \Gamma(S^+(C_-)) \\
F &\in \Gamma(S_{(1/2)}^\pm), \\
F^{(B)} &= e^{B(i)} \wedge F_{(i)} = e^{B(i)} \wedge \sum_n dA_{(i)}^{(n-1)}. \\
\Gamma^A &= \begin{cases} \gamma^a \otimes 1 & \text{para } A = a \\ \gamma^{(10)} \otimes \gamma^{\bar{a}} \gamma^{(10)} & \text{para } A = \bar{a} + d \end{cases} \\
S_{(1/2)} &\simeq S(C_+) \otimes S(C_-) \\
F_\# &: S(C_-) \rightarrow S(C_+). \\
F_\#^T &= (C F_\# C^{-1})^T, \\
\mathbb{F}^{(B, \phi)} &= \sum_n \frac{1}{n!} F_{a_1 \dots a_n}^{(B, \phi)} \gamma^{a_1 \dots a_n}.
\end{aligned}$$

$$\begin{aligned}
F_{(i)} &= e^{-B(i)} \wedge F^{(B)} = e^{-\phi} e^{-B(i)} \wedge F^{(B, \phi)} \\
&= e^{-\phi} e^{-B(i)} \wedge \sum_n \left[\frac{1}{32(n!)} (-)^{[n/2]} \text{tr}(\gamma_{(n)} (\Lambda^+)^{-1} F_{\#} \Lambda^-) \right]. \\
\Gamma^{(-)} F &= -F, \\
\delta \psi_{\bar{a}}^+ &= D_{\bar{a}} \epsilon^+ + \frac{1}{16} F_{\#} \gamma_{\bar{a}} \epsilon^-, \\
\delta \psi_a^- &= D_a \epsilon^- + \frac{1}{16} F_{\#}^T \gamma_a \epsilon^+ \\
\delta \rho^+ &= \gamma^a D_a \epsilon^+, \\
\delta \rho^- &= \gamma^{\bar{a}} D_{\bar{a}} \epsilon^-, \\
\delta \hat{E}_a^+ &= (\delta \log \Phi) \hat{E}_a^+ - (\delta \Lambda_{a\bar{b}}^+) \hat{E}^{-\bar{b}}, \\
\delta \hat{E}_{\bar{a}}^- &= (\delta \log \Phi) \hat{E}_{\bar{a}}^- - (\delta \Lambda_{\bar{a}b}^-) \hat{E}^{+b}, \\
\delta \Lambda_{a\bar{a}}^+ &= \bar{\epsilon}^+ \gamma_a \psi_{\bar{a}}^+ + \bar{\epsilon}^- \gamma_{\bar{a}} \psi_a^-, \\
\delta \Lambda_{\bar{a}a}^- &= \bar{\epsilon}^+ \gamma_a \psi_{\bar{a}}^+ + \bar{\epsilon}^- \gamma_{\bar{a}} \psi_a^-, \\
\delta \log \Phi &= -2\delta\phi + \frac{1}{2} \delta \log(-g) = \bar{\epsilon}^+ \rho^+ + \bar{\epsilon}^- \rho^- \\
\tilde{\delta} e_{\mu}^{+a} &= \bar{\epsilon}^+ \gamma_{\mu} \psi^{+a} + \bar{\epsilon}^- \gamma^a \psi_{\mu}^-, \\
\tilde{\delta} e_{\mu}^{-\bar{a}} &= \bar{\epsilon}^+ \gamma^{\bar{a}} \psi_{\mu}^+ + \bar{\epsilon}^- \gamma_{\mu} \psi^{-\bar{a}}, \\
\tilde{\delta} g_{\mu\nu} &= 2\bar{\epsilon}^+ \gamma_{(\mu} \psi_{\nu)}^+ + 2\bar{\epsilon}^- \gamma_{(\mu} \psi_{\nu)}^-, \\
\tilde{\delta} e_{\mu}^{+a} &= \delta e_{\mu}^{+a} - (\bar{\epsilon}^+ \gamma^a \psi^{+b} - \bar{\epsilon}^+ \gamma^b \psi^{+a}) e_{\mu b}^+, \\
\tilde{\delta} e_{\mu}^{-\bar{a}} &= \delta e_{\mu}^{-\bar{a}} - (\bar{\epsilon}^- \gamma^{\bar{a}} \psi^{-\bar{b}} - \bar{\epsilon}^- \gamma^{\bar{b}} \psi^{-\bar{a}}) e_{\mu \bar{b}}^-. \\
\delta G_{a\bar{a}} &= \delta G_{\bar{a}a} = 2(\bar{\epsilon}^+ \gamma_a \psi_{\bar{a}}^+ + \bar{\epsilon}^- \gamma_{\bar{a}} \psi_a^-) \\
\frac{1}{16} (\delta A_{\#}) &= (\gamma^a \epsilon^+ \bar{\psi}_{\bar{a}}^- - \rho^+ \bar{\epsilon}^-) \mp (\psi_{\bar{a}}^+ \bar{\epsilon}^- \gamma^{\bar{a}} + \epsilon^+ \bar{\rho}^-) \\
R_{a\bar{b}} + \frac{1}{16} \Phi^{-1} \langle F, \Gamma_{a\bar{b}} F \rangle &= 0 \\
S &= 0 \\
\frac{1}{2} \Gamma^A D_A F &= dF = 0 \\
S_B &= \frac{1}{2\kappa^2} \int \left(\Phi S + \frac{1}{4} \langle F, \Gamma^{(-)} F \rangle \right) \\
S_F &= -\frac{1}{2\kappa^2} \int \left[2\Phi [\bar{\psi}^{+\bar{a}} \gamma^b D_b \psi_{\bar{a}}^+ + \bar{\psi}^{-a} \gamma^{\bar{b}} D_{\bar{b}} \psi_a^- \right. \\
&\quad + 2\bar{\rho}^+ D_{\bar{a}} \psi^{+\bar{a}} + 2\bar{\rho}^- D_a \psi^{-a} \\
&\quad - \bar{\rho}^+ \gamma^a D_a \rho^+ - \bar{\rho}^- \gamma^{\bar{a}} D_{\bar{a}} \rho^- \\
&\quad \left. - \frac{1}{8} (\bar{\rho}^+ F_{\#} \rho^- + \bar{\psi}_{\bar{a}}^+ \gamma^a F_{\#} \gamma^{\bar{a}} \psi_a^-) \right]. \\
\gamma^b D_b \psi_{\bar{a}}^+ - D_{\bar{a}} \rho^+ &= +\frac{1}{16} \gamma^b F_{\#} \gamma_{\bar{a}} \psi_b^-, \\
\gamma^{\bar{b}} D_{\bar{b}} \psi_a^- - D_a \rho^- &= +\frac{1}{16} \gamma^{\bar{b}} F_{\#}^T \gamma_a \psi_{\bar{b}}^+, \\
\gamma^a D_a \rho^+ - D^{\bar{a}} \psi_{\bar{a}}^+ &= -\frac{1}{16} F_{\#} \rho^-, \\
\gamma^{\bar{a}} D_{\bar{a}} \rho^- - D^a \psi_a^- &= -\frac{1}{16} F_{\#}^T \rho^+,
\end{aligned}$$

$$\begin{aligned}
\omega_{(k)} &= \frac{1}{k!} \omega_{\mu_1 \dots \mu_k} dx^{\mu_1} \wedge \dots \wedge dx^{\mu_k} \\
\omega_{(k)} \wedge \eta_{(l)} &= \frac{1}{(k+l)!} \left(\frac{(k+l)!}{k! l!} \omega_{[\mu_1 \dots \mu_k} \eta_{\mu_{k+1} \dots \mu_{k+l}]} \right) dx^{\mu_1} \wedge \dots \wedge dx^{\mu_{k+l}} \\
* \omega_{(k)} &= \frac{1}{(10-k)!} \left(\frac{1}{k!} \sqrt{-g} \epsilon_{\mu_1 \dots \mu_{10-k} \nu_1 \dots \nu_k} \omega^{\nu_1 \dots \nu_k} \right) dx^{\mu_1} \wedge \dots \wedge dx^{\mu_{10-k}} \\
&\quad \{ \gamma^\mu, \gamma^\nu \} = 2g^{\mu\nu}, \gamma^{\mu_1 \dots \mu_k} = \gamma^{[\mu_1} \dots \gamma^{\mu_k]} \\
&\quad C \gamma^\mu C^{-1} = -(\gamma^\mu)^T, C^T = -C \\
&\quad C \gamma^{\mu_1 \dots \mu_k} C^{-1} = (-)^{[(k+1)/2]} (\gamma^{\mu_1 \dots \mu_k})^T, \\
&\quad \bar{\epsilon} \gamma^{\mu_1 \dots \mu_k} \chi = (-)^{[(k+1)/2]} \bar{\chi} \gamma^{\mu_1 \dots \mu_k} \epsilon, \\
\gamma^{(10)} &= \gamma^0 \gamma^1 \dots \gamma^9 = \frac{1}{10!} \epsilon_{\mu_1 \dots \mu_{10}} \gamma^{\mu_1 \dots \mu_{10}} \\
\gamma_{\mu_1 \dots \mu_k} \gamma^{(10)} &= (-)^{[k/2]} \frac{1}{(10-k)!} \sqrt{-g} \epsilon_{\mu_1 \dots \mu_k \nu_1 \dots \nu_{10-k}} \gamma^{\nu_1 \dots \nu_{10-k}}, \\
\gamma^{(k)} \gamma^{(10)} &= (-)^{[k/2]} * \gamma^{(10-k)} \\
\Psi &= \sum_k \frac{1}{k!} \Psi_{\mu_1 \dots \mu_k} \gamma^{\mu_1 \dots \mu_k} \\
F &= \{ (x, \{\hat{e}_a\}) : x \in M \text{ and } \{\hat{e}_a\} T_x M \}. \\
v^a \mapsto v'^a &= A^a{}_b v^b, \hat{e}_a \mapsto \hat{e}'_a = \hat{e}_b (A^{-1})^b{}_a. \\
\mathcal{L}_v w &= -\mathcal{L}_w v = [v, w], \\
\mathcal{L}_v \alpha_{\nu_1 \dots \nu_q}^{\mu_1 \dots \mu_p} &= v^\mu \partial_\mu \alpha_{\nu_1 \dots \nu_q}^{\mu_1 \dots \mu_p} \\
&\quad + (\partial_\mu v^{\mu_1}) \alpha_{\nu_1 \dots \nu_q}^{\mu_2 \dots \mu_p} + \dots + (\partial_\mu v^{\mu_p}) \alpha_{\nu_1 \dots \nu_q}^{\mu_1 \dots \mu_{p-1} \mu} \\
&\quad - (\partial_{\nu_1} v^\mu) \alpha_{\mu \nu_2 \dots \nu_q}^{\mu_1 \dots \mu_p} - \dots - (\partial_{\nu_q} v^\mu) \alpha_{\nu_1 \dots \nu_{q-1} \mu}^{\mu_1 \dots \mu_p}. \\
T(v, w) &= \nabla_v w - \nabla_w v - [v, w]. \\
T^\mu{}_{\nu\lambda} &= \omega_\nu{}^\mu{}_\lambda - \omega_\lambda{}^\mu{}_\nu, \\
T^a{}_{bc} &= \omega_b{}^a{}_c - \omega_c{}^a{}_b + [\hat{e}_b, \hat{e}_c]^a. \\
(i_v T) \alpha &= \mathcal{L}_v^\nabla \alpha - \mathcal{L}_v \alpha, \\
\mathcal{R}(u, v) w &= [\nabla_u, \nabla_v] w - \nabla_{[u, v]} w, \\
\mathcal{R}_{\mu\nu}{}^\lambda{}_\rho v^\rho &= [\nabla_\mu, \nabla_\nu] v^\lambda - T^\rho{}_{\mu\nu} \nabla_\rho v^\lambda. \\
\mathcal{R}_{\mu\nu} &= \mathcal{R}_{\lambda\mu\nu}{}^\lambda. \\
\mathcal{R} &= g^{\mu\nu} \mathcal{R}_{\mu\nu} \\
P &= \{ (x, \{\hat{e}_a\}) \in F : g(\hat{e}_a, \hat{e}_b) = \delta_{ab} \}, \\
g|_x &\in GL(d, \mathbb{R}) / O(d) \\
\nabla_{\partial/\partial x^\mu} \hat{e}_a &= \omega_\mu{}^b{}_a \hat{e}_b.
\end{aligned}$$

$$\begin{aligned}
\nabla_\mu A^\alpha{}_\beta &:= \partial_\mu A^\alpha{}_\beta - \Gamma_{\beta\mu}^\rho A^\alpha{}_\rho + \Gamma_{\rho\mu}^\alpha A^\rho{}_\beta. \\
\Gamma_{\mu\nu}^\rho &:= \left\{ \begin{smallmatrix} \rho \\ \mu\nu \end{smallmatrix} \right\} + K_{\mu\nu}^\rho + L^\rho{}_{\mu\nu}, \\
\left\{ \begin{smallmatrix} \rho \\ \mu\nu \end{smallmatrix} \right\} &:= \frac{1}{2} g^{\rho\lambda} (\partial_\mu g_{\lambda\nu} + \partial_\nu g_{\mu\lambda} - \partial_\lambda g_{\mu\nu}), \\
K^\rho{}_{\mu\nu} &:= \frac{1}{2} (T_\mu{}^\rho{}_\nu + T_\nu{}^\rho{}_\mu - T^\rho{}_{\mu\nu}), \\
L^\rho{}_{\mu\nu} &:= \frac{1}{2} (Q^\rho{}_{\mu\nu} - Q_\mu{}^\rho{}_\nu - Q_\nu{}^\rho{}_\mu). \\
R_{\nu\rho\sigma}^\mu &:= \partial_\rho \Gamma_{\nu\sigma}^\mu - \partial_\sigma \Gamma_{\nu\rho}^\mu + \Gamma_{\tau\rho}^\mu \Gamma_{\nu\sigma}^\tau - \Gamma_{\tau\sigma}^\mu \Gamma_{\nu\rho}^\tau, \\
T_{\nu\rho}^\mu &:= 2\Gamma_{[\rho\nu]}^\mu \equiv \Gamma_{\rho\nu}^\mu - \Gamma_{\nu\rho}^\mu, \\
Q_{\mu\nu\rho} &:= \nabla_\mu g_{\nu\rho} \equiv \partial_\mu g_{\nu\rho} - 2\Gamma_{(\nu|\mu}^\lambda g_{\rho)\lambda} \neq 0.
\end{aligned}$$

$$\begin{aligned}
R_{\nu\rho\sigma}^\mu &= -R_{\nu\sigma\rho}^\mu, \\
T_{\nu\rho}^\mu &= -T_{\rho\nu}^\mu, \\
Q_{\mu\nu\rho} &= Q_{\mu\rho\nu}.
\end{aligned}$$

