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GRAVEDAD CUÁNTICA RELATIVISTA

RELATIVISTIC QUANTUM GRAVITY

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Gravedad cuántica relativista

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RESUMEN

La teoría cuántica de campos relativistas o curvos, esbozada por este autor en trabajos anteriores, define la gravedad, como una característica fenomenológica de las llamadas “partículas supermasivas”. Este efecto gravitacional cuántico, puede ser ordinario o extraordinario, según sea el caso. Cabe precisar, que el efecto gravitacional ordinario, es el resultado de la deformación del espacio – tiempo cuántico, sin causar pliegues de campo, es decir, distintas capas geométricas y en D – dimensiones, conservándose la misma dimensión en \mathbb{R}^4 , en tanto que, el efecto gravitacional extraordinario, es el resultado de la deformación con ruptura de simetría del espacio – tiempo cuántico, a razón de la existencia de un agujero negro cuántico en dimensión D, es decir, en multidimensiones, por lo que, el tejido del espacio – tiempo se hipergeometriza, en una superficie de Riemann y más concretamente, en un espacio superplanckiano. Sin embargo, en este artículo, nos ocuparemos específicamente del efecto gravitacional ordinario, el mismo que puede ser endógeno o exógeno, según sea el caso, es decir, causado por la propia partícula supermasiva a propósito de su masa o por la interacción de una partícula con el gravitón, esto es, por la permeabilización de un espacio – tiempo cuántico por un campo gravitónico de gauge y local.

Palabras clave: gravedad cuántica, espacio de Hilbert – Einstein, Dimensión en \mathbb{R}^4 , efecto gravitacional ordinario.

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Relativistic quantum gravity

ABSTRACT

The quantum theory of relativistic or curved fields, outlined by this author in previous works, defines gravity as a phenomenological characteristic of the so-called "supermassive particles". This quantum gravitational effect can be ordinary or extraordinary, as the case may be. It should be noted that the ordinary gravitational effect is the result of the deformation of quantum space-time, without causing field folds, that is, different geometric and D-dimensional layers, preserving the same dimension in \mathbb{R}^4 , while the extraordinary gravitational effect is the result of the deformation with symmetry breaking of quantum space-time, due to the existence of a quantum black hole in dimension D, that is, in multidimensions, so that the fabric of space-time is hypergeometrized, on a Riemann surface and more specifically, in a superplanckian space. However, in this article, we will deal specifically with the ordinary gravitational effect, which can be endogenous or exogenous, as the case may be, that is, caused by the supermassive particle itself in terms of its mass or by the interaction of a particle with the graviton, that is, by the permeabilization of a quantum space-time by a gauge and local gravitonic field.

Keywords: quantum gravity, Hilbert-Einstein space, Dimension in, ordinary gravitational effect. \mathbb{R}^4

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INTRODUCCIÓN

En este trabajo, nos proponemos demostrar el comportamiento de un espacio – tiempo cuántico por efecto gravitacional ordinario. Para estos efectos, consideramos el propagador de una partícula supermasiva, el mismo que se desplaza en una superficie de Riemann y por ende, en un espacio de Hilbert – Einstein, en el que, la partícula supermasiva, a propósito de la gravedad causada, deforma el espacio – tiempo cuántico, provocando un rompimiento de simetría de gauge susceptible de reparación, afectando por ende, las trayectorias inicial y final de las partículas repercutidas, su momento angular y las líneas geométricas de colisión y aniquilación. Asimismo, se considera la deformación del espacio – tiempo cuántico, cuando la energía – stress – momentum de una partícula supermasiva, curva el espacio – tiempo cuántico, esto a propósito de la brecha de masa, que supera el estado de vacío, es decir, cuando el salto de energía es superior a cero. Nos ocuparemos también de las hiperpartículas, esto es, de aquellas partículas que aunque tengan o no tengan masa, deforman el espacio – tiempo cuántico, específicamente cuando superan la velocidad de la luz, liberando cargas de energía infinitas. En este trabajo, no nos ocuparemos de la partícula oscura ni de la partícula estrella respectivamente, pues aquellas, son propias de la supergravedad cuántica.

En este trabajo, se intenta extraer las ecuaciones de campo de Einstein a un espacio – tiempo cuántico específico, en dimensión \mathbb{R}^4 .

RESULTADOS Y DISCUSIÓN.

En este apartado, pasamos a diseñar el modelo matemático de gravedad cuántica relativo a un espacio – tiempo cuántico relativista o curvo.



Cálculos preliminares de un campo de gauge y sus transformaciones e invariancias, respecto de un espacio planckiano.

$$X^\mu(\tau, \sigma) \sim x^\mu + p^\mu \tau + \frac{i}{\sqrt{2}} \sum_{n \in \mathbb{Z}^*} \frac{1}{n} (\alpha_n^\mu e^{-in(\tau-\sigma)} + \bar{\alpha}_n^\mu e^{-in(\tau+\sigma)}),$$

$$[x^\mu, p^\nu] = i\eta^{\mu\nu}$$

$$[\alpha_m^\mu, \alpha_n^\nu] = m\eta^{\mu\nu}\delta_{m+n,0}$$

$$\begin{aligned} H_{\text{closed}} &= -\frac{m^2}{2} + N + \bar{N} - 2 \\ H_{\text{open}} &= -m^2 + N - 1 \end{aligned}$$

$$\begin{aligned} N &= \sum_{n \in \mathbb{N}} n N_n, \quad N_n = \frac{1}{n} \alpha_{-n} \cdot \alpha_n, \\ \bar{N} &= \sum_{n \in \mathbb{N}} n \bar{N}_n, \quad \bar{N}_n = \frac{1}{n} \bar{\alpha}_{-n} \cdot \bar{\alpha}_n. \end{aligned}$$

$$H|\psi\rangle = 0$$

$$(N - \bar{N})|\psi\rangle = 0$$

$$p^\mu|k\rangle = k^\mu|k\rangle, \forall n > 0: \alpha_n^\mu|k\rangle = 0$$

$$|\psi\rangle = \prod_{n>0} \prod_{\mu=0}^{D-1} (\alpha_{-n}^\mu)^{N_{n,\mu}} |k\rangle,$$

$$\text{closed : } m^2 = -4, \text{ open : } m^2 = -1$$

$$\alpha_{-1}^\mu|k\rangle$$

$$m^2 = 0$$

$$|A\rangle = \int d^D k A_\mu(k) \alpha_{-1}^\mu |k\rangle$$

$$k^2 A_\mu = 0$$

$$k^\mu A_\mu = 0$$

$$A_\mu \rightarrow A_\mu + k_\mu \lambda$$

$$\alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu |k\rangle$$

$$m^2 = 0$$



$$\left(\alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu + \alpha_{-1}^\nu \bar{\alpha}_{-1}^\mu - \frac{1}{D} \eta^{\mu\nu} \alpha_{-1} \cdot \bar{\alpha}_{-1} \right) |p\rangle$$

$$(\alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu - \alpha_{-1}^\nu \bar{\alpha}_{-1}^\mu) |p\rangle, \frac{1}{D} \eta_{\mu\nu} \alpha_{-1}^\mu \bar{\alpha}_{-1}^\nu |p\rangle$$

$$| \text{ boson } \rangle = Q | \text{ fermion } \rangle.$$

$$D(N=0) = 26, D(N=1) = 10$$

Superficies de Riemann y amplitudes en base a la función de Green.

$$\chi = 2 - 2g - b$$

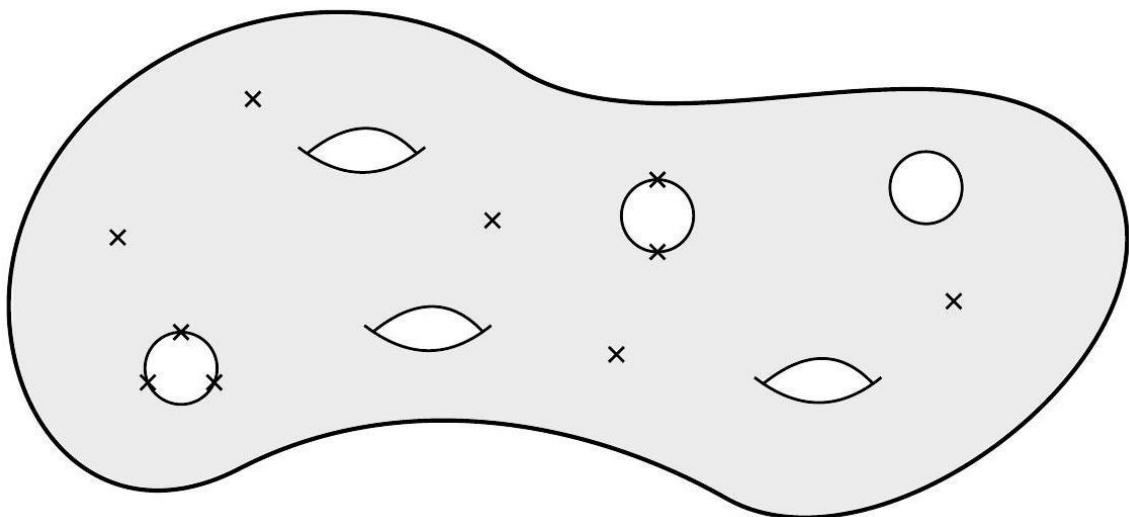


Figura 1. Modelamiento de curvaturas locales.

$$A_n(k_1, \dots, k_n) = \sum_{g \geq 0} g_s^{n-2+2g} A_{g,n}$$

$$A_{g,n} = \int \prod_{i=1}^n dz_i \int dg_{ab} d\Psi e^{-S_{\text{cft}}[g_{ab}, \Psi]} \prod_{i=1}^n V_i(k_i, z_i)$$

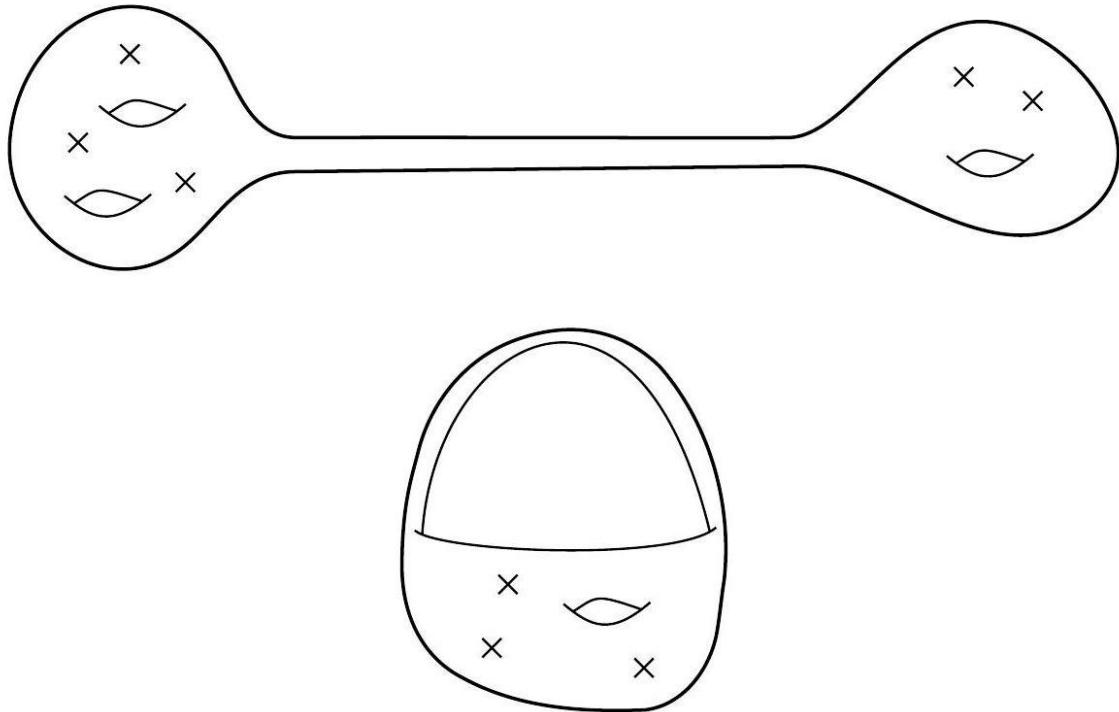
$$\dim_{\mathbb{R}} \mathcal{M}_{g,n} = 6g - 6 + 2n.$$

$$A_{g,n} = \int_{\mathcal{M}_{g,n}} \prod_{\lambda=1}^{6g-6+2n} dt_\lambda F(t)$$

$$F(t) = \left\langle \prod_{i=1}^n V_i \times \text{ghosts} \times \text{super-ghosts} \right\rangle_{\Sigma_{g,n}}$$

Divergencias de Feynman – acción de los osciladores y propagadores de campo.

$$\frac{1}{k^2 + m^2} = \int_0^\infty ds e^{-s(k^2 + m^2)}$$



Figuras 2 y 3. Curvatura bilateral y curvatura local. La gravedad cuántica ocupa la segunda demostración, en tanto que la supergravedad cuántica, ocupa la primera demostración.

$$S_0 = \frac{1}{2} K_\Psi(\Psi, \Psi) + \frac{1}{2} K_\Phi(\Phi, \Phi)$$

$$S_{\text{int}} = \sum_{m,n} \mathcal{V}_{m,n}(\Phi^m, \Psi^n)$$

$$|\Psi\rangle = \sum_n \int \frac{d^D k}{(2\pi)^D} \psi_\alpha(k) |k, \alpha\rangle$$

$$\psi_\alpha = \{T, G_{\mu\nu}, B_{\mu\nu}, \Phi, \dots\}.$$

Simetrías difeomorfistas y tensores de simetría – Acción Nambu – Goto – Polyakov – Regge – Weyl – DeWitt.

$$\sigma \in [0,2\pi), \sigma \sim \sigma + 2\pi.$$

$$S_{\text{NG}}[X^\mu] = \frac{1}{2\pi\alpha'} \int d^2\sigma \sqrt{\det G_{\mu\nu}(X)} \frac{\partial X^\mu}{\partial \sigma^a} \frac{\partial X^\nu}{\partial \sigma^b}$$

$$S_{\text{P}}[g, X^\mu] = \frac{1}{4\pi\alpha'} \int d^2\sigma \sqrt{g} g^{ab} G_{\mu\nu}(X) \frac{\partial X^\mu}{\partial \sigma^a} \frac{\partial X^\nu}{\partial \sigma^b}$$

$$\chi_g := \chi(\Sigma_g) := 2 - 2g = \frac{1}{4\pi} \int_{\Sigma_g} d^2\sigma \sqrt{g} R$$

$$\sigma'^a = f^a(\sigma^b), g'(\sigma') = f^*g(\sigma), \Psi'(\sigma') = f^*\Psi(\sigma)$$

$$g'_{ab}(\sigma') = \frac{\partial \sigma^c}{\partial \sigma'^a} \frac{\partial \sigma^d}{\partial \sigma'^b} g_{cd}(\sigma), X'^\mu(\sigma') = X^\mu(\sigma)$$

$$\delta_\xi \sigma^a = \xi^a, \delta_\xi \Psi = \mathcal{L}_\xi \Psi, \delta_\xi g_{ab} = \mathcal{L}_\xi g_{ab}$$

$$\mathcal{L}_\xi g_{ab} = \xi^c \partial_c g_{ab} + g_{ac} \partial_b \xi^c + g_{bc} \partial_a \xi^c = \nabla_a \xi_b + \nabla_b \xi_a.$$

$$\Gamma_g := \pi_0\left(\text{Diff}(\Sigma_g)\right) = \frac{\text{Diff}(\Sigma_g)}{\text{Diff}_0(\Sigma_g)}$$

$$g'_{ab}(\sigma) = e^{2\omega(\sigma)} g_{ab}(\sigma), \Psi'(\sigma) = \Psi(\sigma)$$

$$\delta_\omega g_{ab} = 2\omega g_{ab}, \delta_\omega \Psi = 0$$

$$\text{Conf}(\Sigma_g) := \frac{\text{Met}(\Sigma_g)}{\text{Weyl}(\Sigma_g)}$$

$$G := \text{Diff}(\Sigma_g) \ltimes \text{Weyl}(\Sigma_g).$$

$$G_0 := \text{Diff}_0(\Sigma_g) \ltimes \text{Weyl}(\Sigma_g).$$

$$g' = f^*(e^{2\omega} g) = e^{2f^*\omega} f^* g$$

$$g_{ab}(\sigma) = e^{2\phi(\sigma)} \hat{g}_{ab}(\sigma)$$

$$\hat{g}_{ab} = \delta_{ab}, \phi$$

$$\hat{g}_{ab} = \delta_{ab}, \phi = 0$$

$$S_m[f^*g, f^*\Psi] = S_m[g, \Psi], S_m[e^{2\omega}g, \Psi] = S_m[g, \Psi]$$

$$X^\mu(\tau, \sigma) \sim X^\mu(\tau, \sigma + 2\pi)$$



$$c_m=D+c_\perp.$$

$$\mathbb{M}_{\text{matter CFT parameters}} \colon D, c_m.$$

$$T_{m,ab}\colon=-\frac{4\pi}{\sqrt{g}}\frac{\delta S_m}{\delta g^{ab}}$$

$$\nabla^a T_{m,ab}=0\,\,\,(\text{on-shell}).$$

$$g^{ab}T_{m,ab}=0\,\,\,(\text{off-shell})$$

$$P^a\colon=\int\,\,\mathrm{d}\sigma T_m^{0a}$$

$$g^{ab}T_{m,ab}\propto R$$

$$Z_g\colon=\int\,\,\frac{\mathrm{d}_gg_{ab}}{\Omega_{\text{gauge}}\,[g]}Z_m[g],Z_m[g]\colon=\int\,\,\mathrm{d}_g\Psi\mathrm{e}^{-S_m[g,\Psi]}$$

$$(\delta\Phi_1,\delta\Phi_2)_g\colon=\int\,\,\mathrm{d}^2\sigma\sqrt{g}\gamma_g(\delta\Phi_1,\delta\Phi_2),|\delta\Phi|_g^2\colon=(\delta\Phi,\delta\Phi)_g$$

$$\int\,\,\mathrm{d}_g\delta\Phi\mathrm{e}^{-\frac{1}{2}(\delta\Phi,\delta\Phi)_g}=\frac{1}{\sqrt{\det\gamma_g}}$$

$$\int\,\,\mathrm{d}\Phi\,\sqrt{\det\gamma_g}$$

$$\int\,\,\mathrm{d}_g\delta\Phi\mathrm{e}^{-\frac{1}{2}(\delta\Phi,\delta\Phi)_g}=1$$

$$\Phi(\sigma)\longrightarrow\Phi'(\sigma)=\Phi(\sigma)+\varepsilon(\sigma)$$

$$\begin{aligned} (\delta f,\delta f)_g &:=\int\,\,\mathrm{d}^2\sigma\sqrt{g}\delta f^2 \\ (\delta V^a,\delta V^a)_g &:=\int\,\,\mathrm{d}^2\sigma\sqrt{g}g_{ab}\delta V^a\delta V^b \\ (\delta T_{ab},\delta T_{ab})_g &:=\int\,\,\mathrm{d}^2\sigma\sqrt{g}G^{abcd}\delta T_{ab}\delta T_{cd} \end{aligned}$$

$$G^{abcd}\colon=G_\perp^{abcd}+ug^{ab}g^{cd}, G_\perp^{abcd}\colon=g^{ac}g^{bd}+g^{ad}g^{bc}-g^{ab}g^{cd}$$

$$G^{abcd}T_{cd}=G_\perp^{abcd}T_{cd}=2T_{ab}, G^{abcd}(\Lambda g_{cd})=2u(\Lambda g_{ab})$$

$$\delta g_{ab}=g_{ab}\delta\Lambda+\delta g^\perp_{ab}, \delta\Lambda=\frac{1}{2}g^{ab}\delta g_{ab}, g^{ab}\delta g^\perp_{ab}=0$$

$$\left|\delta g_{ab}\right|^2_g=4u\left|\delta\Lambda\right|^2_g+\left|\delta g^\perp_{\mu\nu}\right|^2_g$$

$$u>0$$

$$\mathrm{d}_g g_{ab} = \mathrm{d}_g \Lambda \,\mathrm{d}_g g_{ab}^\perp$$



$$\begin{aligned} G^{abcd}\delta g_{ab}\delta g_{cd} &= (G_\perp^{abcd} + ug^{ab}g^{cd})(g_{ab}\delta\Lambda + \delta g_{ab}^\perp)(g_{cd}\delta\Lambda + \delta g_{cd}^\perp) \\ &= (2ug^{cd}\delta\Lambda + G_\perp^{abcd}\delta g_{ab}^\perp)(g_{cd}\delta\Lambda + \delta g_{cd}^\perp) \\ &= 4u(\delta\Lambda)^2 + G_\perp^{abcd}\delta g_{ab}^\perp\delta g_{cd}^\perp \\ &= 4u\delta\Lambda^2 + 2g^{ac}g^{bd}\delta g_{ab}^\perp\delta g_{cd}^\perp \end{aligned}$$

$$G^{abcd}=g^{ac}g^{bd}+cg^{ab}g^{cd}$$

$$(\delta X^\mu,\delta X^\mu)_g=\int\mathrm{~d}^2\sigma\sqrt{g}G_{\mu\nu}(X)\delta X^\mu\delta X^\nu$$

Reparación de gauge Faddeev-Popov en superficies de Riemann – Roch. Acción de Moduli.

Métrica de David-Distler-Kawai y acción de Liouville - Faddeev-Popov.

$$(f^*g,f^*\Psi)\sim(g,\Psi),(\mathrm{e}^{2\omega}g,\Psi)\sim(g,\Psi)$$

$$\mathrm{d}(\text{ fields }) = \text{ Jacobian } \times \mathrm{d}(\text{ gauge }) \times \mathrm{d}(\text{ physical }).$$

$$Z=\int_{\mathbb{R}^2}\mathrm{~d}x\mathrm{~d}y\mathrm{e}^{-(x-y)^2}$$

$$r=x-y,y=a$$

$$Z=\int_{\mathbb{R}}\mathrm{~d}a\int_0^\infty\mathrm{~e}^{-r^2}=\frac{\sqrt{\pi}}{2}\mathrm{Vol}(\mathbb{R})$$

$$\mathcal{M}_g:=\frac{\mathrm{Met}(\Sigma_g)}{G}$$

$$\mathcal{T}_g:=\frac{\mathrm{Met}(\Sigma_g)}{G_0}$$

$$\mathcal{M}_g=\frac{\mathcal{T}_g}{\Gamma_g}$$

$$M_g:=\dim_{\mathbb{R}}\mathcal{M}_g=\dim_{\mathbb{R}}\mathcal{T}_g=\begin{cases} 0 & g=0 \\ 2 & g=1 \\ 6g-6 & g\geq 2 \end{cases}$$

$$\int_{\mathcal{M}_g}\mathrm{~d}^{\mathbf{M}_g}t=\frac{1}{\Omega_{\Gamma_g}}\int_{\mathcal{T}_g}\mathrm{~d}^{\mathbf{M}_g}t$$

$$\Omega_G:=\int_G\mathrm{~d}g=\int_G\mathrm{~d}(hg),$$

$$\alpha=\alpha^iT_i.$$

$$\Omega_G=\int\mathrm{~d}\alpha:=\int\prod_i\mathrm{~d}\alpha^i$$

$$g_{ab}=\hat{g}_{ab}^{(f,\phi)}(t):=\mathrm{e}^{2f^*\phi}f^*\hat{g}_{ab}(t)=f^*\big(\mathrm{e}^{2\phi}\hat{g}_{ab}(t)\big)$$



$$g_{ab}(\sigma)=\hat{g}_{ab}^{(f,\phi)}(\sigma;t)\!:=\mathrm{e}^{2\phi(\sigma)}\hat{g}'_{ab}(\sigma;t), \hat{g}'_{ab}(\sigma;t)=\frac{\partial\sigma'^c}{\partial\sigma^a}\frac{\partial\sigma'^d}{\partial\sigma^b}\hat{g}_{cd}(\sigma';t)$$

$$\delta g_{ab}=2\phi g_{ab}+\nabla_a\xi_b+\nabla_b\xi_a+\delta t_i\partial_ig_{ab}$$

$$\partial_i\!:=\frac{\partial}{\partial t_i}.$$

$$\Omega_{\text{Diff}_0}[g]:=\int\,\,\mathrm{d}_g\xi\\ \Omega_{\text{Weyl}}[g]:=\int\,\,\mathrm{d}_g\phi$$

$$\Omega_{\text{Diff}}\,[g]=\Omega_{\text{Diff}_0}[g]\Omega_{\Gamma_g}$$

$$\Omega_{\text{Diff}_0}[g]\!:=\Omega_{\text{Diff}_0}\!\left[\mathrm{e}^{2\phi}\hat{g}\right]=\Omega_{\text{Diff}_0}[\hat{g}],\Omega_{\text{Weyl}}[g]\!:=\Omega_{\text{Weyl}}\!\left[\mathcal{L}_\xi\hat{g}\right]=\Omega_{\text{Weyl}}[\hat{g}]$$

$$\Omega_{\text{gauge}}\,[g]\!:=\Omega_{\text{gauge}}\,[\hat{g}]$$

$$\Omega_{\text{Diff}_0}\!\left[\mathrm{e}^{2\phi}\hat{g}\right]=\int\,\,\mathrm{d}_{\mathrm{e}^{2\phi}\mathcal{L}_\xi\hat{g}}\xi=\int\,\,\mathrm{d}_{\mathrm{e}^{2\phi}\hat{g}}\xi=\int\,\,\mathrm{d}_{\hat{g}}\xi=\Omega_{\text{Diff}_0}[\hat{g}]\\\Omega_{\text{Weyl}}\!\left[\mathcal{L}_\xi\hat{g}\right]=\int\,\,\mathrm{d}_{\mathrm{e}^{2\phi}\mathcal{L}_\xi\hat{g}}\phi=\int\,\,\mathrm{d}_{\mathrm{e}^{2\phi}\hat{g}}\phi=\Omega_{\text{Weyl}}[\hat{g}]$$

$$|\delta\phi|^2=\int\,\,\mathrm{d}^2\sigma\sqrt{g}\delta\phi^2=\int\,\,\mathrm{d}^2\sigma\sqrt{\hat{g}}\mathrm{e}^{2\phi}\delta\phi^2$$

$$F_{ab}\!:=g_{ab}-\hat{g}_{ab}^{(f,\phi)}(t),$$

$$\delta g_{ab}=2\tilde{\Lambda}g_{ab}+(P_1\xi)_{ab}+\delta t_i\mu_{iab}$$

$$(P_1\xi)_{ab}=\nabla_a\xi_b+\nabla_b\xi_a-g_{ab}\nabla_c\xi^c,\\\mu_{iab}=\partial_ig_{ab}-\frac{1}{2}g_{ab}g^{cd}\partial_ig_{cd},\\\tilde{\Lambda}=\Lambda+\frac{1}{2}\delta t_ig^{ab}\partial_ig_{ab},\Lambda=\phi+\frac{1}{2}\nabla_c\xi^c.$$

$$Z_g=\int\,\,\mathrm{d}^{\mathbf{M}_g}t\,\,\mathrm{d}_g\tilde{\Lambda}\,\,\mathrm{d}_g(P_1\xi)\Omega_{\text{gauge}}\,[g]^{-1}Z_m[g]$$

$$(P_1\xi,\tilde{\Lambda})\rightarrow (\xi,\phi)$$

$$\mathrm{d}_g(P_1\xi)\mathrm{d}_g\tilde{\Lambda}\stackrel{?}{=}\mathrm{d}_g\xi\,\mathrm{d}_g\phi\Delta_{\text{FP}}[g]$$

$$\Delta_{\text{FP}}[g]=\det\frac{\partial(P_1\xi,\tilde{\Lambda})}{\partial(\xi,\phi)}=\det\begin{pmatrix} P_1 & 0 \\ \star & 1 \end{pmatrix}=\det P_1$$

$$(T,P_1v)_g=\left(P_1^\dagger T,v\right)_g,$$

$$\left(P_1^\dagger T\right)_a=-2\nabla^bT_{ab}$$



$$\dim \ker P_1^\dagger - \dim \ker P_1 = -3\chi_g = 6g-6$$

Deformaciones de Teichmüller.

$$(\delta g, P_1 \xi)_g = 0 \implies (P_1^\dagger \delta g, \xi)_g = 0$$

$$P_1^\dagger \delta g = 0$$

$$\delta g \in \ker P_1^\dagger$$

$$\ker P_1^\dagger = \text{Span}\{\phi_i\}, i=1,\dots, \dim \ker P_1^\dagger$$

$$\dim_{\mathbb{R}} \ker P_1^\dagger = M_g = \begin{cases} 0 & g=0 \\ 2 & g=1 \\ 6g-6 & g>1 \end{cases}$$

$$\Pi := P_1 \frac{1}{P_1^\dagger P_1} P_1^\dagger$$

$$\delta t_i \mu_i = \delta t_i (1-\Pi) \mu_i + \delta t_i \Pi \mu_i = \delta t_i (1-\Pi) \mu_i + \delta t_i P_1 \zeta_i.$$

$$\zeta_i := \frac{1}{P_1^\dagger P_1} P_1^\dagger \mu_i.$$

$$(1-\Pi)\mu_i = \phi_j(M^{-1})_{jk}(\phi_k,\mu_i)_g$$

$$M_{ij} := (\phi_i, \phi_j)_g.$$

$$\delta g_{ab} = (P_1 \tilde{\xi})_{ab} + 2\tilde{\Lambda} g_{ab} + Q_{iab} \delta t_i.$$

$$\tilde{\xi} = \xi + \zeta_i \delta t_i, Q_{iab} = \phi_{jab}(M^{-1})_{jk}(\phi_k,\mu_i)_g.$$

$$|\delta g|_g^2 = |\delta \tilde{\Lambda}|_g^2 + |P_1 \tilde{\xi}|_g^2 + |Q_i \delta t_i|_g^2$$

$$\mathrm{d}_g g_{ab} = \mathrm{d}_g \tilde{\Lambda} \, \mathrm{d}_g (P_1 \tilde{\xi}) \mathrm{d}_g (Q_i \delta t_i).$$

$$(\tilde{\xi}, \tilde{\Lambda}, Q_i \delta t_i) \longrightarrow (\xi, \Lambda, \delta t_i),$$

$$\mathrm{d}_g \tilde{\Lambda} \, \mathrm{d}_g (P_1 \tilde{\xi}) \mathrm{d}_g (Q_i \delta t_i) = \mathrm{d}^{M_g} t \, \mathrm{d}_g \Lambda \, \mathrm{d}_g (P_1 \xi) \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}},$$

$$Z_g = \int_{\mathcal{T}_g} \mathrm{d}^{M_g} t \frac{1}{\Omega_{\text{gauge}}[\hat{g}]} \int \mathrm{d}_g \Lambda \, \mathrm{d}_g (P_1 \xi) \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} Z_m[g]$$



Vectores Conformal Killing.

$$\int \mathrm{d}_g \xi = \Omega_{\mathrm{Diff}_0}[\hat{g}]$$

$$\xi^{(0)} \in \mathcal{K}_g := \ker P_1$$

$$(P_1 \xi^{(0)})_{ab} = \nabla_a \xi_b^{(0)} + \nabla_b \xi_a^{(0)} - g_{ab} \nabla_c \xi^{(0)c} = 0$$

$$K_g := \dim_{\mathbb{R}} \mathcal{K}_g = \dim_{\mathbb{R}} \ker P_1 = \begin{cases} 6 & g=0 \\ 2 & g=1 \\ 0 & g>1 \end{cases}$$

$$g=0: \mathcal{K}_0=\mathrm{SL}(2,\mathbb{C}), g=1: \mathcal{K}_1=\mathrm{U}(1)\times \mathrm{U}(1).$$

$$\xi = \xi^{(0)} + \xi'$$

$$(\xi^{(0)}, \xi')_g = 0$$

$$(P_1 \xi, \Lambda) \rightarrow (\xi', \phi).$$

$$\mathrm{d}_g \Lambda \mathrm{d}_g (P_1 \xi) = \mathrm{d}_g \phi \mathrm{d}_g \xi' \Delta_{\mathrm{FP}}[g]$$

$$\Delta_{\mathrm{FP}}[g] = \det' \frac{\partial(P_1 \xi, \Lambda)}{\partial(\xi', \phi)} = \det' P_1 = \sqrt{\det' P_1 P_1^\dagger}$$

$$Z_g = \int_{\mathcal{T}_g} \mathrm{d}^M g t \Omega_{\text{gauge}} [\hat{g}]^{-1} \int \mathrm{d}_g \phi \mathrm{d}_g \xi' \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)_g}} \Delta_{\mathrm{FP}}[g] Z_m[g]$$

$$\Delta_{\mathrm{FP}}[g] = \det' \frac{\partial(P_1 \xi, \Lambda)}{\partial(\xi', \phi)} = \det' \begin{pmatrix} P_1 & 0 \\ \frac{1}{2} \nabla & 1 \end{pmatrix} = \det' P_1$$

$$\sqrt{\det' P_1^\dagger P_1} = \det' P_1$$

$$\begin{aligned} 1 &= \int \mathrm{d}_g \delta \Lambda \mathrm{d}_g (P_1 \delta \xi) e^{-|\delta \Lambda|_g^2 - |P_1 \delta \xi'|_g^2} \\ &= \Delta_{\mathrm{FP}}[g] \int \mathrm{d}_g \delta \phi \mathrm{d}_g \delta \xi' e^{-|\delta \phi + \frac{1}{2} \nabla_c \delta \xi|^2_g - |P_1 \delta \xi'|_g^2} \\ &= \Delta_{\mathrm{FP}}[g] \int \mathrm{d}_g \delta \phi \mathrm{d}_g \delta \xi' e^{-|\delta \phi|_g^2 - (\delta \xi', P_1^\dagger P_1 \delta \xi')_g} \\ &= \Delta_{\mathrm{FP}}[g] (\det' P_1^\dagger P_1)^{-1/2} \end{aligned}$$

$$\Omega'_{\mathrm{Diff}_0}[g] := \Omega'_{\mathrm{Diff}_0}[\hat{g}] = \int \mathrm{d}_g \xi'$$

$$\ker P_1 = \mathrm{Span}\{\psi_i\}, i=1, \dots, K_g$$



$$\xi' \rightarrow \xi$$

$${\rm d}_g \xi' = \frac{1}{\sqrt{\det(\psi_i,\psi_j)}_g} \frac{{\rm d}_g \xi}{\Omega_{\text{ckv}}[g]},$$

$$\Omega_{\text{Diff}_0}[g] = \sqrt{\det(\psi_i,\psi_j)}_g \Omega_{\text{ckv}}[g] \Omega'_{\text{Diff}_0}[g]$$

$$\xi^{(0)}=\alpha_i\psi_i$$

$$\xi \rightarrow (\xi', \alpha_i).$$

$$\begin{aligned} 1 &= \int \, {\rm d}\xi \, {\rm e}^{-|\xi|_g^2} = J \int \, {\rm d}\xi^{(0)} {\rm d}\xi' {\rm e}^{-|\xi'|_g^2 - |\xi^{(0)}|_g^2} \\ &= J \int \, \prod_i \, {\rm d}\alpha_i {\rm e}^{-\alpha_i \alpha_j (\psi_i, \psi_j)}_g \int \, {\rm d}\xi' {\rm e}^{-|\xi'|_g^2} \\ &= J \left(\det(\psi_i, \psi_j)_g \right)^{-1/2} \end{aligned}$$

$${\rm d}\xi = \sqrt{\det(\psi_i,\psi_j)}_g \, {\rm d}\xi' \prod_i \, {\rm d}\alpha_i$$

$$\Omega_{\text{ckv}}[g] = \int \, \prod_i \, {\rm d}\alpha_i$$

$$Z_g = \int_{\mathcal{T}_g} \, {\rm d}^{\mathbf{M}_g} t \Omega_{\text{gauge}} \, [\hat{g}]^{-1} \int \, {\rm d}_g \phi \, {\rm d}_g \xi \, \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)}_g} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)}_g} \Delta_{\text{FP}}[g] Z_m[g]$$

$$(f^*\hat{g}, f^*\phi, f^*\Psi) \rightarrow (\hat{g}, \phi, \Psi)$$

$$Z_g = \int_{\mathcal{T}_g} \, {\rm d}^{\mathbf{M}_g} t \frac{\Omega_{\text{Diff}_0}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \int \, {\rm d}_g \phi \, \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)}_g} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)}_g} \Delta_{\text{FP}}[g] Z_m[g]$$

$$g_{ab} := g_{ab}^{(\phi)} = {\rm e}^{2\phi} \hat{g}_{ab}$$

$$Z_g = \frac{1}{\Omega_{\Gamma_g}} \int_{\mathcal{T}_g} \, {\rm d}^{\mathbf{M}_g} t \frac{\Omega_{\text{Diff}}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \int \, {\rm d}_g \phi \, \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)}_g} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)}_g} \Delta_{\text{FP}}[g] Z_m[g]$$

$$Z_g = \int_{\mathcal{M}_g} \, {\rm d}^{\mathbf{M}_g} t \frac{\Omega_{\text{Diff}}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \int \, {\rm d}_g \phi \, \frac{\det(\phi_i, \mu_j)_g}{\sqrt{\det(\phi_i, \phi_j)}_g} \frac{\Omega_{\text{ckv}}[g]^{-1}}{\sqrt{\det(\psi_i, \psi_j)}_g} \Delta_{\text{FP}}[g] Z_m[g]$$



Transformaciones de Weyl – difeoformismos gravitacionales y anomalías cuánticas – nodos tensoriales de campo – parámetros de Ricci.

$$\frac{\Delta_{\text{FP}}[e^{2\phi}\hat{g}]}{\sqrt{\det(\phi_i, \phi_j)_{e^{2\phi}\hat{g}}}} = e^{\frac{c_{\text{gh}}}{6}S_L[\hat{g}, \phi]} \frac{\Delta_{\text{FP}}[\hat{g}]}{\sqrt{\det(\hat{\phi}_i, \hat{\phi}_j)_{\hat{g}}}}$$

$$Z_m[e^{2\phi}\hat{g}] = e^{\frac{c_m}{6}S_L[\hat{g}, \phi]} Z_m[\hat{g}]$$

$$S_L[\hat{g}, \phi] := \frac{1}{4\pi} \int d^2\sigma \sqrt{\hat{g}} (\hat{g}^{ab} \partial_a \phi \partial_b \phi + \hat{R}\phi)$$

$$c_{\text{gh}} = -26$$

$$\det(\phi_i, \mu_j)_{e^{2\phi}\hat{g}} = \det(\hat{\phi}_i, \hat{\mu}_j)_{\hat{g}}, \det(\psi_i, \psi_j)_{e^{2\phi}\hat{g}} = \det(\psi_i, \psi_j)_{\hat{g}},$$

$$\Omega_{\text{ckv}}[e^{2\phi}\hat{g}] = \Omega_{\text{ckv}}[\hat{g}].$$

$$\langle g^{\mu\nu} T_{\mu\nu} \rangle = \frac{c}{12} R$$

$$Z_g = \int_{\mathcal{M}_g} d^M g t \frac{\Omega_{\text{Diff}}[\hat{g}]}{\Omega_{\text{gauge}}[\hat{g}]} \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \Delta_{\text{FP}}[\hat{g}] Z_m[\hat{g}] \int d_g \phi e^{-\frac{c_L}{6} S_L[\hat{g}, \phi]}$$

$$c_L := 26 - c_m$$

$$c_L = 0 \implies c_m = 26$$

$$c_\perp = 26 - D$$

$$\int d_g \phi = \Omega_{\text{Weyl}}[\hat{g}]$$

$$\Omega_{\text{gauge}}[\hat{g}] = \Omega_{\text{Diff}}[\hat{g}] \times \Omega_{\text{Weyl}}[\hat{g}]$$

$$Z_g = \int_{\mathcal{M}_g} d^M g t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \Delta_{\text{FP}}[\hat{g}] Z_m[\hat{g}]$$

Ambigüedades, ultralocalidad y constante cosmológica – Métrica Weil-Petersson.

$$S_\mu[g] = \int d^2\sigma \sqrt{g}$$

$$\sqrt{\det \gamma_g} = e^{-\mu_\gamma S_\mu[g]},$$

$$\Omega_\Phi = \lim_{\lambda \rightarrow 0} \int d_g \Phi e^{-\lambda(\Phi, \Phi)_g}$$



$$\int \mathrm{d}_g \Phi \mathrm{e}^{-\lambda(\Phi,\Phi)_g} = \mathrm{e}^{-\mu(\lambda) S_\mu[g]}$$

$$\Omega_\Phi=\int \mathrm{d}_g \Phi = \mathrm{e}^{-\mu(0) S_\mu[g]}$$

$$\lim_{\epsilon\rightarrow 0}\frac{1}{\epsilon}\int \mathrm{d}^2\sigma \sqrt{g}$$

$$\begin{aligned} Z_g &= \int_{\mathcal{M}_g} \mathrm{d}^{\mathrm{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \Delta_{\text{FP}}[\hat{g}] Z_m[\hat{g}] \\ &= \int_{\mathcal{M}_g} \mathrm{d}^{\mathrm{M}_g} t \sqrt{\frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}^2}{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\det' \hat{P}_1^\dagger \hat{P}_1}{\det(\psi_i, \psi_j)_{\hat{g}}} \frac{Z_m[\hat{g}]}{\Omega_{\text{ckv}}[\hat{g}]} \end{aligned}$$

$$\mathrm{d}(\text{WP}) = \int_{\mathcal{M}_g} \mathrm{d}^{\mathrm{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}}$$

$$\sigma'^a = \hat{f}^a(\sigma^b), \hat{g}'(\sigma') = f^* \hat{g}(\sigma), \phi'(\sigma') = f^* \phi(\sigma), \Psi'(\sigma') = f^* \Psi(\sigma)$$

$$g_{ab} = f^*(\mathrm{e}^{2\phi} \hat{g}_{ab}(t))$$

$$g'_{ab}(\sigma) = \mathrm{e}^{2\omega(\sigma)} g_{ab}(\sigma), \phi'(\sigma) = \phi(\sigma) - \omega(\sigma), \Psi'(\sigma) = \Psi(\sigma)$$

$$g'_{ab} = F^* g_{ab}, g'_{ab} = f'^*(\mathrm{e}^{2\phi'} \hat{g}'_{ab}), g_{ab} = f^*(\mathrm{e}^{2\phi} \hat{g}_{ab})$$

$$\hat{F} = f'^{-1} \circ F \circ f, \phi' = \hat{F}^*(\phi - \omega), \hat{g}'_{ab} = \hat{F}^*(\mathrm{e}^{2\omega} \hat{g}_{ab})$$

$$g'_{ab} = \tilde{f}^*(\mathrm{e}^{2\phi} \hat{g}_{ab})$$

$$\begin{aligned} g'_{ab} &= F^* g_{ab} = F^*\left(f^*(\mathrm{e}^{2\phi} \hat{g}_{ab})\right) = F^*\left(f^*(\mathrm{e}^{2(\phi-\omega)} \mathrm{e}^{2\omega} \hat{g}_{ab})\right) \\ &= f'^*(\mathrm{e}^{2\phi'} \hat{g}'_{ab}) \end{aligned}$$

Campos fantasma. Métrica de Grassmann y Métrica de Faddeev-Popov.

$$\Delta_{\text{FP}}[g] = \int \mathrm{d}'_g b \mathrm{d}'_g c \mathrm{e}^{-S_{\text{gh}}[g,b,c]}$$

$$\begin{aligned} S_{\text{gh}}[g, b, c] &:= \frac{1}{4\pi} \int \mathrm{d}^2\sigma \sqrt{g} g^{ab} g^{cd} b_{ac} (P_1 c)_{bd} \\ &= \frac{1}{4\pi} \int \mathrm{d}^2\sigma \sqrt{g} g^{ab} (b_{ac} \nabla_b c^c + b_{bc} \nabla_a c^c - b_{ab} \nabla_c c^c) \end{aligned}$$

$$(P_1 c)_{ab} = \nabla_a c_b + \nabla_b c_a - g_{ab} \nabla_c c^c = 0, (P_1^\dagger b)_a = -2 \nabla^b b_{ab} = 0$$

$$T_{ab}^{\text{gh}} = -b_{ac} \nabla_b c^c - b_{bc} \nabla_a c^c + c^c \nabla_c b_{ab} + g_{ab} b_{cd} \nabla^c c^d$$



$$g^{ab}T_{ab}^{\text{gh}}=0$$

$$S_{\text{gh}}[\text{e}^{2\omega}g,b,c]=S_{\text{gh}}[g,b,c]$$

$$N_{\text{gh}}(b)=-1,N_{\text{gh}}(c)=1.$$

$$N_{\text{gh}}(\Psi)=0$$

$$Z_g = \int_{\mathcal{M}_g} \text{d}^{\text{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \int \text{d}_{\hat{g}} \Psi \text{d}'_{\hat{g}} b \text{d}'_{\hat{g}} c \text{e}^{-S_m[\hat{g}, \Psi] - S_{\text{gh}}[\hat{g}, b, c]}$$

$$\langle 0|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1|0\rangle=1$$

$$A_{0,2}(k,k')=\frac{\mathcal{C}_{S^2}}{\text{Vol}\mathcal{K}_{0,2}}\langle \mathcal{V}_k(\infty,\infty)c_0\bar{c}_0\mathcal{V}_{k'}(0,0)\rangle_{S^2}$$

Weyl ghost.

$$F_{ab}^\perp = \sqrt{g} g_{ab} - \sqrt{\hat{g}} \hat{g}_{ab} = 0$$

$$S'_{\text{gh}}[g,b,c,c_w] = \frac{1}{4\pi} \int \text{d}^2\sigma \sqrt{g} g^{ab} (b_{ac} \nabla_b c^c + b_{bc} \nabla_a c^c + 2b_{ab} c_w)$$

$$\nabla_a c_b + \nabla_b c_a + 2g_{ab} c_w = 0, \nabla^a b_{ab} = 0, g^{ab} b_{ab} = 0$$

$$c_w = -\frac{1}{2} \nabla_a c^a$$

$$T'^{\text{gh}}_{ab} = -(b_{ac} \nabla_b c^c + b_{bc} \nabla_a c^c + 2b_{ab} c_w) - \nabla_c (b_{ab} c^c) + \frac{1}{2} g_{ab} g^{cd} (b_{ce} \nabla_d c^e + b_{de} \nabla_c c^e + 2b_{cd} c_w)$$

$$g^{ab} T'^{\text{gh}}_{ab} = -g^{ab} \nabla_c (b_{ab} c^c)$$

$$\begin{aligned} g^{ab} (b_{ac} \delta \nabla_b c^c + b_{bc} \delta \nabla_a c^c) &= 2g^{ab} b_{ac} \delta \nabla_b c^c = 2g^{ab} b_{ac} \delta \Gamma_{bd}^c c^d \\ &= g^{ab} b_{ac} g^{ce} (\nabla_b \delta g_{de} + \nabla_d \delta g_{be} - \nabla_e \delta g_{bd}) c^d \\ &= b^{ab} (\nabla_a \delta g_{bc} + \nabla_c \delta g_{ab} - \nabla_b \delta g_{ac}) c^c \\ &= b^{ab} \nabla_c \delta g_{ab} c^c \end{aligned}$$

$$\int \text{d}_g c_w \text{e}^{-(c_w g^{ab} b_{ab})_g} = \delta(g^{ab} b_{ab})$$

Modo Cero de los campos fantasma.

$$Z_g = \int_{\mathcal{M}_g} \text{d}^{\text{M}_g} t \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \int \text{d}_{\hat{g}} \Psi \text{d}'_{\hat{g}} b \text{d}'_{\hat{g}} c \prod_{i=1}^{\text{M}_g} (b, \hat{\mu}_i)_{\hat{g}} \text{e}^{-S_m[\hat{g}, \Psi] - S_{\text{gh}}[\hat{g}, b, c]}$$

$$(b, \hat{\mu}_i)_{\hat{g}} = \int \text{d}^2\sigma \sqrt{\hat{g}} G_{\perp}^{abcd} b_{ab} \hat{\mu}_{i,cd} = \int \text{d}^2\sigma \sqrt{\hat{g}} g^{ac} g^{bd} b_{ab} \hat{\mu}_{i,cd}.$$



$$b = b_0 + b', b_0 = b_{0i}\phi_i$$

$$1=\int \mathrm{d}_{\hat{g}}be^{-|b|^2_{\hat{g}}}=J\int \mathrm{d}_{\hat{g}}b'\prod_i \mathrm{d}b_{0i}\mathrm{e}^{-|b'|^2_{\hat{g}}-|b_{0i}\phi_i|^2}=J\sqrt{\det(\phi_i,\phi_j)}.$$

$$\int \mathrm{d}^{\mathbf{M}_g}b_{0i}\prod_jb_0(\sigma_j^0)=\int \mathrm{d}^{\mathbf{M}_g}b_{0i}\prod_j[b_{0i}\phi_i(\sigma_j^0)]=\det\phi_i(\sigma_j^0).$$

$$\frac{\mathrm{d}_{\hat{g}}b'}{\sqrt{\det(\phi_i,\phi_j)}_{\hat{g}}}=\frac{\mathrm{d}_{\hat{g}}b}{\det\phi_i(\sigma_j^0)}\prod_{j=1}^{\mathbf{M}_g}b(\sigma_j^0).$$

$$\mathrm{d}_{\hat{g}}b'\frac{\det(\phi_i,\hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i,\phi_j)}_{\hat{g}}}=\mathrm{d}_{\hat{g}}b\prod_{j=1}^{\mathbf{M}_g}(b,\hat{\mu}_j)_{\hat{g}}$$

$$\prod_{j=1}^{\mathbf{M}_g}b(\sigma_j^0)=\prod_{j=1}^{\mathbf{M}_g}[b_{0i}\phi_i(\sigma_j^0)]=\det\phi_i(\sigma_j^0)\prod_{j=1}^{\mathbf{M}_g}b_{0i}\\ \det(\phi_i,\hat{\mu}_j)_{\hat{g}}\prod_{j=1}^{\mathbf{M}_g}b_{0i}=\prod_{j=1}^{\mathbf{M}_g}[b_{0i}(\phi_i,\hat{\mu}_j)_{\hat{g}}]=\prod_{j=1}^{\mathbf{M}_g}(b_{0i}\phi_i,\hat{\mu}_j)_{\hat{g}}=\prod_{j=1}^{\mathbf{M}_g}(b,\hat{\mu}_j)_{\hat{g}}$$

$$Z_g=\int_{\mathcal{M}_g}\mathrm{d}^{\mathbf{M}_g}t\frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\det\psi_i(\sigma_j^0)}\int \mathrm{d}_{\hat{g}}\Psi\,\mathrm{d}_{\hat{g}}b\,\mathrm{d}_{\hat{g}}c\prod_{j=1}^{\mathbf{K}_g^c}\frac{\epsilon_{ab}}{2}c^a(\sigma_j^0)c^b(\sigma_j^0)\\ \times\prod_{i=1}^{\mathbf{M}_g}(\hat{\mu}_i,b)_{\hat{g}}\mathrm{e}^{-S_m[\hat{g},\Psi]-S_{\text{gh}}[\hat{g},b,c]}$$

$$c\rightarrow c+c_0,P_1c_0=0.$$

$$b\rightarrow b+b_0,P_1^\dagger b_0=0.$$

$$\Phi_0=\ln~g_s$$

$$g_s^{2g-2}=\mathrm{e}^{-\Phi_0\chi_g}=\exp\left(-\frac{\Phi_0}{4\pi}\int \mathrm{d}^2\sigma\sqrt{g}R\right)=\mathrm{e}^{-\Phi_0 S_{\text{EH}}[g]}$$

Amplitudes en un espacio de Moduli – Operadores de vértice. Métrica perturbativa de Green.

$$Z_g=\int \frac{\mathrm{d}_g g_{ab}}{\Omega_{\text{gauge}}[g]}Z_m[g],Z_m[g]=\int \mathrm{d}_g\Psi\mathrm{e}^{-S_m[g,\Psi]}$$

$$V_\alpha(k_i)\!:=\!\int \mathrm{d}^2\sigma\sqrt{g(\sigma)}V_\alpha(k;\sigma).$$

$$S_{\text{EH}}[g]\!:=\!\frac{1}{4\pi}\int \mathrm{d}^2\sigma\sqrt{g}R+\frac{1}{2\pi}\oint \mathrm{d}s k=\chi_{g,n}$$



$$\chi_{g,n}\!:=\chi\big(\Sigma_{g,n}\big)=2-2g-n$$

$$g_s^{-\chi_{g,n}}=\mathrm{e}^{-\Phi_0 S_\mathrm{EH}[g]}, \Phi_0\!:=\ln~g_s.$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}}\!:=\int~\frac{\mathrm{d}_g g_{ab}}{\Omega_{\text{gauge}}\left[g\right]}\mathrm{d}_g\Psi\mathrm{e}^{-S_m[g,\Psi]-\Phi_0 S_\mathrm{EH}[g]}\prod_{i=1}^n\left(\int~~\mathrm{d}^2\sigma_i\sqrt{g(\sigma_i)}V_{\alpha_i}(k_i;\sigma_i)\right)$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}}\!:=A_{g,n}(k_1,\ldots,k_n)_{\alpha_1,\ldots,\alpha_n}\!:=A_{g,n}\Big(V_{\alpha_1}(k_1),\ldots,V_{\alpha_n}(k_n)\Big).$$

$$A_n(k_1,\ldots,k_n)_{\alpha_1,\ldots,\alpha_n}=\sum_{g=0}^\infty A_{g,n}(k_1,\ldots,k_n)_{\alpha_1,\ldots,\alpha_n}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = \int~\frac{\mathrm{d}_g g_{ab}}{\Omega_{\text{gauge}}\left[g\right]}\mathrm{e}^{-\Phi_0 S_\mathrm{EH}[g]}\int~\prod_{i=1}^n~~\mathrm{d}^2\sigma_i\sqrt{g}\left\langle\prod_{i=1}^n~V_{\alpha_i}(k_i;\sigma_i)\right\rangle_{m,g}$$

$$H=0.$$

$$S=1+\mathrm{i} T,$$

$$A_n(k_1,\ldots,k_n)=G_n(k_1,\ldots,k_n)\prod_{i=1}^n~\big(k_i^2+m_i^2\big).$$

$$G_n(k_1,\ldots,k_n)\propto\delta^{(D)}(k_1+\cdots+k_n).$$

$$T_2=G_2(k,k')(k^2+m^2)^2\sim (k^2+m^2)\delta^{(D)}(k+k')\mathop{\rightarrow}\limits_{k^2\rightarrow -m^2}0$$

$$G_2(k,k')=\frac{\delta^{(D)}(k+k')}{k^2+m^2}$$

$$A_2(k,k')=2k^0(2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}-\boldsymbol{k'}).$$

$$\left[a(\boldsymbol{k}),a^\dagger(\boldsymbol{k}')\right]=2k^0(2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}-\boldsymbol{k'}).$$

$$\phi(\alpha)=\int~~\mathrm{d}^{D-1}\boldsymbol{k}\alpha(\boldsymbol{k})^*a^\dagger(\boldsymbol{k}).$$

Reparaciones de Gauge. Métrica de Faddeev-Popov.

$$\begin{aligned}\delta_\xi V_{\alpha_i}(k_i)&=\delta_\xi\int~\mathrm{d}^2\sigma\sqrt{g}V_{\alpha_i}(k_i;\sigma)=0\\\delta_\omega V_{\alpha_i}(k_i)&=\delta_\omega\int~\mathrm{d}^2\sigma\sqrt{g}V_{\alpha_i}(k_i;\sigma)=0\end{aligned}$$



$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbb{M}_g} t \frac{\det(\phi_i, \hat{\mu}_j)_{\hat{g}}}{\sqrt{\det(\phi_i, \phi_j)_{\hat{g}}}} \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}}$$

$$\times \int \prod_{i=1}^n d^2\sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle_{m, \hat{g}}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbb{M}_g} t \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\sqrt{\det(\psi_i, \psi_j)_{\hat{g}}}} \int d_{\hat{g}} b \, d'_{\hat{g}} c \prod_{i=1}^{\mathbb{M}_g} (b, \hat{\mu}_i)_{\hat{g}} e^{-S_{\text{gh}}[\hat{g}, b, c]} \\ \times \int \prod_{i=1}^n d^2\sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle_{m, \hat{g}}$$

$$A_{g,n} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbb{M}_g} t \frac{\Omega_{\text{ckv}}[\hat{g}]^{-1}}{\det \psi_i(\sigma_j^0)} \int d_{\hat{g}} b \, d_{\hat{g}} c \prod_{j=1}^{K_g^c} \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0) \prod_{i=1}^{\mathbb{M}_g} (\hat{\mu}_i, b)_{\hat{g}} e^{-S_{\text{gh}}[\hat{g}, b, c]} \\ \times \int \prod_{i=1}^n d^2\sigma_i \sqrt{\hat{g}} \left\langle \prod_{i=1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle$$

$$1 = \Delta(\sigma_j^0) \int d\xi \prod_{j=1}^{K_g^c} \delta^{(2)}\left(\sigma_j - \sigma_j^{0(\xi)}\right), \sigma_j^{0(\xi)} = \sigma_j^0 + \delta_\xi \sigma_j^0, \delta_\xi \sigma_j^0 = \xi(\sigma_j^0)$$

$$\Delta(\sigma_j^0) = \det \psi_i(\sigma_j^0)$$

$$\xi(\sigma_j^0) = \alpha_i \psi_i(\sigma_j^0)$$

$$1 = \int \prod_{j=1}^{K_g^c} d^2 \delta \sigma_j e^{-\Sigma_j(\delta \sigma_j, \delta \sigma_j)} = \Delta \int \prod_{j=1}^{K_g} d \alpha_i e^{-\Sigma_{j,i,i'}(\alpha_i \psi_i(\sigma_j), \alpha_{i'} \psi_{i'}(\sigma_j))} \\ = \Delta \left(\det \psi_i(\sigma_j) \right)^{-1}$$

$$\chi_{g,n} = 2 - 2g - n < 0.$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}} = g_s^{-\chi_{g,n}} \int_{\mathcal{M}_g} d^{\mathbb{M}_g} t \int d_{\hat{g}} b \, d_{\hat{g}} c \prod_{j=1}^{K_g^c} \frac{\epsilon_{ab}}{2} c^a(\sigma_j^0) c^b(\sigma_j^0) \prod_{i=1}^{\mathbb{M}_g} (\hat{\mu}_i, b)_{\hat{g}} e^{-S_{\text{gh}}[\hat{g}, b, c]} \\ \times \int \prod_{i=K_g^c+1}^n d^2\sigma_i \sqrt{\hat{g}} \left\langle \prod_{j=1}^{K_g^c} \hat{V}_{\alpha_j}(k_j; \sigma_j^0) \prod_{i=K_g^c+1}^n \hat{V}_{\alpha_i}(k_i; \sigma_i) \right\rangle$$



$$A_{g,n}(\{k_i\})_{\{\alpha_i\}}=g_s^{-\chi_{g,n}}\int_{\mathcal{M}_g}\mathrm{d}^{\mathsf{M}_g}t\int_{i=\mathsf{K}_g^c+1}^n\mathrm{d}^2\sigma_i\sqrt{\hat{g}}\left.\left\langle \prod_{j=1}^{\mathsf{K}_g^c}\frac{\epsilon_{ab}}{2}c^a(\sigma_j^0)c^b(\sigma_j^0)\prod_{i=1}^{\mathsf{M}_g}(\hat{\mu}_i,b)_{\hat{g}}\right\rangle_{\mathrm{gh},\hat{g}}\right.$$

$$\times \left\langle \prod_{j=1}^{\mathsf{K}_g^c} \hat{V}_{\alpha_j}(k_j;\sigma_j^0) \prod_{i=\mathsf{K}_g^c+1}^n \hat{V}_{\alpha_i}(k_i;\sigma_i) \right\rangle_{m,\hat{g}}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}}=g_s^{-\chi_{g,n}}\int_{\mathcal{M}_g}\mathrm{d}^{\mathsf{M}_g}t\int_{i=\mathsf{K}_g^c+1}^n\mathrm{d}^2\sigma_i\sqrt{\hat{g}}\left.\left\langle \prod_{i=1}^{\mathsf{M}_g}\hat{B}_i\prod_{j=1}^{\mathsf{K}_g^c}\hat{V}_{\alpha_j}(k_j;\sigma_j^0) \prod_{i=\mathsf{K}_g^c+1}^n \hat{V}_{\alpha_i}(k_i;\sigma_i) \right\rangle\right.$$

$$\hat{V}_{\alpha_j}(k_j;\sigma_j^0)\!:=\!\frac{\epsilon_{ab}}{2}c^a(\sigma_j^0)c^b(\sigma_j^0)\hat{V}_{\alpha_j}(k_j;\sigma_j^0),\hat{B}_i\!:=\!(\hat{\mu}_i,b)_{\hat{g}}$$

$$A_{g,n}(\{k_i\})_{\{\alpha_i\}}=g_s^{-\chi_{g,n}}\int_{\mathcal{M}_g\times\mathbb{C}^{n-K_g^c}}\left\langle \prod_{i=1}^{\mathsf{M}_g}\hat{B}_i\,\mathrm{d} t_i\prod_{j=1}^{\mathsf{K}_g^c}\hat{V}_{\alpha_i}(k_i;\sigma_j^0)\prod_{i=K_g^c+1}^n\hat{V}_{\alpha_i}(k_i;\sigma_i)\mathrm{d}^2\sigma_i\sqrt{\hat{g}}\right\rangle.$$

$$\langle V_k(z,\bar z)V_{k'}(z',\bar z')\rangle_{S^2}=\frac{{\rm i}(2\pi)^D\delta^{(D)}(k+k')}{|z-z'|^4}$$

$$A_{0,2}(k,k') = \frac{C_{S^2}}{\text{Vol}\mathcal{K}_{0,0}} \int \; \mathrm{d}^2z \; \mathrm{d}^2z' \langle V_k(z,\bar z)V_{k'}(z',\bar z')\rangle_{S^2}$$

$$A_{0,2}(k,k') = \frac{C_{S^2}}{\text{Vol}\mathcal{K}_{0,2}} \langle V_k(\infty,\infty)V_{k'}(0,0)\rangle_{S^2}$$

$$A_{0,2}(k,k') = (2\pi)^{D-1}\delta^{(D-1)}(\boldsymbol{k}+\boldsymbol{k'})\frac{C_{S_2}2\pi {\rm i}\delta(0)}{\text{Vol}\mathcal{K}_{0,2}}$$

$$A_{0,2}(k,k') = \frac{C_{S^2}}{\text{Vol}\mathcal{K}_{0,2}} \langle V_k(\infty,\infty)V_{k'}(0,0)\rangle_{S^2}$$

$$\text{Vol}\mathcal{K}_{0,2}=\int\frac{\mathrm{d}^2z}{|z|^2}=2\int_0^{2\pi}\mathrm{d}\sigma\int_0^\infty\frac{\mathrm{d} r}{r}$$

$$\text{Vol}\mathcal{K}_{0,2}=4\pi\int_0^\infty\frac{\mathrm{d} r}{r}=4\pi\int_{-\infty}^\infty\mathrm{d}\tau=4\pi\!\lim_{\varepsilon\rightarrow 0}\!\int_{-\infty}^\infty\mathrm{d}\tau\mathrm{e}^{\mathrm{i}\varepsilon\tau}=4\pi\times 2\pi\!\lim_{\varepsilon\rightarrow 0}\!\delta(\varepsilon)$$

$$\text{Vol}_\varepsilon\mathcal{K}_{0,2}=8\pi^2\delta(\varepsilon)$$

$$\text{Vol}_{M,E}\mathcal{K}_{0,2}=8\pi^2\mathrm{i}\delta(E)$$

$$X^0(z,\bar z)=x^0+\frac{\mathrm{i}}{2}\alpha' k^0\mathrm{ln}\;|z|^2=x^0+\mathrm{i}\alpha' k^0\tau,$$

$$X_M^0=x_M^0+\alpha' k_M^0t$$



$$\mathrm{Vol}_M\mathcal{K}_{0,2}\rightarrow \frac{8\pi^2\mathrm{i}\delta(0)}{\alpha'k_M^0}=\frac{C_{S_2}2\pi\mathrm{i}\delta(0)}{2k_M^0}$$

$$A_{0,2}(k,k')=2k^0(2\pi)^{D-1}\delta^{(D-1)}(\pmb{k}+\pmb{k}')$$

$$Z_0 \sim \frac{\delta^{(D)}(0)}{\mathrm{Vol}\mathcal{K}_0}$$

$${\bf Cuantizaci\'on y simetr\'ia BRST - Transformaciones BRST. Par\'ametros de Grassmann.}$$

$$Z_g = \int_{\mathcal{M}_g} \frac{\mathrm{d}^{\mathbf{M}_g}t}{\Omega_{\mathrm{ckv}}[g]} \int \; \mathrm{d}_g g_{ab} \; \mathrm{d}_g\Psi \; \mathrm{d}_g b \; \mathrm{d}'_gc \delta\big(\sqrt{g}g_{ab}-\sqrt{\hat{g}}\hat{g}_{ab}\big) \prod_{i=1}^{\mathbf{M}_g} \; (\phi_i,b)_g \mathrm{e}^{-S_m[g,\Psi]-S_{\mathrm{gh}}[g,b,c]}$$

$$Z_g = \int_{\mathcal{M}_g} \frac{\mathrm{d}^{\mathbf{M}_g}t}{\Omega_{\mathrm{ckv}}[g]} \int \; \mathrm{d}_g g_{ab} \; \mathrm{d}_gB^{ab} \; \mathrm{d}_g\Psi \; \mathrm{d}_gb \; \mathrm{d}'_c c \prod_{i=1}^{\mathbf{M}_g} \; (\phi_i,b)_g \mathrm{e}^{-S_m[g,\Psi]-S_{\mathrm{gf}}[g,\hat{g},B]-S_{\mathrm{gh}}[g,b,c]}$$

$$S_{\mathrm{gf}}[g,\hat{g},B] = -\frac{\mathrm{i}}{4\pi}\int \;\; \mathrm{d}^2\sigma B^{ab}\big(\sqrt{g}g_{ab}-\sqrt{\hat{g}}\hat{g}_{ab}\big)$$

$$\begin{gathered}\delta_\epsilon g_{ab}=\mathrm{i}\epsilon\mathcal{L}_cg_{ab},\delta_\epsilon\Psi=\mathrm{i}\epsilon\mathcal{L}_c\Psi,\\\delta_\epsilon c^a=\mathrm{i}\epsilon\mathcal{L}_cc^a,\delta_\epsilon b_{ab}=\epsilon B_{ab},\delta_\epsilon B_{ab}=0,\end{gathered}$$

$$\delta_\epsilon g_{ab}=\mathrm{i}\epsilon\mathcal{L}_cg_{ab}+\mathrm{i}\epsilon g_{ab}c_w,\delta_\epsilon c_w=\mathrm{i}\epsilon\mathcal{L}_cc_w$$

$$B_{ab}=\mathrm{i}T_{ab}\colon=\mathrm{i}\left(T_{ab}^m+T_{ab}^{\mathrm{gh}}\right),$$

$$\delta_\epsilon\Psi=\mathrm{i}\epsilon\mathcal{L}_c\Psi,\delta_\epsilon c^a=\mathrm{i}\epsilon\mathcal{L}_cc^a,\delta_\epsilon b_{ab}=\mathrm{i}\epsilon T_{ab}$$

$$\delta_\epsilon c^a=\epsilon c^b\partial_b c^a$$

$$Q_B=\int \;\; \mathrm{d}\sigma j_B^0$$

$$Q_B^2=0$$

$$N_{\mathrm{gh}}(Q_B)=1$$

$$\delta_\epsilon\Psi=\mathrm{i}[\epsilon Q_B,\Psi]_\pm$$

$$T_{ab}=[Q_B,b_{ab}]$$

$$|\psi\rangle\in\mathcal{H}(Q_B)\colon=\frac{\ker Q_B}{\mathrm{Im} Q_B},$$

$$Q_B|\psi\rangle=0,\nexists|\chi\rangle\colon|\psi\rangle=Q_B|\chi\rangle.$$

$$|\psi\rangle\sim|\psi\rangle+Q_B|\Lambda\rangle.$$

$$\int \;\; \mathrm{d}\sigma b_{ab}|\psi\rangle=0$$



$$b^+ := \int \mathrm{d}\sigma b_{00}, b^- := \int \mathrm{d}\sigma b_{01}$$

$$\mathcal{H}^-(Q_B) = \mathcal{H}(Q_B) \cap \ker b^-, \mathcal{H}^0(Q_B) = \mathcal{H}^-(Q_B) \cap \ker b^+.$$

$$P_\sigma |\psi\rangle = 0$$

$$P_\sigma |\psi\rangle = \int \mathrm{d}\sigma T_{01} |\psi\rangle = \int \mathrm{d}\sigma \{Q_B, b_{01}\} |\psi\rangle = Q_B \int \mathrm{d}\sigma b_{01} |\psi\rangle,$$

$$b^- |\psi\rangle = 0$$

$$\mathcal{H}^- := \mathcal{H}_\downarrow \oplus \mathcal{H}_\uparrow, \mathcal{H}_\downarrow := \mathcal{H}^0 := \mathcal{H}^- \cap \ker b^+.$$

$$b^+ |\downarrow\rangle = 0, b^+ |\uparrow\rangle = |\downarrow\rangle$$

$$Q_B |\psi_\downarrow\rangle = H |\psi_\uparrow\rangle, Q_B |\psi_\uparrow\rangle = 0$$

$$\begin{aligned} \mathrm{Im} Q_B &= \{|\psi_\uparrow\rangle \in \mathcal{H}_\uparrow | H |\psi_\uparrow\rangle \neq 0\} \\ \ker Q_B &= \{|\psi_\uparrow\rangle \in \mathcal{H}_\uparrow\} \cup \{|\psi_\downarrow\rangle \in \mathcal{H}_\downarrow | H |\psi_\downarrow\rangle = 0\} \end{aligned}$$

$$H |\psi\rangle = 0$$

$$b^+ |\psi\rangle = 0$$

$$Q_B \int \mathrm{d}\sigma b_{00} |\psi\rangle = 0$$

Coordenadas geométricas complejas de un campo de gauge. Métrica de Liouville y simetrías tensoriales.

$$\mathrm{d}s^2 = g_{ab} \mathrm{d}\sigma^a \mathrm{d}\sigma^b = e^{2\phi(\tau,\sigma)} (\mathrm{d}\tau^2 + \mathrm{d}\sigma^2),$$

$$g_{ab} = e^{2\phi} \delta_{ab}, \hat{g}_{ab} = \delta_{ab}.$$

$$\begin{aligned} z &= \tau + i\sigma, & \bar{z} &= \tau - i\sigma \\ \tau &= \frac{z + \bar{z}}{2}, & \sigma &= \frac{z - \bar{z}}{2i} \end{aligned}$$

$$\mathrm{d}s^2 = 2g_{z\bar{z}} \mathrm{d}z \mathrm{d}\bar{z} = e^{2\phi(z,\bar{z})} |\mathrm{d}z|^2.$$

$$\begin{aligned} g_{z\bar{z}} &= \frac{e^{2\phi}}{2}, g_{zz} = g_{\bar{z}\bar{z}} = 0 \\ g^{z\bar{z}} &= 2e^{-2\phi}, g^{zz} = g^{\bar{z}\bar{z}} = 0 \end{aligned}$$

$$\hat{g}_{z\bar{z}} = \frac{1}{2}, \hat{g}^{z\bar{z}} = 2$$

$$w = w(z), \bar{w} = \bar{w}(\bar{z}).$$

$$e^{2\phi(z,\bar{z})} = \left| \frac{\partial w}{\partial z} \right|^2 e^{2\phi(w,\bar{w})}$$



$$\mathrm{d} s^2 = \mathrm{e}^{2\phi(w,\bar w)}|\mathrm{d} w|^2$$

$$\mathrm{d}^2\sigma := \mathrm{d}\tau\;\mathrm{d}\sigma = \frac{1}{2}\;\mathrm{d}^2z,\; \mathrm{d}^2z := \mathrm{d} z\;\mathrm{d}\bar{z}$$

$$\delta^{(2)}(z)=\frac{1}{2}\delta^{(2)}(\sigma)$$

$$\int\;\mathrm{d}^2z\delta^{(2)}(z)=\int\;\mathrm{d}^2\sigma\delta^{(2)}(\sigma)=1$$

$$\begin{gathered}\partial_z=\frac{1}{2}(\partial_\tau-\mathrm{i}\partial_\sigma),\partial_{\bar{z}}=\frac{1}{2}(\partial_\tau+\mathrm{i}\partial_\sigma)\\ \mathrm{d} z=\mathrm{d}\tau+\mathrm{i}\;\mathrm{d}\sigma,\;\mathrm{d}\bar{z}=\mathrm{d}\tau-\mathrm{i}\;\mathrm{d}\sigma\end{gathered}$$

$$\begin{gathered}\epsilon_{01}=\epsilon^{01}=1\\\epsilon_{z\bar{z}}=\frac{\mathrm{i}}{2},\epsilon^{z\bar{z}}=-2\mathrm{i}\end{gathered}$$

$$V^z=V^0+\mathrm{i} V^1,V^{\bar{z}}=V^0-\mathrm{i} V^1$$

$$V=V^0\partial_0+V^1\partial_1=V^z\partial_z+V^{\bar{z}}\partial_{\bar{z}}$$

$$V^w=\frac{\partial w}{\partial z}V^z,V^{\bar{w}}=\frac{\partial \bar{w}}{\partial \bar{z}}V^{\bar{z}}$$

$$\begin{gathered}T\Sigma_g\simeq T\Sigma_g^+\oplus T\Sigma_g^-\\ V^z\partial_z\in T\Sigma_g^+,V^{\bar{z}}\partial_{\bar{z}}\in T\Sigma_g^-\end{gathered}$$

$$\omega_z=\frac{1}{2}(\omega_0-\mathrm{i}\omega_1),\omega_{\bar{z}}=\frac{1}{2}(\omega_0+\mathrm{i}\omega_1)$$

$$\omega=\omega_0\;\mathrm{d}\sigma^0+\omega_1\;\mathrm{d}\sigma^1=\omega_z\;\mathrm{d} z+\omega_{\bar{z}}\;\mathrm{d}\bar{z}$$

$$\begin{gathered}T^*\Sigma_g\simeq\Omega^{1,0}\big(\Sigma_g\big)\oplus\Omega^{0,1}\big(\Sigma_g\big)\\ \omega_z\;\mathrm{d} z\in\Omega^{1,0}\big(\Sigma_g\big),\omega_{\bar{z}}\;\mathrm{d}\bar{z}\in\Omega^{0,1}\big(\Sigma_g\big)\end{gathered}$$

$$V_z=g_{z\bar{z}}V^{\bar{z}},V_{\bar{z}}=g_{z\bar{z}}V^z$$

$$T\overbrace{z^{\cdots z}}^{q++p-} \underbrace{z\cdots z}_{p_++q_-}=(g^{z\bar{z}})^{p_-}(g_{z\bar{z}})^{q_-}T\overbrace{z^{\cdots z}z\bar{z}\cdots \bar{z}}^{q_+q_-} \underbrace{z\cdots z}_{p_+}\underbrace{\bar{z}\cdots \bar{z}}_{p_-}.$$

$$T\overbrace{w^{\cdots w}}^q \underbrace{w\cdots w}_p=\left(\frac{\partial w}{\partial z}\right)^n T\overbrace{z^{\cdots z}}^q \underbrace{z\cdots z}_p,n:=q-p.$$

$$T^{zz}=2(T^{00}+\mathrm{i} T^{01})\in\mathcal{T}^2,T^{\bar{z}\bar{z}}=2(T^{00}-\mathrm{i} T^{01})\in\mathcal{T}^{-2},T^z{}_z=0$$

$$T_{zz}=g_{z\bar{z}}g_{z\bar{z}}T^{\bar{z}\bar{z}}=\frac{1}{2}(T^{00}-\mathrm{i} T^{01})$$

$$T^{zz}=\left(\frac{\partial z}{\partial \tau}\right)^2T^{00}+\left(\frac{\partial z}{\partial \sigma}\right)^2T^{11}+2\frac{\partial z}{\partial \tau}\frac{\partial z}{\partial \sigma}T^{01}=T^{00}-T^{11}+2\mathrm{i} T^{01}$$



$$\int \mathrm{d}^2z (\partial_z v^z + \partial_{\bar{z}} v^{\bar{z}}) = -\mathrm{i}\oint \left(\mathrm{d}z v^{\bar{z}} - \mathrm{d}\bar{z} v^z \right) = -2\mathrm{i}\oint_{\partial R} (v_z \,\mathrm{d}z - v_{\bar{z}} \,\mathrm{d}\bar{z})$$

$$\psi_i(z,\bar z)=\psi^z_i\partial_z+\psi^{\bar z}_i\partial_{\bar z}, \phi_i(z,\bar z)=\phi_{i,zz}(\,\mathrm{d}z)^2+\phi_{i,\bar z}(\,\mathrm{d}\bar z)^2$$

$$(P_1\xi)_{zz}=2\nabla_z\xi_z=\partial_z\xi^{\bar z}, (P_1\xi)_{\bar z\bar z}=2\nabla_{\bar z}\xi_{\bar z}=\partial_{\bar z}\xi^z\\ \left(P_1^\dagger T\right)_z=-2\nabla^zT_{zz}=-4\partial_{\bar z}T_{zz}, \left(P_1^\dagger T\right)_{\bar z}=-2\nabla^{\bar z}T_{\bar z\bar z}=-4\partial_zT_{\bar z\bar z}$$

$$\psi^z=\psi^z(z), \psi^{\bar z}=\psi^{\bar z}(\bar z), \phi_{zz}=\phi_{zz}(z), \phi_{\bar z\bar z}=\phi_{\bar z\bar z}(\bar z)$$

$$\ker P_1=\text{Span}\{\psi_K(z)\}\oplus\text{Span}\{\bar{\psi}_K(\bar{z})\}, K=1,\dots,\mathsf{K}_g^c\\\ker P_1^\dagger=\text{Span}\{\phi_I(z)\}\oplus\text{Span}\{\bar{\phi}_I(\bar{z})\}, I=1,\dots,\mathsf{M}_g^c$$

$$m_I=t_{2I-1}+{\rm i} t_{2I}, \bar m_I=t_{2I-1}-{\rm i} t_{2I}, I=1,\dots,\mathsf{M}_g^c$$

$$\mathrm{d}^{\mathsf{M}_g}t=\mathrm{d}^{2\mathsf{M}_g^c}m$$

$$(T_1,T_2)=2\int \mathrm{d}^2\sigma \sqrt{\hat g}\hat g^{ac}g^{bd}T_{1,ab}T_{2,cd}=4\int \mathrm{d}^2z\big(T_{1,zz}T_{2,\bar z\bar z}+T_{1,\bar z\bar z}T_{2,zz}\big)\\ (\xi_1,\xi_2)=\int \mathrm{d}^2\sigma \sqrt{\hat g}\hat g_{ab}\xi^a\xi^b=\frac{1}{4}\int \mathrm{d}^2z\big(\xi_1^z\xi_2^{\bar z}+\xi_1^{\bar z}\xi_2^z\big)$$

$$\mu_{izz}=\partial_i\bar{g}_{zz}, \mu_{i\bar{z}\bar{z}}=\partial_i\bar{g}_{\bar{z}\bar{z}}$$

$$Z_g=\int_{\mathcal{M}_g}\mathrm{d}^{2\mathsf{M}_g^c}m\frac{\left|\det(\phi_I,\mu_J)\right|^2}{\left|\det(\phi_I,\bar{\phi}_J)\right|}\frac{\det' P_1^\dagger P_1}{\left|\det(\psi_I,\bar{\psi}_J)\right|}\frac{Z_m[\delta]}{\Omega_{\text{ckv}}[\delta]}$$

$$c:=c^z, \bar c:=c^{\bar z}, b:=b_{zz}, \bar b:=b_{\bar z\bar z}$$

$$S_{\text{gh}}[g,b,c]=\frac{1}{2\pi}\int \mathrm{d}^2z\big(b\partial_{\bar z}c+\bar b\partial_z\bar c\big)$$

$$\partial_z c=0, \partial_z b=0, \partial_{\bar z} \bar c=0, \partial_{\bar z} \bar b=0$$

$$\bigwedge_{i=1}^{\mathsf{M}_g}B_i\,\mathrm{d}t_i=\bigwedge_{I=1}^{\mathsf{M}_g^c}B_I\bar{B}_I\,\mathrm{d}m_I\wedge\bar{m}_I, B_I:=(\mu_I,b)$$

$$Z_g=\int_{\mathcal{M}_g}\mathrm{d}^{2\mathsf{M}_g^c}m\frac{\Omega_{\text{ckv}}[\delta]^{-1}}{\left|\det\psi_I(z_j^0)\right|^2}\int \mathrm{d}(b,\bar b)\mathrm{d}(c,\bar c)\prod_{j=1}^{\mathsf{K}_g^c}c(z_j^0)\bar c(\bar z_j^0)\prod_{I=1}^{\mathsf{M}_g^c}|(\mu_I,b)|^2\mathrm{e}^{-S_{\text{gh}}[b,c]}Z_m[\delta].$$

$$\textbf{Simetr\'ia conforme en dimensi\'on }\mathbb{R}^4.\textbf{ Ecuaciones de Killing y m\'etrica Beltrami-Laplace}.$$

$$x^\mu \longrightarrow x'^\mu = x'^\mu(x)$$

$$g_{\mu\nu}(x)\longrightarrow g'_{\mu\nu}(x')=\frac{\partial x^\rho}{\partial x'^\mu}\frac{\partial x^\sigma}{\partial x'^\nu}g_{\rho\sigma}(x)=\Omega(x')^2g_{\mu\nu}(x')$$

$$\frac{u\cdot v}{|u||v|}=\frac{u'\cdot v'}{|u'||v'|}$$



$$\Omega\colon=\mathrm{e}^\omega.$$

$$\delta x^{\mu} = \xi^{\mu}$$

$$\delta g_{\mu\nu} = \mathcal{L}_\xi g_{\mu\nu} = \nabla_\mu \xi_\nu + \nabla_\nu \xi_\mu = \frac{2}{d} g_{\mu\nu} \nabla_\rho \xi^\rho$$

$$\Omega^2=1+\frac{2}{d}\nabla_\rho\xi^\rho$$

$$\mathrm{ISO}(\mathcal{M})\subset \mathrm{CISO}(\mathcal{M})$$

$$\delta g_{\mu\nu} = \mathcal{L}_\xi g_{\mu\nu} = \nabla_\mu \xi_\nu + \nabla_\nu \xi_\mu = 0$$

$$\eta=\mathrm{diag}(\underbrace{-1,\ldots,-1}_q,\underbrace{1,\ldots,1}_p)$$

$$\bigl(\eta_{\mu\nu}\Delta+(D-2)\partial_\mu\partial_\nu\bigr)\partial\cdot\epsilon=0$$

$$\begin{gathered}\xi^\mu=a^\mu\\\xi^\mu={\omega^\mu}_\nu x^\nu\\\xi^\mu=\lambda x^\mu\\\xi^\mu=b^\mu x^2-2b\cdot xx^\mu\end{gathered}$$

$$\mathrm{ISO}(\mathbb{R}^{p,q})=\mathrm{SO}(p,q), \mathrm{CISO}(\mathbb{R}^{p,q})=\mathrm{SO}(p+1,q+1)$$

$$\dim \mathrm{SO}(p+1,q+1)=\frac{1}{2}(p+q+2)(p+q+1)$$

Planos complejos.