$$\begin{aligned}
R_{[\nu\rho\sigma]}^\mu &= \nabla_{[\nu} T_{\rho\sigma]}^\mu + T_{\alpha[\nu}^\mu T_{\rho\sigma]}^\alpha, \\
\nabla_{[\alpha} R_{\nu|\rho\sigma]}^\mu &= -R_{\nu\tau[\alpha}^\mu T_{\rho\sigma]}^\tau
\end{aligned}$$

$$\partial_\mu := \left(\frac{\partial}{\partial x^\mu} \right)_p,$$

$$dx^\mu \partial_\nu = \delta_\nu^\mu.$$

$$e_A := e_A^\mu \partial_\mu, e^A := e_\mu^A dx^\mu,$$

$$g_{\mu\nu} = \eta_{AB} e^A e^B, \eta_{AB} = g_{\mu\nu} e_A^\mu e_B^\nu.$$

$$\begin{aligned}
[e_A, e_B] &:= e_A e_B - e_B e_A \\
&= (e_A^\mu \partial_\mu)(e_B^\nu \partial_\nu) - (e_B^\nu \partial_\nu)(e_A^\mu \partial_\mu) \\
&= [e_A^\mu e_\nu^C (\partial_\mu e_B^\nu) - e_B^\nu e_\mu^C (\partial_\nu e_A^\mu)] e_C \\
&= e_A^\mu e_B^\nu [\partial_\nu e_\mu^C - \partial_\mu e_\nu^C] e_C \\
&= f_{AB}^C e_C,
\end{aligned}$$

$$f_{AB}^C := e_A^\mu e_B^\nu [\partial_\nu e_\mu^C - \partial_\mu e_\nu^C]$$

$$d\omega = \partial_\mu \omega_\nu dx^\mu \wedge dx^\nu$$

$$dx^\mu \wedge dx^\nu = dx^\mu \otimes dx^\nu - dx^\nu \otimes dx^\mu$$

$$d\omega(u, v) = u\omega(v) - v\omega(u) - \omega([u, v]_{\mathcal{L}})$$

$$d\omega(u, v) := \partial_\mu \omega_\nu (u^\mu v^\nu - u^\nu v^\mu),$$

$$u\omega(v) := u^\mu v^\nu \partial_\mu \omega_\nu + u^\mu \omega_\nu \partial_\mu v^\nu,$$

$$\omega([u, v]_{\mathcal{L}}) := \omega_\nu (u^\mu \partial_\mu v^\nu - v^\mu \partial_\mu u^\nu),$$

$$\{de^C(e_A, e_B)\}e_C = \{e_A[e^C(e_B)] - e_B[e^C(e_A)]$$

$$-e^C([e_A, e_B]_{\mathcal{L}})\}e_C$$

$$= -e^C([e_A, e_B]_{\mathcal{L}}^L e_L)e_C$$

$$= -[e_A, e_B]_{\mathcal{L}}.$$

$$\nabla_{e_A} e_B = \gamma_{AB}^C e_C,$$

$$\gamma_{\lambda\nu\mu} := e_\mu^A e_\lambda^B \nabla_A (e_\nu)_B$$

$$= -e_\mu^A (e_\nu)_B \nabla_A e_\lambda^B$$

$$= -e_\mu^A e_\nu^B \nabla_A (e_\lambda)_B = -\gamma_{\nu\lambda\mu},$$

$$\gamma_{AB}^C = \omega_B^C(e_A) \Leftrightarrow \omega_B^C = \gamma_{AB}^C e^A$$

$$[\nabla_\mu, \partial_\nu] = \nabla_\mu \partial_\nu - \nabla_\nu \partial_\mu$$

$$= (\Gamma_{\mu\nu}^\lambda - \Gamma_{\nu\mu}^\lambda) \partial_\lambda$$

$$= T_{\mu\nu}^\lambda \partial_\lambda,$$

$$T(v, u) := \nabla_v u - \nabla_u v - [v, u]_{\mathcal{L}}.$$

$$T(e_A, e_B) = \nabla_{e_A} e_B - \nabla_{e_B} e_A - [e_A, e_B]_{\mathcal{L}}$$

$$= [\omega_B^C(e_A) - \omega_A^C(e_B) + de^C(e_A, e_B)]e_C$$

$$= [(\omega_D^C \wedge e^D + de^C)(e_A, e_B)]e_C.$$

$$T = \Omega^C \otimes e_C$$



$$\begin{aligned}
de^C &:= -\omega_A^C \wedge e^A \\
&= -\frac{1}{2}(\gamma_{AB}^C - \gamma_{BA}^C)e^A \wedge e^B \\
&= -\frac{1}{2}e_A^\mu e_B^\nu (\partial_\nu e_\mu^C - \partial_\mu e_\nu^C)e^A \wedge e^B \\
&= -\frac{1}{2}f_{AB}^C e^A \wedge e^B,
\end{aligned}$$

$$\eta_{AB} = \eta_{\mu\nu} e_A^\mu e_B^\nu.$$

$$\Lambda_v^\mu: x^\mu \longrightarrow x'^\mu = \Lambda_v^\mu(x) x^\nu,$$

$$\eta_{\mu\nu} x^\mu x^\nu = -t^2 + x^2 + y^2 + z^2.$$

$$\Lambda^\alpha{}_\beta = \mathcal{G} \cdot \left[\begin{array}{cc} \gamma & -\gamma \mathcal{R}^i{}_j \frac{v^j}{c} \\ -\gamma \mathcal{R}^i{}_j \frac{v^j}{c} & \mathcal{R}^i{}_j \left(\delta_j^i + (\gamma - 1) \frac{v^i v^j}{v^2} \right) \end{array} \right],$$

$$\mathbb{1}:=\mathrm{diag}(1,1,1,1)$$

$$\mathbb{P}:=\mathrm{diag}(1,-1,-1,-1)$$

$$\mathbb{T}:=\mathrm{diag}(-1,1,1,1)$$

$$\bar{e}^A{}_\mu = \Lambda^A_B e^B{}_\mu,$$

$$g_{\mu\nu} = \eta_{AB} \bar{e}^A{}_\mu \bar{e}^B{}_\nu \quad \eta_{AB} = g_{\mu\nu} \bar{e}^\mu{}_A \bar{e}^\nu{}_B.$$

$$\Lambda^\alpha{}_\beta = \delta^\alpha_\beta + \omega^\alpha{}_\beta + \mathcal{O}\left[\left(\omega^\alpha{}_\beta\right)^2\right].$$

$$(J_{AB})_D^C:=2i\eta_{[B|D}\delta_{A]}^C=i(\eta_{BD}\delta_A^C-\eta_{AD}\delta_B^C).$$

$$\Lambda = e^{\frac{i}{2}\omega_{AB}J^{AB}}.$$

$$\omega_\mu\colon J_{AB}\in\mathfrak{L}\longrightarrow \omega_\mu\colon=\frac{1}{2}\omega^{AB}{}_\mu J_{AB},$$

$$\mathcal{D}_\mu\colon=\partial_\mu-\omega_\mu=\partial_\mu-\frac{i}{2}\omega^{AB}{}_\mu J_{AB},$$

$$\begin{aligned}
\mathcal{D}_\mu e^C &= \partial_\mu e^C - \frac{i}{2} \omega^{AB}{}_\mu [i(\eta_{BD} \delta_A^C - \eta_{AD} \delta_B^C)] e^D \\
&= \partial_\mu e^C + \frac{1}{2} [\omega^A{}_{D\mu} \delta_A^C + \omega^B{}_{D\mu} \delta_B^C] e^D \\
&= \partial_\mu e^C + \omega^C{}_{D\mu} e^D. \\
\mathcal{D}_\mu (e_\lambda^C dx^\lambda) &= \mathcal{D}_\mu (e_\lambda^C) dx^\lambda + e_\lambda^C \mathcal{D}_\mu (dx^\lambda) \\
&= \mathcal{D}_\mu (e_\lambda^C) dx^\lambda + e_\lambda^C (\delta_\mu^\lambda + e_E^\lambda e_\mu^D \omega^E{}_{D\rho} dx^\rho) \\
&= \mathcal{D}_\mu (e_\lambda^C) dx^\lambda + e_\mu^C, \\
\mathcal{D}_\mu (e_\lambda^C dx^\lambda) &= \partial_\mu (e_\lambda^C dx^\lambda) + \omega^C{}_{D\mu} e_\lambda^D dx^\lambda \\
&= \partial_\mu (e_\lambda^C) dx^\lambda + e_\mu^C + \omega^C{}_{D\mu} e_\lambda^D dx^\lambda.
\end{aligned}$$

$$\mathcal{D}_\mu(e_\lambda^C) = \partial_\mu(e_\lambda^C) + \omega^C{}_{D\mu} e_\lambda^D.$$

$$\tilde{\nabla}_\mu X_B^A\colon=\partial_\mu+\omega^A{}_{C\mu}X^C{}_B-\omega^C{}_{B\mu}X_C^A.$$

$$\begin{aligned}
\nabla V &= (\nabla_\mu V^\nu) \mathrm{d}x^\mu \otimes \partial_\nu \\
&= (\partial_\mu V + \Gamma_{\mu\lambda}^\nu V^\lambda) \mathrm{d}x^\mu \otimes \partial_\nu.
\end{aligned}$$

$$\begin{aligned}
\tilde{\nabla} V &= (\tilde{\nabla}_\mu V^A) dx^\mu \otimes e_A \\
&= (\partial_\mu V^A + \omega^A_{B\mu} V^B) dx^\mu \otimes e_A \\
&= [\partial_\mu (e_\lambda^A V^\lambda) + \omega^A_{B\mu} e_\lambda^B V^\lambda] dx^\mu \otimes (e_A^v \partial_v) \\
&= [\partial_\mu V^v + (e_A^v \partial_\mu e_\lambda^A + \omega^A_{B\mu} e_A^v e_\lambda^B) V^\lambda] dx^\mu \otimes \partial_v \\
&= [\partial_\mu V^v + (e_A^v \mathcal{D}_\mu e_\lambda^A) V^\lambda] dx^\mu \otimes \partial_v.
\end{aligned}$$

$$\begin{aligned}
\Gamma_{\mu\nu}^\lambda &\equiv e_A^\lambda \mathcal{D}_\mu e^A_\nu \\
\omega_{B\mu}^A &= e_\lambda^A e_B^v \Gamma_{\mu\nu}^\lambda + e_\sigma^A \partial_\mu e_B^\sigma \equiv e_\nu^A \nabla_\mu e_B^v \\
\omega^{AB} &= \omega^{AB}{}_\mu dx^\mu,
\end{aligned}$$

$$\begin{aligned}
\nabla_\mu e_\nu^A &= \partial_\mu e_\nu^A - \Gamma_{\mu\nu}^\lambda e_\lambda^A + \omega_{B\mu}^A e_\nu^B = 0; \\
0 &= \nabla_\lambda g_{\mu\nu} = \partial_\lambda g_{\mu\nu} - \Gamma_{\lambda\mu}^\sigma g_{\sigma\nu} - \Gamma_{\lambda\nu}^\sigma g_{\mu\sigma} \\
&= \partial_\lambda (e_\mu^A e_\nu^B \eta_{AB}) - e_A^\sigma g_{\sigma\nu} \mathcal{D}_\lambda e_\mu^A - e_A^\sigma g_{\mu\sigma} \mathcal{D}_\lambda e_\nu^A \\
&= -e_\nu^A e_\mu^D (\omega_{AD\lambda} - \omega_{DA\lambda}), \\
\partial_\mu x'^A &= \partial_\mu (\Lambda_B^A(x) x^B) \\
&= (\partial_\mu x^B) \Lambda_B^A(x) + x^B (\partial_\mu \Lambda_B^A(x)), \\
\partial_\mu x'^A &= e'^A{}_\mu \partial'_C x'^A = e'^A{}_\mu = e_\mu^C \Lambda_C^A(x). \\
e_\mu^A &= \partial_\mu x^A + \dot{\omega}^A_{B\mu} x^B \equiv \mathcal{D}_\mu x^A, \\
\dot{\omega}^A_{B\mu} &:= \Lambda_C^A(x) \partial_\mu \Lambda_B^C(x)
\end{aligned}$$

$$\omega^A_{B\mu} = \underbrace{\Lambda_C^A(x) \omega'^C{}_{D\mu} \Lambda_C^D}_{\text{non inertial}} + \underbrace{\Lambda_C^A \partial_\mu \Lambda_B^C(x)}_{\text{inertial}}.$$

$$f_{AB}^C = \dot{\omega}_{BA}^C - \dot{\omega}_{AB}^C.$$

as

$$\dot{\omega}_{BC}^A = \frac{1}{2} (f_B{}^A{}_C + f_C{}^A{}_B - f_{BC}^A).$$

$$R_{B\mu\nu}^A = \partial_\nu \dot{\omega}_{B\mu}^A - \partial_\mu \dot{\omega}_{B\nu}^A + \dot{\omega}_{E\nu}^A \dot{\omega}_{B\mu}^E$$