Métrica de Riemann.

$$\overline{\mathbb{C}}=\mathbb{C}\cup\{\infty\}.$$

$$\lim_{r\rightarrow\infty}re^{\mathrm{i}\theta}\colon=\infty$$

$$z=\mathrm{e}^{\mathrm{i}\phi}\cot\frac{\theta}{2}$$

$$\begin{gathered}z=x+\mathrm{i}y,\quad \bar{z}=x-\mathrm{i}y\\x=\frac{z+\bar{z}}{2},\quad y=\frac{z-\bar{z}}{2\mathrm{i}}\end{gathered}$$

$$\mathrm{d}s^2=\mathrm{d}x^2+\mathrm{d}y^2=\mathrm{d}z\;\mathrm{d}\bar{z}$$

$$\partial\colon=\partial_z=\frac{1}{2}\big(\partial_x-\mathrm{i}\partial_y\big),\bar{\partial}\colon=\partial_{\bar{z}}=\frac{1}{2}\big(\partial_x+\mathrm{i}\partial_y\big)$$

$$\partial\phi(z_1)\partial\phi(z_2)\colon=\partial_{z_1}\partial_{z_2}\phi(z_1)\phi(z_2).$$



$$w = \frac{1}{z}$$

Métrica de Lorentz.

$$t \in \mathbb{R}, \sigma \in [0, L), \sigma \sim \sigma + L,$$

$$ds^2 = -dt^2 + d\sigma^2 = -d\sigma^+ d\sigma^-,$$

$$d\sigma^\pm = dt \pm d\sigma$$

$$\tau = it,$$

$$ds^2 = d\tau^2 + d\sigma^2.$$

$$w = \tau + i\sigma, \bar{w} = \tau - i\sigma$$

$$ds^2 = dw d\bar{w}.$$

$$w = i(t + \sigma) = i\sigma^+, \bar{w} = i(t - \sigma) = i\sigma^-.$$

$$z = e^{2\pi w/L}, \bar{z} = e^{2\pi \bar{w}/L},$$

$$ds^2 = \left(\frac{L}{2\pi}\right)^2 \frac{dz d\bar{z}}{|z|^2}.$$

Álgebra de Witt – Métrica de Cauchy-Riemann- Killing – Laurent – Möbius. Definiciones cuánticas de un CFT y sus operadores y propagadores. Métrica de Virasoro. Funciones de Correlación.

$$ds^2 = dz d\bar{z},$$

$$z \rightarrow z' = f(z), \bar{z} \rightarrow \bar{z}' = \bar{f}(\bar{z})$$

$$ds^2 = dz' d\bar{z}' = \left| \frac{df}{dz} \right|^2 dz d\bar{z}.$$

$$\delta z = v(z), \delta \bar{z} = \bar{v}(\bar{z}),$$

$$\bar{\partial}v = 0, \partial\bar{v} = 0.$$

$$v(z) = \sum_{n \in \mathbb{Z}} v_n z^{n+1}, \bar{v}(\bar{z}) = \sum_{n \in \mathbb{Z}} \bar{v}_n \bar{z}^{n+1},$$

$$\ell_n = -z^{n+1} \partial_z, \bar{\ell}_n = -\bar{z}^{n+1} \partial_{\bar{z}}, n \in \mathbb{Z}.$$

$$[\ell_m, \ell_n] = (m - n) \ell_{m+n}, [\bar{\ell}_m, \bar{\ell}_n] = (m - n) \bar{\ell}_{m+n}, [\ell_m, \bar{\ell}_n] = 0.$$

$$\delta g_{z\bar{z}} = \partial v + \bar{\partial} \bar{v}, \delta g_{zz} = \delta g_{\bar{z}\bar{z}} = 0$$

$$\lim_{|z| \rightarrow 0} v(z) < \infty \implies \forall n < -1: v_n = 0$$



$$v(1/w) = \frac{\mathrm{d}z}{\mathrm{d}w} \sum_n \ v_n w^{-n-1}$$

$$\lim_{|z|\rightarrow\infty}v(z)=\lim_{|w|\rightarrow 0}\frac{\mathrm{d}z}{\mathrm{d}w}v(1/w)=-\lim_{|w|\rightarrow 0}\frac{v(1/w)}{w^2}<\infty\implies \forall n>1\colon v_n=0$$

$$\{\ell_{-1},\ell_0,\ell_1\}\cup\{\bar{\ell}_{-1},\bar{\ell}_0,\bar{\ell}_1\}$$

$$\ell_{-1}=-\partial_z, \ell_0=-z\partial_z, \ell_1=-z^2\partial_z$$

$$[\ell_0,\ell_{\pm 1}]=\mp \ell_{\pm 1}, [\ell_1,\ell_{-1}]=2\ell_0$$

$$\mathrm{PSL}(2,\mathbb{C})\colon=\mathrm{SL}(2,\mathbb{C})/\mathbb{Z}_2\sim\mathrm{SO}(3,1)$$

$$\mathcal{K}_0=\mathrm{PSL}(2,\mathbb{C})$$

$$g=\begin{pmatrix} a&b\\c&d\end{pmatrix}, a,b,c,d\in\mathbb{C}, \det g=ad-bc=1$$

$$\mathrm{K}_0\colon=\dim \mathrm{SL}(2,\mathbb{C})=6.$$

$$f_g(z)=\frac{az+b}{cz+d}$$

$$v(z)=\beta+2\alpha z+\gamma z^2, \bar{v}(\bar{z})=\bar{\beta}+2\bar{\alpha}\bar{z}+\bar{\gamma}\bar{z}^2$$

$$a=1+\alpha, b=\beta, c=-\gamma, d=1-\alpha$$

$$\text{translation:} \quad f_g(z)=z+a, \qquad a\in\mathbb{C},$$

$$\text{rotation:} \quad f_g(z)=\zeta z, \qquad |\zeta|=1,$$

$$\text{dilatation:} \quad f_g(z)=\lambda z, \qquad \lambda\in\mathbb{R},$$

$$\text{SCT:} \quad f_g(z)=\frac{z}{cz+1}, \qquad c\in\mathbb{C}.$$

$$\text{inversion: } I^+(z)\colon=I(z)\colon=\frac{1}{z}$$

$$I^-(z)\colon=-I(z)=I(-z)=-\frac{1}{z}$$

$$g_{\infty,0,1}(z)=\frac{1}{1-z}$$

$$\forall f \text{ meromorphic}: \mathcal{O}(z,\bar{z})=\left(\frac{\mathrm{d}f}{\mathrm{d}z}\right)^h\left(\frac{\mathrm{d}\bar{f}}{\mathrm{d}\bar{z}}\right)^{\bar{h}}\mathcal{O}'(f(z),\bar{f}(\bar{z})),$$



$$\forall f\in \mathrm{PSL}(2,\mathbb{C})\colon \mathcal{O}(z,\bar{z})=\left(\frac{\mathrm{d}f}{\mathrm{d}z}\right)^h\left(\frac{\mathrm{d}\bar{f}}{\mathrm{d}\bar{z}}\right)^{\bar{h}}\mathcal{O}'(f(z),\bar{f}(\bar{z}))$$

$$\Delta := h + \bar{h}, s := h - \bar{h}.$$

$$\mathcal{O}(z,\bar{z})\mathrm{d}z^h\;\mathrm{d}\bar{z}^{\bar{h}}$$

$$f\circ\mathcal{O}(z,\bar{z})\colon=f'(z)^h\bar{f}'(\bar{z})^{\bar{h}}\mathcal{O}'(f(z),\bar{f}(\bar{z}))$$

$$\delta z = \nu(z), \delta \bar{z} = \bar{\nu}(\bar{z})$$

$$\delta \mathcal{O}(z,\bar{z})=(h\partial\nu+\nu\partial)\mathcal{O}(z,\bar{z})+(\bar{h}\bar{\partial}\bar{\nu}+\bar{\nu}\bar{\partial})\mathcal{O}(z,\bar{z})$$

$$\nabla^\nu T_{\mu\nu}=0, g^{\mu\nu}T_{\mu\nu}=0$$

$$g^{\mu\nu}T_{\mu\nu}=4T_{z\bar{z}}=T_{xx}+T_{yy}=0$$

$$T_{z\bar{z}}=0$$

$$\partial_z T_{\bar{z}\bar{z}}=0, \partial_{\bar{z}} T_{zz}=0$$

$$T(z)\colon=T_{zz}(z),\bar{T}(\bar{z})\colon=T_{\bar{z}\bar{z}}(\bar{z})$$

$$J_\nu(z)\colon=J_\nu^{\bar{z}}(z)=-T(z)\nu(z),\bar{J}_\nu(\bar{z})\colon=J_\nu^z(\bar{z})=-\bar{T}(\bar{z})\bar{\nu}(\bar{z})$$

$$Z=\int\; \mathrm{d}\Psi \mathrm{e}^{-S[\Psi]}$$

$$\begin{gathered}[L_m,L_n]=(m-n)L_{m+n}+\frac{c}{12}m(m-1)(m+1)\delta_{m+n}\\ [\bar{L}_m,\bar{L}_n]=(m-n)\bar{L}_{m+n}+\frac{\bar{c}}{12}m(m-1)(m+1)\delta_{m+n}\\ [L_m,\bar{L}_n]=0,[c,L_m]=0,[\bar{c},\bar{L}_m]=0\end{gathered}$$

$$\left\langle \prod_{i=1}^n\mathcal{O}_i(z_i,\bar{z}_i)\right\rangle =\int\;\mathrm{d}\Psi \mathrm{e}^{-S[\Psi]}\prod_{i=1}^n\mathcal{O}_i(z_i,\bar{z}_i),$$

$$\left\langle \prod_{i=1}^n\mathcal{O}_i(z_i,\bar{z}_i)\right\rangle =\prod_{i=1}^n\left(\frac{\mathrm{d}f}{\mathrm{d}z}(z_i)\right)^{h_i}\left(\frac{\mathrm{d}\bar{f}}{\mathrm{d}\bar{z}}(\bar{z}_i)\right)^{\bar{h}_i}\times\left\langle \prod_{i=1}^n\mathcal{O}_i\left(f(z_i),\bar{f}(\bar{z}_i)\right)\right\rangle .$$

$$\delta\left\langle \prod_{i=1}^n\mathcal{O}_i(z_i,\bar{z}_i)\right\rangle =\sum_{i=1}^n\left(h_i\partial_i\nu(z_i)+\nu(z_i)\partial_i+\text{ c.c. }\right)\left\langle \prod_{i=1}^n\mathcal{O}_i(z_i,\bar{z}_i)\right\rangle =0,$$



$$\begin{aligned}\langle \mathcal{O}_i(z_i, \bar{z}_i) \rangle &= \delta_{h_i,0} \delta_{\bar{h}_i,0}, \\ \langle \mathcal{O}_i(z_i, \bar{z}_i) \mathcal{O}_j(z_j, \bar{z}_j) \rangle &= \delta_{h_i,h_j} \delta_{\bar{h}_i,\bar{h}_j} \frac{g_{ij}}{z_{ij}^{2h_i} \bar{z}_{ij}^{2\bar{h}_i}}, \\ \langle \mathcal{O}_i(z_i, \bar{z}_i) \mathcal{O}_j(z_j, \bar{z}_j) \mathcal{O}_k(z_k, \bar{z}_k) \rangle &= \frac{C_{ijk}}{z_{ij}^{h_i+h_j-h_k} z_{jk}^{h_j+h_k-h_i} z_{ki}^{h_i+h_k-h_j}} \\ &\times \frac{1}{\bar{z}_{ij}^{\bar{h}_i+\bar{h}_j-\bar{h}_k} \bar{z}_{jk}^{\bar{h}_j+\bar{h}_k-\bar{h}_i} \bar{z}_{ki}^{\bar{h}_i+\bar{h}_k-\bar{h}_j}},\end{aligned}$$

$$z_{ij}=z_i-z_j.$$

$$\left\langle \prod_{i=1}^4 \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle = f(x, \bar{x}) \prod_{i < j} \frac{1}{z_{ij}^{(h_i+h_j)-h/3}} \times \text{c.c.}$$

$$h:=\sum_{i=1}^4 h_i, \bar{h}:=\sum_{i=1}^4 \bar{h}_i.$$

$$\chi\!:=\!\frac{z_{12}z_{34}}{z_{13}z_{24}}.$$

Operadores formales y cuantización radial. Ordenadores radiales y conmutadores. Productos de expansión.

$$z = e^{\tau + i\sigma} = x + iy.$$

$$\tau \rightarrow \tau + T$$

$$z \rightarrow e^T z.$$

$$H=\frac{2\pi}{L}(L_0+\bar{L}_0)$$

$$\text{on-shell state: } h + \bar{h} = 0.$$

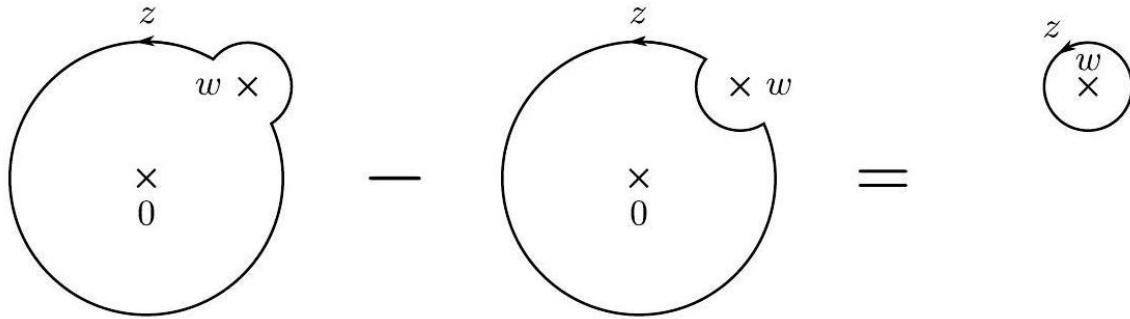
$$R(A(z)B(w)) = \begin{cases} A(z)B(w) & |z| > |w| \\ (-1)^F B(w)A(z) & |w| > |z| \end{cases}$$

$$[A(z), B(w)]_{\pm, |z|=|w|} = \lim_{\delta \rightarrow 0} (A(z)B(w)|_{|z|=|w|+\delta} \pm B(w)A(z)|_{|z|=|w|-\delta})$$

$$A = \oint_{C_0} \frac{dz}{2\pi i} a(z), B = \oint_{C_0} \frac{dz}{2\pi i} b(z)$$

$$\begin{aligned}[A, B]_{\pm} &= \oint_{C_0} \frac{dw}{2\pi i} \oint_{C_w} \frac{dz}{2\pi i} a(z)b(w) \\ [A, b(w)]_{\pm} &= \oint_{C_w} \frac{dz}{2\pi i} a(z)b(w)\end{aligned}$$





$$\partial_\mu j^\mu = \partial j^z + \bar{\partial} j^{\bar{z}} = 2(\partial j_{\bar{z}} + \bar{\partial} j_z) = 0$$

$$Q = \frac{1}{2\pi i} \oint_{C_0} (j_z dz - j_{\bar{z}} d\bar{z})$$

$$j(z) := j_z(z), \bar{j}(\bar{z}) := j_{\bar{z}}(\bar{z})$$

$$Q = Q_L + Q_R, Q_L := \frac{1}{2\pi i} \oint_{C_0} j(z) dz, Q_R := -\frac{1}{2\pi i} \oint_{C_0} \bar{j}(\bar{z}) d\bar{z}$$

$$\delta_\epsilon \mathcal{O}(z, \bar{z}) = -[\epsilon Q, \mathcal{O}(z, \bar{z})] = -\epsilon \oint_{C_z} \frac{dw}{2\pi i} j(w) \mathcal{O}(z, \bar{z}) + \epsilon \oint_{C_z} \frac{d\bar{w}}{2\pi i} \bar{j}(\bar{w}) \mathcal{O}(z, \bar{z})$$

$$Q = \frac{1}{2\pi} \int d\sigma j^0$$

$$Q = \frac{1}{2\pi} \int (d\sigma j^0 - d\tau j^1) = -\frac{1}{2\pi} \int \epsilon_{\mu\nu} j^\mu dx^\nu$$

$$Q = -\frac{1}{2\pi} \oint \epsilon_{z\bar{z}} (j^z d\bar{z} - j^{\bar{z}} dz) = -\frac{i}{4\pi} \oint (j^z d\bar{z} - j^{\bar{z}} dz) = -\frac{1}{2\pi i} \oint (j_z dz - j_{\bar{z}} d\bar{z})$$

$$\mathcal{O}_i(z_i, \bar{z}_i) \mathcal{O}_j(z_j, \bar{z}_j) = \sum_k \frac{c_{ij}^k}{z_{ij}^{h_i+h_j-h_k} \bar{z}_{ij}^{\bar{h}_i+\bar{h}_j-\bar{h}_k}} \mathcal{O}_k(z_j, \bar{z}_j),$$

$$C_{ijk} = g_{k\ell} c_{ij}^\ell.$$

Identidad OPE.

$$\phi(z)1 = \sum_{n \in \mathbb{N}} \frac{(z-w)^n}{n!} \partial^n \phi(w)$$

$$A(z)B(w) := \sum_{n=-\infty}^N \frac{\{AB\}_n(z)}{(z-w)^n}$$

$$A(z)B(w) \sim \sum_{n=1}^N \frac{\{AB\}_n(z)}{(z-w)^n} =: A(z)B(w)$$

$$\phi_i(z_i)\phi_j(z_j) \sim \sum_k \theta(h_i + h_j - h_k) \frac{c_{ij}^k}{(z-w)^{h_i+h_j-h_k}} \phi_k(w)$$

$$T(z)\phi(w) \sim \frac{h\phi(w)}{(z-w)^2} + \frac{\partial\phi(w)}{z-w}$$

$$\begin{aligned}\delta\phi(z) &= \oint_{C_z} \frac{dw}{2\pi i} v(w) T(w) \phi(z) \sim \oint_{C_z} \frac{dw}{2\pi i} v(w) \left(\frac{h\phi(z)}{(w-z)^2} + \frac{\partial\phi(z)}{w-z} \right) \\ &= h\partial v(z)\phi(z) + v(z)\partial\phi(z)\end{aligned}$$

$$T(z)T(w) \sim \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}$$

$$\delta T = 2\partial v T + v\partial T + \frac{c}{12}\partial^3 v$$

$$T'(w) = \left(\frac{dz}{dw}\right)^{-2} \left(T(z) - \frac{c}{12}S(w,z)\right) = \left(\frac{dz}{dw}\right)^{-2} T(z) + \frac{c}{12}S(z,w)$$

$$S(w,z) = \frac{w^{(3)}}{w'} - \frac{3}{2} \left(\frac{w''}{w'} \right)^2$$

$$S(u,z) = S(w,z) + \left(\frac{dw}{dz}\right)^2 S(u,w)$$

$$\begin{aligned}\delta T(z) &= \oint_{C_z} \frac{dw}{2\pi i} v(w) T(w) T(z) \sim \oint_{C_z} \frac{dw}{2\pi i} v(w) \left(\frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} \right) \\ &= \frac{c}{2 \times 3!} \partial^3 v(z) + 2\partial v(z)T(z) + v(z)\partial T(z)\end{aligned}$$

Conjugaciones BPZ y Modo de Expansión.

$$t^\dagger = -i\tau^\dagger = t \implies \tau^\dagger = -\tau.$$

$$z \xrightarrow{\tau \rightarrow -\tau} e^{-\tau+i\sigma} = \frac{1}{z^*} = I(\bar{z}),$$

$$\mathcal{O}(z, \bar{z})^\ddag := (\bar{I} \circ \mathcal{O}(z, \bar{z}))^\dagger,$$

$$\mathcal{O}(z, \bar{z})^\ddag = \left[\frac{1}{\bar{z}^{2h} z^{2\bar{h}}} \mathcal{O}\left(\frac{1}{\bar{z}}, \frac{1}{z}\right) \right]^\dagger = \frac{1}{z^{2h} \bar{z}^{2\bar{h}}} \mathcal{O}^\dagger\left(\frac{1}{z}, \frac{1}{\bar{z}}\right)$$

$$(\lambda \mathcal{O}_1 \cdots \mathcal{O}_n)^\ddag = \lambda^* \mathcal{O}_n^\ddag \cdots \mathcal{O}_1^\ddag, \lambda \in \mathbb{C},$$

$$\mathcal{O}(z, \bar{z})^t := I^\pm \circ \mathcal{O}(z, \bar{z}) = \frac{(\mp 1)^{h+\bar{h}}}{z^{2h} \bar{z}^{2\bar{h}}} \mathcal{O}\left(\pm \frac{1}{z}, \pm \frac{1}{\bar{z}}\right)$$

$$(\lambda \mathcal{O}_1 \cdots \mathcal{O}_n)^t = \lambda \mathcal{O}_1^t \cdots \mathcal{O}_n^t, \lambda \in \mathbb{C}$$

$$1^\ddag = 1^t = 1$$



$$\mathcal{O}(z,\bar{z})=\sum_{m,n}\frac{\mathcal{O}_{m,n}}{z^{m+h}\bar{z}^{n+\bar{h}}}$$

$$m+h\in \mathbb{Z}+\nu, n+\bar{h}\in \mathbb{Z}+\bar{\nu}, \nu, \bar{\nu}=\begin{cases} 0 & \text{periodic} \\ 1/2 & \text{anti-periodic} \end{cases}$$

$$\mathcal{O}\left(\mathrm{e}^{2\pi \mathrm{i}} z, \bar{z}\right) = \mathrm{e}^{2\pi \mathrm{i} \nu} \mathcal{O}(z, \bar{z}), \mathcal{O}\left(z, \mathrm{e}^{2\pi \mathrm{i}} \bar{z}\right) = \mathrm{e}^{2\pi \mathrm{i} \bar{\nu}} \mathcal{O}(z, \bar{z})$$

$$\nu, \bar{\nu} = \begin{cases} 0 & \text{NS} \\ 1/2 & \text{R} \end{cases}$$

$$\nu, \bar{\nu} = \begin{cases} 0 & \text{untwisted} \\ 1/2 & \text{twisted} \end{cases}$$

$$\left(\mathcal{O}^\ddag\right)_{-m,-n}=\left(\mathcal{O}_{m,n}\right)^\dagger.$$

$$\mathcal{O}^\ddag=\mathcal{O}\implies\left(\mathcal{O}_{m,n}\right)^\dagger=\mathcal{O}_{-m,-n}.$$

$$\phi(z)=\sum_{n\in\mathbb{Z}+h+\nu}\frac{\phi_n}{z^{n+h}}.$$

$$\phi_n=\oint_{C_0}\frac{\mathrm{d}z}{2\pi\mathrm{i}}z^{n+h-1}\phi(z)$$

$$\phi^\ddag=\phi\implies (\phi_n)^\dagger=\phi_{-n}.$$

$$\phi_n^t=(I^\pm\circ\phi)_n=(-1)^h(\pm 1)^n\phi_{-n}.$$

$$\begin{aligned}\phi_n^t&=(I^\pm\circ\phi)_n=\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}z^{n+h-1}I^\pm\circ\phi(z)\\&=\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}z^{n+h-1}\left(\mp\frac{1}{z^2}\right)^h\phi\left(\pm\frac{1}{z}\right)\\&=(\mp 1)^h\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}z^{n-h-1}\phi\left(\pm\frac{1}{z}\right)\\&=(\mp 1)^h\oint\frac{\mathrm{d}w}{2\pi\mathrm{i}}\left(\pm\frac{1}{w}\right)^{n-h}w^{-1}\phi(w)\\&=(\mp 1)^h(\pm 1)^{n-h}\oint\frac{\mathrm{d}w}{2\pi\mathrm{i}}w^{-n+h-1}\phi(w)\end{aligned}$$

$$\frac{\mathrm{d}z}{z}=\mp\frac{\mathrm{d}w}{w^2z}=-\frac{\mathrm{d}w}{w}$$

$$T(z)=\sum_{n\in\mathbb{Z}}\frac{L_n}{z^{n+2}}, L_n=\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}T(z)z^{n+1}$$

$$L_n^\dagger=L_{-n}.$$

$$[L_m,\phi_n]=(m(h-1)-n)\phi_{m+n}$$

$$[L_m,\phi(z)]=z^m(z\partial+(n+1)h)\phi(z).$$



$$[L_0, L_{-n}] = n L_{-n}, [L_0, \phi_{-n}] = n \phi_{-n}$$

$$L_n^\pm = L_n \pm \bar{L}_n$$

$$j(z) = \sum_n \frac{j_n}{z^{n+1}}$$

$$Q_L=j_0$$

Espacios de Hilbert y correspondencia operador – estado.

$$L_0|0\rangle = 0, L_{\pm 1}|0\rangle = 0$$

$$\langle \mathcal{O} \rangle := \langle 0 | \mathcal{O} | 0 \rangle.$$

$$|\mathcal{O}\rangle = \lim_{z,\bar{z} \rightarrow 0} \mathcal{O}(z,\bar{z})|0\rangle = \mathcal{O}(0,0)|0\rangle$$

$$|\phi\rangle = \lim_{z \rightarrow 0} \phi(z)|0\rangle = \phi(0)|0\rangle.$$

$$\forall n \geq -h+1: \phi_n|0\rangle = 0$$

$$|\phi\rangle = \phi_{-h}|0\rangle = \oint \frac{dz}{2\pi i} \frac{\phi(z)}{z} |0\rangle.$$

$$\phi(z)|0\rangle = e^{zL_{-1}}\phi(0)e^{-zL_{-1}}|0\rangle = e^{zL_{-1}}|\phi\rangle.$$

$$\forall n \geq -1: L_n|0\rangle = 0$$

$$\forall |\phi\rangle \in \mathcal{H}: \langle \Omega | L_0 | \Omega \rangle \leq \langle \phi | L_0 | \phi \rangle.$$

$$L_0|\Omega\rangle := a_\Omega|\Omega\rangle,$$

$$\langle 0 | = | 0 \rangle^\dagger = | 0 \rangle^t$$

$$\langle 0 | L_0 = 0, \langle 0 | L_{\pm 1} = 0.$$

$$\begin{aligned} \langle \mathcal{O}^\ddagger | &= \lim_{w,\bar{w} \rightarrow 0} \langle 0 | \mathcal{O}(w,\bar{w})^\ddagger = \lim_{w,\bar{w} \rightarrow 0} \frac{1}{w^{2h}\bar{w}^{2\bar{h}}} \langle 0 | \mathcal{O}\left(\frac{1}{\bar{w}}, \frac{1}{w}\right)^\dagger \\ &= \lim_{z,\bar{z} \rightarrow \infty} z^{2h}\bar{z}^{2\bar{h}} \langle 0 | \mathcal{O}^\dagger(z,\bar{z}) \\ &= \langle 0 | I \circ \mathcal{O}^\dagger(0,0) \end{aligned}$$

$$\begin{aligned} \langle \phi^\ddagger | &= \lim_{\bar{w} \rightarrow 0} \langle 0 | \phi(w)^\ddagger = \lim_{\bar{w} \rightarrow 0} \frac{1}{w^{2h}} \langle 0 | \phi^\dagger\left(\frac{1}{w}\right) \\ &= \lim_{z \rightarrow \infty} z^{2h} \langle 0 | \phi^\dagger(z) \\ &= \langle 0 | I \circ \phi^\dagger(0). \end{aligned}$$

$$\langle \phi^\ddagger | = \langle 0 | (\phi^\dagger)_h.$$



$$\begin{aligned}\langle \phi | &:= \lim_{w \rightarrow 0} \langle 0 | \phi(w)^t \\ &= (\pm 1)^h \lim_{z \rightarrow \infty} z^{2h} \langle 0 | \phi(z) \\ &= \langle 0 | I^\pm \circ \phi(0)\end{aligned}$$

$$\langle \phi | = (\pm 1)^h \langle 0 | \phi_h.$$

$$\langle \phi^\pm | = (\pm 1)^h \langle \phi |.$$

$$\forall n \leq h-1: \langle 0 | \phi_n = 0,$$

$$\forall n \leq 1: \langle 0 | L_n = 0.$$

$$\langle 0 | T(z) | 0 \rangle = 0.$$

$$\begin{aligned}\langle \phi_i^\pm | \phi_j \rangle &= \langle 0 | \bar{I} \circ \phi_j(0) \phi_i(0) | 0 \rangle = \lim_{\substack{z \rightarrow \infty \\ w \rightarrow 0}} z^{2h_i} \langle 0 | \phi_i^\dagger(z) \phi_j(w) | 0 \rangle, \\ \langle \phi_i | \phi_j \rangle &= \langle 0 | I \circ \phi_j(0) \phi_i(0) | 0 \rangle = (\pm 1)^{h_i} \lim_{\substack{z \rightarrow \infty \\ w \rightarrow 0}} z^{2h_i} \langle 0 | \phi_i(z) \phi_j(w) | 0 \rangle.\end{aligned}$$

$$\langle \phi_i | \phi_j \rangle = \langle I \circ \phi_i(0) \phi_j(0) \rangle, \langle \phi_i^\pm | \phi_j \rangle = \langle I \circ \phi_i^\dagger(0) \phi_j(0) \rangle.$$

$$\langle \phi_i | \phi_j(z) | \phi_k \rangle = (\pm 1)^{h_i} \lim_{w \rightarrow \infty} w^{2h_i} \langle \phi_i(w) \phi_j(z) \phi_k(0) \rangle.$$

$$\langle \phi_i^c | \phi_j \rangle = \delta_{ij}$$

Módulos Verma.

$$L_0 |\phi\rangle = h |\phi\rangle, \forall n \geq 1: L_n |\phi\rangle = 0.$$

$$|\phi_{\{n_i\}}\rangle := \prod_i L_{-n_i} |\phi\rangle,$$

$$L_0 = h + \sum_i n_i.$$

$$\langle 0 | : \mathcal{O} : | 0 \rangle = 0.$$

$$\langle \Omega | {}^* \mathcal{O}_*^* | \Omega \rangle = 0.$$

$$:A(z)B(w): \stackrel{?}{=} A(z)B(w) - \langle A(z)B(w) \rangle.$$

$$:A(z)B(z): \stackrel{?}{=} \lim_{w \rightarrow z} (A(z)B(w) - \langle A(z)B(w) \rangle).$$

$$:A(z)B(w): := A(z)B(w) - A(z)B(w) = \sum_{n \in \mathbb{N}} (z-w)^n \{AB\}_{-n}(z)$$

$$:AB(z): := :A(z)B(z): := \lim_{w \rightarrow z} :A(z)B(w): = \{AB\}_0(z)$$

$$:AB(z): = \oint_{C_z} \frac{dw}{2\pi i} \frac{A(z)B(w)}{z-w}$$



$$AB(z) := AB(z).$$

$$\begin{aligned} :AB(z): &= \sum_m \frac{:AB:_m}{z^{m+h_A+h_B}} \\ :AB:_m &= \sum_{n \leq -h_A} A_n B_{m-n} + \sum_{n > -h_A} B_{m-n} A_n \end{aligned}$$

$$:AB(z): \neq :BA(z):, :A(BC)(z): \neq :(AB)C(z):.$$

$$\begin{aligned} A_1(z)A_2A_3(w) &=: A_1(z)A_2A_3(w) : + \sqrt{A_1(z)A_2}A_3(w) : \\ A_1(z)A_2A_3(w) &=: A_1(z)A_2(w)A_3(w) : + A_1(z)A_3(w)A_2(w) :. \end{aligned}$$

$$\begin{aligned} A(z)B(w)^n &:= nA(z)B(w)B(w)^{n-1}: \\ A(z)e^{B(w)} &:= A(z)B(w)e^{B(w)}: \\ e^{A(z)}e^{B(w)} &:= \exp(A(z)B(w))e^{A(z)}e^{B(w)}:. \end{aligned}$$

$$\begin{aligned} \prod_{i=1}^n :e^{A_i}: &= : \exp\left(\sum_{i=1}^n A_i\right) : \exp \sum_{i < j} \langle A_i A_j \rangle, \\ \left\langle \prod_{i=1}^n :e^{A_i}: \right\rangle &= \exp \sum_{i < j} \langle A_i A_j \rangle. \end{aligned}$$

$$A(z)e^{B(w)} := A(z) \sum_n \frac{1}{n!} :B(w)^n: = A(z)B(w) \sum_n \frac{1}{(n-1)!} :B(w)^{n-1}:.$$

$$\begin{aligned} :e^{A(z)}e^{B(w)}: &= \sum_{m,n} \frac{1}{m!n!} :A(z)^m: :B(w)^n: \\ &= \sum_{m,n,k} \frac{k!}{m!n!} \binom{m}{k} \binom{n}{k} (A(z)B(w))^k :A(z)^{m-k}: :B(w)^{n-k}: \\ &= \sum_{m,n,k} \frac{1}{k!(m-k)!(n-k)!} (A(z)B(w))^k :A(z)^{m-k}: :B(w)^{n-k}:. \end{aligned}$$

$$\begin{aligned} {}^\star AB(z) &= {}^\star \sum_m \frac{{}^\star AB(z)}{z^{m+h_A+h_B}} {}^\star_m, \\ {}^\star AB {}^\star_m &= \sum_{n \leq 0} A_n B_{m-n} + \sum_{n > 0} B_{m-n} A_n. \end{aligned}$$

$$:AB:_m = {}^\star AB {}^\star_m + \sum_{n=0}^{h_A-1} [B_{m+n}, A_{-n}].$$



$$\begin{aligned}
:AB:_m &= \sum_{n \leq -h_A} A_n B_{m-n} + \sum_{n > -h_A} B_{m-n} A_n \\
&= \sum_{n \geq h_A} A_{-n} B_{m+n} + \sum_{n > 0} B_{m-n} A_n + \sum_{n=0}^{h_A-1} B_{m+n} A_{-n} \\
&= \sum_{n \geq 0} A_{-n} B_{m+n} + \sum_{n > 0} B_{m-n} A_n + \sum_{n=0}^{h_A-1} [B_{m+n}, A_{-n}] \\
&= {}^*AB^*_m + \sum_{n=0}^{h_A-1} [B_{m+n}, A_{-n}].
\end{aligned}$$

$$\phi(z) = \left(\frac{L}{2\pi}\right)^h z^{-h} \phi_{\text{cyl}}(w)$$

$$\phi_{\text{cyl}} = \left(\frac{2\pi}{L}\right)^h \sum_{n \in \mathbb{Z}} \phi_n e^{-\frac{2\pi}{L}w} = \left(\frac{2\pi}{L}\right)^h \sum_{n \in \mathbb{Z}} \frac{\phi_n}{z^n}.$$

$$T_{\text{cyl}}(w) = \left(\frac{2\pi}{L}\right)^2 \left(T(z)z^2 - \frac{c}{24}\right)$$

$$(L_0)_{\text{cyl}} = L_0 - \frac{c}{24}$$

$$H = (L_0)_{\text{cyl}} + (\bar{L}_0)_{\text{cyl}} = L_0 + \bar{L}_0 - \frac{c + \bar{c}}{24}$$

Sistemas CFT – Escalar libre – Acción covariante.

$$S = \frac{\epsilon}{4\pi\ell^2} \int d^2x \sqrt{g} g^{\mu\nu} \partial_\mu X \partial_\nu X$$

$$\epsilon: = \begin{cases} +1 & \text{spacelike} \\ -1 & \text{timelike} \end{cases}, \sqrt{\epsilon}: = \begin{cases} +1 & \text{spacelike} \\ i & \text{timelike} \end{cases}$$

$$X(\tau, \sigma) \sim X(\tau, \sigma + 2\pi).$$

$$T_{\mu\nu} = -\frac{\epsilon}{\ell^2} \left[\partial_\mu X \partial_\nu X - \frac{1}{2} g_{\mu\nu} (\partial X)^2 \right],$$

$$T_\mu^\mu = 0.$$

$$\Delta X = 0,$$

$$0 = \int dX \frac{\delta}{\delta X(\sigma)} \left(e^{-S[X]} X(\sigma') \right)$$

$$\langle \partial^2 X(\sigma) X(\sigma') \rangle = -2\pi\epsilon\ell^2 \delta^{(2)}(\sigma - \sigma').$$

$$\langle X(\sigma) X(\sigma') \rangle = -\frac{\epsilon\ell^2}{2} \ln |\sigma - \sigma'|^2.$$



$$\langle X(\sigma)X(\sigma')\rangle=G(r), r=|\sigma-\sigma'|.$$

$$\Delta G(r)=\frac{1}{r}\partial_r(rG'(r)).$$

$$-2\pi\epsilon\ell^2=2\pi\int_0^r{\rm d}r'r'\times\frac{1}{r'}\partial_{r'}\big(r'G'(r')\big)=2\pi rG'(r).$$

$$G'(r)=-\epsilon\ell^2\ln r$$

$$\ln r=\frac{1}{2}\ln r^2=\frac{1}{2}\ln|\sigma-\sigma'|^2.$$

$$X \longrightarrow X+a, a \in \mathbb{R}.$$

$$J^\mu := 2\pi {\rm i} \epsilon \frac{\partial \mathcal{L}}{\partial (\partial_\mu X)} = \frac{{\rm i}}{\ell^2} g^{\mu\nu} \partial_\nu X, \nabla_\mu J^\mu = 0.$$

$$p=\frac{1}{2\pi}\int\,\,{\rm d}\sigma J^0=\frac{{\rm i}}{2\pi\ell^2}\int\,\,{\rm d}\sigma\partial^0X.$$

$$\tilde{J}^\mu := -{\rm i} \epsilon^{\mu\nu} J_\nu = \frac{1}{\ell^2} \epsilon^{\mu\nu} \partial_\nu X$$

$$\nabla_\mu \tilde{J}^\mu \propto \epsilon^{\mu\nu} \big[\nabla_\mu, \nabla_\nu\big] X = 0$$

$$w=\frac{1}{2\pi}\int\,\,{\rm d}\sigma \tilde{J}^0=\frac{1}{2\pi\ell^2}\int_0^{2\pi}\,\,{\rm d}\sigma \partial_1 X=\frac{1}{2\pi\ell^2}(X(\tau,2\pi)-X(\tau,0))$$

$$J_a^\mu=\frac{{\rm i}}{2\pi\ell^2}\eta_{ab}\partial^\mu X^b$$

Acción en un plano complejo.

$$S=\frac{\epsilon}{2\pi\ell^2}\int\,\,{\rm d}z\,{\rm d}\bar{z}\partial_z X\partial_{\bar{z}} X$$

$$\partial_z\partial_{\bar{z}} X=0.$$

$$X(z,\bar{z})=X_L(z)+X_R(\bar{z})$$

$$X(z)\colon=X_L(z), X(\bar{z})\colon=X_R(\bar{z})$$

$$J\colon=J_z=\frac{{\rm i}}{\ell^2}\partial_z X, \bar{J}\colon=J_{\bar{z}}=\frac{{\rm i}}{\ell^2}\partial_{\bar{z}} X$$

$$\bar{\partial}J=0, \partial\bar{J}=0$$

$$p=p_L+p_R, p_L=\frac{1}{2\pi{\rm i}}\oint\,{\rm d}z J, p_R=-\frac{1}{2\pi{\rm i}}\oint\,{\rm d}\bar{z}\bar{J}$$

$$\tilde{J}_z=\frac{{\rm i}}{\ell^2}\partial_z X=J, \tilde{J}_{\bar{z}}=-\frac{{\rm i}}{\ell^2}\partial_{\bar{z}} X=-\bar{J}$$



$$w=p_L-p_R$$

$$\begin{aligned} p_L &= \frac{p+w}{2}, p_R = \frac{p-w}{2} \\ p^2 + w^2 &= p_L^2 + p_R^2, 2pw = p_L^2 - p_R^2 \end{aligned}$$

$$T := T_{zz} = -\frac{\epsilon}{\ell^2} \partial_z X \partial_z X, \bar{T} := T_{\bar{z}\bar{z}} = -\frac{\epsilon}{\ell^2} \partial_{\bar{z}} X \partial_{\bar{z}} X, T_{z\bar{z}} = 0$$

$$V_k(z,\bar{z}) := e^{i\epsilon k X(z,\bar{z})}.$$

$$V_{k_L,k_R}(z,\bar{z}) := e^{2i\epsilon(k_L X(z) + k_R X(\bar{z}))}.$$

OPE y Modos de Expansión.

$$X(z)X(w) \sim -\frac{\epsilon \ell^2}{2} \ln(z-w).$$

$$\partial X(z)X(w) \sim -\frac{\epsilon \ell^2}{2} \frac{1}{z-w},$$

$$\partial X(z)\partial X(w) \sim -\frac{\epsilon \ell^2}{2} \frac{1}{(z-w)^2}$$

$$T(z)\partial X(w) \sim \frac{\partial X(w)}{(z-w)^2} + \frac{\partial(\partial X(w))}{z-w}.$$

$$T(z)T(w) \sim \frac{1}{2} \frac{1}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}$$

$$c=1.$$

$$h=n$$

$$T(z)\partial^n X(w) \sim \dots + \frac{n\partial^n X(w)}{(z-w)^2} + \frac{\partial(\partial^n X(w))}{z-w}$$

$$T(z)\partial^2 X(w) \sim \frac{2\partial X(w)}{(z-w)^3} + \frac{2\partial^2 X}{(z-w)^2} + \frac{\partial(\partial^2 X(w))}{z-w}.$$

$$J(z)V_k(w,\bar{w}) \sim \frac{\ell^2 k}{2} \frac{V_k(w,\bar{w})}{z-w}.$$

$$T(z)V_k(w,\bar{w}) \sim \frac{h_k V_k(w,\bar{w})}{(z-w)^2} + \frac{\partial V_k(w,\bar{w})}{z-w}$$

$$(h_k, \bar{h}_k) = \left(\frac{\epsilon \ell^2 k^2}{4}, \frac{\epsilon \ell^2 k^2}{4} \right), \Delta_k = \frac{\epsilon \ell^2 k^2}{2}, s_k = 0$$

$$V_k(z,\bar{z})V_{k'}(w,,\bar{w}) \sim \frac{V_{k+k'}(w,\bar{w})}{(z-w)^{-\epsilon kk' \ell^2/2}}$$



$$T(z)\partial X(w) = -\frac{\epsilon}{\ell^2} : \partial X(z)\partial X(z) : \partial X(w) \sim -\frac{2\epsilon}{\ell^2} : \partial X(z)\partial X(z) : \partial X(w) \sim \frac{\partial X(z)}{(z-w)^2}$$

$$\begin{aligned} T(z)\partial X(w) &= \frac{1}{\ell^4} : \partial X(z)\partial X(z) : : \partial X(w)\partial X(w) : \\ &\sim \frac{1}{\ell^4} [: \partial X(z)\partial X(z) : : \partial X(w)\partial X(w) : + : \partial X(z)\partial X(z) : : \partial X(w)\partial X(w) : \\ &\quad + : \partial X(z)\partial X(z) : : \partial X(w)\partial X(w) : + \text{perms}] \\ &\sim 2 \times \frac{1}{4} \frac{1}{(z-w)^4} - 4 \times \frac{1}{2\ell^2} \frac{1}{(z-w)^2} : \partial X(z)\partial X(w) : \\ &\sim \frac{1}{2} \frac{1}{(z-w)^4} - \frac{2}{\ell^2} \frac{1}{(z-w)^2} (: \partial X(w)\partial X(w) : + (z-w) : \partial^2 X(w)\partial X(w) :). \end{aligned}$$

$$\begin{aligned} T(z)\partial^n X(w) &\sim \partial_w^{n-1} \frac{\partial X(z)}{(z-w)^2} \\ &\sim n! \frac{\partial X(z)}{(z-w)^{n+1}} \\ &\sim \frac{n!}{(z-w)^{n+1}} \left(\dots + \frac{1}{(n-1)!} (z-w)^{n-1} \partial^{n-1}(\partial X(w)) \right. \\ &\quad \left. + \frac{1}{n!} (z-w)^n \partial^n(\partial X(w)) \right) \end{aligned}$$

$$\partial X(z)V_k(w, \bar{w}) \sim i\epsilon k \partial X(z)X(w)V_k(w, \bar{w}) \sim i\epsilon k \left(-\frac{\epsilon \ell^2}{2} \frac{1}{z-w} \right) V_k(w, \bar{w})$$

$$\begin{aligned} T(z)V_k(w, \bar{w}) &\sim -\frac{\epsilon}{\ell^2} : \partial X(z)\partial X(z) : : e^{i\epsilon k X(w, \bar{w})} : \\ &\sim \frac{i\epsilon k}{2} \frac{1}{z-w} \partial X(z) : e^{i\epsilon k X(w, \bar{w})} : - \frac{\epsilon}{\ell^2} \partial X(z) : \partial X(z) e^{i\epsilon k X(w, \bar{w})} : \\ &\sim \frac{i\epsilon k}{2} \frac{1}{z-w} (: \partial X(z) e^{i\epsilon k X(w, \bar{w})} : + \partial X(z) : e^{i\epsilon k X(w, \bar{w})} :) \\ &\quad + \frac{i\epsilon k}{2} \frac{\partial X(z) e^{i\epsilon k X(w, \bar{w})} :}{z-w} \\ &\sim \frac{\epsilon k^2 \ell^2}{4} \frac{V_k(w, \bar{w})}{(z-w)^2} + i\epsilon k \frac{: \partial X(w) e^{i\epsilon k X(w, \bar{w})} :}{z-w}. \end{aligned}$$

$$\begin{aligned} V_k(z, \bar{z})V_{k'}(w, \bar{w}) &\sim \exp(-kk'X(z, \bar{z})X(w, \bar{w})) : e^{i\epsilon k X(z, \bar{z})} e^{i\epsilon k' X(w, \bar{w})} : \\ &\sim (z-w)^{\epsilon kk' \ell^2 / 2} V_{k+k'}(w, \bar{w}) \end{aligned}$$

$$\partial X = -i\sqrt{\frac{\ell^2}{2}} \sum_{n \in \mathbb{Z}} \alpha_n z^{-n-1}, \bar{\partial} X = -i\sqrt{\frac{\ell^2}{2}} \sum_{n \in \mathbb{Z}} \bar{\alpha}_n \bar{z}^{-n-1}$$

$$\alpha_n = i\oint \frac{dz}{2\pi i} z^{n-1} \partial X(z), \bar{\alpha}_n = i\oint \frac{dz}{2\pi i} z^{n-1} \bar{\partial} X(z).$$

$$X(z) = \frac{x_L}{2} - i\sqrt{\frac{\ell^2}{2}} \alpha_0 \ln z + i\sqrt{\frac{\ell^2}{2}} \sum_{n \neq 0} \frac{\alpha_n}{n} z^{-n}$$



$$X(\bar{z})=\frac{x_R}{2}-\mathrm{i}\sqrt{\frac{\ell^2}{2}}\bar{\alpha}_0\ln|\bar{z}|+\mathrm{i}\sqrt{\frac{\ell^2}{2}}\sum_{n\neq0}\frac{\bar{\alpha}_n}{n}\bar{z}^{-n}$$

$$p_L=\frac{\alpha_0}{\sqrt{2\ell^2}}, p_R=\frac{\bar{\alpha}_0}{\sqrt{2\ell^2}}$$

$$X(z)=\frac{x_L}{2}-\mathrm{i}\ell^2p_L\ln|z|+\mathrm{i}\sqrt{\frac{\ell^2}{2}}\sum_{n\neq0}\frac{\alpha_n}{n}z^{-n},$$

$$p=\frac{1}{\sqrt{2\ell^2}}(\alpha_0+\bar{\alpha}_0), w=\frac{1}{\sqrt{2\ell^2}}(\alpha_0-\bar{\alpha}_0)$$

$$\alpha_0=\sqrt{\frac{\ell^2}{2}}(p+w), \bar{\alpha}_0=\sqrt{\frac{\ell^2}{2}}(p-w)$$

$$x_L=x+q, x_R=x-q$$

$$x=\frac{1}{2}(x_L+x_R), q=\frac{1}{2}(x_L-x_R)$$

$$X(z,\bar{z})=x-\mathrm{i}\frac{\ell^2}{2}\Big(p\ln|z|^2+w\ln\frac{z}{\bar{z}}\Big)+\mathrm{i}\sqrt{\frac{\ell^2}{2}}\sum_{n\neq0}\frac{1}{n}(\alpha_nz^{-n}+\bar{\alpha}_n\bar{z}^{-n})$$

$$X(\tau,\sigma)=x-\mathrm{i}\ell^2p\tau+\ell^2w\sigma+\cdots$$

$$\begin{aligned} p_L &= \frac{1}{2\pi\mathrm{i}}\oint\mathrm{d} z J = \frac{\mathrm{i}}{\ell^2}\frac{1}{2\pi\mathrm{i}}\oint\mathrm{d} z\partial X = \frac{\mathrm{i}}{\ell^2}\frac{1}{2\pi\mathrm{i}}\oint\mathrm{d} z\partial X \\ &= \frac{1}{\sqrt{2\ell^2}}\frac{1}{2\pi\mathrm{i}}\oint\mathrm{d} z\sum_n\alpha_nz^{-n-1} = \frac{1}{\sqrt{2\ell^2}}\alpha_0. \end{aligned}$$

$$X(\tau,\sigma+2\pi)\sim X(\tau,\sigma)$$

$$X\bigl(\mathrm{e}^{2\pi\mathrm{i}}z,\mathrm{e}^{-2\pi\mathrm{i}}\bar{z}\bigr)\sim X(z,\bar{z}).$$

$$X\bigl(\mathrm{e}^{2\pi\mathrm{i}}z,\mathrm{e}^{-2\pi\mathrm{i}}\bar{z}\bigr)=X(z,\bar{z})-\mathrm{i}\sqrt{\frac{\ell^2}{2}}(\alpha_0-\bar{\alpha}_0)$$

$$\alpha_0=\bar{\alpha}_0\implies p_L=p_R=\frac{p}{2}, w=0.$$

$$N_n=\frac{\epsilon}{n}\alpha_{-n}\alpha_n, \bar{N}_n=\frac{\epsilon}{n}\bar{\alpha}_{-n}\bar{\alpha}_n.$$

$$N=\sum_{n>0}nN_n.$$



$$L_m=\frac{\epsilon}{2}\sum_n:\alpha_n\alpha_{m-n}$$

$$m\neq 0: L_m=\frac{\epsilon}{2}\sum_{n\neq 0,m}:\alpha_n\alpha_{m-n}:+\epsilon\alpha_0\alpha_m,$$

$$L_0=\frac{\epsilon}{2}\sum_n:\alpha_n\alpha_{-n}:=N+\frac{\epsilon}{2}\alpha_0^2=N+\epsilon\ell^2p_L^2$$

$$\hat{L}_0 := N$$

$$\bar{L}_0 = \bar{N} + \epsilon \ell^2 p_R^2, \hat{L}_0 := \bar{N}$$

$$\begin{aligned}L_0^+ &= N + \bar{N} + \epsilon \ell^2(p_L^2 + p_R^2) = N + \bar{N} + \frac{\epsilon \ell^2}{2}(p^2 + w^2) \\L_0^- &= N - \bar{N} + \epsilon \ell^2(p_L^2 - p_R^2) = N - \bar{N} + \epsilon \ell^2 w p\end{aligned}$$

Conmutadores.

$$[\alpha_m,\alpha_n]=\epsilon m\delta_{m+n,0}, [\bar{\alpha}_m,\bar{\alpha}_n]=\epsilon m\delta_{m+n,0}, [\alpha_m,\bar{\alpha}_n]=0$$

$$[p,w]=[p,p]=[w,w]=0, [p,\alpha_n]=[p,\bar{\alpha}_n]=[w,\alpha_n]=[w,\bar{\alpha}_n]=0$$

$$[x_L,p_L]={\rm i}\epsilon, [x_R,p_R]={\rm i}\epsilon$$

$$[x,p]=[q,w]={\rm i}\epsilon, [x,w]=[q,p]=0$$

$$[L_m,\alpha_n]=-n\alpha_{m+n}$$

$$[L_0,\alpha_{-n}]=n\alpha_{-n}$$

$$[N_m,\alpha_{-n}]=\alpha_{-m}\delta_{m,n}$$

$$|k\rangle := \lim_{z,\bar{z}\rightarrow 0} V_k(z,\bar{z}) |0\rangle = {\rm e}^{{\rm i}\epsilon kx}|0\rangle,$$

$$p|k\rangle=k|k\rangle.$$

$$\forall n>0: \alpha_n|k\rangle=0,$$

$$N_n|k\rangle=0.$$

$$L_0^+|k\rangle=2\epsilon\ell^2k^2|k\rangle, L_0^-|k\rangle=0,$$

$$\begin{gathered}\mathcal{F}(k)=\text{Span}\{|k;\{N_n\}\rangle\},\\ |k;\{N_n\}\rangle:=\prod_{n\geq 1}\frac{(\alpha_{-n})^{N_n}}{\sqrt{n^{N_n}N_n!}}|k\rangle, N_n\in\mathbb{N}^*\end{gathered}$$

$$\mathcal{H}=\int_{\mathbb{R}}\mathrm{d}k\mathcal{F}(k)$$



$$\begin{aligned}\lim_{z,\bar{z}\rightarrow 0}e^{i\epsilon kX(z,\bar{z})}|0\rangle &= \lim_{z,\bar{z}\rightarrow 0}\exp i\epsilon k\left[x-i\frac{\ell^2}{2}p\ln|z|^2+i\sqrt{\frac{\ell^2}{2}}\sum_{n\neq 0}\frac{1}{n}(\alpha_n z^{-n}+\bar{\alpha}_n\bar{z}^{-n})\right]|0\rangle \\ &= \lim_{z,\bar{z}\rightarrow 0}\exp\left[i\epsilon kx-\epsilon k\sqrt{\frac{\ell^2}{2}}\sum_{n\neq 0}\frac{1}{n}(\alpha_n z^{-n}+\bar{\alpha}_n\bar{z}^{-n})\right]|0\rangle.\end{aligned}$$

$$\begin{aligned}p|k\rangle &= \frac{1}{\ell^2}\frac{1}{2\pi i}\oint (\mathrm{d}z i\partial X(z) + \mathrm{d}\bar{z} i\bar{\partial} X(\bar{z}))V_k(0,0)|0\rangle \\ &= \frac{1}{\ell^2}\frac{1}{2\pi i}\oint \left(\frac{\mathrm{d}z}{z}\frac{\ell^2 k}{2} + \frac{\mathrm{d}\bar{z}}{\bar{z}}\frac{\ell^2 k}{2}\right)V_k(0,0)|0\rangle \\ &= kV_k(0,0)|0\rangle\end{aligned}$$

$$x^\dagger = x \; p^\dagger = p, \alpha_n^\dagger = \alpha_{-n}.$$

$$L_n^\dagger=L_{-n}$$

$$\langle k|=|k\rangle^\ddag=\langle 0|\mathrm{e}^{-\mathrm{i}\epsilon kx},\langle k|p=\langle k|k.$$

$$\alpha_n^t=-(\pm 1)^n\alpha_{-n}$$

$$p^t=-p,\langle -k|=|k\rangle^t.$$

$$\langle k \mid k' \rangle = 2\pi \delta(k-k')$$

$$\langle k^c|=\frac{1}{2\pi}\langle k|.$$

$$|k\rangle^\ddag=-|k\rangle^t$$

Sistema Ghost.

Acción covariante y holomórfica o antiholomórfica.

$$S=\frac{1}{4\pi}\int \mathrm{d}^2x\sqrt{g}g^{\mu\nu}b_{\mu\mu_1\cdots\mu_{\lambda-1}}\nabla_\nu c^{\mu_1\cdots\mu_{\lambda-1}}$$

$$b_{\mu_1\cdots\mu_n}\longrightarrow \mathrm{e}^{-\mathrm{i}\theta}b_{\mu_1\cdots\mu_n}, c^{\mu_1\cdots\mu_{n-1}}\longrightarrow \mathrm{e}^{\mathrm{i}\theta}c^{\mu_1\cdots\mu_{n-1}}.$$

$$b(z)\colon=b_{z\cdots z}(z), \bar{b}(\bar{z})\colon=b_{\bar{z}\cdots\bar{z}}(\bar{z}), c(z)\colon=c^{z\cdots z}(\bar{z}), \bar{c}(\bar{z})\colon=c^{\bar{z}\cdots\bar{z}}(z)$$

$$S=\frac{1}{2\pi}\int \mathrm{d}^2z(b\bar{\partial}c+\bar{b}\partial\bar{c})$$

$$\partial\bar{b}=0, \bar{\partial}b=0, \partial\bar{c}=0, \bar{\partial}c=0$$

$$h(b)=\lambda, h(c)=1-\lambda, h(\bar{b})=\lambda, h(\bar{c})=1-\lambda,$$



$$\begin{aligned}T &= -\lambda :b\partial c: + (1-\lambda):\partial bc: \\&= -\lambda :\partial(bc): + :\partial bc:\\&= (1-\lambda):\partial(bc): - :b\partial c:.\end{aligned}$$

$$b(z)c(w)=-\epsilon c(w)b(z), b(z)b(w)=-\epsilon b(w)b(z), c(z)c(w)=-\epsilon c(w)c(z)$$

$$\epsilon = \begin{cases} +1 & \text{anticommuting} \\ -1 & \text{commuting} \end{cases}$$

$$\delta b=-{\rm i} b,\delta c={\rm i} c,\delta \bar{b}=-{\rm i} \bar{b},\delta \bar{c}={\rm i} \bar{c}$$

$$j(z)=-:b(z)c(z):, \bar{j}(\bar{z})=-:\bar{b}(\bar{z})\bar{c}(\bar{z}):$$

$$N_{\mathrm{gh}}=N_{\mathrm{gh},L}+N_{\mathrm{gh},R}, N_{\mathrm{gh},L}=\oint\frac{{\rm d} z}{2\pi {\rm i}} j(z), N_{\mathrm{gh},R}=-\oint\frac{{\rm d}\bar{z}}{2\pi {\rm i}} \bar{j}(\bar{z})$$

$$N_{\mathrm{gh}}(c)=1,N_{\mathrm{gh}}(b)=-1,N_{\mathrm{gh}}(\bar{c})=1,N_{\mathrm{gh}}(\bar{b})=-1$$

$$\int \,{\rm d}'b\,{\rm d}'c\frac{\delta}{\delta b(z)}\big[b(w){\rm e}^{-S[b,c]}\big]=0$$

$$\delta^{(2)}(z-w)+\frac{1}{2\pi}\langle b(w)\bar{\partial}c(z)\rangle=0$$

$$\langle c(z)b(w)\rangle=\frac{1}{z-w}$$

$$\begin{aligned}T(z)b(w) &\sim \lambda \frac{b(w)}{(z-w)^2}+\frac{\partial b(w)}{z-w}\\T(z)c(w) &\sim (1-\lambda)\frac{c(w)}{(z-w)^2}+\frac{\partial c(w)}{z-w}\end{aligned}$$

$$T(z)T(w)\sim \frac{c_\lambda/2}{(z-w)^4}+\frac{2T(w)}{(z-w)^2}+\frac{\partial T(w)}{z-w}$$

$$c_\lambda=2\epsilon(-1+6\lambda-6\lambda^2)=-2\epsilon(1+6\lambda(\lambda-1))$$

$$q_\lambda=\epsilon(1-2\lambda)$$

$$c_\lambda=\epsilon\bigl(1-3q_\lambda^2\bigr)$$

$$\begin{aligned}j(z)b(w) &\sim -\frac{b(w)}{z-w}\\j(z)c(w) &\sim \frac{c(w)}{z-w}\end{aligned}$$

$$j(z)\mathcal{O}(w)\sim N_{\mathrm{gh}}(\mathcal{O})\frac{\mathcal{O}(w)}{z-w}$$

$$j(z)j(w)\sim \frac{\epsilon}{(z-w)^2}$$



$$T(z)j(w) \sim \frac{q_\lambda}{(z-w)^3} + \frac{j(w)}{(z-w)^2} + \frac{\partial j(w)}{z-w}.$$

$$j(z) = \frac{dw}{dz} j'(w) + \frac{q_\lambda}{2} \frac{d}{dz} \ln \frac{dw}{dz} = \frac{dw}{dz} j'(w) + \frac{q_\lambda}{2} \frac{\partial_z^2 w}{\partial_z w}.$$

$$j(z) = \frac{dw}{dz} \left(j^{\text{cyl}}(w) - \frac{q_\lambda}{2} \right),$$

$$N_{\text{gh}} = N_{\text{gh}}^{\text{cyl}} - q_\lambda, N_{\text{gh},L} = N_{\text{gh},L}^{\text{cyl}} - \frac{q_\lambda}{2}, N_{\text{gh},R} = N_{\text{gh},R}^{\text{cyl}} - \frac{q_\lambda}{2}.$$

$$N^c - N^b = -\frac{\epsilon q_\lambda}{2} \chi_g = (1-2\lambda)(g-1)$$

$$\begin{aligned} T(z)b(w) &= (-\lambda : b(z)\partial c(z) : + (1-\lambda) : \partial b(z)c(z) :)b(w) \\ &\sim -\lambda : b(z)\partial c(z) : b(w) + (1-\lambda) : \partial b(z)c(z) : b(w) \\ &\sim -\lambda b(z)\partial_z \frac{1}{z-w} + (1-\lambda)\partial b(z) \frac{1}{z-w} \\ &\sim \lambda(b(w) + (z-w)\partial b(w)) \frac{1}{(z-w)^2} + (1-\lambda) \frac{\partial b(w)}{z-w} \end{aligned}$$

$$\begin{aligned} T(z)c(w) &= (-\lambda : b(z)\partial c(z) : + (1-\lambda) : \partial b(z)c(z) :)c(w) \\ &\sim \epsilon \lambda : \partial c(z)b(z) : c(w) - \epsilon (1-\lambda) : c(z)\partial b(z) : c(w) \\ &\sim \lambda \frac{\partial c(z)}{z-w} - (1-\lambda) c(z) \partial_z \frac{1}{z-w} \\ &\sim \lambda \frac{\partial c(w)}{z-w} + (1-\lambda) (c(w) + (z-w)\partial c(w)) \frac{1}{(z-w)^2} \\ &\sim (1-\lambda) \frac{c(w)}{(z-w)^2} + \frac{\partial c(w)}{(z-w)^2} \end{aligned}$$

$$j(z)b(w) = -: b(z)c(z) : b(w) \sim -: b(z)c(z) : b(w) \sim -\frac{b(z)}{z-w} \sim -\frac{b(w)}{z-w}.$$

$$j(z)c(w) = -: b(z)c(z) : c(w) \sim \epsilon : c(z)b(z) : c(w) \sim \frac{c(z)}{z-w} \sim \frac{c(w)}{z-w}.$$

$$\begin{aligned} j(z)j(w) &=: b(z)c(z) : : b(w)c(w) : \\ &\sim : b(z)c(z) : : b(w)c(w) : + : b(z)c(z) : : b(w)c(w) : + : b(z)c(z) : : b(w)c(w) : \\ &\sim \frac{\epsilon}{(z-w)^2} + \frac{\epsilon : c(z)b(w) :}{z-w} + \frac{: b(z)c(w) :}{z-w} \sim \frac{\epsilon}{(z-w)^2}. \end{aligned}$$

$$b(z) = \sum_{n \in \mathbb{Z} + \lambda + \nu} \frac{b_n}{z^{n+\lambda}}, c(z) = \sum_{n \in \mathbb{Z} + \lambda + \nu} \frac{c_n}{z^{n+1-\lambda}},$$

$$b_n = \oint \frac{dz}{2\pi i} z^{n+\lambda-1} b(z), c_n = \oint \frac{dz}{2\pi i} z^{n-\lambda} c(z).$$

$$b \rightarrow -b, c \rightarrow -c.$$

$$N_n^b = : b_{-n} c_n : , N_n^c = \epsilon : c_{-n} b_n :$$



$$N^b = \sum_{n>0} n N_n^b, N^c = \sum_{n>0} n N_n^c$$

$$L_m = \sum_n (n - (1-\lambda)m) :b_{m-n} c_n: = \sum_n (\lambda m - n) :b_n c_{m-n}:.$$

$$L_0 = -\sum_n n :b_n c_{-n}: = \sum_n n :b_{-n} c_n:.$$

$$j_m = -\sum_n :b_{m-n} c_n: = -\sum_n :b_n c_{m-n}:.$$

$$N_{\text{gh},L} = j_0 = -\sum_n :b_{-n} c_n:.$$

$$b_n^\pm = b_n \pm \bar{b}_n, c_n^\pm = \frac{1}{2}(c_n \pm \bar{c}_n).$$

$$b_n^- b_n^+ = 2b_n \bar{b}_n, c_n^- c_n^+ = \frac{1}{2}c_n \bar{c}_n$$

$$\begin{aligned} T &= -\lambda :b \partial c: + (1-\lambda) :\partial b c: \\ &= \sum_{m,n} \left(\lambda :b_m c_n: \frac{n+1-\lambda}{z^{m+\lambda} z^{m+2-\lambda}} - (1-\lambda) :b_m c_n: \frac{m+\lambda}{z^{m+\lambda+1} z^{m+1-\lambda}} \right) \\ &= \sum_{m,n} (\lambda(n+1-\lambda) - (1-\lambda)(m+\lambda)) \frac{:b_m c_n:}{z^{m-n+2}} \\ &= \sum_{m,n} (\lambda(n+1-\lambda) - (1-\lambda)(m-n+\lambda)) \frac{:b_{m-n} c_n:}{z^{m+2}} \\ &= \sum_{m,n} (n-m+\lambda m) \frac{:b_{m-n} c_n:}{z^{m+2}} = \sum_m \frac{L_m}{z^{m+2}}. \end{aligned}$$

$$j = -:bc: = \sum_{m,n} \frac{:b_m c_n:}{z^{m+\lambda} z^{n+1-\lambda}} = \sum_{m,n} \frac{:b_{m-n} c_n:}{z^{m+1}} = \sum_m \frac{j_m}{z^{m+1}}.$$

$$[b_m, c_n]_\epsilon = \delta_{m+n,0}, [b_m, b_n]_\epsilon = 0, [c_m, c_n]_\epsilon = 0$$

$$[b_m^+, c_n^+]_\epsilon = \delta_{m+n}, [b_m^-, c_n^-]_\epsilon = \delta_{m+n}$$

$$[N_m^b, b_{-n}] = b_{-n} \delta_{m,n}, [N_m^c, c_{-n}] = c_{-n} \delta_{m,n}$$

$$[L_m, b_n] = (m(\lambda-1) - n) b_{m+n}, [L_m, c_n] = -(m\lambda + n) c_{m+n}$$

$$[L_0, b_0] = 0, [L_0, c_0] = 0$$

$$[j_m, j_n] = m \delta_{m+n,0}$$

$$[L_m, j_n] = -n j_{m+n} + \frac{q_\lambda}{2} m(m+1) \delta_{m+n,0}$$



$$[N_{gh}, b(w)] = -b(w), [N_{gh}, c(w)] = c(w)$$

$$\begin{aligned}[b_m, c_n]_\epsilon &= \epsilon \oint_{C_0} \frac{dw}{2\pi i} w^{-1} \oint_{C_w} \frac{dz}{2\pi i} z^{-1} w^{n+\lambda} z^{m-\lambda+1} b(z) c(w) \\ &\sim \epsilon \oint_{C_0} \frac{dw}{2\pi i} w^{-1} \oint_{C_w} \frac{dz}{2\pi i} z^{-1} w^{n+\lambda} z^{m-\lambda+1} \frac{\epsilon}{z-w} \\ &= \oint_{C_0} \frac{dw}{2\pi i} w^{m+n-1} = \delta_{m+n,0}\end{aligned}$$

$$[N_{gh}, b(w)] = \oint \frac{dz}{2\pi i} j(z) b(w) \sim -\oint \frac{dz}{2\pi i} \frac{b(w)}{z-w} = -b(w)$$

$$\forall n > -\lambda: b_n |0\rangle = 0, \forall n > \lambda - 1: c_n |0\rangle = 0$$

$$|n_1, \dots, n_{\lambda-1}\rangle = c_1^{n_1} \cdots c_{\lambda-1}^{n_{\lambda-1}} |0\rangle$$

$$L_0 |n_1, \dots, n_{\lambda-1}\rangle = - \left(\sum_{j=1}^{\lambda-1} j n_j \right) |n_1, \dots, n_{\lambda-1}\rangle$$

Estado de energía – momentum. Métrica Grassmann.

$$\{| \downarrow \rangle, | \uparrow \rangle\},$$

$$|\downarrow\rangle := c_1 \cdots c_{\lambda-1} |0\rangle, |\uparrow\rangle := c_0 c_1 \cdots c_{\lambda-1} |0\rangle.$$

$$|\Omega\rangle = \omega_\downarrow |\downarrow\rangle + \omega_\uparrow |\uparrow\rangle, \omega_\downarrow, \omega_\uparrow \in \mathbb{C}.$$

$$b_0 |\uparrow\rangle = |\downarrow\rangle, c_0 |\downarrow\rangle = |\uparrow\rangle, b_0 |\downarrow\rangle = 0, c_0 |\uparrow\rangle = 0.$$

$$\forall n > 0: b_n |\downarrow\rangle = b_n |\uparrow\rangle = 0, c_n |\downarrow\rangle = b_n |\downarrow\rangle = 0.$$

$$|0\rangle = b_{1-\lambda} \cdots b_{-1} |\downarrow\rangle = b_{1-\lambda} \cdots b_{-1} b_0 |\uparrow\rangle.$$

$$L_0 |\downarrow\rangle = a_\lambda |\downarrow\rangle, L_0 |\uparrow\rangle = a_\lambda |\uparrow\rangle,$$

$$a_\lambda = - \sum_{n=1}^{\lambda-1} n = - \frac{\lambda(\lambda-1)}{2} = \frac{c_\lambda}{24} + \frac{2}{24}.$$

$$\{| \downarrow\downarrow \rangle, | \uparrow\downarrow \rangle, | \downarrow\uparrow \rangle, | \uparrow\uparrow \rangle\},$$

$$|\downarrow\downarrow\rangle := c_1 \bar{c}_1 \cdots c_{\lambda-1} \bar{c}_{\lambda-1} |0\rangle$$

$$|\uparrow\downarrow\rangle := c_0 |\downarrow\downarrow\rangle, |\downarrow\uparrow\rangle := \bar{c}_0 |\downarrow\downarrow\rangle, |\uparrow\uparrow\rangle := c_0 \bar{c}_0 |\downarrow\downarrow\rangle.$$

$$c_0 |\downarrow\downarrow\rangle = |\uparrow\downarrow\rangle, \bar{c}_0 |\downarrow\downarrow\rangle = |\downarrow\uparrow\rangle, c_0 |\downarrow\uparrow\rangle = -\bar{c}_0 |\uparrow\downarrow\rangle = |\uparrow\uparrow\rangle$$

$$b_0 |\uparrow\uparrow\rangle = |\downarrow\uparrow\rangle, \bar{b}_0 |\uparrow\uparrow\rangle = -|\uparrow\downarrow\rangle, b_0 |\uparrow\downarrow\rangle = \bar{b}_0 |\downarrow\uparrow\rangle = |\downarrow\downarrow\rangle,$$



$$b_0|\downarrow\downarrow\rangle=\bar{b}_0|\downarrow\downarrow\rangle=0,\quad c_0|\uparrow\downarrow\rangle=\bar{b}_0|\uparrow\downarrow\rangle=0,\\ b_0|\downarrow\uparrow\rangle=\bar{c}_0|\downarrow\uparrow\rangle=0,\quad c_0|\uparrow\uparrow\rangle=\bar{c}_0|\uparrow\uparrow\rangle=0.$$

$$\{| \downarrow\downarrow\rangle, |+\rangle, |- \rangle, | \uparrow\uparrow\rangle\},$$

$$c_0^{\pm}|\downarrow\downarrow\rangle=\frac{1}{2}|\pm\rangle,\quad c_0^{\mp}|\pm\rangle=\pm|\uparrow\uparrow\rangle,\\ b_0^{\pm}|\pm\rangle=\pm 2|\downarrow\downarrow\rangle,\quad b_0^{\mp}|\uparrow\uparrow\rangle=\pm|\pm\rangle.$$

$$b_0^{+}|\downarrow\downarrow\rangle=b_0^{-}|\downarrow\downarrow\rangle=0,\quad c_0^{-}|-\rangle=b_0^{+}|-\rangle=0,\\ c_0^{+}|+\rangle=b_0^{-}|+\rangle=0\quad c_0^{+}|\uparrow\uparrow\rangle=c_0^{-}|\uparrow\uparrow\rangle=0,$$

$$c_0^{-}c_0^{+}|\downarrow\downarrow\rangle=\frac{1}{2}|\uparrow\uparrow\rangle,b_0^{+}b_0^{-}|\uparrow\uparrow\rangle=2|\downarrow\downarrow\rangle.$$

$$2c_0^{+}|\pm\rangle=(c_0+\bar{c}_0)|\uparrow\downarrow\rangle\pm(c_0+\bar{c}_0)|\downarrow\uparrow\rangle=\bar{c}_0|\uparrow\downarrow\rangle\pm c_0|\downarrow\uparrow\rangle=(-1\pm 1)|\uparrow\uparrow\rangle\\ b_0^{+}|\pm\rangle=(b_0+\bar{b}_0)|\uparrow\downarrow\rangle\pm(b_0+\bar{b}_0)|\downarrow\uparrow\rangle=b_0|\uparrow\downarrow\rangle\pm\bar{b}_0|\downarrow\uparrow\rangle=(1\pm 1)|\downarrow\downarrow\rangle\\ 2c_0^{\pm}|\downarrow\downarrow\rangle=(c_0\pm\bar{c}_0)|\downarrow\downarrow\rangle=c_0|\downarrow\downarrow\rangle\pm\bar{c}_0|\downarrow\downarrow\rangle=|\uparrow\downarrow\rangle\pm|\downarrow\uparrow\rangle=|\pm\rangle\\ b_0^{\pm}|\uparrow\uparrow\rangle=(b_0\pm\bar{b}_0)|\uparrow\uparrow\rangle=b_0|\uparrow\uparrow\rangle\pm\bar{b}_0|\uparrow\uparrow\rangle=|\downarrow\uparrow\rangle\mp|\uparrow\downarrow\rangle=\mp|\mp\rangle$$

$$L_0=\sum_nn_{\star}^{\star}b_{-n}c_n{}^{\star}+a_{\lambda}=N^b+N^c+a_{\lambda}$$

$$\hat{L}_0=N^b+N^c.$$

$$L_m=\sum_n\left(n-(1-\lambda)m\right)_{\star}^{\star}b_{m-n}c_n^{\star}+a_{\lambda}\delta_{m,0}$$

$$N_{\text{gh},L}=j_0=\sum_n{}_{\star}^{\star}b_{-n}c_n^{\star}-\left(\frac{q_{\lambda}}{2}+\frac{1}{2}\right)\\ =\sum_{n>0}\left(N_n^c-N_n^b\right)+\frac{1}{2}\left(N_0^c-N_0^b\right)-\frac{q_{\lambda}}{2}$$

$$j_m=\sum_n{}^{\star}b_{m-n}c_n{}^{\star}-\left(\frac{q_{\lambda}}{2}+\frac{1}{2}\right)\delta_{m,0}.$$

$$\widehat{N}_{\text{gh},L}:=\sum_{n>0}\left(N_n^c-N_n^b\right)$$

$$j_0|\downarrow\rangle=(\lambda-1)|\downarrow\rangle=\left(-\frac{q_{\lambda}}{2}-\frac{1}{2}\right)|\downarrow\rangle,\\ j_0|\uparrow\rangle=\lambda|\uparrow\rangle=\left(-\frac{q_{\lambda}}{2}+\frac{1}{2}\right)|\uparrow\rangle.$$

$$j_0|0\rangle=0$$

$$j_0^{\text{cyl}}|\downarrow\rangle=-\frac{1}{2}|\downarrow\rangle,j_0^{\text{cyl}}|\uparrow\rangle=\frac{1}{2}|\uparrow\rangle.$$



$$\begin{aligned}
L_0 &= - \sum_n n : b_n c_{-n} : = - \sum_{n \leq -\lambda} n b_n c_{-n} + \epsilon \sum_{n > -\lambda} n c_{-n} b_n \\
&= \sum_{n \geq \lambda} n b_{-n} c_n + \epsilon \sum_{n > -\lambda} n c_{-n} b_n \\
&= \sum_{n \geq \lambda} n b_{-n} c_n + \epsilon \sum_{n > 0} n c_{-n} b_n + \epsilon \sum_{n=-\lambda+1}^0 n c_{-n} b_n \\
&= \sum_{n \geq \lambda} n b_{-n} c_n + \epsilon \sum_{n > 0} n c_{-n} b_n + \epsilon \sum_{n=0}^{\lambda-1} n b_{-n} c_n + a_\lambda \\
&= \sum_{n > 0} n b_{-n} c_n + \epsilon \sum_{n > 0} n c_{-n} b_n + a_\lambda \\
&= \sum_n n * b_{-n} c_n^* + a_\lambda
\end{aligned}$$

$$\sum_{n=-\lambda+1}^0 c_{-n} b_n = - \sum_{n=0}^{\lambda-1} n c_n b_{-n} = - \sum_{n=0}^{\lambda-1} n (-\epsilon b_{-n} c_n + 1) = \epsilon \sum_{n=0}^{\lambda-1} n b_{-n} c_n + a_\lambda.$$

$$\begin{aligned}
j_0 &= - \sum_n : b_{-n} c_n : = - \sum_{n \geq \lambda} b_{-n} c_n + \epsilon \sum_{n > -\lambda} c_{-n} b_n \\
&= - \sum_{n \geq \lambda} b_{-n} c_n + \epsilon \sum_{n > 0} c_{-n} b_n + \epsilon \sum_{n=1}^{\lambda-1} c_n b_{-n} + \epsilon c_0 b_0 \\
&= - \sum_{n \geq \lambda} b_{-n} c_n + \epsilon \sum_{n > 0} c_{-n} b_n - \sum_{n=1}^{\lambda-1} b_{-n} c_n + \epsilon (\lambda - 1) + \epsilon c_0 b_0 \\
&= - \sum_{n > 0} b_{-n} c_n + \epsilon \sum_{n > 0} c_{-n} b_n + \epsilon (\lambda - 1) + \epsilon c_0 b_0.
\end{aligned}$$

$$\epsilon (\lambda - 1) = - \frac{q_\lambda}{2} - \frac{\epsilon}{2}.$$

$$\begin{aligned}
\epsilon c_0 b_0 + \epsilon (\lambda - 1) &= \frac{\epsilon}{2} c_0 b_0 + \frac{1}{2} (-b_0 c_0 + \epsilon) + \epsilon (\lambda - 1) \\
&= \frac{1}{2} (\epsilon c_0 b_0 - b_0 c_0) + \epsilon \left(\lambda - \frac{1}{2} \right).
\end{aligned}$$

Estructura del espacio de Hilbert. Métrica Grassmann odd.

$$\mathcal{H}_{\text{gh}} = \mathcal{H}_{\text{gh},0} \oplus c_0 \mathcal{H}_{\text{gh},0}, \mathcal{H}_{\text{gh},0} := \mathcal{H}_{\text{gh}} \cap \ker b_0$$

$$\begin{aligned}
\mathcal{H}_{\text{gh},0} &= \text{Span}\{\downarrow; \{N_n^b\}; \{N_n^c\}\}, \\
|\downarrow; \{N_n^b\}; \{N_n^c\}\rangle &= \prod_{n \geq 1} (b_{-n})^{N_n^b} (c_{-n})^{N_n^c} |\downarrow\rangle, N_n^b, N_n^c \in \mathbb{N}^*
\end{aligned}$$

$$\psi = \psi_\downarrow + \psi_\uparrow, \psi_\downarrow \in \mathcal{H}_{\text{gh},0}, \psi_\uparrow \in c_0 \mathcal{H}_{\text{gh},0}$$

$$\mathcal{H}_{\text{gh}} = \mathcal{H}_{\text{gh},0} \oplus c_0 \mathcal{H}_{\text{gh},0} \oplus \bar{c}_0 \mathcal{H}_{\text{gh},0} \oplus c_0 \bar{c}_0 \mathcal{H}_{\text{gh},0}$$



$$\mathcal{H}_{\mathrm{gh},0}\colon=\mathcal{H}_{\mathrm{gh}}\cap \ker b_0\cap \ker \bar{b}_0.$$

$$|\!\downarrow\downarrow;\{N_n^b\};\{N_n^c\};\{\bar N_n^b\};\{\bar N_n^c\}\rangle = \prod_{n\geq 1} (b_{-n})^{N_n^b}(\bar b_{-n})^{\bar N_n^b}(c_{-n})^{N_n^c}(\bar c_{-n})^{\bar N_n^c}|\!\downarrow\downarrow\rangle,\\ N_n^b,\bar N_n^b,N_n^c,\bar N_n^c\in\mathbb{N}^*.$$

$$\psi=\psi_{\downarrow\downarrow}+\psi_{\uparrow\downarrow}+\psi_{\downarrow\uparrow}+\psi_{\uparrow\uparrow},$$

$$\mathcal{H}_{\mathrm{gh}}=\mathcal{H}_{\mathrm{gh},0}\oplus c_0^+\mathcal{H}_{\mathrm{gh},0}\oplus c_0^-\mathcal{H}_{\mathrm{gh},0}\oplus c_0^-c_0^+\mathcal{H}_{\mathrm{gh},0}$$

$$\mathcal{H}_{\mathrm{gh},0}\colon=\mathcal{H}_{\mathrm{gh}}\cap \ker b_0^-\cap \ker b_0^+$$

$$\mathcal{H}_{\mathrm{gh},\pm}\colon=\mathcal{H}_{\mathrm{gh}}\cap \ker b_0^\pm=\mathcal{H}_{\mathrm{gh},0}\oplus c_0^\mp\mathcal{H}_{\mathrm{gh},0},$$

$$\mathcal{H}_{\mathrm{gh}}=\mathcal{H}_{\mathrm{gh},\pm}\oplus c_0^\pm\mathcal{H}_{\mathrm{gh},\pm}.$$

$$b_n^\dagger=\epsilon b_{-n}, c_n^\dagger=c_{-n}.$$

$$b_n^t=(-1)^\lambda b_{-n}, c_n^t=(-1)^{1-\lambda} c_{-n},$$

$$|\downarrow\rangle^\ddag=\langle 0|c_{1-\lambda}\cdots c_{-1}, |\uparrow\rangle^\ddag=\langle 0|c_{1-\lambda}\cdots c_{-1}c_0.$$

$$\langle\downarrow|:=|\downarrow\rangle^t=(-1)^{(1-\lambda)^2}\langle 0|c_{-1}\cdots c_{1-\lambda}\\ \langle\uparrow|:=|\uparrow\rangle^t=(-1)^{\lambda(1-\lambda)}\langle 0|c_0c_{-1}\cdots c_{1-\lambda}.$$

$$\langle\downarrow|=(-1)^{a_\lambda+(1-\lambda)(2-\lambda)}|\downarrow\rangle^\ddag, \langle\uparrow|=(-1)^{a_\lambda}|\uparrow\rangle^\ddag,$$

$$\langle\downarrow|=(-1)^{(1-\lambda)^2+\frac{1}{2}(2-\lambda)(1-\lambda)}|\downarrow\rangle^\ddag=(-1)^{-a_\lambda+(1-\lambda)(2-\lambda)}|\downarrow\rangle^\ddag$$

$$\sum_{i=1}^{\lambda-2} i = \frac{1}{2}(2-\lambda)(1-\lambda)=-a_\lambda+1-\lambda.$$

$$\langle\uparrow|=(-1)^{\lambda(1-\lambda)-\frac{1}{2}\lambda(1-\lambda)}|\uparrow\rangle^\ddag=(-1)^{\frac{1}{2}\lambda(1-\lambda)}|\uparrow\rangle^\ddag.$$

$$\langle\uparrow|b_0=\langle\downarrow|,\langle\downarrow|c_0=\langle\uparrow|,\langle\downarrow|b_0=0,\langle\uparrow|c_0=0.$$

$$|\downarrow\rangle=\begin{pmatrix}0\\1\end{pmatrix}, |\uparrow\rangle=\begin{pmatrix}1\\0\end{pmatrix},$$

$$b_0=\begin{pmatrix}0&0\\1&0\end{pmatrix}, c_0=\begin{pmatrix}0&1\\0&0\end{pmatrix}$$

$$\langle\downarrow|\downarrow\rangle=\langle\uparrow|\uparrow\rangle=0$$

$$\langle\uparrow|\downarrow\rangle=\langle\downarrow|c_0|\downarrow\rangle=\langle 0|c_{1-\lambda}\cdots c_{-1}c_0c_1\cdots c_{\lambda-1}|0\rangle=1.$$

$$\langle 0^c|=\langle 0|c_{1-\lambda}\cdots c_{-1}c_0c_1\cdots c_{\lambda-1}, \langle\downarrow^c|=\langle\uparrow|.$$



Cuantización BRST. Reparametrización e invariancia.

$$c = c_m + c_{\text{gh}} = c_m - 26, T(z) = T^m(z) + T^{\text{gh}}(z), \mathcal{H} = \mathcal{H}_m \otimes \mathcal{H}_{\text{gh}}.$$

$$Q_B |\psi\rangle = 0$$

$$|\psi\rangle \sim |\psi\rangle + Q_B |\Lambda\rangle.$$

$$\begin{aligned} j_B(z) &=: c(z) \left(T^m(z) + \frac{1}{2} T^{\text{gh}}(z) \right) : + \kappa \partial^2 c(z) \\ &= c(z) T^m(z) + : b(z) c(z) \partial c(z) : + \kappa \partial^2 c(z) \end{aligned}$$

$$Q_B = Q_{B,L} + Q_{B,R}, Q_{B,L} = \oint \frac{dz}{2\pi i} j_B(z), Q_{B,R} = \oint \frac{d\bar{z}}{2\pi i} \bar{j}_B(\bar{z}).$$

$$T(z)j_B(w) \sim \left(\frac{c_m}{2} - 4 - 6\kappa \right) \frac{c(w)}{(z-w)^4} + (3-2\kappa) \frac{\partial c(w)}{(z-w)^3} + \frac{j_B(w)}{(z-w)^2} + \frac{\partial j_B(w)}{z-w}.$$

$$c_m = 26, \kappa = \frac{3}{2}.$$

$$T(z)j_B(w) \sim \frac{j_B(w)}{(z-w)^2} + \frac{\partial j_B(w)}{z-w}.$$

$$\begin{aligned} j_B(z)b(w) &\sim \frac{2\kappa}{(z-w)^3} + \frac{j(w)}{(z-w)^2} + \frac{T(w)}{z-w} \\ j_B(z)c(w) &\sim \frac{:c(w)\partial c(w)}{z-w} \end{aligned}$$

$$j_B(z)\phi(w) \sim h \frac{c(w)\phi(w)}{(z-w)^2} + \frac{:h\partial c(w)\phi(w) + c(w)\partial\phi(w):}{z-w}$$

$$j_B(z)j(w) \sim \frac{2\kappa+1}{(z-w)^3} - \frac{2\partial c(w)}{(z-w)^2} - \frac{j_B(w)}{z-w}$$

$$j_B(z)j_B(w) \sim -\frac{c_m - 18}{2} \frac{:c(w)\partial c(w):}{(z-w)^3} - \frac{c_m - 18}{4} \frac{:c(w)\partial^2 c(w):}{(z-w)^2} - \frac{c_m - 26}{12} \frac{:c(w)\partial^3 c(w):}{z-w}.$$

$$\begin{aligned} Q_B &= \sum_m :c_m \left(L_{-m}^m + \frac{1}{2} L_{-m}^{\text{gh}} \right) : \\ &= \sum_m c_{-m} L_m^m + \frac{1}{2} \sum_{m,n} (n-m) :c_{-m} c_{-n} b_{m+n}: \end{aligned}$$

$$\begin{aligned} Q_B &= \sum_m {}^* c_m \left(L_{-m}^m + \frac{1}{2} L_{-m}^{\text{gh}} \right) {}^* - \frac{c_0}{2} \\ &= \sum_n c_m L_{-m}^m + \frac{1}{2} \sum_{m,n} (n-m) {}^* c_{-m} c_{-n} b_{m+n} {}^* - c_0, \end{aligned}$$

$$Q_B = c_0 L_0 - b_0 M + \hat{Q}_B$$



$$\begin{aligned}\hat{Q}_B &= \sum_{m\neq 0} c_{-m} L_m^m - \frac{1}{2} \sum_{\substack{m,n\neq 0 \\ m+n\neq 0}} (m-n)_*^* c_{-m} c_{-n} b_{m+n*}^*, \\ M &= \sum_{m\neq 0} m c_{-m} c_m\end{aligned}$$

$$[L_0,M]=\left[\hat{Q}_B,M\right]=\left[\hat{Q}_B,L_0\right]=0,\hat{Q}_B^2=L_0M$$

$$\begin{aligned}\{Q_B,b(z)\}&=T(z)\\\{Q_B,c(z)\}&=c(z)\partial c(z)\\[Q_B,\phi(z)]&=h\partial c(z)\phi(z)+c(z)\partial\phi(z)\end{aligned}$$

$$[Q_B,\phi(z)]=\partial(c(z)\phi(z))$$

$$\{Q_B,c(z)\phi(z)\}=(1-h)c(z)\partial c(z)\phi(z)$$

$$[Q_B,j(z)]=-j_B(z)$$

$$\left[N_{\rm gh},Q_B\right]=Q_B$$

$$\{Q_B,Q_B\}=0$$

$$[Q_B,T(z)]=0$$

$$c_m = 26.$$

$$L_n=\{Q_B,b_n\}$$

$$[Q_B,L_n]=0$$

Cohomología BRST para sectores holomórficos y antiholomórficos.

$$k_\parallel^2=\epsilon_0(k^0)^2+(k^1)^2.$$

$$\mathcal{H}:=\mathcal{H}_{\parallel}\otimes\mathcal{H}_{\perp}, \mathcal{H}_{\parallel}:=\int~{\rm d} k^0\mathcal{F}_0(k^0)\otimes\int~{\rm d} k^1\mathcal{F}_1(k^1)\otimes\mathcal{H}_{\rm gh}$$

$$|\psi\rangle=|\psi_{\parallel}\rangle\otimes|\psi_{\perp}\rangle,$$

$$\begin{aligned}|\psi_{\parallel}\rangle&=c_0^{N_0^c}\prod_{m>0}(\alpha_{-m}^0)^{N_m^0}(\alpha_{-m}^1)^{N_m^1}(b_{-m})^{N_m^b}(c_{-m})^{N_m^c}|k^0,k^1,\downarrow\rangle\\|k^0,k^1,\downarrow\rangle&:=|k^0\rangle\otimes|k^1\rangle\otimes|\downarrow\rangle,N_m^0,N_m^1\in\mathbb{N},N_m^b,N_m^c=0,1.\end{aligned}$$

$$\mathcal{H}_0=\mathcal{H}\cap\ker b_0$$

$$\mathcal{H}=\mathcal{H}_0\oplus c_0\mathcal{H}_0.$$

$$L_0=L_0^m+L_0^{\text{gh}}=(L_0^m-1)+N^b+N^c,$$

$$L_0=\left(L_0^\perp-m_{\parallel,L}^2\ell^2-1\right)+\hat{L}_0^\parallel$$



$$m_{\parallel,L}^2=-p_{\parallel,L}^2,\hat{L}_0^{\parallel}=N^0+N^1+N^b+N^c\in \mathbb{N}.$$

$$\text{on-shell: } L_0|\psi\rangle=0$$

$$\mathcal{H}_{\rm abs}(Q_B)\!:=\!\{|\psi\rangle\in\mathcal{H}|Q_B|\psi\rangle=0,\nexists|\chi\rangle\in\mathcal{H}||\psi\rangle=Q_B|\chi\rangle\}.$$

$$\{Q,\Delta\}=1$$

$$|\psi\rangle=\{Q_B,\Delta\}|\psi\rangle=Q_B(\Delta|\psi\rangle)$$

$$\Delta\!\!:=\!\frac{b_0}{L_0}$$

$$L_0 = \{Q_B,b_0\}$$

$$Q_B|\psi\rangle=0, L_0|\psi\rangle\neq 0$$

$$|\psi\rangle=Q_B\left(\frac{b_0}{L_0}|\psi\rangle\right)$$

$$\mathcal{H}_{\rm abs}(Q_B)\subset\ker L_0$$

$$|\psi\rangle=\frac{L_0}{L_0}|\psi\rangle=\frac{1}{L_0}\{Q_B,b_0\}|\psi\rangle=\frac{1}{L_0}Q_B(b_0|\psi\rangle)$$

$$|\psi\rangle\in\ker L_0\colon P_0|\psi\rangle=|\psi\rangle, |\psi\rangle\in(\ker L_0)^\perp\colon P_0|\psi\rangle=0.$$

$$\{Q_B,\Delta(1-P_0)\}=(1-P_0).$$

$$\mathcal{H}_0\!:=\mathcal{H}\cap\ker b_0=\mathcal{H}_m\otimes\mathcal{H}_{\mathrm{gh},0}$$

$$|\psi\rangle\in\mathcal{H}_0\implies b_0|\psi\rangle=0.$$

$$b_0|\psi\rangle=Q_B|\psi\rangle=0, L_0|\psi\rangle\neq 0\implies |\psi\rangle=0.$$

$$\mathcal{H}_{\rm rel}(Q_B)\!:=\mathcal{H}_0(Q_B)=\{|\psi\rangle\in\mathcal{H}_0|Q_B|\psi\rangle=0,\nexists|\chi\rangle\in\mathcal{H}||\psi\rangle=Q_B|\chi\rangle\}.$$

$$Q_B|\psi\rangle=0, L_0|\psi\rangle=0, b_0|\psi\rangle=0$$

$$Q_B=c_0L_0-b_0M+\hat{Q}_B,\hat{Q}_B^2=L_0M$$

$$|\psi\rangle\in\mathcal{H}_0\cap\ker L_0\implies Q_B|\psi\rangle=\hat{Q}_B|\psi\rangle,\hat{Q}_B^2|\psi\rangle=0$$

$$\mathcal{H}_0(Q_B)=\mathcal{H}_0\big(\hat{Q}_B\big).$$

$$X_L^{\pm}=\frac{1}{\sqrt{2}}\bigg(X_L^0\pm\frac{\mathrm{i}}{\sqrt{\epsilon_0}}X_L^1\bigg)$$



$$\begin{aligned}\alpha_n^{\pm} &= \frac{1}{\sqrt{2}} \left(\alpha_n^0 \pm \frac{i}{\sqrt{\epsilon_0}} \alpha_n^1 \right), n \neq 0, \\ x_L^{\pm} &= \frac{1}{\sqrt{2}} \left(x_L^0 \pm \frac{i}{\sqrt{\epsilon_0}} x_L^1 \right), p_L^{\pm} = \frac{1}{\sqrt{2}} \left(p_L^0 \pm \frac{i}{\sqrt{\epsilon_0}} p_L^1 \right),\end{aligned}$$

$$[\alpha_m^+, \alpha_n^-] = \epsilon_0 m \delta_{m+n,0}, [x_L^{\pm}, p_L^{\mp}] = i \epsilon_0.$$

$$\begin{aligned}2p_L^+p_L^- &= (p_L^0)^2 + \epsilon_0(p_L^1)^2 = \epsilon_0 p_{\parallel,L}^2, \\ x^+p^- + x^-p^+ &= x^0p^0 + \epsilon_0 x^1p^1, \\ \sum_n \alpha_n^+ \alpha_{m-n}^- &= \frac{1}{2} \sum_n (\alpha_n^0 \alpha_{m-n}^0 + \epsilon_0 \alpha_n^1 \alpha_{m-n}^1).\end{aligned}$$

$$N_n^{\pm} = \frac{\epsilon_0}{n} \alpha_{-n}^{\pm} \alpha_n^{\mp}, N^{\pm} = \sum_{n>0} n N_n^{\pm}.$$

$$N^+ + N^- = N^0 + N^1.$$

$$L_0 = (L_0^\perp - m_{\parallel,L}^2 \ell^2 - 1) + \hat{L}_0^\parallel$$

$$m_{\parallel,L}^2 = -2\epsilon_0 p_L^+ p_L^-, \hat{L}_0^\parallel = N^+ + N^- + N^b + N^c$$

$$L_m^0 + L_m^1 = \epsilon_0 \sum_n : \alpha_n^+ \alpha_{m-n}^- : = \epsilon_0 \sum_{n \neq 0, m} : \alpha_n^+ \alpha_{m-n}^- : + \epsilon_0 (\alpha_0^- \alpha_m^+ + \alpha_m^+ \alpha_m^-).$$

$$\begin{aligned}[\alpha_m^+, \alpha_n^{\pm}] &= \frac{1}{2} \left[\left(\alpha_m^0 + \frac{i}{\sqrt{\epsilon_0}} \alpha_m^1 \right), \left(\alpha_n^0 \pm \frac{i}{\sqrt{\epsilon_0}} \alpha_n^1 \right) \right] \\ &= \frac{1}{2} \left([\alpha_m^0, \alpha_n^0] \mp \frac{1}{\epsilon_0} [\alpha_m^1, \alpha_n^1] \right) = \frac{\epsilon_0}{2} m \delta_{m+n,0} (1 \mp 1),\end{aligned}$$

$$\begin{aligned}[x_L^-, p_L^{\pm}] &= \frac{1}{2} \left[\left(x_L^0 - \frac{i}{\sqrt{\epsilon_0}} x_L^1 \right), \left(p_L^0 \pm \frac{i}{\sqrt{\epsilon_0}} p_L^1 \right) \right] \\ &= \frac{1}{2} ([x_L^0, p_L^0] \pm \epsilon_0 [x_L^1, p_L^1]) = \frac{\epsilon_0}{2} (1 \pm 1)\end{aligned}$$

$$\begin{aligned}\sum_n \alpha_n^+ \alpha_{m-n}^- &= \frac{1}{2} \sum_n \left(\alpha_n^0 + \frac{i}{\sqrt{\epsilon_0}} \alpha_n^1 \right) \left(\alpha_{m-n}^0 - \frac{i}{\sqrt{\epsilon_0}} \alpha_{m-n}^1 \right) \\ &= \frac{1}{2} \sum_n \left(\alpha_n^0 \alpha_{m-n}^0 + \epsilon_0 \alpha_n^1 \alpha_{m-n}^1 + \frac{i}{\sqrt{\epsilon_0}} (\alpha_{m-n}^0 \alpha_n^1 - \alpha_n^0 \alpha_{m-n}^1) \right)\end{aligned}$$

$$N^0 + N^1 = \sum_n n(N_n^0 + N_n^1) = \sum_n n(N_n^+ + N_n^-) = N^+ + N^-.$$

$$\hat{Q}_B = \sum_{m \neq 0} c_{-m} \left(L_m^\perp + \epsilon_0 \sum_n \alpha_n^+ \alpha_{m-n}^- \right) + \frac{1}{2} \sum_{m,n} (n-m) : c_{-m} c_{-n} b_{m+n} :$$

$$\deg := N^+ - N^- + \widehat{N}^c - \widehat{N}^b$$



$$\forall m\neq 0:\deg(\alpha_m^+)=\deg(c_m)=1,\deg(\alpha_m^-)=\deg(b_m)=-1$$

$$\hat{Q}_B=Q_0+Q_1+Q_2, \deg(Q_j)=j$$

$$Q_1=\sum_{m\neq 0}c_{-m}L_m^\perp+\sum_{\substack{m,n\neq 0\\ m+n\neq 0}}{}^*c_{-m}\left(\epsilon_0\alpha_n^+\alpha_{m-n}^-+\frac{1}{2}(m-n)c_{-m}b_{m+n}\right){}^*\nonumber\\ Q_0=\sum_{n\neq 0}\alpha_0^+c_{-n}\alpha_n^-, Q_2=\sum_{n\neq 0}\alpha_0^-c_{-n}\alpha_n^+$$

$$Q_0^2=Q_2^2=0, \{Q_0,Q_1\}=\{Q_1,Q_2\}=0, Q_1^2+\{Q_0,Q_2\}=0$$

$${\mathcal H}_0\big(\hat{Q}_B\big)\simeq {\mathcal H}_0(Q_0)$$

$$\widehat{\Delta}\!:=\!\frac{B}{\widehat{L}_0^{\parallel}}, B\!:=\!\epsilon_0\sum_{n\neq 0}\frac{1}{\alpha_0^+}\alpha_{-n}^+b_n.$$

$$\widehat{L}_0^{\parallel}=\{Q_0,B\}\implies\{Q_0,\widehat{\Delta}\}=1.$$

$$\widehat{L}_0^{\parallel}|\psi\rangle=0,\implies N^{\pm}|\psi\rangle=N^c|\psi\rangle=N^b|\psi\rangle=0,$$

$$0=Q_0\widehat{L}_0^{\parallel}|\psi\rangle=\widehat{L}_0^{\parallel}Q_0|\psi\rangle.$$

$$Q_0|\psi\rangle=0$$

$$L_0=L_0^\perp-m_{\parallel,L}^2\ell^2-1=0$$

$$\{Q_0,Q_1\}|\psi_0\rangle=0\implies Q_0(Q_1|\psi_0\rangle)=0.$$

$$Q_1|\psi_0\rangle=:-Q_0|\psi_1\rangle\implies |\psi_1\rangle=-\frac{B}{\widehat{L}_0^{\parallel}}Q_1|\psi_0\rangle.$$

$$Q_1|\psi_0\rangle=\left\{Q_0,\frac{B}{\widehat{L}_0^{\parallel}}\right\}Q_1|\psi_0\rangle=Q_0\left(\frac{B}{\widehat{L}_0^{\parallel}}Q_1|\psi_0\rangle\right).$$

$$\{Q_0,Q_1\}|\psi_1\rangle=Q_0(Q_1|\psi_1\rangle+Q_2|\psi_0\rangle).$$

$$Q_1|\psi_1\rangle+Q_2|\psi_0\rangle=Q_0|\psi_2\rangle, |\psi_2\rangle=-\frac{B}{\widehat{L}_0^{\parallel}}(Q_1|\psi_1\rangle+Q_2|\psi_0\rangle)$$

$$\{Q_0,Q_1\}|\psi_1\rangle=Q_0Q_1|\psi_1\rangle-Q_1^2|\psi_0\rangle=Q_0Q_1|\psi_1\rangle+\{Q_0,Q_2\}|\psi_0\rangle.$$

$$|\psi_{k+1}\rangle=-\frac{B}{\widehat{L}_0^{\parallel}}(Q_1|\psi_k\rangle+Q_2|\psi_{k-1}\rangle)$$

$$|\psi\rangle=\sum_{k\in\mathbb{N}}|\psi_k\rangle.$$

$$\hat{Q}_B|\psi\rangle=0.$$



$$N_{\mathrm{gh}}(\psi)=N_{\mathrm{gh}}(\psi_0)=1$$

$$Q_1|\psi_0\rangle=Q_2|\psi_0\rangle=0$$

$$\begin{aligned}\hat{Q}_B|\psi\rangle &= \sum_{k\in\mathbb{N}}\hat{Q}_B|\psi_k\rangle \\&= Q_0|\psi_0\rangle + \underbrace{Q_1|\psi_0\rangle+Q_0|\psi_1\rangle}_{=0} + \underbrace{Q_2|\psi_0\rangle+Q_1|\psi_1\rangle+Q_0|\psi_2\rangle}_{=0}+\cdots\\&=0\end{aligned}$$

$$\mathcal{H}_{\text{abs}}(Q_B) = \mathcal{H}_{\text{rel}}(Q_B) \oplus c_0 \mathcal{H}_{\text{rel}}(Q_B)$$

$$\begin{gathered} |\psi\rangle=|k^0,k^1,\downarrow\rangle\otimes|\psi_\perp\rangle,\qquad|\psi_\perp\rangle\in\mathcal{H}_\perp\\ \big(L_0^\perp-m_{\parallel,L}^2\ell^2-1\big)|\psi\rangle=0,\quad p_{L,\parallel}^2=-m_{\parallel,L}^2\ell^2\end{gathered}$$

$$(L_0^m-1)|\psi\rangle=0,\forall n>0:L_n^m|\psi\rangle=0.$$

$$Q_B=c_0L_0-b_0M+\hat{Q}_B+\bar{c}_0\bar{L}_0-\bar{b}_0\bar{M}+\hat{\bar{Q}}_B$$

$$Q_B=c_0^+L_0^+-b_0^+M^++c_0^-L_0^--b_0^-M^-+\hat{Q}_B^+,$$

$$L_0^+=\left(L_0^{\perp+}-\frac{m_{\parallel}^2\ell^2}{2}-2\right)+\hat{L}_0^{\parallel+}, L_0^-=L_0^{\perp-}+\hat{L}_0^{\parallel-}$$

$$M^\pm:=\frac{1}{2}(M\pm\bar{M}).$$

$$L_0^+|\psi\rangle=L_0^-|\psi\rangle=0.$$

$$\hat{L}_0^{\parallel\pm}=N^0\pm\bar{N}^0+N^1\pm\bar{N}^1+N^b\pm\bar{N}^b+N^c\pm\bar{N}^c=0.$$

$$(L_0^m+\bar{L}_0^m-2)|\psi\rangle=0,(L_0^m-\bar{L}_0^m)|\psi\rangle=0$$

$$\forall n>0:L_n^m|\psi\rangle=\bar{L}_n^m|\psi\rangle=0$$

$$|\downarrow\downarrow\rangle=c(0)\bar{c}(0)|0\rangle=c_1\bar{c}_1|0\rangle.$$

$$\hat{L}_0|\psi_{\ell,\bar{\ell}}\rangle=\ell|\psi_{\ell,\bar{\ell}}\rangle,\hat{\bar{L}}_0|\psi_{\ell,\bar{\ell}}\rangle=\bar{\ell}|\psi_{\ell,\bar{\ell}}\rangle.$$

$$k^2=-m^2,m^2:=\frac{2}{\ell^2}(N+\bar{N}-2),$$

$$m^2\ell^2=-4<0.$$

$$\mathcal{V}(k,z,\bar{z})=c(z)\bar{c}(\bar{z})\mathrm{e}^{\mathrm{i} k\cdot X(z,\bar{z})}.$$



Campo gravitónico de gauge y campos cuánticos gravitacionales e interacciones locales y no locales de las partículas supermasivas en simetrías axiales.

Expansión de campo.

$$\Psi[X(\sigma), c(\sigma)] := \langle X(\sigma), c(\sigma) | \Psi \rangle$$

$$|\Psi\rangle = \sum_{\alpha} \int \frac{d^D k}{(2\pi)^D} \psi_{\alpha}(k) |\phi_{\alpha}(k)\rangle$$

$$\psi_{\alpha}(x) = \int \frac{d^D k}{(2\pi)^D} e^{ik \cdot x} \psi_{\alpha}(k)$$

$$|\Psi\rangle = \sum_r \psi_r |\phi_r\rangle$$

Campo escalar.

$$\phi(x) = \int \frac{d^D k}{(2\pi)^D} \phi(k) e^{ik \cdot x}$$

$$|\phi\rangle = \int \frac{d^D k}{(2\pi)^D} \phi(k) |k\rangle, \phi(k) = \langle k | \phi \rangle$$

$$\phi(x) = \langle x | \phi \rangle = \int \frac{d^D k}{(2\pi)^D} \langle x | k \rangle \langle k | \phi \rangle, \langle x | k \rangle = e^{ik \cdot x}$$

Warm-up.

$$(-\Delta + m^2) \phi(x) = 0.$$

$$k^2 = -m^2.$$

$$\phi(x) = \int dk \phi(k) e^{ikx}$$

$$\phi(x) = \langle x | \phi \rangle, \phi(k) = \langle k | \phi \rangle,$$

$$|\phi\rangle = \int dx \phi(x) |x\rangle = \int dk \phi(k) |k\rangle$$

$$K(x, x') := \langle x | K | x' \rangle = \delta(x - x') (-\Delta_x + m^2),$$

$$\int dx' K(x, x') \phi(x') = 0 \Leftrightarrow K|\phi\rangle = 0.$$

$$S = \frac{1}{2} \int dx \phi(x) (-\Delta + m^2) \phi(x) = \frac{1}{2} \int dx dx' \phi(x) K(x, x') \phi(x')$$

$$S = \frac{1}{2} \langle \phi | K | \phi \rangle$$



$$Q_B|\psi\rangle=0$$

$$\Phi\in \mathcal{H}$$

$$Q_B|\Phi\rangle=0$$

$$\langle A,B\rangle\colon=\langle A\mid B\rangle,$$

$$S=\frac{1}{2}\langle\Phi,Q_B\Phi\rangle=\frac{1}{2}\langle\Phi|Q_B|\Phi\rangle.$$

$$\langle A,B\rangle=(-1)^{|A||B|}\langle B,A\rangle,\langle Q_BA,B\rangle=-(-1)^{|A|}\langle A,Q_BB\rangle,$$

$$N_{\mathrm{gh}}(\Phi)=1$$

$$|\Phi|=1$$

$$|\Phi\rangle^\ddagger=|\Phi\rangle^t$$

$$\begin{aligned}\langle\Phi,Q_B\Phi\rangle&=(-1)^{|\Phi|(|Q_B\Phi|)}\langle Q_B\Phi,\Phi\rangle=(-1)^{|\Phi|(1+|\Phi|)}\langle Q_B\Phi,\Phi\rangle\\&=\langle Q_B\Phi,\Phi\rangle=-(-1)^{|\Phi|}\langle\Phi,Q_B\Phi\rangle\end{aligned}$$

$$|\Phi\rangle=|\Phi_\downarrow\rangle+c_0|\widetilde{\Phi}_\downarrow\rangle,$$

$$\Phi_\downarrow,\widetilde{\Phi}_\downarrow\in\mathcal{H}_0\implies b_0|\Phi_\downarrow\rangle=b_0|\widetilde{\Phi}_\downarrow\rangle=0.$$

$$N_{\mathrm{gh}}(\Phi_\downarrow)=1,N_{\mathrm{gh}}\big(\widetilde{\Phi}_\downarrow\big)=0.$$

$$Q_B=c_0L_0-b_0M+\hat{Q}_B$$

$$S=\frac{1}{2}\langle\Phi_\downarrow|c_0L_0|\Phi_\downarrow\rangle+\frac{1}{2}\langle\widetilde{\Phi}_\downarrow|c_0M|\widetilde{\Phi}_\downarrow\rangle+\langle\widetilde{\Phi}_\downarrow|c_0\hat{Q}_B|\Phi_\downarrow\rangle.$$

$$0=-M|\widetilde{\Phi}_\downarrow\rangle+\hat{Q}_B|\Phi_\downarrow\rangle,0=c_0L_0|\Phi_\downarrow\rangle+c_0\hat{Q}_B|\widetilde{\Phi}_\downarrow\rangle.$$

$$|\Phi\rangle=|\Phi_\downarrow\rangle+|\Phi_\uparrow\rangle,|\Phi_\downarrow\rangle=\Pi_s|\Phi\rangle,|\Phi_\uparrow\rangle=\bar{\Pi}_s|\Phi\rangle.$$

$$\Pi_sQ_B|\Phi\rangle=-b_0M|\Phi_\uparrow\rangle+\hat{Q}_B|\Phi_\downarrow\rangle,\bar{\Pi}_sQ_B|\Phi\rangle=c_0L_0|\Phi_\downarrow\rangle+\hat{Q}_B|\Phi_\uparrow\rangle,$$

$$\left[\Pi_s,\hat{Q}_B\right]=\left[\Pi_s,M\right]=\left[\Pi_s,L_0\right]=0$$

$$\begin{aligned}S&=\frac{1}{2}\langle\Phi,Q_B\Phi\rangle\\&=\frac{1}{2}\langle\Pi_s\Phi+\bar{\Pi}_s\Phi,Q_B\Phi\rangle\\&=\frac{1}{2}\langle\Pi_s\Phi,\bar{\Pi}_sQ_B\Phi\rangle+\frac{1}{2}\langle\bar{\Pi}_s\Phi,\Pi_sQ_B\Phi\rangle\\&=\frac{1}{2}\langle\Phi_\downarrow,c_0L_0\Phi_\downarrow+\hat{Q}_B\Phi_\uparrow\rangle+\frac{1}{2}\langle\Phi_\uparrow,-b_0M\Phi_\uparrow+\hat{Q}_B\Phi_\downarrow\rangle\\&=\frac{1}{2}\langle\Phi_\downarrow,c_0L_0\Phi_\downarrow\rangle+\frac{1}{2}\langle\Phi_\downarrow,\hat{Q}_B\Phi_\uparrow\rangle-\frac{1}{2}\langle\Phi_\uparrow,b_0M\Phi_\uparrow\rangle+\frac{1}{2}\langle\Phi_\uparrow,\hat{Q}_B\Phi_\downarrow\rangle.\end{aligned}$$



Invariancia de Gauge.

$$|\phi\rangle \sim |\psi\rangle + Q_B|\lambda\rangle.$$

$$|\Phi\rangle \rightarrow |\Phi'\rangle = |\Phi\rangle + \delta_\Lambda|\Phi\rangle, \delta_\Lambda|\Phi\rangle = Q_B|\Lambda\rangle N_{\text{gh}}(\Lambda) = 0$$

Gauge de Siegel y singularidad.

$$b_0|\Phi\rangle = 0$$

$$|\Lambda\rangle = -\Delta|\Phi\rangle, \Delta = \frac{b_0}{L_0}$$

$$0 = \{Q_B, b_0\}|\psi\rangle = L_0|\psi\rangle$$

$$b_0|\Phi'\rangle = b_0|\Phi\rangle + b_0Q_B|\Lambda\rangle = 0$$

$$b_0|\Phi\rangle = b_0\{Q_B, \Delta\}|\Phi\rangle = b_0Q_B\Delta|\Phi\rangle,$$

$$b_0Q_B(\Delta|\Phi\rangle + |\Lambda\rangle) = 0$$

$$A'_\mu = A_\mu + \partial_\mu \lambda$$

$$\partial^\mu A'_\mu = 0 \Rightarrow \Delta\lambda = -\partial^\mu A_\mu$$

$$\lambda = -\frac{k^\mu}{k^2} A_\mu$$

$$|\tilde{\Phi}_\downarrow\rangle = 0 \Rightarrow |\Phi\rangle = |\Phi_\downarrow\rangle.$$

$$S = \frac{1}{2}\langle\Phi|c_0L_0|\Phi\rangle,$$

$$L_0|\Phi\rangle = 0$$

$$\hat{Q}_B|\Phi\rangle = 0$$

$$\begin{aligned} S &= \frac{1}{2}\langle\Phi|Q_B\{c_0, b_0\}|\Phi\rangle = \frac{1}{2}\langle\Phi|Q_Bb_0c_0|\Phi\rangle \\ &= \frac{1}{2}\langle\Phi|\{b_0, Q_B\}c_0|\Phi\rangle - \frac{1}{2}\langle\Phi|b_0Q_Bc_0|\Phi\rangle \\ &= \frac{1}{2}\langle\Phi|c_0L_0|\Phi\rangle \end{aligned}$$

Campo expansivo, paridad y número fantasma.

$$n_r := N_{\text{gh}}(\phi_r), |\phi_r| = n_r \bmod 2.$$

$$\langle\phi_r^c|\phi_s\rangle = \delta_{rs}$$

$$|\phi_r\rangle = |\phi_{\downarrow,r}\rangle + |\phi_{\uparrow,r}\rangle, b_0|\phi_{\downarrow,r}\rangle = c_0|\phi_{\uparrow,r}\rangle = 0$$

$$|\psi_\uparrow\rangle = c_0|\tilde{\psi}_\downarrow\rangle, b_0|\tilde{\psi}_\downarrow\rangle = 0, N_{\text{gh}}(\psi_\uparrow) = N_{\text{gh}}(\tilde{\psi}_\downarrow) + 1.$$



$$|\Phi\rangle = \sum_r \psi_r |\phi_r\rangle$$

$$\forall r\colon |\Phi| = |\psi_r||\phi_r|.$$

$$G(\psi_r) = 1 - n_r$$

$$\Phi = \sum_{n \in \mathbb{Z}} \Phi_n, N_{\text{gh}}(\Phi_n) = n$$

$$\Phi = \Phi_+ + \Phi_-, \Phi_+ = \sum_{n > 1} \Phi_n, \Phi_- = \sum_{n \leq 1} \Phi_n$$

$$\forall n_r \neq n \colon \psi_r = 0$$

$$|\Phi_n\rangle = \delta(N_{\text{gh}} - n)|\Psi\rangle = \sum_r \delta(n_r - n)\psi_r |\phi_r\rangle$$

$$\begin{aligned} |\Phi\rangle &= |\Phi_\downarrow\rangle + |\Phi_\uparrow\rangle = |\Phi_\downarrow\rangle + c_0 |\widetilde{\Phi}_\downarrow\rangle, \\ |\Phi_\uparrow\rangle &= c_0 |\widetilde{\Phi}_\downarrow\rangle, |\widetilde{\Phi}_\downarrow\rangle = b_0 |\Phi_\uparrow\rangle, \end{aligned}$$

$$b_0 |\Phi_\downarrow\rangle = 0, c_0 |\Phi_\uparrow\rangle = 0, b_0 |\widetilde{\Phi}_\downarrow\rangle = 0$$

$$|\Phi_\downarrow\rangle = \sum_r \psi_{\downarrow,r} |\phi_{\downarrow,r}\rangle, |\Phi_\uparrow\rangle = \sum_r \psi_{\uparrow,r} |\phi_{\uparrow,r}\rangle.$$

$$Z = \int d\Phi_{\text{cl}} e^{-S[\Phi_{\text{cl}}]} = \int d\Phi_{\text{cl}} e^{-\frac{1}{2}\langle \Phi_{\text{cl}} | Q_B | \Phi_{\text{cl}} \rangle}$$

$$Z = \int \prod_s d\psi_s e^{-S[\{\psi_s\}]}$$

Transformaciones de gauge bajo la métrica de Faddeev-Popov.

$$F(\Phi_{\text{cl}}) := b_0 |\Phi_{\text{cl}}\rangle = 0.$$

$$\delta F = b_0 Q_B |\Lambda_{\text{cl}}\rangle,$$

$$\det \frac{\delta F}{\delta \Lambda_{\text{cl}}} = \det b_0 Q_B.$$

$$\det b_0 Q_B = \int dB' dC e^{-S_{\text{FP}}}, S_{\text{FP}} = -\langle B' | b_0 Q_B | C \rangle.$$

$$N_{\text{gh}}(B') = 3, N_{\text{gh}}(C) = 0.$$

$$|B'\rangle = \delta(N_{\text{gh}} - 3) \sum_r b'_r |\phi_r\rangle, |C\rangle = \delta(N_{\text{gh}}) \sum_r c_r |\phi_r\rangle,$$

$$|b_r| = |c_r| = 1.$$



$$|B'|=0, |C|=1.$$

$$\begin{array}{ll} \delta |C\rangle = Q_B |\Lambda_{-1}\rangle, & N_{\rm gh}(\Lambda_{-1})=-1 \\ \delta |B'\rangle = b_0 |\Lambda'\rangle, & N_{\rm gh}(\Lambda')=4 \end{array}$$

$$|\Lambda\rangle=|\Lambda_0\rangle+Q_B|\Lambda_{-1}\rangle,$$

$$|\Phi'_{\rm cl}\rangle\longrightarrow|\Phi_{\rm cl}\rangle+Q_B|\Lambda_0\rangle$$

$$\begin{array}{l} |B'\rangle=|B'_\downarrow\rangle+c_0|B\rangle, |B\rangle:=|\tilde{B}'_\downarrow\rangle, \\ |\Lambda'\rangle=|\Lambda'_\downarrow\rangle+c_0|\widetilde{\Lambda}'_\downarrow\rangle. \end{array}$$

$$\delta |B'_\downarrow\rangle=|\widetilde{\Lambda}'_\downarrow\rangle, \delta |B\rangle=0$$

$$|F'\rangle=c_0|B'\rangle=0\implies |B'_\downarrow\rangle=0.$$

$$\det\frac{\delta F'}{\delta \Lambda'}=\text{det}c_0b_0=\text{det}c_0\text{det}b_0=\frac{1}{2}\det\{b_0,c_0\}=\frac{1}{2}$$

$$\Delta_{\rm FP}=\int~{\rm d}B~{\rm d}C {\rm e}^{-S_{\rm FP}[B,C]}, S_{\rm FP}=\langle B|Q_B|C\rangle$$

$$b_0|B\rangle=0, N_{\rm gh}(B)=2, |B|=1$$

$$S_{\rm FP}=\frac{1}{2}(\langle B|Q_B|C\rangle+\langle C|Q_B|B\rangle).$$

$$|B\rangle=|B_\downarrow\rangle+c_0|\tilde{B}_\downarrow\rangle$$

$$\delta |B\rangle=Q_B|\Lambda_1\rangle$$

$$|\Phi\rangle=\sum_n|\Phi_n\rangle$$

$$b_0|\Phi\rangle=0\implies b_0|\Phi_n\rangle=0.$$

$$|\beta\rangle=\sum_{n\in\mathbb{Z}}|\beta_n\rangle.$$

$$Z=\int~{\rm d}\Phi~{\rm d}\beta {\rm e}^{-S[\Phi,\beta]}$$

$$\begin{aligned} S[\Phi,\beta] &= \frac{1}{2}\langle\Phi|Q_B|\Phi\rangle+\langle\beta|b_0|\Phi\rangle \\ &= \sum_{n\in\mathbb{Z}}\left(\frac{1}{2}\langle\Phi_{2-n}|Q_B|\Phi_n\rangle+\langle\beta_{4-n}|b_0|\Phi_n\rangle\right) \end{aligned}$$

$$\delta |\Phi\rangle=Q_B|\Lambda\rangle$$

$$|\Lambda\rangle=\sum_{n\in\mathbb{Z}}|\Lambda_n\rangle.$$



Acción de espacio – tiempo curvo.

$$|\Phi\rangle = \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \left(T(k) + A_\mu(k) \alpha_{-1}^\mu + i \sqrt{\frac{\alpha'}{2}} B(k) b_{-1} c_0 + \dots \right) |k, \downarrow\rangle$$

$$\begin{aligned} (\alpha' k^2 - 1)T(k) &= 0, k^2 A_\mu(k) + ik_\mu B(k) = 0, \\ k^\mu A_\mu(k) + iB(k) &= 0. \end{aligned}$$

$$k^2 A_\mu(k) - k_\mu k \cdot A(k) = 0$$

$$(\alpha' \Delta + 1)T = 0, B = \partial^\mu A_\mu, \Delta A_\mu = \partial_\mu B$$

$$\begin{aligned} Q_B &= c_0 L_0 - b_0 M + \hat{Q}_B \\ M &\sim 2c_{-1}c_1, \hat{Q}_B \sim c_1 L_{-1}^m + c_{-1} L_1^m, \\ L_1^m &\sim \alpha_0 \cdot \alpha_1, L_{-1}^m \sim \alpha_0 \cdot \alpha_{-1}. \end{aligned}$$

$$\begin{aligned} Q_B |\Phi\rangle &= \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} (T(k) c_0 L_0 |k, \downarrow\rangle + A_\mu(k) (c_0 L_0 + \eta_{\nu\rho} c_{-1} \alpha_1^\nu \alpha_0^\rho) \alpha_{-1}^\mu |k, \downarrow\rangle \\ &\quad + i \sqrt{\frac{\alpha'}{2}} B(k) (-2b_0 c_{-1} c_1 + \eta_{\nu\rho} c_1 \alpha_{-1}^\nu \alpha_0^\rho) b_{-1} c_0 |k, \downarrow\rangle) \\ &= \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} (T(k) (\alpha' k^2 - 1) c_0 |k, \downarrow\rangle \\ &\quad + A_\mu(k) (\alpha' k^2 c_0 \alpha_{-1}^\mu + \sqrt{2\alpha'} \eta_{\nu\rho} \eta^{\mu\nu} k^\rho c_{-1}) |k, \downarrow\rangle \\ &\quad + i \sqrt{\frac{\alpha'}{2}} B(k) (2c_{-1} + \sqrt{2\alpha'} \eta_{\nu\rho} k^\rho \alpha_{-1}^\nu c_0) |k, \downarrow\rangle) \\ &= \frac{1}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} (T(k) (\alpha' k^2 - 1) c_0 |k, \downarrow\rangle \\ &\quad + \alpha' (A_\mu(k) k^2 + ik^\mu B(k) c_0 \alpha_{-1}^\mu) |k, \downarrow\rangle \\ &\quad + \sqrt{2\alpha'} (k^\mu A_\mu(k) + iB(k)) c_{-1} |k, \downarrow\rangle) \\ |\Lambda\rangle &= \frac{i}{\sqrt{2\alpha'}} \int \frac{d^D k}{(2\pi)^D} (\lambda(k) b_{-1} |k, \downarrow\rangle + \dots). \end{aligned}$$

$$Q_B |\Lambda\rangle = \frac{i}{\sqrt{\alpha'}} \int \frac{d^D k}{(2\pi)^D} \lambda(k) \left(-\sqrt{\frac{\alpha'}{2}} k^2 b_{-1} c_0 + k_\mu \alpha_{-1}^\mu \right) |k, \downarrow\rangle,$$

$$\delta A_\mu = -ik_\mu \lambda, \delta B = k^2 \lambda.$$

$$|T\rangle = \int \frac{d^D k}{(2\pi)^D} T(k) c_1 |k, 0\rangle,$$



$$\begin{aligned}\langle T| &= \int \frac{\mathrm{d}^D k}{(2\pi)^D} T(k) \langle -k, 0 | c_{-1} \\ \langle T^\dagger | &= \int \frac{\mathrm{d}^D k}{(2\pi)^D} T(k)^* \langle k, 0 | c_{-1}\end{aligned}$$

$$T(k)^*=T(-k)$$

$$S[T] = \frac{1}{2} \int \frac{\mathrm{d}^D k}{(2\pi)^D} T(-k) \left(k^2 - \frac{1}{\alpha'} \right) T(k)$$

$$S[A]=\frac{1}{2}\int \frac{\mathrm{d}^D k}{(2\pi)^D} A_\mu(-k) k^2 A^\mu(k)$$

$$\frac{L_0^+}{2}=\frac{1}{2}(k^2+m^2)$$

$$\begin{aligned}\langle T|c_0L_0|T\rangle &= \frac{1}{\alpha'} \int \frac{\mathrm{d}^D k}{(2\pi)^D} \frac{\mathrm{d}^D k'}{(2\pi)^D} T(k) T(k') \langle -k', 0 | c_{-1} c_0 L_0 c_1 | k, 0 \rangle \\ &= \frac{1}{\alpha'} \int \frac{\mathrm{d}^D k}{(2\pi)^D} \frac{\mathrm{d}^D k'}{(2\pi)^D} T(k) T(k') (\alpha' k^2 - 1) \langle -k', 0 | c_{-1} c_0 c_1 | k, 0 \rangle \\ &= \frac{1}{\alpha'} \int \frac{\mathrm{d}^D k}{(2\pi)^D} \mathrm{d}^D k' T(k) T(k') (\alpha' k^2 - 1) \delta^{(D)}(k + k')\end{aligned}$$

$$Q_B|\Psi\rangle=0$$

$$S=\frac{1}{2}\langle\Psi,Q_B\Psi\rangle.$$

$$N_{\mathrm{gh}}(\Psi)=2$$

$$N_{\mathrm{gh}}(\langle\cdot,\cdot\rangle)=1.$$

$$\langle A,B\rangle=\langle A|c_0^-|B\rangle.$$

$$S=\frac{1}{2}\langle\Psi|c_0^-Q_B\Psi\rangle.$$

$${\mathcal H}={\mathcal H}^-\oplus c_0^-{\mathcal H}^-,\mathcal{H}^-:=\mathcal{H}\cap\ker b_0^-,$$

$$|\Psi\rangle=|\Psi_-\rangle+c_0^-|\widetilde{\Psi}_-\rangle,\Psi_-,\widetilde{\Psi}_-\in\mathcal{H}^-$$

$$c_0^-|\Psi\rangle=c_0^-|\Psi_-\rangle.$$

$$b_0^-|\Psi\rangle=0$$

$$L_0^-|\Psi\rangle=0,$$

$$\Psi\in\mathcal{H}^-\cap\ker L_0^-.$$

$$|\Psi\rangle\longrightarrow|\Psi'\rangle=|\Psi\rangle+\delta_\Lambda|\Psi\rangle,\delta_\Lambda|\Psi\rangle=Q_B|\Lambda\rangle,$$

$$N_{\mathrm{gh}}(\Lambda)=1,L_0^-|\Lambda\rangle=0,b_0^-|\Lambda\rangle=0$$



$$b_0^+|\Psi\rangle=0.$$

$$S=\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+|\Psi\rangle=\frac{1}{4}\langle\Psi|c_0\bar{c}_0L_0^+|\Psi\rangle.$$

$$L_0^+|\Psi\rangle=0.$$

$$c_0^-Q_B=(c_0-\bar c_0)(c_0L_0+\bar c_0\bar L_0)=c_0\bar c_0(L_0+\bar L_0).$$

Off-shell.

$$A_{0,3} = \left\langle \prod_{i=1}^3 \mathcal{V}_i(z_i) \right\rangle_{S^2} \propto (z_1 - z_2)^{h_3 - h_1 - h_2} \times \text{perms} \times \text{c.c.}$$

$$z \rightarrow f_g(z) = \frac{az+b}{cz+d} \in \mathrm{SL}(2,\mathbb{C})$$

$$z=f_i(w_i), z_i=f_i(0)$$

$$f \circ \mathcal{V}(w) = f'(w)^h \overline{f'(w)}^{\bar{h}} \mathcal{V}(f(w))$$

$$\begin{aligned} A_{0,3} &= \left\langle \prod_{i=1}^3 f_i \circ \mathcal{V}_i(0) \right\rangle_{S^2} = \left(\prod_{i=1}^3 f'_i(0)^{h_i} \overline{f'_i(0)}^{\bar{h}_i} \right) \left\langle \prod_{i=1}^3 \mathcal{V}_i(f_i(0)) \right\rangle_{S^2} \\ &\propto \left(\prod_{i=1}^3 f'_i(0)^{h_i} \overline{f'_i(0)}^{\bar{h}_i} \right) (f_1(0) - f_2(0))^{h_3 - h_1 - h_2} \times \text{perms} \times \text{c.c.} \end{aligned}$$

$$f_i \rightarrow \frac{af_i+b}{cf_i+d}$$

$$f'_i \rightarrow \frac{f'_i}{(cf_i+d)^2}, f_i - f_j \rightarrow \frac{f_i - f_j}{(cf_i+d)(cf_j+d)}$$

$$\mathcal{V}_{0,3}(\mathscr{V}_1, \mathscr{V}_2, \mathscr{V}_3) := \begin{array}{ccc} & \mathscr{V}_1 & \\ & \searrow & \\ & 0 & \\ & \swarrow & \\ \mathscr{V}_2 & & \mathscr{V}_3 \end{array} = A_{0,3}(\mathscr{V}_1, \mathscr{V}_2, \mathscr{V}_3).$$



$$A_{0,4} = \int d^2 z_4 \left\langle \prod_{i=1}^3 c\bar{c} V_i(z_i) V_4(z_4) \right\rangle_{S^2}$$

$$z_4 \longrightarrow z_1, z_2, z_3$$

$$A_{0,4} \propto \prod_{\substack{i,j=1 \\ i < j}}^3 |z_i - z_j|^{2+k_i \cdot k_j} \int d^2 z_4 \prod_{i=1}^3 |z_4 - z_i|^{k_i \cdot k_4}$$

$$\mathcal{F}_{0,4}^{(s)} = \int \frac{d^2 q}{|q|^2} \langle c\bar{c}V_1(z_1)c\bar{c}V_2(z_2)c\bar{c}V_3(0)|qy_4 \left| V_4(qy_4) \right\rangle$$

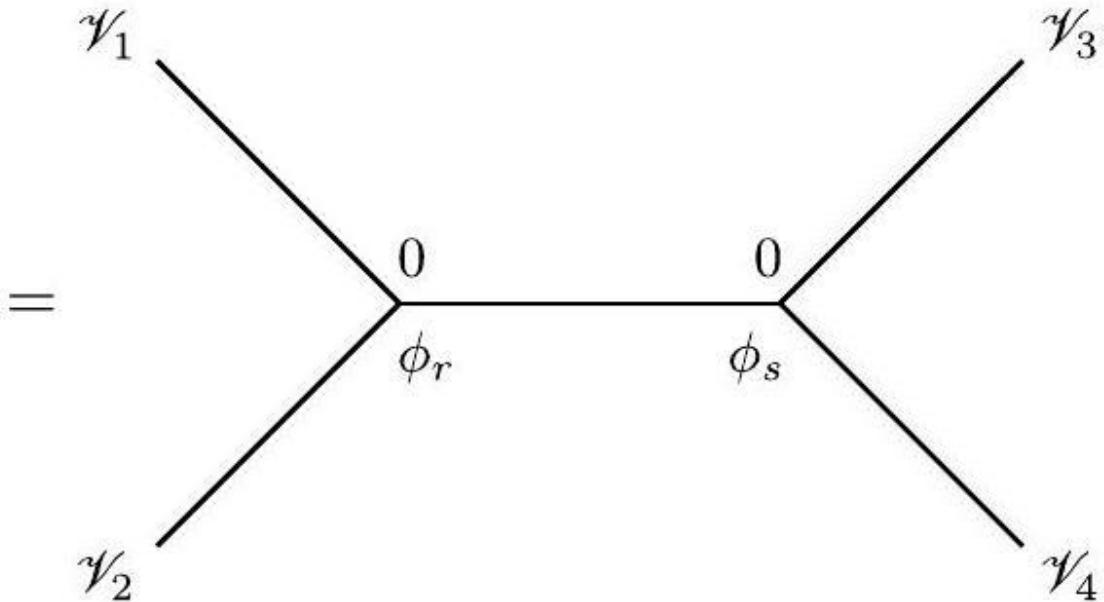
$$\mathcal{F}_{0,4}^{(s)} = - \int \frac{d^2 q}{|q|^2} \left\langle c\bar{c}V_1(z_1)c\bar{c}V_2(z_2)\oint_{|w|=|q|^{1/2}} dw w(w)\oint_{|w|=|q|^{1/2}} d\bar{w} \overline{b(w)} c\bar{c}V_4(qy_4)c\bar{c}V_3(0) \right\rangle$$

$$\mathcal{F}_{0,4}^{(s)} = - \int \frac{d^2 q}{|q|^2} \langle c\bar{c}V_1(z_1)c\bar{c}V_2(z_2)\oint dw w b(w)\oint d\bar{w} \overline{b(w)} q^{L_0} \bar{q}^{\bar{L}_0} c\bar{c}V_4(y_4)c\bar{c}V_3(0) \rangle$$

$$\mathcal{F}_{0,4}^{(s)} = \langle c\bar{c}V_1(z_1)c\bar{c}V_2(z_2)\phi_r(0) \rangle \langle c\bar{c}V_3(z_3)c\bar{c}V_4(y_4)\phi_s(0) \rangle \int \frac{d^2 q}{|q|^2} \langle \phi_r^c q^{L_0} \bar{q}^{\bar{L}_0} b_0 \bar{b}_0 \phi_s^c \rangle$$

$$\Delta(\phi_r^c, \phi_s^c) := \langle \phi_r^c | \Delta | \phi_s^c \rangle := - \int \frac{d^2 q}{|q|^2} \langle \phi_r^c q^{L_0} \bar{q}^{\bar{L}_0} b_0 \bar{b}_0 \phi_s^c \rangle$$

$$\mathcal{F}_{0,4}^{(s)} = \mathcal{V}_{0,3}(c\bar{c}V_1(z_1), c\bar{c}V_2(z_2), \phi_r(0)) \times \Delta(\phi_r^c, \phi_s^c) \times \mathcal{V}_{0,3}(c\bar{c}V_3(z_3), c\bar{c}V_4(y_4), \phi_s(0))$$



$$q = e^{-s+i\theta}, s \in \mathbb{R}_+, \theta \in [0, 2\pi)$$



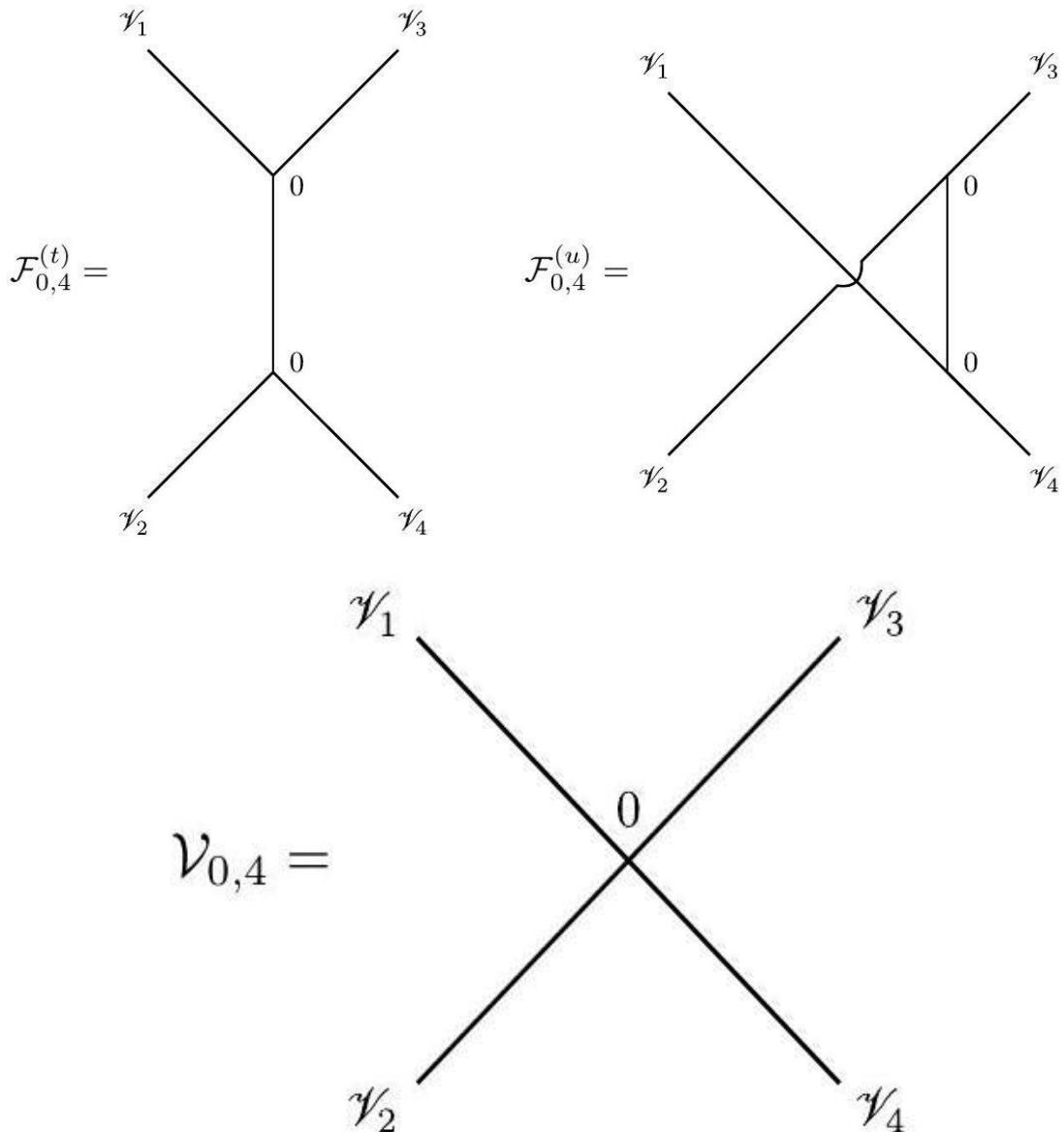
$$\int \frac{d^2 q}{|q|^2} q^{L_0} \bar{q}^{\bar{L}_0} = 2 \int_0^\infty ds \int_0^{2\pi} d\theta e^{-s(L_0 + \bar{L}_0)} e^{i\theta(L_0 - \bar{L}_0)} = \frac{2}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0}$$

$$\Delta = -\frac{2b_0\bar{b}_0}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0} = \frac{b_0^+}{L_0^+} b_0^- \delta_{L_0^-, 0}$$

$$L_0 |\phi_\alpha(k)\rangle = \bar{L}_0 |\phi_\alpha(k)\rangle = \frac{\alpha'}{4} (k^2 + m_\alpha^2) |\phi_\alpha(k)\rangle$$

$$\Delta_{\alpha\beta}(k) = \int \frac{d^2 q}{|q|^2} \langle \phi_\alpha^c(k) q^{L_0} \bar{q}^{\bar{L}_0} b_0 \bar{b}_0 \phi_\beta^c(-k) \rangle = \frac{M_{\alpha\beta}(k)}{k^2 + m_\alpha^2}$$

$$M_{\alpha\beta}(k) := \frac{2}{\alpha'} \langle \phi_\alpha^c(k) | b_0^+ b_0^- | \phi_\beta^c(-k) \rangle.$$



$$A_{0,4} = \mathcal{F}_{0,4}^{(s)} + \mathcal{F}_{0,4}^{(t)} + \mathcal{F}_{0,4}^{(u)} + \mathcal{V}_{0,4}.$$



Estados Off-shell.

$$\mathcal{H} = \mathcal{H}_m \otimes \mathcal{H}_{\text{gh}}.$$

$$\mathcal{H}=\text{Span}\{|\phi_r\rangle\}.$$

$$n_r:=N_{\text{gh}}(\phi_r)\in\mathbb{Z}$$

$$|\phi_r|:=N_{\text{gh}}(\phi_r) \bmod 2 = \begin{cases} 0 & N_{\text{gh}}(\phi_r) \text{ even} \\ 1 & N_{\text{gh}}(\phi_r) \text{ odd} \end{cases}$$

$$\langle \phi_r^c \mid \phi_s \rangle = \delta_{rs}$$

$$n_r^c:=N_{\text{gh}}(\phi_r^c)$$

$$n_r^c+n_r=6$$

$$\langle \phi_r \mid \phi_s \rangle = 0$$

$$\langle \phi_r \mid \phi_s^c \rangle = (-1)^{|\phi_r|} \delta_{rs}.$$

$$1=\sum_r |\phi_r\rangle\langle\phi_r^c|=\sum_r (-1)^{|\phi_r|}|\phi_r^c\rangle\langle\phi_r|.$$

$$\mathcal{H}=\mathcal{H}_{\pm}\oplus c_0^{\pm}\mathcal{H}_{\pm}$$

$$\mathcal{H}_{\pm}:=\mathcal{H}\cap\ker b_0^{\pm}=\mathcal{H}_0\oplus c_0^{\mp}\mathcal{H}_0,\mathcal{H}_0:=\mathcal{H}\cap\ker b_0^{-}\cap\ker b_0^{+}.$$

$$L_0^-|\phi\rangle=0,b_0^-|\phi\rangle=0$$

$$\mathcal{H}\sim\mathcal{H}_{\downarrow\downarrow}\oplus\mathcal{H}_{\downarrow\uparrow}\oplus\mathcal{H}_{\uparrow\downarrow}\oplus\mathcal{H}_{\uparrow\uparrow},$$

$$\mathcal{H}_{\downarrow\uparrow}\sim\mathcal{H}_0,\mathcal{H}_{\downarrow\uparrow}\sim\bar{c}_0\mathcal{H}_{\downarrow\downarrow}\mathcal{H}_{\uparrow\downarrow}\sim c_0\mathcal{H}_{\downarrow\downarrow}\mathcal{H}_{\uparrow\uparrow}\sim c_0\bar{c}_0\mathcal{H}_{\downarrow\downarrow}$$

$$\phi_r=\phi_{\downarrow\downarrow,r}+\phi_{\downarrow\uparrow,r}+\phi_{\uparrow\downarrow,r}+\phi_{\uparrow\uparrow,r}$$

$$\begin{aligned} b_0|\phi_{\downarrow\downarrow,r}\rangle&=\bar{b}_0|\phi_{\downarrow\downarrow,r}\rangle=0,& b_0|\phi_{\downarrow\uparrow,r}\rangle&=\bar{c}_0|\phi_{\downarrow\uparrow,r}\rangle=0,\\ c_0|\phi_{\uparrow\downarrow,r}\rangle&=\bar{b}_0|\phi_{\uparrow\downarrow,r}\rangle=0,& c_0|\phi_{\uparrow\uparrow,r}\rangle&=\bar{c}_0|\phi_{\uparrow\uparrow,r}\rangle=0.\end{aligned}$$

$$|\phi_{\downarrow\uparrow,r}\rangle=\bar{c}_0|\phi_{\downarrow\downarrow,r}\rangle\,|\phi_{\uparrow\downarrow,r}\rangle=c_0|\phi_{\downarrow\downarrow,r}\rangle\,|\phi_{\uparrow\uparrow,r}\rangle=c_0\bar{c}_0|\phi_{\downarrow\downarrow,r}\rangle.$$

$$\phi_r^c=\phi_{\downarrow\downarrow,r}^c+\phi_{\downarrow\uparrow,r}^c+\phi_{\uparrow\downarrow,r}^c+\phi_{\uparrow\uparrow,r}^c$$

$$\begin{aligned} \langle\phi_{\downarrow\downarrow,r}^c|c_0&=\langle\phi_{\downarrow\downarrow,r}^c|\bar{c}_0=0,& \langle\phi_{\downarrow\uparrow,r}^c|c_0&=\langle\phi_{\downarrow\uparrow,r}^c|\bar{b}_0=0,\\ \langle\phi_{\uparrow\downarrow,r}^c|b_0&=\langle\phi_{\uparrow\downarrow,r}^c|\bar{c}_0=0,& \langle\phi_{\uparrow\downarrow,r}^c|b_0&=\langle\phi_{\uparrow\downarrow,r}^c|\bar{b}_0=0.\end{aligned}$$

$$\langle\phi_{\downarrow\uparrow,r}^c|=\langle\phi_{\downarrow\downarrow,r}^c|\bar{b}_0\,\langle\phi_{\uparrow\downarrow,r}^c|=\langle\phi_{\downarrow\downarrow,r}^c|b_0\,\langle\phi_{\uparrow\uparrow,r}^c|=\langle\phi_{\downarrow\downarrow,r}^c|\bar{b}_0b_0.$$

$$\langle\phi_{x,r}^c\mid\phi_{y,s}\rangle=\delta_{xy}\delta_{rs}$$

$$b_0^-|\phi_r\rangle=0\implies b_0|\phi_r\rangle=\bar{b}_0|\phi_r\rangle$$



$$\phi_{\uparrow\downarrow,r} + \phi_{\uparrow\uparrow,r} = \phi_{\downarrow\downarrow,r} + \phi_{\downarrow\uparrow,r}$$

$$\phi_r = 2(\phi_{\downarrow\downarrow,r} + \phi_{\downarrow\uparrow,r})$$

Amplitudes Off-shell.

$$\chi_{g,n} := \chi(\Sigma_{g,n}) = 2 - 2g - n.$$

$$M_{g,n} := \dim_{\mathbb{R}} \mathcal{M}_{g,n} = 6g - 6 + 2n, \text{ for } \begin{cases} g \geq 2, \\ g = 1, n \geq 1, \\ g = 0, n \geq 3. \end{cases}$$

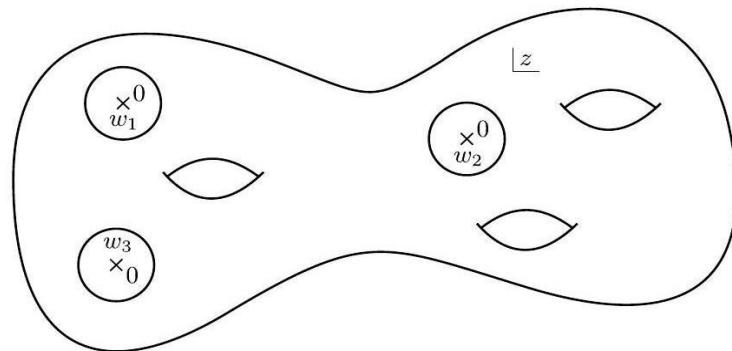
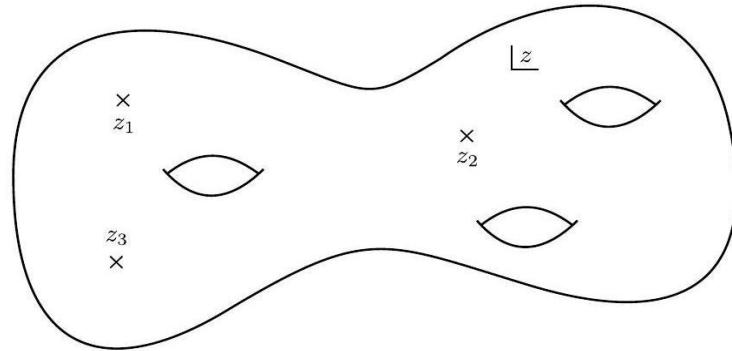
$$A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) = \int_{\mathcal{M}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)$$

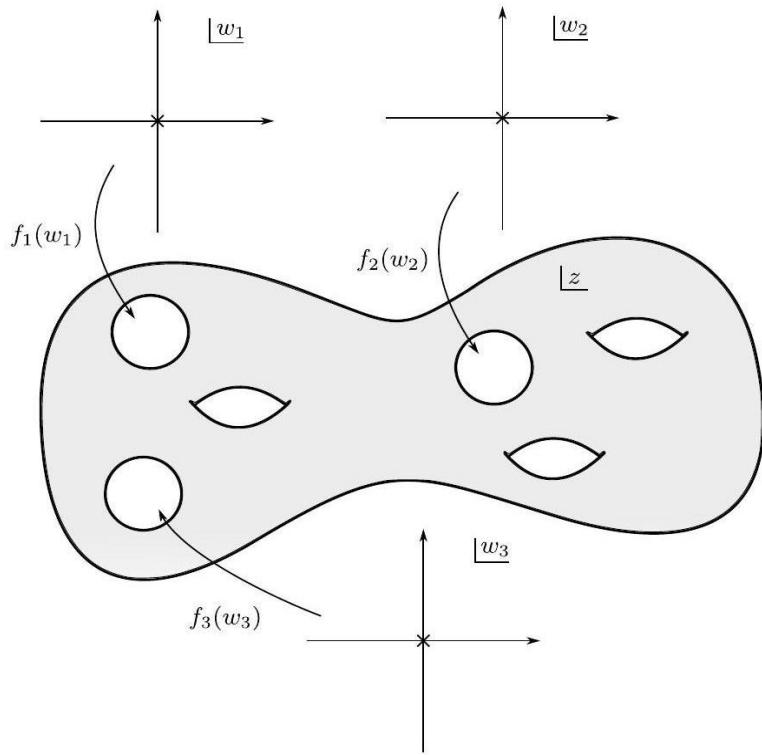
$$\omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) = \left\langle \text{ghosts} \times \prod_{i=1}^n \mathcal{V}_i \right\rangle_{\Sigma_{g,n}} \times \bigwedge_{\lambda=1}^{M_{g,n}} dt_\lambda$$

Coordenadas locales.

$$z = f_i(w_i), z_i = f_i(0)$$

$$A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)_{\mathcal{S}_{g,n}} = \int_{\mathcal{S}_{g,n}} \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) \Big|_{\mathcal{S}_{g,n}}$$





Figuras 4, 5 y 6. Curvatura local provocada por una partícula supermasiva.

$$\forall \mathcal{S}_{g,n}: A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)_{\mathcal{S}_{g,n}} = A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) \text{ (on-shell).}$$

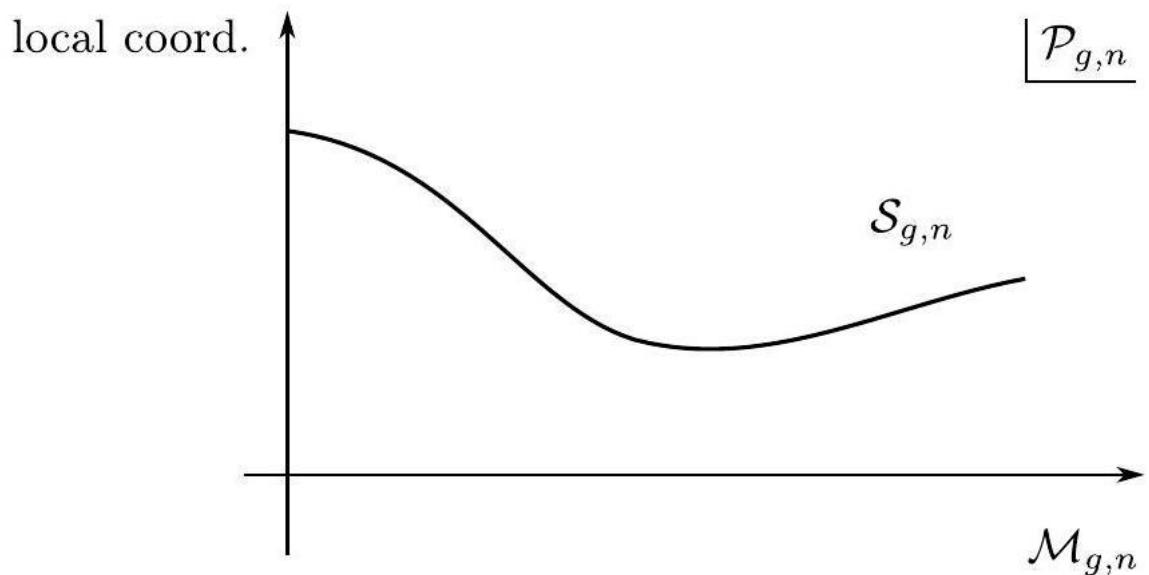


Figura 7. Fluctuaciones en el sistema de coordenadas de una partícula supermasiva al deformar o curvar el espacio – tiempo cuántico.

Geometría en espacios de Moduli y superficies de Riemann.

$$\text{disks} = n.$$

$$\backslash \# \text{spheres} = 2g - 2 + n.$$

$$\text{circles} = \frac{n + 3(2g - 2 + n)}{2} = 3g - 3 + 2n$$

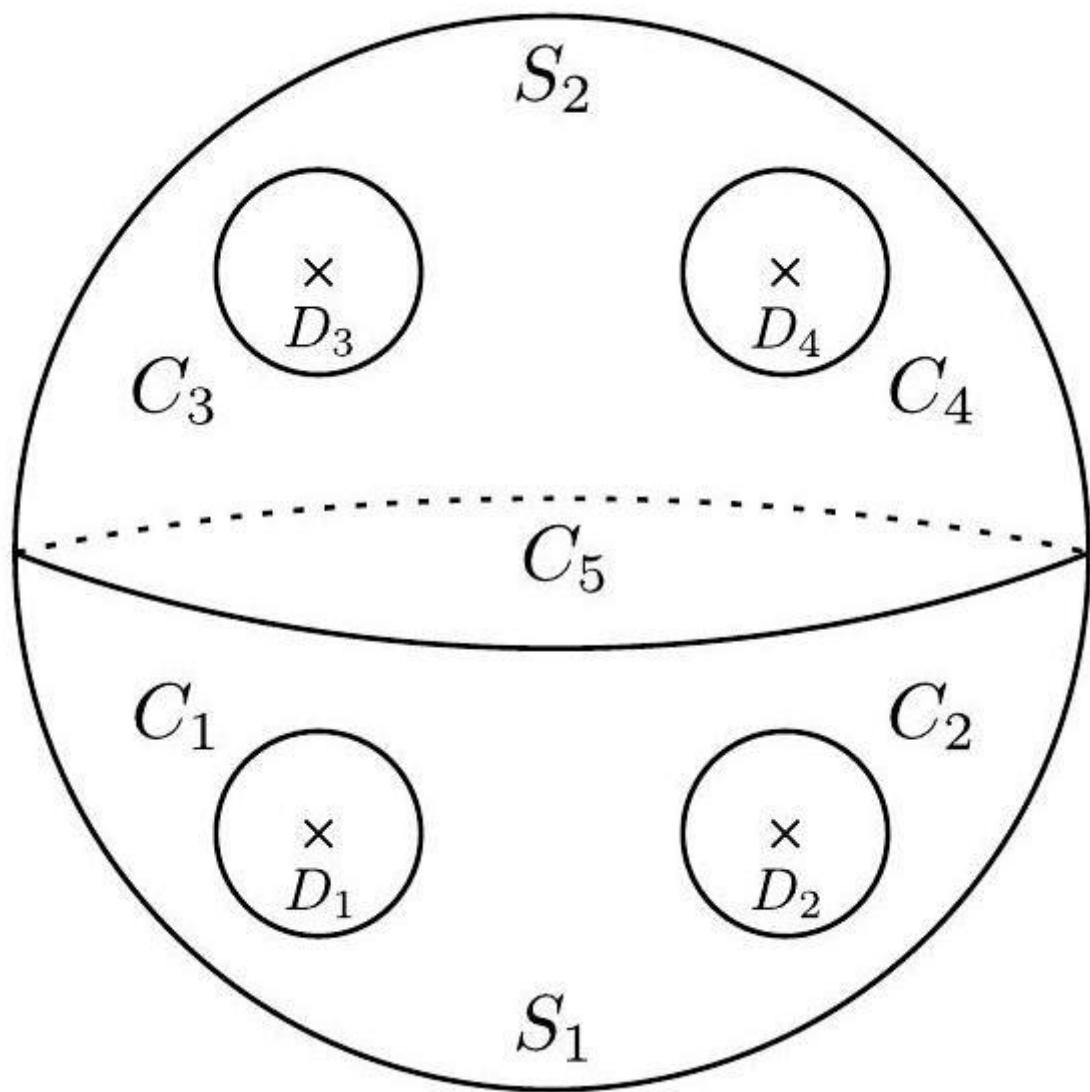
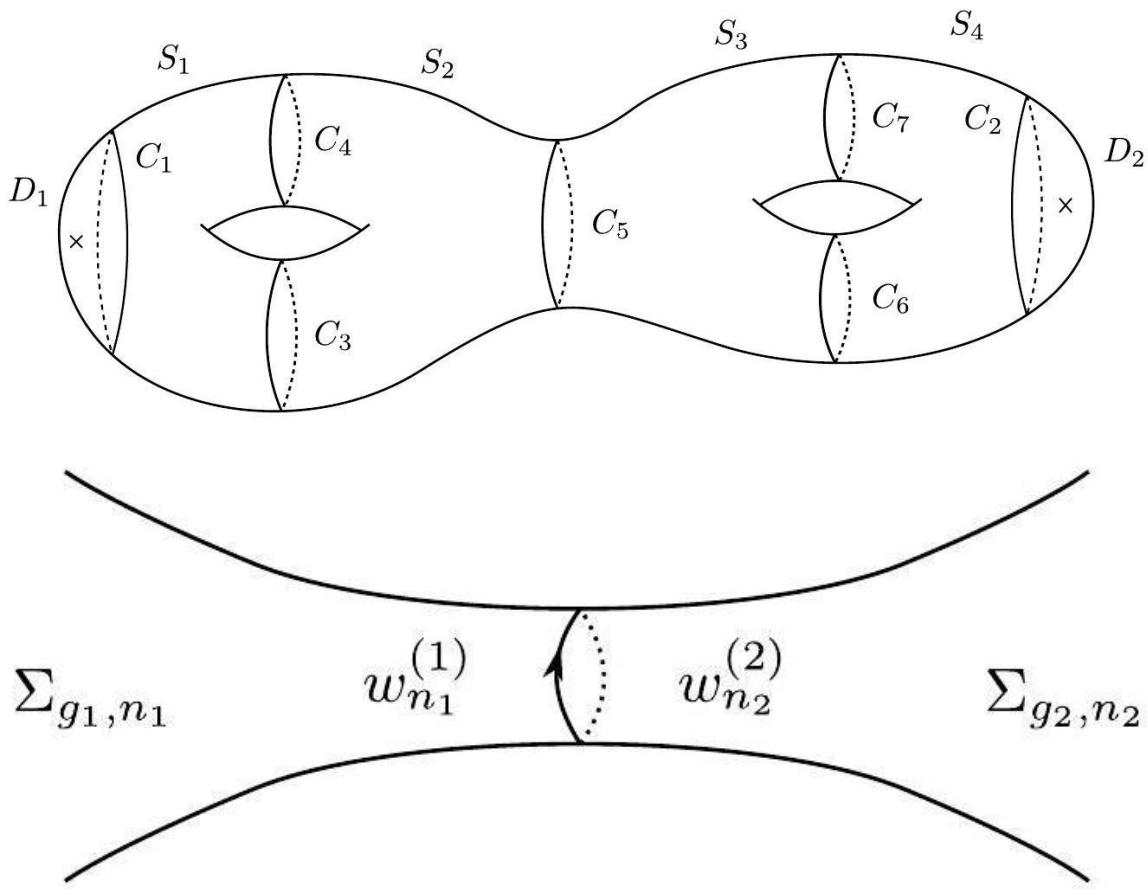


Figura 8. Propagadores de una partícula repercutida en relación al propagador de la partícula supermasiva.



Figuras 9 y 10. Entrelazamiento cuántico por curvatura local en dimensión \mathbb{R}^4 , es decir, en la misma dimensión.

$$C_{\Lambda(ab)} := S_a \cap S_b, C_{i(a)} := S_a \cap D_i, \{C_\alpha\} = \{C_\Lambda, C_i\}$$

$$\Lambda = 1, \dots, M_{g,n}^c, M_{g,n}^c = 3g - 3 + n$$

$$\begin{aligned} \text{on } C_{\Lambda(ab)}: z_a &= F_{ab}(z_b) \\ \text{on } C_{i(a)}: z_a &= f_{ai}(w_i) \end{aligned}$$

$$w_i \rightarrow \tilde{w}_i = e^{i\alpha_i} w_i$$

$$z = f_i(w_i), z = \tilde{f}_i(\tilde{w}_i)$$

$$f_i(w_i) = \tilde{f}_i(e^{i\alpha_i} w_i)$$

$$\hat{\mathcal{P}}_{g,n} = \mathcal{P}_{g,n}/U(1)^n$$

$$\text{on } C_{\Lambda(ab)}: z_a - z_{a,m} = \frac{q_\Lambda}{z_b - z_{b,n}}$$

$$\text{on } C_{i(a)}: z_a - z_{a,m} = w_i + \sum_{N=1}^{\infty} p_{i,N} w_i^N$$

$$\{x_s\}=\left\{q_{\Lambda}, p_{i,N}\right\}$$

$$z_a=F_{ab}(z_b;x_s), z_a=f_{ai}(w_i;x_s)$$

$$\text{on }C_\alpha \colon \sigma_\alpha = F_\alpha(\tau_\alpha;x_s)$$

Espacio tangencial.

$$\delta x_s = \epsilon V_s$$

$$\epsilon V_s \partial_s f = f(x_s + \epsilon V_s) - f(x_s)$$

$$F_\alpha \longrightarrow F_\alpha + \epsilon \delta F_\alpha$$

$$\sigma_\alpha = F_\alpha(\tau_\alpha)$$

$$\sigma'_\alpha = F_\alpha(\tau_\alpha) + \epsilon \delta F_\alpha(\tau_\alpha) = \sigma_\alpha + \epsilon \delta F_\alpha(\tau_\alpha) = \sigma_\alpha + \epsilon \delta F_\alpha\big(F_\alpha^{-1}(\sigma_\alpha)\big)$$

$$\sigma'_\alpha = \sigma_\alpha + \epsilon v^{(\alpha)}(\sigma_\alpha), v^{(\alpha)} = \delta F_\alpha \circ F_\alpha^{-1}$$

$$V^{(\alpha)} \sim \big(v^{(\alpha)}, \mathcal{C}_{(\alpha)}\big).$$

$$V \sim (\nu,\mathcal{C}), \mathcal{C} \subseteq \bigcup_\alpha \mathcal{C}_\alpha$$

$$v|_{\mathcal{C}_\alpha}=v^{(\alpha)}$$

$$x_s \longrightarrow x_s + \epsilon \delta x_s$$

$$C_\alpha \colon F_\alpha \longrightarrow F_\alpha + \epsilon \delta F_\alpha, \delta F_\alpha = \frac{\partial F_\alpha}{\partial x_s} \delta x_s$$

$$\sigma'_\alpha = \sigma_\alpha + \epsilon v^{(\alpha)}_s(\sigma_\alpha) \delta x_s, v^{(\alpha)}_s(\sigma_\alpha) = \frac{\partial F_\alpha}{\partial x_s}\big(F_\alpha^{-1}(\sigma_\alpha)\big)$$

Plumbing fixture.

$$D_q^{(1)}=\Big\{\Big|w_{n_1}^{(1)}\Big|\leq |q|^{1/2}\Big\}, D_q^{(2)}=\Big\{\Big|w_{n_2}^{(2)}\Big|\leq |q|^{1/2}\Big\}$$

$$\Sigma_{g,n}=\Sigma_{g_1,n_1}\qquad\qquad\qquad\Sigma_{g_2,n_2},\begin{cases} g=g_1+g_2\\ n=n_1+n_2-2\end{cases}$$

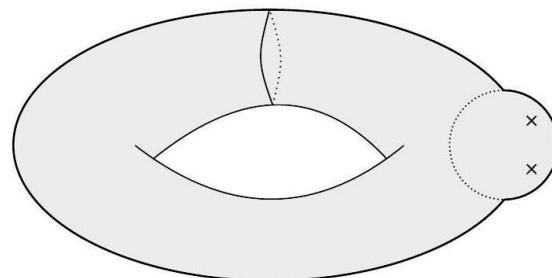
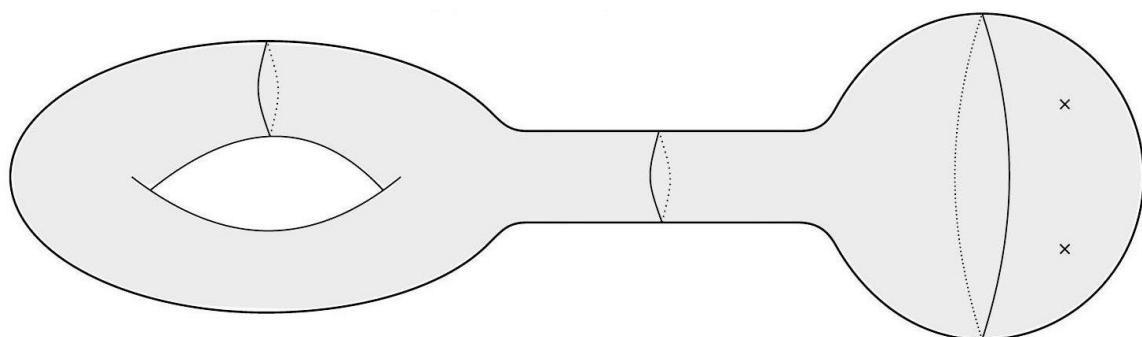
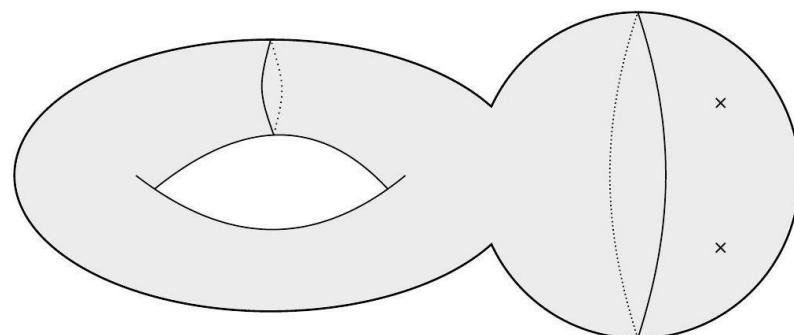
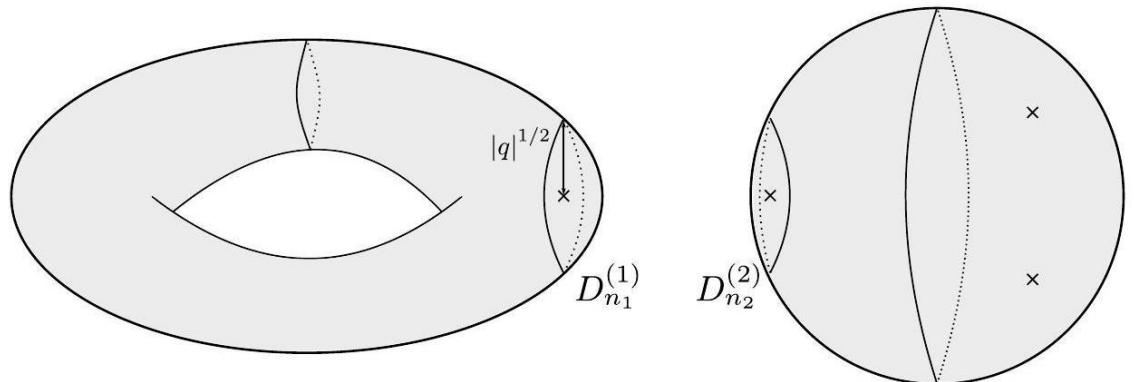
$$w_{n_1}^{(1)}w_{n_2}^{(2)}=q, |q|\leq 1,$$

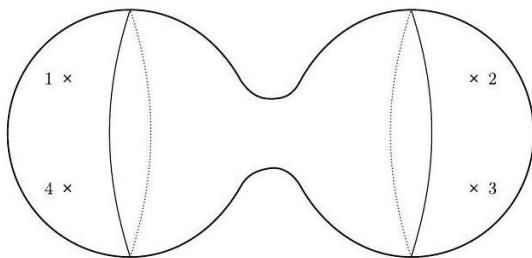
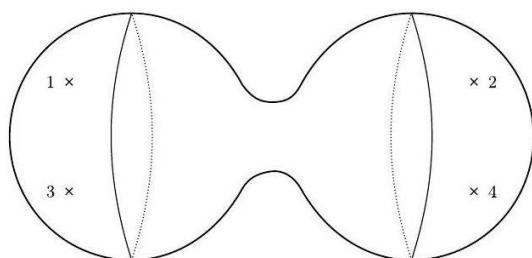
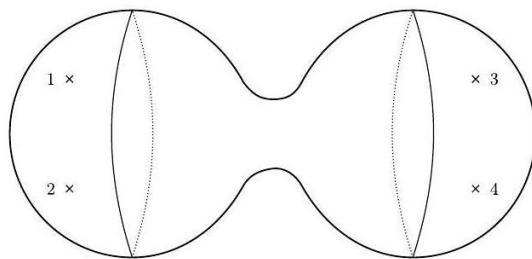
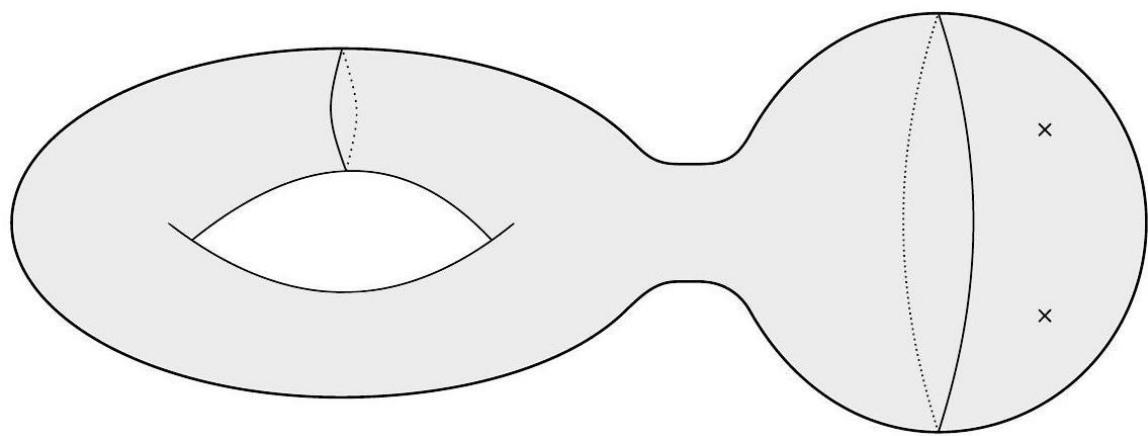
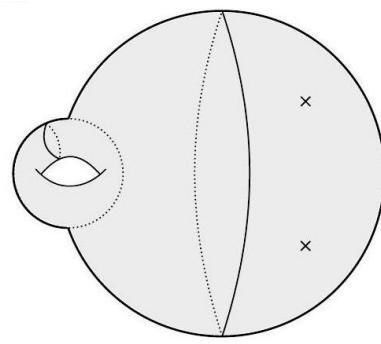
$$q=\mathrm{e}^{-s+\mathrm{i}\theta}, s\in\mathbb{R}_+, \theta\in[0,2\pi).$$

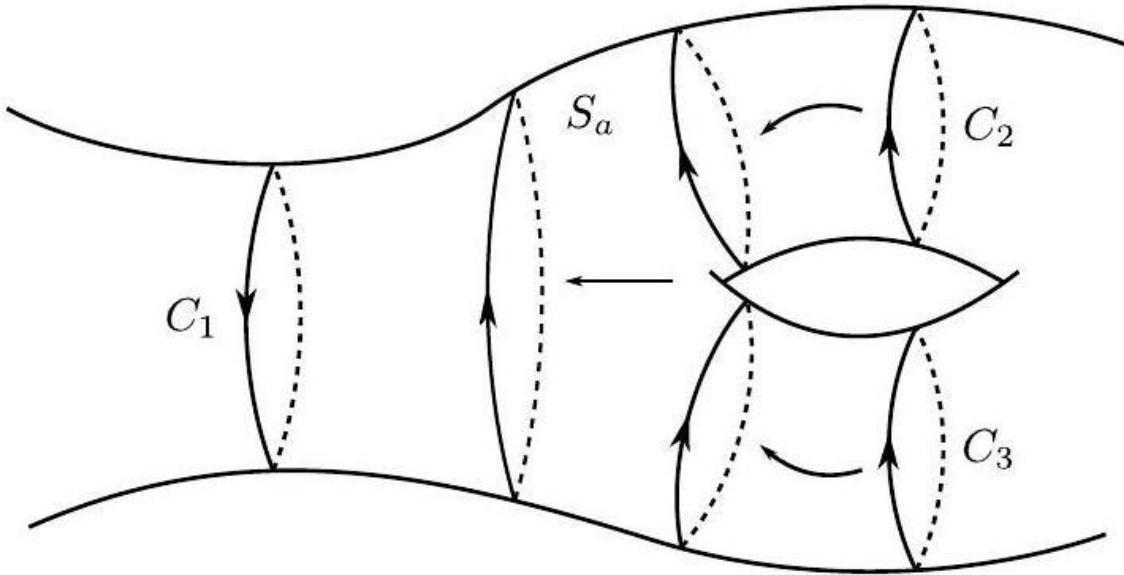


$$\left| w_{n_2}^{(2)} \right| = \frac{|q|}{\left| w_{n_1}^{(1)} \right|} > |q|^{1/2}$$

$$z_a^{(1)} = f_{an_1}^{(1)}(w_{n_1}^{(1)}), z_b^{(2)} = f_{bn_2}^{(2)}(w_{n_2}^{(2)})$$







Figuras 11, 12, 13, 14, 15, 16, 17, 18, 19 y 20. Puente de Einstein – Rosen a escala cuántica por curvatura local en dimensión \mathbb{R}^4 , es decir, en la misma dimensión.