—

$$\begin{aligned}
\dot{\omega}_{E\mu}^A \dot{\omega}_{B\nu}^E &\equiv 0, \\
T_{\nu\mu}^A &= \partial_\nu e_\mu^A - \partial_\mu e_\nu^A + \dot{\omega}_{E\nu}^A e_\mu^E - \dot{\omega}_{E\mu}^A e_\nu^E.
\end{aligned}$$

$$g_{\mu\nu}(\varphi(p)) = \eta_{\mu\nu}, \partial_\lambda g_{\mu\nu}(\varphi(p)) = 0.$$

$$\frac{d^2 \xi^\alpha}{ds^2} = 0,$$

$$\frac{d^2 x^\lambda}{ds^2} + \overset{\circ}{\Gamma}_{\mu\nu}^\lambda \frac{dx^\mu}{ds} \frac{dx^\nu}{ds} = 0,$$

$$\overset{\circ}{\Gamma}_{\mu\nu}^\lambda := \frac{\partial x^\lambda}{\partial \xi^\sigma} \frac{\partial^2 \xi^\sigma}{\partial x^\mu \partial x^\nu},$$

$$\left[\overset{\circ}{\nabla}_\mu, \overset{\circ}{\nabla}_\nu \right] v^\alpha = \overset{\circ}{R}_{\beta\mu\nu}^\alpha v^\beta.$$

$$\overset{\circ}{R}_{\mu\nu\alpha\beta} = -\overset{\circ}{R}_{\nu\mu\alpha\beta},$$

$$\overset{\circ}{R}_{\mu\nu\alpha\beta} = \overset{\circ}{R}_{\alpha\beta\mu\nu}.$$

$$T^\mu = \frac{\partial x^\mu}{\partial t}, S^\mu = \frac{\partial x^\mu}{\partial s}.$$

$$V^\mu = T^\nu \overset{\circ}{\nabla}_\nu T^\mu,$$

$$A^\mu = T^\nu \overset{\circ}{\nabla}_\nu V^\mu.$$



$$A^\mu = \overset{\circ}{R}{}^\mu_{\lambda\alpha\beta} T^\lambda T^\alpha S^\beta,$$

$$S_{\text{GR}} := \frac{c^4}{16\pi G} \int d^4x \sqrt{-g} (\mathcal{L}_{\text{GR}} + \mathcal{L}_{\text{m}}),$$

$$\overset{\circ}{G}_{\mu\nu} := \overset{\circ}{R}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \overset{\circ}{R} = \frac{8\pi G}{c^4} T_{\mu\nu},$$

$$T^{\mu\nu} = -\frac{1}{2\sqrt{-g}} \frac{\delta \mathcal{L}_m}{\delta g_{\mu\nu}}$$

$$\mathrm{d}e^C + \overset{\circ}{\omega}^A_B \wedge e^B = 0,$$

$$\overset{\circ}{\omega}_{AB} + \overset{\circ}{\omega}_{BA} = \mathrm{d}g_{AB},$$

$$\mathrm{d}\overset{\circ}{\omega}^A_B + \overset{\circ}{\omega}^A_C \wedge \overset{\circ}{\omega}^C{}_B = \frac{1}{2} \overset{\circ}{r}_{BCD} e^C \wedge e^D,$$

$$\overset{\circ}{\omega}^A_{B\mu} := e^A_\nu \overset{\circ}{\nabla}_\mu e^\nu_B,$$

$$\overset{\circ}{f}^A_{BC} := \dot{\gamma}^A_{BC} - \dot{\gamma}^A_{CB},$$

$$\mathrm{d}g_{AB} = \partial_C g_{AB} e^C,$$

$$\dot{\gamma}^A_{BC} = \frac{1}{2} \Big(\overset{\circ}{f}^A_{BC} - g_{CL} g^{AM} \overset{\circ}{f}^L_{BM} - g_{BL} g^{AM} \overset{\circ}{f}^L_{CM} \Big)$$

$$+\overset{\circ}{\Gamma}^A_{BC}.$$

$$\overset{\circ}{r}^A_{BCD} = \partial_D \overset{\circ}{\gamma}^A_{BC} - \partial_C \overset{\circ}{\gamma}^A_{BD} + \overset{\circ}{\gamma}^A_{CM} \overset{\circ}{\gamma}^M_{DB} \\ - \overset{\circ}{\gamma}^A_{DM} \overset{\circ}{\gamma}^M_{CB} - \overset{\circ}{\gamma}^A_{MB} \overset{\circ}{\gamma}^M_{CD}.$$

$$x^A \longrightarrow \bar{x}^A = x^A + \varepsilon^A(x^\mu),$$

$$P_A := \partial_A.$$

$$[P_A,P_B]\equiv [\partial_A,\partial_B]=0.$$

$$\delta \bar{x}^A = \varepsilon (x^\mu)^B \partial_B x^A = \varepsilon (x^\mu)^A.$$

$$\delta_\varepsilon \Psi = \epsilon^A(x^\mu) \partial_A \Psi.$$

$$\partial_\varepsilon (\partial_\mu \Psi) = \varepsilon^A \partial_A (\partial_\mu \Psi).$$

$$\partial_\varepsilon (\partial_\mu \Psi) = \underbrace{\varepsilon^A(x^\mu) \partial_A (\partial_\mu \Psi)}_{\text{correct}} + \underbrace{(\partial_\mu \varepsilon^A(x^\mu)) \partial_A \Psi}_{\text{spurious}},$$

$$e'_\mu \Psi \equiv \partial_\mu \Psi = \partial_\mu + B^A_\mu \partial_A \Psi,$$

$$\delta_\varepsilon B^A_\mu = -\partial_\mu \varepsilon^A(x^\mu).$$

$$\partial_\varepsilon (e'_\mu \Psi) = \underbrace{\varepsilon^A(x^\mu) \partial_A (\partial_\mu \Psi)}_{\text{correct}},$$

$$e_\mu = \Psi = \overset{\text{correct}}{e^A_\mu \partial_A \Psi}, e^A_\mu = \partial_\mu x^A + B^A_\mu,$$

$$B^A_\mu \longrightarrow \Lambda^A_B(x) B^B_\mu.$$

$$e_\mu \Psi = \partial_\mu + \overset{\circ}{\omega}^A_{B\mu} x^B \partial_A \Psi + B^A_\mu \partial_A \Psi,$$

$$e^A_\mu = \partial_\mu x^A + \overset{\circ}{\omega}^A_{B\mu} x^B + B^A_\mu = \dot{\mathcal{D}}_\mu x^A + B^A_\mu.$$

$$\delta_\varepsilon B^A_\mu = -\dot{\mathcal{D}}_\mu \varepsilon^A(x^\mu).$$

$$e'^A_\mu \longrightarrow e^A_\mu,$$

$$\eta_{\mu\nu} \longrightarrow g_{\mu\nu}.$$

$$\partial_\mu \Psi \rightarrow \mathcal{D}'_\mu \Psi = \partial_\mu \Psi \\ + \frac{1}{2} e'^A{}_\mu (f'^C{}_B{}_A + f'^C{}_A{}_B - f'^C{}_{BA}) S^B_C \Psi,$$



$$\partial_\mu \Psi \rightarrow \mathcal{D}_\mu \Psi = \partial_\mu \Psi + \frac{1}{2} e^A{}_\mu (f_B^C{}_A + f_A C_B^C - f_{BA}^C) S_C^B \Psi,$$

$$\underbrace{\left\{ \begin{array}{l} e'^A \rightarrow e^A \\ \partial_\mu \rightarrow \mathcal{D}_\mu \end{array} \right\}}_{\text{grav. coupling prescription in TG}} \Leftrightarrow \underbrace{\eta_{\mu\nu} \rightarrow g_{\mu\nu}}_{\text{grav. coupling prescription in GR}}$$

$$\frac{du'^A}{d\sigma} = 0,$$

$$\frac{du'^A}{d\sigma} = \underbrace{\Lambda_B^A(x) \frac{du^B}{d\sigma}}_{\text{correct}} + \underbrace{\frac{d\Lambda_B^A(x)}{d\sigma} u^B}_{\text{spurious}}.$$

$$\frac{du'^A}{d\sigma} = 0 \rightarrow \frac{du^B}{d\sigma} + \dot{\omega}_{B\mu}^A u^B u^\mu = 0.$$

$$\dot{\gamma}(\tau) := \frac{d\gamma^\mu}{d\tau} \partial_\mu.$$

$$\frac{dY^\mu}{d\tau} := \nabla_\gamma Y^\mu \equiv \frac{dY^\mu}{d\tau} + \Gamma_{\alpha\beta}^\mu Y^\alpha \frac{d\gamma^\beta}{d\tau} = 0,$$

$$\nabla_\gamma \dot{\gamma} \equiv \frac{d^2 x^\mu}{d\tau^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = 0,$$

$$\frac{d^2 x^\mu}{d\tau^2} + \overset{\circ}{\Gamma}_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = -K_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau},$$

$$\frac{d^2 x^\mu}{d\tau^2} + \dot{\Gamma}_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = -L_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau}.$$

$$\frac{d^2 x^\mu}{d\tau^2} + \Gamma_{\alpha\beta}^\mu \frac{dx^\alpha}{d\tau} \frac{dx^\beta}{d\tau} = -\left(\frac{d\lambda}{d\tau}\right)^2 \frac{d^2 \tau}{d\lambda^2} \frac{d\gamma^\mu}{d\tau}.$$

$$[e_\mu, e_\nu] = \hat{T}_{\nu\mu}^A \partial_A,$$

$$\hat{T}_{\mu\nu}^A = \partial_\nu B_\mu^A - \partial_\mu B_\nu^A + \dot{\omega}_{B\nu}^A B_\mu^B - \dot{\omega}_{B\mu}^A B_\nu^B$$

$$\dot{\mathcal{D}}_\mu (\dot{\mathcal{D}}_\nu x^A) - \dot{\mathcal{D}}_\nu (\dot{\mathcal{D}}_\mu x^A) \equiv 0$$

$$\hat{T}_{\mu\nu}^A = \partial_\nu e_\mu^A - \partial_\mu e_\nu^A + \dot{\omega}_{B\nu}^A e_\mu^B - \dot{\omega}_{B\mu}^A e_\nu^B.$$

$$\hat{T}_{\mu\nu}^\lambda = e_A^\lambda \hat{T}_{\mu\nu}^\lambda := \Gamma_{\nu\mu}^\lambda - \Gamma_{\mu\nu}^\lambda.$$

$$e_{(r)\mu}^A := \lim_{G \rightarrow 0} e_\mu^A.$$

$$\hat{T}_{BC}^A(e_\mu^A, \dot{\omega}_{B\mu}^A) = \dot{\omega}_{BC}^A - \dot{\omega}_{BC}^A - f_{BC}^A(e_{(r)}) = 0,$$

$$\dot{\omega}_{BC}^A = \frac{1}{2} e_{(r)\mu}^C [f_B^A{}_C(e_{(r)}) + f_C^A{}_B(e_{(r)}) - f_{BC}^A(e_{(r)})].$$

$$\dot{\omega}_{AB}^C - \dot{\omega}_{BA}^C = f_{AB}^C + T_{AB}^C.$$

$$\frac{1}{2} (f_B^C{}_A + f_A C_B^C - f_{BA}^C) = \dot{\omega}_{BA}^C - \hat{K}_{BA}^C,$$

$$\hat{K}_{BA}^C = \frac{1}{2} (\hat{T}_{B\ A}^C + \hat{T}_{A\ B}^C - \hat{T}_{BA}^C),$$

$$\dot{\omega}_{B\mu}^C - \hat{K}_{B\mu}^C = \overset{\circ}{\omega}_{B\mu}^C,$$

$$\hat{S}_A^{\mu\nu} := \hat{K}^{\mu\nu}{}_A - e_A{}^\nu \hat{T}^\mu + e_A{}^\mu \hat{T}^\nu,$$

$$\hat{T} := \frac{1}{2} \hat{S}_A^{\mu\nu} \hat{T}_{\mu\nu}^A$$

$$= \frac{1}{4} \hat{T}_{\mu\nu}^\rho \hat{T}_\rho^{\mu\nu} + \frac{1}{2} \hat{T}_{\mu\nu}^\rho \hat{T}_\rho^{\nu\mu} - \hat{T}_\mu \hat{T}^\mu,$$

$$\hat{R} = \overset{\circ}{R} + \hat{T} + \frac{2}{e} \partial_\mu (e \hat{T}^\mu) = 0,$$

$$\overset{\circ}{R} = -\hat{T} - \underbrace{\frac{2}{e^2 \partial_\mu (e T^\mu)}}_{\text{boundary term}}.$$

$$S_{\text{TEGR}} = -\frac{c^4}{16\pi G} \int d^4x e \underbrace{\mathcal{L}_{\text{TEGR}}}_{-\hat{T}} + \int d^4x e \mathcal{L}_m,$$

$$\hat{G}_{\mu\nu} := \frac{1}{e} \partial_\lambda (e \hat{S}_{\mu\nu}{}^\lambda) - \frac{4\pi G}{c^4} t_{\mu\nu} = \frac{4\pi G}{c^4} T_{\mu\nu},$$

$$t_{\mu\nu} = \frac{c^4}{4\pi G} \hat{S}_{\lambda\nu}{}^\rho \Gamma_{\rho\mu}^\lambda - g_{\mu\nu} \frac{c^4}{16\pi G} \hat{T}$$

$$\hat{S}_A{}^{\mu\nu} = -\frac{8\pi G}{c^4 e} \frac{\partial \mathcal{L}_{\text{TEGR}}}{\partial (\partial_\nu e^A{}_\mu)}.$$

$$\hat{G}_{\mu\nu} := \frac{1}{e} e^A{}_\mu g_{\nu\rho} \partial_\sigma (e \hat{S}_A{}^{\rho\sigma}) - \hat{S}_B{}^\sigma{}_\nu \hat{T}_{\sigma\mu}^B$$

$$\hat{T}_{\text{gen}} := -\frac{c_1}{4} \hat{T}_{\alpha\mu\nu} \hat{T}^{\alpha\mu\nu} - \frac{c_2}{2} \hat{T}_{\alpha\mu\nu} \hat{T}^{\mu\alpha\nu} + c_3 \hat{T}_\alpha \hat{T}^\alpha,$$

$$\hat{T}_{[\mu\nu]} = e_{[\mu}^A g_{\nu]\rho} \hat{T}_A{}^\rho = 0.$$

$$\mathcal{L}_{\text{TEGR}}(e_\mu^A, 0), \mathcal{L}_{\text{TEGR}}(e_\mu^A, \dot{\omega}_{B\mu}^A),$$

$$\mathcal{L}_{\text{TEGR}}(e_\mu^A, \dot{\omega}_{B\mu}^A) + \partial_\mu \left[\frac{ec^4}{8\pi G} \hat{T}^\mu(e_\mu^A, \dot{\omega}_{B\mu}^A) \right] = \mathcal{L}_{\text{TEGR}}(e_\mu^A, 0) + \partial_\mu \left[\frac{ec^4}{8\pi G} \hat{T}^\mu(e_\mu^A, 0) \right],$$

$$\hat{T}^\mu(e_\mu^A, \dot{\omega}_{B\mu}^A) = \hat{T}^\mu(e_\mu^A, 0) - \dot{\omega}^\mu.$$

$$\mathcal{L}_{\text{TEGR}}(e_\mu^A, \dot{\omega}_{B\mu}^A) = \mathcal{L}_{\text{TEGR}}(e_\mu^A, 0) + \partial_\mu \left[\frac{ec^4}{8\pi G} \dot{\omega}^\mu \right].$$

$$g_{\nu\lambda\mu}{}^{\vec{\nabla}}{}^\lambda = \vec{\nabla}_\mu v_\nu - v^\lambda \mu_{\nu\lambda}{}^Q$$

$$T^\lambda \vec{\nabla}_\lambda v \cdot w = T^\lambda v^\mu w^\nu \mathring{Q}_{\lambda\mu\nu}$$

$$T^\lambda \vec{\nabla}_\lambda \left(\frac{v \cdot w}{|v||w|} \right) \neq 0$$

$$a^\mu := u^\lambda \mathring{\nabla}_\lambda u^\mu$$

$$\tilde{a}_\mu := u^\lambda \mathring{\nabla}_\lambda u_\mu = a_\mu + \lambda v_\mu u^\lambda u^\nu$$

$$u_\mu a^\mu = u_\mu u^\lambda \mathring{\nabla}_\lambda u^\mu$$

$$= u^\lambda \mathring{\nabla} (u_\mu u^\mu) - u^\mu u^\lambda \mathring{\nabla}_\lambda u_\mu$$

$$= \mathring{Q}_{\lambda\mu\nu} u^\lambda u^\mu u^\nu + 2u_\mu a^\mu - \tilde{a}_\mu u^\mu,$$

$$a^\mu u_\mu = \tilde{a}_\mu u^\mu - \mathring{Q}_{\lambda\mu\nu} u^\lambda u^\mu u^\nu.$$

$$(\tilde{a}_\mu - a_\mu) u^\mu = \mathring{Q}_{\lambda\mu\nu} u^\lambda u^\mu u^\nu$$

$$a^\mu = 0, \tilde{a}_\mu = \mathring{Q}_{\lambda\nu\mu} u^\lambda u^\nu;$$

$$\mathring{Q}_{(\lambda\mu\nu)} = 0, \mathring{Q}_{(\lambda\mu)v} = 0$$

$$S_{\text{STTEGR}} := \int d^4x \sqrt{-g} \left[\frac{c^4}{16\pi G} \underbrace{\mathcal{L}_{\text{STTEGR}}}_{\mathring{Q}} + \mathcal{L}_m \right],$$

$$\begin{aligned}
\dot{Q} &:= g^{\mu\nu} \left(\dot{L}^\alpha{}_{\beta\mu} \dot{L}^\beta{}_{\nu\alpha} - \dot{L}^\alpha{}_{\beta\alpha} \dot{L}^\beta{}_{\mu\nu} \right) \\
&= \frac{1}{4} \left(\dot{Q}_\alpha \dot{Q}^\alpha - \dot{Q}_{\alpha\beta\gamma} \dot{Q}^{\alpha\beta\gamma} \right) \\
&\quad + \frac{1}{2} \left(\dot{Q}_{\alpha\beta\gamma} \dot{Q}^{\beta\alpha\gamma} - \dot{Q}_\alpha \dot{Q}^\alpha \right), \\
\dot{Q} &= \dot{R} + \dot{\nabla}_\mu \left(\dot{Q}^\mu - \bar{\dot{Q}}^\mu \right), \\
\dot{\nabla}_\mu \left(\dot{Q}^\mu - \bar{\dot{Q}}^\mu \right) &\equiv \frac{1}{\sqrt{-g}} \partial_\mu \left[\sqrt{-g} \left(\dot{Q}^\mu - \bar{\dot{Q}}^\mu \right) \right] \\
\dot{Q}_{\text{gen}} &:= c_1 \dot{Q}_{\alpha\beta\gamma} \dot{Q}^{\alpha\beta\gamma} + c_2 \dot{Q}_{\alpha\beta\gamma} \dot{Q}^{\beta\alpha\gamma} + c_3 \dot{Q}_\alpha \dot{Q}^\alpha \\
&\quad + c_4 \dot{Q}_\alpha \dot{Q}^\alpha + c_5 \dot{Q}_\alpha \dot{Q}^\alpha, \\
\dot{P}_{\mu\nu}^\alpha &:= \frac{1}{2\sqrt{-g}} \frac{\partial(\sqrt{-g}\dot{Q})}{\partial\dot{Q}_\alpha^{\mu\nu}} \\
&= \frac{1}{4} \dot{Q}^\alpha{}_{\mu\nu} - \frac{1}{4} \dot{Q}_{(\mu}{}^\alpha{}_{\nu)} - \frac{1}{4} g_{\mu\nu} \dot{Q}^{\alpha\beta}{}_\beta \\
&\quad + \frac{1}{4} \left[\dot{Q}_\beta^{\beta\alpha} g_{\mu\nu} + \frac{1}{2} \delta_{(\mu}^\alpha \dot{Q}_{\nu)}{}^\beta{}_\beta \right].
\end{aligned}$$

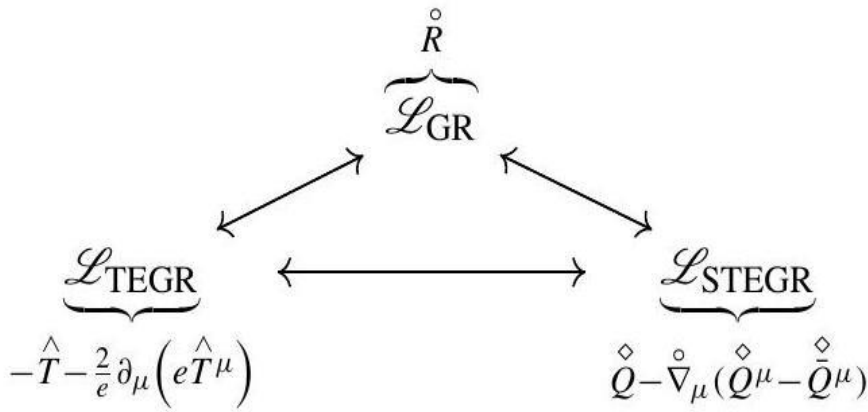
$$\dot{Q} := \dot{Q}_{\alpha\mu\nu} \dot{P}^{\alpha\mu\nu}.$$

$$\begin{aligned}
\frac{1}{\sqrt{-g}} \dot{q}_{\mu\nu} &:= \frac{1}{\sqrt{-g}} \frac{\partial(\sqrt{-g}\dot{Q})}{\partial g^{\mu\nu}} - \frac{1}{2} \dot{Q} g_{\mu\nu} \\
&= \frac{1}{4} \left(2 \dot{Q}_{\alpha\beta\mu} \dot{Q}^{\alpha\beta}{}_\nu - \dot{Q}_{\mu\alpha\beta} \dot{Q}_\nu{}^{\alpha\beta} \right) \\
&\quad - \frac{1}{4} \left(2 \dot{Q}_\alpha{}^\beta{}_\beta \dot{Q}^\alpha{}_{\mu\nu} - \dot{Q}_\mu{}^\beta{}_\beta \dot{Q}_\nu{}^\beta{}_\beta \right) \\
&\quad - \frac{1}{2} \left(\dot{Q}_{\alpha\beta\mu} - \dot{Q}_\nu{}^{\beta\alpha}{}_\alpha \dot{Q}^\alpha{}_{\mu\nu} \right). \\
\dot{G}_{\mu\nu} &:= -2 \nabla_\alpha \left(\sqrt{-g} \dot{P}_{\mu\nu}^\alpha \right) \\
&\quad + \dot{q}_{\mu\nu} - \frac{\sqrt{-g}}{2} \dot{Q} g_{\mu\nu} = \frac{8\pi G}{c^4} T_{\mu\nu},
\end{aligned}$$

$$\nabla_\mu \nabla_\nu \left(\sqrt{-g} \dot{P}_\alpha^{\mu\nu} \right) = 0,$$

$$\begin{aligned}
\Gamma_{\mu\nu}^\alpha &:= (e^{-1})^\alpha{}_\beta \partial_\mu e^\beta{}_\nu. \\
T_{\mu\nu}^\alpha &:= (e^{-1})^\alpha{}_\beta \partial_{[\nu} e_{\mu]}^\beta = 0, \\
\partial_\mu e^\beta{}_\nu &= \partial_\nu e^\beta{}_\mu \Leftrightarrow e_\beta^\alpha \equiv e_\beta'^\alpha := \partial_\beta \xi^\alpha, \\
\Gamma_{\mu\nu}^\alpha &= \frac{\partial x^\alpha}{\partial \xi^\lambda} \partial_\mu \partial_\nu \xi^\lambda. \\
\xi^\alpha &:= M_\beta^\alpha x^\beta + \xi_0^\alpha, \\
\dot{\nabla}_\mu &= \partial_\mu, \dot{L}_{\mu\nu}^\lambda = -\dot{\Gamma}_{\mu\nu}^\lambda.
\end{aligned}$$





$$\begin{aligned} & \nabla_\lambda R_{\beta\mu\nu}^\alpha + \nabla_\mu R_{\beta\nu\lambda}^\alpha + \nabla_\nu R_{\beta\lambda\mu}^\alpha \\ &= T_{\mu\lambda}^\rho R_{\beta\nu\rho}^\alpha + T_{\nu\lambda}^\rho R_{\beta\mu\rho}^\alpha + T_{\nu\mu}^\rho R_{\beta\lambda\rho}^\alpha, \end{aligned}$$

$$\overset{\circ}{\nabla}_\lambda \overset{\circ}{R}_{\beta\mu\nu}^\alpha + \overset{\circ}{\nabla}_\mu \overset{\circ}{R}_{\beta\nu\lambda}^\alpha + \overset{\circ}{\nabla}_\nu \overset{\circ}{R}_{\beta\lambda\mu}^\alpha = 0.$$

$$\partial_\lambda \overset{\circ}{R}_{\beta\mu\nu}^\lambda + \partial_\mu \overset{\circ}{R}_{\beta\nu\lambda}^\lambda + \partial_\nu \overset{\circ}{R}_{\beta\lambda\mu}^\lambda = 0.$$

$$\partial_\lambda \overset{\circ}{R}_{\beta\mu\nu}^\lambda - \partial_\mu \overset{\circ}{R}_{\beta\lambda\nu}^\lambda + \partial_\nu \overset{\circ}{R}_{\beta\lambda\mu}^\lambda = 0.$$

$$-\partial_\lambda \overset{\circ}{R}_\nu^\lambda - \partial_\beta \overset{\circ}{R}_\nu^\beta + \partial_\nu \overset{\circ}{R} = 0,$$

$$\partial_\mu \overset{\circ}{R}_\nu^\mu - \frac{1}{2} \partial_\nu \overset{\circ}{R} = 0.$$

$$\partial_\mu \left(\overset{\circ}{R}^{\mu\nu} - \frac{1}{2} g^{\mu\nu} \overset{\circ}{R} \right) = 0 \Rightarrow \overset{\circ}{\nabla}_\mu \left(\overset{\circ}{R}^{\mu\nu} - \frac{1}{2} g^{\mu\nu} \overset{\circ}{R} \right) = 0,$$

$$\overset{\circ}{\nabla}_\mu \overset{\circ}{G}^{\mu\nu} = 0, \Leftrightarrow \overset{\circ}{\nabla}_\mu T^{\mu\nu} = 0.$$

$$\overset{\circ}{\nabla}_\lambda R_{\beta\mu\nu}^\alpha + \overset{\circ}{\nabla}_\nu R_{\beta\lambda\mu}^\alpha + \overset{\circ}{\nabla}_\mu R_{\beta\nu\lambda}^\alpha = 0,$$

$$\begin{aligned} \hat{\mathcal{K}}_{\beta\mu\nu}^\alpha &:= \overset{\circ}{\nabla}_\mu \hat{K}_{\beta\nu}^\alpha - \overset{\circ}{\nabla}_\nu \hat{K}_{\beta\mu}^\alpha \\ &\quad + \hat{K}_{\sigma\mu}^\alpha \hat{K}_{\beta\nu}^\sigma - \hat{K}_{\sigma\nu}^\alpha \hat{K}_{\beta\mu}^\sigma, \end{aligned}$$

$$\hat{\mathcal{K}}_{\beta\mu\nu}^\alpha = -\hat{\mathcal{K}}_{\beta}{}^\alpha{}_{\mu\nu}, \hat{\mathcal{K}}_{\beta\mu\nu}^\alpha = -\hat{\mathcal{K}}_{\nu\mu}^\alpha{}_{\beta}.$$

$$\begin{aligned} & \hat{\nabla}_\lambda \overset{\circ}{R}_{\beta\mu\nu}^\lambda + \hat{\nabla}_\mu \overset{\circ}{R}_{\beta\nu\lambda}^\lambda + \hat{\nabla}_\nu \overset{\circ}{R}_{\beta\lambda\mu}^\lambda \\ & + \hat{\nabla}_\lambda \hat{\mathcal{K}}_{\beta\mu\nu}^\lambda + \hat{\nabla}_\mu \hat{\mathcal{K}}_{\beta\nu\lambda}^\lambda + \hat{\nabla}_\nu \hat{\mathcal{K}}_{\beta\lambda\mu}^\lambda = 0. \end{aligned}$$

$$\hat{\nabla}_\mu \left(\overset{\circ}{R}_\nu^\mu + \hat{\mathcal{K}}_\nu^\mu \right) - \frac{1}{2} \hat{\nabla}_\nu (\overset{\circ}{R} + \hat{\mathcal{K}}) = 0,$$

$$\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \overset{\circ}{R} = -\hat{\mathcal{K}}_{\mu\nu} + \frac{1}{2} g_{\mu\nu} \hat{\mathcal{K}},$$

$$\hat{\mathcal{K}}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \hat{\mathcal{K}} = 0,$$

$$\begin{aligned} \hat{\mathcal{K}}_{\mu\nu} &= \overset{\circ}{\nabla}_\alpha \hat{K}_{\mu\nu}^\alpha - \overset{\circ}{\nabla}_\nu \hat{K}_{\mu\alpha}^\alpha + \hat{K}^\sigma{}_{\mu\nu} \hat{K}^\alpha{}_{\sigma\alpha} - \hat{K}^\sigma{}_{\mu\alpha} \hat{K}^\alpha{}_{\sigma\nu} \\ &= \overset{\circ}{\nabla}_\alpha \hat{K}_{\mu\nu}^\alpha + \overset{\circ}{\nabla}_\nu \hat{T}_\mu - \hat{K}_{\sigma\mu\nu} \hat{T}^\sigma - \hat{K}_{\mu\alpha}^\sigma \hat{K}_{\sigma\nu}^\alpha \\ &= \overset{\circ}{\nabla}_\alpha \hat{S}_\nu{}^\alpha{}_\mu + \overset{\circ}{\nabla}_\alpha \hat{T}^\alpha g_{\mu\nu} - \hat{K}_{\sigma\nu}^\alpha \hat{S}_\alpha{}^\sigma{}_\mu, \end{aligned}$$

$$\hat{K}^\alpha{}_{\mu\alpha} = -\hat{T}_\mu,$$

$$\hat{K}^\alpha{}_{\alpha\mu} = 0,$$

$$\hat{K}_{\nu\lambda}^\mu = \hat{S}_\lambda^{\mu\nu} + \delta_\lambda^\nu \hat{T}^\mu - \delta_\lambda^\mu \hat{T}^\nu.$$

$$\hat{\mathcal{K}} = 2\overset{\circ}{\nabla}_\lambda \hat{T}^\lambda + \hat{T} = \frac{2}{e} \partial_\lambda (e \hat{T}^\lambda) + \hat{T}.$$

$$\overset{\circ}{\nabla}_\alpha \hat{S}_{\nu\mu}{}^\alpha + \hat{K}_{\sigma\nu}^\alpha \hat{S}_\alpha{}^\sigma{}_\mu + \frac{1}{2} g_{\mu\nu} \hat{T} = 0,$$