$$z_a^{(1)} = f_{an_1}^{(1)}(w_{n_1}^{(1)}) = f_{an_1}^{(1)}\left(\frac{q}{w_{n_2}^{(2)}}\right) = f_{an_1}^{(1)}\left(\frac{q}{f_{bn_2}^{(2)-1}(z_b^{(2)})}\right)$$

$$z_a = F_{ab}(z_b), F_{ab} = f_{an_1}^{(1)} \circ (q \cdot I) \circ f_{bn_2}^{(2)-1},$$

$$\mathcal{M}_{g_1, n_1} + \mathcal{M}_{g_2, n_2} = 6g_1 - 6 + 2n_1 + 6g_2 - 6 + 2n_2 = \mathcal{M}_{g, n} - 2.$$

$$\Sigma_{g,n} = \Sigma_{g_1, n_1}, \begin{cases} g = g_1 + 1 \\ n = n_1 - 2 \end{cases}$$

$$w_{n_1-1}^{(1)} w_{n_1}^{(1)} = q$$

$$\mathcal{M}_{g_1, n_1} = \mathcal{M}_{g, n} - 2.$$

$$\mathcal{M}_{g_1, n_1} \mathcal{M}_{g_2, n_2} \subset \mathcal{M}_{g, n}$$

$$\mathcal{M}_{g_1, n_1} \subset \mathcal{M}_{g, n}$$

$$\mathcal{F}_{g,n} := \mathcal{M}_{g-1, n+2} \bigcup \left(\bigcup_{\substack{n_1+n_2=n+2 \\ g_1+g_2=g}} \mathcal{M}_{g_1, n_1} \# \mathcal{M}_{g_2, n_2} \right)$$

$$\mathcal{V}_{g,n} := \mathcal{M}_{g,n} - \mathcal{F}_{g,n}$$

$$\mathcal{F}_{g,n}^{1PR} := \bigcup_{\substack{n_1+n_2=n+2 \\ g_1+g_2=g}} \mathcal{M}_{g_1, n_1} \mathcal{M}_{g_2, n_2}$$

$$\mathcal{V}_{g,n}^{1PI} := \mathcal{M}_{g,n} - \mathcal{F}_{g,n}^{1PR}$$

$$\mathcal{V}^{1\text{PI}}_{g,n}=\mathcal{V}_{g,n}\;U\left(\bigcup_{g'}\;\#\mathcal{M}_{g-g',n+g'}\right)$$

$$\mathcal{V}_{0,3} = \mathcal{M}_{0,3}, \mathcal{F}_{0,3} = \emptyset.$$

$$\mathcal{F}_{0,4}=\mathcal{V}_{0,3}\mathcal{V}_{0,3}$$

$$\begin{array}{ll} \mathcal{F}_{0,5}=\mathcal{M}_{0,4}\mathcal{M}_{0,3}\\ \mathcal{F}_{1,1}=\mathcal{V}_{0,3} \end{array}\qquad\qquad\qquad \begin{array}{l} \mathcal{V}_{0,3}=\mathcal{V}_{0,3}\mathcal{V}_{0,3}\mathcal{V}_{0,3}+\mathcal{V}_{0,4}\mathcal{V}_{0,3}\\ \mathcal{F}_{0,4}=\mathcal{V}_{0,3}+\mathcal{V}_{0,4} \end{array}$$

$$r\big(\Sigma_{g_1,n_1}\#\Sigma_{0,3}\big)=r\big(\Sigma_{g_1,n_1}\big)+1, r\big(\#\Sigma_{g_1,n_1}\big)=r\big(\Sigma_{g_1,n_1}\big)+1.$$

$$r\big(\Sigma_{g,n}\big)=3g+n-2\in\mathbb{N}^*$$

$$r\big(\Sigma_{0,3}\big)=1$$

$$r\big(\Sigma_{g_1,n_1}\#\Sigma_{g_2,n_2}\big)=r\big(\Sigma_{g_1,n_1}\big)+r\big(\Sigma_{g_2,n_2}\big).$$

$$\mathbf{Stubs}.$$

$$q=\mathrm{e}^{-s+\mathrm{i}\theta}, s\in [s_0,\infty), \theta\in[0,2\pi), s_0\geq 0.$$

$$s_0 < s'_0 \colon \mathcal{F}_{g,n}(s'_0) \subset \mathcal{F}_{g,n}(s_0) \; \mathcal{V}_{g,n}(s_0) \subset \mathcal{V}_{g,n}(s'_0)$$

$$w_1=\lambda \tilde w_1, w_2=\lambda \tilde w_2$$

$$\tilde w_1\tilde w_2=\mathrm{e}^{-\tilde s+\mathrm{i}\tilde\theta}$$

$$\tilde s=s+2\ln|\lambda|, \tilde\theta=\theta+\mathrm{i}\ln\frac{\lambda}{\bar\lambda}$$

$$\tilde s\in[s_0,\infty), s_0:=2\ln|\lambda|.$$

$$\mathbf{Amplitudes\;y\;estados\;de\;superficie\;Off-shell}.$$

$$\omega_{i_1\cdots i_p}:=\omega_p\left(\partial_{s_1},\ldots,\partial_{s_p}\right),$$

$$\omega_{i_1i_2\cdots i_p}=-\omega_{i_2i_1\cdots i_p},$$

$$\omega_p\big(V^{(1)},\dots,V^{(p)}\big)=\omega_p\left(V^{(1)}_{s_1}\partial_{s_1},\dots,V^{(p)}_{s_p}\partial_{s_p}\right)=\omega_{i_1\cdots i_p}V^{(1)}_{s_1}\cdots V^{(p)}_{s_p},$$

$$\omega_p(\mathcal{V}_1,\dots,\mathcal{V}_n)\colon=\omega_p(\otimes_i\mathcal{V}_i).$$

$$\omega_0=(2\pi\mathrm{i})^{-\mathsf{M}^c_{g,n}}\left\langle\prod_{i=1}^n\;f_i\circ\mathcal{V}_i(0)\right\rangle_{\Sigma_{g,n}}$$

$$B(V)\colon=\oint\limits_c\frac{\mathrm{d} z}{2\pi\mathrm{i}}b(z)\nu(z)+\oint\limits_c\frac{\mathrm{d}\bar{z}}{2\pi\mathrm{i}}\bar{b}(\bar{z})\bar{\nu}(\bar{z})$$



$$B(V)\!:=\!\sum_{\alpha}\oint_{c_\alpha}\frac{{\rm d} z}{2\pi{\rm i}}b(z)v(z)+\,{\rm c.c.}$$

$$T(V)\!:=\!\oint_c\frac{{\rm d} z}{2\pi{\rm i}}T(z)v(z)+\oint_c\frac{{\rm d} \bar{z}}{2\pi{\rm i}}\bar{T}(\bar{z})\bar{v}(\bar{z})$$

$$T(V)=\{Q_B,B(V)\}$$

$$\begin{aligned} B &= B_s \; {\rm d} x_s, B_s \!:= B(\partial_s) \\ B_s &= \sum_{\alpha} \oint_{c_\alpha} \frac{{\rm d} \sigma_\alpha}{2\pi{\rm i}} b(\sigma_\alpha) \frac{\partial F_\alpha}{\partial x_s}(F_\alpha^{-1}(\sigma_\alpha)) + \sum_{\alpha} \oint_{c_\alpha} \frac{{\rm d} \bar{\sigma}_\alpha}{2\pi{\rm i}} \bar{b}(\bar{\sigma}_\alpha) \frac{\partial \bar{F}_\alpha}{\partial x_s}(\bar{F}_\alpha^{-1}(\bar{\sigma}_\alpha)) \end{aligned}$$

$$\omega_p\big(V^{(1)},\ldots,V^{(p)}\big)(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:= (2\pi{\rm i})^{-{\mathsf M}_{g,n}^c}\left\langle B\big(V^{(1)}\big)\cdots B\big(V^{(p)}\big)\prod_{i=1}^n\mathcal{V}_i\right\rangle_{\Sigma_{g,n}}$$

$$\begin{aligned} \omega_p &= \omega_{p,s_1\cdots s_p}\,{\rm d} x_{s_1}\wedge\cdots\wedge\,{\rm d} x_{s_p} \\ &= (2\pi{\rm i})^{-{\mathsf M}_{g,n}^c}\left\langle B_{s_1}\,{\rm d} x_{s_1}\wedge\cdots\wedge B_{s_p}\,{\rm d} x_{s_p}\prod_{i=1}^n\mathcal{V}_i\right\rangle_{\Sigma_{g,n}} \end{aligned}$$

$$x_s=x_s(t_1,\dots,t_q)$$

$$\begin{aligned} \forall p\leq q\colon \omega_p\big|_{\mathcal{S}} &= (2\pi{\rm i})^{-{\mathsf M}_{g,n}^c}\biggl\langle B_{r_1}\frac{\partial x_{s_1}}{\partial t_{r_1}}{\rm d} t_{r_1}\wedge\cdots\wedge B_{r_p}\frac{\partial x_{s_p}}{\partial t_{r_p}}{\rm d} t_{r_p}\prod_{i=1}^n\mathcal{V}_i\biggr\rangle_{\Sigma_{g,n}}, \\ \forall p>q\colon \omega_p\mid \mathcal{S} &= 0. \end{aligned}$$

$$B_r\!:=\frac{\partial x_s}{\partial t_r}B_s.$$

$$\begin{aligned} A_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)_{\mathcal{S}_{g,n}} &:= \int_{\mathcal{S}_{g,n}}\omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\Bigg|_{\mathcal{S}_{g,n}} \\ \omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\Big|_{\mathcal{S}_{g,n}} &= (2\pi{\rm i})^{-{\mathsf M}_{g,n}^c}\left\langle\bigwedge_{\lambda=1}^{{\mathsf M}_{g,n}}B_s\frac{\partial x_s}{\partial t_\lambda}{\rm d} t_\lambda\prod_{i=1}^nf_i\circ\mathcal{V}_i(0)\right\rangle \end{aligned}$$

$$A_n(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\sum_{g\geq 0}A_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n).$$

$$\mathcal{R}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\int_{\mathcal{R}_{g,n}}\omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n),$$

$$\mathcal{R}_n\!:=\sum_{g\geq 0}\mathcal{R}_{g,n},$$

$$\mathcal{R}_n(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\sum_{g\geq 0}\mathcal{R}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)=\sum_{g\geq 0}\int_{\mathcal{R}_{g,n}}\omega_{\mathbb{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n).$$



$$\begin{aligned}\langle \Sigma_{g,n} | B_{s_1} \cdots B_{s_p} | \otimes_i \mathcal{V}_i \rangle &:= \omega_{s_1 \cdots s_p}(\mathcal{V}_1, \dots, \mathcal{V}_n), \\ \langle \omega_p^{g,n} | \otimes_i \mathcal{V}_i \rangle &:= \omega_p(\mathcal{V}_1, \dots, \mathcal{V}_n), \\ \langle A_{g,n} | \otimes_i \mathcal{V}_i \rangle &:= A_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n).\end{aligned}$$

$$\langle \mathcal{R}_{g,n} | \otimes_i \mathcal{V}_i \rangle := \mathcal{R}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n).$$

$$\langle \omega_p^{g,n} | = \langle \Sigma_{g,n} | B_{s_1} \, \mathrm{d}x^{s_1} \cdots B_{s_p} \, \mathrm{d}x^{s_p}, \langle A_{g,n} | = \int_{\mathcal{M}_{g,n}} \langle \omega_p^{g,n} |.$$

$$\begin{array}{lll} C_1: w_1=z_1-y_1, & C_3: w_3=z_2-y_3, & C_5: z_1=z_2 \\ C_2: w_2=z_1-y_2, & C_4: w_4=z_2-y_4. \end{array}$$

$$y_4\rightarrow y_4+\delta y_4,\bar{y}_4\rightarrow \bar{y}_4+\delta \bar{y}_4$$

$$z'_2=z_2+\delta y_4, \bar{z}'_2=\bar{z}_2+\delta \bar{y}_4.$$

$$\nu=1, \bar{\nu}=0,$$

$$\nu=0, \bar{\nu}=1.$$

$$B\big(\partial_{y_4}\big)=\oint_{c_4}\mathrm{d} z_2 b(z_2)(+1), B\big(\partial_{\bar{y}_4}\big)=\oint_{c_4}\mathrm{d}\bar{z}_2 \bar{b}(\bar{z}_2)(+1)$$

$$\begin{aligned}\omega_2\big(\partial_{y_4},\partial_{\bar{y}_4}\big) &= \frac{1}{2\pi\mathrm{i}}\Bigg\langle B\big(\partial_{y_4}\big)B\big(\partial_{\bar{y}_4}\big)\prod_{i=1}^4\mathcal{V}_i\Bigg\rangle_{\Sigma_{0,4}} \\ &= \frac{1}{2\pi\mathrm{i}}\Bigg\langle \oint_{c_4}\mathrm{d} z_2 b(z_2)\oint_{c_4}\mathrm{d}\bar{z}_2 \bar{b}(\bar{z}_2)\prod_{i=1}^4\mathcal{V}_i\Bigg\rangle_{\Sigma_{0,4}}\end{aligned}$$

$$\omega_2\big(\partial_{y_4},\partial_{\bar{y}_4}\big)=\frac{1}{2\pi\mathrm{i}}\leftg\langle \prod_{i=1}^3c\bar{c}V_i(y_i,\bar{y}_i)\oint_{c_4}\mathrm{d} z_2 b(z_2)\oint_{c_4}\mathrm{d}\bar{z}_2 \bar{b}(\bar{z}_2)\bar{c}(\bar{y}_4)c(y_4)V_4(y_4,\bar{y}_4)\rightg\rangle_{\Sigma_{0,4}}$$

$$\oint_{c_4}\mathrm{d} z_2 b(z_2)c(y_4)\sim \oint_{c_4}\mathrm{d} z_2\frac{1}{z_2-y_4}$$

$$A_{0,4}=\frac{1}{2\pi\mathrm{i}}\int\mathrm{d} y_4\wedge\mathrm{d}\bar{y}_4\leftg\langle \prod_{i=1}^3c\bar{c}V_i(y_i)V_4(y_4)\rightg\rangle_{\Sigma_{0,4}}$$

$$z_a\rightarrow z_a+\phi(z_a)$$

$$C_i\colon \nu^{(i)}=\phi|_{C_i}.$$

$$B(V)=\sum_{i=1}^3\oint_{c_i}\mathrm{d} z_ab(z_a)\phi(z_a)+\text{ c.c.}$$

$$w_i\rightarrow (1+\mathrm{i}\alpha_i)w_i, \bar{w}_i\rightarrow (1-\mathrm{i}\alpha_i)\bar{w}_i,$$



$$v=\mathrm{i} w_i,\bar{v}=-\mathrm{i}\bar{w}_i.$$

$$B(V)=\mathrm{i}\oint_{~c_i}\mathrm{d} w_iw_ib(w_i)-\mathrm{i}\oint_{~c_i}\mathrm{d}\bar{w}_i\bar{w}_i\bar{b}(\bar{w}_i),$$

$$B(V)\mathcal{V}_i(0)=\mathrm{i}\oint_{~c_i}\mathrm{d} w_iw_ib(w_i)\mathcal{V}_i(0)-\mathrm{i}\oint_{~c_i}\mathrm{d}\bar{w}_i\bar{w}_i\bar{b}(\bar{w}_i)\mathcal{V}_i(0),$$

$$\mathrm{i}\big(b_0-\bar{b}_0\big)|\mathcal{V}_i\rangle,$$

$$b^-_0|\mathcal{V}_i\rangle=0$$

$$\mathcal{V}_i \in \mathcal{H}^{-},$$

$$w_i\longrightarrow f(w_i),f(0)=0$$

$$f(w_i) = \sum_{m \geq 0} \; p_m w_i^{m+1}$$

$$\nu_m=w_i^{m+1},\bar{\nu}_m=0$$

$$B\big(\partial_{p_m}\big)=\oint_{~c_i}\mathrm{d} w_ib(w_i)w_i^{m+1}$$

$$\forall m\geq 0\colon b_m|\mathcal{V}_i\rangle=0,\bar{b}_m|\mathcal{V}_i\rangle=0$$

$$\omega_p\left(\sum_i\;Q_B^{(i)}\otimes_i\mathcal{V}_i\right)=(-1)^p\;\mathrm{d}\omega_{p-1}(\otimes\mathcal{V}_i)$$

$$Q_B^{(i)}=1_{i-1}\otimes Q_B\otimes 1_{n-i}$$

$$Q_B\mathcal{V}_i(z,\bar{z})=\frac{1}{2\pi\mathrm{i}}\oint\mathrm{d} w j_B(w)\mathcal{V}_i(z,\bar{z})+\text{ c.c.}$$

$$\begin{aligned}\omega_p\left(\sum_i\;Q_B^{(i)}\otimes_i\mathcal{V}_i\right)&=\omega_p(Q_B\mathcal{V}_1,\mathcal{V}_2,\dots,\mathcal{V}_n)+(-1)^{|\mathcal{V}_1|}\omega_p(\mathcal{V}_1,Q_B\mathcal{V}_2,\dots,\mathcal{V}_n)\\&+\cdots+(-1)^{|\mathcal{V}_1|+\cdots+|\mathcal{V}_{n-1}|}\omega_p(\mathcal{V}_1,\mathcal{V}_2,\dots,Q_B\mathcal{V}_n)\end{aligned}$$

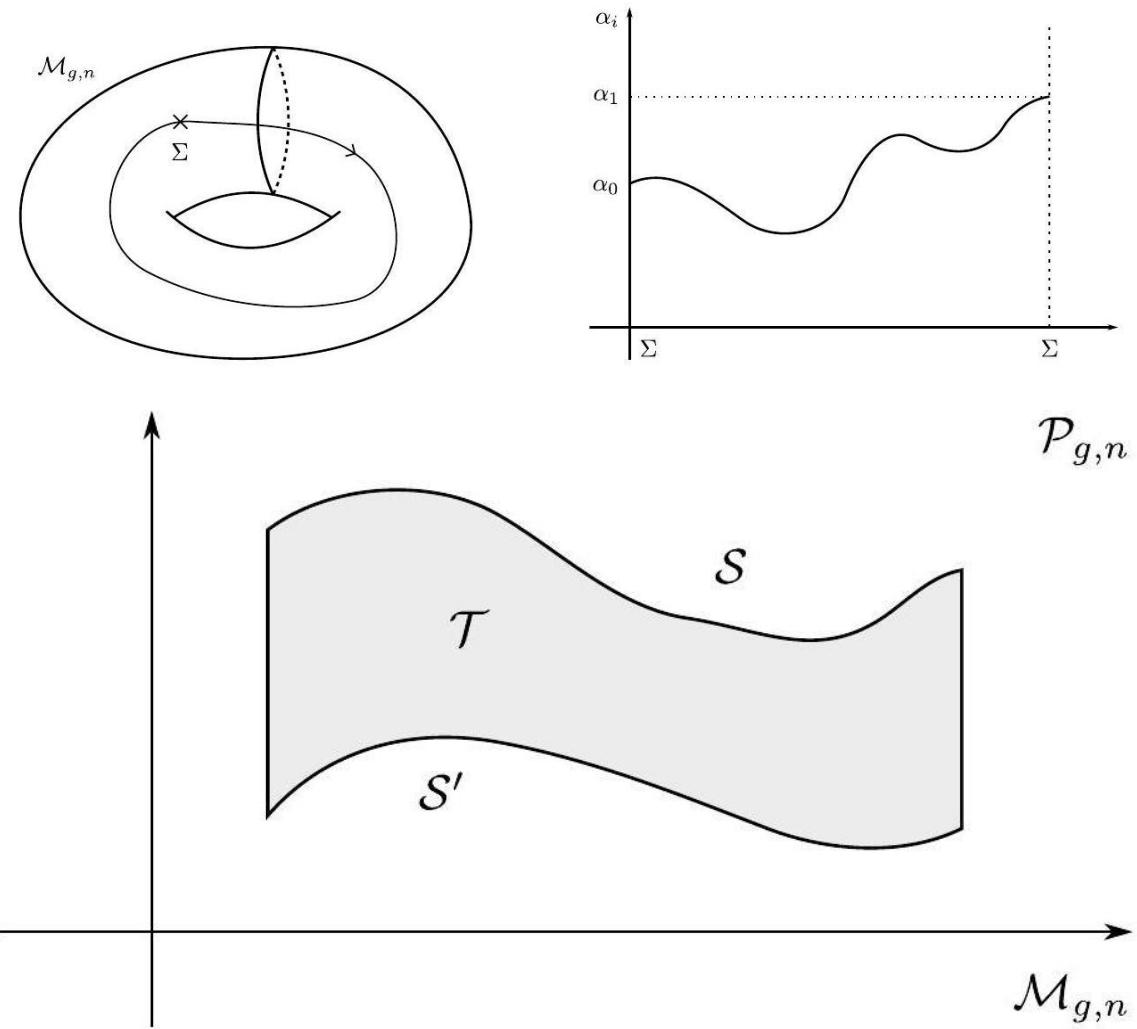
$$T_s=\{Q_B,B_s\}$$

$$\mathrm{d}x_s\{Q_B,B_s\}=\mathrm{d}x_sT_s=\mathrm{d}x_s\partial_s=\mathrm{d}$$

$$N_{\mathrm{gh}}\left(\omega_p(\mathcal{V}_1,\dots,\mathcal{V}_n)\right)=\sum_{i=1}^n\;N_{\mathrm{gh}}(\mathcal{V}_i)-p=6-6g$$

$$N_{\mathrm{gh}}\left(\omega_{\mathsf{M}_{g,n}}\right)=6-6g\;\Rightarrow\;\sum_{i=1}^n\;N_{\mathrm{gh}}(\mathcal{V}_i)=2n$$





Figuras 21 y 22. Deformación local del espacio – tiempo cuántico y fluctuaciones de materia y energía de la partícula supermasiva.

$$L_0^- |\mathcal{V}_i\rangle = 0$$

$$b_0^- |\mathcal{V}_i\rangle = 0$$

$$\mathcal{V}_i \in \mathcal{H}^- \cap \ker L_0^-,$$

$$\mathcal{V}_i(0) \rightarrow (e^{i\alpha_i})^h (e^{-i\alpha_i})^{\bar{h}} \mathcal{V}_i(0)$$

$$|\mathcal{V}_i\rangle \rightarrow e^{i\alpha_i(L_0 - \bar{L}_0)} |\mathcal{V}_i\rangle$$

$$\forall i: Q_B |\mathcal{V}_i\rangle = 0 \Rightarrow d\omega_{p-1}(\mathcal{V}_1, \dots, \mathcal{V}_n) = 0.$$

$$\int_S \omega_{M_{g,n}} - \int_{S'} \omega_{M_{g,n}} = \int_{\partial T} \omega_{M_{g,n}-1} = \int_T d\omega_{M_{g,n}-1} = 0$$

$$|\mathcal{V}_1\rangle = Q_B |\Lambda\rangle, Q_B |\mathcal{V}_i\rangle = 0$$



$$\omega_{M_{g,n}}(Q_B \Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = d\omega_{M_{g,n-1}}(\Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n),$$

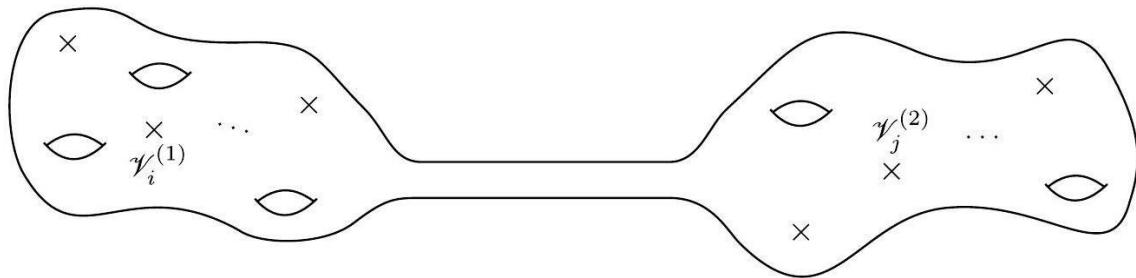
$$\int_{\mathcal{S}} \omega_{M_{g,n}}(Q_B \Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = \int_{\mathcal{S}} d\omega_{M_{g,n-1}}(\Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = \int_{\partial \mathcal{S}} \omega_{M_{g,n-1}}(\Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n)$$

$$\int_{\mathcal{S}} \omega_{M_{g,n}}(Q_B \Lambda, \mathcal{V}_2, \dots, \mathcal{V}_n) = 0$$

Factorización de amplitudes y diagramas de Feynman.

$$w_{n_1}^{(1)} w_{n_2}^{(2)} = q.$$

$$A_{g,n} = \int_{\Sigma_{g,n}} \omega_{M_{g,n}}^{g,n} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1-1}^{(1)}, \mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-1}^{(2)} \right).$$



$$w'_{n_1}^{(1)} = w_{n_1}^{(1)} + \frac{w_{n_1}^{(1)}}{q} \delta q$$

$$v_q = \frac{w_{n_1}^{(1)}}{q}, B_q = \frac{1}{q} \oint c_q dw_{n_1}^{(1)} b(w_{n_1}^{(1)}) w_{n_1}^{(1)}$$

$$w'_{n_1}^{(1)} w_{n_2}^{(2)} = q + \delta q$$

$$w'_{n_1}^{(1)} = \frac{w_{n_1}^{(1)}}{q} (q + \delta q) = w_{n_1}^{(1)} + \frac{q}{w_{n_1}^{(1)}} \delta q$$

$$\omega_{M_{g,n}} \left(V_1^{(1)}, \dots, V_{M_{g_1,n_1}^{(1)}}^{(1)}, \partial_q, \partial_{\bar{q}}, V_1^{(2)}, \dots, V_{M_{g_2,n_2}^{(2)}}^{(2)} \right)$$

$$= (2\pi i)^{-M_{g,n}^c} \left\langle \prod_{\lambda=1}^{M_{g_1,n_1}} B(V_\lambda^{(1)}) B(\partial_q) B(\partial_{\bar{q}}) \prod_{\kappa=1}^{M_{g_2,n_2}} B(V_\kappa^{(2)}) \prod_{i=1}^{n_1-1} \mathcal{V}_i^{(1)} \prod_{j=1}^{n_2-1} \mathcal{V}_j^{(2)} \right\rangle_{\Sigma_{g,n}}$$

$$\left\langle \Sigma_{n_1} \mid \mathcal{V}_{n_1}^{(1)} \right\rangle := \omega_{M_{g_1,n_1}} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1}^{(1)} \right) = (2\pi i)^{-M_{g_1,n_1}^c} \left\langle \prod_{\lambda=1}^{M_{g_1,n_1}} B(V_\lambda^{(1)}) \prod_{i=1}^{n_1-1} \mathcal{V}_i^{(1)} \right\rangle_{\Sigma_{g_1,n_1}}$$

$$\left\langle \Sigma_{n_2} \mid \mathcal{V}_{n_2}^{(2)} \right\rangle := \omega_{M_{g_2,n_2}} \left(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2}^{(2)} \right) = (2\pi i)^{-M_{g_2,n_2}^c} \left\langle \prod_{\lambda=1}^{M_{g_2,n_2}} B(V_\lambda^{(2)}) \prod_{j=1}^{n_2-1} \mathcal{V}_j^{(2)} \right\rangle_{\Sigma_{g_2,n_2}}$$



$$\left\langle \Sigma_{n_1}\right|=\left\langle 0|I\circ\Sigma_{n_1}(0),\left\langle \Sigma_{n_2}\right|= \left\langle 0|I\circ\Sigma_{n_2}(0),$$

$$\left\langle \Sigma_{n_1}\mid \mathcal{V}_{n_1}^{(1)}\right\rangle =\left\langle I\circ\Sigma_{n_1}(0)\mathcal{V}_{n_1}(0)\right\rangle _{w_{n_1}^{(1)}},\left\langle \Sigma_{n_2}\mid \mathcal{V}_{n_2}^{(2)}\right\rangle =\left\langle I\circ\Sigma_{n_2}(0)\mathcal{V}_{n_2}(0)\right\rangle _{w_{n_2}^{(2)}}$$

$$w_1=I\left(w_{n_1}^{(1)}\right), w_2=I\left(w_{n_2}^{(2)}\right)$$

$$w_{n_1}^{(1)}=\frac{q}{w_{n_2}^{(2)}}=\frac{q}{I(w_2)}=qw_2:=f(w_2)$$

$$\omega_{\mathsf{M}_{g,n}}=\frac{1}{2\pi\mathrm{i}}\big\langle I\circ\Sigma_{n_1}(0)B_qB_{\bar q}f\circ\Sigma_{n_2}(0)\big\rangle_{w_{n_1}^{(1)}}=\frac{1}{2\pi\mathrm{i}}\langle\Sigma_{n_1}\big|B_qB_{\bar q}q^{L_0}\bar q^{\bar L_0}\big|\Sigma_2\rangle,$$

$$\langle\Sigma_{n_1}\big|B_qB_{\bar q}\,\Big|\mathcal{V}_{n_1}^{(1)}\Big\rangle=\frac{1}{q\bar q}\langle\Sigma_{n_1}\big|b_0\bar b_0\,\Big|\mathcal{V}_{n_1}^{(1)}\Big\rangle$$

$$B_q\mathcal{V}_{n_1}^{(1)}(z,\bar z)=\frac{1}{q}\oint_{c_q}\,\mathrm{d} w_{n_1}^{(1)}b\left(w_{n_1}^{(1)}\right)w_{n_1}^{(1)}\mathcal{V}_{n_1}^{(1)}(z,\bar z)\rightarrow\frac{1}{q}b_0\left|\mathcal{V}_{n_1}^{(1)}\right\rangle$$

$$\omega_{\mathsf{M}_{g,n}}=\frac{1}{2\pi\mathrm{i}}\frac{1}{q\bar q}\langle\Sigma_{n_1}\big|b_0\bar b_0q^{L_0}\bar q^{\bar L_0}\big|\Sigma_{n_2}\rangle.$$

$$\mathcal{F}_{g,n}\left(\mathcal{V}_i^{(1)}\mid \mathcal{V}_j^{(2)}\right):=\frac{1}{2\pi\mathrm{i}}\int\bigwedge_{\lambda=1}^{\mathsf{M}_{g_1,n_1}}\mathrm{d} t_\lambda^{(1)}\bigwedge_{\kappa=1}^{\mathsf{M}_{g_2,n_2}}\mathrm{d} t_\kappa^{(2)}\wedge\frac{\mathrm{d} q}{q}\wedge\frac{\mathrm{d} \bar q}{\bar q}\langle\Sigma_{n_1}\big|b_0\bar b_0q^{L_0}\bar q^{\bar L_0}\big|\Sigma_{n_2}\rangle.$$

$$\mathcal{F}_{g,n}\left(\mathcal{V}_i^{(1)}\mid \mathcal{V}_j^{(2)}\right)$$

$$=\frac{1}{2\pi\mathrm{i}}\int\,\frac{\mathrm{d}^D k}{(2\pi)^D}\frac{\mathrm{d}^D k'}{(2\pi)^D}(-1)^{|\phi_\alpha|}\times\int\,\frac{\mathrm{d} q}{q}$$

$$\wedge\frac{\mathrm{d} \bar q}{\bar q}\langle\phi_\alpha(k)^c\big|b_0\bar b_0q^{L_0}\bar q^{\bar L_0}\big|\phi_\beta(k')^c\big\rangle$$

$$\times\int\bigwedge_{\lambda=1}^{\mathsf{M}_{g_1,n_1}}\mathrm{d} t_\lambda^{(1)}\langle\Sigma_{n_1}\mid\phi_\alpha(k)\rangle\bigwedge_{\kappa=1}^{\mathsf{M}_{g_2,n_2}}\mathrm{d} t_\kappa^{(2)}\langle\phi_\beta(k')\mid\Sigma_{n_2}\rangle$$

$$A_{g_1,n_1}\left(\mathcal{V}_1^{(1)},\ldots,\mathcal{V}_{n_1-1}^{(1)},\phi_\alpha(k)\right)=\int_{\mathcal{S}_{g_1,n_1}}\omega_{\mathsf{M}_{g_1,n_1}}\left(\mathcal{V}_1^{(1)},\ldots,\mathcal{V}_{n_1-1}^{(1)},\phi_\alpha(k)\right)$$

$$=\int_{\mathcal{S}_{g_1,n_1}}\bigwedge_{\lambda=1}^{\mathsf{M}_{g_1,n_1}}\mathrm{d} t_\lambda^{(1)}\langle\Sigma_{n_1}\mid\phi_\alpha(k)\rangle$$



$$\begin{aligned} A_{g_2, n_2} \left(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-1}^{(2)}, \phi_\beta(k') \right) &= \int_{\mathcal{S}_{g_2, n_2}} \omega_{M_{g_2, n_2}} \left(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-2}^{(2)}, \phi_\beta(k') \right) \\ &= \int_{\mathcal{S}_{g_2, n_2}} \bigwedge_{\lambda=1}^{M_{g_2, n_2}} dt_\lambda^{(2)} \langle \Sigma_{n_2} \mid \phi_\beta(k') \rangle \end{aligned}$$

$$\Delta_{\alpha\beta}(k, k') := \Delta(\phi_\alpha(k)^c, \phi_\beta(k')^c) := \frac{1}{2\pi i} \int \frac{dq}{q} \wedge \frac{d\bar{q}}{\bar{q}} \langle \phi_\alpha(k)^c | b_0 \bar{b}_0 q^{L_0} \bar{q}^{\bar{L}_0} | \phi_\beta(k')^c \rangle$$

$$\begin{aligned} \mathcal{F}_{g,n} \left(\mathcal{V}_i^{(1)} \mid \mathcal{V}_j^{(2)} \right) &= \int \frac{d^D k}{(2\pi)^D} \frac{d^D k'}{(2\pi)^D} A_{g_1, n_1} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1-1}^{(1)}, \phi_\alpha(k) \right) \Delta_{\alpha\beta}(k, k') \\ &\quad \times A_{g_2, n_2} \left(\mathcal{V}_1^{(2)}, \dots, \mathcal{V}_{n_2-1}^{(2)}, \phi_\beta(k') \right) \end{aligned}$$

$$A_{g_1, n_1} \sim \delta^{(D)} \left(k_1^{(1)} + \dots + k_{n_1-1}^{(1)} + k \right), A_{g_2, n_2} \sim \delta^{(D)} \left(k_1^{(2)} + \dots + k_{n_2-1}^{(2)} + k' \right).$$

$$\mathcal{F}_{g,n} \sim \delta^{(D)} \left(k_1^{(1)} + \dots + k_{n_1-1}^{(1)} + k_1^{(2)} + \dots + k_{n_2-1}^{(2)} \right)$$

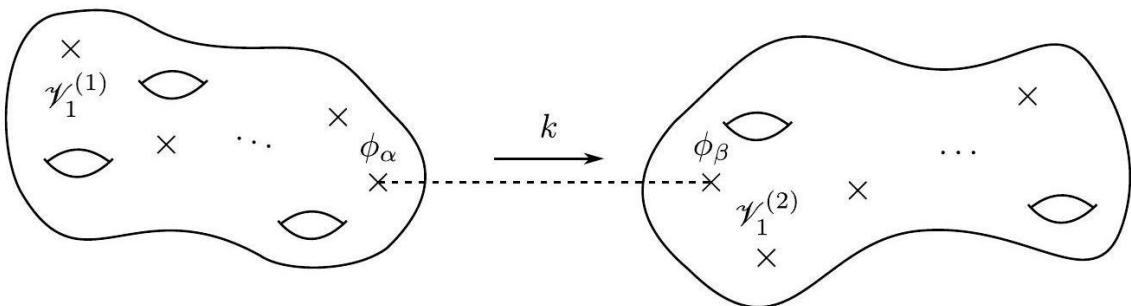
$$N_{\text{gh}}(\phi_\alpha) = 2n_1 - \sum_{i=1}^{n_1-1} N_{\text{gh}} \left(\mathcal{V}_i^{(1)} \right), N_{\text{gh}}(\phi_\beta) = 2n_2 - \sum_{j=1}^{n_2-1} N_{\text{gh}} \left(\mathcal{V}_j^{(2)} \right)$$

$$N_{\text{gh}}(\phi_\alpha) + N_{\text{gh}}(\phi_\beta) = 4$$

$$N_{\text{gh}}(\phi_\alpha) = N_{\text{gh}}(\phi_\beta) = 2$$

$$w_{n_1-1} w_{n_1} = q$$

$$A_{g,n} = \int_{\mathcal{S}_{g,n}} \omega_{M_{g,n}}^{g,n} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1-2}^{(1)} \right)$$



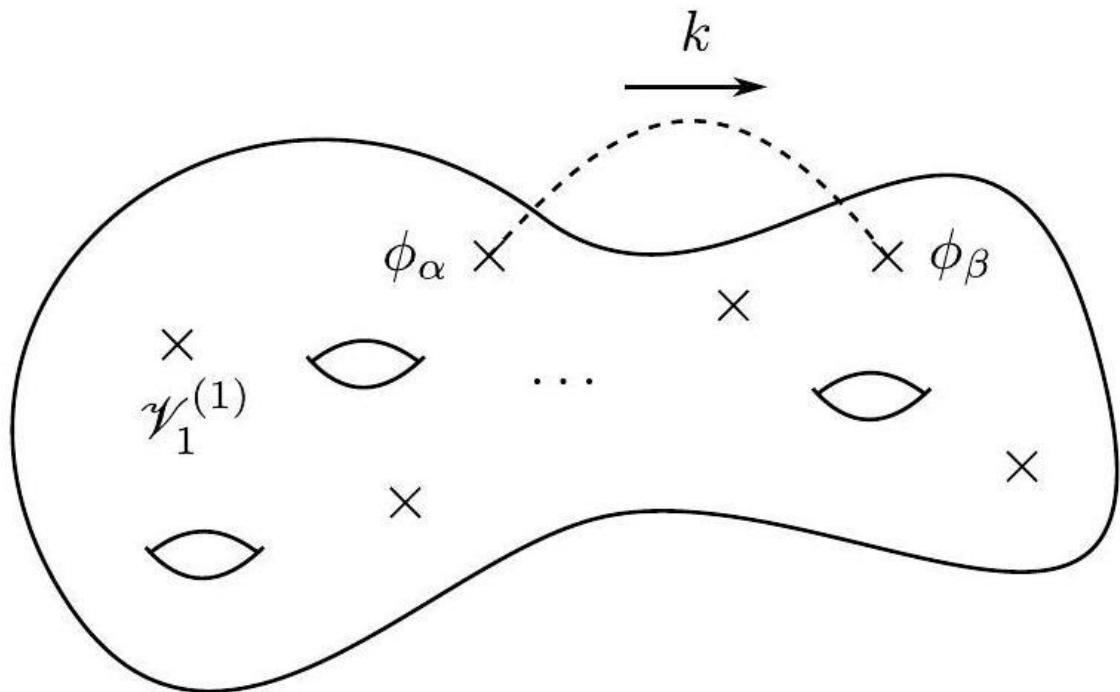
$$\omega_{M_{g,n}} \left(V_1^{(1)}, \dots, V_{M_{g_1, n_1}}^{(1)}, \partial_q, \partial_{\bar{q}} \right) = (2\pi i)^{-M_{g,n}^c} \left\langle \prod_{\lambda=1}^{M_{g_1, n_1}} B(V_\lambda^{(1)}) B(\partial_q) B(\partial_{\bar{q}}) \prod_{i=1}^{n_1-2} \mathcal{V}_i^{(1)} \right\rangle_{\Sigma_{g,n}}$$

$$\left\langle \Sigma_{n_1-1, n_1} \mid \mathcal{V}_{n_1-1}^{(1)} \otimes \mathcal{V}_{n_1}^{(1)} \right\rangle := \omega_{M_{g_1, n_1}} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1}^{(1)} \right)$$

$$\mathcal{F}_{g,n} \left(\mathcal{V}_i^{(1)} \mid \right) = \int \frac{d^D k}{(2\pi)^D} \frac{d^D k'}{(2\pi)^D} A_{g_1, n_1} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1-2}^{(1)}, \phi_\alpha(k), \phi_\beta(k') \right) \Delta_{\alpha\beta}(k, k')$$

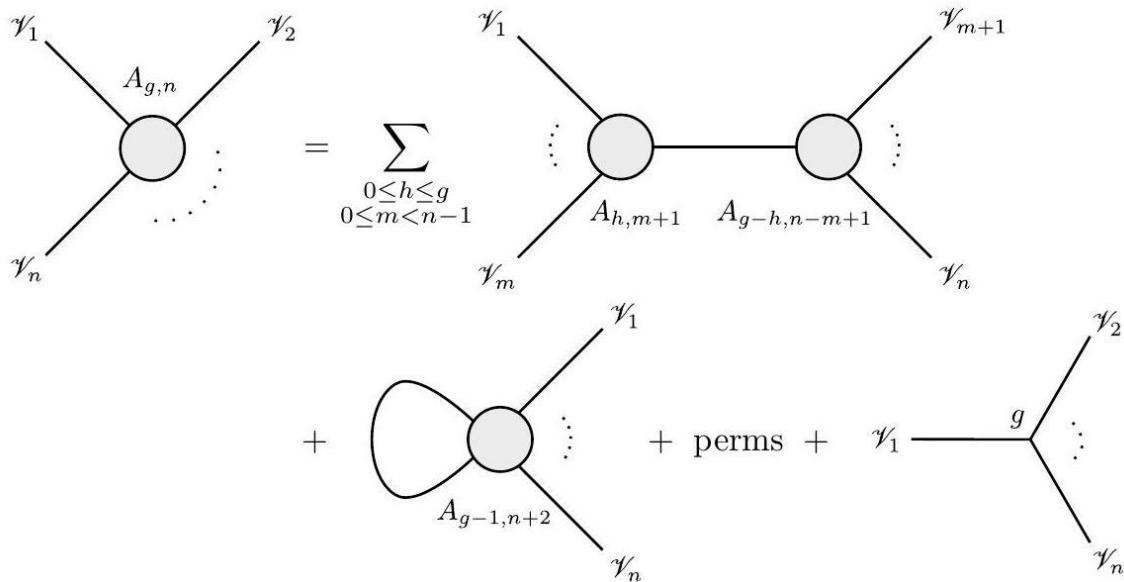
$$N_{\text{gh}}(\phi_\alpha) + N_{\text{gh}}(\phi_\beta) = 2n_1 - \sum_{i=1}^{n_1-2} N_{\text{gh}} \left(\mathcal{V}_i^{(1)} \right) = 4$$

$$A_{g_1, n_1} \left(\mathcal{V}_1^{(1)}, \dots, \mathcal{V}_{n_1-2}^{(1)}, \phi_\alpha(k), \phi_\beta(-k) \right) \sim \delta^{(D)} \left(k_1^{(1)} + \dots + k_{n_1-2}^{(1)} \right)$$



$$r(A_{g,n}) := r(\Sigma_{g,n}) = 3g + n - 2$$

$$\mathcal{V}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) := \mathcal{V}_1 \xrightarrow{g} / \mathcal{V}_n := \int_{\mathcal{R}_{g,n}} \omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n),$$



$$\mathcal{V}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n) := \mathcal{V}_{g,n}(\otimes_i \mathcal{V}_i) := \{\mathcal{V}_1, \dots, \mathcal{V}_n\}_g.$$

Escalar QFT.

Propagador.

$$\Delta = \frac{1}{2\pi i} \int \frac{dq}{q} \wedge \frac{d\bar{q}}{\bar{q}} b_0 \bar{b}_0 q^{L_0} \bar{q}^{\bar{L}_0}$$

$$q = e^{-s+i\theta}, s \in \mathbb{R}_+, \theta \in [0, 2\pi)$$

$$\frac{dq}{q} \wedge \frac{d\bar{q}}{\bar{q}} = -2i ds \wedge d\theta$$

$$\Delta = \frac{1}{2\pi} b_0^+ b_0^- \int_0^\infty ds e^{-sL_0^+} \int_0^{2\pi} d\theta e^{i\theta L_0^-}$$

$$\int_0^\infty ds e^{-sL_0^+} = \frac{1}{L_0^+}, \int_0^{2\pi} d\theta e^{i\theta L_0^-} = 2\pi \delta_{L_0^-, 0}$$

$$\Delta = \frac{b_0^+}{L_0^+} b_0^- \delta_{L_0^-, 0}$$

$$\Delta = -2b_0 \bar{b}_0 \frac{1}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0}$$

$$L_0^+ |\phi_\alpha(k)\rangle = \frac{\alpha'}{2} (k^2 + m_\alpha^2) |\phi_\alpha(k)\rangle, L_0^- |\phi_\alpha(k)\rangle = 0$$



$$\begin{aligned}\Delta_{\alpha\beta}(k, k') &:= \langle \phi_\alpha(k)^c | \Delta | \phi_\beta(k')^c \rangle := (2\pi)^D \delta^{(D)}(k + k') \Delta_{\alpha\beta}(k), \\ \Delta_{\alpha\beta}(k) &:= \frac{M_{\alpha\beta}(k)}{k^2 + m_\alpha^2}, M_{\alpha\beta}(k) := \frac{2}{\alpha'} \langle \phi_\alpha^c(k) | b_0^+ b_0^- | \phi_\beta^c(-k) \rangle,\end{aligned}$$

$$b_0^+ |\phi_\alpha^c\rangle \neq 0, b_0^- |\phi_\alpha^c\rangle \neq 0$$

$$c_0^+ |\phi_\alpha^c\rangle = 0, c_0^- |\phi_\alpha^c\rangle = 0$$

$$|\phi_\alpha^c\rangle = |\phi_1\rangle + b_0^\pm |\phi_2\rangle, c_0^\pm |\phi_1\rangle = c_0^\pm |\phi_2\rangle = 0$$

$$c_0^\pm |\phi_\alpha^c\rangle = 0 \Rightarrow |\phi_2\rangle = 0$$

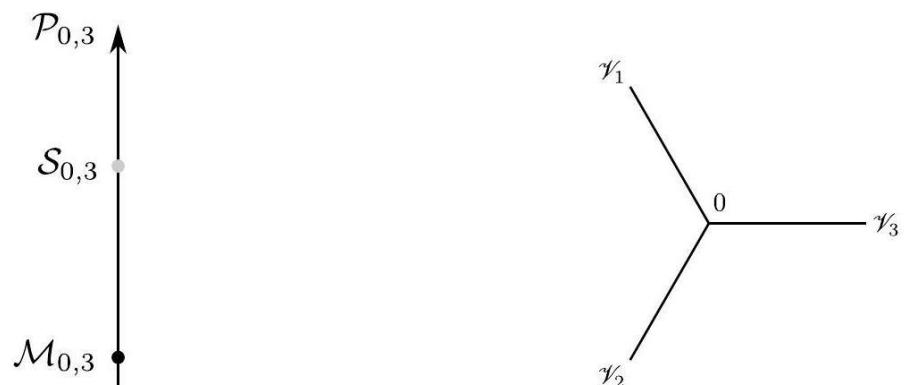
$$b_0^+ |\phi_\alpha\rangle = 0, b_0^- |\phi_\alpha\rangle = 0$$

$$b_0^+ |\mathcal{V}_i\rangle = 0$$

$$\Delta^{-1} = c_0^+ c_0^- L_0^+ \delta_{L_0^-, 0}$$

Vértices fundamentales.

$$\mathcal{V}_{0,3}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3) := A_{0,3}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3) = \omega_0^{0,3}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3).$$



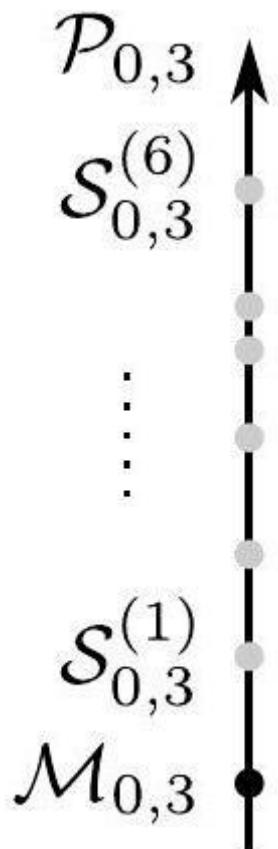
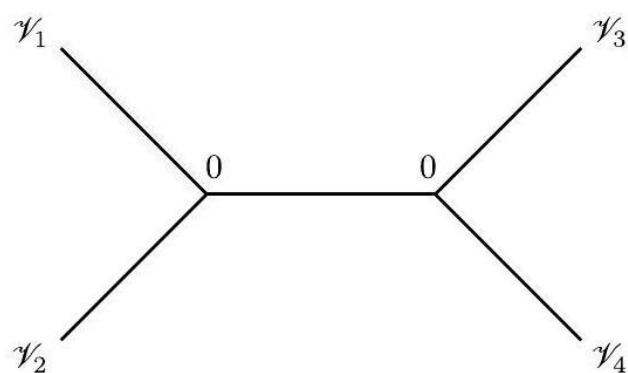


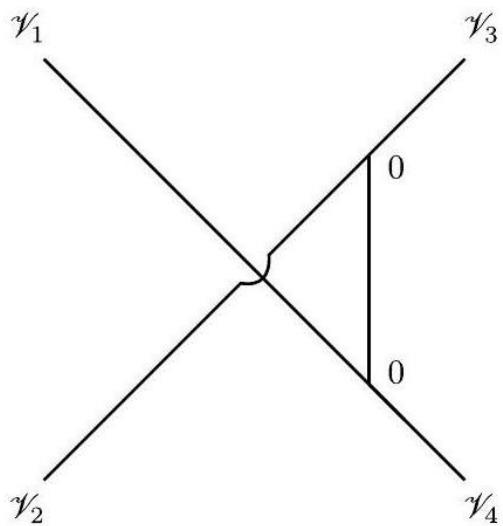
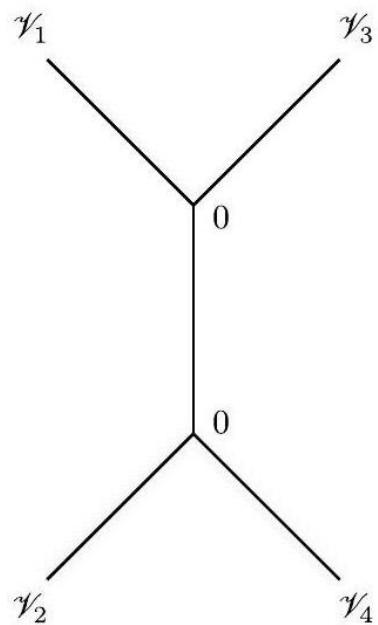
Figura 23, 23A y 23B. Operadores de vértice.

$$\mathcal{V}_{0,4}(\mathcal{V}_1, \mathcal{V}_2, \mathcal{V}_3, \mathcal{V}_4) := \int_{\mathcal{R}_{0,4}} \omega_2^{0,4}(\mathcal{V}_1, \dots, \mathcal{V}_4)$$

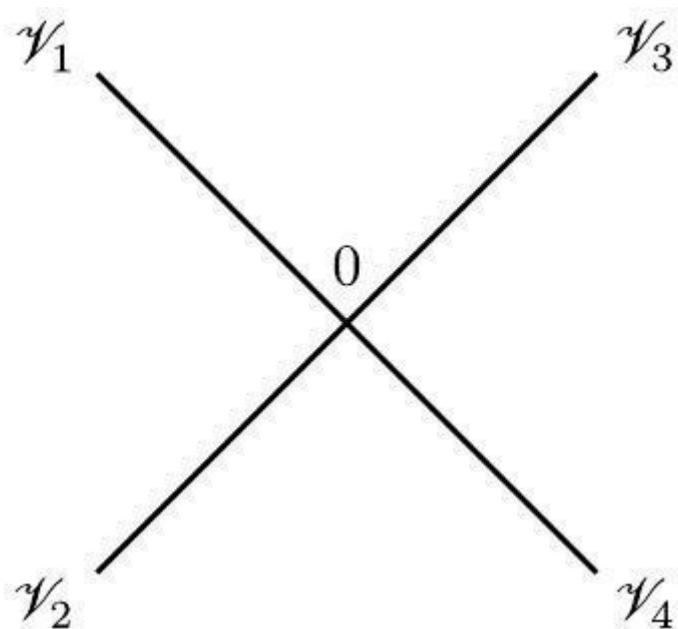
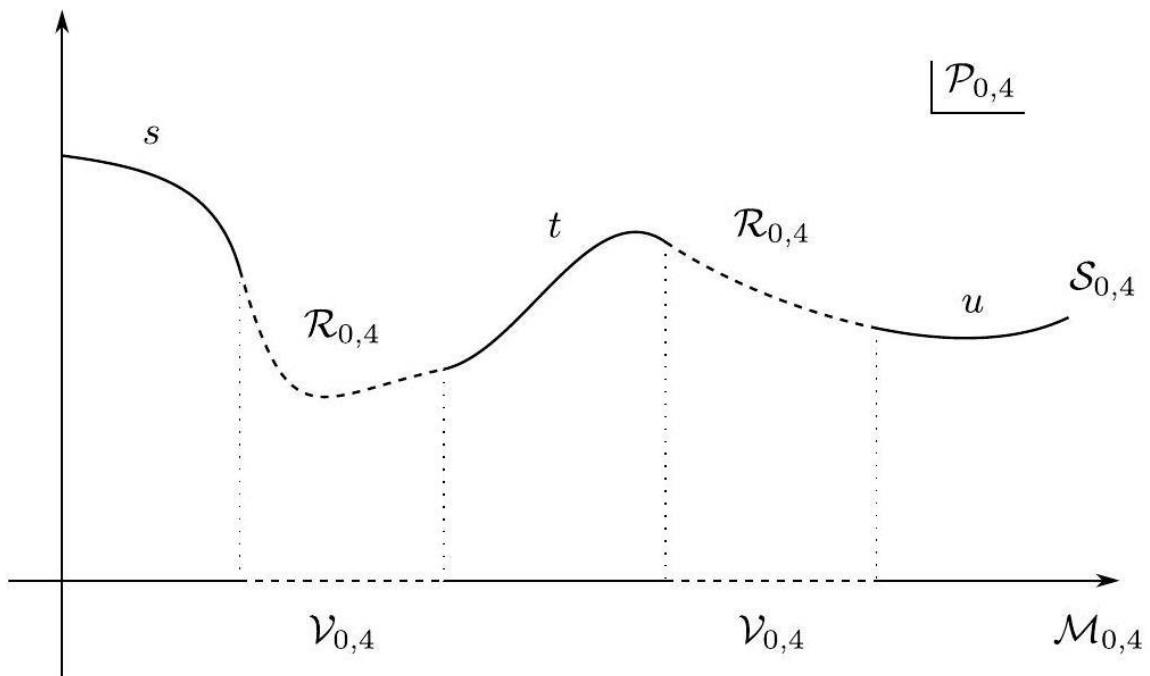
$$A_{0,4} = \mathcal{F}_{0,4}^{(s)} + \mathcal{F}_{0,4}^{(t)} + \mathcal{F}_{0,4}^{(u)} + \mathcal{V}_{0,4}$$

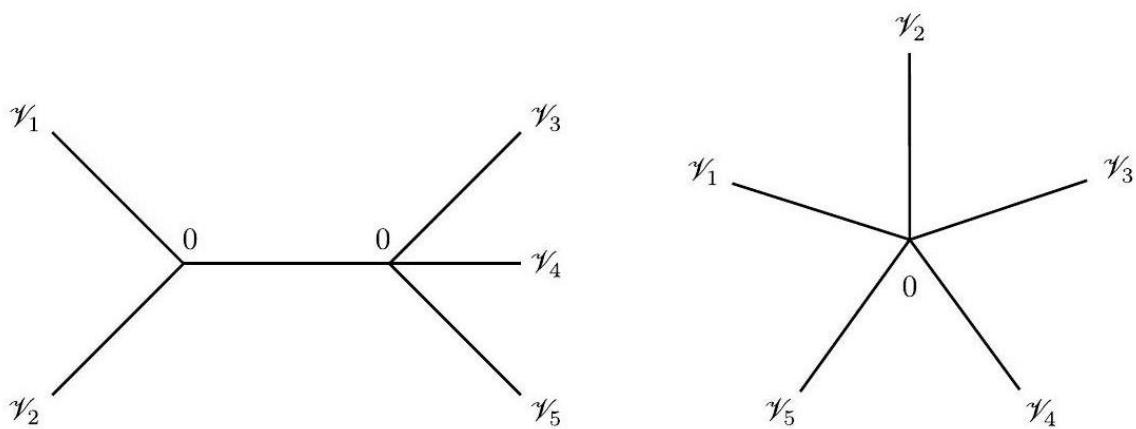
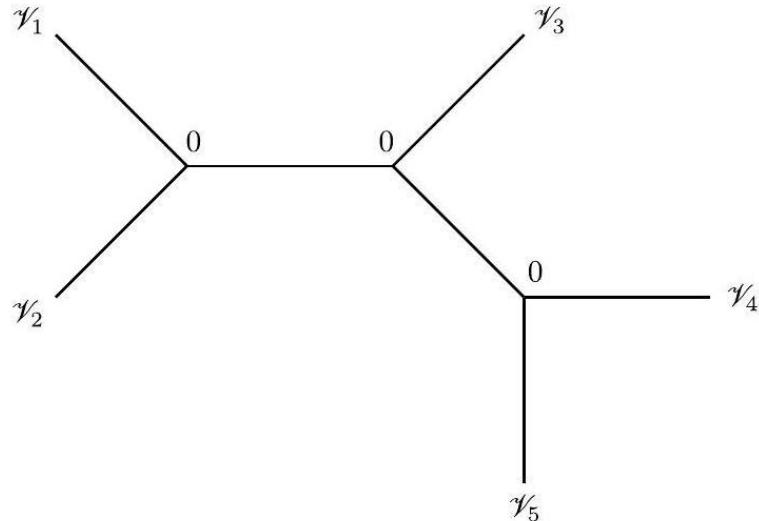
$$\mathcal{V}_{0,5}(\mathcal{V}_1, \dots, \mathcal{V}_5) := \int_{\mathcal{R}_{0,5}} \omega_4^{0,5}(\mathcal{V}_1, \dots, \mathcal{V}_5)$$





local coord.





Figuras 24, 25, 26, 27, 28, 29 y 30. Diagramas de Feynman de una partícula supermasiva.

$$\mathcal{V}_{0,2} := \Delta^{-1} \Delta \Delta^{-1} = \Delta^{-1}$$

$$\mathcal{V}_{0,2}(\mathcal{V}_1, \mathcal{V}_2) := \langle \mathcal{V}_1 | c_0^+ c_0^- L_0^+ \delta_{L_0^-, 0} | \mathcal{V}_2 \rangle.$$

$$\mathcal{V}_n(\mathcal{V}_1, \dots, \mathcal{V}_n) := \sum_{g \geq 0} (\hbar g_s^2)^g \mathcal{V}_{g,n}(\mathcal{V}_1, \dots, \mathcal{V}_n)$$

$$q = e^{-s+i\theta}, s \in [s_0, \infty), \theta \in [0, 2\pi).$$

$$\int_{s_0}^{\infty} ds e^{-s L_0^+} = \frac{e^{-s_0 L_0^+}}{L_0^+}$$

$$\Delta(s_0) = b_0^+ \frac{e^{-s_0 L_0^+}}{L_0^+} b_0^- \delta_{L_0^-, 0}.$$



$$\Delta_{\alpha \beta}(k)\!:=\!\frac{{\rm e}^{-\frac{\alpha' s_0}{2}(k^2+m_\alpha^2)}}{k^2+m_\alpha^2}M_{\alpha \beta}(k)$$

$$\mathcal{V}_{g,n}^{\text{1PI}}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\mathcal{V}_1\;\widehat{\mathcal{V}_n}:=\int_{\mathcal{R}_{g,n}^{\text{1PI}}}^{g/\text{1PI}}\omega_{\mathsf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n),$$

$$\langle \mathcal{V}^{g,n}\mid \otimes_i \mathcal{V}_i\rangle\!:=\mathcal{V}_{g,n}(\otimes_i \mathcal{V}_i).$$

$$\mathcal{V}_{g,n+1}(\mathcal{V}_0,\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=\langle \mathcal{V}_0|c_0^-|\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\rangle.$$

$$\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\!:=[\mathcal{V}_1,\ldots,\mathcal{V}_n]_g$$

$$\ell_{g,0}\!:=\![\cdot]_g\in\mathcal{H}.$$

$$N_{\mathrm{gh}}\Big(\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\Big)=3-2n+\sum_{i=1}^n\,N_{\mathrm{gh}}(\mathcal{V}_i)=3+\sum_{i=1}^n\,\big(N_{\mathrm{gh}}(\mathcal{V}_i)-2\big),$$

$$\left|\ell_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)\right|=1+\sum_{i=1}^n\,|\mathcal{V}_i|\,{\rm mod}2,$$

$$0=\sum_{g_1,g_2\geq 0\atop g_1+g_2=g}\sum_{n_1,n_2\geq 0\atop n_1+n_2=n}\frac{n!}{n_1!\,n_2!}\mathcal{V}_{g_1,n_1+1}\Big(\Psi^{n_1},\ell_{g_2,n_2}(\Psi^{n_2})\Big)+(-1)^{|\phi_s|}\mathcal{V}_{g-1,n+2}(\phi_s,b_0^-\phi_s^c,\Psi^n)$$

$$\mathcal{V}_{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)=\frac{1}{N}\sum_{a=1}^N\int_{\mathcal{R}_{g,n}^{(a)}}\omega_{\mathsf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n)$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)=\mathcal{V}_{0,3}(\mathcal{V}_3,\mathcal{V}_1,\mathcal{V}_2)+\cdots$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)=\omega_0^{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)\Big|_{\mathcal{S}_{0,3}}=\langle f_1\circ\mathcal{V}_1(0)f_2\circ\mathcal{V}_2(0)f_3\circ\mathcal{V}_3(0)\rangle,$$

$$\mathcal{V}_{0,3}(\mathcal{V}_3,\mathcal{V}_1,\mathcal{V}_2)=\langle f_1\circ\mathcal{V}_3(0)f_2\circ\mathcal{V}_1(0)f_3\circ\mathcal{V}_2(0)\rangle\neq\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3),$$

$$g(z_1)=z_2, g(z_2)=z_3, g(z_3)=z_1$$

$$\mathcal{V}_{0,3}(\mathcal{V}_3,\mathcal{V}_1,\mathcal{V}_2)=\langle g\circ f_1\circ\mathcal{V}_3(0)g\circ f_2\circ\mathcal{V}_1(0)g\circ f_3\circ\mathcal{V}_2(0)\rangle.$$

$$g\circ f_1=f_2, g\circ f_2=f_3, g\circ f_3=f_1$$

$$\mathcal{V}_{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)=\frac{1}{6}\sum_{a=1}^6\,\omega_0^{0,3}(\mathcal{V}_1,\mathcal{V}_2,\mathcal{V}_3)\Bigg|_{\mathcal{S}_{0,3}^{(a)}}$$



Espacio de Hilbert en CFT.

$$|\Psi\rangle = \sum_r \psi_r |\phi_r\rangle.$$

$$|\Psi\rangle = \sum_r (\psi_{\downarrow\downarrow,r} |\phi_{\downarrow\downarrow,r}\rangle + \psi_{\downarrow\uparrow,r} |\phi_{\downarrow\uparrow,r}\rangle + \psi_{\uparrow\downarrow,r} |\phi_{\uparrow\downarrow,r}\rangle + \psi_{\uparrow\uparrow,r} |\phi_{\uparrow\uparrow,r}\rangle),$$

$$\begin{aligned} b_0 |\phi_{\downarrow\downarrow,r}\rangle &= \bar{b}_0 |\phi_{\downarrow\downarrow,r}\rangle = 0, & b_0 |\phi_{\downarrow\uparrow,r}\rangle &= \bar{c}_0 |\phi_{\downarrow\uparrow,r}\rangle = 0, \\ c_0 |\phi_{\uparrow\downarrow,r}\rangle &= \bar{b}_0 |\phi_{\uparrow\downarrow,r}\rangle = 0, & c_0 |\phi_{\uparrow\uparrow,r}\rangle &= \bar{c}_0 |\phi_{\uparrow\uparrow,r}\rangle = 0. \end{aligned}$$

$$\langle \phi_r^c | \phi_s \rangle = \delta_{rs}$$

$$\begin{aligned} \langle \phi_{\downarrow\downarrow,r}^c | c_0 &= \langle \phi_{\downarrow\downarrow,r}^c | \bar{c}_0 = 0, \langle \phi_{\downarrow\uparrow,r}^c | c_0 &= \langle \phi_{\downarrow\uparrow,r}^c | \bar{b}_0 = 0 \\ \langle \phi_{\uparrow\downarrow,r}^c | b_0 &= \langle \phi_{\uparrow\downarrow,r}^c | \bar{c}_0 = 0, \langle \phi_{\uparrow\downarrow,r}^c | b_0 &= \langle \phi_{\uparrow\downarrow,r}^c | \bar{b}_0 = 0, \\ \langle \phi_{x,r}^c | \phi_{y,s} \rangle &= \delta_{xy} \delta_{rs}, \end{aligned}$$

$$G(\psi_r) = 2 - n_r$$

$$n_r = N_{\text{gh}}(\phi_r), n_r^c = N_{\text{gh}}(\phi_r^c) = 6 - n_r$$

Gauge fixed y transformaciones de gauge bajo propagador.

$$\Delta = b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-, 0}, \Delta_{rs} = \langle \phi_r^c | b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-, 0} | \phi_s^c \rangle.$$

$$S_{0,2} = \frac{1}{2} \langle \Psi | K | \Psi \rangle = \frac{1}{2} \psi_r K_{rs} \psi_s$$

$$K = c_0^- c_0^+ L_0^+ \delta_{L_0^-, 0} K_{rs} = \langle \phi_r | c_0^- c_0^+ L_0^+ \delta_{L_0^-, 0} | \phi_s \rangle$$

$$K = \frac{1}{2} c_0 \bar{c}_0 L_0^+ \delta_{L_0^-, 0}$$

$$\ker K|_{\mathcal{H}} \neq \emptyset$$

$$K_{rs} = \frac{1}{2} \begin{pmatrix} \langle \phi_{\downarrow\downarrow,r} | \\ \langle \phi_{\downarrow\uparrow,r} | \\ \langle \phi_{\uparrow\downarrow,r} | \\ \langle \phi_{\uparrow\uparrow,r} | \end{pmatrix}^t \begin{pmatrix} c_0 \bar{c}_0 L_0^+ & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 \end{pmatrix} \begin{pmatrix} |\phi_{\downarrow\downarrow,s}\rangle \\ |\phi_{\downarrow\uparrow,s}\rangle \\ |\phi_{\uparrow\downarrow,s}\rangle \\ |\phi_{\uparrow\uparrow,s}\rangle \end{pmatrix}.$$

$$L_0^- |\Psi\rangle = 0, b_0^- |\Psi\rangle = 0, b_0^+ |\Psi\rangle = 0$$

$$|\Psi\rangle = \sum_r \psi_{\downarrow\downarrow,r} |\phi_{\downarrow\downarrow,r}\rangle.$$



$$\begin{aligned}
K_{rs}\Delta_{st} &= \langle \phi_r | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} | \phi_s \rangle \langle \phi_s^c | b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-,0} | \phi_t^c \rangle \\
&= \langle \phi_r | c_0^- c_0^+ L_0^+ \delta_{L_0^-,0} b_0^+ b_0^- \frac{1}{L_0^+} \delta_{L_0^-,0} | \phi_t^c \rangle \\
&= \langle \phi_r | \{c_0^-, b_0^-\} \{c_0^+, b_0^+\} | \phi_t^c \rangle \\
&= \langle \phi_r | \phi_t^c \rangle = \delta_{rt}.
\end{aligned}$$

$$1=\sum_r|\phi_r\rangle\langle\phi_r^c|=\sum_r|\phi_{\downarrow\downarrow,r}\rangle\langle\phi_{\downarrow\downarrow,r}^c|.$$

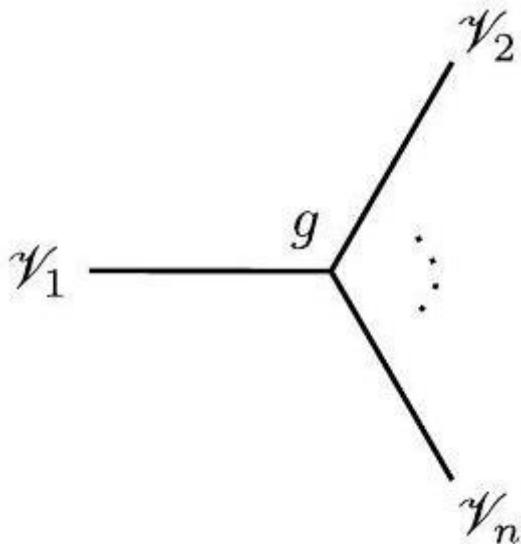
$$S_{0,2}=\frac{1}{2}\mathcal{V}_{0,2}(\Psi^2)=\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+\delta_{L_0^-,0}|\Psi\rangle.$$

$$S=\int~{\rm d}^Dx\left(\frac{1}{2}\phi(x)(-\partial^2+m^2)\phi(x)+\frac{\lambda}{n!}\phi(x)^n\right)$$

$$\phi_k(x) = {\rm e}^{{\rm i} k \cdot x}$$

$$\begin{aligned}
V_n(k_1, \dots, k_n) &= \frac{\lambda}{n!} \int ~{\rm d}^Dx n! \prod_{i=1}^n \phi_{k_i}(x) = \lambda \int ~{\rm d}^Dx {\rm e}^{{\rm i}(k_1 + \dots + k_n)x} \\
&= \lambda (2\pi)^D \delta^{(D)}(k_1 + \dots + k_n)
\end{aligned}$$

$$\mathcal{V}_{g,n}(\mathcal{V}_1,\dots,\mathcal{V}_n)=\int_{\mathcal{R}_{g,n}}\omega_{\mathbf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\dots,\mathcal{V}_n)=$$



$$S_{g,n} = \hbar^g \frac{g_s^{2g-2+n}}{n!} \mathcal{V}_{g,n}(\Psi^n)$$

$$\begin{aligned}
&\mathcal{V}_{g,m}(\mathcal{V}_1, \dots, \mathcal{V}_{m-1}, \phi_r) \langle \phi_r^c | b_0^+ b_0^- \frac{1}{L_0^+} | \phi_s^c \rangle \mathcal{V}_{g',n}(\mathcal{W}_1, \dots, \mathcal{W}_{n-1}, \phi_s) \\
&= \mathcal{V}_{g,m}(\mathcal{V}_1, \dots, \mathcal{V}_{m-1}, \phi_{\downarrow\downarrow,r}) \langle \phi_{\downarrow\downarrow,r}^c | b_0^+ b_0^- \frac{1}{L_0^+} | \phi_{\downarrow\downarrow,s}^c \rangle \mathcal{V}_{g',n}(\mathcal{W}_1, \dots, \mathcal{W}_{n-1}, \phi_{\downarrow\downarrow,s})
\end{aligned}$$



$$S=\sum_{g,n\geq 0}\hbar^g\frac{g_s^{2g-2+n}}{n!}\mathcal{V}_{g,n}(\Psi^n)\!:=\!\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+\delta_{L_0^-,0}|\Psi\rangle+\sum'_{g,n\geq 0}\hbar^g\frac{g_s^{2g-2+n}}{n!}\mathcal{V}_{g,n}(\Psi^n)$$

$$\mathcal{V}_{g,n}(\Psi^n)\!:=\langle\Psi|c_0^-|\ell_{g,n-1}(\Psi^{n-1})\rangle$$

$$S=\sum_{g,n\geq 0}\hbar^g\frac{g_s^{2g-2+n}}{n!}\langle\Psi|c_0^-|\ell_{g,n-1}(\Psi^{n-1})\rangle$$

$$\ell_{0,1}(\Psi_{\rm cl})=c_0^+L_0^+|\Psi_{\rm cl}\rangle$$

$$S=\sum_{\substack{g,n\geq 0\\ \chi_{g,n}\leq 0}}\hbar^g\frac{g_s^{2g-2+n}}{n!}\mathcal{V}_{g,n}(\Psi^n)$$

$$S_{\rm cl}=\frac{1}{2}\langle\Psi_{\rm cl}|c_0^-c_0^+L_0^+|\Psi_{\rm cl}\rangle+\sum_{n\geq 3}\frac{g_s^n}{n!}\mathcal{V}_{0,n}(\Psi_{\rm cl}^n)$$

$$S=\sum_{g,n\geq 0}\hbar^g g_s^{2g-2}\frac{1}{n!}\mathcal{V}_{g,n}(\Psi^n)\!:=\!\frac{1}{2g_s^2}\langle\Psi|c_0^-c_0^+L_0^+\delta_{L_0^-,0}|\Psi\rangle+\frac{1}{g_s^2}\sum'_{g,n\geq 0}\frac{(\hbar g_s^2)^g}{n!}\mathcal{V}_{g,n}(\Psi^n)$$

$$\frac{S}{\hbar}=\sum_{g,n\geq 0}(\hbar g_s^2)^{g-1}\frac{1}{n!}\mathcal{V}_{g,n}(\Psi^n)$$

$$S\rightarrow \alpha S\implies g_s^2\rightarrow \frac{g_s^2}{\alpha}$$

$$L_0^+|\Psi\rangle=0$$

$$\Psi_{\rm cl}\in {\mathcal H}^{-}\cap \ker L_0^-, N_{\rm gh}(\Psi_{\rm cl})=2$$

$$S_{0,2}=\frac{1}{2}\langle\Psi|c_0^-Q_B|\Psi\rangle.$$

$$S_{\rm cl}=\frac{1}{2}\langle\Psi_{\rm cl}|c_0^-Q_B|\Psi_{\rm cl}\rangle+\frac{1}{g_s^2}\sum_{n\geq 3}\frac{g_s^n}{n!}\mathcal{V}_{0,n}(\Psi_{\rm cl}^n)$$

$$\mathcal{V}_{0,2}(\Psi_{\rm cl}^2)\!:=\langle\Psi_{\rm cl}|c_0^-Q_B|\Psi_{\rm cl}\rangle$$

$$S_{\rm cl}=\frac{1}{g_s^2}\sum_{n\geq 2}\frac{g_s^n}{n!}\mathcal{V}_{0,n}(\Psi_{\rm cl}^n)=\frac{1}{g_s^2}\sum_{n\geq 2}\frac{g_s^n}{n!}\langle\Psi_{\rm cl}|c_0^-|\ell_{0,n-1}(\Psi_{\rm cl}^{n-1})\rangle,$$

$$\ell_{0,1}(\Psi_{\rm cl})=Q_B|\Psi_{\rm cl}\rangle.$$



$$\mathcal{F}_{\text{cl}}(\Psi_{\text{cl}}) := \sum_{n \geq 1} \frac{g_s^{n-1}}{n!} \ell_{0,n}(\Psi_{\text{cl}}^n) = Q_B |\Psi_{\text{cl}}\rangle + \sum_{n \geq 2} \frac{g_s^{n-1}}{n!} \ell_{0,n}(\Psi_{\text{cl}}^n) = 0$$

$$\delta S_{\text{cl}} = \frac{1}{g_s^2} \sum_{n \geq 2} \frac{g_s^n}{n!} n \{\delta\Psi_{\text{cl}}, \Psi_{\text{cl}}^{n-1}\}_0 = \frac{1}{g_s^2} \sum_{n \geq 2} \frac{g_s^n}{(n-1)!} \langle \delta\Psi_{\text{cl}} | c_0^- \Big| \ell_{0,n-1}(\Psi_{\text{cl}}^{n-1}) \rangle.$$

$$\delta_\Lambda S_{\text{cl}} = 0$$

$$\delta_\Lambda \Psi_{\text{cl}} = \sum_{n \geq 0} \frac{g_s^n}{n!} \ell_{0,n+1}(\Psi_{\text{cl}}^n, \Lambda) = Q_B |\Lambda\rangle + \sum_{n \geq 1} \frac{g_s^n}{n!} \ell_{0,n+1}(\Psi_{\text{cl}}^n, \Lambda)$$

$$[\delta_{\Lambda_2}, \delta_{\Lambda_1}] \Psi_{\text{cl}} = \delta_{\Lambda(\Lambda_1, \Lambda_2, \Psi_{\text{cl}})} |\Psi_{\text{cl}}\rangle + \sum_{n \geq 0} \frac{g_s^{n+2}}{n!} \ell_{0,n+3}(\Psi_{\text{cl}}^n, \Lambda_2, \Lambda_1, \mathcal{F}_{\text{cl}}(\Psi_{\text{cl}})),$$

$$\Lambda(\Lambda_1, \Lambda_2, \Psi_{\text{cl}}) = \sum_{n \geq 0} \frac{g_s^{n+1}}{n!} \ell_{0,n+2}(\Lambda_1, \Lambda_2, \Psi_{\text{cl}}^n) = g_s \ell_{0,2}(\Lambda_1, \Lambda_2) + \sum_{n \geq 1} \frac{g_s^{n+1}}{n!} \ell_{0,n+2}(\Lambda_1, \Lambda_2, \Psi_{\text{cl}}^n)$$

$$\forall n \geq 4: \mathcal{V}_{0,4}(\mathcal{V}_1, \dots, \mathcal{V}_n) = 0, \ell_{g,n-1}(\mathcal{V}_1, \dots, \mathcal{V}_{n-1}) = 0, (\text{cubic theory}),$$

$$\begin{aligned} \delta_\Lambda S_{\text{cl}} &= \sum_{n \geq 2} \frac{g_s^{n-2}}{n!} n \mathcal{V}_{0,n}(\delta\Psi_{\text{cl}}, \Psi_{\text{cl}}^{n-1}) = \sum_{m,n \geq 0} \frac{g_s^{m+n-1}}{m! n!} \mathcal{V}_{0,n+1}(\ell_{0,m+1}(\Psi_{\text{cl}}^m, \Lambda), \Psi_{\text{cl}}^n) \\ &= \sum_{m \geq 0} \sum_{n=0}^m \frac{g_s^{m-1}}{(m-n)! n!} \langle \ell_{m-n+1}(\Psi_{\text{cl}}^{m-n}, \Lambda) | c_0^- | \ell_{0,n}(\Psi_{\text{cl}}^n) \rangle \\ &= \langle \ell_{0,n}(\Psi_{\text{cl}}^n) | c_0^- | \ell_{0,m-n+1}(\Psi_{\text{cl}}^{m-n}, \Lambda) \rangle \\ &= \mathcal{V}_{0,m-n+2}(\ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n}, \Lambda) \\ &= -\mathcal{V}_{0,m-n+2}(\Lambda, \ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n}) \\ &= \langle \Lambda | c_0^- | \ell_{0,m-n+1}(\ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n}) \rangle. \end{aligned}$$

$$0 = \sum_{n=0}^m \frac{m!}{(m-n)! n!} \ell_{0,m-n+1}(\ell_{0,n}(\Psi_{\text{cl}}^n), \Psi_{\text{cl}}^{m-n})$$

Formalismo BV.

$$\begin{aligned} S &= \frac{1}{g_s^2} \sum_{g \geq 0} \hbar^g g_s^{2g} \sum_{n \geq 0} \frac{g_s^n}{n!} \mathcal{V}_{g,n}(\Psi^n) \\ &= \frac{1}{2} \langle \Psi | c_0^- Q_B | \Psi \rangle + \sum'_{g,n \geq 0} \frac{\hbar^g g_s^{2g-2+n}}{n!} \mathcal{V}_{g,n}(\Psi^n) \\ &= \frac{1}{g_s^2} \sum_{g,n \geq 0} \frac{\hbar^g g_s^{2g-2+n}}{n!} \langle \Psi | c_0^- | \ell_{g,n-1}(\Psi^{n-1}) \rangle \end{aligned}$$

$$\Psi \in \mathcal{H}^- \cap \ker L_0^-.$$

$$(S, S) - 2\hbar \Delta S = 0$$



$$|\Psi\rangle=\sum_r\psi_r|\phi_r\rangle$$

$$\Psi = \Psi_+ + \Psi_-$$

$$\Psi_{-}=\sum_r\sum_{n_r\leq 2}|\phi_r\rangle\psi^r,\Psi_{+}=\sum_r\sum_{n_r^c>2}b_0^{-}|\phi_r^c\rangle\psi_r^{*}$$

$$G(\psi^r)\geq 0, G(\psi_r^*)<0$$

$$G(\psi_r^*)=-1-G(\psi^r)$$

$$\begin{aligned}G(\psi_r^*) &= 2-N_{\mathrm{gh}}(b_0^{-}\phi_r^c)=2+1-n_r^c \\&=3-(6-n_r)=-3+\left(2-G(\psi^r)\right)=-1-G(\psi^r)\end{aligned}$$

$$\frac{\partial_RS}{\partial\psi^r}\frac{\partial_L S}{\partial\psi_r^*}+\hbar\frac{\partial_R\partial_L S}{\partial\psi^r\partial\psi_r^*}=0$$

$$\sum_{g_1,g_2\geq 0\atop g_1+g_2=g}\sum_{n_1,n_2\geq 0\atop n_1+n_2=n}\frac{\partial_RS_{g_1,n_1}}{\partial\psi^r}\frac{\partial_L S_{g_2,n_2}}{\partial\psi_r^*}+\hbar\frac{\partial_R\partial_L S_{g-1,n}}{\partial\psi^r\partial\psi_r^*}=0$$

$$\begin{aligned}\mathcal{V}_n^{1\text{PI}}(\mathcal{V}_1,\ldots,\mathcal{V}_n)&:=\mathcal{V}_1\rightarrow\sum_{\mathcal{V}_n}^{\mathcal{V}_2}:=\sum_{g\geq 0}(\hbar g_s^2)^g\mathcal{V}_{g,n}^{1\text{PI}}(\mathcal{V}_1,\ldots,\mathcal{V}_n),\\\mathcal{V}_{g,n}^{1\text{PI}}(\mathcal{V}_1,\ldots,\mathcal{V}_n)&:=\int_{\mathcal{R}_{g,n}^{1\text{PI}}}\omega_{\mathsf{M}_{g,n}}^{g,n}(\mathcal{V}_1,\ldots,\mathcal{V}_n),\end{aligned}$$

$$S_{1\text{PI}}=\frac{1}{g_s^2}\sum_{n\geq 0}\frac{g_s^n}{n!}\mathcal{V}_n^{1\text{PI}}(\Psi^n)\!:=\!\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+|\Psi\rangle+\frac{1}{g_s^2}\sum_{n\geq 0}'\frac{g_s^n}{n!}\mathcal{V}_n^{1\text{PI}}(\Psi^n)$$

$$S_{1\text{PI}}=\frac{1}{2}\langle\Psi|c_0^-Q_B|\Psi\rangle+\frac{1}{g_s^2}\sum_{n\geq 0}'\frac{g_s^n}{n!}\mathcal{V}_n^{1\text{PI}}(\Psi^n)$$

Independencia background.

$$\delta S_{\rm cft} = \frac{\lambda}{2\pi} \int \;\; {\rm d}^2 z \varphi(z,\bar{z})$$

$$S_1[\Psi_1]=\frac{1}{g_s^2}\Biggl(\frac{1}{2}\langle\Psi_1|c_0^-Q_B|\Psi_1\rangle+\sum_{n\geq 0}'\frac{1}{n!}\mathcal{V}_n^{1\text{PI}}(\Psi_1^n)\Biggr)$$

$$\mathcal{F}_1(\Psi_1)=Q_B|\Psi_1\rangle+\sum_n\frac{1}{n!}\ell_n(\Psi_1^n)=0$$



Deformación de CFT.

$$S_{\text{cft},2}[\psi_1] = S_{\text{cft},1}[\psi_1] + \frac{\lambda}{2\pi} \int d^2z \varphi(z, \bar{z})$$

$$\begin{aligned} \left\langle \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_2 &= \left\langle \exp \left(-\frac{\lambda}{2\pi} \int d^2z \varphi(z, \bar{z}) \right) \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 \\ &\approx \left\langle \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 - \frac{\lambda}{2\pi} \int_{\Sigma} d^2z \left\langle \varphi(z, \bar{z}) \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 \end{aligned}$$

$$\delta \left\langle \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 = -\frac{\lambda}{2\pi} \int_{\Sigma - \cup_i D_i} d^2z \left\langle \varphi(z, \bar{z}) \prod_i \mathcal{O}_i(z_i, \bar{z}_i) \right\rangle_1 .$$

$$\delta L_n = \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} z^{n+1} \varphi(z, \bar{z}), \quad \delta \bar{L}_n = \lambda \oint_{|z|=1} \frac{dz}{2\pi i} \bar{z}^{n+1} \varphi(z, \bar{z}).$$

$$\delta Q_B = \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} c(z) \varphi(z, \bar{z}) + \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} \bar{c}(\bar{z}) \varphi(z, \bar{z}).$$

$$\{Q_B, \delta Q_B\} = O(\lambda^2)$$

$$\delta L_0^- |\mathcal{O}\rangle = \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} z \sum_{p,q} z^{p-1} \bar{z}^{q-1} |\mathcal{O}_{p,q}\rangle - \lambda \oint_{|z|=1} \frac{d\bar{z}}{2\pi i} \sum_{p,q} z^{p-1} \bar{z}^{q-1} |\mathcal{O}_{p,q}\rangle,$$

$$\varphi(z, \bar{z}) \mathcal{O}(0,0) = \sum_{p,q} z^{p-1} \bar{z}^{q-1} \mathcal{O}_{p,q}(0,0)$$

$$S_2[\Psi_1] = S_1[\Psi_1] + \delta S_1[\Psi_1]$$

$$\delta S_1[\Psi_1] = \frac{1}{g_s^2} \left(\frac{1}{2} \langle \Psi_1 | c_0^- \delta Q_B | \Psi_1 \rangle + \sum_{n \geq 0} \frac{1}{n!} \delta \mathcal{V}_n(\Psi_1^n) \right)$$

$$\mathcal{F}_2(\Psi_1) = \mathcal{F}_1(\Psi_1) + \lambda \delta \mathcal{F}_1(\Psi_1) = 0$$

$$\lambda \delta \mathcal{F}_1(\Psi_1) = \delta Q_B |\Psi_1\rangle + \sum_n \frac{1}{n!} \delta \ell_n(\Psi_1^n)$$

$$|\Psi_1\rangle = \lambda |\Psi_0\rangle, |\Psi_0\rangle = c_1 \bar{c}_1(0) |\varphi\rangle$$

$$|\Psi_1\rangle = \lambda |\Psi_0\rangle + |\Psi'\rangle,$$

$$S_1[\Psi_1] = S_1[\Psi_0] + S'[\Psi']$$

$$S'[\Psi'] = \frac{1}{g_s^2} \left(\frac{1}{2} \langle \Psi' | c_0^- Q_B | \Psi' \rangle + \sum_n \frac{1}{n!} (\mathcal{V}_n(\Psi'^n) + \lambda \mathcal{V}_{n+1}(\Psi_0, \Psi'^n)) \right)$$

$$\mathcal{F}'(\Psi') := \mathcal{F}_1(\Psi') + \lambda \delta \mathcal{F}'(\Psi') = 0$$



$$\delta\mathcal{F}'(\Psi') = \sum_n \frac{1}{n!} \ell_{n+1}(\Psi_0, \Psi'^n)$$

$$\begin{aligned}\mathcal{F}_1(\Psi_1) + \lambda \delta\mathcal{F}_1(\Psi_1) &= (1 + \lambda \mathcal{M}(\Psi')) (\mathcal{F}_1(\Psi') + \lambda \delta\mathcal{F}'(\Psi')) \\ |\Psi_1\rangle &= |\Psi'\rangle + \lambda |\delta\Psi'\rangle\end{aligned}$$

$$\left. \frac{d}{d\lambda} \mathcal{F}_1(\Psi' + \lambda \delta\Psi') \right|_{\lambda=0} + \delta\mathcal{F}_1(\Psi_1) - \delta\mathcal{F}'(\Psi') = \mathcal{M}(\Psi') \mathcal{F}_1(\Psi')$$

$$\begin{aligned}\lambda Q_B |\delta\Psi'\rangle + \lambda \sum_n \frac{1}{n!} \ell_{n+1}(\delta\Psi', \Psi'^n) + \delta Q_B |\Psi'\rangle \\ + \sum_n \frac{1}{n!} \delta \ell_n(\Psi'^n) - \lambda \sum_n \frac{1}{n!} \ell_{n+1}(\Psi_0, \Psi'^n) = 0\end{aligned}$$

$$\begin{aligned}\Delta := \lambda \langle A | c_0^- Q_B | \delta\Psi' \rangle + \lambda \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \delta\Psi', \Psi'^n) + \langle A | c_0^- \delta Q_B | \Psi' \rangle + \sum_n \frac{1}{n!} \delta \mathcal{V}_{n+1}(A, \Psi'^n) \\ - \lambda \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n)\end{aligned}$$

$$\langle A | c_0^- \delta Q_B | B \rangle = \lambda \mathcal{V}'_{0,3}(\Psi_0, B, A), \delta \mathcal{V}_n(\Psi'^n) = \lambda \mathcal{V}'_{n+1}(\Psi_0, \Psi'^n),$$

$$\langle A | c_0^- | \delta\Psi' \rangle = \sum_n \frac{1}{n!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A)$$

$$\mathcal{V}'_n = \sum_{g \geq 0} \mathcal{V}'_{g,n}, \mathcal{B}_n = \sum_{g \geq 0} \mathcal{B}_{g,n}$$

$$\begin{aligned}\Delta = - \sum_n \frac{1}{n!} \mathcal{B}_{n+2} \left(\Psi_0, \Psi'^n, Q_B A \right) + \sum_{m,n} \frac{1}{m! n!} \mathcal{B}_{n+2} \left(\Psi_0, \Psi'^m, \ell_{n+1}(A, \Psi'^n) \right) \\ + \sum_n \frac{1}{n!} \mathcal{V}'_{n+2}(A, \Psi_0, \Psi'^n) - \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n)\end{aligned}$$

$$\begin{aligned}\mathcal{B}_{n+2}(\Psi_0, \Psi'^n, Q_B A) &= \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) + n \mathcal{B}_{n+2}(\Psi_0, \Psi'^{n-1}, Q_B \Psi', A) \\ &= \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) - \sum_m \frac{n}{m!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^{n-1}, \ell_m(\Psi'^m), A)\end{aligned}$$



$$\begin{aligned}
\Delta = & \sum_n \frac{1}{n!} \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) - \sum_{m,n} \frac{1}{m! n!} \mathcal{B}_{n+3}(\Psi_0, \Psi'^n, \ell_m(\Psi'^m), A) \\
& + \sum_{m,n} \frac{1}{m! n!} \mathcal{B}_{n+2}(\Psi_0, \Psi'^m, \ell_{n+1}(A, \Psi'^n)) + \sum_n \frac{1}{n!} \mathcal{V}'_{n+2}(A, \Psi_0, \Psi'^n) \\
& - \sum_n \frac{1}{n!} \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n) \\
& \partial \mathcal{B}_{n+2}(\Psi_0, \Psi'^n, A) \\
& = -\mathcal{V}'_{n+2}(A, \Psi_0, \Psi'^n) + \mathcal{V}_{n+2}(A, \Psi_0, \Psi'^n) \\
& + \sum_{\substack{m_1, m_2 \\ m_1 + m_2 = n}} \frac{n!}{m_1! m_2!} \mathcal{B}_{m_1+3}(\Psi_0, \Psi'^{m_1}, \ell_{m_2}(\Psi'^{m_2}), A) \\
& - \sum_{\substack{m_1, m_2 \\ m_1 + m_2 = n}} \frac{n!}{m_1! m_2!} \mathcal{B}_{m_1+2}(\Psi_0, \Psi'^{m_1}, \ell_{m_2+1}(A, \Psi'^{m_2}))
\end{aligned}$$