$$\partial_\lambda R_{\beta\mu\nu}^\alpha + \partial_\nu R_{\beta\lambda\mu}^\alpha + \partial_\mu R_{\beta\nu\lambda}^\alpha = 0,$$

$$\overset{\circ}{\mathcal{L}}^\alpha{}_{\beta\mu\nu} = \overset{\circ}{\nabla}_\mu \overset{\circ}{L}^\alpha{}_{\beta\nu} - \overset{\circ}{\nabla}_\nu \overset{\circ}{L}^\alpha{}_{\beta\mu} + \overset{\circ}{L}^\alpha{}_\sigma \overset{\circ}{L}^\sigma{}_{\beta\nu} - \overset{\circ}{L}^\alpha{}_{\sigma\nu} \overset{\circ}{L}^\sigma{}_{\beta\mu},$$

$$\overset{\circ}{\mathcal{L}}^\alpha{}_{\beta\mu\nu} = -\overset{\circ}{\mathcal{L}}^\beta{}_{\alpha\mu\nu}, \quad \overset{\circ}{\mathcal{L}}^\alpha{}_{\beta\mu\nu} = -\overset{\circ}{\mathcal{L}}^\alpha{}_{\beta\nu\mu},$$

$$\begin{aligned} &\partial_\lambda \overset{\circ}{R}_{\beta\mu\nu} + \partial_\mu \overset{\circ}{R}_{\beta\nu\lambda}^\lambda + \partial_\nu \overset{\circ}{R}_{\beta\lambda\mu} \\ &+ \partial_\lambda \overset{\circ}{\mathcal{L}}_{\beta\mu\nu}^\lambda + \partial_\mu \overset{\circ}{\mathcal{L}}_{\beta\nu\lambda}^\lambda + \partial_\nu \overset{\circ}{\mathcal{L}}_{\beta\lambda\mu}^\lambda = 0. \end{aligned}$$

$$\overset{\circ}{R}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \overset{\circ}{R} = -\overset{\circ}{\mathcal{L}}_{\mu\nu} + \frac{1}{2} g_{\mu\nu} \overset{\circ}{\mathcal{L}},$$

$$\overset{\circ}{\mathcal{L}}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \overset{\circ}{\mathcal{L}} = 0,$$

$$\begin{aligned} \overset{\circ}{\mathcal{L}}_{\mu\nu} &= \overset{\circ}{\nabla}_\alpha \overset{\circ}{\mathcal{L}}^\alpha{}_{\mu\nu} - \overset{\circ}{\nabla}_\nu \overset{\circ}{\mathcal{L}}^\alpha{}_{\mu\alpha} + \overset{\circ}{\mathcal{L}}^\sigma{}_{\mu\nu} \overset{\circ}{\mathcal{L}}^\alpha{}_{\sigma\alpha} - \overset{\circ}{L}^\sigma{}_{\mu\alpha} \overset{\circ}{\mathcal{L}}^\alpha{}_{\sigma\nu} \\ &= \overset{\circ}{\nabla}_\alpha \overset{\circ}{\mathcal{L}}^\alpha{}_{\mu\nu} + \frac{1}{2} \circ_\nu{}_\nu \overset{\circ}{Q}_\mu - \frac{1}{2} \overset{\circ}{Q}_\alpha \overset{\circ}{\mathcal{L}}^\alpha{}_{\mu\nu} \\ &\quad - \frac{1}{4} \left[\overset{\circ}{Q}_\mu{}^\sigma{}_\alpha \overset{\circ}{Q}_\nu{}^\alpha{}_\sigma + 2 \overset{\circ}{Q}^\alpha{}_{\sigma\nu} \left(\overset{\circ}{Q}^\sigma{}_{\alpha\mu} - \overset{\circ}{Q}_\alpha{}^\sigma{}_\mu \right) \right], \end{aligned}$$

$$\overset{\circ}{L}^\alpha{}_{\mu\alpha} = -\frac{1}{2} \overset{\circ}{Q}_\mu$$

$$\overset{\circ}{L}^\alpha{}_{\mu\nu} = 2 \overset{\circ}{P}_{\mu\nu}^\alpha + \frac{1}{2} g_{\mu\nu} \left(\overset{\circ}{Q}^\alpha - \overset{\circ}{Q}^\alpha \right)$$

$$-\frac{1}{4} \left(\delta_\mu^\alpha \overset{\circ}{Q}_\nu + \delta_\nu^\alpha \overset{\circ}{Q}_\mu \right)$$

$$\overset{\circ}{\mathcal{L}} = \overset{\circ}{\nabla}_\alpha \left(\overset{\circ}{Q}^\alpha - \overset{\circ}{Q}^\alpha \right) + \frac{1}{4} \overset{\circ}{Q}_{\alpha\beta\gamma} \overset{\circ}{Q}^{\alpha\beta\gamma} - \frac{1}{2} \overset{\circ}{Q}_{\alpha\beta\gamma} \overset{\circ}{Q}^{\gamma\beta\alpha}$$

$$-\frac{1}{4} \overset{\circ}{Q}_\alpha \overset{\circ}{Q}^\alpha + \frac{1}{2} \overset{\circ}{Q}_\alpha \overset{\circ}{Q}^\alpha$$

$$= \overset{\circ}{\nabla}_\alpha \left(\overset{\circ}{Q}^\alpha - \overset{\circ}{Q}^\alpha \right) - \overset{\circ}{Q}.$$

$$\partial_\alpha \overset{\circ}{Q}^\alpha = \overset{\circ}{\nabla}_\alpha \overset{\circ}{Q}^\alpha + \overset{\circ}{L}^\alpha{}_{\sigma\alpha} \overset{\circ}{Q}^\sigma = \overset{\circ}{\nabla}_\alpha \overset{\circ}{Q}^\alpha - \frac{1}{2} \overset{\circ}{Q}_\alpha \overset{\circ}{Q}^\alpha,$$

$$\begin{aligned} &2 \partial_\alpha \overset{\circ}{P}^\alpha{}_{\mu\nu} + \frac{1}{2} \overset{\circ}{Q}_{\alpha\mu\nu} \left(\overset{\circ}{Q}^\alpha - \overset{\circ}{Q}^\alpha \right) + \frac{1}{2} g_{\mu\nu} \partial_\alpha \left(\overset{\circ}{Q}^\alpha - \overset{\circ}{Q}^\alpha \right) \\ &+ \frac{1}{2} \overset{\circ}{L}^\sigma{}_{\mu\nu} \overset{\circ}{Q}_\sigma + \frac{1}{4} \overset{\circ}{Q}_\mu{}^\alpha{}_\sigma \overset{\circ}{Q}_\nu{}^\sigma{}_\alpha + \frac{1}{2} \overset{\circ}{Q}^\alpha{}_{\sigma\mu} \left(\overset{\circ}{Q}^\sigma{}_{\nu\alpha} - \overset{\circ}{Q}_\alpha{}^\sigma{}_\nu \right) \\ &- \frac{1}{2} g_{\mu\nu} \overset{\circ}{\nabla}_\alpha \left(\overset{\circ}{Q}^\alpha - \overset{\circ}{Q}_\alpha \right) + \frac{1}{2} g_{\mu\nu} \overset{\circ}{Q}, \end{aligned}$$

$$\frac{2}{\sqrt{-g}} \partial_\alpha (\sqrt{-g} \dot{P}_{\mu\nu}^\alpha) - \frac{1}{\sqrt{-g}} \dot{q}_{\mu\nu} + \frac{1}{2} g_{\mu\nu} \dot{Q} = 0. \dots$$

$$ds^2 = -e^{v(t,r)} dt^2 + e^{\lambda(t,r)} dr^2 + r^2 d\varphi^2,$$

$$-e^{v(t,r)} \approx -1 + \frac{2M}{r},$$

$$ds^2 = -\left(1 - \frac{2M}{r}\right) dt^2 + \left(1 - \frac{2M}{r}\right)^{-1} dr^2 + r^2 d\varphi^2,$$

$$\dot{f}(t,r) := \frac{df(t,r)}{dt}, f'(t,r) := \frac{df(t,r)}{dr}.$$

$$\dot{G}_{\mu\nu} \equiv \dot{R}_{\mu\nu} = 0,$$

$$\dot{G}_{tr} \equiv \frac{\dot{\lambda}(t,r)}{r} = 0, \Rightarrow \lambda = \lambda(r).$$

$$\dot{G}_{rr} \equiv -e^{\lambda(r)} + rv'(t,r) + 1 = 0,$$

$$\dot{G}_{tt} \equiv e^{-\lambda(r)}(r\lambda'(r) - 1) + 1 = 0.$$

$$[e^{-\lambda(r)}r]' = 1 \Rightarrow e^{-\lambda(r)} = 1 - \frac{C_1}{r},$$

$$\lambda'(r) + v'(r) = 0, \Rightarrow \lambda(r) + v(r) = C_2,$$

$$-e^{v(r)} = 1 - \frac{2M}{r}, e^{\lambda(r)} = \frac{1}{1 - \frac{2M}{r}}.$$

$$e_\mu^A = \begin{pmatrix} \sqrt{-e^{v(r)}} & 0 & 0 & 0 \\ 0 & \sqrt{e^{\lambda(r)}} & 0 & 0 \\ 0 & 0 & r & 0 \\ 0 & 0 & 0 & r \sin \theta \end{pmatrix}.$$

$$\hat{T}_{tr}^t = -\frac{1}{2} v'(r) = -\frac{M}{r^2} \left(1 - \frac{2M}{r}\right)^{-1},$$

$$\hat{T}_{r\varphi}^\varphi = \frac{1}{r}.$$

$$\hat{K}_{ttr} = \frac{1}{2} e^{v(r)} v'(r) = \frac{M}{r^2}$$

$$\hat{K}_{\varphi r \varphi} = r$$

$$\hat{S}_{\hat{t}}^{tr} = \frac{2e^{-\lambda(r)}\sqrt{e^{-v(r)}}}{r} = \frac{2}{r} \sqrt{1 - \frac{2M}{r}},$$

$$\hat{S}_{\hat{\varphi}}^{r\varphi} = -\frac{e^{-\lambda(r)}(rv'(r)+2)}{2r^2} = \frac{M-r}{r^3}.$$

$$\hat{T} = -\frac{2e^{-\lambda(r)}(rv'(r)+1)}{r^2} = -\frac{2}{r^2},$$

$$\dot{Q}_{r\mu\nu} = \begin{pmatrix} -e^{v(r)}v'(r) & 0 & 0 \\ 0 & e^{\lambda(r)}\lambda'(r) & 0 \\ 0 & 0 & 2r \end{pmatrix}$$

$$= \begin{pmatrix} -\frac{2M}{r^2} & 0 & 0 \\ 0 & -\frac{2M}{r^2\left(1 - \frac{2M}{r}\right)^2} & 0 \\ 0 & 0 & 2r \end{pmatrix},$$

$$\begin{aligned}
\mathring{P}_{tr}^t &= \frac{r\lambda'(r) - rv'(r) + 4}{8r} = \frac{1}{8}(\lambda'(r) + v'(r)), \\
\mathring{P}_{rr}^r &= \frac{e^{v(r)-\lambda(r)}}{r} = \frac{1}{r}\left(1 - \frac{2M}{r}\right)^2, \\
\mathring{P}_{\varphi\varphi}^r &= -\frac{1}{4}re^{-\lambda(r)}(rv'(r) + 2) = \frac{M-r}{2}, \\
\mathring{P}_{r\varphi}^\varphi &= \frac{1}{8}(\lambda'(r) + v'(r)) = 0, \\
\frac{\mathring{q}_{\mu\nu}}{\sqrt{-g}} &= \begin{pmatrix} \frac{2e^{v(r)-\lambda(r)}v'(r)}{r} & 0 & 0 \\ 0 & \frac{2rv'(r) + 2}{r^2} & 0 \\ 0 & 0 & -\frac{rv'(r) + 2}{e^{\lambda(r)}} \end{pmatrix} \\
&= \begin{pmatrix} \frac{4M}{r^3}\left(1 - \frac{2M}{r}\right) & 0 & 0 \\ 0 & \frac{2}{r^2\left(1 - \frac{2M}{r}\right)} & 0 \\ 0 & 0 & \frac{2M}{r} - 2 \end{pmatrix}. \\
&e_\mu^A g_{\nu\rho} \partial_\sigma \hat{S}_A^{\rho\sigma} + e^{-1} e_\mu^A g_{\nu\rho} \partial_\sigma e \\
&\quad - \hat{S}_B^{\sigma}{}_\nu T^B{}_{\sigma\mu} + \frac{1}{2} g_{\mu\nu} \hat{T} = 0. \\
&\partial_\sigma \hat{S}_{\mu\nu}{}^\sigma - \hat{S}_{\alpha\nu}{}^\sigma \Gamma^\alpha{}_{\nu\sigma} - \hat{S}_{\alpha\mu}{}^\sigma \Gamma^\alpha{}_{\nu\sigma} - \hat{S}_\mu{}^{\rho\sigma} \Gamma^\alpha{}_{\rho\sigma} g_{\nu\alpha} \\
&\quad - \hat{S}_\alpha{}^\sigma{}_\nu \hat{T}^\alpha{}_{\sigma\mu} + \Gamma_{\alpha\sigma}^\alpha \hat{S}_{\mu\nu}{}^\sigma + \frac{1}{2} \hat{T} g_{\mu\nu} = 0. \\
&\mathring{\nabla}_\sigma \hat{S}_{\mu\nu}{}^\sigma - \hat{K}^\alpha{}_{\mu\sigma} \hat{S}_{\alpha\nu}{}^\sigma - \hat{K}^\alpha{}_{\nu\sigma} \hat{S}_{\mu\alpha}{}^\sigma \\
&\quad + \hat{K}^\sigma{}_{\alpha\sigma} \hat{S}_{\mu\nu}{}^\alpha - \hat{K}_{\nu\rho\sigma} \hat{S}_\mu{}^{\rho\sigma} + \hat{T}_\sigma \hat{S}_{\mu\nu}{}^\sigma \\
&\quad - \hat{S}_\alpha{}^\sigma{}_\nu \hat{T}^\alpha{}_{\sigma\mu} + \frac{1}{2} g_{\mu\nu} \hat{T} = 0.
\end{aligned}$$

$$-\hat{K}^\alpha{}_{\nu\sigma} \hat{S}_{\mu\alpha}{}^\sigma - \hat{K}_{\nu\rho\sigma} \hat{S}_\mu{}^{\rho\sigma} = 0,$$

$$\hat{K}^\sigma{}_{\alpha\sigma} \hat{S}_{\mu\nu}{}^\alpha + \hat{T}_\sigma \hat{S}_{\mu\nu}{}^\sigma = 0.$$