Membranas, Simetrías y espacios planckianos en dimensión \mathbb{R}^4 . Singularidad.

$$\begin{aligned}
\partial \mathcal{B}_{0,3} &= \mathcal{V}_{0,3} - \mathcal{V}'_{0,3} \\
T(z)T(w) &\sim \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w} \\
G(z)G(w) &\sim \frac{2c/3}{(z-w)^3} + \frac{2T(w)}{(z-w)} \\
T(z)G(w) &\sim \frac{3}{2} \frac{G(w)}{(z-w)^2} + \frac{\partial G(w)}{(z-w)}
\end{aligned}$$

$$h(\beta) = \left(\frac{3}{2}, 0\right), h(\gamma) = \left(-\frac{1}{2}, 0\right)$$

$$\gamma(z)\beta(w) \sim \frac{1}{z-w}, \beta(z)\gamma(w) \sim -\frac{1}{z-w}$$

$$T^{\text{gh}} = -2b\partial c + c\partial b, T^{\beta\gamma} = \frac{3}{2}\beta\partial\gamma + \frac{1}{2}\gamma\partial\beta$$

$$N_{\text{gh}}(b) = N_{\text{gh}}(\beta) = -1, N_{\text{gh}}(c) = N_{\text{gh}}(\gamma) = 1$$

$$\gamma = \eta e^\phi, \beta = \partial\xi e^{-\phi}$$

$$\delta(\gamma) = e^{-\phi}, \delta(\beta) = e^\phi$$

$$\int \, dc_0 = 0 \implies \int \, dc_0 c_0 = 1$$



$$\int \; \mathrm{d}\gamma_0 = \infty \implies \int \; \mathrm{d}\gamma_0 \delta(\gamma_0) = 1$$

$$T^{\beta\gamma}=T^{\eta\xi}+T^\phi,$$

$$T^{\eta\xi}=-\eta\partial\xi,T^\phi=-\frac{1}{2}(\partial\phi)^2-\partial^2\phi.$$

$$\xi(z)\eta(w)\sim \frac{1}{z-w}, {\rm e}^{q_1\phi(z)}{\rm e}^{q_2\phi(w)}\sim \frac{{\rm e}^{(q_1+q_2)\phi(w)}}{(z-w)^{q_1q_2}}, \partial\phi(z)\partial\phi(w)\sim -\frac{1}{(z-w)^2}.$$

$$N_{\mathrm{gh}}(\eta)=1,N_{\mathrm{gh}}(\xi)=-1,N_{\mathrm{gh}}(\phi)=0.$$

$$N_{\mathrm{pic}}\left({\rm e}^{q\phi}\right)=q,N_{\mathrm{pic}}\left(\xi\right)=1,N_{\mathrm{pic}}\left(\eta\right)=-1.$$

$$N_{\mathrm{pic}}=2(g-1)=-\chi_g.$$

$$h\bigl({\rm e}^{q\phi}\bigr)=-\frac{q}{2}(q+2)$$

$$h\bigl({\rm e}^{\phi}\bigr)=\frac{3}{2}, h\bigl({\rm e}^{-\phi}\bigr)=\frac{1}{2}.$$

$$\begin{gathered}j_B=c\big(T^{\mathrm{m}}+T^{\beta\gamma}\big)+\gamma G+bc\partial c-\frac{1}{4}\gamma^2b\\\bar J_B=\bar c\bar T^{\mathrm{m}}+\bar b\bar c\bar\partial\bar c\end{gathered}$$

$$\mathcal{X}(z)=\{Q_B,\xi(z)\}=c\partial\xi+{\rm e}^\phi G-\frac{1}{4}\partial\eta{\rm e}^{2\phi}b-\frac{1}{4}\partial\big(\eta{\rm e}^{2\phi}b\big),$$

$$\mathcal{X}_0=\frac{1}{2\pi\mathrm{i}}\oint\;\frac{\mathrm{d} z}{z}\mathcal{X}(z)$$

$$\mathcal{H}_{\text{small}}=\{| \psi\rangle|\eta_0|\psi\rangle=0\}.$$

$$\mathcal{H}_{\text{small}}=\mathcal{H}_{\text{large}}\cap\ker\eta_0$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1{\rm e}^{-2\phi(z)}|k'\rangle=(2\pi)^D\delta^{(D)}(k+k')$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1{\rm e}^{-2\phi(z)}{\rm e}^{-\bar{\phi}(\bar{w})}|k'\rangle=-(2\pi)^D\delta^{(D)}(k+k')$$

$$\mathcal{H}_T=\mathcal{H}_{\mathrm{NS}}\oplus\mathcal{H}_{\mathrm{R}}$$

$$\widehat{\mathcal{H}}_T=\mathcal{H}_{-1}\oplus\mathcal{H}_{-1/2}, \widetilde{\mathcal{H}}_T=\mathcal{H}_{-1}\oplus\mathcal{H}_{-3/2}$$

$$|p\rangle={\rm e}^{p\phi}(0)|0\rangle.$$

$$\forall n\geq-p-\frac{1}{2} \colon \beta_n|p\rangle=0\\ \forall n\geq p+\frac{3}{2} \colon \gamma_n|p\rangle=0$$

$$\widehat{\mathcal{H}}_T=\mathrm{Span}\{|\phi_r\rangle\}, \widetilde{\mathcal{H}}_T=\mathrm{Span}\{|\phi_r^c\rangle\}$$



$$\langle \phi_r^c\mid \phi_s\rangle=\delta_{rs}.$$

$$1 = \sum_r |\phi_r\rangle\langle\phi_r^c|$$

$$1 = \sum_r (-1)^{|\phi_r|} |\phi_r^c\rangle\langle\phi_r|$$

$${\mathcal G}=\begin{cases} 1 & \text{NS sector}\,, \\ {\mathcal X}_0 & \text{R sector}\,. \end{cases}$$

$$\left[\mathcal{G},L_0^\pm\right]=\left[\mathcal{G},b_0^\pm\right]=\left[\mathcal{G},Q_B\right]=0.$$

$$\mathrm{M}_{g,m,n}\!:=\dim\!\mathcal{M}_{g,m,n}=6g-6+2m+2n$$

$$N_{\rm gh}=6-6g, N_{\rm pic}=2g-2$$

$$n_{\rm pco} := 2g-2+m+\frac{n}{2}$$

$$A_{g,m,n}\big(\mathcal{V}_i^{\rm NS},\mathcal{V}_j^{\rm R}\big)=\int_{\mathcal{S}_{g,m,n}}\Omega_{\mathbf{M}_{g,m,n}}\big(\mathcal{V}_i^{\rm NS},\mathcal{V}_j^{\rm R}\big),$$

$$\Omega_{\mathbf{M}_{g,m,n}}=(-2\pi{\rm i})^{-\mathbf{M}_{g,m,n}^c}\left\langle\bigwedge_{\lambda=1}^{\mathbf{M}_{g,m,n}}\mathcal{B}_\lambda{\rm d} t_\lambda\prod_{A=1}^{n_{\rm pco}}\mathcal{X}(y_A)\prod_{i=1}^m\mathcal{V}_i^{\rm NS}\prod_{j=1}^n\mathcal{V}_j^{\rm R}\right\rangle_{\Sigma_{g,n}}$$

$$\begin{aligned} \mathcal{B}_\lambda = \sum_\alpha \oint_{c_\alpha} \frac{{\rm d}\sigma_\alpha}{2\pi{\rm i}} b(\sigma_\alpha) \frac{\partial F_\alpha}{\partial t_\lambda}\big(F_\alpha^{-1}(\sigma_\alpha)\big) & + \sum_\alpha \oint_{c_\alpha} \frac{{\rm d}\bar\sigma_\alpha}{2\pi{\rm i}} \bar b(\bar\sigma_\alpha) \frac{\partial \bar F_\alpha}{\partial t_\lambda}\big(\bar F_\alpha^{-1}(\bar\sigma_\alpha)\big) \\ & - \sum_A \frac{1}{\mathcal{X}(y_A)} \frac{\partial y_A}{\partial t_\lambda} \partial\xi(y_A) \end{aligned}$$

$$\mathcal{X}(y_A)-\partial\xi(y_A){\rm d} y_A$$

$$n_{\rm pco}^{(1)}+n_{\rm pco}^{(2)}=2(g_1+g_2)-2+(m_1+m_2-2)+\frac{n_1+n_2}{2}=n_{\rm pco}^{\rm (NS)}$$

$$\Delta_{\rm NS}=b_0^+b_0^-\frac{1}{L_0+\bar L_0}\delta(L_0^-)$$

$$n_{\rm pco}^{(1)}+n_{\rm pco}^{(2)}=2(g_1+g_2)-2+(m_1+m_2)+\frac{n_1+n_2-2}{2}-1=n_{\rm pco}^{\rm (R)}-1$$

$$\Delta_{\rm R}=b_0^+b_0^-\frac{{\mathcal X}_0}{L_0+\bar L_0}\delta(L_0^-)$$

$$\mathcal{X}_0=\frac{1}{2\pi{\rm i}}\oint\frac{{\rm d} w_n^{(1)}}{w_n^{(1)}}\mathcal{X}\left(w_n^{(1)}\right)=\frac{1}{2\pi{\rm i}}\oint\frac{{\rm d} w_n^{(2)}}{w_n^{(2)}}\mathcal{X}\left(w_n^{(2)}\right)$$



$$\Delta = b_0^+ b_0^- \frac{\mathcal{G}}{L_0 + \bar L_0} \delta(L_0^-)$$

$$\Delta_{\mathrm{NS}} \sim \frac{1}{k^2+m^2}, \Delta_{\mathrm{R}} \sim \frac{\mathrm{i}\,\partial + m}{k^2+m^2}$$

$$\mathcal{C}(x_i,y_j,z_q)=\left\langle \prod_{i=1}^{n+1}\;\xi(x_i)\prod_{j=1}^n\;\eta(y_j)\prod_{k=1}^m\;{\rm e}^{q_k\phi(z_k)}\right\rangle$$

$$=\frac{\prod_{j'=1}^n\vartheta_\delta(-y_{j'}+\sum_ix_i-\sum_jy_j+\sum_kq_kz_k)}{\prod_{i'=1}^{n+1}\vartheta_\delta(-x_{i'}+\sum_ix_i-\sum_jy_j+\sum_kq_kz_k)}\times\frac{\prod_{i< i'}E(x_i,x_{i'})\prod_{j< j'}E(y_j,y_{j'})}{\prod_{i,j}E(x_i,y_j)\prod_{k,\ell}E(z_k,z_\ell)^{q_kq_\ell}}$$

$$\sum_k\;q_k=0.$$

$$E(x,y) = \frac{\vartheta_1(x-y)}{\vartheta_1'(0)} \sim_{x \rightarrow y} x - y$$

$$\vartheta_\delta\left(\sum_{i=2}^{n+1}\;x_i-\sum_{j=1}^n\;y_j+\sum_{k=1}^m\;q_kz_k\right)=0$$

$$\left\langle A^{(p)}\right|Q=0,$$

$$Q=Q_B\otimes 1^{\otimes n-1}+\cdots+1^{\otimes n-1}\otimes Q_B$$

$$\left\langle A^{(p)}\right|\eta=0$$

$$\eta=\eta_0\otimes 1^{\otimes n-1}+\cdots+1^{\otimes n-1}\otimes \eta_0.$$

$$\left\langle A^{(p)}\right|=\left\langle \alpha^{(p)}\right|Q,$$

$$\left\langle \alpha^{(p)}\right|\eta Q=0$$

$$\left\langle A^{(p-1)}\right|=\left\langle \alpha^{(p)}\right|\eta$$

$$\Psi=\Psi_{-1}+\Psi_{-1/2}$$

$$\eta_0|\Psi\rangle=0$$

$$\Delta=b_0^+b_0^-\frac{\mathcal{G}}{L_0+\bar L_0}\delta(L_0^-), \mathcal{G}=\begin{cases} 1 & \text{NS} \\ \mathcal{X}_0 & \text{R} \end{cases}$$

$$b_0^-|\Psi\rangle=L_0^-|\Psi\rangle=0,b_0^+|\Psi\rangle=0$$

$$X=G_0\delta(\beta_0)+b_0\delta'(\beta_0), Y=-c_0\delta'(\gamma_0)$$

$$[Q_B,X]=0$$

$$XYX=X$$



$$XY\left|\Psi_{-1/2}\right\rangle = \left|\Psi_{-1/2}\right\rangle$$

$$Bc_0^- |\Psi\rangle = |\Psi\rangle,$$

$$B=b_0^-\int_0^{2\pi}\frac{\mathrm{d}\theta}{2\pi}e^{\mathrm{i}\theta L_0^-}=\delta(b_0^-)\delta(L_0^-)$$

$$S_{0,2}=-\frac{1}{2}\langle\Psi_{-1}|c_0^-Q_B|\Psi_{-1}\rangle-\frac{1}{2}\langle\Psi_{-1/2}|c_0^-YQ_B|\Psi_{-1/2}\rangle.$$

$$\delta|\Psi\rangle=Q_B|\Lambda\rangle$$

$$\Lambda=\Lambda_{-1}+\Lambda_{-1/2}$$

$$\widetilde{\Psi}=\widetilde{\Psi}_{-1}+\widetilde{\Psi}_{-3/2}$$

$$b_0^-|\widetilde{\Psi}\rangle=L_0^-|\widetilde{\Psi}\rangle=0,b_0^+|\widetilde{\Psi}\rangle=0$$

$$S_{0,2}=\frac{1}{2}\langle\widetilde{\Psi}|c_0^-c_0^+L_0^+\mathcal{G}|\widetilde{\Psi}\rangle-\langle\widetilde{\Psi}|c_0^-c_0^+L_0^+|\Psi\rangle.$$

$$K=c_0^-c_0^+L_0^+\begin{pmatrix}-\mathcal{G}&1\\1&0\end{pmatrix}$$

$$\Delta=b_0^-b_0^+\frac{1}{L_0^+}\begin{pmatrix}0&1\\1&\mathcal{G}\end{pmatrix}$$

$$\delta|\Psi\rangle=Q_B|\Lambda\rangle,\delta|\widetilde{\Psi}\rangle=Q_B|\widetilde{\Lambda}\rangle,$$

$$Q_B|\Psi\rangle=0,Q_B|\widetilde{\Psi}\rangle=0$$

$$Q_B(|\Psi\rangle-\mathcal{G}|\widetilde{\Psi}\rangle)=0,Q_B|\widetilde{\Psi}\rangle=|J(\Psi)\rangle,$$

$$Q_B|\Psi\rangle=\mathcal{G}|J(\Psi)\rangle$$

$$S_{0,2}=-\frac{1}{2}\big\langle\langle\Psi_0,\eta_0Q_B\Psi_0\rangle\big\rangle,$$

$$\delta|\Psi_0\rangle=Q_B|\Lambda_0\rangle+\eta_0|\Omega_1\rangle,$$

$$Q_B\eta_0|\Psi_0\rangle=0$$

$$\xi_0|\Psi_0\rangle=0$$

$$|\Psi_0\rangle=\xi_0|\Psi_{-1}\rangle$$

$$Q_B|\Psi_{-1}\rangle=0$$

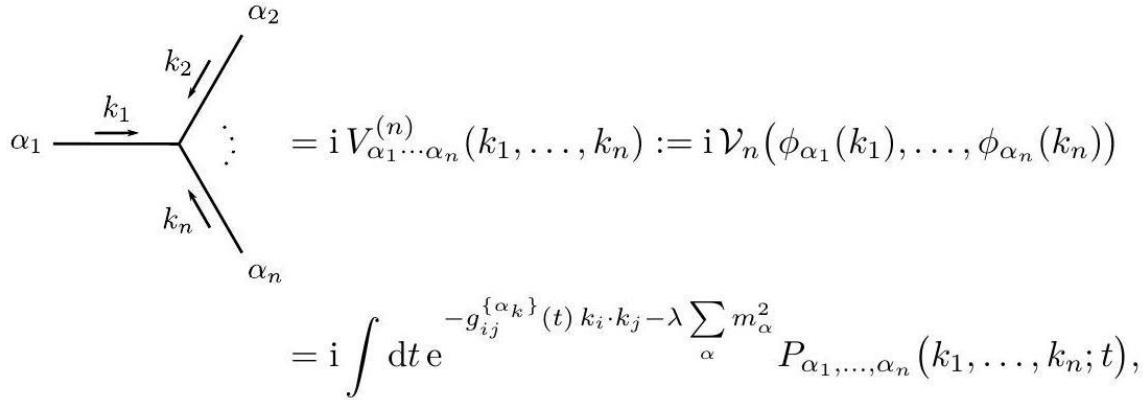
Momentum en dimensión \mathbb{R}^4 – espacio SFT.

$$|\Psi\rangle=\sum_j\int\frac{\mathrm{d}^Dk}{(2\pi)^D}\psi_\alpha(k)|\phi_\alpha(k)\rangle$$



$$S = - \int d^D k \psi_\alpha(k) K_{\alpha\beta}(k) \psi_\beta(-k) - \sum_{n \geq 0} \int d^D k_1 \cdots d^D k_n V_{\alpha_1 \cdots \alpha_n}^{(n)}(k_1, \dots, k_n) \psi_{\alpha_1}(k_1) \cdots \psi_{\alpha_n}(k_n)$$

$$\alpha - \vec{k}\beta = K_{\alpha\beta}(k)^{-1} = \frac{-iM_{\alpha\beta}}{k^2 + m_\alpha^2} Q_\alpha(k),$$



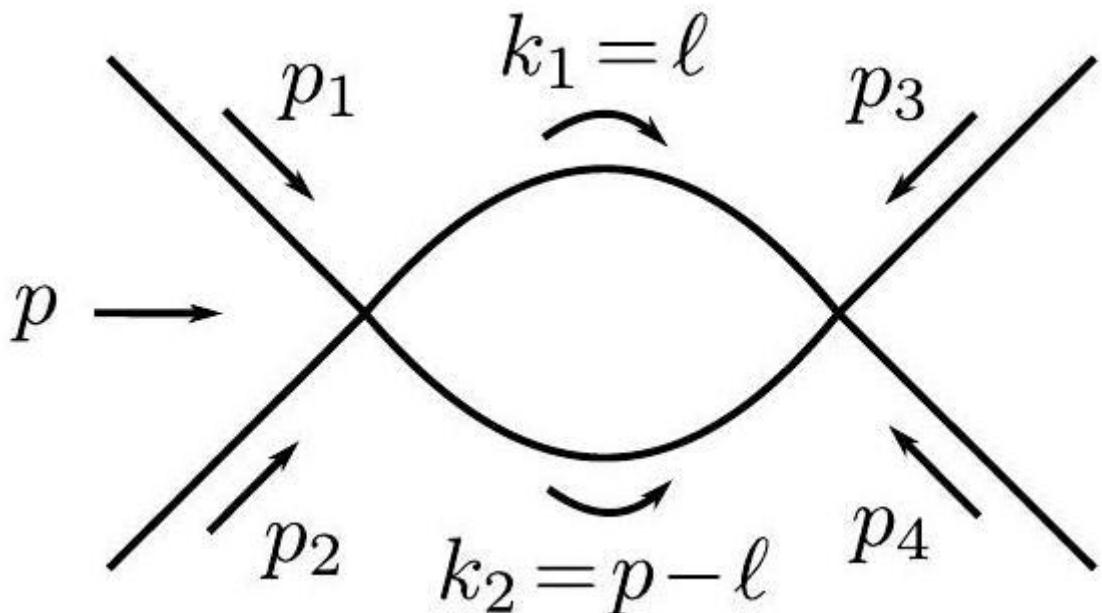
$$\begin{aligned} &= i V_{\alpha_1 \cdots \alpha_n}^{(n)}(k_1, \dots, k_n) := i \mathcal{V}_n(\phi_{\alpha_1}(k_1), \dots, \phi_{\alpha_n}(k_n)) \\ &= i \int dt e^{-g_{ij}^{\{\alpha_k\}}(t) k_i \cdot k_j - \lambda \sum_{\alpha} m_{\alpha}^2} P_{\alpha_1, \dots, \alpha_n}(k_1, \dots, k_n; t), \end{aligned}$$

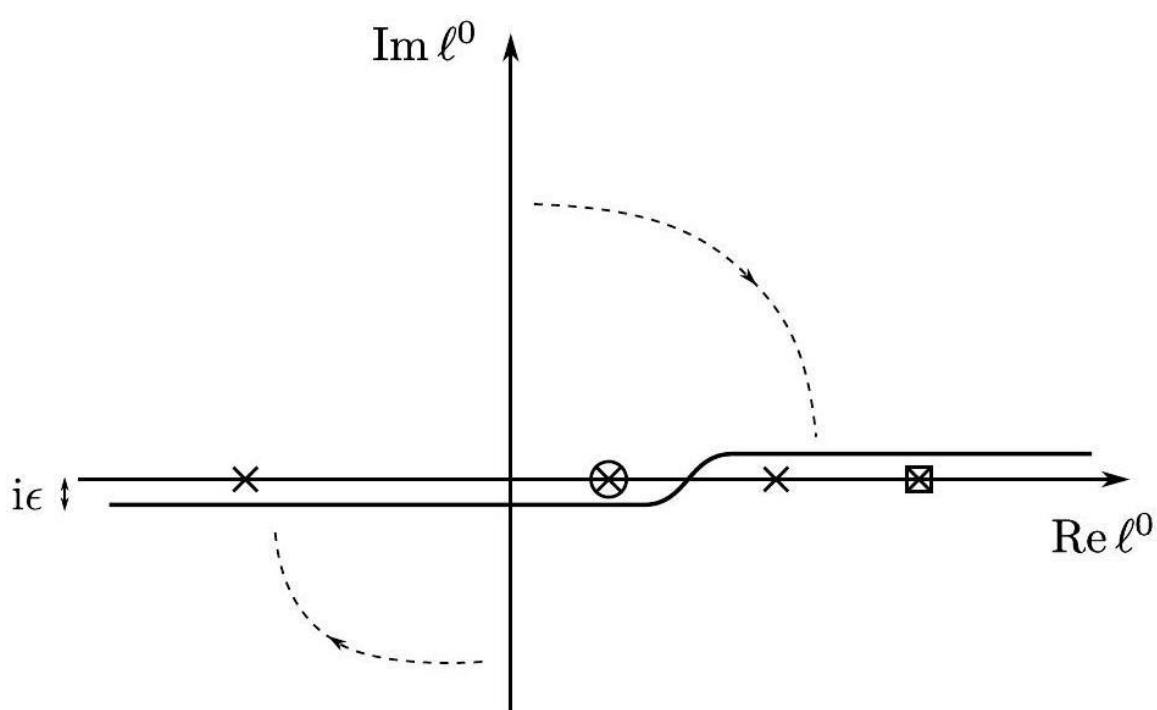
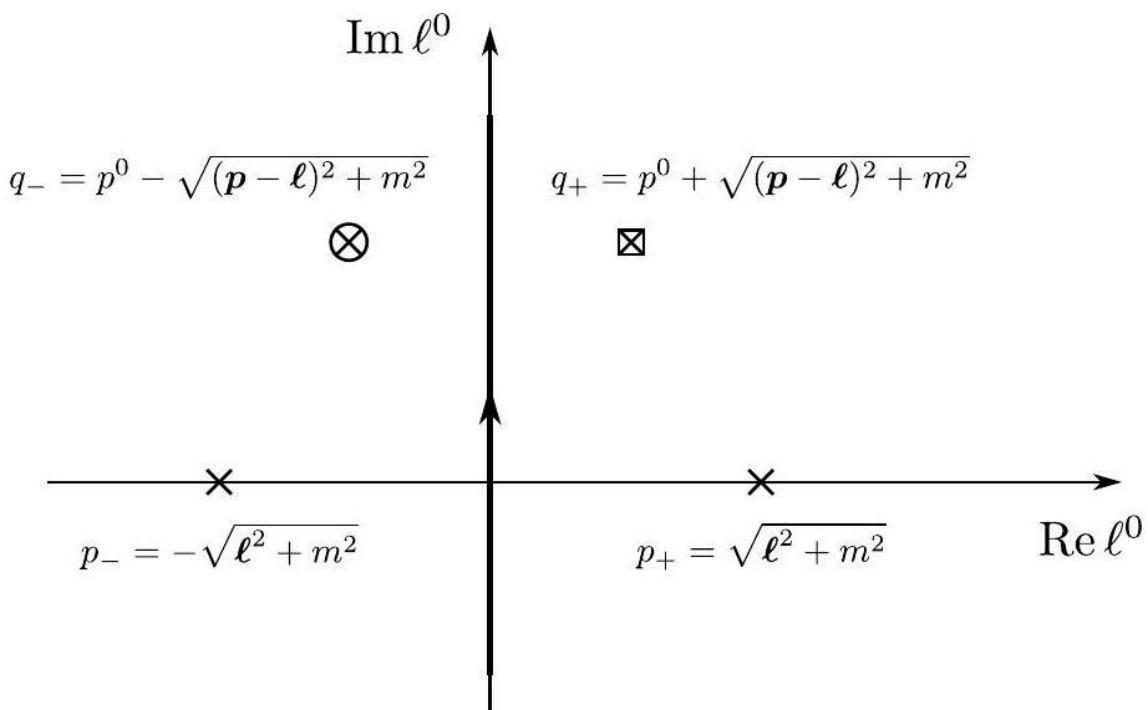
$$\lim_{k^0 \rightarrow \pm i\infty} V^{(n)} = 0, \lim_{k^0 \rightarrow \pm \infty} V^{(n)} = \infty$$

$$\begin{aligned} F_{g,n}(p_1, \dots, p_n) \sim & \int dT \prod_s d^D \ell_s e^{-G_{rs}(T) \ell_r \cdot \ell_s} - 2H_{ri}(T) \ell_r \cdot p_i - F_{ij}(T) p_i \cdot p_j \\ & \times \prod_a \frac{1}{k_a^2 + m_a^2} \mathcal{P}(p_i, \ell_r; T) \end{aligned}$$

Rotación de Wick.

$$p_{\pm} = \pm \sqrt{\ell^2 + m^2}, q_{\pm} = p^0 \pm \sqrt{(\mathbf{p} - \ell)^2 + m^2}.$$





Figuras 31, 32 y 33. Coordenadas de una partícula supermasiva en aniquilación.

Sistema de Coordenadas.

$$x^\mu = (x^0, x^i), \mu = 0, \dots, D-1 = d \quad i = 1, \dots, d$$

$$\sigma^\alpha = (\sigma^0, \sigma^\alpha), \alpha = 0, \dots, p-1, \alpha = 1, \dots, p.$$

$$\eta_{\mu\nu} = \text{diag}(-1, \underbrace{1, \dots, 1}_d).$$



$$\delta_{\mu\nu}=\mathrm{diag}(\underbrace{1,\ldots,1}_D).$$

$$\epsilon_{01}=-\epsilon^{01}=1$$

$$t = - \mathrm{i} \tau.$$

$$V_M^0=-\mathrm{i} V_E^0,V_{M,0}=\mathrm{i} V_{E,0}.$$

$$x^\pm=x^0\pm x^1.$$

$$\tau\in\mathbb{R},\sigma\in[0,L),\sigma\sim\sigma+L,$$

$$\mathcal{L}=\frac{1}{2\pi}\int_0^L\,\mathrm{d}\sigma=\frac{L}{2\pi}$$

$$w=\tau+\mathrm{i}\sigma,\bar w=\tau-\mathrm{i}\sigma$$

$$\mathrm{d}s^2=\mathrm{d}\tau^2+\mathrm{d}\sigma^2=\mathrm{d}w\;\mathrm{d}\bar{w}$$

$$w=\mathrm{i}\sigma^+,\bar w=\mathrm{i}\sigma^-.$$

$$z=\mathrm{e}^{2\pi w/L},\bar z=\mathrm{e}^{2\pi \bar w/L}.$$

$$\epsilon_{z\bar z}=\frac{\mathrm{i}}{2},\epsilon^{z\bar z}=-2\mathrm{i}$$

Sistema de Operadores.

$$[A,B]:=[A,B]_-=AB-BA,\{A,B\}:=[A,B]_+=AB+BA.$$

$$|A|=\begin{cases} +1 & \text{Grassmann odd}\\ 0 & \text{Grassmann even}\end{cases}$$

$$AB=(-1)^{|A||B|}BA.$$

Sistema QFT.

$$p^\mu := \left(E,p^i\right)$$

$$(N_L,N_R), N=N_L+N_R$$

$$\delta\phi(x)=\phi'(x)-\phi(x).$$

$$J_a^\mu=\lambda\frac{\partial\mathcal{L}}{\partial\left(\partial_\mu\phi\right)}\frac{\delta\phi}{\delta\alpha^a},\nabla_\mu J_a^\mu=0$$

$$Q_a=\frac{1}{\lambda}\oint_{\Sigma}\mathrm{d}^{D-1}x\sqrt{h}J_a^0$$

$$\delta_{\alpha^a}\phi(x)=\mathrm{i}\alpha^a[Q_a,\phi(x)]$$



$$J_a^\mu = i\lambda \frac{\partial \mathcal{L}}{\partial(\partial_\mu \phi)} \frac{\delta \phi}{\delta \alpha^a}$$

$$\delta_{\alpha^a} \phi(x) = -\alpha^a [Q_a, \phi(x)]$$

Espacio curvo y gravedad.

$$\nabla_\mu = \partial_\mu + \Gamma_\mu$$

$$\nabla_\mu A^\nu = \partial_\mu A^\nu + \Gamma_{\mu\rho}^\nu A^\rho$$

$$\Delta = g^{\mu\nu} \nabla_\mu \nabla_\nu = \frac{1}{\sqrt{g}} \nabla_\mu (\sqrt{g} g^{\mu\nu} \nabla_\nu)$$

$$T_{\mu\nu} = -\frac{2\lambda}{\sqrt{g}} \frac{\delta S}{\delta g^{\mu\nu}}$$

Análisis Complejo - Cauchy-Riemann.

$$\oint_{C_z} \frac{dw}{2\pi i} \frac{f(w)}{(w-z)^n} = \frac{f^{(n-1)}(z)}{(n-1)!},$$

$$\bar{\partial} \frac{1}{z} = 2\pi \delta^{(2)}(z)$$

QFT, espacios curvos y gravedad. Función de Green.

$$D_x G(x, y) = \frac{\delta(x-y)}{\sqrt{g}} - P(x, y),$$

$$\nabla_\mu v^\mu = \frac{1}{\sqrt{g}} \partial_\mu (\sqrt{g} v^\mu)$$

$$\delta x^\mu = \xi^\mu,$$

$$\delta g_{\mu\nu} = \mathcal{L}_\xi g_{\mu\nu} = \nabla_\mu \xi_\nu + \nabla_\nu \xi_\mu$$

$$\int_V d^D x \nabla_\mu v^\mu = \oint_{\partial V} d\Sigma_\mu v^\mu, d\Sigma_\mu := \epsilon n_\mu d^{D-1}\Sigma,$$

$$d^{D-1}\Sigma = \sqrt{g} d^{D-1}x, n_\mu = \delta_\mu^0$$

$$Q_S = \frac{1}{\lambda} \int_S d\Sigma_\mu J_a^\mu$$

Dimensiones D. Teorema de Stokes.

$$\int d^2x \partial_\mu v^\mu = \oint \epsilon_{\mu\nu} dx^\nu v^\mu = \oint (v^0 d\sigma - v^1 d\tau)$$



$$\begin{aligned}\chi_{g;b} &:= \frac{1}{4\pi} \int d^2\sigma \sqrt{g} R + \frac{1}{2\pi} \oint ds k \\ &= 2 - 2g - b\end{aligned}$$

$$T_{ab} = -\frac{4\pi}{\sqrt{g}} \frac{\delta S}{\delta g^{ab}}$$

Plano Complejo.

$$z = x + iy, z = x - iy$$

$$\begin{aligned}ds^2 &= dx^2 + dy^2 = dz d\bar{z}, g_{z\bar{z}} = \frac{1}{2}, g_{zz} = g_{\bar{z}\bar{z}} = 0 \\ \epsilon_{z\bar{z}} &= \frac{i}{2}, \epsilon^{z\bar{z}} = -2i \\ \partial_z &:= \partial_x = \frac{1}{2}(\partial_x - i\partial_y), \bar{\partial}_z := \partial_{\bar{x}} = \frac{1}{2}(\partial_x + i\partial_y) \\ V^z &= V^x + iV^y, V^{\bar{z}} = V^x - iV^y \\ d^2x &= dx dy = \frac{1}{2} d^2z, d^2z = dz d\bar{z} \\ \delta(z) &= \frac{1}{2} \delta^{(2)}(x), 1 = \int d^2z \delta^{(2)}(z) = \int d^2x \delta^{(2)}(x) \\ \int_R d^2z (\partial_z v^z + \partial_{\bar{z}} v^{\bar{z}}) &= -i \oint_{\partial R} (dz v^{\bar{z}} - d\bar{z} v^z) = -2i \oint_{\partial R} (v_z dz - v_{\bar{z}} d\bar{z})\end{aligned}$$

Propiedades generales.

$$f \circ \phi(z) = \left(\frac{df}{dz} \right)^h \phi(f(z))$$

$$\phi(z) = \sum_n \frac{\phi_n}{z^{n+h}}, \phi_n = \oint_{c_0} \frac{dz}{2\pi i} z^{n+h-1} \phi(z),$$

$$\forall n \geq -h+1: \phi_n |0\rangle = 0$$

$$\forall n \leq h-1: \langle 0 | \phi_n = 0.$$

$$|\phi\rangle := \phi(0)|0\rangle = \phi_{-h}|0\rangle.$$

$$\langle \phi^\pm | := \langle 0 | I \circ \phi^\pm(0) = \lim_{z \rightarrow \infty} z^{2h} \langle 0 | \phi^\pm(z), \langle \phi | := \langle 0 | I^\pm \circ \phi(0) = (\pm 1)^h \lim_{z \rightarrow \infty} z^{2h} \langle 0 | \phi(z).$$

$$T(z) = \sum_n \frac{L_n}{z^{n+2}}.$$

$$T(z)\phi(w) \sim \frac{h\phi(w)}{(z-w)^2} + \frac{\partial\phi(w)}{z-w}$$

$$T(z)T(w) \sim \frac{c/2}{(z-w)^4} + \frac{2T(w)}{(z-w)^2} + \frac{\partial T(w)}{z-w}$$



Conjugaciones Hermitianas y BPZ.

Hermitiana.

$$(\lambda A_1 \cdots A_n | 0 \rangle)^\dagger = \lambda^* \langle 0 | A_n^\dagger \cdots A_1^\dagger.$$

BPZ.

$$\phi_n^t = (I^\pm \circ \phi)_n = (-1)^h (\pm 1)^n \phi_{-n}$$

$$(\lambda A_1 \cdots A_n | 0 \rangle)^t = \lambda \langle 0 | (A_1)^t \cdots (A_n)^t.$$

$$\langle A, B \rangle = (-1)^{|A||B|} \langle B, A \rangle.$$

$$\forall A: \langle A | B \rangle = 0 \Rightarrow |B\rangle = 0.$$

$$\langle \phi_r^c | \phi_s \rangle = \delta_{rs}$$

$$\langle \phi_r | \phi_s^c \rangle = (-1)^{|\phi_r|} \delta_{rs}.$$

Campo escalar.

$$i\partial X^\mu = \sum_n \frac{\alpha_n^\mu}{z^{n+1}}, i\bar{\partial} X^\mu = \sum_n \frac{\bar{\alpha}_n^\mu}{\bar{z}^{n+1}}$$

$$[\alpha_m^\mu, \alpha_n^\nu] = m\delta_{m+n,0}\eta^{\mu\nu}, [\bar{\alpha}_m^\mu, \bar{\alpha}_n^\nu] = m\delta_{m+n,0}\eta^{\mu\nu}, [\alpha_m^\mu, \bar{\alpha}_n^\nu] = 0.$$

$$\alpha_0^\mu = \bar{\alpha}_0^\mu = \sqrt{\frac{\alpha'}{2}} p^\mu.$$

$$[x^\mu, p^\nu] = \eta^{\mu\nu}.$$

$$V_k(z, \bar{z}) = : e^{ik \cdot X(z, \bar{z})} :, h = \bar{h} = \frac{\alpha'^2 k^2}{4}.$$

$$p^\mu |k\rangle = k^\mu |k\rangle, \forall n > 0: \alpha_n^\mu |k\rangle = 0, \bar{\alpha}_n^\mu |k\rangle = 0.$$

$$|k\rangle = V_k(0,0)|0\rangle = e^{ik \cdot x}|0\rangle.$$

$$\langle k | p^\mu = \langle k | k^\mu, \langle k | = |k\rangle^\dagger, \langle -k | = |k\rangle^t.$$

Reparametrizaciones fantasma.

$$\epsilon = 1, \lambda = 2, c_{\text{gh}} = -26, q_{\text{gh}} = -3, a_{\text{gh}} = -1.$$

$$h(b) = 2, h(c) = -1$$



$$b(z)=\sum_{n\in\mathbb{Z}}\frac{b_n}{z^{n+2}}, c(z)=\sum_{n\in\mathbb{Z}}\frac{c_n}{z^{n-1}},$$

$$b_n=\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}z^{n+1}b(z), c_n=\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}z^{n-2}c(z).$$

$$\{b_m,c_n\}=\delta_{m+n,0}, \{b_m,b_n\}=0, \{c_m,c_n\}=0.$$

$$T=-2:b\partial c:-\partial bc:\\ L_m=\sum_n(n+m):b_{m-n}c_n:=\sum_n(2m-n):b_nc_{m-n}:\\$$

$$L_0=-\sum_nn:b_nc_{-n}:=\sum_nn:b_{-n}c_n:$$

$$[L_m,b_n]=(m-n)b_{m+n}, [L_m,c_n]=-(2m+n)c_{m+n}$$

$$[L_0,b_0]=0, [L_0,c_0]=0$$

$$j=-:bc:, N_{\text{gh},L}=\oint\frac{\mathrm{d}z}{2\pi\mathrm{i}}j(z)$$

$$N_{\text{gh}}(c)=1, N_{\text{gh}}(b)=-1.$$

$$j_m=-\sum_n:b_{m-n}c_n:=-\sum_n:b_nc_{m-n}:, N_{\text{gh},L}=j_0=-\sum_n:b_{-n}c_n:$$

$$[j_m,j_n]=m\delta_{m+n,0}, [L_m,j_n]=-nj_{m+n}-\frac{3}{2}m(m+1)\delta_{m+n,0}$$

$$\left[N_{\text{gh}}, b(w)\right]=-b(w), \left[N_{\text{gh}}, c(w)\right]=c(w)$$

$$N^b=\sum_{n>0}nN_n^b, N^c=\sum_{n>0}nN_n^c,$$

$$N_n^b=:b_{-n}c_n:, N_n^c=:c_{-n}b_n:$$

$$[N_m^b,b_{-n}]=b_{-n}\delta_{m,n}, [N_m^c,c_{-n}]=c_{-n}\delta_{m,n}$$

$$c(z)b(w)\sim\frac{1}{z-w}, b(z)c(w)\sim\frac{1}{z-w}, b(z)b(w)\sim 0, c(z)c(w)\sim 0,$$

$$T(z)b(w)\sim\frac{2b(w)}{(z-w)^2}+\frac{\partial b(w)}{z-w}, T(z)c(w)\sim\frac{-c(w)}{(z-w)^2}+\frac{\partial c(w)}{z-w}.$$

$$j(z)b(w)\sim-\frac{b(w)}{z-w}, j(z)c(w)\sim\frac{c(w)}{z-w}. j(z)\mathcal{O}(w)\sim N_{\text{gh}}(\mathcal{O})\frac{\mathcal{O}(w)}{z-w},$$

$$j(z)j(w)\sim\frac{1}{(z-w)^2}.$$

$$T(z)j(w)\sim\frac{-3}{(z-w)^3}+\frac{j(w)}{(z-w)^2}+\frac{\partial j(w)}{z-w}.$$

$$N^c-N^b=3-3g$$



$$N_{\text{gh},L}=N_{\text{gh},L}^{\text{cyl}}+\frac{3}{2}$$

$$\forall n>-2 \colon b_n|0\rangle=0, \forall n>1 \colon c_n|0\rangle=0$$

$$|\downarrow\rangle = c_1|0\rangle, |\uparrow\rangle = c_0c_1|0\rangle.$$

$$L_0|\downarrow\rangle=a_{\text{gh}}|\downarrow\rangle,L_0|\uparrow\rangle=a_{\text{gh}}|\uparrow\rangle,a_{\text{gh}}=-1.$$

$$\begin{aligned} L_m &= \sum_n \ (n - (1-\lambda)m)^{*} b_{m-n} {c_n}^{*} + a_{\text{gh}} \delta_{m,0}, \\ j_m &= \sum_n \ {^{*}b_{m-n}} {c_n}^{*} + \delta_{m,0}. \end{aligned}$$

$$\begin{aligned} L_0 &= \sum_n \ n_{\star}^{*} b_{-n} {c_n}_{\star}^{*} + a_{\text{gh}} = \hat{L}_0 - 1, \\ N_{\text{gh},L} = j_0 &= \sum_n^{*} b_{-n} {c_n}^{*} + 1 = \hat{N}_{\text{gh},L} + \frac{1}{2} \left(N_0^c - N_0^b \right) - \frac{3}{2}, \\ \hat{L}_0 = N^b + N^c, \hat{N}_{\text{gh},L} &:= \sum_{n>0} \left(N_n^c - N_n^b \right). \end{aligned}$$

$$N_{\text{gh}}|0\rangle=0,N_{\text{gh}}|\downarrow\rangle=|\downarrow\rangle,N_{\text{gh}}|\uparrow\rangle=2|\uparrow\rangle.$$

$$N_{\text{gh}}^{\text{cyl}}|\downarrow\rangle=-\frac{1}{2}|\downarrow\rangle,N_{\text{gh}}^{\text{cyl}}|\uparrow\rangle=\frac{1}{2}|\uparrow\rangle.$$

$$b_n^\dagger=b_{-n}, c_n^\dagger=c_{-n}$$

$$b_n^t=(\pm 1)^nb_{-n}, c_n^t=-(\pm 1)^nc_{-n}$$

$$|\downarrow\rangle^\ddag=\langle 0|c_{-1}, |\uparrow\rangle^\ddag=\langle 0|c_{-1}c_0.$$

$$\langle \downarrow | := |\downarrow\rangle^t = \mp \langle 0|c_{-1}, \langle \uparrow | := |\uparrow\rangle^t = \pm \langle 0|c_0c_{-1}$$

$$\langle \downarrow | = \mp |\downarrow\rangle^\ddag, \langle \uparrow | = \mp |\uparrow\rangle^\ddag.$$

$$\langle \uparrow|\downarrow\rangle=\langle \downarrow|c_0|\downarrow\rangle=\langle 0|c_{-1}c_0c_1|0\rangle=1,$$

$$\langle 0^c|=\langle 0|c_{-1}c_0c_1$$

$$b_n^\pm=b_n\pm\bar{b}_n, c_n^\pm=\frac{1}{2}(c_n\pm\bar{c}_n)$$

$$\{b_m^+, c_n^+\}=\delta_{m+n}, \{b_m^-, c_n^-\}=\delta_{m+n}$$

$$b_n^-b_n^+=2b_n\bar{b}_n, c_n^-c_n^+=\frac{1}{2}c_n\bar{c}_n$$



Campo gravitónico. Operadores, propagadores, funciones fantasma, supersimetrías, antisimetría y supermembranas y sistema de coordenadas.

$$(\alpha_n)^t = -(\pm 1)^n \alpha_{-n}, (b_n)^t = (\pm 1)^n b_{-n}, (c_n)^t = -(\pm 1)^n c_{-n}$$

$$L_n^\pm = L_n \pm \bar{L}_n, b_n^\pm = b_n \pm \bar{b}_n, c_n^\pm = \frac{1}{2}(c_n \pm \bar{c}_n)$$

$$\langle A, B \rangle = \langle A | c_0^- | B \rangle,$$

$$\langle A, B \rangle = \langle A | B \rangle.$$

$$|k, 0\rangle := |k\rangle \otimes |0\rangle, |k, \downarrow\rangle := |k\rangle \otimes |\downarrow\rangle.$$

$$\langle k, \downarrow | c_0 | k, \downarrow \rangle = \langle k', 0 | c_{-1} c_0 c_1 | k, 0 \rangle = (2\pi)^D \delta^{(D)}(k + k')$$

$$\langle k, \downarrow \downarrow | c_0 \bar{c}_0 | k, \downarrow \downarrow \rangle = \langle k', 0 | c_{-1} \bar{c}_{-1} c_0 \bar{c}_0 c_1 \bar{c}_1 | k, 0 \rangle = (2\pi)^D \delta^{(D)}(k + k')$$

$$|\delta\Phi|^2 = G(\Phi)(\delta\Phi, \delta\Phi),$$

$$|\Phi|^2 = G(\Phi)(\Phi, \Phi),$$

$$d\Phi \sqrt{\det G(\Phi)}.$$

$$(\delta\Phi_1, \delta\Phi_2) = G(\Phi)(\delta\Phi_1, \delta\Phi_2), (\Phi_1, \Phi_2) = G(\Phi)(\Phi_1, \Phi_2).$$

$$|\delta\Phi_a|^2 = \int dx \rho(x) \gamma_{ab}(\Phi(x)) \delta\Phi_a(x) \delta\Phi_b(x),$$

$$G_{ab}(x, y)(\Phi) = \delta(x - y) \rho(x) \gamma_{ab}(\Phi(x)).$$

$$\int d\delta\Phi e^{-G(\Phi)(\delta\Phi, \delta\Phi)} = \frac{1}{\sqrt{\det G(\Phi)}}$$

$$\int d\Phi e^{-G(\Phi)(\Phi, \Phi)}$$

$$\int d\Phi \sqrt{\det G(\Phi)} F(\Phi)$$

$$S = -\frac{1}{2} \text{tr} \ln G(\Phi)$$

$$Z = \int d\Phi e^{-S_{cl}(\Phi)}$$

$$G(\delta\Phi, D\delta\Phi) = G(D^\dagger \delta\Phi, \delta\Phi)$$

$$\Phi \rightarrow \Phi + \varepsilon$$



$$\int \mathrm{d}\Phi \mathrm{e}^{-\frac{1}{2}|\Phi+\varepsilon|^2}=\int \mathrm{d}\Phi \mathrm{e}^{-\frac{1}{2}|\Phi|^2}$$

$$\int \mathrm{d}\Phi \mathrm{e}^{-\frac{1}{2}|\Phi+\varepsilon|^2}=\int \mathrm{d}\widetilde{\Phi} \mathrm{det}\frac{\delta \Phi}{\delta \widetilde{\Phi}} \mathrm{e}^{-\frac{1}{2}|\widetilde{\Phi}|^2}=\int \mathrm{d}\widetilde{\Phi} \mathrm{e}^{-\frac{1}{2}|\widetilde{\Phi}|^2}$$

$$\mathrm{d}\Phi\sqrt{\det G(\Phi)}=\mathrm{d}\widetilde{\Phi}\sqrt{\det \widetilde{G}(\widetilde{\Phi})}, G(\Phi)(\delta\Phi,\delta\Phi)=\widetilde{G}(\widetilde{\Phi})(\delta\widetilde{\Phi},\delta\widetilde{\Phi})$$

$$\mathrm{d}\Phi=J(\Phi,\widetilde{\Phi})\mathrm{d}\widetilde{\Phi}, J(\Phi,\widetilde{\Phi})=\left|\mathrm{det}\frac{\partial \Phi}{\partial \widetilde{\Phi}}\right|=\sqrt{\frac{\det \widetilde{G}(\widetilde{\Phi})}{\det G(\Phi)}}$$

$$\int \mathrm{d}\delta\Phi \mathrm{e}^{-|\delta\Phi|^2}=1$$

$$J(\widetilde{\Phi})^{-1}=\int \mathrm{d}\delta\widetilde{\Phi} \mathrm{e}^{-\tilde{G}(\delta\widetilde{\Phi},\delta\widetilde{\Phi})}$$

$$J=\mathrm{det}\frac{\partial \tilde{x}^\mu}{\partial x^\mu}=\mathrm{det}\frac{\partial \tilde{v}^\mu}{\partial v^\mu}$$

$$\tilde{v}^\mu = v^\nu \frac{\partial \tilde{x}^\mu}{\partial x^\nu}$$

$$D\Phi_0=0$$

$$\Phi=\Phi_0+\Phi', (\Phi_0,\Phi')=0$$

$$Z[D]=\int \mathrm{d}\Phi \sqrt{\det G} \mathrm{e}^{-\frac{1}{2}(\Phi,D\Phi)}=\left(\int \mathrm{d}\Phi_0\right)\int \mathrm{d}\Phi' \mathrm{e}^{-\frac{1}{2}(\Phi',D\Phi')}$$

$$Z[D,J]=\int \mathrm{d}\Phi \sqrt{\det G} \mathrm{e}^{-\frac{1}{2}(\Phi,D\Phi)-(J,\Phi)}$$

$$Z[D,J]=\int \mathrm{d}\Phi_0 \mathrm{e}^{-(J,\Phi_0)}\int \mathrm{d}\Phi' \mathrm{e}^{-\frac{1}{2}(\Phi',D\Phi')-(J,\Phi')}$$

$$\int \mathrm{d}x=\infty$$

$$\int \mathrm{d}\theta=0$$

$$\int \mathrm{d}\theta\theta=\int \mathrm{d}\theta\delta(\theta)=1$$

$$\int \mathrm{d}x\delta(x)=1$$

$$\theta_0(x)=\theta_{0i}\psi_i(x), \ker D=\mathrm{Span}\{\psi_i\}$$



$$d\theta = \frac{1}{\sqrt{\det(\psi_i, \psi_j)}} d\theta' \prod_{i=1}^n d\theta_{0i}$$

$$d\theta \prod_{i=1}^n \theta(x_i) = \frac{\det \psi_i(x_j)}{\sqrt{\det(\psi_i, \psi_j)}} d\theta'$$

$$\begin{aligned} 1 &= \int d\theta e^{-|\theta|^2} = \int d\theta' d\theta_0 e^{-|\theta|^2 - |\theta_0|^2} \\ &= J \int d\theta' \prod_i d\theta_{0i} e^{-|\theta'|^2 - |\theta_{0i}\psi_i|^2} = J \sqrt{\det(\psi_i, \psi_j)} \end{aligned}$$

$$\begin{aligned} \int d\theta_0 \prod_{j=1}^n \theta(x_j) &= \int d\theta_0 \prod_{j=1}^n \theta_0(x_j) = \frac{1}{\sqrt{\det(\psi_i, \psi_j)}} \int d^n \theta_{0i} \prod_{j=1}^n [\theta_{0i} \psi_i(x_j)] \\ &= \frac{\det \psi_i(x_j)}{\sqrt{\det(\psi_i, \psi_j)}} \int \prod_i d\theta_{0i} \theta_{0i} = \frac{\det \psi_i(x_j)}{\sqrt{\det(\psi_i, \psi_j)}} \end{aligned}$$

Cuantización BRST.

$$\delta\phi^i = \epsilon^a \delta_a \phi^i = \epsilon^a R_a^i(\phi)$$

$$\epsilon^a \delta_a S_m = 0$$

$$[\delta_a, \delta_b] = f_{ab}^c \delta_c$$

$$R_a^i(\phi) = (T_a^{\mathbf{R}})^i{}_j \phi^j,$$

$$R_{a\mu}^b = \delta_a^b \partial_\mu + f_{ab}^c A_\mu^b$$

$$Z = \Omega_{\text{gauge}}^{-1} \int d\phi^i e^{-S_m}$$

$$F^A(\phi^i) = 0$$

$$\begin{aligned} S_{\text{gh}} &= b_A c^a \delta_a F^A(\phi^i), \\ S_{\text{gf}} &= -i B_A F^A(\phi^i) \end{aligned}$$

$$Z = \int d\phi^i db_A dc^a dB_A e^{-S_{\text{tot}}}$$

$$S_{\text{tot}} = S_m + S_{\text{gf}} + S_{\text{gh}}$$

$$\delta_\epsilon S_{\text{tot}} = 0$$

$$\delta_\epsilon \phi^i = i \epsilon c^a \delta_a \phi^i, \delta_\epsilon c^a = -\frac{i}{2} \epsilon f_{bc}^a c^b c^c, \delta_\epsilon b_A = \epsilon B_A, \delta_\epsilon B_A = 0$$



$$\delta_\epsilon \delta_{\epsilon'} = 0$$

$$\delta_\epsilon \phi^i = \mathrm{i} [\epsilon Q_B, \phi^i]$$

$$S_{\text{gf}}+S_{\text{gh}}=\{Q_B,b_A F^A\}$$

$$\delta S=\{Q_B,b_A\delta F^A\}$$

$$[Q_B,\{Q_B,b_A\delta F^A\}]=0\implies Q_B^2=0$$

$$|\psi\rangle \text{ closed} \iff |\psi\rangle \in \ker Q_B \iff Q_B|\psi\rangle = 0.$$

$$|\psi\rangle \text{ exact} \iff |\psi\rangle \in \text{Im} Q_B \iff \exists |\chi\rangle : |\psi\rangle = Q_B|\chi\rangle.$$

$$|\psi\rangle \in \mathcal{H}(Q_B) \iff |\psi\rangle \in \ker Q_B, \nexists |\chi\rangle : |\psi\rangle = Q_B|\chi\rangle.$$

$$\mathcal{H}(Q_B)=\frac{\ker Q_B}{\text{Im} Q_B}$$

$$|\psi\rangle\simeq|\psi\rangle+Q_B|\chi\rangle.$$

$$\delta_F\langle\psi_f\mid\psi_i\rangle=\langle\psi_f\big|\{Q_B,b_A\delta F^A\}|\psi_i\rangle$$

$$Q_B|\psi\rangle=0$$

$$\langle\psi|Q_B|\chi\rangle=0$$

$$|\psi\rangle \text{ physical} \iff |\psi\rangle \in \mathcal{H}(Q_B).$$

$$\frac{\delta F^A}{\delta \phi^i}B_A=-\frac{\delta S_m}{\delta \phi^i}$$

$$\delta_\epsilon b_A=-\epsilon\left(\frac{\delta F^A}{\delta \phi^i}\right)^{-1}\frac{\delta S_m}{\delta \phi^i}$$

$$\{Q_B,b_AB_BM^{AB}\}=\mathrm{i} B_AM^{AB}B_B$$

$$S[\phi,b,c,B] = S_0[\phi] + Q_B\Psi[\phi,b,c,B]$$

Formalismo Batalin-Vilkovisky.

$$\mathcal{F}_i(\phi)=\frac{\partial S_0}{\partial \phi^i}$$

$$[T_a,T_b]=F_{ab}^c(\phi)T_c+\lambda_{ab}^i\mathcal{F}_i(\phi)$$

$$m_1=m_0-\text{rank}R_a^i$$

$$\delta A_p=\mathrm{d}\lambda_{p-1}$$

$$\delta\lambda_{p-1}=\mathrm{d}\lambda_{p-2}$$



$$\delta\lambda_{p-2}=\mathrm{d}\lambda_{p-3}$$

$$\psi^r = \{\phi^i, B_A, b_A, c^a\}$$

$$S[\psi^r,\psi_r^*]=S_0[\phi]+Q_B\psi^r\psi_r^*$$

$$\psi_r^*=\{\phi_i^*,B^{A*},b^{A*},c_a^*\}$$

$$\psi_r^*=\frac{\partial\Psi}{\partial\psi^r}$$

$$\delta\phi^i=\epsilon_0^{a_0}R_{a_0}^i(\phi^i)$$

$$\delta c_0^{a_0}=\epsilon_1^{a_1}R_{a_1}^{a_0}(\phi^i,c^{a_0})$$

$$\delta c_n^{a_n}=\epsilon_{n+1}^{a_{n+1}}R_{a_{n+1}}^{a_n}(\phi^i,c_0^{a_0},\ldots,c_n^{a_n})$$

$$\psi^r=\left\{c_n^{a_n}\right\}_{n=-1,\dots,\ell}, c_{-1}\colon=\phi$$

$$N_{\mathrm{gh}}(\phi^i)=0,N_{\mathrm{gh}}(c_n^{a_n})=n+1$$

$$|c_n|=\left|\epsilon_n^{a_n}\right|+n+1$$

$$N_{\mathrm{gh}}(\psi_r^*)=-1-N_{\mathrm{gh}}(\psi^r), |\psi_r^*|=-|\psi^r|.$$

$$\omega=\sum_r\;\; \mathrm{d}\psi^r\wedge\;\mathrm{d}\psi_r^*$$

$$(\psi^r,\psi_s^*)=\delta_{rs}, (\psi^r,\psi^s)=0, (\psi_r^*,\psi_s^*)=0$$

$$(A,B) = \frac{\partial_R A}{\partial \psi^r} \frac{\partial_L B}{\partial \psi_r^*} - \frac{\partial_R A}{\partial \psi_r^*} \frac{\partial_L B}{\partial \psi^r}$$

$$(A,B) = -(-1)^{(|A|+1)(|B|+1)}(B,A)$$

$$N_{\mathrm{gh}}((A,B))=N_{\mathrm{gh}}(A)+N_{\mathrm{gh}}(B)+1, |(A,B)|=|A|+|B|+1 \bmod 2$$

$$(A,BC)=(A,B)C+(-1)^{|B|C}(A,C)B$$

$$N_{\mathrm{gh}}(S)=0, |S|=0$$

$$S[\psi^r,\psi_r^*=0]=S_0\big[\phi^i\big],\left.\frac{\partial_L\partial_RS}{\partial c_{n-1,a_{n-1}}^*\partial c_n^{a_n}}\right|_{\psi^*=0}=R_{a_n}^{a_{n-1}}$$

$$\begin{aligned}\delta_\theta\psi^r &= \theta s\psi^r = -\theta(S,\psi_r) = \theta\frac{\partial_RS}{\partial\psi_r^*} \\ \delta_\theta\psi_r^* &= \theta s\psi_r^* = -\theta(S,\psi_r^*) = -\theta\frac{\partial_RS}{\partial\psi^r}\end{aligned}$$

$$\delta_\theta F = \theta sF = -\theta(S,F)$$



$$(S,S)=0$$

$$(S,(S,F))=0$$

$$s^2 = 0$$

$$s\mathcal{O}=0$$

$$S'=S+(S,\delta F)$$

$$\psi'=\psi-\frac{\delta F}{\delta \psi^*}, \psi'^*=\psi^*+\frac{\delta F}{\delta \psi}$$

$$S'[\psi,\psi^*]=S\left[\psi-\frac{\delta F}{\delta \psi^*},\psi^*+\frac{\delta F}{\delta \psi}\right]$$

$$S'[\psi,\psi^*]=S[\psi,\psi^*]+(S,\psi)\frac{\delta F}{\delta \psi}+(S,\psi^*)\frac{\delta F}{\delta \psi^*}=S[\psi,\psi^*]-\frac{\partial_RS}{\partial \psi^*}\frac{\delta F}{\delta \psi}+\frac{\partial_RS}{\partial \psi}\frac{\delta F}{\delta \psi^*}$$

$$(\psi'^r,\psi'^*_s)=\delta_{rs}, (S',S')=0$$

$$G'=G+(\delta F,G)$$

$$S_\Psi[\psi^r]=S\left[\psi^r,\frac{\partial\Psi}{\partial\psi^r}\right],\psi^*_r=\frac{\partial\Psi}{\partial\psi^r}$$

$$N_{\mathrm{gh}}(\Psi)=-1, |\Psi|=1$$

$$|B|=-|\bar c|, N_{\mathrm{gh}}(B)=N_{\mathrm{gh}}(\bar c)+1,\\ s\bar c=B, sB=0$$

$$\bar S=S[\psi^r,\psi^*_r]-B\bar c^*$$

$$(\bar S,\bar S)=(S,S)=0$$

$$\left(B_{0a_0},\bar{c}_{0a_0}\right)\!:=\left(B_{0a_0}^0,\bar{c}_{0a_0}^0\right)$$

$$|B_0|=|\epsilon_0|, |\bar c_0|=-|\epsilon_0|, N_{\mathrm{gh}}(B_0)=0, N_{\mathrm{gh}}(c_0)=-1.$$

$$\left(B_{1a_1}^0,\bar{c}_{1a_1}^0\right),\left(\bar{B}_1^{1a_1},c_1^{1a_1}\right)$$

$$\delta_\theta \psi^r = \theta s \psi^r = \theta \left.\frac{\partial_RS}{\partial\psi^r}\right|_{\psi^*_r=\partial_r\Psi}$$

$$s^2 \propto \text{ eom.}$$

$${\bf Cu\'antica\;BV}.$$

$$Z=\int~\mathrm{d}\psi^r~\mathrm{d}\psi^*_r\mathrm{e}^{-W[\psi^r,\psi^*_r]/\hbar}$$

$$\delta_\theta F = \theta \sigma F = (W,F)-\hbar \Delta F$$



$$\Delta = \frac{\partial_R}{\partial \psi_r^*} \frac{\partial_L}{\partial \psi^r}$$

$$(W,W)-2\hbar\Delta W=0,$$

$$\Delta {\mathrm{e}}^{-W/\hbar}=0$$

$$\delta W=\frac{1}{2}(W,W),$$

$${\rm sdet} J \sim 1 + \Delta W$$

$$W=S+\sum_{p\geq 1}\,\hbar^pW_p$$

$$\sigma \mathcal{O}=0$$

$$\delta \langle \mathcal{O} \rangle = 0$$

$$Z=\int~{\mathrm{d}}\psi^r{\mathrm{e}}^{-W_\Psi[\psi^r]},W_\Psi[\psi^r]=W\left[\psi^r,\frac{\partial\Psi}{\partial\psi^r}\right]$$

$$Z=\int~{\mathrm{d}}\psi^r{\mathrm{e}}^{-W_\Psi[\psi^r]}\left(\frac{\partial_RS}{\partial\psi_r^*}\right)_{\psi^*=\partial_\psi\Psi}\frac{\partial(\delta\Psi)}{\partial\psi^r}$$

Gravedad cuántica endógena en planos cuánticos relativistas.

Cálculos preliminares.

$$|\dot{H}| \ll H^2,$$

$${\mathrm{d}} s^2=-{\mathrm{d}} t^2+a^2(t){\mathrm{d}} \boldsymbol{x}^2.$$

$$M_P\equiv\sqrt{\frac{\hbar c}{G}}=1.2\times10^{19}{\mathrm{GeV}}/{c^2},$$

$$M_{\mathrm{pl}}^{-2}\equiv 8\pi G=(2.4\times10^{18}{\mathrm{GeV}})^{-2}$$

$$\mathcal{R}_{\boldsymbol{k}}=\int~{\mathrm{d}}^3x\mathcal{R}(\boldsymbol{x})e^{i\boldsymbol{k}\cdot\boldsymbol{x}}.$$

$$\langle \mathcal{R}_{\boldsymbol{k}}\mathcal{R}_{\boldsymbol{k}'}\rangle=(2\pi)^3P_{\mathcal{R}}(k)\delta(\boldsymbol{k}+\boldsymbol{k}').$$

$$\Delta_{\mathcal{R}}^2(k)\equiv\frac{k^3}{2\pi^2}P_{\mathcal{R}}(k).$$

$$\varepsilon\equiv-\frac{\dot{H}}{H^2},\tilde{\eta}\equiv\frac{\dot{\varepsilon}}{H\varepsilon}$$

$$\epsilon\equiv\frac{M_{\mathrm{pl}}^2}{2}\bigg(\frac{V'}{V}\bigg)^2,\eta\equiv M_{\mathrm{pl}}^2\frac{V''}{V}$$



$$\ell_s^2 \equiv \alpha', M_s^2 \equiv \frac{1}{\alpha'}$$

$$2\kappa^2 = (2\pi)^7(\alpha')^4$$

Métrica de Friedmann-Robertson-Walker (FRW).

$$ds^2 = -dt^2 + a^2(t)dx^2.$$

$$ds^2 = a^2(\tau)[-d\tau^2 + dx^2],$$

$$\Delta\tau = \int_0^t \frac{dt'}{a(t')} = \int_0^a \frac{d\ln a}{aH}, \text{ where } H \equiv \frac{1}{a} \frac{da}{dt}$$

$$\frac{d}{dt}(aH)^{-1} = -\frac{1}{a}\left[\frac{\dot{H}}{H^2} + 1\right] < 0 \Rightarrow \varepsilon \equiv -\frac{\dot{H}}{H^2} < 1$$

$$a(t) \propto e^{Ht}$$

$$\begin{aligned} 3M_{\text{pl}}^2 H^2 &= \rho \\ 6M_{\text{pl}}^2 (\dot{H} + H^2) &= -(\rho + 3P) \end{aligned}$$

$$2M_{\text{pl}}^2 \dot{H} = -(\rho + P)$$

$$\varepsilon = \frac{3}{2}\left(1 + \frac{P}{\rho}\right)$$

Acción de Goldstone.

$$U(t, \mathbf{x}) \equiv t + \pi(t, \mathbf{x})$$

$$\delta\psi_m(t, \mathbf{x}) \equiv \psi_m(t + \pi(t, \mathbf{x})) - \psi_m(t).$$

$$g_{ij} = a^2(t)\delta_{ij}$$

$$g_{ij} = a^2(t)e^{2\mathcal{R}(t, \mathbf{x})}\delta_{ij}$$

$$\mathcal{R} = -H\pi + \dots,$$

$$S = \int d^4x \sqrt{-g} \mathcal{L} \left[U, (\partial_\mu U)^2, \square U, \dots \right]$$

$$S_\pi^{(2)} = \int d^4x \sqrt{-g} \frac{M_{\text{pl}}^2 |\dot{H}|}{c_s^2} \left[\dot{\pi}^2 - \frac{c_s^2}{a^2} (\partial_i \pi)^2 + 3\varepsilon H^2 \pi^2 \right]$$

$$S_{\mathcal{R}}^{(2)} = \frac{1}{2} \int d^4x a^3 y^2(t) \left[\dot{\mathcal{R}}^2 - \frac{c_s^2}{a^2} (\partial_i \mathcal{R})^2 \right]$$

$$y^2 \equiv 2M_{\text{pl}}^2 \frac{\varepsilon}{c_s^2}$$



$$v \equiv y\mathcal{R} = \int d^3k [v_k(t)a_k e^{ik \cdot x} + c.c.]$$

$$\ddot{v}_k+3H\dot{v}_k+\frac{c_s^2k^2}{a^2}v_k=0$$

$$\omega_k(t) \equiv \frac{c_s k}{a(t)}$$

$$\hat{v}_{\boldsymbol{k}} = v_{\boldsymbol{k}}(t)\hat{a}_{\boldsymbol{k}} + h.c.$$

$$|v_k|^2=\frac{1}{a^3}\frac{1}{2\omega_k}.$$

$$|v_k|^2=\frac{1}{2}\frac{1}{a_\star^3}\frac{1}{c_sk/a_\star},$$

$$\frac{c_s k}{a_\star}=H.$$

$$|v_k|^2=\frac{1}{2}\frac{H^2}{(c_sk)^3},$$

Perturbaciones de Curvatura.

$$P_{\mathcal{R}}(k) \equiv |\mathcal{R}_k|^2 = \frac{1}{4} \frac{H^4}{M_{\text{pl}}^2 |\dot{H}| c_s} \frac{1}{k^3}$$

$$\Delta_{\mathcal{R}}^2(k) \equiv \frac{k^3}{2\pi^2} P_{\mathcal{R}}(k) = \frac{1}{8\pi^2} \frac{H^4}{M_{\text{pl}}^2 |\dot{H}| c_s}.$$

$$n_s - 1 \equiv \frac{d \ln \Delta_{\mathcal{R}}^2}{d \ln k} = -2\varepsilon - \tilde{\eta} - \kappa$$

$$\tilde{\eta} \equiv \frac{\dot{\varepsilon}}{H\varepsilon} \text{ and } \kappa \equiv \frac{\dot{c}_s}{Hc_s}$$

Ondas gravitacionales.

$$g_{ij} = a^2(t)(\delta_{ij} + 2h_{ij})$$

$$S_h^{(2)} = \frac{1}{2} \int d^4x a^3 y^2 \left[(\dot{h}_{ij})^2 - \frac{1}{a^2} (\partial_k h_{ij})^2 \right]$$

$$y^2 \equiv \frac{1}{4} M_{\text{pl}}^2$$

$$\Delta_h^2(k) \equiv \frac{k^3}{2\pi^2} P_h(k) = \frac{2}{\pi^2} \frac{H^2}{M_{\text{pl}}^2},$$



$$n_t \equiv \frac{d\ln~\Delta_h^2}{d\ln~k} = -2\varepsilon$$

$$r\equiv\frac{\Delta_h^2}{\Delta_{\mathcal{R}}^2}$$

$$\frac{H}{M_{\rm pl}}=\pi\Delta_{\mathcal{R}}(k_\star)\sqrt{\frac{r}{2}},$$

$$H=3\times 10^{-5}\left(\frac{r}{0.1}\right)^{1/2} M_{\rm pl}$$

$$E_{\text{inf}}\equiv\left(3H^2M_{\text{pl}}^2\right)^{1/4}=8\times10^{-3}\left(\frac{r}{0.1}\right)^{1/4}M_{\text{pl}}$$

$$\mathbf{Anisotropías\;CMB}.$$

$$e+p\rightarrow H+\gamma$$

$$\Delta T(\boldsymbol{n})\equiv T(\boldsymbol{n})-\bar{T}.$$

$$C(\theta) \equiv \Bigl\langle \frac{\Delta T}{\bar{T}}(\boldsymbol{n}) \frac{\Delta T}{\bar{T}}(\boldsymbol{n}') \Bigr\rangle,$$

$$\frac{\Delta T(\boldsymbol{n})}{\bar{T}}=\sum_{\ell=0}^\infty\sum_{m=-\ell}^{+\ell}a_{\ell m}Y_{\ell m}(\boldsymbol{n})$$

$$\langle a_{\ell m}a_{\ell'm'}^*\rangle=C_\ell\delta_{\ell\ell'}\delta_{mm'}$$

$$C_\ell=2\pi\int_{-1}^1\mathrm{d}\cos~\theta C(\theta)P_\ell(\cos~\theta)$$

$$\hat{C}_\ell=\frac{1}{2\ell+1}\sum_m~|a_{\ell m}|^2$$

$$\mathrm{var}\big(\hat{C}_\ell\big)\equiv\big\langle\hat{C}_\ell\hat{C}_\ell\big\rangle-\big\langle\hat{C}_\ell\big\rangle^2=\frac{2}{2\ell+1}C_\ell^2$$

$$C_\ell=\int~\mathrm{d}\ln~k\Delta_{\mathcal{R}}^2(k)T_\ell^2(k)$$

$$\mathbf{Polarización\;CMB}.$$

$$(Q\pm iU)'(\boldsymbol{n})=e^{\mp 2i\psi}(Q\pm iU)(\boldsymbol{n}).$$

$$(Q\pm iU)(\boldsymbol{n})=\sum_{\ell m}~a_{\pm 2,\ell m\pm 2}Y_{\ell m}(\boldsymbol{n})$$



$$E(\mathbf{n}) \equiv a_{E,\ell m} Y_{\ell m}(\mathbf{n}), \quad a_{E,\ell m} \equiv -\frac{a_{2,\ell m} + a_{-2,\ell m}}{2}$$

$$B(\mathbf{n}) \equiv a_{B,\ell m} Y_{\ell m}(\mathbf{n}), \quad a_{B,\ell m} \equiv -\frac{a_{2,\ell m} - a_{-2,\ell m}}{2i}$$

$$\langle a_{X,\ell m} a_{Y,\ell' m'}^* \rangle = C_\ell^{XY} \delta_{\ell \ell'} \delta_{mm'}, X,Y \equiv \{T,E,B\}.$$

Estructura Escalar.

$$P_\delta(z, k) = T_\delta^2(z, k) P_{\mathcal{R}}(k)$$

$$\delta_g(z, \mathbf{x}) = b(z) \delta(z, \mathbf{x})$$

$$D_V(\bar{z}) \equiv \left[(1+\bar{z})^2 D_A^2(\bar{z}) \frac{c\bar{z}}{H(\bar{z})} \right]^{1/3}$$

Modelo CDM.

$$\Delta_{\mathcal{R}}^2(k) = A_s \left(\frac{k}{k_*} \right)^{n_s-1}.$$

$$A_s = (2.196^{+0.051}_{-0.060}) \times 10^{-9}.$$

$$n_s = 0.9603 \pm 0.0073$$

Fluctuaciones Escalares.

$$\Delta_{\mathcal{R}}^2(k) = A_s \left(\frac{k}{k_*} \right)^{n_s-1 + \frac{1}{2}\alpha_s \ln(k/k_*)}$$

Fluctuaciones Tensoriales.

$$r < 0.12 \text{ (95% limit).}$$

$$\langle 0 | \hat{\mathcal{R}}_{\mathbf{k}_1} \hat{\mathcal{R}}_{\mathbf{k}_2} | 0 \rangle = (2\pi)^3 P_{\mathcal{R}}(k_1) \delta(\mathbf{k}_1 + \mathbf{k}_2),$$

$$\langle \Omega | \hat{\mathcal{R}}_{\mathbf{k}_1} \cdots \hat{\mathcal{R}}_{\mathbf{k}_n} | \Omega \rangle \propto \int [D\mathcal{R}] \mathcal{R}_{\mathbf{k}_1} \cdots \mathcal{R}_{\mathbf{k}_n} e^{iS[\mathcal{R}]},$$

$$\langle \Omega | \hat{\mathcal{R}}_{\mathbf{k}_1} \hat{\mathcal{R}}_{\mathbf{k}_2} \hat{\mathcal{R}}_{\mathbf{k}_3} | \Omega \rangle = (2\pi)^3 B_{\mathcal{R}}(k_1, k_2, k_3) \delta(\mathbf{k}_1 + \mathbf{k}_2 + \mathbf{k}_3).$$

$$f_{\text{NL}} \equiv \frac{5}{18} \frac{B_{\mathcal{R}}(k, k, k)}{P_{\mathcal{R}}^2(k)}$$



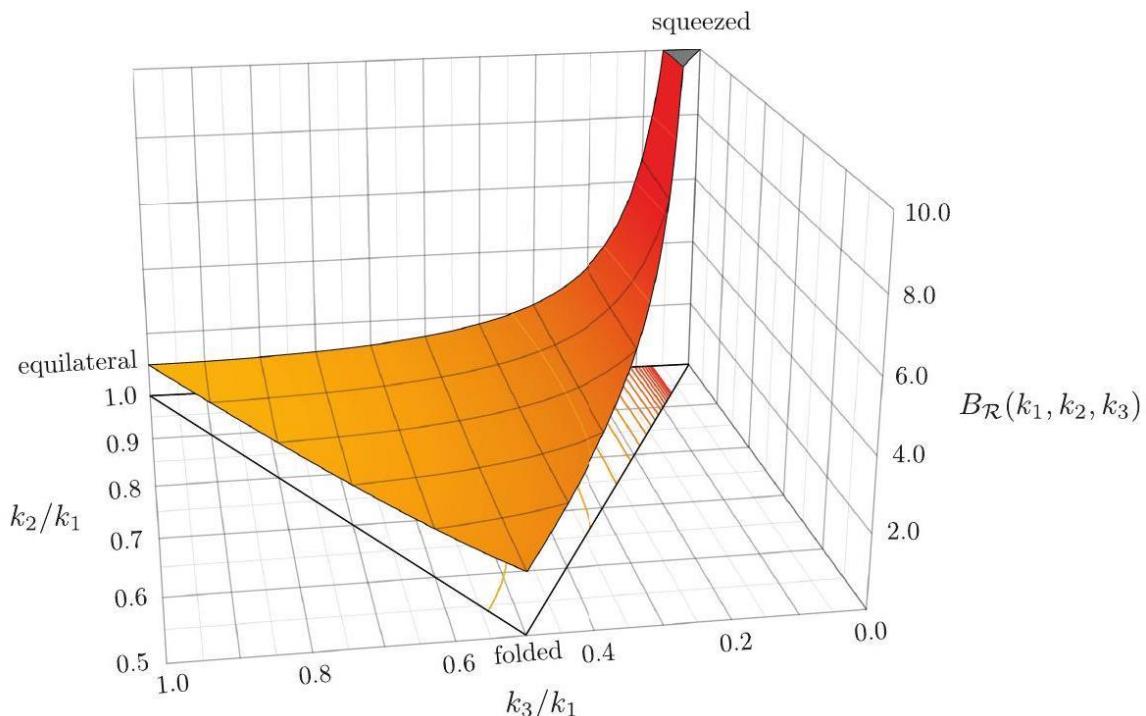
$$B_{\text{local}} \equiv \frac{6}{5}(P_1 P_2 + \text{perms.})$$

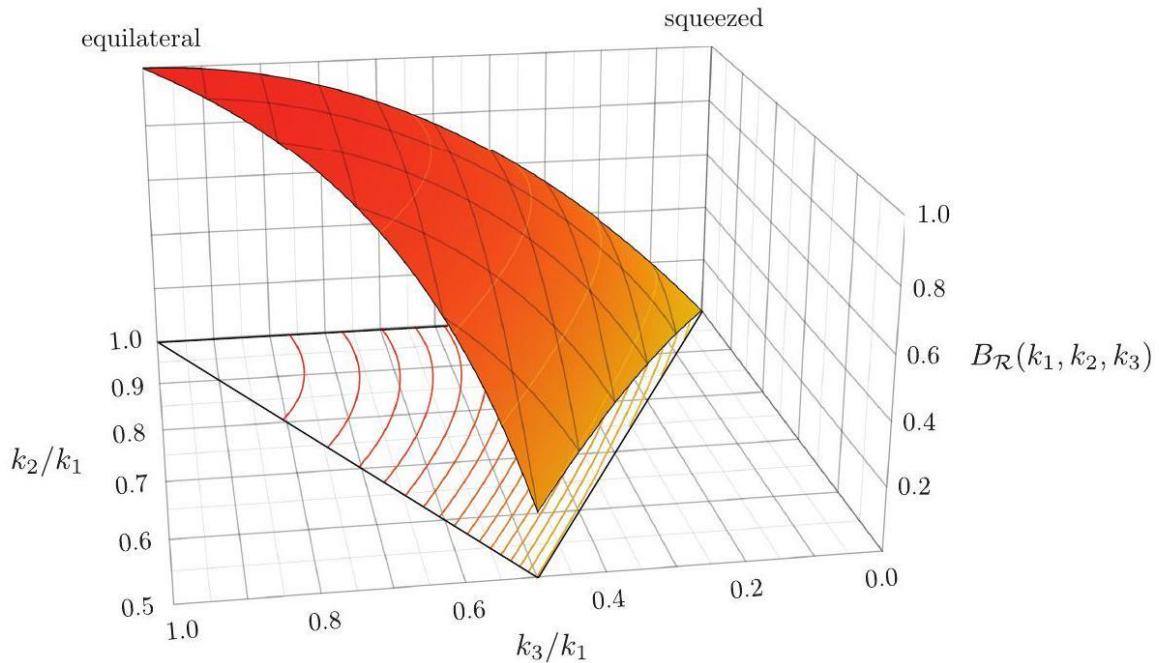
$$B_{\text{equil}} \equiv \frac{3}{5}(6(P_1^3 P_2^2 P_3)^{1/3} - 3P_1 P_2 - 2(P_1 P_2 P_3)^{2/3} + \text{perms.})$$

$$B_{\text{ortho}} \equiv \frac{3}{5}(18(P_1^3 P_2^2 P_3)^{1/3} - 9P_1 P_2 - 8(P_1 P_2 P_3)^{2/3} + \text{perms.})$$

$$\mathcal{R}(\boldsymbol{x}) \equiv \mathcal{R}_g(\boldsymbol{x}) + \frac{3}{5} f_{\text{NL}}^{\text{local}} [\mathcal{R}_g^2(\boldsymbol{x}) - \langle \mathcal{R}_g^2 \rangle]$$

$$\lim_{k_1 \rightarrow 0} \frac{B_{\mathcal{R}}(k_1, k_2, k_3)}{P_{\mathcal{R}}(k_1)P_{\mathcal{R}}(k_2)} = (1 - n_s) \ll 1$$





Figuras 34 y 35. Curvatura local por deformación del espacio – tiempo cuántico, a propósito de una particular supermasiva, esto en dimensión \mathbb{R}^4 .

$$S_{\pi}^{(3)} = \int d^4x \sqrt{-g} \frac{M_{\text{pl}}^2 \dot{H}}{c_s^2} (1 - c_s^2) \left(\frac{\dot{\pi}(\partial_i \pi)^2}{a^2} + \frac{A}{c_s^2} \dot{\pi}^3 \right)$$

$$\begin{aligned} f_{\text{NL}}^{\text{local}} &= 2.7 \pm 5.8, \\ f_{\text{NL}}^{\text{equil}} &= -42 \pm 75, \\ f_{\text{NL}}^{\text{ortho}} &= -25 \pm 39. \end{aligned}$$

Adiabaticity.

$$\delta_I(t, \mathbf{x}) \equiv \frac{\bar{\rho}_I(t + \pi(t, \mathbf{x})) - \bar{\rho}_I(t)}{\bar{\rho}_I(t)} \approx \frac{\dot{\bar{\rho}}_I}{\bar{\rho}_I} \pi(t, \mathbf{x}).$$

$$\mathcal{S} \equiv \delta_c - \frac{3}{4} \delta_\gamma$$

$$\alpha \equiv \frac{P_{\mathcal{S}}}{P_{\mathcal{R}}}$$

$$\alpha_0 < 0.036$$

$$\alpha_{+1} < 0.0025.$$



Curvatura isotrópica.

$$n_t = 2 \frac{\dot{H}}{H^2}$$

$$n_t = -\frac{r}{8}$$

$$\alpha_s = 16\varepsilon^2 - 6\varepsilon\tilde{\eta} + \tilde{\eta}\chi$$

$$\mathcal{N}^{\text{CMB}} \sim \left(\frac{\ell_{\max}}{\ell_{\min}} \right)^2 \sim 10^6$$

$$\mathcal{N}_{\text{linear}}^{\text{LSS}} \sim \left(\frac{k_{\max}}{k_{\min}} \right)^3 \sim 10^9$$

$$\mathcal{N}_{\text{linear}}^{\text{Euclid}} \sim \left(\frac{k_{\max}}{k_{\min}} \right)^3 \sim 10^6,$$

$$|f_{\text{NL}}^{\text{local}}| \gtrsim \mathcal{O}(1)$$

$$f_{\text{NL}}^{\text{equil}} \sim \mathcal{O}(1)$$

$$M_{\text{pl}} \equiv \frac{1}{\sqrt{8\pi G_N}} = 2.4 \times 10^{18} \text{GeV}$$