$$\mathring{\nabla} \hat{S}_{\mu\nu}{}^\sigma + \hat{K}_{\mu\sigma}^\alpha \hat{S}_\alpha{}_\nu - \hat{T}^\alpha{}_{\sigma\mu} \hat{S}_\alpha{}_\nu + \frac{1}{2} \hat{T} g_{\mu\nu} = \mathring{\nabla}_\sigma \hat{S}_{\mu\nu}{}^\sigma + \hat{K}^\alpha{}_{\sigma\mu} \hat{S}_\alpha{}_\nu + \frac{1}{2} g_{\mu\nu} \hat{T} = 0$$

$$\frac{\partial_\alpha \sqrt{-g}}{\sqrt{-g}} = \mathring{\Gamma}_{\alpha\sigma}^\sigma = -\mathring{L}_{\alpha\sigma}^\sigma = \frac{1}{2} \mathring{Q}_\alpha.$$

$$\begin{aligned}
&2\partial_\alpha P^\alpha{}_{\mu\nu} + \frac{1}{2} \mathring{Q}_{\alpha\mu\nu} (\mathring{Q}^\alpha - \mathring{Q}^\alpha) + \frac{1}{2} g_{\mu\nu} \partial_\alpha (\mathring{Q}^\alpha - \mathring{Q}^\alpha) \\
&+ \frac{1}{2} \mathring{L}^\sigma{}_{\mu\nu} \mathring{Q}_\sigma + \frac{1}{4} \mathring{Q}_\mu{}^\alpha{}_\sigma \mathring{Q}_\nu{}^\sigma{}_\alpha + \frac{1}{2} \mathring{Q}^\alpha{}_{\sigma\mu} (\mathring{Q}^\sigma{}_{\nu\alpha} - \mathring{Q}_\alpha{}^\sigma{}_\nu) \\
&- \frac{1}{2} g_{\mu\nu} \mathring{\nabla}_\alpha (\mathring{Q}^\alpha - \mathring{Q}^\alpha) + \frac{1}{2} g_{\mu\nu} \mathring{Q} = 0.
\end{aligned}$$

$$\mathbb{P}(x_n, t_n \mid x_1, t_1; \dots; x_{n-1}, t_{n-1}) = \mathbb{P}(x_n, t_n \mid x_{n-1}, t_{n-1})$$

$$\hat{\mathcal{E}}(t)(\hat{\rho}(0)) = \hat{\rho}(t)$$

$$\hat{\mathcal{E}}(t_1+t_2)=\hat{\mathcal{E}}(t_1)\hat{\mathcal{E}}(t_2),$$



$$\begin{aligned}
\frac{d\hat{p}}{dt}(t) &= \hat{L}\hat{p}(t) \\
\hat{H} &= \hat{H}_S + \hat{H}_B + \hat{H}_{SB} \\
\hat{H}_S &= \frac{\hat{p}^2}{2M} + V(\hat{Q}), \\
\hat{H}_B^{(k)} &= \sum_{k=1}^N \left(\frac{\hat{p}_k^2}{2m_k} + \frac{m_k \omega_k^2 \hat{q}_k^2}{2} \right), \\
\hat{H}_{SB}^{(k)} &= \sum_{k=1}^N \left(g_k \hat{q}_k \hat{Q} + \frac{g_k^2}{2m_k \omega_k^2} \hat{Q}^2 \right). \\
\frac{d^2 \hat{Q}}{dt^2}(t) + V'(\hat{Q}(t)) + \frac{1}{M} \int_0^t k(t-t') \hat{P}(t') dt' + k(t) \hat{Q}(0) &= \hat{f}(t) \\
k(t) &= \sum_{i=1}^N \frac{g_k^2}{m_k \omega_k^2} \cos(\omega_k t) \\
\hat{f}(t) &= \sum_{i=1}^N \left(g_k \hat{q}_k(0) \cos(\omega_k t) + \frac{g_k \hat{p}_k(0)}{m_k \omega_k} \sin(\omega_k t) \right) \\
\langle \hat{f}(t) \rangle &= \text{Tr}(\hat{f}(t) \rho_{\text{th}}) = 0 \\
\langle \{\hat{f}(t), \hat{f}(t')\} \rangle &= \sum_{k=1}^N \frac{\hbar g_k^2}{m_k \omega_k} \coth\left(\frac{\hbar \omega_k}{2k_B T}\right) \cos(\omega_k(t-t')), \\
\hat{\Theta} \hat{q} \hat{\Theta}^{-1} &= \hat{q}, \hat{\Theta} \hat{p} \hat{\Theta}^{-1} = -\hat{p}, \\
\hat{A}_R(t) &= \hat{\Theta} \hat{A}(-t) \hat{\Theta}^{-1} \\
\hat{q}_R(t) &= \hat{\Theta} \hat{q}(-t) \hat{\Theta}^{-1} = \hat{q}(t), \hat{p}_R(t) = \hat{\Theta} \hat{p}(-t) \hat{\Theta}^{-1} = -\hat{p}(t). \\
M \frac{d^2 \hat{Q}}{dt^2}(-t) + V'(\hat{Q}(-t)) + \frac{1}{M} \int_0^{-t} k(-t-t') \hat{P}(t') dt' + k(-t) \hat{Q}(0) &= \hat{f}(-t) \\
\frac{d\hat{p}_R}{dt} &= -\hat{\Theta} \frac{d\hat{p}}{dt} \hat{\Theta}^{-1} = -\hat{\Theta} \hat{L} \hat{\Theta}^{-1} \hat{\Theta} \hat{p} \hat{\Theta}^{-1} = -\hat{L}_R \hat{p}_R. \\
\hat{\rho}(t) &= \hat{\mathcal{E}}(t) \hat{\rho}(0), \hat{\mathcal{E}}(t) = \exp(-\hat{L}t). \\
\hat{\Theta} \hat{\mathcal{E}}(t) \hat{\Theta}^{-1} &= \exp(-\hat{\Theta} \hat{L} \hat{\Theta}^{-1} t) = \exp(-\hat{L}_R t). \\
\hat{\Theta} \hat{\mathcal{E}}(t) \hat{\Theta}^{-1} &= \exp(\hat{L}t) = \hat{\mathcal{E}}(t)^{-1}. \\
\hat{\Theta} \hat{U}(t) \hat{\Theta}^{-1} &= \hat{U}(-t) \\
\frac{d\hat{p}}{dt} &= \hat{L}(t) \hat{p}(t) \\
\frac{d\hat{p}_R}{dt}(t) &= -\hat{\Theta} \frac{d\hat{p}}{dt}(-t) \hat{\Theta}^{-1} = -\hat{L}_R(-t) \hat{p}_R(t) \\
\hat{\rho}(t) &= \hat{\mathcal{E}}(t, 0) \hat{\rho}(0), \hat{\mathcal{E}}(t_2, t_1) = \hat{T} \exp\left(-\int_{t_1}^{t_2} \hat{L}(t') dt'\right), \\
\hat{\Theta} \hat{\mathcal{E}}(t_1, t_2) \hat{\Theta}^{-1} &= \hat{T} \exp\left(-\int_{t_1}^{t_2} \hat{\Theta} \hat{L}(t') \hat{\Theta}^{-1} dt'\right) = \hat{T} \exp\left(-\int_{t_1}^{t_2} \hat{L}_R(t') dt'\right). \\
\hat{\Theta} \hat{\mathcal{E}}(t_1, t_2) \hat{\Theta}^{-1} &= \mathcal{T} \exp\left(-\int_{t_2}^{t_1} \hat{L}(t') dt'\right) = \mathcal{E}(t_2, t_1)^{-1}. \\
\int_0^\infty k(t') dt' &< \infty
\end{aligned}$$

$$\begin{aligned}
& \int_0^{\tau_B} k(t') dt' \approx \int_0^{\infty} k(t') dt' \\
\int_0^t k(t-t') \hat{P}(t') dt' &= \int_0^t k(t') \hat{P}(t-t') dt' \approx \hat{P}(t) \int_0^{\tau_B} k(t') dt' = \hat{P}(t) \int_0^{\infty} k(t') dt' \\
\int_0^t k(t') dt' &= \text{sgn}(t) \int_0^{|t|} k(t') dt', \\
\int_0^t k(t-t') \hat{P}(t') dt' &\approx \text{sgn}(t) \hat{P}(t) \int_0^{\infty} k(t') dt' \\
M \frac{d^2 \hat{Q}}{dt^2} + V'(\hat{Q}(t)) + \text{sgn}(t) \gamma \hat{P}(t) &= \hat{f}(t). \\
\int_0^{\infty} k(t') dt' &= M\gamma \\
\langle \{\hat{f}(t), \hat{f}(t')\} \rangle &= \frac{\gamma M \hbar}{\pi} \int_0^{\Lambda} \omega \coth\left(\frac{\hbar \omega}{2k_B T}\right) \cos(\omega(t-t')) d\omega \\
\langle \{\hat{f}(t), \hat{f}(t')\} \rangle &= 2\gamma M k_B T \delta(t-t'). \\
\text{Tr}_S(\hat{Y} \hat{\mu}(t)) &= \text{Tr}_S(\hat{\rho}_S \hat{Y}(t)) \\
\dot{\hat{\mu}}(t) &= -\frac{i}{\hbar} [\hat{H}_S, \hat{\mu}(t)] + \frac{i}{2\hbar} [\{\gamma \text{sgn}(t) \hat{P}, \hat{\mu}(t)\}, \hat{Q}] + \frac{i}{2\hbar} [\{\hat{f}(t), \hat{\mu}(t)\}, \hat{Q}] \\
\hat{\rho}(t) &= \langle \hat{\mu}(t) \rangle := \text{Tr}_B(\hat{\mu}(t) \hat{\rho}_{\text{th}}), \\
\dot{\hat{\rho}}(t) &= -\frac{i}{\hbar} [\hat{H}_S, \hat{\rho}(t)] + \frac{i \text{sgn}(t)}{2\hbar} [\{\gamma \hat{P}, \hat{\rho}(t)\}, \hat{Q}] - \frac{\Gamma(t)}{\hbar^2} [[\hat{\rho}(t), \hat{Q}], \hat{Q}]. \\
\Gamma(t) &= \int_0^t \langle \{\hat{f}(t), \hat{f}(t')\} \rangle dt' \\
\int_0^t \langle \{\hat{f}(t), \hat{f}(t')\} \rangle dt' &= \text{sgn}(t) \int_0^{|t|} \langle \{\hat{f}(|t|), \hat{f}(t')\} \rangle dt' \\
\lim_{t \rightarrow \infty} \int_0^{|t|} \langle \{\hat{f}(|t|), \hat{f}(t')\} \rangle dt' &= 2\gamma M k_B T. \\
\dot{\hat{\rho}}(t) &= -\frac{i}{\hbar} [\hat{H}_S, \hat{\rho}(t)] + \frac{i \text{sgn}(t)}{2\hbar} [\{\gamma \hat{P}, \hat{\rho}(t)\}, \hat{Q}] - \text{sgn}(t) \frac{2\gamma M k_B T}{\hbar^2} [[\hat{\rho}(t), \hat{Q}], \hat{Q}]. \\
\frac{d\hat{\rho}}{dt}(t) &= (i\hat{\mathcal{L}}_H + \text{sgn}(t)\hat{\mathcal{L}}_D)\hat{\rho}(t) \\
\hat{\mathcal{L}}_H \hat{\rho}(t) &= -\frac{1}{\hbar} [\hat{H}_S, \hat{\rho}(t)] \\
\hat{\mathcal{L}}_D \hat{\rho}(t) &= \frac{i}{2\hbar} [\{\gamma \hat{P}, \hat{\rho}(t)\}, \hat{Q}] - \frac{2\gamma M k_B T}{\hbar^2} [[\hat{\rho}(t), \hat{Q}], \hat{Q}]. \\
\frac{d\hat{\rho}_R}{dt}(t) &= -\hat{\Theta} \frac{d\hat{\rho}}{dt}(t) \hat{\Theta}^{-1} = (i\hat{\mathcal{L}}_H + \text{sgn}(t)\hat{\mathcal{L}}_D)\hat{\rho}_R(t) \\
\hat{\rho}(t) &= \hat{\mathcal{E}}(t) \hat{\rho}(0), \hat{\mathcal{E}}(t) = \exp(i\hat{\mathcal{L}}_H t + \hat{\mathcal{L}}_D |t|) \\
\rho(p, q, t) &= \frac{1}{\sqrt{\pi N(t)}} \exp\left(\frac{-(q + \text{sgn}(t)A(t)p)^2}{N(t)} - B(t)p^2\right)
\end{aligned}$$

$$N(t) = \frac{mk_B T}{\hbar^2} (1 - e^{-2\gamma|t|}) + \frac{e^{-2\gamma|t|}}{\sigma^2}$$

$$A(t) = \frac{i\hbar}{2\sigma^2 m \gamma} e^{-\gamma|t|} (1 - e^{-\gamma|t|}) - \frac{ik_B T}{2\hbar \gamma} (1 - e^{-\gamma|t|})^2$$

$$B(t) = \frac{\hbar^2}{4\sigma^2 m^2 \gamma^2} (1 - e^{-\gamma|t|})^2 + \frac{\sigma^2}{4} + \frac{k_B T}{m \gamma^2} (2\gamma|t| - 3 + 4e^{-\gamma|t|} - e^{-2\gamma|t|}).$$

$$\psi(x, 0) = \frac{1}{(\sigma^2 \pi)^{1/4}} \exp\left(-\frac{x^2}{2\sigma^2}\right).$$

$$S_{\text{vN}}(\xi) = \frac{1-\xi}{2\xi} \log\left(\frac{1+\xi}{1-\xi}\right) - \log\left(\frac{2\xi}{1+\xi}\right)$$

$$\hat{H} = \hat{H}_S + \hat{H}_B + \hat{H}_{\text{SB}}$$

$$\hat{H}_I(t) = e^{\frac{i}{\hbar}(\hat{H}_S + \hat{H}_B)t} \hat{H}_{\text{SB}} e^{-\frac{i}{\hbar}(\hat{H}_S + \hat{H}_B)t}$$

$$\frac{d}{dt}\hat{\rho}_I(t) = -\frac{i}{\hbar}[\hat{H}_I(t), \hat{\rho}_I(t)]$$

$$\frac{d}{dt}\hat{\rho}_S(t) = -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left([\hat{H}_I(t), [\hat{H}_I(s), \hat{\rho}_I(s)]] \right) ds$$

$$\frac{d}{dt}\hat{\rho}_S(t) = -\frac{1}{\hbar^2} \int_0^t \text{Tr}_B \left([\hat{H}_I(t), [\hat{H}_I(s), \hat{\rho}_S(t) \otimes \hat{\rho}_B]] \right) ds$$

$$\hat{H}_I(t) = \sum_{\alpha,\omega} e^{-i\omega t} \hat{A}_\alpha(\omega) \otimes \hat{B}_\alpha(t)$$