Acción efectiva.

$$\mathcal{L}[\phi, \Psi] = \mathcal{L}_l[\phi] + \mathcal{L}_h[\Psi] + \mathcal{L}_{lh}[\phi, \Psi]$$

$$e^{iS_{\text{eff}}[\phi]} = \int [\mathcal{D}\Psi] e^{iS[\phi, \Psi]}$$

$$\phi(-\square + M^2)^{-1}\phi = \frac{\phi}{M^2} \left(1 + \frac{\square}{M^2} + \dots \right) \phi,$$

$$\mathcal{L}_{\text{eff}}[\phi] = \mathcal{L}_l[\phi] + \sum_i c_i(g) \frac{\mathcal{O}_i[\phi]}{M^{\delta_i-4}},$$

$$\mathcal{L}[\phi, \Psi] = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 - \frac{1}{4!}\lambda\phi^4 - \frac{1}{2}(\partial\Psi)^2 - \frac{1}{2}M^2\Psi^2 - \frac{1}{4}g\phi^2\Psi^2.$$

$$\mathcal{L}_{\text{eff}}[\phi] = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m_R^2\phi^2 - \frac{1}{4!}\lambda_R\phi^4 - \sum_{i=1}^{\infty} \left(\frac{c_i(g)}{M^{2i}}\phi^{4+2i} + \frac{d_i(g)}{M^{2i}}(\partial\phi)^2\phi^{2i} + \dots \right)$$



$$m_R^2 = \text{---} + \text{---} + \dots ,$$

$$\lambda_R = \text{X} + \text{---} + \dots ,$$

$$m_R^2 = m^2 + \frac{g}{32\pi^2} (\Lambda^2 - M^2 L) + \dots$$

$$\lambda_R = \lambda - \frac{3g^2}{32\pi^2} L + \dots$$

$$L \rightarrow \frac{1}{\epsilon} + \gamma - \ln(4\pi)$$

$$c_1 = \text{---} + \dots \sim \mathcal{O}(g^3) + \dots .$$

Parametrizaciones de simetría.

$$\mathcal{L}_{\text{eff}}[\phi] = \mathcal{L}_l[\phi] + \sum_i c_i \frac{\mathcal{O}_i[\phi]}{\Lambda^{\delta_i-4}},$$

$$\Delta m_e = \alpha \Lambda$$

$$\Delta m_e = \alpha m_e \ln(\Lambda/m_e)$$

$$m_{\pi^+}^2 - m_{\pi^0}^2 = \frac{3\alpha}{4\pi} \Lambda^2$$

$$\frac{m_{K_L^0} - m_{K_S^0}}{m_{K_L^0}} = \frac{G_F^2 f_K^2}{6\pi^2} \sin^2 \theta_c \Lambda^2,$$

$$\Delta m_H^2 \sim \frac{y_t^2}{(4\pi)^2} \Lambda^2$$

$$\langle T_{\mu\nu} \rangle = -\rho_{\text{vac}} g_{\mu\nu}$$

$$\Delta \rho_{\text{vac}} \sim \Lambda^4$$



$$\Delta m^2 \propto \Lambda^2$$

$$\phi\mapsto \phi+\,\mathrm{const}\,.$$

$$\Delta m^2 \propto m^2$$

$$\lim_{g\rightarrow 0}c_i(g)=0$$

$$S_{\text{EH}} = \frac{M_{\text{pl}}^2}{2} \int \text{d}^4x \sqrt{-g} R$$

$$S_{\text{EH}} = \int \text{d}^4x \left[(\partial h)^2 + \frac{1}{M_{\text{pl}}} h(\partial h)^2 + \frac{1}{M_{\text{pl}}^2} h^2 (\partial h)^2 + \dots \right]$$

$$S_{\text{YM}} = \int \text{d}^4x [(\partial A)^2 + g A^2 \partial A + g^2 A^4]$$

$$S_g = \int \text{d}^4x \sqrt{-g} \left[M_\Lambda^4 + \frac{M_{\text{pl}}^2}{2} R + c_1 R^2 + c_2 R_{\mu\nu} R^{\mu\nu} + \frac{1}{M^2} (d_1 R^3 + \dots) + \dots \right]$$

$$S_g = \int \text{d}^{10}X \sqrt{-G} \left[\frac{M_{10}^2}{2} R + \frac{\zeta(3)}{3 \cdot 2^5} \frac{1}{M^6} \mathcal{R}^4 + \dots \right]$$

$$M^2=\frac{4}{\alpha'}$$

$$S_{\text{eff}}[\phi,g] = S_g + S_{\text{eff}}[\phi] + S_{g,\phi}$$

$$S_{g,\phi} = \int \text{d}^4x \sqrt{-g} \left[\sum_i c_i \frac{\mathcal{O}_i[g,\phi]}{\Lambda^{\delta_i - 4}} \right]$$

$$S_{g,\phi}^{(4)} = \int \text{d}^4x \sqrt{-g} \xi \phi^2 R$$

$$g_{\mu\nu} \mapsto \bar{g}_{\mu\nu} \equiv e^{2\omega(\phi)} g_{\mu\nu},$$

Dinámicas y perturbaciones inflacionarias.

$$S = \int \text{d}^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right]$$

$$3M_{\text{pl}}^2 H^2 = \frac{1}{2} \dot{\phi}^2 + V \text{ and } \ddot{\phi} + 3H\dot{\phi} = -V',$$

$$\varepsilon = -\frac{\dot{H}}{H^2} = \frac{\frac{1}{2}\dot{\phi}^2}{M_{\text{pl}}^2 H^2}$$



$$\epsilon \equiv \frac{M_{\rm pl}^2}{2}\bigg(\frac{V'}{V}\bigg)^2 \ll 1\,, |\eta| \equiv M_{\rm pl}^2\frac{|V''|}{V} \ll 1.$$

$$\pi=\frac{\delta\phi}{\dot{\phi}}$$

$$\mathcal{R}(t,\boldsymbol{x})=-H\pi(t,\boldsymbol{x})=-\frac{H}{\dot{\phi}}\delta\phi(t,\boldsymbol{x})$$

$$\Delta_{\mathcal{R}}^2 = \frac{1}{24\pi^2}\frac{1}{\epsilon}\frac{V}{M_{\rm pl}^4}\;,\Delta_h^2 = \frac{2}{3\pi^2}\frac{V}{M_{\rm pl}^4}$$

$$\begin{array}{l} n_s-1=2\eta-6\epsilon \\ r=16\epsilon \end{array}$$

$$N_\star = \int_{\phi_{\rm end}}^{\phi_\star} \frac{{\rm d}\phi}{M_{\rm pl}} \frac{1}{\sqrt{2\epsilon}}$$

$$V(\phi)=\mu^{4-p}\phi^p$$

$$\epsilon = \frac{p^2}{2}\bigg(\frac{M_{\rm pl}}{\phi}\bigg)^2\,, \eta = p(p-1)\bigg(\frac{M_{\rm pl}}{\phi}\bigg)^2\,.$$

$$N_\star \approx \frac{1}{2p}\bigg(\frac{\phi_\star}{M_{\rm pl}}\bigg)^2\,,$$

$$n_s-1=-\frac{(2+p)}{2N_\star}\,,r=\frac{4p}{N_\star}.$$

$$\begin{array}{lll} p=1: & n_s\approx 0.975, & r\approx 0.07, \quad \phi_\star\approx 11M_{\rm pl}, \\ p=2: & n_s\approx 0.967, & r\approx 0.13, \quad \phi_\star\approx 15M_{\rm pl}, \\ p=3: & n_s\approx 0.958, & r\approx 0.20, \quad \phi_\star\approx 19M_{\rm pl}, \\ p=4: & n_s\approx 0.950, & r\approx 0.27, \quad \phi_\star\approx 22M_{\rm pl}. \end{array}$$

$$V(\phi) = \frac{V_0}{2}\bigg[1-\cos\left(\frac{\phi}{f}\right)\bigg],$$

$$\begin{aligned} n_s-1 &= -\alpha \frac{e^{N_\star \alpha}+1}{e^{N_\star \alpha}-1} \stackrel{\alpha \ll 1}{\rightarrow} -\frac{2}{N_\star}, \\ r &= 8\alpha \frac{1}{e^{N_\star \alpha}-1} \stackrel{\alpha \ll 1}{\rightarrow} +\frac{8}{N_\star}, \end{aligned}$$

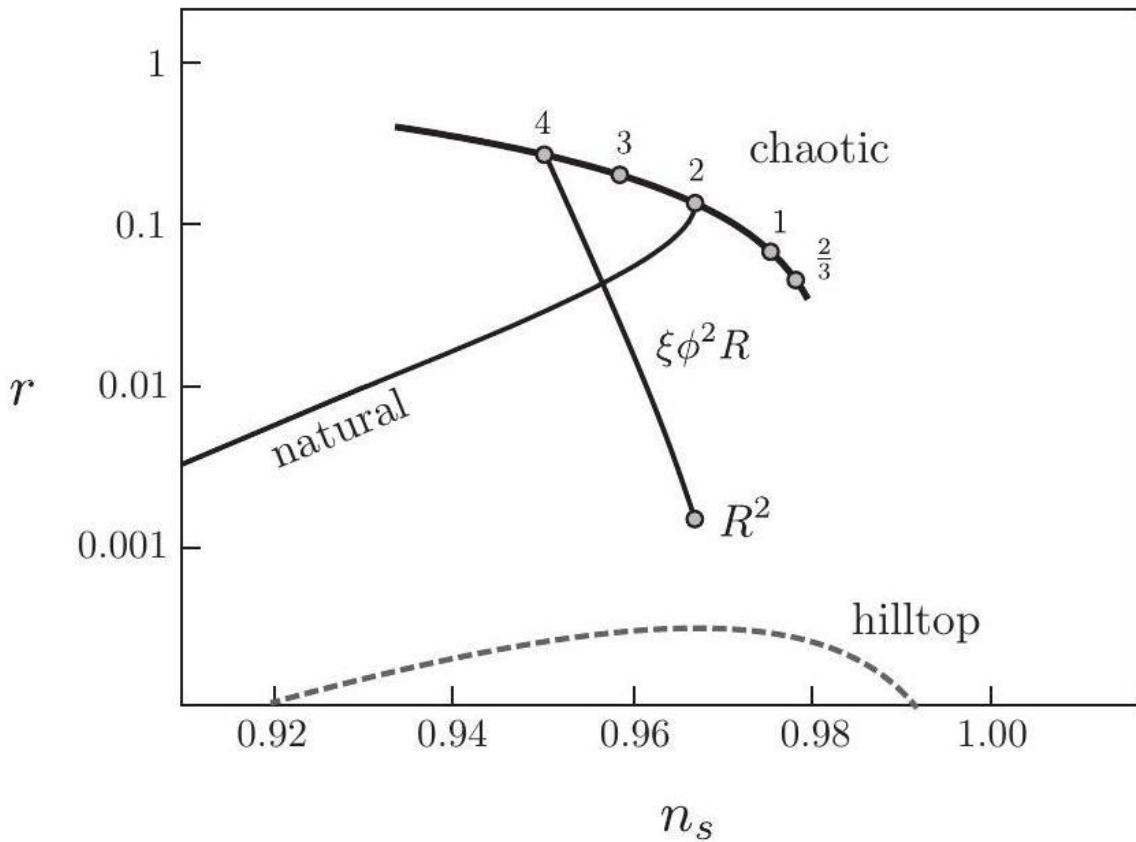
$$V(\phi)=V_0+\frac{1}{2}m^2\phi^2+\cdots$$

$$V(\phi)=V_0\left[1+\frac{1}{2}\eta_0\frac{\phi^2}{M_{\rm pl}^2}+\cdots\right], \text{ where } \eta\approx\eta_0<0$$



$$n_s - 1 = 2\eta_0$$

$$r = 2(1 - n_s)^2 e^{-N_*(1-n_s)} \left(\frac{\phi_{\text{end}}}{M_{\text{pl}}}\right)^2 \approx 10^{-3} \left(\frac{\phi_{\text{end}}}{M_{\text{pl}}}\right)^2.$$



$$V(\phi) = V_0 \left[1 + \lambda_0 \frac{\phi}{M_{\text{pl}}} + \frac{1}{2} \eta_0 \frac{\phi^2}{M_{\text{pl}}^2} + \frac{1}{3!} \mu_0 \frac{\phi^3}{M_{\text{pl}}^3} + \dots \right]$$

$$V(\phi) \approx V_0 \left[1 + \lambda_0 \frac{\phi}{M_{\text{pl}}} + \frac{1}{3!} \mu_0 \frac{\phi^3}{M_{\text{pl}}^3} + \dots \right]$$

$$n_s - 1 = -4 \sqrt{\frac{\lambda_0 \mu_0}{2}} \cot \left(N_* \sqrt{\frac{\lambda_0 \mu_0}{2}} \right),$$

$$r = 16 \lambda_0^2$$

$$V(\phi, \Psi) = V(\phi) + V(\Psi) + \frac{1}{2} g \phi^2 \Psi^2,$$

$$V(\Psi) \equiv \frac{1}{4\lambda} (M^2 - \lambda \Psi^2)^2$$



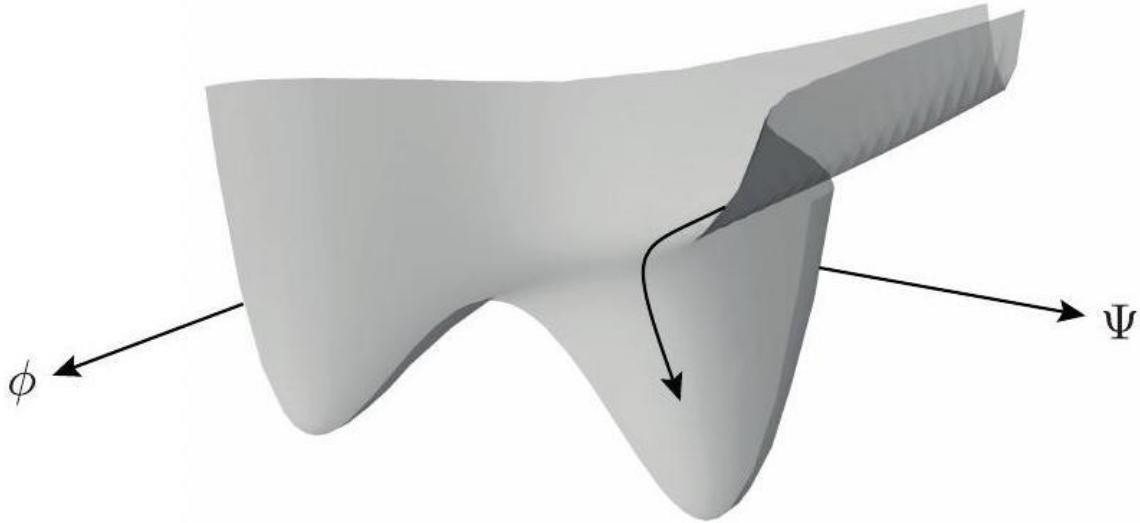


Figura 36. Membranas en dimensión \mathbb{R}^4

$$M_\Psi^2(\phi) = -M^2 + g\phi^2$$

$$S = \frac{M_{\text{pl}}^2}{2} \int d^4x \sqrt{-g} \left(R + \frac{\alpha}{2M_{\text{pl}}^2} R^2 \right)$$

$$S = \int d^4x \sqrt{-\tilde{g}} \left(\frac{M_{\text{pl}}^2}{2} \tilde{R} - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right)$$

$$V(\phi) = \frac{M_{\text{pl}}^4}{4\alpha} \left(1 - \exp \left[-\sqrt{\frac{2}{3}} \frac{\phi}{M_{\text{pl}}} \right] \right)^2.$$

$$\eta = -\frac{4}{3} e^{-\sqrt{2/3}\phi/M_{\text{pl}}}, \epsilon = \frac{3}{4}\eta^2$$

$$\alpha = 2.2 \times 10^8$$

$$n_s - 1 \approx -\frac{2}{N_\star}, r \approx \frac{12}{N_\star^2}$$

$$S = \int d^4x \sqrt{-g} \left[\frac{M_{\text{pl}}^2}{2} \left(1 + \xi \frac{\varphi^2}{M_{\text{pl}}^2} \right) R - \frac{1}{2} (\partial\varphi)^2 - \frac{\lambda}{4} \varphi^4 \right]$$

$$S = \int d^4x \sqrt{-\tilde{g}} \left[\frac{M_{\text{pl}}^2}{2} \tilde{R} - \frac{1}{2} k(\varphi) (\partial\varphi)^2 - V(\varphi) \right]$$

$$k(\varphi) = \frac{1 + (6\xi + 1)\psi^2}{(1 + \psi^2)^2}$$

$$V(\varphi) = \frac{\lambda M_{\text{pl}}^4}{4\xi^2} \frac{\psi^4}{(1 + \psi^2)^2}, \psi^2 \equiv \frac{\xi \varphi^2}{M_{\text{pl}}^2}$$

$$\frac{\phi}{M_{\mathrm{pl}}}=\sqrt{\frac{6\xi+1}{\xi}}\sinh^{-1}\left(\sqrt{6\xi+1}\psi\right)-\sqrt{6}\sinh^{-1}\left(\sqrt{6\xi}\frac{\psi}{\sqrt{1+\psi^2}}\right).$$

$$\frac{\phi}{M_{\mathrm{pl}}}\approx \sqrt{\frac{3}{2}}\ln{(1+\psi^2)},$$

$$V(\phi)=\frac{\lambda M_{\mathrm{pl}}^4}{4\xi^2}\Bigg(1-\exp\left[-\sqrt{\frac{2}{3}}\frac{\phi}{M_{\mathrm{pl}}}\right]\Bigg)^2.$$

$$\xi=47000\sqrt{\lambda}$$

$$\begin{gathered} n_s-1=-\frac{32\xi}{16\xi N_\star-1}\stackrel{\xi\gg1}{\rightarrow}-\frac{2}{N_\star},\\ r=+\frac{12}{N_\star^2}\frac{6\xi+1}{6\xi}\stackrel{\xi\gg1}{\rightarrow}+\frac{12}{N_\star^2}. \end{gathered}$$

$${\bf Inflación K.}$$

$$S = \int \; {\rm d}^4x \sqrt{-g} \biggl[\frac{M_{\mathrm{pl}}^2}{2} R + P(X,\phi) \biggr]$$

$$3M_{\mathrm{pl}}^2H^2=2P_{,X}X-P\;\;\text{and}\;\;\frac{d}{dt}\big(a^3P_{,X}\dot{\phi}\big)=a^3P_{,\phi}$$

$$\varepsilon=-\frac{\dot{H}}{H^2}=\frac{3XP_{,X}}{2XP_{,X}-P}$$

$$c_s^2=\frac{dP}{d\rho}=\frac{P_{,X}}{P_{,X}+2XP_{,XX}}.$$

$$\begin{array}{l} n_s-1=-2\varepsilon-\tilde{\eta}-\kappa \\ r=16\varepsilon c_s \end{array}$$

$$S_{\mathrm{eff}}[\phi]=\int \; {\rm d}^4x \sqrt{-g} \biggl[\frac{M_{\mathrm{pl}}^2}{2} R + \mathcal{L}_l[\phi] + \sum_i \; c_i \frac{\mathcal{O}_i[\phi]}{\Lambda^{\delta_i-4}} \biggr]$$

$$\Delta m^2 \sim \Lambda^2.$$

$$\Delta\eta\sim\frac{\Lambda^2}{H^2}\gtrsim 1$$

$$\Delta m^2 \sim H^2$$

$$\Delta\eta\sim 1.$$

$$\phi\mapsto\phi+\,{\mathrm{const}}\,.$$



Operadores en altas dimensiones.

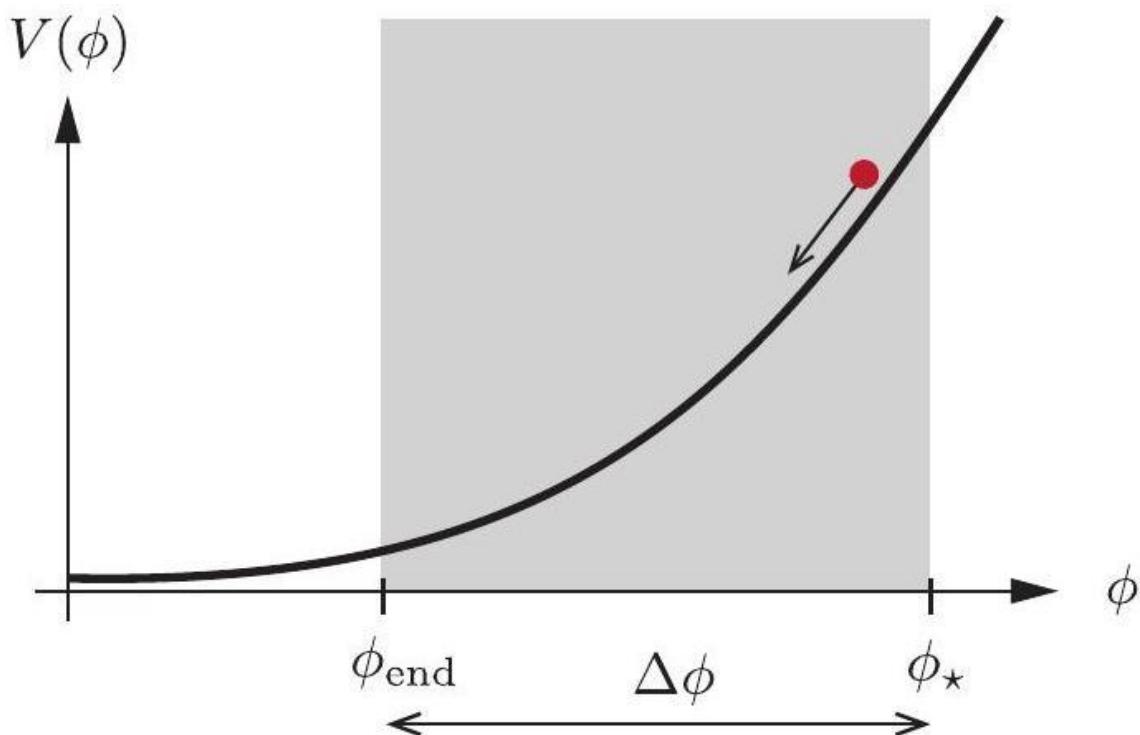
$$\mathcal{O}_6 = c V_l(\phi) \frac{\phi^2}{\Lambda^2}$$

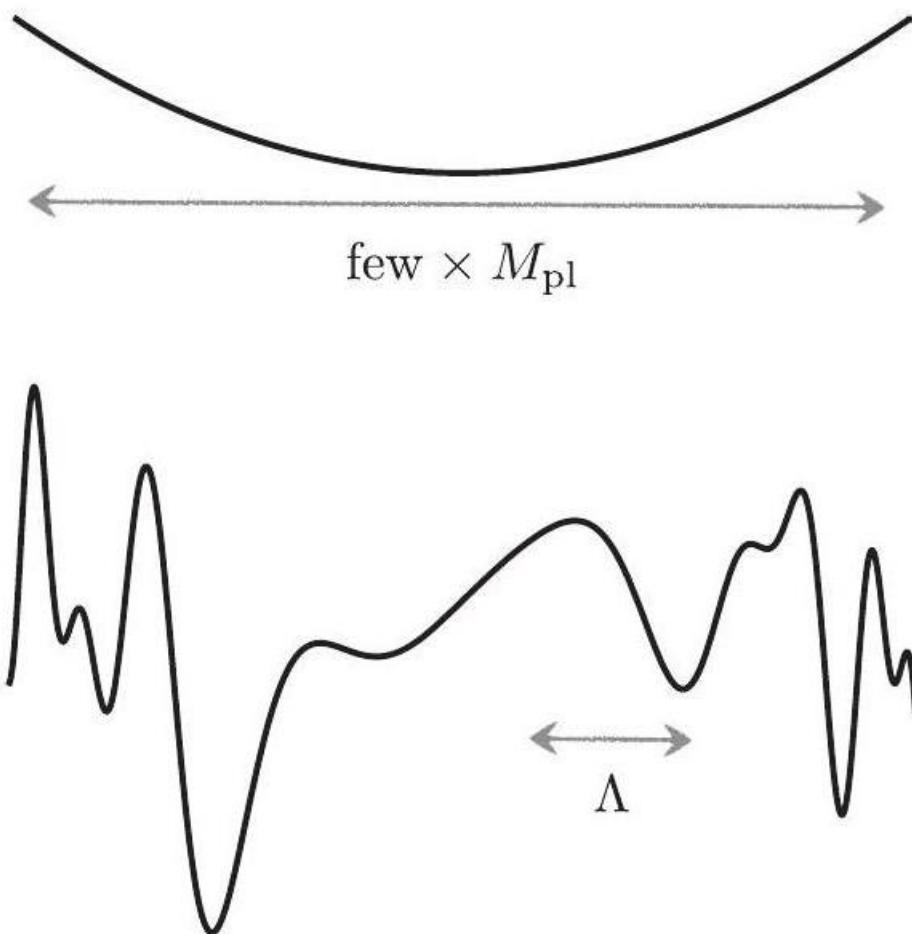
$$\Delta\eta \approx 2c \left(\frac{M_{\text{pl}}}{\Lambda} \right)^2$$

$$\mathcal{O}_\delta = c \langle V \rangle \left(\frac{\phi}{\Lambda} \right)^{\delta-4}$$

$$\Delta\eta \approx c(\delta-4)(\delta-5) \left(\frac{M_{\text{pl}}}{\Lambda} \right)^2 \left(\frac{\phi}{\Lambda} \right)^{\delta-6}.$$

Ondas gravitacionales en campos planckianos y campos inflacionarios.





Figuras 37 y 38. Fluctuaciones gravitacionales en la curvatura local.

$$r = 8 \left(\frac{1}{M_{\text{pl}}} \frac{d\phi}{dN} \right)^2, \text{ where } dN \equiv H dt.$$

$$\frac{\Delta\phi}{M_{\text{pl}}} = \int_0^{N_*} dN \sqrt{\frac{r(N)}{8}}.$$

$$N_{\text{eff}} \equiv \int_0^{N_*} dN \sqrt{\frac{r(N)}{r_*}}$$

$$\frac{\Delta\phi}{M_{\text{pl}}} = N_{\text{eff}} \sqrt{\frac{r_*}{8}}.$$

$$\frac{d\ln r}{dN} = - \left[n_s - 1 + \frac{r}{8} \right],$$

$$\frac{\Delta\phi}{M_{\text{pl}}} \gtrsim \left(\frac{r}{0.01} \right)^{1/2}$$



$$\frac{\Delta \phi}{M_{\rm pl}} \gtrsim 0.25 \times \Big(\frac{r}{0.01}\Big)^{1/2}$$

$$\frac{\Delta V}{V}=c_1\frac{V''}{M_{\rm pl}^2}+c_2\frac{V}{M_{\rm pl}^4},$$

$$\mathcal{L}_{\text{eff}}[\phi]=\mathcal{L}_l[\phi]+\sum_{i=1}^\infty\,\Big(\frac{c_i}{\Lambda^{2i}}\phi^{4+2i}+\frac{d_i}{\Lambda^{2i}}(\partial\phi)^2\phi^{2i}+\frac{e_i}{\Lambda^{4i}}(\partial\phi)^{2(i+1)}+\cdots\Big),$$

$$\phi \mapsto \phi + \; {\mathrm {const}}\, .$$

$$\frac{\Delta \phi}{M_{\rm pl}} = \left(c_s P_{,X}\right)^{-1/2}\sqrt{\frac{r}{8}}\Delta N.$$

$$\Delta \mathcal{L}=P\left(X-V(\phi)\frac{\phi^2}{M_{\rm pl}^2}\right)=P(X)-P_{,X}V(\phi)\frac{\phi^2}{M_{\rm pl}^2}+\cdots$$

$$c_s^2=\frac{P_{,X}}{P_{,X}+2XP_{,XX}}$$

$$P=X+\frac{1}{2}\frac{X^2}{\Lambda^4}+\cdots$$

$$\left|f_{\rm NL}^{\rm equil}\right|\sim\frac{1}{c_s^2}-1\approx\frac{X}{\Lambda^4}+\cdots\ll1.$$

$${\mathcal O}_5=\frac{\psi X}{\Lambda}$$

$$\Lambda \gtrsim 10^5 H$$

$$\Lambda \gtrsim \Big(\frac{r}{0.01}\Big)^{1/2} M_{\rm pl}$$

$$S_{\rm P}=-\frac{1}{4\pi\alpha'}\int\;\;{\rm d}^2\sigma\sqrt{-h}h^{ab}\partial_aX^M(\sigma)\partial_bX^N(\sigma)\eta_{MN}$$

$$S_{\rm P}=-\frac{1}{4\pi\alpha'}\int\;\;{\rm d}^2\sigma\partial^aX^M\partial_aX_M$$

$$S_\sigma=-\frac{1}{4\pi\alpha'}\int\;\;{\rm d}^2\sigma\sqrt{-h}\big([h^{ab}G_{MN}(X)+\epsilon^{ab}B_{MN}(X)]\partial_aX^M\partial_bX^N\!+\!\alpha'\Phi(X)R(h)\big)$$

$$S_{\rm B}=\frac{1}{2\kappa_D^2}\!\int\;\;{\rm d}^DX\sqrt{-G}e^{-2\Phi}\!\left(R+4(\partial\Phi)^2-\frac{1}{2}|H_3|^2-\frac{2(D-26)}{3\alpha'}+\mathcal{O}(\alpha')\right)$$



$$e^{-\Phi \chi}=e^{-\Phi(2-2g)}\equiv g_s^{2g-2}$$

$$S=S_{\rm P}+S_{\rm F}=-\frac{1}{4\pi\alpha'}\int\,\,{\rm d}^2\sigma(\partial^aX^M\partial_aX_M-i\bar{\psi}^M\rho^a\partial_a\psi_M)$$

$$\{\rho^a,\rho^b\}=-2\eta^{ab}$$

$$\psi^M\equiv\begin{pmatrix}\psi_-^M\\ \psi_+^M\end{pmatrix}$$

$$S_{\rm F}=\frac{i}{2\pi\alpha'}\int\,\,{\rm d}^2\sigma(\psi_-^M\partial_+\psi_-^N+\psi_+^M\partial_-\psi_+^N)\eta_{MN}$$

$$S_{\rm NS} = \frac{1}{2\kappa^2} \int \,\,\, {\rm d}^{10} X \sqrt{-G} e^{-2\Phi} \left(R + 4 (\partial \Phi)^2 - \frac{1}{2} |H_3|^2 \right)$$

$$2\kappa^2=(2\pi)^7(\alpha')^4$$

$$S_{\rm IIA}=S_{\rm NS}+S_{\rm R}^{(\rm IIA)}+S_{\rm CS}^{(\rm IIA)},$$

$$\begin{aligned} S_{\rm R}^{(\rm IIA)} &= -\frac{1}{4\kappa^2}\int\,\,{\rm d}^{10} X \sqrt{-G} \left(|F_2|^2 + \left| \tilde{F}_4 \right|^2 \right), \\ S_{\rm CS}^{(\rm IIA)} &= -\frac{1}{4\kappa^2}\int\,\,B_2 \wedge F_4 \wedge F_4, \end{aligned}$$

$$S_{\rm IIB}=S_{\rm NS}+S_{\rm R}^{(\rm IIB)}+S_{\rm CS}^{(\rm IIB)},$$

$$\begin{aligned} S_{\rm R}^{(\rm IIB)} &= -\frac{1}{4\kappa^2}\int\,\,{\rm d}^{10} X \sqrt{-G} \left(|F_1|^2 + \left| \tilde{F}_3 \right|^2 + \frac{1}{2} \left| \tilde{F}_5 \right|^2 \right), \\ S_{\rm CS}^{(\rm IIB)} &= -\frac{1}{4\kappa^2}\int\,\,C_4 \wedge H_3 \wedge F_3, \end{aligned}$$

$$\tilde{F}_5=\star\tilde{F}_5$$

$$G_{E,MN}\equiv e^{-\Phi/2}G_{MN}$$

$$\begin{array}{lcl} G_3 & \equiv & F_3 - \tau H_3 \\ \tau & \equiv & C_0 + i e^{-\Phi} \end{array}$$

$$S_{\rm IIB}=\frac{1}{2\kappa^2}\int\,\,{\rm d}^{10} X \sqrt{-G_E} \Bigg[R_E - \frac{|\partial\tau|^2}{2({\rm Im}(\tau))^2} - \frac{|G_3|^2}{2{\rm Im}(\tau)} - \frac{\left|\tilde{F}_5\right|^2}{4} \Bigg] - \frac{i}{8\kappa^2}\int\,\,\frac{C_4\wedge G_3\wedge\bar{G}_3}{{\rm Im}(\tau)}$$



Membranas en dimensión \mathbb{R}^4 .

$$S_{\text{CS}} = \mu_p \int_{\Sigma_{p+1}} C_{p+1}$$

$$S_{\text{D}} = -T_p \int \text{d}^{p+1}\sigma \sqrt{-\det(G_{ab})}$$

$$G_{ab} \equiv \frac{\partial X^M}{\partial \sigma^a} \frac{\partial X^N}{\partial \sigma^b} G_{MN}$$

$$S_{\text{BI}} = -Q_p \int \text{d}^{p+1}\sigma \sqrt{-\det(\eta_{ab} + 2\pi\alpha' F_{ab})} = -Q_p \int \text{d}^{p+1}\sigma \left(1 + \frac{(2\pi\alpha')^2}{4} F_{ab} F^{ab} + \dots \right)$$

$$S_{\text{DBI}} = -g_s T_p \int \text{d}^{p+1}\sigma e^{-\Phi} \sqrt{-\det(G_{ab} + \mathcal{F}_{ab})}$$

$$\mathcal{F}_{ab} \equiv B_{ab} + 2\pi\alpha' F_{ab}$$

$$T_p \equiv \frac{1}{(2\pi)^p g_s (\alpha')^{(p+1)/2}}$$

$$S_{\text{CS}} = i\mu_p \int_{\Sigma_{p+1}} \sum_n C_n \wedge e^{\mathcal{F}}$$

$$S_{\text{D}p} = S_{\text{DBI}} + S_{\text{CS}}$$

$$r_+^{7-p} = d_p g_s N (\alpha')^{\frac{1}{2}(7-p)}$$

$$e^\Phi = g_s \left(1 + \left(\frac{r_+}{r} \right)^{7-p} \right)^{\frac{1}{4}(3-p)}$$

$$1 \ll g_s N \ll N$$

$$\mathcal{M}_{10}=\mathcal{M}_4\times X_6$$

$$G_{MN} \text{d}X^M \text{d}X^N = \eta_{\mu\nu} \text{d}x^\mu \text{d}x^\nu + g_{mn} \text{d}y^m \text{d}y^n$$

$$G_{MN} \text{d}X^M \text{d}X^N = e^{2A(y)} g_{\mu\nu} \text{d}x^\mu \text{d}x^\nu + e^{-2A(y)} g_{mn} \text{d}y^m \text{d}y^n,$$

$$\mathcal{L}_\Phi = -K_{i\bar J}\partial^\mu\phi^i\partial_\mu\bar\phi^j-V_F,$$

$$V_F(\phi^i,\bar\phi^i)=e^{K/M_{\text{pl}}^2}\left[K^{i\bar J}D_iW\overline{D_jW}-\frac{3}{M_{\text{pl}}^2}|W|^2\right]$$

$$G_{MN} \text{d}X^M \text{d}X^N = e^{-6u(x)} g_{\mu\nu} \text{d}x^\mu \text{d}x^\nu + e^{2u(x)} \hat{g}_{mn} \text{d}y^m \text{d}y^n,$$

$$\int_{X_6} \text{d}^6y \sqrt{\hat{g}} \equiv \mathcal{V}$$



$$S_{\text{EH}}^{(10)}=\frac{1}{2\kappa^2}\int \;\; {\rm d}^{10}X \sqrt{-G} e^{-2\Phi} R_{10}$$

$$\bar g_{MN}=e^{2\omega(x)}g_{MN}$$

$$e^{2\omega}\bar R=R-2(D-1)\nabla^2\omega-(D-2)(D-1)g^{MN}\nabla_M\omega\nabla_N\omega.$$

$$e^{2\omega}\bar\nabla^2=\nabla^2+(D-2)g^{MN}\nabla_M\omega\nabla_N.$$

$$S_{\text{EH}}^{(10)}=\frac{1}{2\kappa^2}\int \;\; {\rm d}^4x \sqrt{-g} \int_{X_6} \;\; {\rm d}^6y \sqrt{\hat g} e^{-2\Phi} \big(R_4+e^{-8u}\hat R_6+12\partial_\mu u\partial^\mu u\big)$$

$$S_{\text{EH}}^{(4)}=\frac{M_{\text{pl}}^2}{2}\int \;\; {\rm d}^4x \sqrt{-g}R_4$$

$$M_{\text{pl}}^2\equiv\frac{\mathcal{V}}{g_s^2\kappa^2}$$

$$K=-3\mathrm{ln}\left(T+\bar{T}\right)$$

$${\bf Métrica de Moduli - Kähler.}$$

$$J\equiv i g_{i\bar J}\,{\rm d} z^i\wedge\,{\rm d}\bar z^{\bar J}$$

$$J=t^I(x)\omega_I$$

$$\delta g_{ij} = \frac{i}{\|\Omega\|^2} \zeta^A(x) (\chi_A)_{i\overline{l}j} \Omega^{\overline{l}\overline{j}}{}_j$$

$$\begin{gathered}B_2=B_2(x)+b^I(x)\omega_I,\\C_2=C_2(x)+c^I(x)\omega_I,\\C_4=\vartheta^I(x)\tilde{\omega}_I.\end{gathered}$$

$$\mathcal{O}=(-1)^{F_L}\Omega_{ws}\sigma$$

$$H^{1,1}=H^{1,1}_+\oplus H^{1,1}_-$$

$$J=t^i(x)\omega_i$$

$$\begin{gathered}B_2=b^\alpha(x)\omega_\alpha,\\C_2=c^\alpha(x)\omega_\alpha,\\C_4=\vartheta^i(x)\tilde{\omega}_i.\end{gathered}$$

$$\tau = C_0 + ie^{-\Phi}$$

$$G_\alpha \equiv c_\alpha - \tau b_\alpha$$

$$\mathcal{V}=\frac{1}{6}\int_{X_6}J\wedge J\wedge J=\frac{1}{6}c_{ijk}t^it^jt^k,$$

$$T_i\equiv\frac{1}{2}c_{ijk}t^jt^k+i\vartheta_i+\frac{1}{4}e^\Phi c_{i\alpha\beta}G^\alpha(G-\bar G)^\beta,$$



$$T_i=\frac{1}{2}c_{ijk}t^jt^k+i\vartheta_i$$

$$\tau_i = \frac{\partial \mathcal{V}}{\partial t^i} = \frac{1}{2} c_{ijk} t^j t^k,$$

$$T_i=\tau_i+i\vartheta_i$$

$${\bf Simetrías de Peccei-Quinn (PQ).}$$

$$a\mapsto a+\text{ const.}$$

$$b_I=\frac{1}{\alpha'}\int_{\Sigma_2^I}B_2,c_I=\frac{1}{\alpha'}\int_{\Sigma_2^I}C_2,\vartheta_I=\frac{1}{(\alpha')^2}\int_{\Sigma_4^I}C_4$$

$$\int_{\Sigma_2^I}\omega^J=\alpha'\delta_I^J,\int_{\Sigma_4^I}\tilde{\omega}^J=(\alpha')^2\delta_I^J$$

$$S_\sigma \supset -\frac{1}{4\pi\alpha'}\int_{\Sigma_2} {\rm d}^2\sigma \epsilon^{ab}\partial_a X^M\partial_b X^N B_{MN}(X),$$

$$S_\sigma \supset -\frac{1}{2\pi\alpha'}\int_{\Sigma_2} B_2 \equiv -\frac{b}{2\pi},$$

$$B_{MN}(X)=B_{MN}\big(X_{(0)}\big)+X^P\partial_P B_{MN}\big(X_{(0)}\big)+\cdots$$

$$-\frac{1}{4\pi\alpha'}\int_{\Sigma_2} {\rm d}^2\sigma \partial_a \left(\epsilon^{ab}X^M\partial_b X^N B_{MN}(X_{(0)})\right)$$

$$S_{\rm inst}=\exp\left(-\frac{1}{2\pi\alpha'}\int_{\Sigma_2}\left(J+iB_2\right)\right)\propto\exp\left(-i\frac{b}{2\pi}\right)$$

$$S_{\mathrm{ED}p}=\exp\left(-T_p\mathrm{Vol}(\Sigma_p)\right),$$

$$\mathcal{L}(a)=-\frac{1}{2}f^2(\partial a)^2-\Lambda^4[1-\cos{(a/2\pi)}]+\cdots,$$

$$B_2=b_\alpha(x)\omega^\alpha.$$

$$\frac{1}{2(2\pi)^7g_s^2(\alpha')^4}\int\;\;{\rm d}^{10}X|\;{\rm d} B_2|^2\supset\frac{1}{2}\int\;\;{\rm d}^4x\sqrt{-g}\gamma^{\alpha\beta}(\partial^\mu b_\alpha\partial_\mu b_\beta),$$

$$\gamma^{\alpha\beta}\equiv\frac{1}{6(2\pi)^7g_s^2(\alpha')^4}\int_{X_6}\omega^\alpha\wedge\star_6\omega^\beta.$$

$$\frac{f^2}{M_{\rm pl}^2}\approx\frac{1}{6}\frac{\alpha'^2}{L^4}$$



Compactificaciones de Calabi-Yau.

$$\mathrm{d} s^2 = e^{2A(y)} \eta_{\mu\nu} \mathrm{d} x^\mu \mathrm{d} x^\nu + e^{-2A(y)} g_{mn} \; \mathrm{d} y^m \; \mathrm{d} y^n$$

$$\tilde F_5=(1+\star_{10})\mathrm{d}\alpha(y)\wedge \mathrm{d} x^0\wedge \mathrm{d} x^1\wedge \mathrm{d} x^2\wedge \mathrm{d} x^3$$

$$\nabla^2e^{4A}=\frac{e^{8A}}{2{\rm Im}(\tau)}|G_3|^2+e^{-4A}(|\partial\alpha|^2+|\partial e^{4A}|^2)+2\kappa^2e^{2A}\mathcal{J}_{\rm loc}$$

$$\mathcal{J}_{\rm loc}\equiv\frac{1}{4}\biggl(\sum_{M=4}^9T^M{}_M-\sum_{M=0}^3T^M{}_M\biggr)_{\rm loc}$$

$$\mathrm{d}\tilde F_5=H_3\wedge F_3+2\kappa^2T_3\rho_3^{\rm loc}$$

$$\frac{1}{2\kappa^2 T_3}\int_{X_6} H_3\wedge F_3 + Q_3^{\rm loc}=0$$

$$\nabla^2(e^{4A}-\alpha)=\frac{e^{8A}}{24{\rm Im}(\tau)}|iG_3-\star_6G_3|^2+e^{-4A}|\partial(e^{4A}-\alpha)|^2+2\kappa^2e^{2A}(\mathcal{J}_{\rm loc}-\mathcal{Q}_{\rm loc})$$

$$\mathcal{J}_{\rm loc}\geq \mathcal{Q}_{\rm loc}$$

$$\star_6 G_3=iG_3,$$

$$e^{4A}=\alpha$$

$$V_{\rm flux}=\frac{1}{2\kappa^2}\int\;\;\mathrm{d}^{10}X\sqrt{-G_E}\biggl[-\frac{|G_3|^2}{2{\rm Im}(\tau)}\biggr]$$

$$K_0=-2\text{ln}\left(\mathcal{V}\right)-\text{ln}\left(-i(\tau-\bar{\tau})\right)-\text{ln}\left(-i\int\;\Omega\wedge\bar{\Omega}\right)$$

$$W_0=\frac{c}{\alpha'}\int\;\;G_3\wedge\Omega$$

$$V_F=e^{K_0}\left[K_0^{I\bar J}D_IW_0\overline{D_JW_0}-3|W_0|^2\right]$$

$$D_I W_0 \equiv \partial_I W_0 + (\partial_I K) W_0 = 0$$

$$\sum_{I,J=T_i}K_0^{I\bar J}\partial_I K_0\partial_{\bar J}K_0=3.$$

$$V_F=e^{K_0}\sum_{I,J\neq T_i}K_0^{I\bar J}D_IW_0\overline{D_JW_0}$$

$$K(T,\bar{T},z_\alpha,\bar{z}_\alpha)=-3\ln\,[T+\bar{T}-\gamma k(z_\alpha,\bar{z}_\alpha)],$$



Correcciones perturbativas Kaluza-Klein (KK) y singularidad.

$$K = -2\ln \left[\mathcal{V} + \frac{\xi}{2g_s^{3/2}} \right], \xi \equiv -\frac{\chi(X_6)\zeta(3)}{2(2\pi)^3}$$

$$\delta K_{(g_s)} = \delta K_{(g_s)}^{\text{KK}} + \delta K_{(g_s)}^{\text{W}}$$

$$\delta K_{(g_s)}^{\text{KK}} = -\frac{1}{128\pi^2} \sum_{i=1}^3 \frac{\mathcal{E}_i^{\text{KK}}(\zeta, \bar{\zeta})}{\text{Re}(\tau)\tau_i},$$

$$\delta K_{(g_s)}^{\text{W}} = -\frac{1}{128\pi^2} \sum_{i=1}^3 \left. \frac{\mathcal{E}_i^{\text{W}}(\zeta, \bar{\zeta})}{\tau_j \tau_k} \right|_{j \neq k \neq i}$$

$$S = \frac{1}{2g_7^2} \int_{\Sigma_4} d^4\sigma \sqrt{g_{\text{ind}}} e^{-4A(y)} \cdot \int d^4x \sqrt{-g} \text{Tr}[F_{\mu\nu}F^{\mu\nu}]$$

$$g_7^2 = 2(2\pi)^5(\alpha')^2$$

$$\frac{1}{g^2} = \frac{T_3 \mathcal{V}_4}{8\pi^2},$$

$$\mathcal{V}_4 \equiv \int_{\Sigma_4} d^4\sigma \sqrt{g_{\text{ind}}} e^{-4A(y)}$$

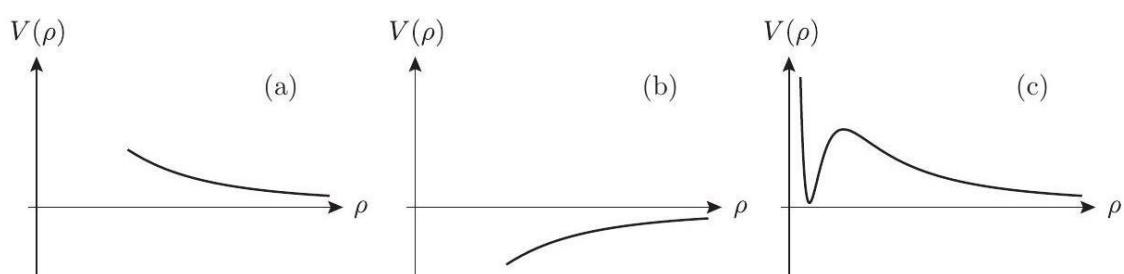
$$|W_{\lambda\lambda}| = 16\pi^2 M_{\text{UV}}^3 \exp\left(-\frac{1}{N_c} \frac{8\pi^2}{g^2}\right) \propto \exp\left(-\frac{T_3 \mathcal{V}_4}{N_c}\right).$$

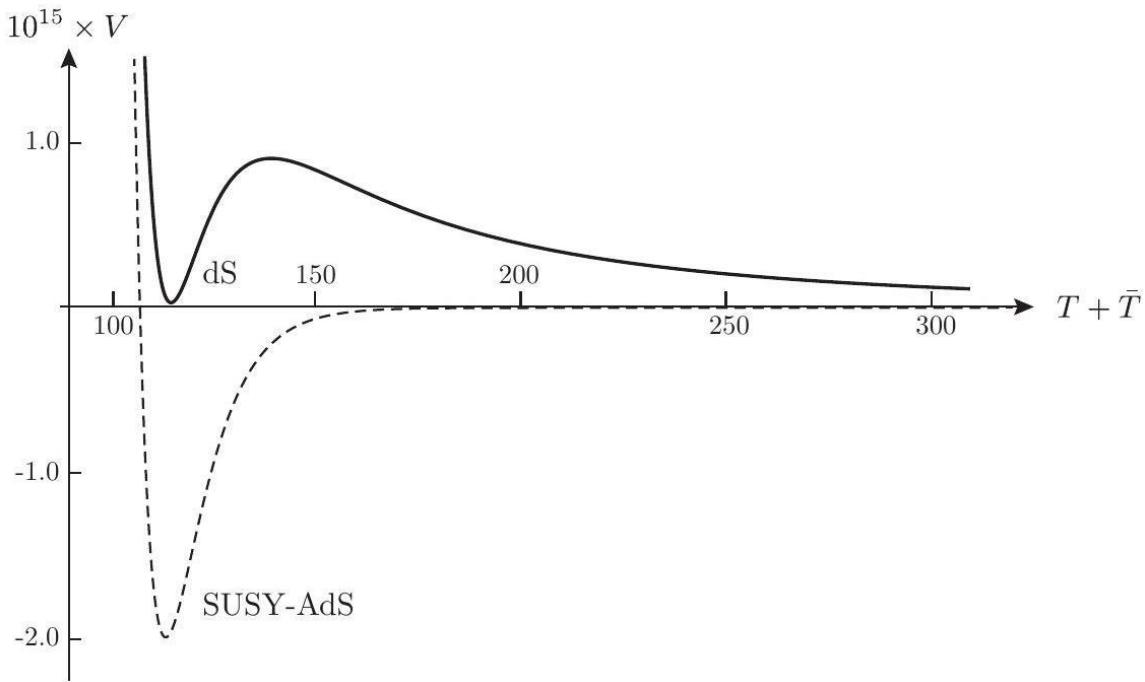
$$W_{\lambda\lambda} = \mathcal{A}e^{-\alpha T}$$

$$W_{\text{ED3}} = \mathcal{A}e^{-\alpha T},$$

$$\chi(\mathcal{O}_D) \equiv \sum_{i=0}^3 (-1)^i h^{0,i}(D) = 1$$

$$h^{0,1}(D) = h^{0,2}(D) = h^{0,3}(D) = 0$$





Figuras 39 y 40. Fluctuaciones del centro de masa – energía de la partícula supermasiva.

$$W = W_0 + \sum_{i=1}^{h_+^{1,1}} \mathcal{A}_i e^{-a_i T_i} + \dots$$

$$V_{(np)} = e^K K^{j\bar{i}} \left[a_j \mathcal{A}_j a_{\bar{i}} \bar{\mathcal{A}}_{\bar{i}} e^{-(a_j T_j + a_{\bar{i}} \bar{T}_{\bar{i}})} - (a_j \mathcal{A}_j e^{-a_j T_j} \bar{W} \partial_{\bar{i}} K + a_{\bar{i}} \bar{\mathcal{A}}_{\bar{i}} e^{-a_{\bar{i}} \bar{T}_{\bar{i}}} W \partial_j K) \right]$$

$$V_{(np)} = \frac{a \mathcal{A} e^{-a(T+\bar{T})}}{2(T+\bar{T})^2} \left[\left(1 + \frac{T+\bar{T}}{3} \right) a \mathcal{A} e^{-a(T+\bar{T})} + W_0 \right]$$

$$W_0 = -\mathcal{A} e^{-a(T+\bar{T})_*} \left(1 + \frac{2}{3} a(T+\bar{T})_* \right)$$

$$W_0 = W_0|_{Z=0} + \ell_A Z^A + m_{AB} Z^A Z^B + \dots$$

$$\delta V_{(\alpha')} = 3\hat{\xi} e^K \frac{(\hat{\xi}^2 + 7\hat{\xi}\mathcal{V} + \mathcal{V}^2)}{(\mathcal{V} - \hat{\xi})(2\mathcal{V} + \hat{\xi})^2} W_0^2 \approx \frac{3}{4} \hat{\xi} W_0^2 \frac{1}{\mathcal{V}^3},$$

$$V_{(np)} + \delta V_{(\alpha')} = e^K \left\{ K^{j\bar{i}} \left[a_j \mathcal{A}_j a_{\bar{i}} \bar{\mathcal{A}}_{\bar{i}} e^{-(a_j T_j + a_{\bar{i}} \bar{T}_{\bar{i}})} - (a_j \mathcal{A}_j e^{-a_j T_j} \bar{W} \partial_{\bar{i}} K + a_{\bar{i}} \bar{\mathcal{A}}_{\bar{i}} e^{-a_{\bar{i}} \bar{T}_{\bar{i}}} W \partial_j K) \right] + \frac{3}{4} \hat{\xi} W_0^2 \frac{1}{\mathcal{V}} \right\}$$

$$\mathcal{V} \rightarrow \infty, \text{ with } a_s \tau_s = \ln \mathcal{V}$$

$$\{T_i\} = \{T_b^\rho\} \cup \{T_s^r\},$$

$$\mathcal{V} \rightarrow \infty, \text{ with } a_s^r \tau_s^r = \ln \mathcal{V} \text{ for all } r = 1, \dots, N_s$$



$$\mathcal{V}=\alpha\tau_b^{3/2}-p^{(3/2)}(\tau_s^r),$$

$$\mathcal{V} = \alpha \Bigg(\tau_b^{3/2} - \sum_{r=1}^{N_s} \lambda_r (\tau_s^r)^{3/2} \Bigg)$$

$$\rho_{\mathcal{S}}=\frac{D}{\mathcal{V}^\alpha}$$

$$\frac{1}{(2\pi)^2\alpha'}\int_{S^3}F_3\equiv M$$

$$R^4=4\pi g_s N(\alpha')^2,$$

$$\nabla^2\Phi=4\pi e\delta(\boldsymbol{x}),$$

$$\Phi(\boldsymbol{x}')=-\frac{e}{|\boldsymbol{x}-\boldsymbol{x}'|}.$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-\frac{1}{4}F_{\mu\nu}F^{\mu\nu}-\frac{\phi}{f}F_{\mu\nu}F_{\rho\sigma}\epsilon^{\mu\nu\rho\sigma}$$

$$\boldsymbol{E}\cdot\boldsymbol{B}=-\frac{eB}{r^2}\cos~\theta$$

$$\nabla^2\phi=\frac{eB}{fr^2}\cos~\theta$$

$${\mathrm d} G_- = - {\mathrm d} \Bigl(\frac{\Phi_- G_+}{\Phi_+} \Bigr)$$

$$V=V_{\mathfrak{F}}(\zeta_1,\ldots,\zeta_{h^{2,1}})$$

$${\mathrm d} \mathcal{N}_{\min}(\zeta) \equiv \sum_i ~\delta(\zeta-\zeta_i)$$

$${\mathrm d} \mathcal{I}_{\min}(\zeta) \equiv \sum_i ~\delta(\zeta-\zeta_i)(-1)^{F_i}$$

$$\sum_{L\leq L_{\max}}{\mathrm d} \mathcal{I}_{\min}=\frac{(2\pi L_{\max})^{b_3}}{\pi^{b_3/2}b_3!}\det(-\mathcal{R}-\omega)$$



Gravedad cuántica y simetrías de gauge local.

$$\mathcal{N}_{dS}=\mathcal{N}_{\text{c.p.}}\times f_{dS}$$

$$\mathcal{N}_{\min} = \mathcal{N}_{\text{c.p.}} \times f_{\min}$$

$$\mathfrak{C}\equiv\bigcup_{\mathfrak{F}_\star}\left\{p_i^{(\mathfrak{F}_\star)}\right\},$$

$$\mathcal{N}_{\text{c.p.}}\gtrsim\mathcal{N}_{\mathfrak{F}}$$

$$V=e^K(F_a\bar F^a-3|W|^2)$$

$$\mathcal{H}=\begin{pmatrix} \partial^2_{a\bar b} V & \partial^2_{ab} V \\ \partial^2_{\bar a\bar b} V & \partial^2_{\bar a b} V \end{pmatrix}$$

$$f_{\min}=P(\lambda_1>0)$$

$$F_a\equiv \mathcal{D}_a W, Z_{ab}\equiv \mathcal{D}_a \mathcal{D}_b W, U_{abc}\equiv \mathcal{D}_a \mathcal{D}_b \mathcal{D}_c W$$

$$\mathcal{H}=\begin{pmatrix} Z_a^{\bar c}\bar Z_{\bar b\bar c}-F_a\bar F_{\bar b}-R_{a\bar b c\bar d}\bar F^cF^{\bar d}&U_{abc}\bar F^c-Z_{ab}\bar W\\ \bar U_{\bar a\bar b\bar c}F^{\bar c}-\bar Z_{\bar a\bar b}W&\bar Z_{\bar a}^cZ_{bc}-F_b\bar F_{\bar a}-R_{b\bar a c\bar d}\bar F^cF^{\bar d} \end{pmatrix}+\mathbb{1}(F^2-2|W|^2)$$

$$F\ll m_{\mathrm{susy}}\,M_{\mathrm{pl}}$$

$$F\gtrsim m_{\mathrm{susy}}\,M_{\mathrm{pl}}$$

$$F_{\mathrm{rms}}\sim Z_{\mathrm{rms}}\sim U_{\mathrm{rms}}$$

$$M=A+A^\dag$$

$$\rho(\lambda) = \frac{1}{2\pi N\sigma^2} \sqrt{4N\sigma^2 - \lambda^2}$$

$$M=AA^\dag$$

$$\rho(\lambda)=\frac{1}{2\pi N\sigma^2\lambda}\sqrt{(\eta_+-\lambda)(\lambda-\eta_-)}$$

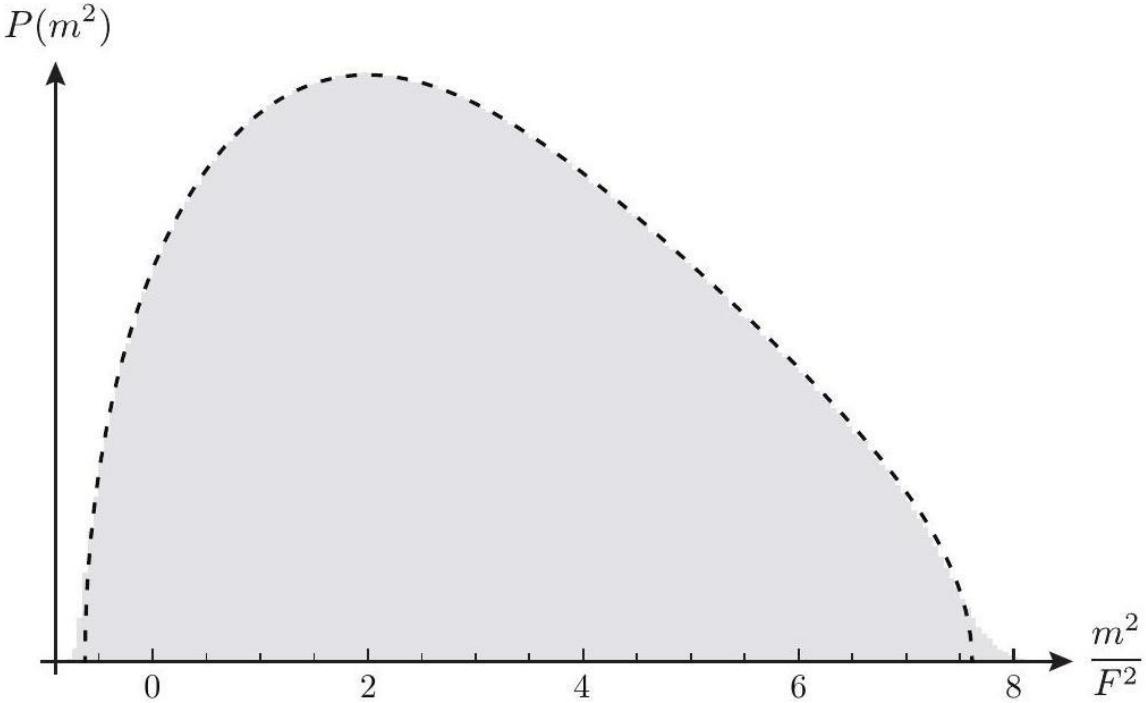
$$\mathcal{H}_Z\equiv\begin{pmatrix} Z_a^{\bar c}\bar Z_{\bar b\bar c}&0\\0&\bar Z_{\bar a}^cZ_{bc} \end{pmatrix}$$

$$\mathcal{H}_{WWW}\equiv\mathcal{H}_{\mathrm{Wigner}}+\mathcal{H}_{\mathrm{Wishart}}^{(I)}+\mathcal{H}_{\mathrm{Wishart}}^{(II)},$$

$$\rho(\mathcal{H}_{WWW})=\rho\!\left(\mathcal{H}_{\mathrm{Wigner}}\right)\boxplus\rho\!\left(\mathcal{H}_{\mathrm{Wishart}}^{(I)}\right)\boxplus\rho\!\left(\mathcal{H}_{\mathrm{Wishart}}^{(II)}\right).$$

$$f_>\equiv\frac{\int_0^\infty\rho(\lambda)}{\int_{-\infty}^\infty\rho(\lambda)}$$





$$P(\lambda_1 > \zeta) = \exp(-N^2\Psi(\zeta))$$

$$f_{\min} \equiv P(\lambda_1 > 0) = \exp(-cN^2),$$

$$f_{\min} \equiv P(\lambda_1 > 0) = \exp(-dN)$$

$$\begin{array}{l} K=K_l(\phi,\bar{\phi})+K_h(\Sigma,\bar{\Sigma})\\ W=W_l(\phi)+W_h(\Sigma)\end{array}$$

$$f_{\min} \equiv P(\lambda_1 > 0) = \exp(-cN_l^2),$$

$$S_{10}[\mathcal{C}] \mapsto S_4[\Phi(t)].$$

$$M_{\text{pl}} \sim g_s^{-1} (M_s/M_{\text{KK}})^3 M_s \gg M_s.$$

$$M_{\text{SUSY}} < H < M_{\text{KK}} < M_s < M_{\text{pl}}$$

$$\Delta \mathcal{L} \supset \frac{1}{M_{AB}^{\delta_A + \delta_B - 4}} \mathcal{O}_A^{(\delta_A)} \mathcal{O}_B^{(\delta_B)}$$

$$\Delta \mathcal{L} \supset \frac{V_0}{M_{AB}^2} \phi^2$$

$$\frac{M_{AB}}{M_{\text{pl}}} \lesssim g_s \left(\frac{\ell_s}{L} \right)^2$$

$$\frac{M_{AB}}{M_{\text{pl}}} \lesssim g_s \left(\frac{\ell_s}{L} \right)^{\frac{1}{2}p-1} \left(\frac{\ell_s}{S} \right)^{\frac{1}{2}(6-p)}$$



$$V(r)=2T_3\left(1-\frac{1}{2\pi^3}\frac{T_3g_s^2\kappa^2}{r^4}\right),$$

$$\eta \approx -\frac{10}{\pi^3} \frac{\mathcal{V}}{r^6},$$

$$\nabla_y^2\big(\delta e^{-4A(y_b;y)}\big)=-\mathcal{C}\left(\frac{\delta(y_b-y)}{\sqrt{g(y)}}-\bar\rho(y)\right)$$

$$\delta e^{-4A(y_b;y)}=\mathcal{C}\left(\mathcal{G}(y_b;y)-\int\,\mathrm{d}^6y'\sqrt{g}\mathcal{G}(y;y')\bar\rho(y')\right)$$

$$\nabla_{y'}^2\mathcal{G}(y;y')=\nabla_{y'}^2\mathcal{G}(y;y')=-\frac{\delta(y-y')}{\sqrt{g}}+\frac{1}{\mathcal{V}}.$$

$$\nabla_{y_b}^2\big(\delta e^{-4A(y_b;y)}\big)=-\mathcal{C}\left(\frac{\delta(y_b-y)}{\sqrt{g(y_b)}}-\frac{1}{\mathcal{V}}\right)$$

$$V(y_b)=2T_3e^{4A(y_b)}\approx 2T_3\big(1-\delta e^{-4A(y_b)}\big)$$

$${\rm Tr}(\eta)\approx -\frac{M_{\rm pl}^2}{T_3}\nabla_{y_b}^2\big(\delta e^{-4A(y_b;y)}\big)=-2$$

$$\mathcal{L} = -K_{\varphi\bar\varphi}\partial_\mu\varphi\partial^\mu\bar\varphi-e^{K/M_{\rm pl}^2}\bigg[K^{\varphi\bar\varphi}D_\varphi W\overline{D_\varphi W}-\frac{3}{M_{\rm pl}^2}|W|^2\bigg].$$

$$K=K(0)+K_{\varphi\bar\varphi}(0)\varphi\bar\varphi+\cdots,$$

$$\mathcal{L}\approx -\partial_\mu\phi\partial^\mu\bar\phi-V(0)\bigg(1+\frac{\phi\bar\phi}{M_{\rm pl}^2}+\cdots\bigg),$$

$$m_\phi^2=\frac{V(0)}{M_{\rm pl}^2}+\cdots=3H^2+\cdots\,\Rightarrow\,\eta=1+\cdots$$

$$K=-3\text{ln}\,[T+\bar T-\gamma k(z_\alpha,\bar z_\alpha)]\equiv -2\text{ln}\;\mathcal{V},$$

$$K=(\phi-\bar\phi)^2$$

$$\Delta\phi^2<\frac{1}{8\pi}\frac{M_{\rm s}^2}{g_{\rm s}}\Bigl(\frac{L}{\ell_s}\Bigr)^{p-1}.$$

$$M_{\rm pl}^2=\frac{1}{\pi}\frac{M_{\rm s}^2}{g_{\rm s}^2}\Bigl(\frac{L}{\ell_s}\Bigr)^6.$$

$$\frac{\Delta\phi^2}{M_{\rm pl}^2}<\frac{g_s}{8}\Bigl(\frac{\ell_s}{L}\Bigr)^{7-p}$$



$$\frac{\Delta\phi^2}{M_{\rm pl}^2}<\frac{g_s}{4\pi}\frac{L}{\ell_s}$$

$$\Delta \Phi = \sqrt{N} \Delta \phi$$

$$\delta M_{\rm pl}^2 \sim \frac{N}{16\pi^2} \Lambda_{\rm UV}^2$$

$$V^{1/4} \ll M_{\mathrm{KK}} \ll M_s \ll M_{\mathrm{pl}}$$

$${\mathcal L}_{\rm eff}[\Phi]={\mathcal L}_l[\Phi]+\sum_i~c_i\frac{{\mathcal O}_i[\Phi]}{\Lambda^{\delta_i-4}},$$

$$\textbf{Campos multiperturbativos y espacios anti - De Sitter.}$$

$$\mathcal{R}=f(\pi_\star,\psi_\star)$$

$$\mu_{p,q}=\frac{1}{2\pi\alpha'}\sqrt{(p-C_0q)^2+e^{-2\Phi}q^2}$$

$$\mathrm{d}s^2=e^{2A(r)}\eta_{\mu\nu}\mathrm{d}x^\mu\mathrm{d}x^\nu+e^{-2A(r)}\big(\mathrm{d}r^2+r^2\;\mathrm{d}\Omega_{S_5}^2\big)$$

$$e^{-4A(r)}=1+\frac{L^4}{r^4},\text{ with }\frac{L^4}{(\alpha')^2}=4\pi g_s N$$

$$\alpha(r)\equiv (\mathcal{C}_4)_{tx^i}=e^{4A(r)}$$

$$\mathrm{d}s_{AdS_5}^2=\frac{L^2}{r^2}\;\mathrm{d}r^2+\frac{r^2}{L^2}\eta_{\mu\nu}\mathrm{d}x^\mu\mathrm{d}x^\nu,$$

$$S_{\mathrm{D}3}=-T_3\int\;\;\mathrm{d}^4\sigma\sqrt{-\mathrm{det}(G^E_{ab})}+\mu_3\int\;\;\mathcal{C}_4$$

$$\mathcal{L}=-T_3e^{4A(r)}\sqrt{1+e^{-4A(r)}g^{\mu\nu}\partial_\mu r\partial_\nu r}+T_3\alpha(r)$$

$$\mathcal{L}\approx-\frac{1}{2}(\partial\phi)^2-T_3\big(e^{4A(\phi)}-\alpha(\phi)\big),$$

$$\sum_{A=1}^4 z_A^2=0$$

$$x\cdot x=\frac{1}{2}\rho^2,y\cdot y=\frac{1}{2}\rho^2,x\cdot y=0$$

$$T^{1,1}=[SU(2)\times SU(2)]/U(1)$$

$$d\Omega_{T^{1,1}}^2\equiv\frac{1}{9}\Bigg(\,\mathrm{d}\psi+\sum_{i=1}^2\cos\,\theta_i\,\mathrm{d}\phi_i\Bigg)^2+\frac{1}{6}\sum_{i=1}^2\,\big(\,\mathrm{d}\theta_i^2+\sin^2\,\theta_i\,\mathrm{d}\phi_i^2\big)$$



$$\mathrm{d} s^2 = \mathrm{d} r^2 + r^2\; \mathrm{d}\Omega_{T^{1,1}}^2$$

$$\mathrm{d} s^2 = k_{\alpha\bar{\beta}}\mathrm{d} z^\alpha\mathrm{d}\overline{z^\beta}$$

$$k(z_\alpha,\bar{z}_\alpha)=\frac{3}{2}\biggl(\sum_{A=1}^4|z^A|^2\biggr)^{2/3}$$

$$\sum_{A=1}^4\,z_A^2=\varepsilon^2$$

$$\begin{array}{l}x\cdot x-y\cdot y=\varepsilon^2\\ x\cdot x+y\cdot y=\rho^2\end{array}$$

$$\mathrm{d} s^2=e^{2A(r)}\eta_{\mu\nu}\mathrm{d} x^\mu\mathrm{d} x^\nu+e^{-2A(r)}\bigl(\mathrm{d} r^2+r^2\;\mathrm{d}\Omega_{T^{1,1}}^2\bigr),$$

$$e^{-4A(r)}=1+\frac{L^4}{r^4}~~\text{with}~~L^4\equiv\frac{27\pi}{4}g_sN(\alpha')^2$$

$$\frac{1}{(2\pi)^2\alpha'}\int_A F_3=M ~~\text{and}~~ \frac{1}{(2\pi)^2\alpha'}\int_B H_3=K$$

$$\mathrm{d} s^2=e^{2A(r)}\eta_{\mu\nu}\mathrm{d} x^\mu\mathrm{d} x^\nu+e^{-2A(r)}\mathrm{d}\tilde{s}^2,$$

$$e^{-4A(r)}=\frac{L^4}{r^4}\Big(1+\frac{3g_sM}{8\pi K}+\frac{3g_sM}{2\pi K}\ln\frac{r}{r_{\rm UV}}\Big)$$

$$L^4\equiv\frac{27\pi}{4}g_sN(\alpha')^2,N\equiv MK$$

$$e^{A_{\text{IR}}}=\exp\left(-\frac{2\pi K}{3g_s M}\right)$$

$$\mathcal{V}_{\mathcal{T}}\equiv\int\;\mathrm{d}\Omega_{T^{1,1}}^2\int_{r_{\text{IR}}}^{r_{\text{UV}}}r^5\;\mathrm{d} r e^{-4A(r)}=2\pi^4g_sN(\alpha')^2r_{\text{UV}}^2$$

$$M_{\mathrm{pl}}^2>\frac{N}{4}\frac{r_{\mathrm{UV}}^2}{(2\pi^3)g_s(\alpha')^2}$$

$$\Delta\phi^2 < T_3r_{\mathrm{UV}}^2=\frac{r_{\mathrm{UV}}^2}{(2\pi^3)g_s(\alpha')^2}$$

$$\frac{\Delta\phi}{M_{\mathrm{pl}}}\leq\frac{2}{\sqrt{N}}$$

$$V(\phi)=T_3\big(e^{4A(\phi)}-\alpha(\phi)\big).$$



$$V_{\mathcal{C}}(\phi)=D_0\left(1-\frac{27}{64\pi^2}\frac{D_0}{\phi^4}\right)$$

$$D_0 \equiv 2 T_3 e^{4A(r_{\rm IR})}$$

$$V_{\mathcal{R}}(\phi) = \frac{1}{12} R \phi^2.$$

$$V(\phi)=V_{\mathcal{C}}(\phi)+V_{\mathcal{R}}(\phi)+\cdots\approx V_0+H^2\phi^2+\cdots\,\Rightarrow\,\eta\approx\frac{2}{3}+\cdots$$

$$V_F=e^K\big[K^{I\bar J}D_IW\overline{D_JW}-3|W|^2\big]$$

$$K=-2\mathrm{ln}\left(\mathcal{V}\right)$$

$$\mathcal{V}=\left(T+\bar{T}-\gamma k(z_\alpha,\bar{z}_\alpha)\right)^{3/2}$$

$$\gamma\equiv\frac{T_3}{6}(T+\bar{T})_{\text{IR}}$$

$$K(Z^I,\bar Z^I)=-3\mathrm{ln}\left[T+\bar{T}-\gamma k(z_\alpha,\bar{z}_\alpha)\right]\equiv -3\mathrm{ln}\left[U(Z^I,\bar Z^I)\right].$$

$$V_F(T,z_\alpha)=\frac{1}{3U^2}\big[\big(T+\bar{T}+\gamma\big(k_\gamma k^{\gamma\bar\delta}k_{\bar\delta}-k\big)\big)\big|W_{,T}\big|^2-3\big(\bar{W}W_{,T}+c.c.\,\big)$$

$$+\underbrace{\big(k^{\alpha\bar\delta}k_{\bar\delta}W_{,T}W_{,\alpha}+c.c.\,\big)+\frac{k^{\alpha\bar\beta}}{\gamma}W_{,\alpha}W_{,\beta}}_{\Delta V_F}$$

$$V_F(r)\approx \frac{V_0}{\left(1-\frac{1}{6}\phi^2\right)^2}\approx V_0+\frac{1}{3}\frac{V_0}{M_{\text{pl}}^2}\phi^2,$$

$$|\Delta W|\propto \exp\left(-\frac{2\pi}{N_c}\mathcal{V}_4\right)$$

$$f(z_\alpha)=0$$

$$W(T,z_\alpha)=W_0+\mathcal{A}(z_\alpha)e^{-aT}, a\equiv \frac{2\pi}{N_c}$$

$$\mathcal{A}(z_\alpha)=\mathcal{A}_0\bigg(\frac{f(z_\alpha)}{f(0)}\bigg)^{1/N_c}$$

$$f(z_1)=\mu-z_1$$

$$V_F(\phi)\approx V_0+\cdots+\lambda\phi^{3/2}+\frac{V_0}{M_{\text{pl}}^2}\phi^2+\cdots$$



$$\delta\Phi(r)=\delta\Phi(r_{\rm UV})\left(\frac{r}{r_{\rm UV}}\right)^\Delta.$$

$$\Phi_-\equiv e^{4A}-\alpha$$

$${\rm d} s^2 = e^{2A(y)} g_{\mu\nu} {\rm d} x^\mu {\rm d} x^\nu + e^{-2A(y)} g_{mn}\;{\rm d} y^m \;{\rm d} y^n,$$

$$\nabla^2\Phi_- = R_4 + \frac{g_s}{96}|\Lambda|^2 + e^{-4A}|\nabla\Phi_-|^2 + \mathcal{S}_{\text{loc}}$$

$$\Lambda \equiv \Phi_+ G_- + \Phi_- G_+$$

$$G_\pm \equiv (\star_6 \pm i) G_3 ~~\text{and}~~ \Phi_\pm \equiv e^{4A} \pm \alpha$$

$${\rm d}\Lambda+\frac{i}{2}\frac{{\rm d}\tau}{{\rm Im}\tau}\wedge(\Lambda+\bar\Lambda)=0$$

$$V(x,\Psi)=V_0+V_{\mathcal C}(x)+V_{\mathcal R}(x)+V_{\mathcal B}(x,\Psi),$$

$$V_{\mathcal C}(x)=D_0\left(1-\frac{27}{64\pi^2}\frac{D_0}{T_3^2r_{\rm UV}^4}\frac{1}{x^4}\right)$$

$$V_{\mathcal R}(x)=\frac{1}{3}\mu^4x^2+\cdots,~\text{where}~\mu^4\equiv(V_0+D_0)\frac{T_3r_{\rm UV}^2}{M_{\rm pl}^2}$$

$$V_{\mathcal B}(x,\Psi)=\mu^4\sum_{LM}c_{LM}x^{\Delta(L)}f_{LM}(\Psi)$$

$$\nabla^2\Phi_h=0$$

$$\nabla^2\Phi_f=\frac{g_s}{96}|\Lambda|^2$$

$$\Delta_h(L)\equiv -2\sqrt{H(j_1,j_2,R)+4}$$

$$H(j_1,j_2,R)\equiv 6\left[j_1(j_1+1)+j_2(j_2+1)-\frac{1}{8}R^2\right]$$

$$\Delta_h=\frac{3}{2},2,3,\sqrt{28}-2,\cdots$$

$$\Delta_f(L)=\delta_i(L)+\delta_j(L)-4$$

$$\begin{gathered}\delta_1(L)\equiv -1+\sqrt{H(j_1,j_2,R+2)+4},\\\delta_2(L)\equiv \sqrt{H(j_1,j_2,R)+4},\\\delta_3(L)\equiv 1+\sqrt{H(j_1,j_2,R-2)+4}.\end{gathered}$$

$$\Delta_f=1,2,\frac{5}{2},\sqrt{28}-\frac{5}{2},\cdots$$

$$\Delta=\{\Delta_h,\Delta_f\}=1,\frac{3}{2},\sqrt{28}-\frac{5}{2},3,\sqrt{28}-2,\frac{7}{2},\sqrt{28}-\frac{3}{2},\cdots$$



$$\mathcal{M} = \begin{pmatrix} A\bar{A} + B\bar{B} & C \\ \bar{C} & \bar{A}A + \bar{B}B \end{pmatrix}$$

$$\tau_{\rm therm}\,\ll\tau_{\rm graviton}\,\ll\tau_{\rm tunnel}\,,$$

$$V(\phi) \approx V_0 \left[1 + \lambda_0 \frac{\phi}{M_{\rm pl}} + \frac{1}{3!} \mu_0 \frac{\phi^3}{M_{\rm pl}^3} + \cdots \right],$$

$$n_s - 1 \approx - \frac{4\pi}{N_{\rm tot}} \cot \left(\pi \frac{N_\star}{N_{\rm tot}} \right) \approx - \frac{4}{N_\star} \Bigg(1 + \mathcal{O} \left(\frac{N_\star^2}{N_{\rm tot}^2} \right) \Bigg),$$

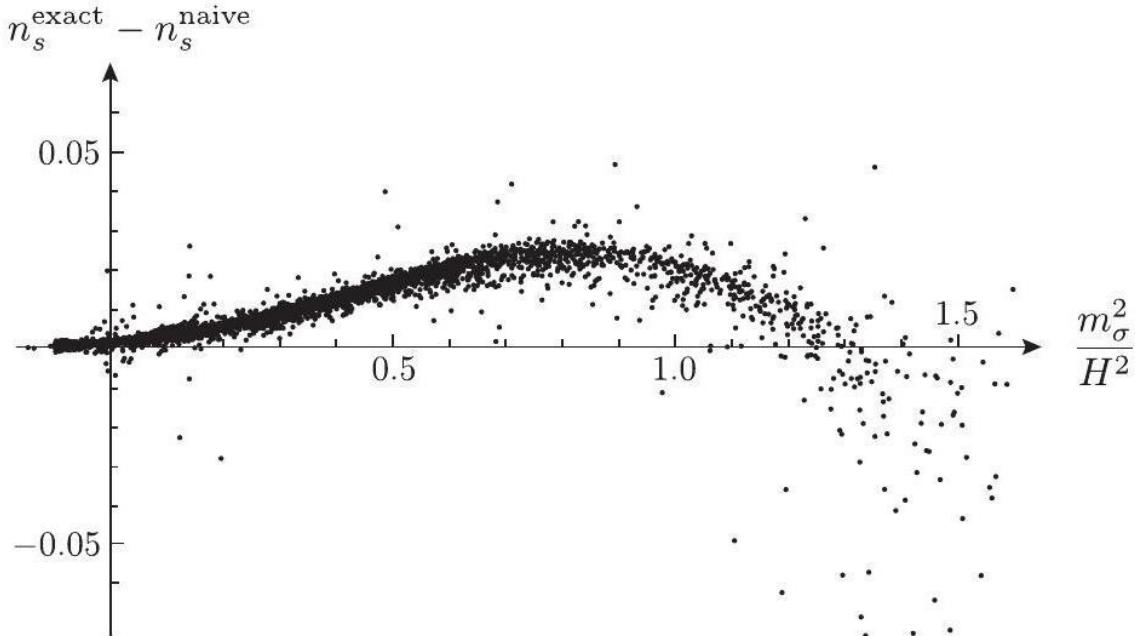
$$N_{\rm tot}=\int_{-\infty}^\infty \frac{1}{\sqrt{2\epsilon}}\frac{{\rm d}\phi}{M_{\rm pl}}=\pi\sqrt{\frac{2}{\lambda_0\lambda_1}}.$$

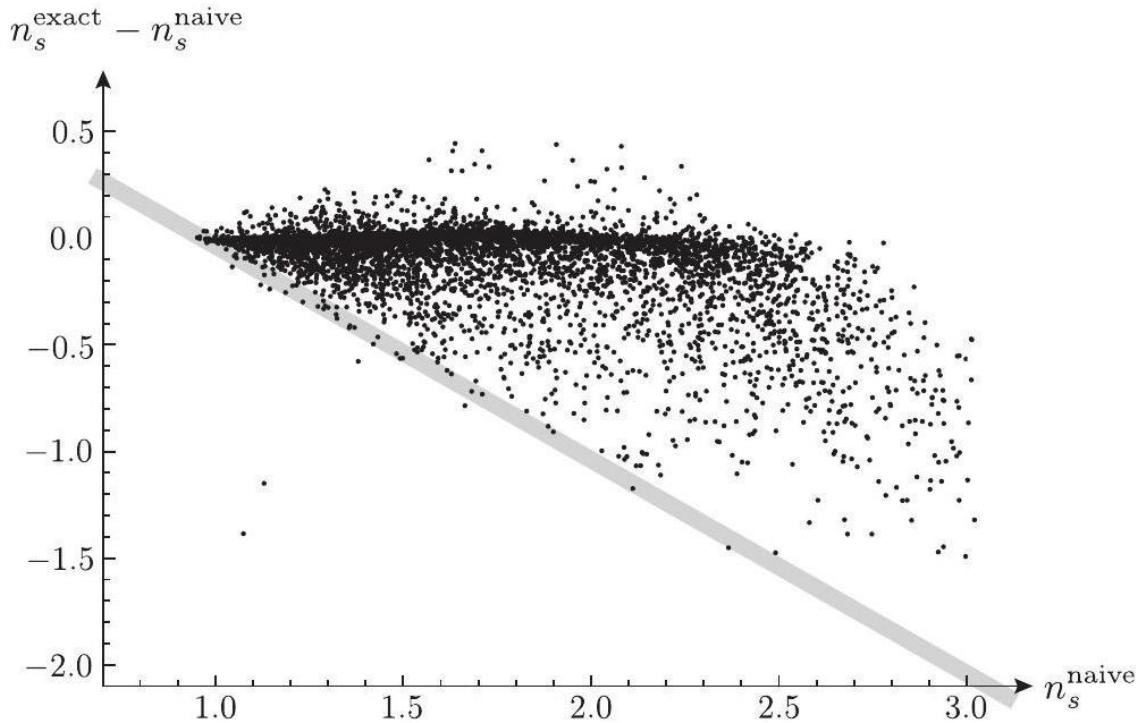
$$N_e(\phi)=\int_{\phi_{\rm end}}^{\phi}\frac{1}{\sqrt{2\epsilon}}\frac{{\rm d}\phi}{M_{\rm pl}}=\frac{N_{\rm tot}}{\pi}\arctan\left(\frac{\eta(\phi)N_{\rm tot}}{2\pi}\right)\bigg|_{\phi_{\rm end}}^{\phi}$$

$$\alpha_s=-\frac{4\pi^2}{N_{\rm tot}^2}\sin^{-2}\left(\pi\frac{N_\star}{N_{\rm tot}}\right) \approx -\frac{4}{N_\star^2} \Bigg(1 + \mathcal{O} \left(\frac{N_\star^2}{N_{\rm tot}^2} \right) \Bigg).$$

$$r<\frac{4}{N}\times 0.01\ll 0.01$$

$$P(N_e) = P(N_\star) \left(\frac{N_\star}{N_e}\right)^3$$





$$T_{(p,q)} \approx \frac{e^{2A_{\text{IR}}}}{2\pi\alpha'} \sqrt{\frac{q^2}{g_s^2} + \left(\frac{bM}{\pi}\right)^2 \sin^2\left(\frac{\pi(p-qC_0)}{M}\right)}$$

$$V_D(\phi)=\frac{g^2\xi^2}{2}\biggl(1+\frac{g^2}{16\pi^2}U(x)\biggr),$$

$$U(x)\equiv (x^2+1)^2\ln{(x^2+1)}+(x^2-1)^2\ln{(x^2-1)}-4x^4\ln{(x)}-4\ln{(2)},$$

$$K(Z^I,\bar Z^I)=-3\mathrm{ln}\left[T+\bar T-\gamma k(z_\alpha,\bar z_\alpha)\right]$$