$$\frac{d}{dt}\hat{\rho}_S(t) =$$

$$-\frac{1}{\hbar^2} \sum_{\alpha,\beta,\omega} [\Gamma_{\alpha\beta}(\omega, t) \hat{A}_\alpha^\dagger(\omega) \hat{A}_\beta(\omega) \hat{\rho}_S(t) + \Gamma_{\beta\alpha}^*(\omega, t) \hat{\rho}_S(t) \hat{A}_\alpha^\dagger(\omega) \hat{A}_\beta(\omega) (\Gamma_{\alpha\beta}(\omega, t) + \Gamma_{\beta\alpha}^*(\omega, t)) \hat{A}_\beta(\omega) \hat{\rho}_S(t)]$$

$$\Gamma_{\alpha\beta}(\omega, t) = \int_0^t e^{i\omega s} \text{Tr}(\hat{B}_\alpha^\dagger(t) \hat{B}_\beta(t-s) \hat{\rho}_B) ds$$

$$\Gamma_{\alpha\beta}(\omega, t) + \Gamma_{\beta\alpha}^*(\omega, t) = \text{sgn}(t) \int_{-|t|}^{|t|} e^{i\omega s} \text{Tr}(\hat{B}_\alpha^\dagger(s) \hat{B}_\beta(0) \hat{\rho}_B) ds.$$

$$\Gamma_{\alpha\beta}(\omega, t) + \Gamma_{\beta\alpha}^*(\omega, t) \approx \text{sgn}(t) \gamma_{\alpha\beta}(\omega)$$

$$\gamma_{\alpha\beta}(\omega) = \int_{-\infty}^{\infty} e^{i\omega s} \text{Tr}(\hat{B}_\alpha^\dagger(s) \hat{B}_\beta(0) \hat{\rho}_B) ds$$

$$\eta_{\alpha\beta}(\omega) = \frac{1}{2i} (\Gamma_{\alpha\beta}(\omega, t) - \Gamma_{\beta\alpha}^*(\omega, t)).$$

$$\frac{d}{dt}\hat{\rho}_S(t) = -\frac{i}{\hbar} [\hat{H}_S, \hat{\rho}_S] + \frac{\text{sgn}(t)}{\hbar^2} \sum_{\alpha,\beta,\omega} \hat{\mathcal{D}}_{\alpha\beta}(\omega) \hat{\rho}_S(t)$$

$$\hat{H}_S = \frac{1}{\hbar} \sum_{\alpha,\beta,\omega} \eta_{\alpha\beta}(\omega) \hat{A}_\alpha^\dagger(\omega) \hat{A}_\beta(\omega)$$

$$\hat{\mathcal{D}}_{\alpha\beta}(\omega) \hat{\rho}_S(t) = \gamma_{\alpha\beta}(\omega) \left(\hat{A}_\beta(\omega) \hat{\rho}_S(t) \hat{A}_\alpha^\dagger(\omega) - \frac{1}{2} \{ \hat{A}_\alpha^\dagger(\omega) \hat{A}_\beta(\omega), \hat{\rho}_S(t) \} \right).$$

$$\hat{\Theta} \gamma_{\alpha\beta}(\omega) \hat{\Theta}^{-1} = \int_{-\infty}^{\infty} e^{i\omega s} \text{Tr}(\hat{\Theta} \hat{B}_\alpha^\dagger(s) \hat{B}_\beta(0) \hat{\Theta}^{-1} \hat{\rho}_B) ds = \gamma_{\alpha\beta}(\omega).$$



$$\begin{aligned}
\frac{d}{dt}\rho_n(t) &= \sum_{n' \neq n} (W_{n,n'}\rho_{n'}(t) - W_{n',n}\rho_n(t)), \\
\langle n|\hat{A}_\alpha^\dagger(\omega)\hat{A}_\beta(\omega)|n'\rangle &= \delta_{n,n'}\langle n|\hat{A}_\alpha|m\rangle\langle m|\hat{A}_\beta|n\rangle, \\
\frac{d}{dt}\rho_n(t) &= \frac{\text{sgn}(t)}{\hbar^2} \sum_{n \neq n'} (W_{n,n'}\rho_{n'}(t) - W_{n',n}\rho_n(t)), \\
W_{n',n} &= \sum_{\alpha,\beta} \gamma_{\alpha\beta}(\varepsilon_n - \varepsilon_{n'})\langle n|\hat{A}_\alpha|n'\rangle\langle n'|\hat{A}_\beta|n\rangle \\
\hat{H} &= \hat{H}_0 + \lambda\hat{H}_{\text{SB}} \\
\frac{d}{dt}\rho_{\varepsilon,k}(t) &= \frac{\lambda^2}{\hbar^2} \sum_{\varepsilon',k'} \int_0^t \Lambda_{\varepsilon,k,\varepsilon',k'}(t-s)(\rho_{\varepsilon',k'}(s) - \rho_{\varepsilon,k}(s))ds, \\
\Lambda_{\varepsilon,k,\varepsilon',k'}(t) &= 2|\langle \varepsilon,k|\hat{H}_{\text{SB}}|\varepsilon',k'\rangle|^2 \cos\left(\frac{\varepsilon - \varepsilon'}{\hbar}t\right) \\
\int_0^t \Lambda_{\varepsilon,k,\varepsilon',k'}(t-s)ds &= 2\hbar \frac{|\langle \varepsilon,k|\hat{H}_{\text{SB}}|\varepsilon',k'\rangle|^2}{\varepsilon - \varepsilon'} \sin\left(\frac{\varepsilon - \varepsilon'}{\hbar}t\right). \\
\frac{\hbar}{\varepsilon - \varepsilon'} \sin\left(\frac{\varepsilon - \varepsilon'}{\hbar}t\right) &\approx \text{sgn}(t)\pi\hbar\delta(\varepsilon - \varepsilon') \\
\frac{d}{dt}\rho_{\varepsilon,k}(t) &= \text{sgn}(t) \sum_{k'} (W_{k,k'}^\varepsilon \rho_{\varepsilon,k'}(t) - W_{k',k}^\varepsilon \rho_{\varepsilon,k}(t)) \\
W_{k,k'}^\varepsilon &= \frac{2\lambda^2}{\hbar} |\langle \varepsilon,k|\hat{H}_{\text{SB}}|\varepsilon,k'\rangle|^2 \eta(\varepsilon) \\
\psi_R(x,a+t) &= \psi^*(x,a-t) \\
\frac{dW}{dt} &= -\frac{p}{M} \frac{dW}{dx} + \sum_{n=0}^{\infty} \frac{(-\hbar^2)^n V^{(2n+1)}(x)}{(2n+1)! 2^{2n}} \frac{d^{2n+1}W}{dp^{2n+1}} + \text{sgn}(t)\gamma \frac{d}{dp}(pW) + \text{sgn}(t)2\gamma M k_B T \frac{d^2W}{dp^2} \\
\frac{dW}{dt} &= -\frac{p}{M} \frac{dW}{dx} + \frac{dV}{dx} \frac{dW}{dp} + \text{sgn}(t)\gamma \frac{d}{dp}(pW) + \text{sgn}(t)2\gamma M k_B T \frac{d^2W}{dp^2} \\
\dot{\hat{\mu}}(t) &= -\frac{i}{\hbar} [\hat{H}_S, \hat{\mu}(t)] + \frac{i}{2\hbar} [\{\gamma \text{sgn}(t)\hat{P}, \hat{\mu}(t)\}, \hat{Q}] + \frac{i}{2\hbar} [\{\hat{f}(t), \hat{\mu}(t)\}, \hat{Q}], \\
\hat{\rho}(t) &= \langle \hat{\mu}(t) \rangle := \text{Tr}_B(\hat{\mu}(t)\hat{\rho}_{\text{th}}), \\
\hat{A}\hat{\mu}(t) &= -\frac{i}{\hbar} [\hat{H}_S, \hat{\mu}(t)] + \frac{i}{2M\hbar} [\{\gamma \text{sgn}(t)\hat{P}, \hat{\mu}(t)\}, \hat{Q}], \\
\hat{B}\hat{\mu}(t) &= \frac{i}{\hbar} [\hat{\mu}(t), \hat{Q}], \\
\hat{\alpha}(t)\hat{\mu}(t) &= \frac{1}{2} \{\hat{f}(t), \hat{\mu}(t)\}, \\
\hat{\eta}(t) &= \exp(-\hat{A}t)\hat{\mu}(t). \\
\dot{\hat{\eta}}(t) &= \hat{B}(t)\hat{\alpha}(t)\hat{\eta}(t), \\
\frac{d}{dt}\langle \hat{\eta}(t) \rangle &= \hat{B}(t)\langle \hat{\alpha}(t)\hat{\mu}(0) \rangle + \int_0^t \hat{B}(t)\hat{B}(t')\langle \hat{\alpha}(t)\hat{\alpha}(t')\hat{\mu}(t') \rangle dt' \\
\frac{d\hat{\rho}}{dt}(t) &= \hat{A}\hat{\rho}(t) + \int_0^t \langle \hat{\alpha}(t)\hat{\alpha}(t')\hat{f} \rangle \hat{B}\hat{B}(t'-t)\hat{\rho}(t')dt' \\
\frac{d\hat{\rho}}{dt}(t) &= \hat{A}\hat{\rho}(t) - \int_0^t \langle \hat{\alpha}(t)\hat{\alpha}(t')\hat{f} \rangle dt' \frac{1}{\hbar} [[\hat{\rho}(t), \hat{Q}], \hat{Q}].
\end{aligned}$$

$$\dot{\hat{\rho}}(t) = -\frac{i}{\hbar} [\hat{H}_S, \hat{\rho}(t)] + \frac{i}{2\hbar} [\{\gamma \text{sgn}(t) \hat{P}, \hat{\rho}(t)\}, \hat{Q}] - \frac{\Gamma(t)}{\hbar^2} [[\hat{\rho}(t), \hat{Q}], \hat{Q}],$$

$$\sum_{k=1}^{m_1} \Pr(1 \mid [1, k]) = 1 \qquad \frac{\sum_{k=1}^{m_{n-1}} \Pr(1 \mid [n-1, k]) = 1}{\sum_{k=1}^{m_n} \Pr(1 \mid [n, k]) = 1} = \begin{cases} 0, & \text{NCHV,} \\ 1, & \text{Q} \end{cases}$$

$$\sum_{i \in V} \Pr(1 \mid i) - \sum_{(i,j) \in E} \Pr(1,1 \mid i,j) \stackrel{\text{NCHV}}{\leqslant} \alpha(G) \stackrel{\text{Q}}{\leqslant} \vartheta(G)$$

$$\alpha(G)=n-1, \vartheta(G)=n, \text{ and } \chi(\bar{G})=n$$

$$p_1\!:=\!\sum_{k=1}^{19}\Pr(1\mid k)=1$$

$$p_2\!:=\!\sum_{k=20}^{38}\Pr(1\mid k)=1$$

$$p_3\!:=\!\sum_{k=39}^{57}\Pr(1\mid k)=\begin{cases} 0, & \text{NCHV} \\ 1, & \text{Q} \end{cases}$$

$$|\boldsymbol{a}\rangle = \begin{pmatrix} a_1 & a_2 & \cdots \\ & & a_d \end{pmatrix}^\dagger$$

$$\leftrightarrow \{|\alpha_1,\Delta t\rangle, |\alpha_2,2\Delta t\rangle, \cdots, |\alpha_d,d\Delta t\rangle\}$$

$$\langle j \mid \boldsymbol{a} \rangle = \sum_{k=1}^b \langle j_k \mid \boldsymbol{a} \rangle$$

$$p_1=0.9939(15), p_2=0.9980(2), p_3=0.9983(2)$$

$$\sum_{i \in V} \Pr(1 \mid i) = 2.9902(4)$$

$$\alpha(G) + \sum_{(i,j) \in E} \Pr(1,1 \mid i,j) = 2.651(4)$$

$$\max_{\mathbf{B}} \vartheta = \text{tr}(\mathbf{B} \mathbf{J})$$

$$\mathbf{B} \geqslant 0, \text{tr}(\mathbf{B}) = 1$$

$$B_{ij}=0, \forall (i,j) \in E(G)$$

$$h_0=\tilde{\alpha}\mathrm{Re}\left(e^{-i\phi}\sum_{k=1}c_ke^{ik\varphi}\right)$$

$$h_1=\tilde{\alpha}\mathrm{Re}\left(e^{-i\phi+\pi/2}\sum_{k=1}c_ke^{ik\varphi}\right)$$

$$h_2=\tilde{\alpha}\mathrm{Re}\left(e^{-i\phi+\pi/2}\sum_{k=4}c_ke^{ik\varphi}\right)$$

$$P(1,1 \mid i,j) = P(1 \mid i)P(1 \mid j, i = 1)$$

$$\hat{\mathcal{H}}^{(N)}/\hbar = \omega_0 \hat{a}^\dagger \hat{a} + \omega_{\rm a} \left(\hat{S}_z + \frac{N}{2} \right) + \frac{2g}{\sqrt{N}} (\hat{a}^\dagger + \hat{a}) \hat{S}_x$$