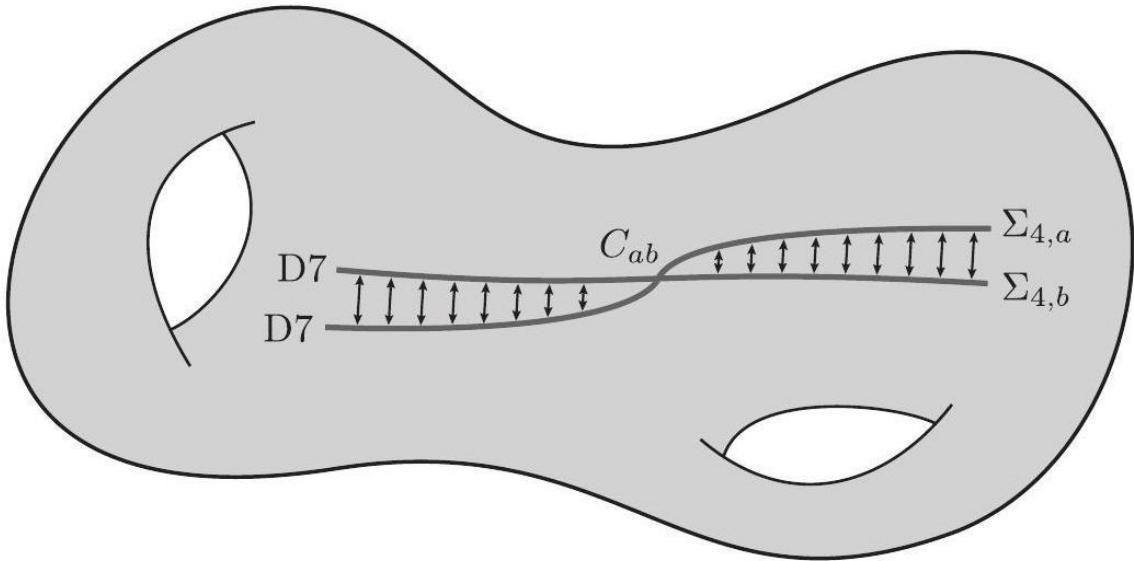
$$k=\frac{1}{2}(z_1+\bar{z}_1)^2=x^2,$$

$$W=W_0+\left[\vartheta_1\bigl(\sqrt{2\pi}(z_1+\mu),\zeta\bigr)\vartheta_1\bigl(\sqrt{2\pi}(z_1-\mu),\zeta\bigr)\right]^{-1/N_c}e^{-2\pi T/N_c}$$

$$W = \Lambda^{3-1/N_c} m^{1/N_c},$$

$$V(\phi)=V_D(\phi)-\frac{m^2}{2}\phi^2+\frac{\lambda}{4}\phi^4$$





$$V(\phi) = V_D(\phi) + V_F(\phi)$$

$$V_D(\phi)=V_0\big(1+\alpha{\rm ln}\;(\phi/\phi_0)\big)$$

$$K \supset -\ln \left(S+\bar{S}-\frac{(Y+\bar{Y})^2}{T+\bar{T}}\right)$$

$$M_{\rm pl} \propto \sqrt{{\rm Vol}(K3)L_1L_2},$$

$$\Delta\phi_1\propto\sqrt{L_1/L_2}$$

$${\rm d}s^2=\Big(\frac{r}{L}\Big)^2\,\eta_{\mu\nu}{\rm d}x^\mu{\rm d}x^\nu+\Big(\frac{L}{r}\Big)^2\,\big({\rm d}r^2+r^2\,{\rm d}\Omega^2_{S^5}\big),$$

$$\mathcal{L}=-\frac{\phi^4}{\lambda}\Biggl(\sqrt{1+\frac{\lambda}{\phi^4}(\partial\phi)^2}-1\Biggr)-V(\phi),$$

$$\lambda \equiv T_3 L^4 = \frac{1}{2\pi^2} N$$

$$\mathcal{L}=-T(\phi)\Biggl(\sqrt{1+\frac{(\partial\phi)^2}{T(\phi)}}-1\Biggr)-V(\phi),$$

$${\rm d}s^2=e^{2A(r)}\eta_{\mu\nu}{\rm d}x^\mu{\rm d}x^\nu+e^{-2A(r)}\big({\rm d}r^2+r^2\,{\rm d}\Omega^2_{X_5}\big),$$

$$e^{-4A(r)}\approx \frac{L^4}{r^4}\;\;\text{with}\;\;L^4\equiv \frac{4\pi^4 g_s}{{\rm Vol}(X_5)}N(\alpha')^2$$

$$\frac{\Delta \phi}{M_{\rm pl}} \leq \frac{2}{\sqrt{N}}$$

$$V(\phi)=V_0-\frac{1}{2}\beta H^2\phi^2,$$

$$\gamma \equiv \left(1 - \frac{\dot{\phi}^2}{T(\phi)}\right)^{-1/2}$$

$$\dot{\phi}^2 < T(\phi).$$

$$\begin{array}{l} \rho \,\, = (\gamma - 1)T + V, \\ P \,\, = (\gamma - 1)\dfrac{T}{\gamma} - V. \end{array}$$

$$3 M_{\rm pl}^2 H^2 = (\gamma - 1) T + V(\phi)$$

$$\dot{\phi}=-\frac{2M_{\rm pl}^2H'}{\gamma},$$

$$\gamma=\sqrt{1+\frac{\left(2M_{\rm pl}^2H'\right)^2}{T}}$$

$$\begin{aligned}\varepsilon \,\, &= -\frac{\dot{H}}{H^2}=\frac{2M_{\rm pl}^2}{\gamma}\bigg(\frac{H'}{H}\bigg)^2 \\ \tilde{\eta} \,\, &= \frac{\dot{\varepsilon}}{H\varepsilon}=\frac{2M_{\rm pl}^2}{\gamma}\bigg[2\bigg(\frac{H'}{H}\bigg)^2-2\frac{H''}{H}+\frac{H'}{H}\frac{\gamma'}{\gamma}\bigg]\end{aligned}$$

$$\frac{V}{\gamma T}\gg 1$$

$$\gamma^2=\frac{2}{3}\epsilon\frac{V}{T}\gg 1$$

$$\begin{gathered}\text{dilatation :} x^\mu \\ \phi(x) \mapsto \tilde{x}^\mu \equiv (1+c)x^\mu \\ \text{SCTs : } \mapsto \phi(\tilde{x})+c \\ x^\mu \mapsto \tilde{x}^\mu \equiv x^\mu + (b \cdot x)x^\mu - \frac{1}{2}\left(x^2 + \frac{\lambda}{\phi^2}\right)b^\mu \\ \phi(x) \mapsto \phi(\tilde{x})(1-(b \cdot x))\end{gathered}$$

$$S=\int\,\,\mathrm{d}^4x\phi^4f((\partial\phi)^2/\phi^4)+\cdots$$

$$f(z)=\alpha[\sqrt{1+\lambda z}-\beta]$$



$$V_{\mathcal{C}}(\phi)=D_0\left(1-\frac{\pi}{4\mathrm{Vol}(X_5)}\frac{D_0}{\phi^4}\right).$$

$$\frac{V(\phi)}{\gamma T(\phi)}=\frac{\lambda}{2\gamma}\frac{m^2}{\phi^2}$$

$$M_{\rm KK} \sim \frac{1}{L} e^{A_{\rm IR}} \sim \frac{r_{\rm IR}}{L^2},$$

$$r_{\rm IR}\sim mL^2$$

$$\phi^2 \gtrsim \phi_{\rm IR}^2 \sim T_3 (m L^2)^2 = \lambda m^2.$$

$$\frac{V(\phi)}{\gamma T(\phi)}\lesssim \frac{1}{2\gamma}\ll 1$$

$$\frac{V(\phi)}{\gamma T(\phi)}=\frac{1}{\gamma}\frac{e^{4A}-\alpha}{e^{4A}}$$

$$|\alpha| \gg \gamma e^{4A}$$

$$N_e^{\text{EFT}} \approx \frac{N^{1/8}}{\sqrt{\beta}}$$

$$S=\int~\mathrm{d}^4x\sqrt{-g}\left[\frac{M_{\mathrm{pl}}^2}{2}R+P(X,\phi)\right]$$

$$P(X,\phi)=-T(\phi)\left(\sqrt{1-\frac{2X}{T(\phi)}}-1\right)-V(\phi).$$

$$c_s^2=\frac{P_{,X}}{P_{,X}+2XP_{,XX}}=\frac{1}{\gamma^2(\phi)}$$

$$c_s^2=1-\frac{2X}{T(\phi)}.$$

$$r=16 c_s \varepsilon$$

$$\frac{\Delta \phi}{M_{\mathrm{pl}}} = \int_0^{N_\star} \sqrt{\frac{r(N)}{8} \frac{1}{c_s P_{,X}}}~\mathrm{d}N$$

$$c_s P_{,X}=1.$$

$$A=c_s^2\left(-1-\frac{2}{3}\frac{XP_{,XXX}}{P_{,XX}}\right).$$

$$A=-1.$$



$$f_{\rm NL}^{\rm equil}=(-0.27+0.08A)\frac{1-c_s^2}{c_s^2},$$

$$f_{\rm NL}^{\rm ortho}=(+0.02+0.02A)\frac{1-c_s^2}{c_s^2}.$$

$$f_{\rm NL}^{\rm equil}=-\frac{35}{108}\biggl(\frac{1}{c_s^2}-1\biggr)\approx-\frac{35}{108}\gamma^2,$$

$$\gamma \lesssim 24~(95\%\,\mathrm{C.L.})$$

$$2\Bigl(\frac{M_{\rm pl}}{\phi}\Bigr)^2=\varepsilon\gamma=\frac{1}{16}r\gamma^2,$$

$$2\Bigl(\frac{M_{\rm pl}}{\phi}\Bigr)^2=\frac{27}{140}r\left|f_{\rm NL}^{\rm equil}\right|$$

$$N<\frac{27}{70}r\left|f_{\rm NL}^{\rm equil}\right|\lesssim 9~(95\%\,\mathrm{C.L.})$$

$$\Delta_{\mathcal{R}}^2(k_\star) = \left(\frac{32}{3\pi}\right)^2 \frac{3}{r^4\left(f_{\rm NL}^{\rm equil}\right)^2} \frac{{\rm Vol}(X_5)}{N} \gtrsim 15 \frac{{\rm Vol}(X_5)}{N},$$

$$N\gtrsim 7\times 10^9 {\rm Vol}(X_5)$$

$$f_{\rm NL}^{\rm equil}\approx -\frac{35}{108}\gamma^2\mathrm{cos}^2~\theta$$

$$\mathcal{L}(\phi) = -\frac{1}{2} (\partial \phi)^2 - \Lambda^4 \left[1 - \cos \left(\frac{\phi}{f} \right) \right] + \cdots$$

$$\mathcal{L}\supset \sum_{i=1}^2\frac{\phi_i}{f_i}\Big(\frac{c_{ia}}{32\pi^2}\text{Tr}\big[F^{(a)}\wedge F^{(a)}\big]+\frac{c_{ib}}{32\pi^2}\text{Tr}\big[F^{(b)}\wedge F^{(b)}\big]\Big)$$

$$V=\Lambda_a^4\left[1-\cos\left(c_{1a}\frac{\phi_1}{f_1}+c_{2a}\frac{\phi_2}{f_2}\right)\right]+\Lambda_b^4\left[1-\cos\left(c_{1b}\frac{\phi_1}{f_1}+c_{2b}\frac{\phi_2}{f_2}\right)\right].$$

$$\frac{c_{1a}}{c_{2a}}=\frac{c_{1b}}{c_{2b}}$$

$$f_\xi=\frac{(c_{2a}^2f_1^2+f_2^2)^{1/2}}{|c_{2b}-c_{2a}|},\text{ with }\xi=\frac{\phi_2f_2-c_{2a}\phi_1f_1}{c_{2a}^2f_1^2+f_2^2}$$

$$\mathcal{L}=\sum_{i=1}^N\Bigl[-\frac{1}{2}(\partial\phi_i)^2-V_i(\phi_i)\Bigr],$$

$$\ddot{\phi}_i+3H\dot{\phi}_i=-\partial_i V_i, \text{ where } \, 3M_{\rm pl}^2 H^2\approx \sum_{i=1}^N\,V_i.$$



$$V(\Phi)=\frac{1}{2}m^2\Phi^2$$

$$\delta M_{\mathrm{pl}}^2 \sim \frac{N}{16\pi^2}\Lambda_{\mathrm{UV}}^2$$

$$W=W_0+\sum_{i=1}^N\mathcal{A}_ie^{-a_iT_i}$$

$$K=-2\mathrm{ln}\,[\mathcal{V}(T_i,\bar{T}_i)],$$

$$\mathcal{L} = -\frac{1}{2}M_{\mathrm{pl}}^2 K_{ij} \partial_\mu \vartheta^i \partial^\mu \vartheta^j - \frac{1}{2}M_{ij} \vartheta^i \vartheta^j + \cdots$$

$$\mathcal{L}=\sum_{i=1}^N\left[-\frac{1}{2}(\partial\phi_i)^2-\frac{1}{2}m_i^2\phi_i^2\right]$$

$$\rho(m^2) = \frac{1}{2\pi N \sigma^2 m^2} \sqrt{(\eta_+ - m^2)(m^2 - \eta_-)}.$$

$$\mathcal{L}=\frac{M_{\mathrm{pl}}^2}{2}\bigg(1+\frac{\zeta(3)\chi(X_6)}{(2\pi)^3}\frac{(\alpha')^3}{\mathcal{V}}+\cdots\bigg)R_4,$$

$$\chi(X_6)=2h^{1,1}-2h^{2,1},$$

$$S_{\mathrm{D5}}=\frac{1}{(2\pi)^5 g_s (\alpha')^3}\int_{\mathcal{M}_4\times\Sigma_2} \mathrm{d}^6\sigma \sqrt{-\mathrm{det}(G_{ab}+B_{ab})}$$

$$V(b)=\frac{\varrho}{(2\pi)^6 g_s (\alpha')^2}\sqrt{(2\pi)^2\ell^4+b^2}$$

$$V(c)=\frac{\varrho}{(2\pi)^6 g_s^2 (\alpha')^2}\sqrt{(2\pi)^2\ell^4+g_s^2 c^2}$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-\mu^3\phi,~\text{with}~\mu^3\equiv\frac{1}{f}\frac{\varrho}{(2\pi)^6 g_s (\alpha')^2},$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-\mu^{4-p}\phi^p$$

$$S \supset \int \; \mathrm{d}^{10}X |C_2 \wedge H_3|^2$$

$$S_{\mathrm{CS}}\supset i\mu_5\int_{\mathcal{M}_4\times\Sigma_2} C_4\wedge \mathcal{F}_2$$

$$S_{\mathrm{CS}}^{(\mathrm{D3})}=i\mu_3\int_{\mathcal{M}_4} C_4,$$



$$\int_{\Sigma_2} \mathcal{F}_2 \neq 0$$

$$W=W_0+\mathcal{A}e^{-2\pi T}\\ K=-3\mathrm{ln}\,(T+\bar{T}+\gamma b^2)$$

$$S \rhd \int_{\Sigma_4} C_2 \wedge \mathcal{F}_2$$

$$W_{\lambda\lambda}=\mathcal{A}e^{-f/N_c},$$

$$f=f_0+f_1+f_{np}$$

$$W_{\lambda\lambda}=\mathcal{A}e^{-(f_0+f_1)/N_c}[1+\mathcal{O}(e^{-S})g(c)],$$

$$V(c)=\mu^3fc+\hat V[T_i(c)]$$

$$\text{Re}(T_i) = \int_{\Sigma_4^i} \text{d}^4y \sqrt{g} e^{-4A(y;c)}$$

$$\mathcal{L}=-\frac{1}{2}(\partial\phi)^2-V_0(\phi)-\Lambda^4\text{cos}\left(\frac{\phi}{f}\right)$$

$$b_\star\equiv\frac{\Lambda^4}{V_0'(\phi_\star)f}<1$$

$$\frac{f^2}{M_{\mathrm{pl}}^2}>\frac{\sqrt{g_{\mathrm{s}}}}{(2\pi)^3\mathcal{V}}$$

$$f\gg(2\Delta_{\mathcal{R}})^{-1/2}\frac{H}{2\pi}$$

$$\frac{f}{M_{\mathrm{pl}}} \gg \sqrt{\frac{r}{16}}\Delta_{\mathcal{R}}^{1/2} \approx 4.5\times 10^{-4}\left(\frac{r}{0.07}\right)^{1/2}$$

$$\Delta_{\mathcal{R}}^2(k)=\Delta_{\mathcal{R}}^2(k_\star)\Big(\frac{k}{k_\star}\Big)^{n_s-1}\Big[1+\mathcal{A}\text{cos}\left(\frac{\phi_k}{f}\right)\Big]$$

$$\mathcal{A}\equiv 3b_\star\left(\frac{2\pi f}{\sqrt{2\epsilon_\star}}\right)^{1/2}$$

$$\mathcal{A}_{p=1}=\frac{6b_\star}{\sqrt{1+(3f\phi_\star)^2}}\sqrt{\frac{\pi}{2}\coth\left(\frac{\pi}{2f\phi_\star}\right)f\phi_\star}$$

$$B_{\mathcal{R}}(k_1,k_2,k_3)=f_{\mathrm{NL}}^{\mathrm{res}}\times\frac{(2\pi\Delta_{\mathcal{R}})^4}{k_1^2k_2^2k_3^2}\Biggl[\sin\left(\frac{\sqrt{2\epsilon_\star}}{f}\ln\frac{K}{k_\star}\right)+\frac{f}{\sqrt{2\epsilon_\star}}\sum_{i\neq j}\cos\left(\frac{\sqrt{2\epsilon_\star}}{f}\ln\frac{K}{k_\star}\right)+\cdots\Biggr]$$



$$f_{\rm NL}^{\rm res}\equiv \frac{3\sqrt{2\pi}}{8} b_\star \left(\frac{\sqrt{2\epsilon_\star}}{f}\right)^{3/2}$$

$$f_{\rm NL}^{\rm res}\ll \frac{3\sqrt{\pi}}{2}(2\Delta_{\mathcal R})^{-3/4}\approx 3\times 10^3$$

$$\mathcal{L} \supset -\frac{\alpha}{4} \frac{\phi}{f} F \tilde{F}$$

$$\left(\frac{\partial^2}{\partial \tau^2}+k^2\mp 2aHk\xi\right)A_{\pm}(\tau,k)=0,\text{ where }\,\xi\equiv\frac{\alpha\dot{\phi}}{2fH}.$$

$$A_+(\tau,k)\simeq \frac{1}{\sqrt{2k}}\Big(\frac{k}{2\xi aH}\Big)^{1/4}e^{\pi\xi-2\sqrt{2\xi k/(aH)}}$$

$$f_{\rm NL}^{\rm equil}\simeq \frac{\Delta_{\mathcal R,0}^6}{\Delta_{\mathcal R}^4}f_3(\xi)e^{6\pi\xi}$$

$$\begin{array}{l} f_3(\xi)=2.8\times 10^{-7}\xi^{-9} \text{ for } \xi\gg 1 \\ f_3(\xi)\approx 7.4\times 10^{-8}\xi^{-8.1} \text{ for } 2<\xi<3 \end{array}$$

$$f>\frac{\alpha}{10\pi}\frac{H}{\Delta_{\mathcal R}}.$$

$$\frac{f}{M_{\mathrm{pl}}} > \frac{\alpha}{10} \sqrt{\frac{r}{2}} \approx 2 \times 10^{-2} \Big(\frac{\alpha}{1} \Big) \Big(\frac{r}{0.07} \Big)^{1/2}.$$

$$\Delta_{\mathcal R}^2(k)=\Delta_{\mathcal R,0}^2(k)\big[1+\Delta_{\mathcal R,0}^2(k)f_2(\xi)e^{4\pi\xi}\big]$$

$$\begin{array}{l} f_2(\xi)=7.5\times 10^{-5}\xi^{-6} \text{ for } \xi\gg 1 \\ f_2(\xi)\approx 3.0\times 10^{-5}\xi^{-5.4} \text{ for } 2<\xi<3 \end{array}$$

$$W=W_0+\mathcal{A} e^{-aT}+\mathcal{B} e^{-bT},$$

$$K=-3\mathrm{ln}\,(T+\bar{T}),$$

$$\delta V=\frac{\varrho}{(T+\bar{T})^2}$$

$$K=K_0+\delta K_{(\alpha')}=-2\mathrm{ln}\,(\mathcal{V})-\frac{\hat{\xi}}{\mathcal{V}}$$

$$K=K_0+\delta K_{(\alpha')}+\delta K_{(g_{\mathrm{s}})}$$

$$\delta K^{\rm KK}_{(g_{\mathrm{s}})}\sim g_{\mathrm{s}}\sum_{i=1}^{h^{1,1}}\frac{\mathcal{C}_i^{\rm KK}(\zeta,\bar{\zeta})M_{\rm KK}^{-2}}{\mathcal{V}}\sim g_{\mathrm{s}}\sum_{i=1}^{h^{1,1}}\frac{\mathcal{C}_i^{\rm KK}(\zeta,\bar{\zeta})\big(a_{ij}t^j\big)}{\mathcal{V}},$$

$$\delta K^{\rm W}_{(g_{\mathrm{s}})}\sim \sum_i\frac{\mathcal{C}_i^{\rm W}(\zeta,\bar{\zeta})M_{\rm W}^{-2}}{\mathcal{V}}\sim \sum_i\frac{\mathcal{C}_i^{\rm W}(\zeta,\bar{\zeta})}{(b_{ij}t^j)\mathcal{V}},$$



$$\delta K^{\text{KK}}_{(g_s)} \sim \sum_i \frac{t_i}{\mathcal{V}} > \delta K_{(\alpha')} \sim \frac{1}{\mathcal{V}}.$$

$$K = -2\mathrm{ln}\left(\mathcal{V}\right) - \frac{\hat{\xi}}{\mathcal{V}} + \frac{\sqrt{\tau}}{\mathcal{V}}$$

$$V=\frac{W_0^2}{\mathcal{V}^3}\bigg[0+\hat{\xi}+0\cdot\sqrt{\tau}+\frac{1}{\sqrt{\tau}}+\frac{1}{\tau^{3/2}}\bigg].$$

$$\delta V_{\text{CW}} \simeq 0 \cdot \Lambda^4 + \Lambda^2 \text{STr}(M^2) + \text{STr}\left[M^4 \text{ln}\left(\frac{M^2}{\Lambda^2}\right)\right]$$

$$\Lambda=M_{\text{KK}}\simeq\frac{M_{\text{pl}}}{\mathcal{V}^{2/3}}\,,\,\,\text{STr}(M^2)\simeq\frac{M_{\text{pl}}^2}{\mathcal{V}^2},$$

$$\delta V_{\text{CW}} \simeq 0 \cdot \frac{1}{\mathcal{V}^{8/3}} + \frac{1}{\mathcal{V}^{10/3}} + \frac{1}{\mathcal{V}^4}$$

$$\delta V_{\text{CW}} \simeq \frac{1}{\mathcal{V}^3}\bigg[0\cdot\sqrt{\tau}+\frac{1}{\sqrt{\tau}}+\frac{1}{\tau^{3/2}}\bigg]$$

$$\mathcal{V}=\alpha\left(\tau_b^{3/2}-\lambda_\phi\tau_\phi^{3/2}-\lambda_s\tau_s^{3/2}\right)$$

$$W=W_0+\mathcal{A}_\phi e^{-a_\phi T_\phi}+\mathcal{A}_s e^{-a_s T_s}$$

$$V=W_0^2\left(\mathfrak{a}\frac{\sqrt{\tau_\phi}e^{-2a_\phi\tau_\phi}}{\mathcal{V}}-\mathfrak{b}\frac{\tau_\phi e^{-a_\phi\tau_\phi}}{\mathcal{V}^2}+\mathfrak{c}\frac{\hat{\xi}}{\mathcal{V}^3}\right),$$

$$V(\phi) \simeq V_0 \big(1 - c_1 \mathcal{V}^{5/3} \phi^{4/3} \mathrm{exp}\left[-c_2 \mathcal{V}^{2/3} \phi^{4/3}\right]\big),$$

$$\phi \equiv \sqrt{4\lambda_\phi/(3\mathcal{V})}\tau_\phi^{3/4},$$

$$a_\phi \langle \tau_\phi \rangle \approx \mathcal{O}(1) \times \ln{(\mathcal{V})} \; \Leftrightarrow \; \langle \phi \rangle \approx \mathcal{O}(1) \times \frac{\ln{(\mathcal{V})}^{3/4}}{\mathcal{V}^{1/2}}$$

$$V(\hat{\phi}) \simeq V_0 \Big(1 - \kappa_1 e^{-\kappa_2 \hat{\phi}}\Big)$$

$$\begin{aligned}\kappa_1&\equiv c_1\mathcal{V}^{5/3}\langle\phi\rangle^{4/3}\approx\mathcal{O}(\mathcal{V}\ln{(\mathcal{V})})\\\kappa_2&\equiv\frac{4}{3}c_2\mathcal{V}^{2/3}\langle\phi\rangle^{1/3}=\mathcal{O}\big(\mathcal{V}^{1/2}\ln{(\mathcal{V})}^{1/4}\big)\end{aligned}$$

$$\mathcal{V}^{2/3}\phi^{4/3}\gg 1 \text{ and } \phi\ll 1$$

$$\delta V_{(g_s)} \sim \frac{1}{\sqrt{\tau_\phi}\mathcal{V}^3} \sim \frac{1}{\phi^{2/3}\mathcal{V}^{10/3}}$$

$$\delta\eta\sim\frac{\delta V''_{(g_s)}}{V_0}\sim\frac{1}{\phi^{8/3}\mathcal{V}^{1/3}}\sim\frac{\mathcal{V}}{\tau_\phi^2}$$



$$\delta\eta \approx a_\phi^2 \frac{\mathcal{V}}{\ln{(\mathcal{V})^2}} \gg 1$$

$$\mathcal{V}=\frac{1}{2}\sqrt{\tau_1}\tau_2,$$

$$\mathcal{V} = \alpha \left(\sqrt{\tau_1} \tau_2 - \lambda_s \tau_s^{3/2} \right)$$

$$W=W_0+\mathcal{A}_se^{-a_sT_s}$$

$$V=a_s^2\mathcal{A}_s^2\frac{\sqrt{\tau_s}}{\mathcal{V}}e^{-2a_s\tau_s}-a_s\mathcal{A}_sW_0\frac{\tau_s}{\mathcal{V}}e^{-a_s\tau_s}+\hat{\xi}W_0^2\frac{1}{\mathcal{V}^3}.$$

$$\delta V_{(g_s)}=\frac{W_0^2}{\mathcal{V}^2}\biggl(\mathfrak{a}\frac{g_s^2}{\tau_1^2}-\mathfrak{b}\frac{1}{\sqrt{\tau_1}\mathcal{V}}+\mathfrak{c}\frac{g_s^2\tau_1}{\mathcal{V}^2}\biggr)$$

$$V(\phi)=V_0\left(1-\frac{4}{3}e^{-\phi/\sqrt{3}}+\frac{1}{3}e^{-4\phi/\sqrt{3}}+\frac{\mathfrak{C}}{3}e^{2\phi/\sqrt{3}}\right),$$

$$\phi\equiv\frac{\sqrt{3}}{2}\ln~\tau_1$$

$$V(\phi)\simeq V_0\left(1-\frac{4}{3}e^{-\phi/\sqrt{3}}\right)$$

$$W=\mathcal{A}_a\text{exp}\;(-S_a+\mathcal{A}_be^{-S_b}),$$

$$\mathcal{V} = \alpha \left(\sqrt{\tau_1} \tau_2 - \lambda_s \tau_s^{3/2} \right)$$

$$\mathcal{V}=\tau_b^{3/2}-\tau_s^{3/2}-(\tau_s+\tau_w)^{3/2}$$

$$W=W_0+\mathcal{A}\text{exp}\;[-aT_s]-\mathcal{B}\text{exp}\;[-bT_s],$$

$$W=W_0+\mathcal{A}\text{exp}\;[-a(T_s+c_1e^{-2\pi T_1})]-\mathcal{B}\text{exp}\;[-b(T_s+c_2e^{-2\pi T_1})],$$

$$V=\frac{F_{\mathrm{poly}}}{\mathcal{V}^{3+p}}\bigl(1-(1+2\pi\hat t_1)e^{-2\pi\hat t_1}\bigr)$$

$$V(\hat{\phi})\simeq V_0\Big(1-\kappa_2\hat{\phi}e^{-\kappa_2\hat{\phi}}\Big), \hat{\phi}\approx\frac{\sqrt{3}}{2}\frac{\hat{t}_1}{\langle\tau_1\rangle},$$

$$V(\phi) \approx V_0\big(1-\kappa_1 e^{-\kappa_2 \phi}\big)$$

$$\eta\simeq -\kappa_1\kappa_2^2e^{-\kappa_2\phi}\;\;\text{and}\;\;\epsilon\simeq \frac{1}{2}\frac{\eta^2}{\kappa_2^2}$$

$$n_s\simeq 1+2\eta$$



$$r\simeq \frac{2}{\kappa_2^2}\,(n_s-1)^2\stackrel{n_s=0.96}{\rightarrow} \frac{3\times 10^{-3}}{\kappa_2^2}$$

$$r\simeq 6(n_s-1)^2\stackrel{n_s=0.96}{\rightarrow}~0.01.$$

$$r\sim 10^{-5}$$

$$S=\int\; {\rm d}^4x\sqrt{-g}\left[\frac{M_{\rm pl}^2}{2}R-\frac{1}{2}(\partial\phi)^2-V(\phi)+\mathcal{O}(\phi,\psi)\right].$$

$$\mathcal{L}_{\text{int}} = -\frac{1}{2} g^2 (\phi - \phi_0)^2 \psi^2$$

$$\phi(t) \approx \phi_0 + \dot{\phi}_0(t-t_0)$$

$$m_\psi^2(t)\equiv g^2(\phi-\phi_0)^2\approx k_\star^4(t-t_0)^2$$

$$\ddot{\psi}_k+3H\dot{\psi}_k+\underbrace{\left(\frac{k^2}{a^2}+k_\star^4(t-t_0)^2\right)}_{\equiv\omega_k^2(t)}\psi_k=0$$

$$|\dot{\omega}_k|>\omega_k^2$$

$$n_k=e^{-\pi k^2/k_\star^2}$$

$$n_\psi(t_0)=\int\; \frac{{\rm d}^3k}{(2\pi)^3}n_k\approx \frac{k_\star^3}{(2\pi)^3}$$

$$n_\psi(t)=\frac{k_\star^3}{(2\pi)^3}\frac{a^3(t_0)}{a^3(t)}\Theta(t-t_0)$$

$$\ddot{\phi}+3H\dot{\phi}+V'=-g^2(\phi-\phi_0)\langle\psi^2\rangle,$$

$$\langle\psi^2\rangle\approx\frac{n_\psi(t)}{g|\phi-\phi_0|}$$

$$\mathcal{L}_{\text{int}}=-\frac{1}{2}g^2\sum_{i=1}^N\;(\phi-\phi_i)^2\psi_i^2$$

$$n_{\psi_i}(t_i)\simeq\frac{\left(g\dot{\phi}(t_i)\right)^{3/2}}{(2\pi)^3}\!\equiv\!\frac{\dot{\phi}^{3/2}(t_i)}{(2\pi)^3}$$

$$\ddot{\phi}+3H\dot{\phi}+V'+\sum_i\frac{g\dot{\phi}^{3/2}(t_i)}{(2\pi)^3}\frac{a^3(t_i)}{a^3(t)}=0.$$



$$\sum_i \frac{g\dot{\phi}^{3/2}(t_i)}{(2\pi)^3}\frac{a^3(t_i)}{a^3(t)}\approx \int \frac{{\rm d} t'}{\Delta}\frac{{\rm d} t'}{\Delta}\frac{\dot{\phi}^{5/2}(t')}{(2\pi)^3}\frac{a^3(t')}{a^3(t)}\approx \frac{1}{3H\Delta}\frac{\dot{\phi}^{5/2}(t)}{(2\pi)^3},$$

$$\ddot{\phi}+3H\dot{\phi}+V'+\frac{1}{24\pi^3}\frac{g^{5/2}}{H\Delta}\dot{\phi}^{5/2}=0$$

$$\dot{\phi}=g\dot{\phi}\approx -(24\pi^3 H\Delta V')^{2/5}.$$

$$3M_{\mathrm{pl}}^2H^2=\rho_\phi+\rho_\psi\approx V(\phi)$$

$$2 M_{\mathrm{pl}}^2 \dot{H} \approx - \rho_\psi,$$

$$\varepsilon = - \frac{\dot{H}}{H^2} \approx \frac{3}{2} \frac{\rho_\psi}{V},$$

$$\rho_\psi(t)=\sum_i~g|\phi-\phi_i|n_{\psi_i}(t)\approx\int\frac{{\rm d} t'}{\Delta}\frac{{\rm d} t'}{\Delta}|\phi(t)-\phi(t')|\frac{\dot{\phi}^{5/2}(t')}{(2\pi)^3}\frac{a^3(t')}{a^3(t)}.$$

$$\rho_\psi(t) \approx \frac{1}{(3H)^2}\frac{1}{g\Delta}\frac{\dot{\phi}^{7/2}(t)}{(2\pi)^3}$$

$$\varepsilon \sim \frac{\epsilon^{7/10}}{g} \biggl(\frac{H}{M_{\mathrm{pl}}} \frac{\Delta^2}{M_{\mathrm{pl}}^2}\biggr)^{1/5}$$

$$\begin{array}{l} t_x\!:\!(x,u_1,u_2)\mapsto(x+1,u_1,u_2)\\ t_{u_1}\!:\!(x,u_1,u_2)\mapsto(x-Mu_2,u_1+1,u_2)\\ t_{u_2}\!:\!(x,u_1,u_2)\mapsto(x,u_1,u_2+1),\end{array}$$

$$\frac{{\rm d}s^2}{\alpha'}=L_{u_1}^2\;{\rm d} u_1^2+\underbrace{L_{u_2}^2\;{\rm d} u_2^2+L_x^2({\rm d} x+Mu_1\;{\rm d} u_2)^2}_{\tilde{T}^2}$$

$$\frac{{\rm d}s_{T^2}^2(u_1)}{\alpha'}=L_{u_2}^2\;{\rm d} u_2^2+L_x^2({\rm d} x+Mu_1\;{\rm d} u_2)^2$$

$$\frac{{\rm d}s_{T^2,\perp}^2}{\alpha'}=L_x^2\;{\rm d} y_1^2+L_{u_2}^2\;{\rm d} y_2^2$$

$$S_{\text{D4}}=-\frac{1}{(2\pi)^4 g_s (\alpha')^2}\int\;\; {\rm d}^4x \sqrt{-g}\sqrt{L_{u_2}^2+L_x^2 M^2 u_1^2}\Big(1-\frac{1}{2}\alpha' L_{u_1}^2 \dot{u}_1^2\Big).$$

$$S_{\text{D4}}=\int\; {\rm d}^4x \sqrt{-g}\Big(\frac{1}{2}\dot{\phi}^2-\mu^{10/3}\phi^{2/3}\Big)$$

$$\frac{\phi^2}{M_{\mathrm{pl}}^2}=\frac{2}{9}(2\pi)^3g_s\frac{M}{L^3}\frac{L_{u_1}}{L_{u_2}}u_1^3\\\frac{\mu}{M_{\mathrm{pl}}}=\frac{M_s}{M_{\mathrm{pl}}}\bigg(\frac{9}{4}\frac{M^2}{(2\pi)^8g_s^2}\bigg(\frac{L_x}{L}\bigg)^3\frac{L_{u_2}}{L_{u_1}}\bigg)^{1/10}$$



$$\Delta u_1^3 \gg \frac{L^3}{M}$$

$$\int_{\Sigma_2} {\mathcal F}_2 = n \in {\mathbb Z}$$

$$T_{{\rm D}3/\Sigma_2}=T_3\sqrt{\ell^2+(cg_s+n)^2},$$

$$\mathcal{M}_{10}=dS_4\times X_6,$$

$$\int_{dS_4\times\Sigma_{p-2}}F_{p+2}=Q_0\gg 1.$$

$${\mathrm d}s_{\rm E}^2=H^{-2}(\,{\mathrm d}\xi^2+\sin^2\,\xi\,{\mathrm d}\Omega_3^2)+{\mathrm d}z^2,$$

$$S=\int\;{\mathrm dz}\int\;\;{\mathrm d}\mathcal{H}_3\;{\mathrm dt}\frac{\sinh^3\left(Ht\right)}{H^3}\biggl(-2\sigma\delta(z-z_b)\sqrt{1-(\partial z_b)^2}-\frac{F_5^2}{2\cdot5!}\biggr)$$

$$\frac{F_5^2}{5!}=\mu^5 Q^2=\mu^5\Bigg(Q_0+\sum_{j=-\infty}^\infty\left[\Theta(z-z_b+j\ell)-\Theta(z+z_b+j\ell)\right]\Bigg)^2$$

$$S=\int\;{\mathrm dt}\;{\mathrm d}^3x a^3(t)\Bigl(-2\sigma\sqrt{1-(\partial z_b)^2}-V(z_b)\Bigr)$$

$$V(z_b) \sim \mu^5 \left(Q_0 - \frac{z_b}{\ell} \right)^2.$$

$$\mathcal{L} \supset -\sum_{i=1}^N \alpha_i \frac{\phi}{f} F_i \tilde{F}_i$$

$$S_{\text{CS}}=\frac{i}{2\pi}\!\int_{\mathcal{M}_4}\!C_0\text{Tr}[\mathcal{F}_2\wedge\mathcal{F}_2]$$

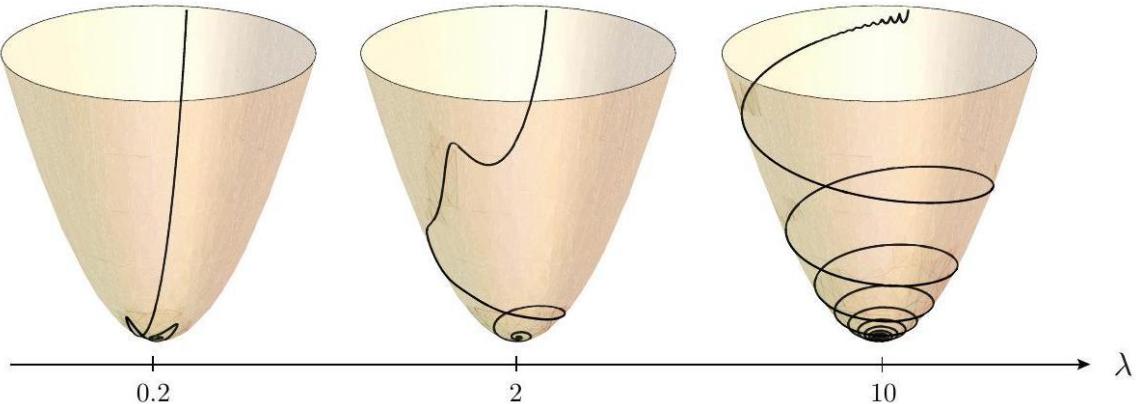
$$S_{\text{eff}}=\int\;{\mathrm d}^4x\bigg\{a^3\left[\gamma_C\dot{C}_0^2-V(C_0)+\gamma_A\frac{\text{Tr}(\dot{A}^2)}{a^2}+\gamma_A\frac{\text{Tr}([A,A]^2)}{a^4}\right]+\kappa C_0\text{Tr}(\dot{A}[A,A])\bigg\}$$

$$C_0(t)\equiv\frac{\phi(t)}{\sqrt{\gamma_c}}$$

$$A_0\equiv 0, A_i(t)\equiv \frac{a(t)\psi(t)}{\sqrt{\gamma_A\nu}}J_i$$

$$S_{\text{eff}}=\int\;{\mathrm d}^4x a^3\left[\frac{1}{2}\phi^2-V(\phi)+\frac{3}{2}(\dot{\psi}+H\psi)^2-\frac{3}{2}g^2\psi^4-\frac{3g\lambda}{f}\phi\psi^2(\dot{\psi}+H\psi)\right]$$





$$\ddot{\phi}+3H\dot{\phi}+V_{,\phi}=-3\frac{g\lambda}{f}\psi^2(\dot{\psi}+H\psi)$$

$$(N_e)_{\rm max} \approx \frac{3}{5} \lambda$$

$$S_{\text{CS}} = \frac{i}{(2\pi)^3} \int_{\mathcal{M}_4 \times \Sigma_4} \frac{\mathcal{C}_0}{24} \left(\text{Tr}[\mathcal{F}_2 \wedge \mathcal{F}_2 \wedge \mathcal{F}_2 \wedge \mathcal{F}_2] + \frac{1}{2} \text{Tr}[\mathcal{F}_2 \wedge \mathcal{F}_2] \hat{\chi}(R) \right)$$

$$\lambda=\left[K+\frac{\chi(\Sigma_4)}{24}\right]\times g_s\times\frac{\ell_s^4}{\mathcal{V}_4}$$

$$\ddot{\delta\phi}+\left(M^2+\frac{k^2}{a^2}\right)\delta\phi+\int^t\mathrm{d}\;\mathrm{d}t'M^2\left(\frac{5}{2}\dot{\delta\phi}(t')-3H\delta\phi(t')\right)\frac{a^3(t')}{a^3(t)}=-g\Delta n_\psi(k,t)$$

$$M^2\equiv\frac{g^{5/2}}{(2\pi)^3}\frac{\dot{\phi}^{3/2}}{\Delta}$$

$$\Delta_{\mathcal{R}}^2 \approx g^{8/3} \frac{H}{M} \biggl(\frac{M}{\Delta}\biggr)^{2/3}$$

$$\Delta_{\mathcal{R}}^2 \approx g^{9/4} \biggl(\frac{H}{\Delta}\biggr)^{1/2} \biggl(\frac{H^2}{\dot{\phi}}\biggr)^{1/4},$$

$$n_s-1=\frac{\dot{H}}{H^2}-\frac{1}{4}\frac{\ddot{\phi}}{H\dot{\phi}}.$$

$$r=g^{-8/2}\frac{HM}{M_{\rm pl}^2}\biggl(\frac{\Delta}{M}\biggr)^{2/3}.$$

$$f_{\rm NL}^{\rm equil}\simeq\frac{M^2}{H^2}$$



Interacciones gravitónicas y transferencia de gravedad al centro de masa de las partículas portadoras, volviéndolas supermasivas. Formalismo de Batalin-Vilkovisky por permeabilidad de un campo gravitónico para inmersión gravitacional de un espacio – tiempo cuántico específico.

Propagadores y campos fantasma.

$$Q = \oint_0 j_B(z) dz + \oint_0 \bar{j}_B(\bar{z}) d\bar{z}$$

$$j_B = c T_m + b c \partial c, \bar{j}_B = \bar{c} \bar{T}_m + \bar{b} \bar{c} \partial \bar{c}$$

$$\oint_0 \frac{dz}{z} = 1, \oint_0 \frac{d\bar{z}}{\bar{z}} = 1$$

$$b(z)c(w) \simeq \frac{1}{z-w}, \bar{b}(\bar{z})\bar{c}(\bar{w}) \simeq \frac{1}{\bar{z}-\bar{w}}$$

$$T_m(z)T_m(w) \simeq \frac{26}{2}\frac{1}{(z-w)^4} + \frac{2}{(z-w)^2}T_m(w) + \frac{1}{z-w}\partial_w T_m(w)$$

$$\bar{T}_m(\bar{z})\bar{T}_m(\bar{w}) \simeq \frac{26}{2}\frac{1}{(\bar{z}-\bar{w})^4} + \frac{2}{(\bar{z}-\bar{w})^2}\bar{T}_m(\bar{w}) + \frac{1}{\bar{z}-\bar{w}}\bar{\partial}_w\bar{T}_m(\bar{w})$$

$$T_{bc}(z) = -\partial bc - 2b\partial c, \bar{T}_{bc}(\bar{z}) = -\bar{\partial}\bar{b}\bar{c} - 2\bar{b}\bar{\partial}\bar{c}$$

$$\partial X^\mu(z)\partial X^\nu(w) = -\frac{\eta^{\mu\nu}}{2(z-w)^2}, \bar{\partial}X^\mu(\bar{z})\bar{\partial}X^\nu(\bar{w}) = -\frac{\eta^{\mu\nu}}{2(\bar{z}-\bar{w})^2}$$

$$T(z) = -\eta_{\mu\nu}\partial X^\mu\partial X^\nu, \bar{T}(\bar{z}) = -\eta_{\mu\nu}\bar{\partial}X^\mu\bar{\partial}X^\nu$$

$$|k\rangle \equiv e^{ik\cdot X}(0)|0\rangle,$$

$$\langle k | c_{-1} \bar{c}_{-1} c_0 \bar{c}_0 c_1 \bar{c}_1 | k' \rangle = -(2\pi)^D \delta^{(D)}(k+k')$$

$$\langle k | c_{-1} c_0 c_1 | k' \rangle' = -(2\pi)^{p+1} \delta^{(p+1)}(k+k')$$

$$\langle k | c_{-1} c_0 c_1 | k' \rangle = -(2\pi)^{p+1} K \delta^{(p+1)}(k+k').$$

$$|s\rangle \in \mathcal{H}_c \text{ iff } b_0^-|s\rangle = 0, L_0^-|s\rangle = 0$$

$$b_0^\pm = b_0 \pm \bar{b}_0, L_0^\pm = L_0 \pm \bar{L}_0, c_0^\pm = (c_0 \pm \bar{c}_0)/2$$

$$\langle A,B \rangle \equiv \langle A | c_0^- | B \rangle \; A,B \in \mathcal{H}_c$$

$$\langle A,B \rangle' \equiv \langle A \mid B \rangle', \langle A,B \rangle \equiv \langle A \mid B \rangle.$$

$$\gamma = \eta e^\phi, \beta = \partial \xi e^{-\phi}, \delta(\gamma) = e^{-\phi}, \delta(\beta) = e^\phi$$

$$[\gamma] = -\frac{1}{2}, [\beta] = \frac{3}{2}, [\eta] = 1, [\xi] = 0, [\phi] = 0, [e^{q\phi}] = -\frac{1}{2}q(q+2)$$



$$\xi(z)\eta(w) \simeq \frac{1}{z-w}, \partial\phi(z)\partial\phi(w) \simeq -\frac{1}{(z-w)^2}, e^{q_1\phi}(z)e^{q_2\phi}(w) \simeq (z-w)^{-q_1q_2}e^{(q_1+q_2)\phi}(w)$$

$$T_{\beta\gamma}=\frac{3}{2}\beta\partial\gamma+\frac{1}{2}\gamma\partial\beta=T_{\xi\eta}+T_\phi,T_{\xi\eta}=-\eta\partial\xi,T_\phi=-\frac{1}{2}\partial\phi\partial\phi-\partial^2\phi.$$

$$T_m(z)T_m(w)\simeq \frac{15}{2}\frac{1}{(z-w)^4}+\frac{2}{(z-w)^2}T_m(w)+\frac{1}{z-w}\partial_wT_m(w)$$

$$T_F(z)T_F(w)\simeq \frac{5}{2}\frac{1}{(z-w)^3}+\frac{1}{2}\frac{1}{z-w}T_m(w)$$

$$T_m(z)T_F(w)\simeq \frac{3}{2}\frac{1}{(z-w)^2}T_F(w)+\frac{1}{z-w}T_F(w)$$

$$\psi^\mu(z)\psi^\nu(w)\simeq -\frac{\eta^{\mu\nu}}{2(z-w)}$$

$$j_B=c\big(T_m+T_{\beta\gamma}\big)+bc\partial c+\gamma T_F-\frac{1}{4}\gamma^2 b$$

$$\mathcal{X}(z)=\{Q,\xi(z)\}=c\partial\xi+e^{\phi}T_F-\frac{1}{4}\partial\eta e^{2\phi}b-\frac{1}{4}\partial\big(\eta e^{2\phi}b\big).$$

$$\begin{array}{l} \mathrm{gh}(c)=\mathrm{gh}(\eta)=\mathrm{gh}(\gamma)=1 \\ \mathrm{gh}(b)=\mathrm{gh}(\xi)=\mathrm{gh}(\beta)=-1 \end{array}$$

$$\mathrm{pic}(\eta) = -1, \mathrm{pic}(\xi) = 1, \mathrm{pic}(e^{q\phi}) = q, \mathrm{pic}(\mathcal{X}) = 1$$

$$\text{GSO odd fields: } \beta, \gamma, \psi^\mu, T_F$$

$$|k;m,n\rangle\equiv e^{ik\cdot X}(0)e^{m\phi}(0)e^{n\bar\phi}(0)|0\rangle, |k;m\rangle\equiv e^{ik\cdot X}(0)e^{m\phi}(0)|0\rangle$$

$$\langle k;0,0|c_{-1}\bar{c}_{-1}c_0\bar{c}_0c_1\bar{c}_1e^{-2\phi}e^{-2\bar{\phi}}|k';0,0\rangle=-(2\pi)^D\delta^{(D)}(k+k'),$$

$$\langle k;0|c_{-1}c_0c_1e^{-2\phi}|k';0\rangle'=(2\pi)^{p+1}\delta^{(p+1)}(k+k'),$$

$$\langle k;0|c_{-1}c_0c_1e^{-2\phi}|k';0\rangle=(2\pi)^D\delta^{(D)}(k+k')$$

$$\mathcal{C}=\left\langle \prod_i f_i\circ V_i(0)\right\rangle_{\Sigma},$$

$$f\circ V(w)=(f'(w))^h\big(\bar{f}'(\bar{w})\big)^{\bar{h}}V(f(w))$$

$$V(w)(dw)^h(d\bar w)^{\bar h}=V(z)(dz)^h(d\bar z)^{\bar h}\,\rightarrow\,V(w)=V(z(w))\Big(\frac{dz}{dw}\Big)^h\Big(\frac{d\bar z}{d\bar w}\Big)^{\bar h}$$



$$\mathcal{C}=\left\langle \prod_i V_i(w_i=0)\right\rangle_{\Sigma}$$

$$\langle \chi_s^c \mid \chi_r \rangle = \delta_{rs} \Leftrightarrow \sum_r \mid \chi_r \rangle \langle \chi_r^c \mid = I$$

$$w_1 w_2 = 1$$

$$w_1 w_2 = -1$$

$$d_{g,n}\equiv \dim_\mathbb{R}\bigl(\mathcal{M}_{g,n}\bigr)=6g-6+2n$$

$$\chi_{g,n}=2-2g-n$$

$$\mathcal{A}_g(V_1,\cdots,V_n)=(g_s)^{-\chi_{g,n}}\int_{\mathcal{F}_{g,n}}\Omega^{(g,n)}_{6g-6+2n}(V_1,\cdots,V_n)$$

$$\sigma_s=F_s(\tau_s,u)$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^i}\right]\equiv\sum_s\left[\oint_{c_s}\frac{\partial F_s}{\partial u^i}d\sigma_sb(\sigma_s)+\oint_{c_{\bar{s}}}\frac{\partial\bar{F}_{\bar{s}}}{\partial u^i}d\bar{\sigma}_{\bar{s}}\bar{b}(\bar{\sigma}_{\bar{s}})\right]$$

$$\Omega_p^{(g,n)}(A_1,\cdots,A_n)\left[\frac{\partial}{\partial u^{j_1}},\cdots,\frac{\partial}{\partial u^{j_p}}\right]\equiv\left(-\frac{1}{2\pi i}\right)^{3g-3+n}\left\langle\mathcal{B}\left[\frac{\partial}{\partial u^{j_1}}\right]\cdots\mathcal{B}\left[\frac{\partial}{\partial u^{j_p}}\right]A_1\cdots A_n\right\rangle_{\Sigma_{g,n}}$$

$$z=F(w,u)$$

$$y(u) = F(0,u)$$

$$\Omega_2(c\bar{c}V)=\left(-\frac{1}{2\pi i}\right)du^1\wedge du^2\left\langle\cdots\mathcal{B}\left[\frac{\partial}{\partial u^1}\right]\mathcal{B}\left[\frac{\partial}{\partial u^2}\right]c\bar{c}V(w=0)\right\rangle$$

$$\begin{aligned}\mathcal{B}\left[\frac{\partial}{\partial u^1}\right]\mathcal{B}\left[\frac{\partial}{\partial u^2}\right]&=\left[\oint b(z)dz\frac{\partial F(w,u)}{\partial u^1}+\oint\bar{b}(\bar{z})d\bar{z}\frac{\partial\overline{F(w,u)}}{\partial u^1}\right]\\&\quad\left[\oint b(z)dz\frac{\partial F(w,u)}{\partial u^2}+\oint\bar{b}(\bar{z})d\bar{z}\frac{\partial\overline{F(w,u)}}{\partial u^2}\right]\end{aligned}$$

$$\Omega_2(c\bar{c}V)=-\left(-\frac{1}{2\pi i}\right)du^1\wedge du^2\left(\frac{\partial y}{\partial u^1}\frac{\partial\bar{y}}{\partial u^2}-\frac{\partial y}{\partial u^2}\frac{\partial\bar{y}}{\partial u^1}\right)\langle\cdots V(z=y)\rangle$$

$$du^1\wedge du^2\left(\frac{\partial y}{\partial u^1}\frac{\partial\bar{y}}{\partial u^2}-\frac{\partial y}{\partial u^2}\frac{\partial\bar{y}}{\partial u^1}\right)=dy\wedge d\bar{y}$$

$$\Omega_2(c\bar{c}V)=\frac{1}{2\pi i}\langle\cdots(dy\wedge d\bar{y}V(z=y))\rangle$$

$$\Omega_p^{(g,n)}(QA_1,A_2,\cdots,A_n)+\cdots+(-1)^{A_1+\cdots+A_{n-1}}\Omega_p^{(g,n)}(A_1,A_2,\cdots,QA_n)=(-1)^p~\mathrm{d}\Omega_{p-1}^{(g,n)}(A_1,\cdots,A_n)$$



$$\int_{\mathcal{F}'_{g,n}} \Omega_{6g-6+2n}^{(g,n)}(V_1, \dots, V_n) - \int_{\mathcal{F}_{g,n}} \Omega_{6g-6+2n}^{(g,n)}(V_1, \dots, V_n) = \int_{\mathcal{R}_{g,n}} d\Omega_{6g-5+2n}^{(g,n)}(V_1, \dots, V_n)$$

$$\Omega_{6g-6+2n}^{(g,n)}(QW_1, V_2, \dots, V_n) = d\Omega_{6g-7+2n}^{(g,n)}(W_1, V_2, \dots, V_n)$$

$$n_{-1,-1}, n_{-1,-1/2}, n_{-1/2,-1}, n_{-1/2,-1/2}$$

$$2g-2 = N_L - n_{-1,-1} - n_{-1,-1/2} - \frac{1}{2}n_{-1/2,-1} - \frac{1}{2}n_{-1/2,-1/2}$$

$$2g-2 = N_R - n_{-1,-1} - n_{-1/2,-1} - \frac{1}{2}n_{-1,-1/2} - \frac{1}{2}n_{-1/2,-1/2}$$

$$p=k+\ell+\bar{\ell}$$

$$\Omega_p^{(g,n)}(A_1, \cdots, A_n) \left[\frac{\partial}{\partial u^{j_1}}, \cdots, \frac{\partial}{\partial u^{j_k}}, \frac{\partial}{\partial y_{\alpha_1}}, \cdots, \frac{\partial}{\partial y_{\alpha_\ell}}, \frac{\partial}{\partial \bar{y}_{\beta_1}}, \cdots, \frac{\partial}{\partial \bar{y}_{\beta_{\bar{\ell}}}} \right]$$

$$\equiv \left(-\frac{1}{2\pi i}\right)^{3g-3+n} \left\langle \mathcal{B}\left[\frac{\partial}{\partial u^{j_1}}\right] \cdots \mathcal{B}\left[\frac{\partial}{\partial u^{j_k}}\right] \left(-\partial\xi(y_{\alpha_1})\right) \cdots \left(-\partial\xi(y_{\alpha_\ell})\right) \left(-\bar{\partial}\bar{\xi}(\bar{y}_{\beta_1})\right) \cdots \left(-\bar{\partial}\bar{\xi}(\bar{y}_{\beta_{\bar{\ell}}})\right) \prod_{\alpha=l+1}^{N_R} \mathcal{X}(y_\alpha) \prod_{\beta=\bar{l}+1}^{N_L} \right.$$

$$-\int_{y_\alpha}^{y'_\alpha} \partial\xi(y) dy = \xi(y_\alpha) - \xi(y'_\alpha)$$

$$\Omega_p^{(g,n)}(A_1, \cdots, A_n) \left[\frac{\partial}{\partial u^{j_1}}, \cdots, \frac{\partial}{\partial u^{j_k}}, \frac{\partial}{\partial y_{\alpha_1}}, \cdots, \frac{\partial}{\partial y_{\alpha_\ell}} \right]$$

$$= \left(-\frac{1}{2\pi i}\right)^{3g-3+n} \left\langle \mathcal{B}\left[\frac{\partial}{\partial u^{j_1}}\right] \cdots \mathcal{B}\left[\frac{\partial}{\partial u^{j_k}}\right] \left(-\partial\xi(y_{\alpha_1})\right) \cdots \left(-\partial\xi(y_{\alpha_\ell})\right) \prod_{\alpha=l+1}^{N_R} \mathcal{X}(y_\alpha) A_1 \cdots A_n \right\rangle_{\Sigma_{g,n}}$$

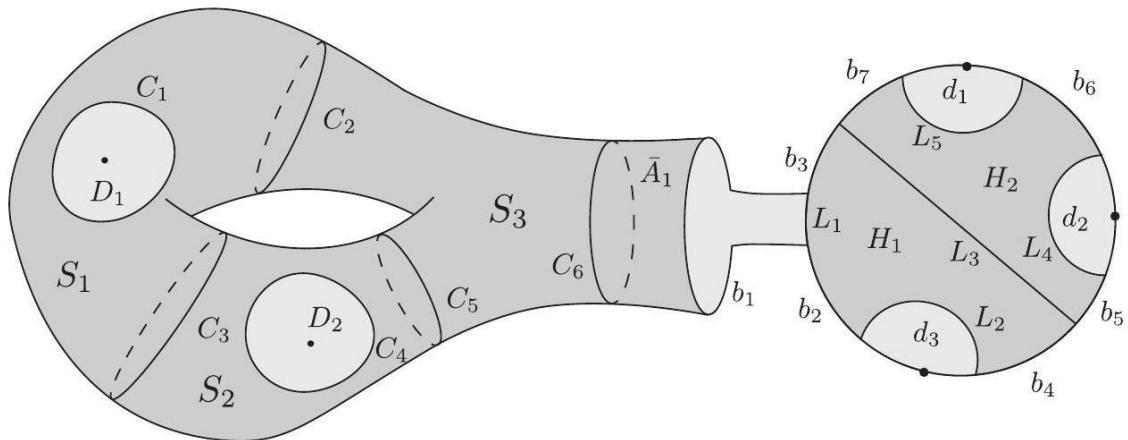


Figura 41. Puente Einstein – Rosen en dimensión \mathbb{R}^4 para entrelazamiento cuántico.

$$d_{g,b,n_c,n_o}\equiv \dim_\mathbb{R}(\mathcal{M}_{g,b,n_c,n_o})=6g-6+3b+2n_c+n_o$$

$$\chi_{g,b,n_c,n_o}=2-2g-n_c-b-\frac{1}{2}n_o$$

$$S+\frac{1}{2}\#H+\frac{1}{2}\#\bar{A}=-\chi_{g,b,n_c,n_o}$$

$$2\#C=3S+n_c+A+\bar{A}$$

$$2\#L=3\#H+n_o+\#\bar{A}$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^i}\right]\equiv\sum_s\left[\oint_{c_s}\frac{\partial F_s}{\partial u^i}d\sigma_sb(\sigma_s)+\oint_{c_{\bar{s}}}\frac{\partial \bar{F}_s}{\partial u^i}d\bar{\sigma}_s\bar{b}(\bar{\sigma}_s)\right]+\sum_m\left[\int_{L_m}\frac{\partial G_m}{\partial u^i}d\sigma_mb(\sigma_m)+\int_{L_m}\frac{\partial \bar{G}_m}{\partial u^i}d\bar{\sigma}_m\bar{b}(\bar{\sigma}_m)\right]$$

$$\sigma_m=G_m(\tau_m,u^i)$$

$$\begin{aligned}&\widehat{\Omega}_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)\Big[\frac{\partial}{\partial u^{j_1}},\cdots,\frac{\partial}{\partial u^{j_p}}\Big]\\&\sim\Big\langle\mathcal{B}\left[\frac{\partial}{\partial u^{j_1}}\right]\cdots\mathcal{B}\left[\frac{\partial}{\partial u^{j_p}}\right]A_1^c\cdots A_{n_c}^c;A_1^o\cdots A_{n_o}^o\Big\rangle_{\Sigma_{g,b,n_c,n_o}}\end{aligned}$$

$$\Omega_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)\equiv N_{g,b,n_c,n_o}\widehat{\Omega}_p^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)$$

$$N_{g,b,n_c,n_o}=\eta_c^{-3g-3+n_c+\frac{3}{2}b+\frac{3}{4}n_o}, \eta_c\equiv -\frac{1}{2\pi i}=\frac{i}{2\pi}$$

$$\mathcal{A}_{g,b}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)=(g_S)^{-\chi_{g,b,n_c,n_o}}\int_{\mathcal{F}_{g,b,n_c,n_o}}\Omega_{d_{g,b,n_c,n_o}}^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)$$

$$\widehat{\Omega}_0^{(0,1,1,0)}(A^c)=\eta_c\langle c_0^-A^c\rangle,$$

$$\Omega_0^{(0,1,1,0)}(A^c)=\eta_cN_{0,1,1,0}\langle c_0^-A^c\rangle.$$

$$c_o\equiv N_{g,b,n_c,n_o+1}/N_{g,b,n_c,n_o}$$

$$z = G(w,u)$$

$$\Omega_1(cV_o) = c_o du \left\langle \cdots \mathcal{B}\left[\frac{\partial}{\partial u}\right] cV_o(w=0) \right\rangle$$

$$\Omega_1(cV_o) = c_o du \left\langle \cdots \oint b(z) dz \frac{\partial G(w,u)}{\partial u} cV_o(z=y) \right\rangle$$



$$\Omega_1(cV_o)=-c_oud u\frac{\partial y}{\partial u}\langle \cdots V_o(z=y)\rangle=-c_ody\langle \cdots V_o(y)\rangle$$

$$\Omega_p^{(g,b,n_c,n_o)}(A^c_1,\cdots,A^c_{n_c};A^o_1,\cdots,A^o_{n_o})\left[\frac{\partial}{\partial u^{j_1}},\cdots,\frac{\partial}{\partial u^{j_k}},\frac{\partial}{\partial y_{\alpha_1}},\cdots,\frac{\partial}{\partial y_{\alpha_\ell}},\frac{\partial}{\partial \bar{y}_{\beta_1}},\cdots,\frac{\partial}{\partial \bar{y}_{\beta_{\bar{\ell}}}}\right]$$

$$\sim N_{g,b,n_c,n_o} \left\langle {\cal B}\left[\frac{\partial}{\partial u^{j_1}}\right] \cdots {\cal B}\left[\frac{\partial}{\partial u^{j_k}}\right] \left(-\partial \xi(y_{\alpha_1})\right) \cdots \left(-\partial \xi(y_{\alpha_\ell})\right) \left(-\bar{\partial} \bar{\xi}(\bar{y}_{\beta_1})\right) \cdots \left(-\bar{\partial} \bar{\xi}(\bar{y}_{\beta_{\bar{\ell}}})\right) \prod_{\alpha=l+1}^{N_R} {\cal X}(y_\alpha) \prod_{\beta=\bar{l}+1}^{N_L} \overline{{\cal X}}(\bar{y}_\beta) \right\rangle$$

$$2(2g-2+b)=N-2n_{-1,-1}-\frac{3}{2}(n_{-1,-1/2}+n_{-1/2,-1})-n_{-1/2,-1/2},-n_{-1}-\frac{1}{2}n_{-1/2}$$

$$|\textbf{b}\rangle = -\frac{d q_r}{q_r} (e^{-\Lambda} q_r)^{L_0+\bar L_0} b_0^+ |B\rangle, q_r \in [0,1]$$

$$dy V_o(y)$$

$$\widehat{\Omega}^{(0,1,1,1)}(c\bar{c}V_c;cV_o)=-\langle cV_oc\bar{c}V_c\rangle$$

$$N_{g,b,n_c,n_o}=\eta_c^{3g-3+n_c+\frac{3}{2}b+\frac{3}{4}n_o}$$

$$\eta_c\equiv-\frac{1}{2\pi i}=\frac{i}{2\pi}$$

$$\widehat{\Omega}_p^{(g,b,n_c,n_o)}\bigl(QA^c_1,\cdots,A_{n_c};A^o_1,\cdots,A^o_{n_o}\bigr)+\cdots$$

$$+ (-1)^{A^c_1+\cdots+A^c_n+A^o_1+\cdots A^o_{n_o-1}}\widehat{\Omega}_p^{(g,b,n_c,n_o)}\bigl(A^c_1,\cdots,A_{n_c};A^o_1,\cdots,Q A^o_{n_o}\bigr)$$

$$= (-1)^p ~\mathrm{d}\widehat{\Omega}_{p-1}^{(g,b,n_c,n_o)}\bigl(A^c_1,\cdots,A_{n_c};A^o_1,\cdots,A^o_{n_o}\bigr)$$

$$\{F,G\}=\frac{\partial_r F}{\partial \psi_m}\frac{\partial_\ell G}{\partial \psi_m^*}-\frac{\partial_r F}{\partial \psi_m^*}\frac{\partial_\ell G}{\partial \psi_m},$$

$$\frac{\partial_r A}{\partial \psi} = (-1)^{\psi(A+1)} \frac{\partial_l A}{\partial \bar{\psi}}$$

$$\{F,G\}= -(-1)^{(F+1)(G+1)}\{G,F\}$$

$$\Delta F=(-1)^{\psi_m}\frac{\partial_\ell}{\partial \psi_m}\frac{\partial_\ell F}{\partial \psi_m^*}=(-1)^{\psi_m F}\frac{\partial_r}{\partial \psi_m}\Big(\frac{\partial_\ell F}{\partial \psi_m^*}\Big)$$

$$\{S_{cl}, S_{cl}\}=0$$

$$\frac{1}{2}\{S,S\}+\Delta S=0$$

$$\Delta e^S=0$$



$$\psi_m^*=\frac{\partial_\ell \Psi}{\partial \psi_m}$$

$$\tilde{\psi}_m^* = \psi_m^* - \frac{\partial_\ell \Psi}{\partial \psi_m}, \tilde{\psi}_m = \psi_m$$

$$\{F,G\}=\frac{\partial_r F}{\partial\psi^i}\omega^{ij}\frac{\partial_l G}{\partial\psi^j},\omega^{ij}=-\omega^{ji}$$

$$\Delta F=\frac{1}{2}(-1)^{\psi^i}\frac{\partial_\ell}{\partial\psi^i}\Big(\omega^{ij}\frac{\partial_\ell F}{\partial\psi^j}\Big)=\frac{1}{2}(-1)^{\psi^{i_F}}\frac{\partial_r}{\partial\psi^i}\Big(\omega^{ij}\frac{\partial_\ell F}{\partial\psi^j}\Big)$$

$${\bf Campos bos\'onicos.}$$

$$\Psi=\sum_i~\varphi_i\psi^i$$

$$|\Psi\rangle=\sum_i~(-1)^{\psi^i\epsilon_0}|\varphi_i\rangle\psi^i$$

$$b^-_0|\Psi\rangle=0,L^-_0|\Psi\rangle=0$$

$$\Psi\in\mathcal{H}_c.$$

$$\langle A,B\rangle\equiv\langle A|c^-_0|B\rangle\,|A\rangle, |B\rangle\in\mathcal{H}_c.$$

$$\langle A,B\rangle=(-1)^{(A+1)(B+1)}\langle B,A\rangle$$

$$\langle \lambda A,B\eta\rangle=\lambda\langle A,B\rangle\eta,$$

$$\partial \mathcal{V}_{g,n} = -\Delta_c \mathcal{V}_{g-1,n+2} - \frac{1}{2} \sum_{\substack{g_1+g_2=g \\ n_1+n_2=n+2}} \left\{ \mathcal{V}_{g_1,n_1}, \mathcal{V}_{g_2,n_2} \right\}_c$$

$$\text{General collection of } \mathcal{V}_{g,n} = \begin{cases} \mathcal{V}_{0,n}, & n \geqslant 3, \\ \mathcal{V}_{1,n}, & n \geqslant 1, \\ \mathcal{V}_{g,n}, & n \geqslant 0, g \geqslant 2. \end{cases}$$

$$w_1 w_2 = e^{i \theta}, 0 \leqslant \theta \leqslant 2 \pi$$

$$\{A_1,\cdots,A_n\}\equiv\sum_{g=0}^\infty~(g_s)^{-\chi_{g,n}}\{A_1,\cdots,A_n\}_g\equiv\sum_{g=0}^\infty~g_s^{2g+n-2}\int_{\mathcal{V}_{g,n}}\Omega_{d_{g,n}}^{(g,n)}(A_1,\cdots,A_n)$$

$$\langle A_0,[A_1,\cdots,A_n]\rangle=\{A_0,\cdots,A_n\},\forall A_0\in\mathcal{H}_c$$

$$[A_1,\cdots,A_n]\equiv\sum_{g=0}^\infty~g_s^{2g+n-1}[A_1,\cdots,A_n]_g$$



$$\begin{aligned} & \sum_{i=1}^N \{A'_1 \dots A'_{i-1}(QA'_i)A'_{i+1} \dots A'_N\} \\ &= -\frac{1}{2} \sum_{\substack{\ell,k \geq 0 \\ \ell+k=N}} \sum_{\substack{\{i_a; a=1, \dots, \ell\} \\ \{j_b; b=1, \dots, k\} \\ \{i_a\} \cup \{j_b\} = \{1, \dots, N\}}} \{A'_{i_1} \dots A'_{i_\ell} \varphi_s\} \{\varphi_r A'_{j_1} \dots A'_{j_k}\} \langle \varphi_s^c, \varphi_r^c \rangle \\ &\quad - \frac{1}{2} \{A'_1 \dots A'_N \varphi_s \varphi_r\} \langle \varphi_s^c, \varphi_r^c \rangle \end{aligned}$$

$$\langle \varphi_r^c, \varphi_s \rangle = \langle \varphi_s, \varphi_r^c \rangle = \delta_{rs}$$

$$\langle \varphi_i, \varphi_j \rangle = \langle \varphi_j, \varphi_i \rangle, \text{ and } \langle \varphi_i^c, \varphi_j^c \rangle = \langle \varphi_j^c, \varphi_i^c \rangle.$$

$$\begin{aligned} Q[A'_1 \dots A'_N] &= - \sum_{i=1}^N [A'_1 \dots A'_{i-1}(QA'_i)A'_{i+1} \dots A'_N] \\ &- \sum_{\substack{\ell,k \geq 0 \\ \ell+k=N}} \sum_{\substack{\{i_a; a=1, \dots, \ell\} \\ \{j_b; b=1, \dots, k\} \\ \{i_a\} \cup \{j_b\} = \{1, \dots, N\}}} [A'_{i_1} \dots A'_{i_\ell} [A'_{j_1} \dots A'_{j_k}]] - \frac{1}{2} [A'_1 \dots A'_N \varphi_s \varphi_r] \langle \varphi_s^c, \varphi_r^c \rangle \end{aligned}$$

$$S = \frac{1}{2} \langle \Psi, Q\Psi \rangle + \sum_{n=1}^{\infty} \frac{1}{n!} \{\Psi^n\}$$

$$\Psi = \sum_r \varphi_r \psi^r \rightarrow \delta \Psi = \sum_r \varphi_r \delta \psi^r$$

$$\delta F = \sum_k \frac{\partial^r F}{\partial \psi^k} \delta \psi^k = \langle F_R, \delta \Psi \rangle \rightarrow \frac{\partial^r F}{\partial \psi^k} = \langle F_R, \varphi_k \rangle$$

$$\delta G = \sum_k \delta \psi^k \frac{\partial^l G}{\partial \psi^k} = \langle \delta \Psi, G_L \rangle \rightarrow \frac{\partial^l G}{\partial \psi^k} = (-1)^{\varphi_k} \langle \varphi_k, G_L \rangle$$

$$\{F, G\} = -\langle F_R, G_L \rangle$$

$$\omega^{rs} = (-1)^{\varphi_s+1} \langle \varphi_r^c, \varphi_s^c \rangle, \omega_{rs} = (-1)^{\varphi_r+1} \langle \varphi_r, \varphi_s \rangle,$$

$$Q\Psi + \sum_{n=1}^{\infty} \frac{1}{n!} [\Psi^n] = 0$$

$$S_{\text{cl}} = \frac{1}{2} \langle \Psi, Q\Psi \rangle + \sum_{n=3}^{\infty} \frac{1}{n!} \{\Psi^n\}_0$$

$$\delta \Psi = Q\Lambda + \sum_{n=1}^{\infty} \frac{1}{n!} [\Lambda, \Psi^n]_0 = Q\Lambda + [\Lambda, \Psi]_0 + \frac{1}{2!} [\Lambda, \Psi, \Psi]_0 + \frac{1}{3!} [\Lambda, \Psi, \Psi, \Psi]_0 + \dots$$



$$\begin{aligned}\Psi \in \mathcal{H}_c &\equiv \mathcal{H}_{-1,-1} \oplus \mathcal{H}_{-1,-1/2} \oplus \mathcal{H}_{-1/2,-1} \oplus \mathcal{H}_{-1/2,-1/2}, \\ \widetilde{\Psi} \in \widetilde{\mathcal{H}}_c &\equiv \mathcal{H}_{-1,-1} \oplus \mathcal{H}_{-1,-3/2} \oplus \mathcal{H}_{-3/2,-1} \oplus \mathcal{H}_{-3/2,-3/2}.\end{aligned}$$

$$\sum_{i=1}^N\{A'_1\dots A'_{i-1}(QA'_i)A'_{i+1}\dots A'_N\}$$

$$= -\frac{1}{2}\sum_{\substack{\ell,k\geqslant 0\\ \ell+k=N}}\sum_{\substack{\{i_a;a=1,\dots,\ell\}\\ \{j_b;b=1,\dots,k\}\\ \{i_a\}\cup\{j_b\}=\{1,\dots,N\}}} \{A'_{i_1}\dots A'_{i_\ell}\varphi_s\}\{\varphi_rA'_{j_1}\dots A'_{j_k}\}\langle\varphi_s^c,\mathcal{G}\varphi_r^c\rangle$$

$$-\frac{1}{2}\{A'_1\dots A'_N\varphi_s\varphi_r\}\langle\varphi_s^c,\mathcal{G}\varphi_r^c\rangle$$

$$\mathcal{X}_0 = \oint \frac{dz}{z} \mathcal{X}(z), \overline{\mathcal{X}}_0 = \oint \frac{d\bar{z}}{\bar{z}} \overline{\mathcal{X}}(\bar{z})$$

$$\mathcal{G} = \begin{cases} \mathbf{1} \text{ on } \mathcal{H}_{-1,-1} \\ \mathcal{X}_0 \text{ on } \mathcal{H}_{-1,-3/2} \\ \overline{\mathcal{X}}_0 \text{ on } \mathcal{H}_{-3/2,-1} \\ \mathcal{X}_0 \overline{\mathcal{X}}_0 \text{ on } \mathcal{H}_{-3/2,-3/2} \end{cases}$$

$$S = -\frac{1}{2}\langle \widetilde{\Psi}, Q\mathcal{G}\widetilde{\Psi} \rangle + \langle \widetilde{\Psi}, Q\Psi \rangle + \sum_{n=1}^{\infty} \frac{1}{n!} \{\Psi^n\}$$

$$\delta F = \langle F_R, \delta \widetilde{\Psi} \rangle + \langle \tilde{F}_R, \delta \Psi \rangle = \langle \delta \widetilde{\Psi}, F_L \rangle + \langle \delta \Psi, \tilde{F}_L \rangle$$

$$\{F,G\}=-\langle F_R,\tilde{G}_L\rangle-\langle \tilde{F}_R,G_L\rangle-\langle \tilde{F}_R,\mathcal{G}\tilde{G}_L\rangle$$

$$\langle A_1,[A_1,\cdots,A_n]\rangle=\{A_1,\cdots,A_n\}$$

$$\delta|\widetilde{\Psi}\rangle = Q|\widetilde{\Lambda}\rangle + \sum_{n=2}^{\infty} \frac{1}{n!} [\Lambda \Psi^n]_0, \delta|\Psi\rangle = Q|\Lambda\rangle + \sum_{n=2}^{\infty} \frac{1}{n!} \mathcal{G}[\Lambda \Psi^n]_0$$

Campos heteróticos.

$$\begin{aligned}\Psi \in \mathcal{H}_c &\equiv \mathcal{H}_{-1} \oplus \mathcal{H}_{-1/2} \\ \widetilde{\Psi} \in \widetilde{\mathcal{H}}_c &\equiv \mathcal{H}_{-1} \oplus \mathcal{H}_{-3/2}\end{aligned}$$

$$\mathcal{G} = \begin{cases} \mathbf{1} \text{ on } \mathcal{H}_{-1} \\ \mathcal{X}_0 \text{ on } \mathcal{H}_{-3/2} \end{cases}$$

$$S = -\frac{1}{2}\langle \widetilde{\Psi}, Q\mathcal{G}\widetilde{\Psi} \rangle + \langle \widetilde{\Psi}, Q\Psi \rangle + \sum_{n=1}^{\infty} \frac{1}{n!} \{\Psi^n\}$$

$$\Omega_{n-3}^{o(0,n)}(A_1,\cdots,A_n)\equiv K^{-1}\widehat{\Omega}_{n-3}^{(0,1,0,n)}(A_1,\cdots,A_n)$$



$$\partial \mathcal{V}^o_{0,n} = -\frac{1}{2}\sum_{\substack{n_1,n_2 \\ n_1+n_2=n+2}} \left\{\mathcal{V}^o_{0,n_1},\mathcal{V}^o_{0,n_2}\right\}_o$$

$$\{A_1,\cdots,A_n\}\equiv g_o^{n-2}\int_{\mathcal{V}_{0,n}^o}\Omega_{n-3}^{o(0,n)}(A_1,\cdots,A_n)$$

$$S_o=\frac{1}{2}\langle\psi_o,Q\psi_o\rangle'+\sum_{n=3}^\infty\frac{1}{n!}\{\psi_o^n\}$$

$$\mathcal{T}=\frac{1}{2\pi^2g_o^2}$$

$$\{A_1\cdots A_iA_{i+1}\cdots A_n\}=(-1)^{A_iA_{i+1}+1}\{A_1\cdots A_{i+1}A_i\cdots A_n\}.$$

$$\begin{aligned} Q[A'_1\cdots A'_N] &= \sum_{i=1}^N\; (-1)^{i-1}[A'_1\ldots A'_{i-1}(QA'_i)A'_{i+1}\ldots A'_N]\\ &\quad + \sum_{\substack{\ell,k\geqslant 0\\ \ell+k=N}}\; \sum_{\substack{\{i_a;a=1,\ldots \ell\}\\ \{j_b;b=1,\ldots k\}\\ \{i_a\}\cup\{j_b\}=\{1,\ldots N\}}} \Big[[A'_{j_1}\ldots A'_{j_k}]A'_{i_1}\ldots A'_{i_\ell}\Big]\end{aligned}$$

$$\langle A_0,[A_1,\cdots,A_n]\rangle'=\{A_0,A_1,\cdots,A_n\}$$

$$\delta F=\langle F_R^o,\delta\psi_o\rangle'=\langle\delta\psi_o,F_L^o\rangle',$$

$$\{F,G\}=-\langle F_R^o,G_L^o\rangle'$$

$$\delta \Phi = Q \Lambda + [\Lambda \Phi] + \frac{1}{2!}[\Lambda,\Phi,\Phi] + \cdots$$

$$[A_1,A_2]=A_1\star A_2-(-1)^{A_1A_2}A_2\star A_1$$

$$A_1\star(A_2\star A_3)=(A_1\star A_2)\star A_3$$

$$[A_1,A_2,\cdots,A_n]=0\;\;\text{for}\;\;n\geqslant 3$$

$$\langle A_1,A_2\star A_3\rangle'=\langle f_1\circ A_1(0)f_2\circ A_2(0)f_3\circ A_3(0)\rangle'_{\rm UHP}$$

$$\begin{gathered}f_1(w_1)=h^{-1}\left(e^{2\pi i/3}\big(h(w_1)\big)^{2/3}\right)\\ f_2(w_2)=h^{-1}\left(\big(h(w_2)\big)^{2/3}\right), h(u)\equiv\frac{1+iu}{1-iu}\\ f_3(w_3)=h^{-1}\left(e^{-2\pi i/3}\big(h(w_3)\big)^{2/3}\right)\end{gathered}$$

$$\begin{gathered}\psi_o\in\mathcal{H}_o\equiv\mathcal{H}_{-1}\oplus\mathcal{H}_{-1/2},\\\tilde{\psi}_o\in\widetilde{\mathcal{H}}_o\equiv\mathcal{H}_{-1}\oplus\mathcal{H}_{-3/2}.\end{gathered}$$



$$S = -\frac{1}{2}\langle \tilde{\psi}_o, Q\mathcal{G}\tilde{\psi}_o\rangle' + \langle \tilde{\psi}_o, Q\psi_o\rangle' + \sum_{n=1}^{\infty}\frac{1}{n!}\{\psi_o^n\}$$

$$\begin{array}{c}\Psi_c\in\mathcal{H}_c,\widetilde{\Psi}_c\in\widetilde{\mathcal{H}}_c,\\\Psi_o\in\mathcal{H}_o,\widetilde{\Psi}_o\in\widetilde{\mathcal{H}}_o\end{array}$$

$$\begin{aligned}\partial \mathcal{V}_{g,b,n_c,n_o}&=-\Delta_c\mathcal{V}_{g-1,b,n_c+2,n_o}-\Delta'_o\mathcal{V}_{g,b-1,n_c,n_o+2}-\Delta_o\mathcal{V}_{g-1,b+1,n_c,n_o+2}\\&\quad-\frac{1}{2}\sum_{\substack{g_1+g_2=g,b_1+b_2=b\\n_{c1}+n_{c2}=n_c+2,n_{o1}+n_{o2}=n_o}}\{\mathcal{V}_{g_1,b_1,n_{c1},n_{o1}},\mathcal{V}_{g_2,b_2,n_{c2},n_{o2}}\}_c\\&\quad-\frac{1}{2}\sum_{\substack{g_1+g_2=g,b_1+b_2=b+1\\n_{c1}+n_{c2}=n_c,n_{o1}+n_{o2}=n_o+2}}\{\mathcal{V}_{g_1,b_1,n_{c1},n_{o1}},\mathcal{V}_{g_2,b_2,n_{c2},n_{o2}}\}_o\end{aligned}$$

$$\begin{aligned}\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\}&=\sum_{g=0}^{\infty}\sum_{b=0}^{\infty}\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\}_{g,b}\\\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\}_{g,b}&=(g_s)^{-\chi_{g,b,n_c,n_o}}\int_{\mathcal{V}_{g,b,n_c,n_o}}\Omega_{d_{g,b,n_c,n_o}}^{(g,b,n_c,n_o)}(A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o)\end{aligned}$$

$$\{\tilde{A}^c\}_D=\Omega_0^{(0,1,1,0)}(\tilde{A}^c)\equiv-\frac{1}{2\pi i}N_{0,1,1,0}\langle\hat{\mathcal{G}}\tilde{A}^c|c_0^-e^{-\Lambda(L_0+\bar{L}_0)}|B\rangle,$$

$$\hat{\mathcal{G}}\equiv\begin{cases}1\text{ on }\mathcal{H}_{-1,-1}\\\frac{1}{2}(\mathcal{X}_0+\overline{\mathcal{X}_0})\text{ on }\mathcal{H}_{-3/2,-3/2}\end{cases}$$

$$\{A^c\}_D=\Omega_0^{(0,1,1,0)}(A^c)\equiv-\frac{1}{2\pi i}N_{0,1,1,0}\langle A^c|c_0^-e^{-\Lambda(L_0+\bar{L}_0)}|B\rangle.$$

$$\begin{aligned}\langle A_0^c|c_0^-|[A_1^c\cdots A_N^c;A_1^o\cdots A_M^o]_c\rangle&=\{A_0^cA_1^c\cdots A_N^c;A_1^o\cdots A_M^o\},\forall |A_0^c\rangle\in\mathcal{H}_c,\\\langle A_0^o|[A_1^c\cdots A_N^c;A_1^o\cdots A_M^o]_o\rangle&=\{A_1^c\cdots A_N^c;A_0^oA_1^o\cdots A_M^o\},\forall |A_0^o\rangle\in\mathcal{H}_o,\\\langle\tilde{A}^c|c_0^-|[]_c\rangle&=\{\tilde{A}^c\}_D,(4.65)\end{aligned}$$

$$\{(Q\tilde{A}^c)\}_D=0$$

$$\sum_{i=1}^N\{A_1^c\cdots A_{i-1}^c(QA_i^c)A_{i+1}^c\cdots A_N^c;A_1^o\cdots A_M^o\}+\sum_{j=1}^M\{A_1^c\cdots A_N^c;A_1^o\cdots A_{j-1}^o(QA_j^o)A_{j+1}^o\cdots A_M^o\}(-1)^{j-1}$$

$$\begin{aligned}&=-\frac{1}{2}\sum_{k=0}^N\sum_{\ell=0}^M\sum_{\substack{\{i_1,\cdots,i_k\}\subset\{1,\cdots,N\}\\ \{j_1,\cdots,j_\ell\}\subset\{1,\cdots,M\}}}(\{A_{i_1}^c\cdots A_{i_k}^c\mathcal{B}^c;A_{j_1}^o\cdots A_{j_\ell}^o\}+\{A_{i_1}^c\cdots A_{i_k}^c;\mathcal{B}^oA_{j_1}^o\cdots A_{j_\ell}^o\})\\&\quad-\{[A_1^c\cdots A_N^c;A_1^o\cdots A_M^o]_c\}_D-\frac{1}{2}\{A_1^c\cdots A_N^c\varphi_s\varphi_r;A_1^o\cdots A_M^o\}\langle\varphi_s^c,\mathcal{G}\varphi_r^c\rangle\end{aligned}$$

$$-\frac{1}{2}(-1)^{\hat{\varphi}_s}\{A_1^c\cdots A_N^c;\hat{\varphi}_s\hat{\varphi}_rA_1^o\cdots A_M^o\}\langle\hat{\varphi}_s^c,\mathcal{G}\hat{\varphi}_r^c\rangle$$



$$\mathcal{B}^c \equiv \mathcal{G}[A_{\bar{l}_1}^c \cdots A_{\bar{l}_{N-k}}^c; A_{\bar{j}_1}^o \cdots A_{\bar{j}_{M-\ell}}^o]_c, \mathcal{B}^o \equiv \mathcal{G}[A_{\bar{l}_1}^c \cdots A_{\bar{l}_{N-k}}^c; A_{\bar{j}_1}^o \cdots A_{\bar{j}_{M-\ell}}^o]_o$$

$$\{i_1, \dots, i_k\} \cup \{\bar{i}_1, \dots, \bar{i}_{N-k}\} = \{1, \dots, N\}, \{j_1, \dots, j_\ell\} \cup \{\bar{j}_1, \dots, \bar{j}_{M-\ell}\} = \{1, \dots, M\},$$

$$S = -\frac{1}{2}\langle \tilde{\Psi}_c, Q\mathcal{G}\tilde{\Psi}_c \rangle + \langle \tilde{\Psi}_c, Q\Psi_c \rangle - \frac{1}{2}\langle \tilde{\Psi}_o, Q\mathcal{G}\tilde{\Psi}_o \rangle + \langle \tilde{\Psi}_o, Q\Psi_o \rangle + \{\tilde{\Psi}_c\}_D + \sum_{N=0}^{\infty} \sum_{M=0}^{\infty} \frac{1}{N! M!} \{(\Psi_c)^N; (\Psi_o)^M\}$$

$$\delta F = \langle F_R^c, \delta \tilde{\Psi}_c \rangle + \langle \tilde{F}_R^c, \delta \Psi_c \rangle + \langle F_R^o, \delta \tilde{\Psi}_o \rangle + \langle \tilde{F}_R^o, \delta \Psi_o \rangle = \langle \delta \tilde{\Psi}_c, F_L^c \rangle + \langle \delta \Psi_c, \tilde{F}_L^c \rangle + \langle \delta \tilde{\Psi}_o, F_L^o \rangle + \langle \delta \Psi_o, \tilde{F}_L^o \rangle$$

$$\{F, G\} = -(\langle F_R^c, \tilde{G}_L^c \rangle + \langle \tilde{F}_R^c, G_L^c \rangle + \langle \tilde{F}_R^c, \mathcal{G}\tilde{G}_L^c \rangle) - (\langle F_R^o, \tilde{G}_L^o \rangle + \langle \tilde{F}_R^o, G_L^o \rangle + \langle \tilde{F}_R^o, \mathcal{G}\tilde{G}_L^o \rangle)$$

$$S = \frac{1}{2}\langle \Psi_c, Q\Psi_c \rangle + \frac{1}{2}\langle \Psi_o, Q\Psi_o \rangle + \{\Psi_c\}_D + \sum_{N=0}^{\infty} \sum_{M=0}^{\infty} \frac{1}{N! M!} \{(\Psi_c)^N; (\Psi_o)^M\}$$

$$\begin{aligned} S_{oc} &= \frac{1}{2}\langle \Psi_o, Q\Psi_o \rangle + \sum_{n=3}^{\infty} \frac{1}{n!} g_s^{(n-2)/2} \int \Omega_{n-3}^{(0,1,0,n)}(\Psi_o, \dots, \Psi_o) \\ &= \frac{1}{2}K\langle \Psi_o, Q\Psi_o \rangle' + \sum_{n=3}^{\infty} \frac{1}{n!} g_s^{(n-2)/2} N_{0,1,0,n} K \int \Omega_{n-3}^{o(0,n)}(\Psi_o, \dots, \Psi_o) \\ S_o &= \frac{1}{2}\langle \psi_o, Q\psi_o \rangle' + \sum_{n=3}^{\infty} \frac{1}{n!} g_o^{n-2} \int \Omega_{n-3}^{o(0,n)}(\psi_o, \dots, \psi_o) \end{aligned}$$

$$\Psi_o = K^{-1/2} \psi_o$$

$$S_{oc} = \frac{1}{2}\langle \psi_o, Q\psi_o \rangle' + \sum_{n=3}^{\infty} \frac{1}{n!} g_s^{(n-2)/2} K^{(2-n)/2} \eta_c^{3(n-2)/4} \int \Omega_{n-3}^{o(0,n)}(\psi_o, \dots, \psi_o)$$

$$g_o = g_s^{1/2} K^{-1/2} \eta_c^{3/4}$$

$$g_o^2 = (g_s/K) \eta_c^{3/2}$$

Partículas supermasivas – campos de gauge – transformaciones de gauge – campos cuánticos gravitacionales y simetría – osciladores y propagadores.