$$\begin{aligned}
g &> \frac{\sqrt{\omega_a \omega_0}}{2} \\
\hat{\mathcal{H}}_{\text{spin}} &= \hat{\mathcal{H}}_{\text{Fe}} + \hat{\mathcal{H}}_{\text{Er}} + \hat{\mathcal{H}}_{\text{Fe-Er}} \\
\hat{\mathcal{H}}_{\text{Fe}} &= \sum_{s=A, B} \sum_{i=1}^{N_0} \mu_0 \mu_B g_{\text{Fe}}^x S_{i,x}^s H_x^{\text{DC}} + J_{\text{Fe}} \sum_{i,i'} S_i^A \cdot \mathbf{S}_{i'}^B \\
&\quad - D_{\text{Fe}}^y \sum_{i,j'} (S_{i,z}^A S_{i',x}^B - S_{i',z}^B S_{i,x}^A) \\
&\quad - \sum_{s=A,B} \sum_{i=1}^{N_0} [A_{\text{Fe}}^x (S_{i,x}^s)^2 + A_{\text{Fe}}^z (S_{i,z}^s)^2], \\
\hat{\mathcal{H}}_{\text{Er}} &= \sum_{s=A, B} \sum_{i=1}^{N_0} \mu_0 \mu_B g_{\text{Er}}^x \alpha_{i,x}^s H_x^{\text{DC}} + J_{\text{Er}} \sum_{i,i'} \mathfrak{Z}_i^A \cdot \mathfrak{Z}_{i'}^B \\
&\quad - \sum_{s=A, B} \sum_{i=1}^{N_0} [A_{\text{Er}}^x (\mathfrak{Z}_{i,x}^s)^2 + A_{\text{Er}}^z (\mathfrak{Z}_{i,z}^s)^2], \\
\hat{\mathcal{H}}_{\text{Fe-Er}} &= \sum_{i=1}^{N_0} \sum_{s,s'=A, B} [J_i^s \cdot S_i^{s'} + D^{s,s'} \cdot (\mathfrak{Z}_i^s \times S_i^{s'})] \\
\langle \mathfrak{B}_{\parallel}^s \rangle &= -\frac{1}{2} \tanh \left(\frac{g_{\text{Er}} \mu_B |\bar{\mathbf{B}}_{\text{Er}}^s|}{2k_B T} \right) \\
\langle S_{\parallel}^s \rangle &= -B_s \left(\frac{S g_{\text{Fe}} \mu_B |\bar{\mathbf{B}}_{\text{Fe}}^s|}{k_B T} \right) \\
\hbar \frac{d}{dt} \delta \mathfrak{B}^s &= -\delta \mathfrak{s}^s \times g_{\text{Er}} \mu_B \bar{\mathbf{B}}_{\text{Er}}^s (g^{A/B}, \mathbf{S}^{A/B}) - \bar{\mathfrak{s}}^s \times g_{\text{Er}} \mu_B \mathbf{B}_{\text{Er}}^s (\delta \mathfrak{s}^{A/B}, \delta \mathbf{S}^{A/B}) \\
\hbar \frac{d}{dt} \delta \mathbf{S}^s &= -\delta \mathbf{S}^s \times g_{\text{Fe}} \mu_B \bar{\mathbf{B}}_{\text{Er}}^s (g^{A/B}, \mathbf{S}^{A/B}) - \bar{\mathbf{S}}^s \times g_{\text{Fe}} \mu_B \mathbf{B}_{\text{Er}}^s (\delta \mathfrak{s}^{A/B}, \delta \mathbf{S}^{A/B}) \\
\frac{\hat{\mathcal{H}}_{\text{Dicke}}}{\hbar} &\sim \omega_0 \hat{a}^\dagger \hat{a} + \omega_a \hat{\Sigma}_x^+ + \frac{i g_z}{\sqrt{N_0}} (\hat{a}^\dagger - \hat{a}) \hat{\Sigma}_z^- + \dots \\
\hbar \frac{d \mathfrak{s}^s}{dt} &= -\mathfrak{B}^s \times g_{\text{Er}} \mu_B \mathbf{B}_{\text{Er}}^s (g^{A/B}, \mathbf{S}^{A/B}) \\
\hbar \frac{d \mathbf{S}^s}{dt} &= -\mathbf{S}^s \times g_{\text{Fe}} \mu_B \mathbf{B}_{\text{Fe}}^s (g^{A/B}, \mathbf{S}^{A/B}) \\
\mathbf{B}_{\text{Er}}^s &= \mathbf{B}^{\text{DC}} + \frac{2 Z_{\text{Er}} J_{\text{Er}}}{\mu_B g_{\text{Er}}} \mathfrak{s}^{\bar{s}} + \frac{2}{\mu_B g_{\text{Er}}} \sum_{s'=A, B} [J \mathbf{S}^s - (\mathbf{D}^{s,s'} \times \mathbf{S}^s) - \mathbf{A}_{\text{Er}} \cdot \mathfrak{s}^s] \\
\mathbf{B}_{\text{Fe}}^s &= \mathbf{B}^{\text{DC}} + \frac{2 Z_{\text{Fe}} J_{\text{Fe}}}{\mu_B g_{\text{Fe}}} \mathbf{s}^{\bar{s}} - \frac{Z_{\text{Fe}}}{\mu_B g_{\text{Fe}}} \mathbf{D}_{\text{Fe}} \times \mathbf{s}^{\bar{s}} + \frac{2}{\mu_B g_{\text{Fe}}} \sum_{s'=A, B} [J \mathbf{S}^s - (\mathbf{D}^{s,s'} \times \mathbf{S}^s) - \mathbf{A}_{\text{Fe}} \cdot \mathfrak{s}^s] \\
\mathcal{H}_{\text{Er}}^s &= g_{\text{Er}} \mu_B \mathfrak{Z}^s \cdot \bar{\mathbf{B}}_{\text{Er}}^s = g_{\text{Er}} \mu_B \mathfrak{Z}_{\parallel}^s |\bar{\mathbf{B}}_{\text{Er}}^s| \\
\mathcal{H}_{\text{Fe}}^s &= g_{\text{Fe}} \mu_B \mathbf{S}^s \cdot \bar{\mathbf{B}}_{\text{Fe}}^s = g_{\text{Fe}} \mu_B S_{\parallel}^s |\bar{\mathbf{B}}_{\text{Fe}}^s|
\end{aligned}$$

$$\langle \mathfrak{h}_{\parallel}^s \rangle = -\frac{1}{2} \tanh \left(\frac{g_{\text{Er}} \mu_{\text{B}} |\overline{\mathbf{B}}_{\text{Er}}^s|}{2k_{\text{B}} T} \right)$$

$$\langle S_{\parallel}^s \rangle = -B_S \left(\frac{S g_{\text{Fe}} \mu_{\text{B}} |\overline{\mathbf{B}}_{\text{Fe}}^s|}{k_{\text{B}} T} \right)$$

$$\hbar \frac{d}{dt} \delta \mathfrak{s}^s = -\delta \mathfrak{s}^s \times g_{\text{Er}} \mu_{\text{B}} \overline{\mathbf{B}}_{\text{Er}}^s (g^{\text{A/B}}, \mathbf{s}^{\text{A/B}}) - \overline{\mathfrak{s}}^s \times g_{\text{Er}} \mu_{\text{B}} \mathbf{B}_{\text{Er}}^s (\delta \mathbf{3}^{\text{A/B}}, \delta \mathbf{s}^{\text{A/B}})$$

$$\hbar \frac{d}{dt} \delta \mathbf{S}^s = -\delta \mathbf{S}^s \times g_{\text{Fe}} \mu_{\text{B}} \overline{\mathbf{B}}_{\text{Er}}^s (\mathfrak{s}^{\text{A/B}}, \mathbf{s}^{\text{A/B}}) - \overline{\mathbf{S}}^s \times g_{\text{Fe}} \mu_{\text{B}} \mathbf{B}_{\text{Er}}^s (\delta \mathbf{3}^{\text{A/B}}, \delta \mathbf{s}^{\text{A/B}})$$

$$C = \sum_i \frac{w_i}{N_i} \sum_{j \in \text{ith mode}} \frac{(f_{ij} - \tilde{f}_{ij})^2}{\tilde{f}_{ij}^2}$$

$$\overline{\mathbf{S}}_0^{\text{A}} = \begin{pmatrix} S \sin \beta_0 \\ 0 \\ -S \cos \beta_0 \end{pmatrix}, \overline{\mathbf{S}}_0^{\text{B}} = \begin{pmatrix} S \sin \beta_0 \\ 0 \\ S \cos \beta_0 \end{pmatrix}$$

$$\beta_0 = -\frac{1}{2} \arctan \left[\frac{z_{\text{Fe}} D_{\text{Fe}}^y}{z_{\text{Fe}} J_{\text{Fe}} - A_{\text{Fe}}^x + A_{\text{Fe}}^z} \right]$$

$$\hat{\mathcal{H}}_{\text{Fe}} \approx \sum_{k=0,\pi} \hbar \omega_k a_k^\dagger a_k + \text{const}$$

$$\omega_k = g_{\text{Fe}} \mu_{\text{B}} / \hbar \sqrt{(b \cos k - a)(d \cos k + c)}$$

$$a = [S/(g_{\text{Fe}}^x \mu_{\text{B}})] [-A_{\text{Fe}}^z - A_{\text{Fe}}^x - (z_{\text{Fe}} J_{\text{Fe}} + A_{\text{Fe}}^z - A_{\text{Fe}}^x) \cos (2\beta_0) + z_{\text{Fe}} D_{\text{Fe}}^y \sin (2\beta_0)]$$

$$b = [S/(g_{\text{Fe}}^x \mu_{\text{B}})] z_{\text{Fe}} J_{\text{Fe}}$$

$$c = [S/(g_{\text{Fe}}^x \mu_{\text{B}})] [(z_{\text{Fe}} J_{\text{Fe}} + 2A_{\text{Fe}}^z - 2A_{\text{Fe}}^x) \cos (2\beta_0) + z_{\text{Fe}} D_{\text{Fe}}^y \sin (2\beta_0)]$$

$$d = [S/(g_{\text{Fe}}^x \mu_{\text{B}})] [-z_{\text{Fe}} J_{\text{Fe}} \cos (2\beta_0) - z_{\text{Fe}} D_{\text{Fe}}^y \sin (2\beta_0)]$$

$$\delta \mathbf{S}_i^{\text{A}} \approx \sqrt{\frac{S}{2N_0}} \begin{bmatrix} -(T_0 - T_{\pi}) \cos \beta_0 \\ (Y_0 - Y_{\pi}) \\ -(T_0 - T_{\pi}) \sin \beta_0 \end{bmatrix}$$

$$\delta \mathbf{S}_i^{\text{B}} \approx \sqrt{\frac{S}{2N_0}} \begin{bmatrix} (T_0 + T_{\pi}) \cos \beta_0 \\ (Y_0 + Y_{\pi}) \\ -(T_0 + T_{\pi}) \sin \beta_0 \end{bmatrix}$$

$$T_k = \left(\frac{b \cos k - a}{d \cos k - a} \right)^{1/4} \frac{a_{-k}^\dagger + a_k}{\sqrt{2}}$$

$$Y_k = \left(\frac{d \cos k + c}{b \cos k - a} \right)^{1/4} \frac{(a_{-k}^\dagger - a_k) \sqrt{2}}{\sqrt{2}}$$

$$\mathfrak{I}_i^s = \frac{1}{N_0} \sum_{j=1}^{N_0} \mathfrak{I}_j^s$$

$$\equiv \frac{1}{N_0} \Sigma_j^s$$

$$\begin{aligned}
\hat{\mathcal{H}}_{\text{Er}} &\approx \mu_0 \mu_{\text{B}} g_{\text{Er}}^x H_x^{\text{DC}} \Sigma_x^+ + z_{\text{Er}} J_{\text{Er}} \sum_{i=1}^{N_0} \mathfrak{s}_i^{\text{A}} \cdot \sum_{i'=1}^{N_0} \frac{\mathfrak{s}_{i'}^{\text{B}}}{N_0} \\
&\quad - \sum_i \sum_{\xi=x,z} A_{\text{Er}}^{\xi} \sum_{s=\text{A}, \text{B}} \left(\frac{1}{N_0} \sum_j 3_{j,\xi}^s \right)^2 \\
&= \mu_0 \mu_{\text{B}} g_{\text{Er}}^x H_x^{\text{DC}} \Sigma_x^+ + \frac{z_{\text{Er}} J_{\text{Er}}}{N_0} \Sigma^{\text{A}} \cdot \Sigma^{\text{B}} - \sum_{s=\text{A}, \text{B}} \sum_{\xi=x,z} \sum_{\text{Er}} \frac{A_{\text{Er}}^{\xi}}{N_0} (\Sigma_{\xi}^s)^2
\end{aligned}$$

$$\begin{aligned}
\hat{\mathcal{H}}_{\text{Fe-Er}} &= 4S(J \sin \beta_0 + D_y \cos \beta_0) \Sigma_x^+ + (-4SD_x \cos \beta_0) \Sigma_y^- \\
&\quad + \sqrt{\frac{S}{N_0}} [(J \cos \beta_0 - D_y \sin \beta_0) T_{\pi} \Sigma_x^+ + JY_0 \Sigma_y^+ \\
&\quad + (D_x \sin \beta_0) T_{\pi} \Sigma_y^- + D_x Y_{\pi} \Sigma_z^- - (J \sin \beta_0 + D_y \cos \beta_0) T_0 \Sigma_z^+]
\end{aligned}$$

$$\begin{aligned}
\hat{\mathcal{H}}_{\text{Dicke}} &\approx \sum_{m=\text{qFM}, \text{qAFM}} \hbar \omega_m a_m^{\dagger} a_m + E_x \Sigma_x^+ + E_y \Sigma_y^- + \mu_0 \mu_{\text{B}} g_{\text{Er}}^x H_x^{\text{DC}} \Sigma_x^+ \\
&\quad + \frac{z_{\text{Er}} J_{\text{Er}}}{N_0} \Sigma^{\text{A}} \cdot \Sigma^{\text{B}} - \sum_{\xi=x,z} \sum_{s=\text{A}, \text{B}} \sum_0 \frac{A_{\text{Er}}^{\xi}}{N_0} (\Sigma_{\xi}^s)^2 + \frac{\hbar g_x}{\sqrt{N_0}} (a_{\text{qAFM}}^{\dagger} + a_{\text{qAFM}}) \Sigma_x^+ \\
&\quad + \frac{i \hbar g_y}{\sqrt{N_0}} (a_{\text{qFM}}^{\dagger} - a_{\text{qFM}}) \Sigma_y^+ + \frac{\hbar g_{y'}}{\sqrt{N_0}} (a_{\text{qAFM}}^{\dagger} + a_{\text{qAFM}}) \Sigma_y^- \\
&\quad + \frac{i \hbar g_z}{\sqrt{N_0}} (a_{\text{qAFM}}^{\dagger} - a_{\text{qAFM}}) \Sigma_z^- + \frac{\hbar g_{z'}}{\sqrt{N_0}} (a_{\text{qFM}}^{\dagger} + a_{\text{qFM}}) \Sigma_z^+ \\
&\quad \hbar g_x = \sqrt{xS} (J \cos \beta_0 - D_y \sin \beta_0) \left(\frac{b+a}{d-c} \right)^{1/4} = h\sqrt{x} \\
&\quad \hbar g_y = \sqrt{xS} J \left(\frac{d+c}{b-a} \right)^{1/4} = h\sqrt{x} \\
&\quad \hbar g_{y'} = \sqrt{xS} D_x \sin \beta_0 \left(\frac{b+a}{d-c} \right)^{1/4} = h\sqrt{x} \\
&\quad \hbar g_z = \sqrt{xS} D_x \left(\frac{d-c}{b+a} \right)^{1/4} = h\sqrt{x} \\
&\quad \hbar g_{z'} = \sqrt{xS} (-J \sin \beta_0 - D_y \cos \beta_0) \left(\frac{b-a}{d+c} \right)^{1/4} = h\sqrt{x} \\
&\quad x = \tanh \left[\frac{|E_x + \mu_0 \mu_{\text{B}} g_{\text{Er}}^x H_x^{\text{DC}}|}{2k_{\text{B}} T} \right] \\
&\quad \varepsilon \propto \frac{2\pi}{\hbar} \sum_{\alpha} \left| \sum_{(v_i, c_i, i=1,2,3)} \langle e_3 h_3 | \mathbf{e} \cdot \hat{\mathbf{p}} | e_1 h_1 e_2 h_2 \rangle_{\alpha} \right|^2 \Gamma
\end{aligned}$$

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