$$|\psi_o\rangle = (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \left(A_\mu(k) c_1 \alpha_{-1}^\mu - i \sqrt{\frac{1}{2}} B(k) c_0 \right) |k\rangle,$$

$$|\epsilon\rangle = (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \frac{i}{\sqrt{2}} \epsilon(k) |k\rangle.$$

$$\phi(x) = (\alpha')^{(p+1)/2} \int \frac{d^{p+1}k}{(2\pi)^{p+1}} \tilde{\phi}(k) e^{ikx}, \text{ so that } ik_\mu \leftrightarrow \partial_\mu$$



$$\mathrm{bpz}(\phi_n) = (-1)^{n+d}\phi_{-n}$$

$$\langle \psi_o|=(\alpha')^{(p+1)/2}\int\frac{d^{p+1}k}{(2\pi)^{p+1}}\langle k|\Bigg(A_\mu(k)c_{-1}\alpha_1^\mu+i\sqrt{\frac{1}{2}}B(k)c_0\Bigg),$$

$$Q=\sum_nc_n(L^m_{-n}-\delta_{n,0})+\frac{1}{2}\sum_{m,n}(m-n)\colon c_mc_nb_{-n-m}\colon$$

$$L^m_n=\frac{1}{2}\sum_m~\alpha_m\cdot\alpha_{n-m},(n\neq 0), L^m_0=\frac{1}{2}\alpha_0^2+\sum_{n\geqslant 1}~\alpha_{-n}\cdot\alpha_n,\alpha_0^\mu=\sqrt{2\alpha'}k^\mu$$

$$Q=c_0(\alpha'k^2-1)+c_0(\alpha_{-1}\cdot\alpha_1+b_{-1}c_1+c_{-1}b_1)+\sqrt{2\alpha'}k_\mu\big(\alpha_{-1}^\mu c_1+c_{-1}\alpha_1^\mu\big)-2b_0c_{-1}c_1+\cdots$$

$$Qc_1\alpha_{-1}^\mu|k\rangle=\big(\alpha'k^2c_0c_1\alpha_{-1}^\mu+\sqrt{2\alpha'}k^\mu c_{-1}c_1\big)|k\rangle$$

$$\begin{aligned} Qc_0|k\rangle &= \big(\sqrt{2\alpha'}k_\mu\alpha_{-1}^\mu c_1c_0-2c_{-1}c_1\big)|k\rangle \\ Q|k\rangle &= \big(\alpha'k^2c_0+\sqrt{2\alpha'}k_\mu\alpha_{-1}^\mu c_1\big)|k\rangle \end{aligned}$$

$$\langle k|c_{-1}c_0c_1|k'\rangle'=-(2\pi)^{p+1}(\alpha')^{-(p+1)/2}\delta^{(p+1)}(k+k').$$

$$\begin{aligned} S_2 &= \frac{1}{2}\langle \psi_o|Q|\psi_o\rangle' \\ &= (\alpha')^{(p+1)/2}\int\frac{d^{p+1}k}{(2\pi)^{p+1}}\Big(-\frac{1}{2}A^\mu(-k)\alpha'p^2A_\mu(k)-\sqrt{\alpha'}A^\mu(-k)ik_\mu B(k) \\ &\quad -\frac{1}{2}B(-k)B(k)\Big) \\ S_2 &= (\alpha')^{-(p+1)/2}\int~d^{p+1}x\left(\frac{\alpha'}{2}A^\mu\Box A_\mu-\sqrt{\alpha'}A^\mu\partial_\mu B-\frac{1}{2}B^2\right) \\ \delta A_\mu(k) &= i\sqrt{\alpha'}k_\mu\epsilon(k), \delta B(k)=-\alpha'k^2\epsilon(k) \\ \delta A_\mu &= \sqrt{\alpha'}\partial_\mu\epsilon, \delta B=\alpha'\Box\epsilon \end{aligned}$$

$$\begin{aligned} S &= (\alpha')^{-(p-1)/2}\int~d^{p+1}x\left(\frac{1}{2}A_\mu\Box A^\mu+\frac{1}{2}(\partial\cdot A)^2\right)= (\alpha')^{-(p-1)/2}\int~d^{p+1} \\ &\quad x\left(-\frac{1}{2}\partial_\mu A_\nu\partial^\mu A^\nu+\frac{1}{2}\partial_\mu A_\nu\partial^\nu A^\mu\right)= (\alpha')^{-(p-1)/2}\int~d^{p+1}x\left(-\frac{1}{4}F^{\mu\nu}F_{\mu\nu}\right) \end{aligned}$$

$$S_{\text{ws}}=-\frac{1}{4\pi\alpha'}\int~dxdy g_{\mu\nu}\big(\partial_xX^\mu\partial_xX^\nu+\partial_yX^\mu\partial_yX^\nu\big)=-\frac{1}{\pi\alpha'}\int~dxdy g_{\mu\nu}\partial X^\mu\bar{\partial}X^\nu,z\equiv x+iy$$



$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}$$

$$S_{\text{ws}}|_h = - \frac{1}{\pi \alpha'} \int \; dx dy h_{\mu\nu} \partial X^\mu \bar{\partial} X^\nu$$

$$-g_s\frac{1}{\pi}\int \; dxdy V$$

$$\mathcal{O}_h=\frac{1}{g_s\alpha'}h_{\mu\nu}c\bar{c}\partial X^\mu\bar{\partial} X^\nu=\frac{1}{g_s}\Big(-\frac{1}{2}h_{\mu\nu}\Big)c\bar{c}i\sqrt{\frac{2}{\alpha'}}\partial X^\mu i\sqrt{\frac{2}{\alpha'}}\bar{\partial} X^\nu$$

$$\begin{aligned} |\Psi\rangle=(\alpha')^{D/2}\frac{1}{g_s}\int\;\frac{d^Dk}{(2\pi)^D}\bigg(&-\frac{1}{2}e_{\mu\nu}(k)\alpha_{-1}^\mu\bar{\alpha}_{-1}^\nu c_1\bar{c}_1+e(k)c_1c_{-1}+\bar{e}(k)\bar{c}_1\bar{c}_{-1}\\ &+i\sqrt{\frac{1}{2}}\big(f_\mu(k)c_0^+c_1\alpha_{-1}^\mu+\bar{f}_\mu(k)c_0^+\bar{c}_1\bar{\alpha}_{-1}^\mu\big)\bigg)|k\rangle\end{aligned}$$

$$e_{\mu\nu}=h_{\mu\nu}+b_{\mu\nu}, \text{ with } h_{\mu\nu}=h_{\nu\mu}, b_{\mu\nu}=-b_{\nu\mu}$$

$$\langle k|c_{-1}\bar{c}_{-1}c_0^-c_0^+c_1\bar{c}_1|k'\rangle=-\frac{1}{2}(\alpha')^{-D/2}(2\pi)^D\delta^{(D)}(k+k').$$

$$\mathrm{bpz}(\phi_n)=(-1)^d\phi_{-n},$$

$$\begin{aligned} \langle\Psi|=(\alpha')^{D/2}\frac{1}{g_s}\int\;\frac{d^Dk}{(2\pi)^D}\langle k|\Big(&-\frac{1}{2}e_{\mu\nu}(k)\alpha_1^\mu\bar{\alpha}_1^\nu c_{-1}\bar{c}_{-1}+e(k)c_{-1}c_1+\bar{e}(k)\bar{c}_{-1}\bar{c}_1\\ &-i\sqrt{\frac{1}{2}}\big(f_\mu(k)c_0^+c_{-1}\alpha_1^\mu+\bar{f}_\mu(k)c_0^+\bar{c}_{-1}\bar{\alpha}_1^\mu\big)\end{aligned}$$

$$S^{(2)}=\frac{1}{2}\langle\Psi|c_0^-Q|\Psi\rangle$$

$$Q=c_0^+\left(\frac{\alpha'}{2}k^2-2\right)+c_0^+\big(\alpha_{-1}\cdot\alpha_1+b_{-1}c_1+c_{-1}b_1+\bar{\alpha}_{-1}\cdot\bar{\alpha}_1+\bar{b}_{-1}\bar{c}_1+\bar{c}_{-1}\bar{b}_1\big)+\sqrt{\frac{\alpha'}{2}}k$$

$$\cdot(\alpha_{-1}c_1+c_{-1}\alpha_1)+\sqrt{\frac{\alpha'}{2}}k\cdot(\bar{\alpha}_{-1}\bar{c}_1+\bar{c}_{-1}\bar{\alpha}_1)-b_0^+(c_{-1}c_1+\bar{c}_{-1}\bar{c}_1)+\cdots$$

$$\begin{aligned} S^{(2)}=(\alpha')^{-D/2}\frac{1}{8g_s^2}\int\;d^Dx\bigg[\frac{\alpha'}{4}e^{\mu\nu}\Box e_{\mu\nu}+2\alpha'\bar{e}\Box e-f^\mu f_\mu-\bar{f}^\mu\bar{f}_\mu\\ -\sqrt{\alpha'}f^\mu\big(\partial^\nu e_{\mu\nu}-2\partial_\mu\bar{e}\big)+\sqrt{\alpha'}\bar{f}^\nu\big(\partial^\mu e_{\mu\nu}+2\partial_\nu e\big)\big]$$

$$|\Lambda\rangle=(\alpha')^{D/2}\frac{1}{g_s}\int\;\frac{d^Dk}{(2\pi)^D}\Big(\frac{i}{\sqrt{2}}\lambda_\mu(k)\alpha_{-1}^\mu c_1-\frac{i}{\sqrt{2}}\bar{\lambda}_\mu(k)\bar{\alpha}_{-1}^\mu\bar{c}_1+\mu(k)c_0^+\Big)|k\rangle.$$

$$\delta|\Psi\rangle=Q|\Lambda\rangle$$



$$\delta e_{\mu\nu} = \sqrt{\alpha'} (\partial_\mu \bar{\lambda}_\nu + \partial_\nu \lambda_\mu)$$

$$\delta f_\mu = -\frac{1}{2}\alpha' \square \lambda_\mu + \sqrt{\alpha'} \partial_\mu \mu$$

$$\begin{aligned}\delta \bar{f}_\nu &= \frac{1}{2}\alpha' \square \bar{\lambda}_\nu + \sqrt{\alpha'} \partial_\nu \mu \\ \delta e &= -\frac{1}{2}\sqrt{\alpha'} \partial \cdot \lambda + \mu \\ \delta \bar{e} &= \frac{1}{2}\sqrt{\alpha'} \partial \cdot \bar{\lambda} + \mu\end{aligned}$$

$$d=\frac{1}{2}(e-\bar{e}), \text{ and } \chi=\frac{1}{2}(e+\bar{e})$$

$$\begin{aligned}\delta d &= -\frac{1}{4}\sqrt{\alpha'}(\partial \cdot \lambda + \partial \cdot \bar{\lambda}) \\ \delta \chi &= -\frac{1}{4}\sqrt{\alpha'}(\partial \cdot \lambda - \partial \cdot \bar{\lambda}) + \mu\end{aligned}$$

$$\chi = 0$$

$$f_\mu = -\frac{1}{2}\sqrt{\alpha'}(\partial^\nu e_{\mu\nu} - 2\partial_\mu \bar{e}), \bar{f}_\nu = \frac{1}{2}\sqrt{\alpha'}(\partial^\mu e_{\mu\nu} + 2\partial_\nu e)$$

$$S^{(2)} = (\alpha')^{-(D-2)/2} \frac{1}{8g_s^2} \int d^D x \left[\frac{1}{4} e_{\mu\nu} \square e^{\mu\nu} + \frac{1}{4} (\partial^\nu e_{\mu\nu})^2 + \frac{1}{4} (\partial^\mu e_{\mu\nu})^2 - 2d\partial^\mu \partial^\nu e_{\mu\nu} - 4d \square d \right]$$

$$\begin{aligned}\delta e_{\mu\nu} &= \sqrt{\alpha'} (\partial_\nu \lambda_\mu + \partial_\mu \bar{\lambda}_\nu) \\ \delta d &= -\frac{1}{4}\sqrt{\alpha'} (\partial \cdot \lambda + \partial \cdot \bar{\lambda})\end{aligned}$$

$$S^{(2)} = (\alpha')^{-(D-2)/2} \frac{1}{8g_s^2} \int d^D x L[h, b, d]$$

$$L[h, b, d] = \frac{1}{4} h^{\mu\nu} \partial^2 h_{\mu\nu} + \frac{1}{2} (\partial^\nu h_{\mu\nu})^2 - 2d\partial^\mu \partial^\nu h_{\mu\nu} - 4d\partial^2 d + \frac{1}{4} b^{\mu\nu} \partial^2 b_{\mu\nu} + \frac{1}{2} (\partial^\nu b_{\mu\nu})^2$$

$$S_{\text{st}} = \frac{1}{2\kappa^2} \int d^D x \sqrt{-g} e^{-2\phi} \left[R - \frac{1}{12} H^2 + 4(\partial\phi)^2 \right]$$

$$g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}, \phi = d + \frac{1}{4}\eta^{\mu\nu}h_{\mu\nu}, H_{\mu\nu\rho} = \partial_\mu b_{\nu\rho} + \dots$$

$$S_{\text{st}}^{(2)} = \frac{1}{2\kappa^2} \int d^D x L[h, b, d]$$

$$\kappa = (\alpha')^{(D-2)/4} (2g_s)$$



$$\begin{array}{lcl} \delta h_{\mu\nu} & = & \partial_\mu \epsilon_\nu + \partial_\nu \epsilon_\mu \\ \delta b_{\mu\nu} = -\partial_\mu \tilde{\epsilon}_\nu + \partial_\nu \tilde{\epsilon}_\mu & & (4.124) \\ \delta d = -\frac{1}{2}\partial\cdot\epsilon & & (4.124) \end{array}$$

$$\epsilon_\mu \equiv \frac{1}{2}\sqrt{\alpha'}(\lambda_\mu+\bar{\lambda}_\mu), \tilde{\epsilon}_\mu \equiv \frac{1}{2}\sqrt{\alpha'}(\lambda_\mu-\bar{\lambda}_\mu)$$

$$b_0^+|\Psi\rangle=0$$

$$S=\frac{1}{2}\langle\Psi|c_0^-c_0^+L_0^+|\Psi\rangle+\sum_{n=1}^\infty\frac{1}{n!}\{\Psi^n\}$$

$$\mathcal{P}_b=-b_0^+b_0^-(L_0+\bar L_0)^{-1}\delta_{L_0,\bar L_0}$$

$$\mathcal{P}_b=-\frac{1}{2\pi}b_0^+b_0^-\int_0^{2\pi}d\theta\int_0^\infty ds e^{-s(L_0+\bar L_0)}e^{i\theta(L_0-\bar L_0)}=\frac{1}{\pi}b_0\bar b_0\int_{|q|\leqslant 1}\frac{d^2q}{|q|^2}q^{L_0}\bar q^{\bar L_0}$$

$$q\equiv e^{-s+i\theta}, d^2q=d\theta ds |q|^2\equiv \frac{i}{2}dq\wedge d\bar{q}$$

$$w_1 w_2 = q$$

$$(6g_1-6+2n_1)+(6g_2-6+2n_2)+2$$

$$\mathcal{P}_b=\int_{|q|\leqslant 1}\left[\left(-\frac{1}{2\pi i}\right)b_0\bar b_0\frac{dq\wedge d\bar q}{|q|^2}\right]q^{L_0}\bar q^{\bar L_0}$$

$$\Omega_{\mathcal{P}_b}=-\frac{1}{2\pi i}b_0\bar b_0\frac{dq\wedge d\bar q}{|q|^2}$$

$$\frac{\partial F}{\partial q}=\frac{1}{w_2}=\frac{w_1}{q}, \frac{\partial \bar F}{\partial q}=0;\, \frac{\partial F}{\partial \bar q}=0, \frac{\partial \bar F}{\partial \bar q}=\frac{\bar w_1}{\bar q}.$$

$$\widehat{\Omega}_b=\mathcal{B}\left[\frac{\partial}{\partial q}\right]\mathcal{B}\left[\frac{\partial}{\partial \bar{q}}\right]dq\wedge d\bar{q}=\frac{1}{q}\oint~w_1b(w_1)dw_1\times\frac{1}{\bar{q}}\oint~\bar{w}_1\bar{b}(\bar{w}_1)d\bar{w}_1\times dq\wedge d\bar{q}=\frac{b_0}{q}\frac{\bar{b}_0}{\bar{q}}dq\wedge d\bar{q}$$

$$\Omega_{\mathcal{P}_b}=-\frac{1}{2\pi i}\hat{\Omega}_b$$

$$\left(-\frac{1}{2\pi i}\right)^{3g_1-3+n_1}\left(-\frac{1}{2\pi i}\right)^{3g_2-3+n_2}\left(-\frac{1}{2\pi i}\right)^1=\left(-\frac{1}{2\pi i}\right)^{3(g_1+g_2)-3+(n_1+n_2-2)}$$

$$S=-\frac{1}{2}\langle\widetilde{\Psi}|c_0^-c_0^+L_0^+\mathcal{G}|\widetilde{\Psi}\rangle+\langle\widetilde{\Psi}|c_0^-c_0^+L_0^+|\Psi\rangle+\sum_{n=1}^\infty\frac{1}{n!}\{\Psi^n\}$$

$$K_0=c_0^-c_0^+L_0^+\left(\begin{matrix} \mathcal{G} & -1 \\ -1 & 0 \end{matrix}\right)$$

$$\mathcal{P}_s=-b_0^+b_0^-(L_0+\bar L_0)^{-1}\delta_{L_0,\bar L_0}\left(\begin{matrix} 0 & 1 \\ 1 & \mathcal{G} \end{matrix}\right)$$



$$\mathcal{P}_o = -b_0 L_0^{-1} = -b_0 \int_0^1 \frac{dq_o}{q_o} q_o^{L_0} = \int_0^1 \left[(-b_0) \frac{dq_o}{q_o} \right] q_o^{L_0}$$

$$\Omega_{\mathcal{P}_o}=-b_0\frac{dq_o}{q_o}$$

$$w_1 w_2 = -q_o, q_o \in [0,1]$$

$$\widehat{\Omega}_o=\mathcal{B}\left[\frac{\partial}{\partial q_o}\right]dq_o$$

$$\mathcal{B}\left[\frac{\partial}{\partial q_o}\right]=\int\;dw_1b(w_1)\frac{w_1}{q_o}+\int\;d\bar{w}_1\bar{b}(\bar{w}_1)\frac{\bar{w}_1}{q_o}=\frac{1}{q_o}\oint\;_cdw_1b(w_1)w_1$$

$$\mathcal{B}\left[\frac{\partial}{\partial q_o}\right]=-\frac{1}{q_o}b_0,\widehat{\Omega}_o=-b_0\frac{dq_o}{q_o}$$

$$\Omega_{\mathcal{P}_o}=\widehat{\Omega}_o$$

$$N_{g_1,b_1,n_{c1},n_{o1}}N_{g_2,b_2,n_{c2},n_{o2}}\sim N_{g_1+g_2,b_1+b_2-1,n_{c1}+n_{c2},n_{o1}+n_{o2}-2}$$

$$N_{g,b,n_c,n_o}\sim N_{g,b+1,n_c,n_o+2}, N_{g,b,n_c,n_o}\sim N_{g+1,b-1,n_c,n_o+2}$$

$$\begin{array}{l} \eta_c N_{g_1,b_1,n_{c1},n_{o1}}N_{g_2,b_2,n_{c2},n_{o2}}\sim N_{g_1+g_2,b_1+b_2,n_{c1}+n_{c2}-2,n_{o1}+n_{o2}}\\ \eta_c N_{g-1,b,n_c+2,n_o}\sim N_{g,b,n_c,n_o} \end{array}$$

Membranas.

$$\mathcal{O}_h=\frac{1}{g_s}h_{\mu\nu}c\bar{c}\partial X^\mu\bar{\partial}X^\nu$$

$$g_{\mu\nu}=\eta_{\mu\nu}+h_{\mu\nu}, \phi=\frac{1}{4}\eta^{\mu\nu}h_{\mu\nu}$$

$$S_p=-\mathcal{T}\int\;d^{p+1}xe^{-\phi}\sqrt{-g}$$

$$S_p\big|_h=-\mathcal{T}\Big\{\frac{1}{2}\eta^{\mu\nu}h_{\mu\nu}^{\parallel}(x_\perp=0)-\phi(x_\perp=0)\Big\}\int\;d^{p+1}x=-\frac{\mathcal{T}}{4}\eta^{\mu\nu}\Big(h_{\mu\nu}^{\parallel}-h_{\mu\nu}^{\perp}\Big)(2\pi)^{p+1}\delta^{(p+1)}(0)$$

$$\Omega_0^{(0,1,1,0)}(A)=-\frac{1}{2\pi i}\eta_c^{-1/2}\langle c_0^-A\rangle=\eta_c^{1/2}\langle c_0^-A\rangle$$

$$\Omega_0^{(0,1,1,0)}(\mathcal{O}_h)=\eta_c^{1/2}\frac{1}{g_s}h_{\mu\nu}\langle c_0^-c\bar{c}\partial X^\mu\bar{\partial}X^\nu\rangle$$

$$\Omega_0^{(0,1,1,0)}(\mathcal{O}_h)=\eta_c^{1/2}\frac{1}{g_s}\Big(h_{\mu\nu}^{\parallel}-h_{\mu\nu}^{\perp}\Big)\frac{K}{2}\eta^{\mu\nu}(2\pi)^{p+1}\delta^{(p+1)}(0)$$

$$K=-g_s\frac{\mathcal{T}}{2\sqrt{\eta}_c}$$



$$g_o^2=-\frac{2}{\mathcal{T}}\eta_c^2=\frac{1}{2\pi^2\mathcal{T}}\,\rightarrow\,\mathcal{T}g_o^2=\frac{1}{2\pi^2}$$

$$-\frac{g_s}{\pi}\int \; dy_Rdy_I V_c(y)$$

$$g_o\int \; dx V_o(x)$$

$$-g_s\frac{\mathcal{T}}{2}\langle c_0^-c\bar{c}V_c\rangle'$$

$$-i\pi \mathcal{T} g_sg_o\langle c\bar{c}V_c cV_o\rangle'$$

$$\frac{i}{2}g_s^2\mathcal{T}\int \; dy\left\langle c\bar{c}V_c^{(1)}(i)(c+\bar{c})V_c^{(2)}(iy)\right\rangle'$$

$$w_1w_2=q, |q|\leqslant 1$$

$$\mathcal{V}_{0,n}^\text{1PI} = \mathcal{V}_{0,n}$$

$$\mathcal{V}_{1,1}^\text{1PI} = \mathcal{F}_{1,1}$$

$$\mathcal{V}_\text{1PI}\equiv\sum_{g,n}\mathcal{V}_{g,n}^\text{1PI}$$

$$\partial \mathcal{V}_\text{1PI} + \frac{1}{2}\{\mathcal{V}_\text{1PI},\mathcal{V}_\text{1PI}\}=0$$

$$\{A_1,\cdots,A_n\}_\text{1PI}=\sum_{g=0}^\infty g_s^{2g+n-2}\int_{\mathcal{V}_{g,n}^\text{1PI}}\Omega_{6g-6+2n}^{(g,n)}(A_1,\cdots,A_n)$$

$$S_\text{1PI}=-\frac{1}{2}\langle\widetilde{\Psi},Q\mathcal{G}\widetilde{\Psi}\rangle+\langle\widetilde{\Psi},Q\Psi\rangle+\sum_{n=1}^\infty\frac{1}{n!}\{\Psi^n\}_\text{1PI}$$

$$\delta |\widetilde{\Psi}\rangle=Q|\widetilde{\Lambda}\rangle+\sum_{n=1}^\infty\frac{1}{n!}[\Lambda\Psi^n]_\text{1PI}, \delta |\Psi\rangle=Q|\Lambda\rangle+\sum_{n=1}^\infty\frac{1}{n!}\mathcal{G}[\Lambda\Psi^n]_\text{1PI}$$

$$\{S_\text{1PI},S_\text{1PI}\}=0$$

$${\bf Acci\'on Wilsoniana.}$$

$$\left[P,L_0^\pm\right]=0,\left[P,b_0^\pm\right]=0,\left[P,c_0^\pm\right]=0,[P,Q]=0,[P,\mathcal{G}]=0$$

$$\{A_1,\cdots,A_n\}_\text{eff}, A_1,\cdots,A_n\in P\mathcal{H}_c.$$

$$\Psi\in P\mathcal{H}_c, \tilde{\Psi}\in P\tilde{\mathcal{H}}_c$$

$$S_\text{eff}=-\frac{1}{2}\langle\widetilde{\Psi},Q\mathcal{G}\widetilde{\Psi}\rangle+\langle\widetilde{\Psi},Q\Psi\rangle+\sum_n\frac{1}{n!}\{\Psi^n\}_\text{eff}$$



$$\mathcal{V}^{\text{1PI}} = \sum_{g,n} g_s^{2g-n+2} \mathcal{V}_{g,n}^{\text{1PI}}$$

$$\partial\mathcal{V}^{\text{1PI}} + \frac{1}{2}\{\mathcal{V}^{\text{1PI}},\mathcal{V}^{\text{1PI}}\} = 0$$

$$S'_{\text{1PI}} - S_{\text{1PI}} = \sum_{g=0}^{\infty} \sum_{n=1}^{\infty} g_s^{2g-n+2} \left(\int_{\mathcal{V}'_{g,n}^{\text{1PI}}} - \int_{\mathcal{V}_{g,n}^{\text{1PI}}} \right) \Omega_{6g-6+2n}^{(g,n)}(\Psi^{\otimes n})$$

$$\langle \Phi | c_0^- | \delta \Psi \rangle = - \sum_{g=0}^{\infty} \sum_{n=1}^{\infty} g_s^{2g-n+2} \frac{1}{(n-1)!} \int_{\mathcal{V}_{g,n}^{\text{1PI}}} \Omega_{6g-5+2n}^{(g,n)}(\Phi, \Psi^{\otimes(n-1)}) [\hat{U}_{g,n}]$$

Teorema del dilatón.

$$D(z, \bar{z}) = \frac{1}{2}(c\partial^2 c - \bar{c}\partial^2 \bar{c})$$

$$|D\rangle = (c_1 c_{-1} - \bar{c}_1 \bar{c}_{-1}) |0\rangle$$

$$b_0^- |D\rangle = 0, L_0^- |D\rangle = 0$$

$$|D\rangle = Q c_0^- |0\rangle = -Q |\chi\rangle, |\chi\rangle = -c_0^- |0\rangle$$

$$L_1 |D\rangle = c_0 c_1 |0\rangle, \bar{L}_1 |D\rangle = -\bar{c}_0 \bar{c}_1 |0\rangle$$

$$\Omega_2(D) = \left(-\frac{1}{2\pi i}\right) du^1 \wedge du^2 \mathcal{B} \left[\frac{\partial}{\partial u^1}\right] \mathcal{B} \left[\frac{\partial}{\partial u^2}\right] D(w=0)$$

$$\Omega_2(D) = -d\Omega_1(\chi)$$

$$\Omega_1(\chi) = \left(-\frac{1}{2\pi i}\right) \left(du^1 \mathcal{B} \left[\frac{\partial}{\partial u^1}\right] + du^2 \mathcal{B} \left[\frac{\partial}{\partial u^2}\right]\right) \chi(w=0)$$

$$\int_M \Omega_2(D) + \int_{\partial M} \Omega_1(\chi)$$

$$\chi(M) = \frac{1}{2\pi} \int_M K^{(2)} + \frac{1}{2\pi} \int_{\partial M} k^{(1)}$$

$$\begin{aligned} K^{(2)} &= -2i\partial\bar{\partial}\rho dz \wedge d\bar{z} \\ k^{(1)} &= d\theta_\gamma - i[dz\partial\log \rho - d\bar{z}\bar{\partial}\log \rho] \end{aligned}$$

$$\Omega_2(D) = -\frac{1}{2\pi} K^{(2)} \text{ and } \Omega_1(\chi) = -\frac{1}{2\pi} k^{(1)}$$

$$[Q, \mathcal{P}_b] = \frac{1}{2\pi} b_0^- \int_0^{2\pi} d\theta \int_0^\infty ds \frac{\partial}{\partial s} e^{-s(L_0 + \bar{L}_0)} e^{i\theta(L_0 - \bar{L}_0)}$$

$$z = F(w; u) = F(w; u^1, u^2)$$



$$y(u) = F(0; u)$$

$$z = F(w; u^1, u^2) = y(u) + a(u)w + \frac{1}{2}b(u)w^2 + \frac{1}{3!}c(u)w^3 + \mathcal{O}(w^4)$$

$$\begin{aligned}\mathcal{B}\left[\frac{\partial}{\partial u^i}\right] &= \oint b(z)dz \frac{\partial F}{\partial u^i} + \oint \bar{b}(\bar{z})d\bar{z} \frac{\partial \bar{F}}{\partial u^i} \\ &= -\oint b(w)dw \left(\frac{\partial F}{\partial w}\right)^{-1} \frac{\partial F}{\partial u^i} - \oint \bar{b}(\bar{w})d\bar{w} \left(\frac{\partial \bar{F}}{\partial \bar{w}}\right)^{-1} \frac{\partial \bar{F}}{\partial u^i}\end{aligned}$$

$$\left(\frac{\partial F}{\partial w}\right)^{-1} \frac{\partial F}{\partial u^i} = \alpha_i + \beta_i w + \gamma_i w^2 + \mathcal{O}(w^3)$$

$$\begin{aligned}\alpha_i &= \frac{1}{a} \frac{\partial y}{\partial u^i} \\ \beta_i &= \frac{1}{a} \frac{\partial a}{\partial u^i} - \frac{b}{a^2} \frac{\partial y}{\partial u^i} \\ \gamma_i &= a \frac{\partial}{\partial u^i} \left(\frac{b}{2a^2} \right) + \left(\frac{b^2}{a^3} - \frac{1}{2} \frac{c}{a^2} \right) \frac{\partial y}{\partial u^i}\end{aligned}$$

$$\mathcal{B}\left[\frac{\partial}{\partial u^i}\right] = -(\alpha_i b_{-1} + \beta_i b_0 + \gamma_i b_1 + \bar{\alpha}_i \bar{b}_{-1} + \bar{\beta}_i \bar{b}_0 + \bar{\gamma}_i \bar{b}_1 + \dots)$$

$$\begin{aligned}\Omega_2(D) &= \left(-\frac{1}{2\pi i}\right) du^1 \wedge du^2 \mathcal{B}\left[\frac{\partial}{\partial u^1}\right] \mathcal{B}\left[\frac{\partial}{\partial u^2}\right] (c_1 c_{-1} - \bar{c}_1 \bar{c}_{-1}) |0\rangle \\ &= \left(\frac{1}{2\pi i}\right) du^1 \wedge du^2 [\alpha_1 \gamma_2 - \alpha_2 \gamma_1 - (\bar{\alpha}_1 \bar{\gamma}_2 - \bar{\alpha}_2 \bar{\gamma}_1)] |0\rangle\end{aligned}$$

$$\Omega_2(D) = \left(\frac{1}{2\pi i}\right) du^1 \wedge du^2 \left[\frac{\partial y}{\partial u^1} \frac{\partial}{\partial u^2} \left(\frac{b}{2a^2} \right) - \frac{\partial y}{\partial u^2} \frac{\partial}{\partial u^1} \left(\frac{b}{2a^2} \right) - (\text{c.c.}) \right] |0\rangle$$

$$\Omega_2(D) = \frac{1}{2\pi i} \left[dy \wedge d \left(\frac{b}{2a^2} \right) - d\bar{y} \wedge d \left(\frac{\bar{b}}{2\bar{a}^2} \right) \right] |0\rangle$$

$$\Omega_2(D) = \frac{1}{2\pi i} dy \wedge d\bar{y} \left[\frac{\partial}{\partial \bar{y}} \left(\frac{b}{2a^2} \right) + \frac{\partial}{\partial y} \left(\frac{\bar{b}}{2\bar{a}^2} \right) \right]$$

$$\begin{aligned}\Omega_1(\chi) &= -\frac{1}{2\pi i} \left(-\frac{1}{2} du^1 \mathcal{B}\left[\frac{\partial}{\partial u^1}\right] - \frac{1}{2} du^2 \mathcal{B}\left[\frac{\partial}{\partial u^2}\right] \right) (c_0 - \bar{c}_0) |0\rangle \\ &= -\frac{1}{2\pi i} \left(\frac{1}{2} du^1 (\beta_1 - \bar{\beta}_1) + \frac{1}{2} du^2 (\beta_2 - \bar{\beta}_2) \right) |0\rangle\end{aligned}$$

$$\Omega_1(\chi) = \frac{1}{2\pi i} \left[-\frac{1}{2} d \left(\ln \frac{a}{\bar{a}} \right) + \frac{b}{2a^2} dy - \frac{\bar{b}}{2\bar{a}^2} d\bar{y} \right]$$

$$\rho^w(w) = \rho \left| \frac{dz}{dw} \right| = \left(\rho(y) + (z-y) \frac{\partial \rho}{\partial z} \Big|_y + (\bar{z}-\bar{y}) \frac{\partial \rho}{\partial \bar{z}} \Big|_y + \dots \right) |a + bw + \dots|$$



$$\rho^w(w) = \rho \left| \frac{dz}{dw} \right| = \left(\rho(y) + aw \frac{\partial \rho}{\partial y} \Big|_y + \bar{a}\bar{w} \frac{\partial \rho}{\partial \bar{y}} \Big|_y + \dots \right) |a| \left(1 + \frac{b}{2a}w + \frac{\bar{b}}{2\bar{a}}\bar{w} + \dots \right)$$

$$= \rho(y)|a| \left[1 + aw \left(\frac{\partial}{\partial y} \log \rho + \frac{b}{2a^2} \right) + \bar{a}\bar{w} \left(\frac{\partial}{\partial \bar{y}} \log \rho + \frac{\bar{b}}{2\bar{a}^2} \right) + \dots \right]$$

$$\frac{b}{2a^2} = -\frac{\partial}{\partial y} \log \rho, \frac{\bar{b}}{2\bar{a}^2} = -\frac{\partial}{\partial \bar{y}} \log \rho$$

$$\begin{aligned} \Omega_2(D) &= \left(\frac{1}{2\pi i} \right) dy \wedge d\bar{y} \left[-2 \frac{\partial}{\partial \bar{y}} \frac{\partial}{\partial y} \log \rho \right] \\ \Omega_1(\chi) &= \frac{1}{2\pi i} (-i) \left[d\theta_a - idy \frac{\partial}{\partial y} \log \rho + id\bar{y} \frac{\partial}{\partial \bar{y}} \log \rho \right] \end{aligned}$$

$$\Omega_2(D) = -\frac{1}{2\pi} K^{(2)} \text{ and } \Omega_1(\chi) = -\frac{1}{2\pi} k^{(1)}$$

$$(2-2g-n)\Omega_{6g-6+2n}^{(g,n)}(A_1, \cdots, A_n)$$

Estructuras algebraicas de campo – álgebra de Lie y álgebra cuántica.

$$\deg(b_n(A_1, \dots, A_n)) = -1 + \sum_{i=1}^n d_{A_i}$$

$$T(V) \equiv V \oplus (V \otimes V) \oplus (V \otimes V \otimes V) \oplus \dots$$

$$\mathbf{b} = \sum_{i=1}^n \sum_{j=0}^{n-i} \mathbb{1}^{\otimes j} \otimes b_i \otimes \mathbb{1}^{n-i-j}, \text{ on } V^{\otimes n}$$

$$\mathbf{b} = \sum_{i=1}^{\infty} \mathbf{b}_i = \mathbf{b}_1 + \mathbf{b}_2 + \dots$$

$$\mathbf{b}_i = \sum_{j=0}^{n-i} \mathbb{1}^{\otimes j} \otimes b_i \otimes \mathbb{1}^{n-i-j}, \text{ on } V^{\otimes n}, \text{ for } i \leq n$$

$$\mathbf{b} = b_1 \otimes \mathbb{1} \otimes \mathbb{1} + \mathbb{1} \otimes b_1 \otimes \mathbb{1} + \mathbb{1} \otimes \mathbb{1} \otimes b_1 + b_2 \otimes \mathbb{1} + \mathbb{1} \otimes b_2 + b_3, \text{ on } V^{\otimes 3},$$

$$\begin{aligned} \mathbf{b}(A \otimes B \otimes C) &= b_1(A) \otimes B \otimes C + (-1)^{d_A} A \otimes b_1(B) \otimes C + (-1)^{d_A+d_B} A \otimes B \otimes b_1(C) \\ &\quad + b_2(A, B) \otimes C + (-1)^{d_A} A \otimes b_2(B, C) + b_3(A, B, C) \end{aligned}$$

$$\mathbf{b}^2 = 0$$

$$\sum_{i=1}^n \mathbf{b}_i \mathbf{b}_{n+1-i} = 0, \text{ on } V^{\otimes n}$$



$$\begin{aligned}V\!:\!0&=\mathbf{b}_1\mathbf{b}_1\\V^{\otimes 2}\!:\!0&=\mathbf{b}_1\mathbf{b}_2+\mathbf{b}_2\mathbf{b}_1\\V^{\otimes 3}\!:\!0&=\mathbf{b}_1\mathbf{b}_3+\mathbf{b}_2\mathbf{b}_2+\mathbf{b}_3\mathbf{b}_1\end{aligned}$$

$$V\!:\!0=b_1\circ b_1$$

$$V^{\otimes 2}\!:\!0=b_1\circ b_2+b_2\circ(b_1\otimes \mathbb{1}+\mathbb{1}\otimes b_1)$$

$$\begin{aligned}V^{\otimes 3}\!:\!0=&b_1\circ b_3+b_2\circ(b_2\otimes \mathbb{1}+\mathbb{1}\otimes b_2),\\&+b_3(b_1\otimes \mathbb{1}\otimes \mathbb{1}+\mathbb{1}\otimes b_1\otimes \mathbb{1}+\mathbb{1}\otimes \mathbb{1}\otimes b_1)\end{aligned}$$

$$0=Q^2A$$

$$0=Q(AB)+(QA)B+(-1)^{d_A}A(QB)$$

$$0=Q(A,B,C)+(AB)C+(-1)^{d_A}A(BC)+(QA,B,C)+(-1)^{d_A}(A,QB,C)+(-1)^{d_A+d_B}(A,B,QC)$$

$$(A,B)=-(-1)^{d_Ad_B}(B,A)$$

$$(QA,B)=-(-1)^{d_A}(A,QB)$$

$$\left(A_1,b_n(A_2,\cdots,A_{n+1})\right)=(-1)^{\#}\big(A_2,b_n(A_3,\cdots,A_{n+1},A_1)\big)=d_{A_2}+d_{A_1}\big(1+d_{A_2}+d_{A_3}+\cdots+d_{A_{n+1}}\big)$$

$$S(\Phi)=\sum_{n=1}^\infty \frac{1}{n+1}(\Phi,b_n(\Phi,\cdots,\Phi))$$

$$(A,B)=(-1)^{A+1}\langle A,B\rangle'$$

$$A=d_A+1 (\mathrm{mod} 2)$$

$$\langle B,A\rangle'=(-1)^{AB}\langle A,B\rangle'$$

$$(B,A)=(-1)^{B+1}\langle B,A\rangle'=(-1)^{AB+B+1}\langle A,B\rangle'=(-1)^{AB+A+B}(A,B)=(-1)^{d_Ad_B+1}(A,B).$$

$$A\star B=(-1)^{A+1}b_2(A,B)=(-1)^{A+1}(AB)$$

$$Q(A\star B)=(-1)^{A+1}Q(AB)=(-1)^A((QA)B)-(A(QB))=(QA)\star B+(-1)^AA\star QB.$$

$$[A_1\cdots A_n]=(-1)^{n(n+1)/2}\sum_\sigma ~(-1)^{s_\sigma}(-1)^{A_{\sigma(1)}+2A_{\sigma(2)}+\cdots+nA_{\sigma(n)}}b_n\big(A_{\sigma(1)},A_{\sigma(2)},\cdots A_{\sigma(n)}\big)$$

$$[A_1,A_2]=A_1\star A_2-(-1)^{A_1A_2}A_2\star A_1$$

$$[A_1\cdots A_n]=\sum_\sigma ~b_n\big(A_{\sigma(1)},A_{\sigma(2)},\cdots A_{\sigma(n)}\big)$$

$$b_1(B_1),b_2(B_1,B_2),b_3(B_1,B_2,B_3),\cdots$$

$$b_1(B_1)=QB_1, b_1(B_1,B_2)=[B_1,B_2]_0, b_3(B_1,B_2,B_3)=[B_1,B_2,B_3]_0, \cdots,$$



$$b_n(B_1,\cdots,B_k,B_{k+1},\cdots,B_n)=(-1)^{d_{B_k}d_{B_{k+1}}}b_n(B_1,\cdots,B_{k+1},B_k,\cdots,B_n)$$

$$\deg\big(b(B_1,\cdots,B_n)\big)=-1+\sum_{i=1}^n~d_{B_i}$$

$$T(W)=\sum_{n=1}^\infty SW^{\otimes n}$$

$$B_1 \wedge B_2 \wedge \cdots \wedge B_n$$

$$B_1 \wedge \cdots \wedge B_n = \epsilon(\sigma;B) B_{\sigma(1)} \wedge \cdots \wedge B_{\sigma(n)}$$

$$\mathbf{b}(B_1 \wedge \cdots \wedge B_n) = \sum_{l=1}^n \sum_{\sigma'} \epsilon(\sigma',B) \big(b_l \big(B_{i_1}, \cdots, B_{i_l} \big) \wedge B_{j_1} \wedge \cdots \wedge B_{j_{n-l}} \big),$$

$$\mathbf{b}(B_1) = b_1(B_1)$$

$$\mathbf{b}(B_1 \wedge B_2) = b_1(B_1) \wedge B_2 + (-1)^{d_{B_1} d_{B_2}} b_1(B_2) \wedge B_1 + b_1(B_1, B_2)$$

$$= b_1(B_1) \wedge B_2 + (-1)^{d_{B_1} B_1 \wedge b_1(B_2) + b_1(B_1, B_2)}$$

$$\mathbf{b}^2=0$$

$$\mathbf{b}^2(B_1) = \mathbf{b}\mathbf{b}(B_1) = \mathbf{b}(QB_1) = Q(QB_1) = 0$$

$$\mathbf{b}^2(B_1 \wedge B_2) = \mathbf{b}(QB_1 \wedge B_2) + (-1)^{d_{B_1}} \mathbf{b}(B_1 \wedge QB_2) + \mathbf{b}\big(b_2(B_1, B_2)\big) = 0$$

$$\begin{aligned} \mathbf{b}^2(B_1 \wedge B_2) &= \left((-1)^{d_{B_1}+1} QB_1 \wedge QB_2 + b_2(QB_1, B_2) \right) \\ &\quad + (-1)^{d_{B_1}} \left(QB_1 \wedge QB_2 + b_2(B_1, QB_2) \right) + QB_2(B_1, B_2) = 0 \end{aligned}$$

$$QB_2(B_1, B_2) + b_2(QB_1, B_2) + (-1)^{d_{B_1}} b_2(B_1, QB_2) = 0$$

$$0 = QB_3(B_1, B_2, B_3) + b_3(QB_1, B_2, B_3) + (-1)^{d_{B_1}} b_3(B_1, QB_2, B_3) + (-1)^{d_{B_1} + d_{B_2}} b_3(B_1, B_2, QB_3)$$

$$+ b_2(b_2(B_1, B_2), B_3) + (-1)^{d_{B_2} d_{B_3}} b_2(b_2(B_1, B_3), B_2)$$

$$+ (-1)^{d_{B_1} (d_{B_2} + d_{B_3})} b_2(b_2(B_2, B_3), B_1)$$

$$0 = \sum_{l=1}^n \sum_{\sigma'} \epsilon(\sigma',B) b_{n-l+1} \big(b_l \big(B_{i_1}, \cdots, B_{i_l} \big), B_{j_1}, \cdots, B_{j_{n-l}} \big).$$

$$(B_1, B_2) = (-1)^{(d_{B_1}+1)(d_{B_2}+1)} (B_2, B_1),$$

$$(QB_1, B_2) = (-1)^{d_{B_1}} (B_1, QB_2)$$

$$\{B_1, \cdots, B_n\}_{L_\infty} \equiv (B_1, b_{n-1}(B_2, \cdots B_n))$$

$$\{B_1, \cdots, B_k, B_{k+1}, \cdots, B_n\}_{L_\infty} = (-1)^{d_{B_k} d_{B_{k+1}}} \{B_1, \cdots, B_{k+1}, B_k, \cdots, B_n\}_{L_\infty}$$



$$S=\frac{1}{2}(\Psi,Q\Psi)+\sum_{n=3}^{\infty}\frac{1}{n!}\{\Psi^n\}_{L_\infty}$$

$$\bar{b}_1(A_1)\equiv b_1(A_1)$$

$$\bar{b}_2(A_1,A_2)\equiv b_2(A_1,A_2)+(-1)^{d_{A_1}d_{A_2}}b_2(A_2,A_1)$$

$$\bar{b}_1\bar{b}_2(A_1,A_2)+\bar{b}_2\big(\bar{b}_1(A_1),A_2\big)+(-1)^{d_{A_1}}\bar{b}_2\left(A_1,\bar{b}_1(A_2)\right)=0$$

$$b_2(b_2(A_1,A_2),A_3)+(-1)^{d_{A_1}}b_2\big(A_1,b_2(A_2,A_3)\big)=0$$

$$\bar{b}_2\big(\bar{b}_2(A_1,A_2),A_3\big)+(-1)^{d_{A_2}d_{A_3}}\bar{b}_2\big(\bar{b}_2(A_1,A_3),A_2\big)+(-1)^{d_{A_1}(d_{A_2}+d_{A_3})}\bar{b}_2\big(\bar{b}_2(A_2,A_3),A_1\big)=0$$

$$\bar{b}_n(A_1,\cdots,A_n)\equiv \sum_{\sigma}\;\epsilon(\sigma,A)b_n\big(A_{\sigma(1)},\cdots,A_{\sigma(n)}\big)$$

$$[A_1\cdots A_n]=(-1)^{d_{A_1}+2d_{A_2}+\cdots +nd_{A_n}}\bar{b}_n(A_1,\cdots,A_n)$$

$$(A_1,A_2)=\langle A_1,A_2\rangle'=(-1)^{d_{A_1}}(A_1,A_2).$$

$$\begin{aligned}\{A_0A_1\cdots A_n\}_{L_\infty}&=\bigg(A_0,\bar{b}_n(A_1A_2\cdots A_n)+=(-1)^{d_{A_0}}\big(A_0,\bar{b}_n(A_1A_2\cdots A_n)\big)\\&=(-1)^{d_{A_0}}\sum_{\sigma}\;\epsilon(\sigma,A)\big(A_0,b_n\big(A_{\sigma(1)},\cdots,A_{\sigma(n)}\big)\big)\end{aligned}$$

$$S(\psi_o)=\sum_{n=1}^{\infty}\frac{1}{n+1}(\psi_o,b_n(\psi_o^n))=\frac{1}{2}f\psi_o,Q\psi_ot+\sum_{n=2}^{\infty}\frac{1}{(n+1)!}\{\psi_o^{n+1}\}_{L_\infty}$$

$$[A_1\cdots A_n]=\bar{b}_n(A_1,\cdots,A_n)$$

$$b_n(B_1,\cdots,B_n)=\sum_{g=0}^{\infty}g_s^{2g+n-1}b_n^g(B_1,\cdots,B_n), n=0,1,\cdots$$

$$\begin{aligned}0&=\sum_{l=1}^n\sum_{\sigma'}\;\epsilon(\sigma',B)b_{n-l+1}\big(b_l\big(B_{i_1},\cdots,B_{i_l}\big),B_{j_1},\cdots,B_{j_{n-l}}\big)\\&\quad+\frac{1}{2}b_{n+2}(B_1,\cdots,B_n,\varphi_s,\varphi_r)(\varphi_s^c,\varphi_r^c)\end{aligned}$$

$$\pmb{\theta}(A_1\wedge\cdots\wedge A_n)=\frac{1}{2}A_1\wedge\cdots\wedge A_n\wedge\varphi_r\wedge\varphi_s(\varphi_r^c,\varphi_s^c)$$

$$\boldsymbol{\pi}_1(\mathbf{b}^2+\mathbf{b}\boldsymbol{\theta})=0$$

$$(\mathbf{b}+\boldsymbol{\theta})^2=0$$

$$\big\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\big\}_{L_\infty}\equiv (-1)^{d_{A_2^o}+2d_{A_3^o}+\cdots+(n-1)d_{A_n^o}}\big\{A_1^c,\cdots,A_{n_c}^c;A_1^o,\cdots,A_{n_o}^o\big\}$$



Transferencia homotópica.

$$\mathbf{M} = \mathbf{Q} + \mathbf{m}, \mathbf{m} = \sum_{n=2}^{\infty} \mathbf{b}_n$$

$$\mathbf{Q}^2 = 0, \text{ and } \mathbf{Q}\mathbf{m} + \mathbf{m}\mathbf{Q} + \mathbf{m}^2 = 0$$

$$P: V \rightarrow \bar{V}, P^2 = P$$

$$(A_1, PA_2) = (PA_1, A_2)$$

$$QP = PQ$$

$$\mathbf{P}(A_1 \otimes A_2 \otimes \cdots \otimes A_n) \equiv PA_1 \otimes PA_2 \otimes \cdots \otimes PA_n$$

$$\mathbf{P}^2 = \mathbf{P}$$

$$\mathbf{Q}\mathbf{P} = \mathbf{P}\mathbf{Q}$$

$$hP = Ph = 0, h^2 = 0, Qh + hQ = 1 - P.$$

$$h(A_1 \otimes A_2 \otimes \cdots \otimes A_n) \equiv (hA_1) \otimes PA_2 \otimes \cdots \otimes PA_n + (-1)^{d_{A_1}} A_1 \otimes hA_2 \otimes \cdots \otimes PA_n :$$

$$: + (-1)^{d_{A_1} + d_{A_2} + \cdots + d_{A_{n-1}}} A_1 \otimes A_2 \otimes \cdots \otimes hA_n$$

$$\mathbf{h}\mathbf{P} = \mathbf{P}\mathbf{h} = 0, \mathbf{h}^2 = 0, \mathbf{Q}\mathbf{h} + \mathbf{h}\mathbf{Q} = \mathbf{1} - \mathbf{P}$$

$$\overline{\mathbf{M}} = \mathbf{P}\mathbf{Q}\mathbf{P} + \mathbf{P}\mathbf{m} \frac{1}{1 + \mathbf{h} \cdot \mathbf{m}} \mathbf{P}$$

$$\frac{1}{1 + \mathbf{h} \cdot \mathbf{m}} = 1 - \mathbf{h} \cdot \mathbf{m} + \mathbf{h}\mathbf{m}\mathbf{h} \cdot \mathbf{m} - \cdots$$

$$\overline{\mathbf{M}}^2 = \mathbf{P} \left\{ \mathbf{Q}, \mathbf{m} \frac{1}{1 + \mathbf{h}\mathbf{m}} \right\} \mathbf{P} + \mathbf{P}\mathbf{m} \frac{1}{1 + \mathbf{h}\mathbf{m}} \mathbf{P}\mathbf{m} \frac{1}{1 + \mathbf{h}\mathbf{m}} \mathbf{P}$$

$$\begin{aligned} \mathbf{P} \left\{ \mathbf{Q}, \mathbf{m} \frac{1}{1 + \mathbf{h}\mathbf{m}} \right\} \mathbf{P} &= \mathbf{P}\{\mathbf{Q}, \mathbf{m}\} \frac{1}{1 + \mathbf{h}\mathbf{m}} \mathbf{P} + \mathbf{P}\mathbf{m} \frac{1}{1 + \mathbf{h}\mathbf{m}} [\mathbf{Q}, \mathbf{1} + \mathbf{h}\mathbf{m}] \frac{1}{1 + \mathbf{h}\mathbf{m}} \mathbf{P} \\ &= -\mathbf{P}\mathbf{m}^2 \frac{1}{1 + \mathbf{h}\mathbf{m}} \mathbf{P} + \mathbf{P}\mathbf{m} \frac{1}{1 + \mathbf{h}\mathbf{m}} ((1 - \mathbf{P})\mathbf{m} + \mathbf{h}^2) \frac{1}{1 + \mathbf{h}\mathbf{m}} \end{aligned}$$

$$\overline{\mathbf{Q}} = \mathbf{P}\mathbf{Q}\mathbf{P}$$

$$\overline{\mathbf{m}}_2 = \mathbf{P}\mathbf{m}_2\mathbf{P}$$

$$\overline{\mathbf{m}}_3 = \mathbf{P}[\mathbf{m}_3 - \mathbf{m}_2\mathbf{h}\mathbf{m}_2]\mathbf{P},$$

$$\overline{\mathbf{m}}_4 = \mathbf{P}[\mathbf{m}_4 - \mathbf{m}_2\mathbf{h}\mathbf{m}_3 - \mathbf{m}_3\mathbf{h}\mathbf{m}_2 + \mathbf{m}_2\mathbf{h}\mathbf{m}_2\mathbf{h}\mathbf{m}_2]\mathbf{P}.$$

$$\bar{Q}A_1 = PQ\bar{A}_1$$

$$\bar{m}_2(A_1 \otimes A_2) = P(\bar{A}_1 \star \bar{A}_2)$$

$$\bar{m}_3(A_1 \otimes A_2 \otimes A_3) = Pm_3(\bar{A}_1, \bar{A}_2, \bar{A}_3) - P(h(\bar{A}_1 \star \bar{A}_2) \star \bar{A}_3 + \bar{A}_1 \star h(\bar{A}_2 \star \bar{A}_3)).$$



$$\mathbf{L} = \mathbf{Q} + \boldsymbol{\ell}, \boldsymbol{\ell} = \sum_{n=2}^{\infty} \mathbf{b}_n$$

$$\mathbf{Q}^2=0,\text{ and }\mathbf{Q}\boldsymbol{\ell}+\boldsymbol{\ell}\mathbf{Q}+\boldsymbol{\ell}^2=0$$

$$P^2=P,PQ=QP,(B_1,PB_2)=(PB_1,B_2)$$

$$\mathbf{P}(B_1\wedge B_2\wedge\dots\wedge B_n)\equiv PB_1\wedge PB_2\wedge\dots\wedge PB_n$$

$$hP=Ph=0, h^2=0, Qh+hQ=1-P.$$

$$\begin{aligned} h(A_1 \wedge \dots \wedge A_n) &\equiv \frac{1}{n!} \sum_{\sigma \in S_n} \epsilon(\sigma; A) \left(h(A_{\sigma(1)} \wedge \dots \wedge A_{\sigma(n-1)}) \wedge A_{\sigma(n)} \right. \\ &\quad \left. + (-1)^{A_{\sigma(1)} + \dots + A_{\sigma(n-1)}} \bar{A}_{\sigma(1)} \wedge \dots \wedge \bar{A}_{\sigma(n-1)} \wedge hA_{\sigma(n)} \right) \end{aligned}$$

$$h(A_1 \wedge A_2) = \frac{1}{2} \left(hA_1 \wedge A_2 + hA_1 \wedge \bar{A}_2 + (-1)^{A_1} (\bar{A}_1 \wedge hA_2 + A_1 \wedge hA_2) \right)$$

$$\mathbf{h}\mathbf{P}=\mathbf{P}\mathbf{h}=0, \mathbf{h}^2=0, \mathbf{Q}\mathbf{h}+\mathbf{h}\mathbf{Q}=\mathbf{1}-\mathbf{P}$$

$$\overline{\mathbf{L}} = \mathbf{P} \mathbf{Q} \mathbf{P} + \mathbf{P} \boldsymbol{\ell} \frac{1}{1 + \mathbf{h} \boldsymbol{\ell}} \mathbf{P}$$

$$\begin{aligned} \bar{Q}A_1 &= Q\bar{A}_1\bar{\ell}_2(A_1 \wedge A_2) = P[\bar{A}_1, \bar{A}_2]\bar{\ell}_3(A_1 \wedge A_2 \wedge A_3) \\ &= P\ell_3(\bar{A}_1 \wedge \bar{A}_2 \wedge \bar{A}_3) \\ &\quad - P([h[\bar{A}_1, \bar{A}_2], \bar{A}_3] + (-1)^{A_2 A_3}[h[\bar{A}_1, \bar{A}_3], \bar{A}_2] + [\bar{A}_1, h[\bar{A}_2, \bar{A}_3]]) \end{aligned}$$

Vórtices – espacios de Moduli.

$$\Delta(\Delta X) = \Delta^2 X = 0, \forall X \in \mathcal{C}$$

$$\begin{aligned} \Delta(XYZ) &= \Delta(XY)Z + (-1)^X X\Delta(YZ) + (-1)^{XY+Y} Y\Delta(XZ) \\ &\quad - \Delta X(YZ) - (-1)^X X(\Delta Y)Z - (-1)^{X+Y} XY\Delta Z \end{aligned}$$

$$\{X, Y\} \equiv (-1)^X \Delta(XY) - (-1)^X (\Delta X)Y - X\Delta Y$$

$$\begin{aligned} \{X, Y\} &= -(-1)^{(X+1)(Y+1)} \{Y, X\} \\ 0 &= (-1)^{(X+1)(Z+1)} \{\{X, Y\}, Z\} + \mathfrak{I}_{\text{cyclic}} \\ \{X, YZ\} &= \{X, Y\}Z + (-1)^{XY+Y} Y\{X, Z\} \end{aligned}$$

$$\mathcal{A}_{g_1,n_1} \times \dots \times \mathcal{A}_{g_r,n_r}$$

$$\deg(\mathcal{A}_{g,n}) = \dim(\mathcal{A}_{g,n})(\bmod 2)$$

$$\mathcal{A}_{g_1,n_1} \times \mathcal{A}_{g_2,n_2} = (-1)^{\mathcal{A}_{g_1,n_1} \cdot \mathcal{A}_{g_2,n_2}} \mathcal{A}_{g_2,n_2} \times \mathcal{A}_{g_1,n_1}$$

$$\left[[\mathcal{A}_{g_1,n_1}, \dots, \mathcal{A}_{g_r,n_r}] \right] \equiv \frac{1}{n_1! \cdots n_r!} \sum_{\sigma \in S_N} \mathbf{P}_\sigma (\mathcal{A}_{g_1,n_1} \times \dots \times \mathcal{A}_{g_r,n_r}) \in \mathcal{C}$$



$$\left[\left[\mathcal{A}_{g_1,n_1},\dots,\mathcal{A}_{g_r,n_r}\right]\right]\cdot\left[\left[\mathcal{B}_{g'_1,n'_1},\dots,\mathcal{B}_{g'_p,n'_p}\right]\right]=\left[\left[\mathcal{A}_{g_1,n_1},\dots,\mathcal{A}_{g_r,n_r},\mathcal{B}_{g'_1,n'_1},\dots,\mathcal{B}_{g'_p,n'_p}\right]\right]$$

$$\Delta X \equiv \Delta_{ij}X \equiv \frac{1}{n_1!\cdots n_r!}\sum_{\sigma \in S_N}\Delta_{ij}\mathbf{P}_\sigma(\mathcal{A}_{n_1,g_1}\times\cdots\times\mathcal{A}_{g_r,n_r})$$

$$\Delta_{12}(\Delta_{12}\Sigma_X)\left[\frac{\partial}{\partial\theta_{34}},\frac{\partial}{\partial\theta_{12}},\{X\}\right]$$

$$\Delta_{12}(\Delta_{34}\Sigma_X)\left[\frac{\partial}{\partial\theta_{12}},\frac{\partial}{\partial\theta_{34}},\{X\}\right]$$

$$\{X,Y\}\equiv (-1)^X\Delta(XY)-(-1)^X(\Delta X)Y-X\Delta Y$$

$$(-1)^X\Delta(XY)=S_{XX}+S_{YY}+S_{XY}$$

$$(-1)^X(\Delta X)Y=S_{XX}, X\Delta Y=S_{YY}$$

$$\partial(XY)=(\partial X)Y+(-1)^X X\partial Y$$

$$\Delta\partial X=-\partial\Delta X$$

$$\partial\{X,Y\}=\{\partial X,Y\}+(-1)^{X+1}\{X,\partial Y\}$$

$$\textbf{Estructuras de campo en CFT}.$$

$$\Bigl\langle \Omega_k^{(g,n)} \Bigr| \in (\mathcal{H}^*)^{\otimes n}$$

$$\Bigl\langle \Omega_k^{(g,n)} \mid A_1 \Bigr\rangle \otimes \cdots \otimes |A_n\rangle = \Omega_k^{(g,n)}(A_1,\cdots,A_n),$$

$$\int_{\mathcal{A}_{g,n}}\Bigl\langle \Omega_k^{(g,n)} \Bigl|\sum_{i=1}^n Q^{(i)}=(-1)^k\int_{\partial\mathcal{A}_{g,n}}\Bigl\langle \Omega_{k-1}^{(g,n)} \Bigr|$$

$$f\left(\mathcal{A}_{g,n}^{(k)}\right)\equiv\frac{1}{n!}\int_{\mathcal{A}_{g,n}^{(k)}}\Bigl\langle \Omega_k^{(g,n)} \mid \Psi \Bigr\rangle_1\cdots |\Psi\rangle_n$$

$$f\left(\left[\left[\mathcal{A}_{g_1,n_1}^{(k_1)}\cdots\mathcal{A}_{g_r,n_r}^{(k_r)}\right]\right]\right)\equiv\prod_{i=1}^r\frac{1}{n_i!}\int_{\mathcal{A}_{g_in_i}^{(k_i)}}\Bigl\langle \Omega_{k_i}^{(g_in_i)} \mid \Psi \Bigr\rangle_1\cdots |\Psi\rangle_{n_i}$$

$$f(XY)=f(X)f(Y), X,Y\in\mathcal{C}$$

$${}_1\langle A| {}_2\langle B\mid R_{12}\rangle=\langle A\mid B\rangle, \langle R_{12}\mid A\rangle_1|B\rangle_2=\langle A\mid B\rangle,$$

$$\begin{aligned}\langle\omega_{12}|&=\langle R_{12}|c_0^{-(2)}\\|S_{12}\rangle&=b_0^{-(1)}|R_{12}\rangle\end{aligned}$$

$$\{F,G\}=(-1)^{G+1}\frac{\partial F}{\partial|\Psi\rangle}\frac{\partial G}{\partial|\Psi\rangle}|S\rangle$$



$$S_{0,2}=\frac{1}{2}\langle\omega_{12}|Q^{(2)}|\Psi\rangle_1|\Psi\rangle_2$$

$$\Delta F = \frac{1}{2}(-1)^{F+1}\left(\frac{\partial}{\partial |\Psi\rangle}\frac{\partial}{\partial |\Psi\rangle}F\right)|S\rangle$$

$$f(\Delta X)=-\Delta f(X)$$

$$f(\{X,Y\})=-\{f(X),f(Y)\},$$

$$\{S_{0,2},f(X)\}=-f(\partial X)$$

$$\mathcal{V}=\sum_{g,n} g_s^{2g+n-2}\mathcal{V}_{g,n}, \text{ with } \begin{cases} n\geqslant 3, \text{ for } g=0 \\ n\geqslant 1, \text{ for } g=1 \\ n\geqslant 0, \text{ for } g\geqslant 2 \end{cases}$$

$$\partial \mathcal{V} + \Delta \mathcal{V} + \frac{1}{2}\{\mathcal{V},\mathcal{V}\}=0$$

$$S=S_{0,2}+f(\mathcal{V})$$

$$\{S_{0,2},S_{0,2}\}=0,\Delta S_{0,2}=0$$

$$0=\frac{1}{2}\{S,S\}+\Delta S=\{S_{0,2},f(\mathcal{V})\}+\frac{1}{2}\{f(\mathcal{V}),f(\mathcal{V})\}+\Delta f(\mathcal{V})$$

$$=-f(\partial \mathcal{V})-\frac{1}{2}f(\{\mathcal{V},\mathcal{V}\})-f(\Delta \mathcal{V})-f\left(\partial \mathcal{V}+\frac{1}{2}\{\mathcal{V},\mathcal{V}\}+\Delta \mathcal{V}\right)=0$$

$$\partial \begin{array}{c} \text{circle} \\ \mathcal{V}_{g,n} \end{array} = -\frac{1}{2} \sum_{\substack{g_1+g_2=g \\ n_1+n_2=n+2}} \begin{array}{c} \text{circle} \quad \text{circle} \\ \mathcal{V}_{g_1,n_1} \quad \mathcal{V}_{g_2,n_2} \end{array} -\frac{1}{2} \begin{array}{c} \text{circle} \\ \text{circle} \\ \mathcal{V}_{g-1,n+2} \end{array}$$

$$\partial \mathcal{V}_{g,n} = -\Delta \mathcal{V}_{g-1,n+2} - \frac{1}{2} \sum_{\substack{g_1+g_2=g \\ n_1+n_2=n+2}} \{\mathcal{V}_{g_1,n_1},\mathcal{V}_{g_2,n_2}\} \equiv \mathcal{O}_{g,n}$$

$$\partial \mathcal{V}_{g,n} = -\partial_p R_1,$$



Compactificación de Deligne-Mumford, nodos de superficie y diagramas de Feynman.

Transformaciones canonicales.

$$\mathcal{F}_{g,n} = \mathcal{V}_{g,n} \oplus R_1 \oplus \cdots R_{3g-3+n}$$

$$\partial\mathcal{V}_{0,4} = -\frac{1}{2}\{\mathcal{V}_{0,3}, \mathcal{V}_{0,3}\}, \text{ and } \partial\mathcal{V}_{1,1} = -\Delta\mathcal{V}_{0,3}.$$

$$\partial\mathcal{V}_{g,n} = \mathcal{O}_{g,n}$$

$$[\mathcal{O}_{g,n}] \in H_{6g-6+2n-1}(\hat{\mathcal{P}}_{g,n})^{S_n}$$

$$\delta S = \Delta\epsilon + \{S, \epsilon\}$$

$$\delta_{\mathcal{W}}\mathcal{V} = \{\mathcal{V}, \mathcal{W}\} + \Delta\mathcal{W} + \partial\mathcal{W},$$

$$\mathcal{V} + \delta_{\mathcal{W}}\mathcal{V}$$

$$\begin{aligned} & \frac{1}{2}\{\mathcal{V} + \delta_{\mathcal{W}}\mathcal{V}, \mathcal{V} + \delta_{\mathcal{W}}\mathcal{V}\} + \Delta(\mathcal{V} + \delta_{\mathcal{W}}\mathcal{V}) + \partial(\mathcal{V} + \delta_{\mathcal{W}}\mathcal{V}) \\ &= \left\{ \partial\mathcal{V} + \Delta\mathcal{V} + \frac{1}{2}\{\mathcal{V}, \mathcal{V}\}, \mathcal{W} \right\} + O((\delta_{\mathcal{W}}\mathcal{V})^2) = 0 \end{aligned}$$

$$S_{\mathcal{V} + \delta_{\mathcal{W}}\mathcal{V}} = S_{\mathcal{V}} + f(\delta_{\mathcal{W}}\mathcal{V}) = S_{\mathcal{V}} + f(\{\mathcal{V}, \mathcal{W}\} + \Delta\mathcal{W} + \partial\mathcal{W})$$

$$= S_{\mathcal{V}} - \{f(\mathcal{V}), f(\mathcal{W})\} - \Delta f(\mathcal{W}) - \{S_{0,2}, \mathcal{W}\} = S_{\mathcal{V}} - \{S_{\mathcal{V}}, f(\mathcal{W})\} - \Delta f(\mathcal{W})$$

$$\frac{d}{dt}\mathcal{V}(t) = \delta_{\mathcal{W}}\mathcal{V}(t)$$

$$M_{\mathcal{V}}(t) \equiv \partial\mathcal{V}(t) + \Delta\mathcal{V}(t) + \frac{1}{2}\{\mathcal{V}(t), \mathcal{V}(t)\}$$

$$\frac{dM_{\mathcal{V}}}{dt} = \{M_{\mathcal{V}}(t), \mathcal{W}\}$$

$$\mathcal{V}(t) = \mathcal{V} + t\delta_{\mathcal{W}}\mathcal{V} + \frac{1}{2}t^2\{\delta_{\mathcal{W}}\mathcal{V}, \mathcal{W}\} + \frac{1}{3!}t^3\{\{\delta_{\mathcal{W}}\mathcal{V}, \mathcal{W}\}, \mathcal{W}\} + \dots$$

$$\exp(\delta_{\mathcal{W}})\mathcal{V} \equiv \mathcal{V} + \delta_{\mathcal{W}}\mathcal{V} + \frac{1}{2}\{\delta_{\mathcal{W}}\mathcal{V}, \mathcal{W}\} + \frac{1}{3!}\{\{\delta_{\mathcal{W}}\mathcal{V}, \mathcal{W}\}, \mathcal{W}\} + \dots$$

$$\exp(\delta_{\mathcal{W}})\mathcal{V} = \mathcal{V}'$$



$$\partial \begin{array}{c} g, b \\ n_c, n_o \end{array} = - \begin{array}{c} \text{wavy line} \\ \text{circle} \\ g-1, b \\ n_c+2, n_o \end{array} - \begin{array}{c} \text{circle} \\ \text{dot} \\ g, b-1 \\ n_c, n_o + 2 \end{array} - \begin{array}{c} \text{circle} \\ \text{double loop} \\ g-1, b+1 \\ n_c, n_o + 2 \end{array}$$

$$-\frac{1}{2} \begin{array}{c} \text{circle} \\ g_1, b_1 \\ n_{c1}, n_{o1} \end{array} \begin{array}{c} \text{wavy line} \\ \text{circle} \\ g_2, b_2 \\ n_{c2}, n_{o2} \end{array}$$

$$-\frac{1}{2} \begin{array}{c} \text{circle} \\ g_1, b_1 \\ n_{c1}, n_{o1} \end{array} \begin{array}{c} \text{dot} \\ \text{circle} \\ g_2, b_2 \\ n_{c2}, n_{o2} \end{array}$$

$$(b_1 + b_2 = b + 1)$$

$$\mathcal{V} = \sum_{g,b,n_c,n_o} (g_s)^{-\chi_{g,b,n_c,n_o}} \mathcal{V}_{g,b,n_c,n_o},$$

$$\partial \mathcal{V} + \Delta \mathcal{V} + \frac{1}{2} \{\mathcal{V}, \mathcal{V}\} = 0$$

$$\partial \mathcal{V}_c + \Delta \mathcal{V}_c + \frac{1}{2} \{\mathcal{V}_c, \mathcal{V}_c\}_c = 0$$

$$\partial \hat{\mathcal{V}}_c + \frac{1}{2} \{\hat{\mathcal{V}}_c, \hat{\mathcal{V}}_c\}_c = 0$$

$$\partial \mathcal{V}_o + \frac{1}{2} \{\mathcal{V}_o, \mathcal{V}_o\}_o = 0$$

$$\partial \mathcal{V}_1 + \{\mathcal{V}_o, \mathcal{V}_1\}_o = 0$$

$$\partial \mathcal{V}_2 + \{\mathcal{V}_o, \mathcal{V}_2\}_o + \frac{1}{2} \{\mathcal{V}_1, \mathcal{V}_1\}_o + \{\mathcal{V}_1, \mathcal{V}_{0,3}\}_c = 0$$

$$\overline{\mathcal{V}} \equiv \mathcal{V}_o + \sum_{n=1}^{\infty} \mathcal{V}_n, \mathcal{V}_n \equiv \sum_{n_o=0}^{\infty} (g_s)^{n+\frac{1}{2}n_o-1} \mathcal{V}_{0,1,n,n_o}$$



$$\partial \overline{\mathcal{V}} + \frac{1}{2} \{ \overline{\mathcal{V}}, \overline{\mathcal{V}} \}_o + \{ \overline{\mathcal{V}}, \hat{\mathcal{V}}_c \}_c = 0$$

$$\tilde{\mathcal{V}}=\sum_{n_c,b,n_o}(g_s)^{n_c+b+\frac{1}{2}n_o-2}\mathcal{V}_{0,b,n_c,n_o}$$

$$\partial \tilde{\mathcal{V}} + \frac{1}{2} \{ \tilde{\mathcal{V}}, \tilde{\mathcal{V}} \} + \Delta'_o \tilde{\mathcal{V}} = 0$$

Vórtice de Witten.

$$\begin{aligned} w_1 w_2 &= -1, & \text{for } |w_1| = 1, \operatorname{Re} w_1 \leq 0 \\ w_2 w_3 &= -1, & \text{for } |w_2| = 1, \operatorname{Re} w_2 \leq 0 \\ w_3 w_1 &= -1, & \text{for } |w_3| = 1, \operatorname{Re} w_3 \leq 0 \end{aligned}$$

$$\varphi = \phi(w_i) dw_i^2 = -\frac{1}{w_i^2} dw_i^2$$

$$ds^2 = |\phi(w_i)| |dw_i|^2$$

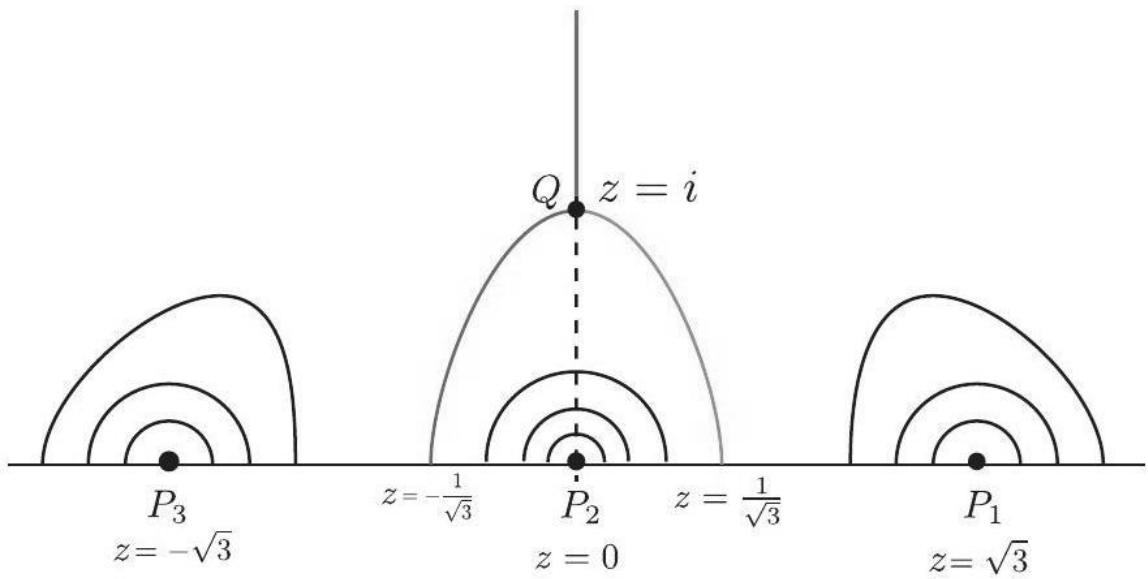
$$h(u)=\frac{1+iu}{1-iu}$$

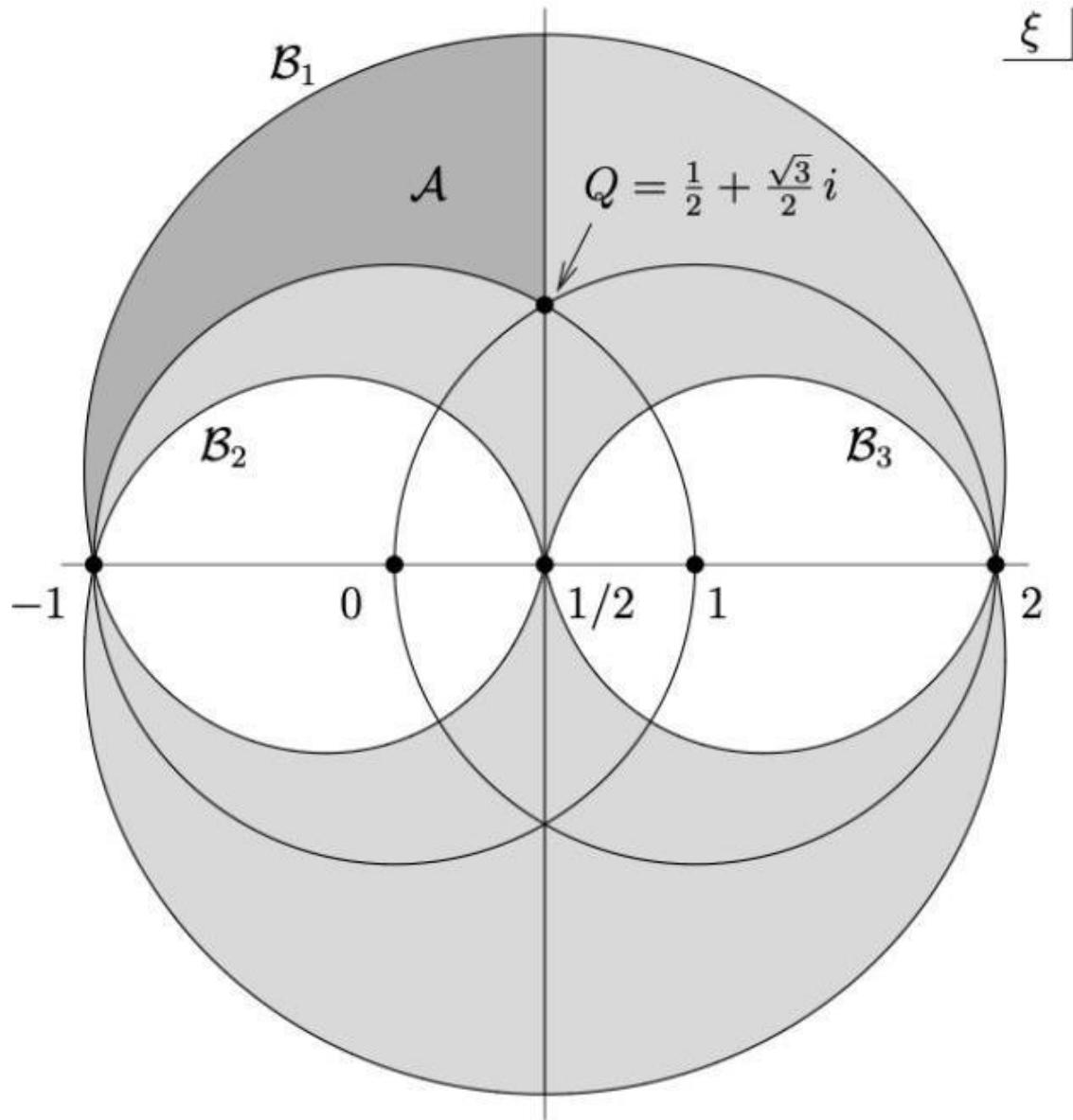
$$w(w_1) = e^{2\pi i/3} (h(w_1))^{2/3}, w(w_2) = (h(w_2))^{2/3}, w(w_3) = e^{-2\pi i/3} (h(w_3))^{2/3},$$

$$z = f_1(w_1) = h^{-1}(w(w_1)) = \sqrt{3} + \frac{8}{3}w_1 + \frac{16}{9}\sqrt{3}w_1^2 + \frac{248}{81}w_1^3 + \mathcal{O}(w_1^4)$$

$$\begin{aligned} z = f_2(w_2) &= h^{-1}(w(w_2)) = \frac{2}{3}w_2 - \frac{10}{81}w_2^3 + \mathcal{O}(w_2^5) \\ z = f_3(w_3) &= h^{-1}(w(w_3)) = -\sqrt{3} + \frac{8}{3}w_3 - \frac{16}{9}\sqrt{3}w_3^2 + \frac{248}{81}w_3^3 + \mathcal{O}(w_3^4) \end{aligned}$$

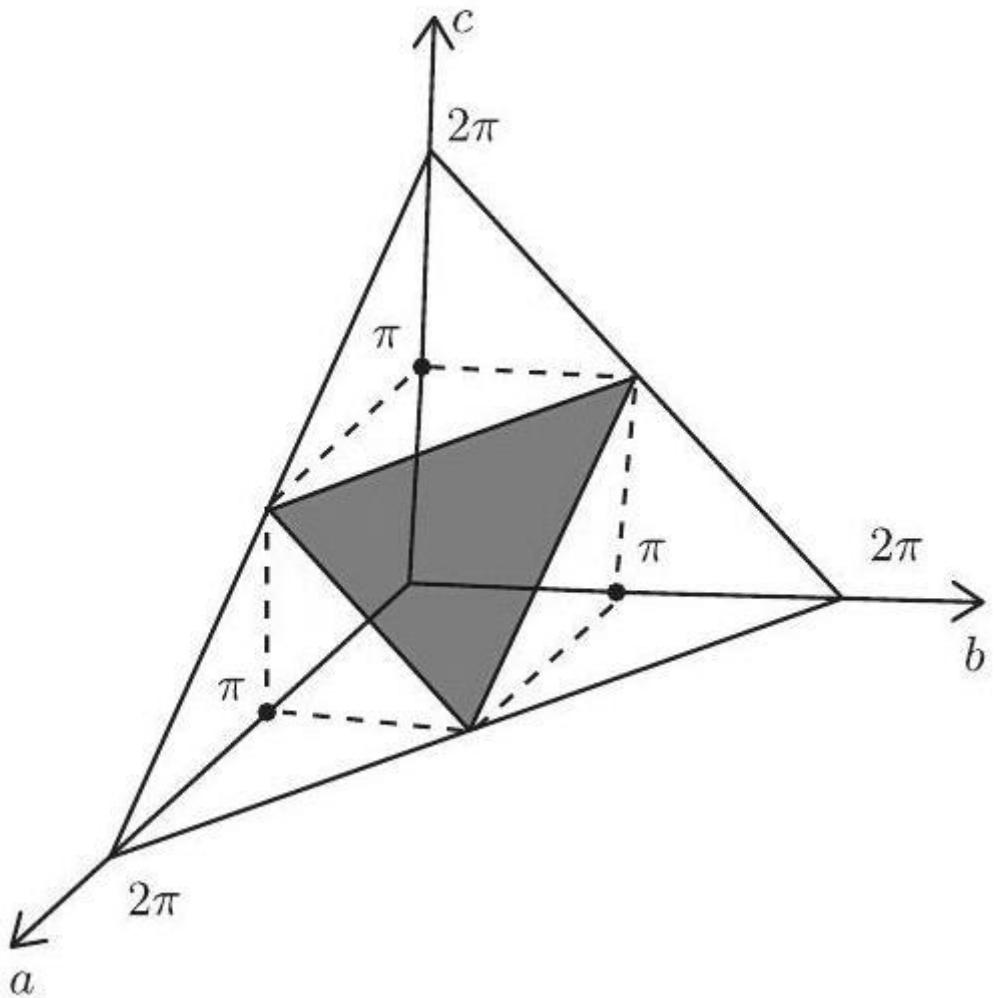






$$a, b, c \leq \pi, a', b', c' \leq \pi$$





Métricas hiperbólicas y curvatura de Gauss – Riemann – Espacios de Hilbert – Einstein deformados.

$$\text{gr}_\infty: \Sigma_{g,n} \rightarrow \hat{\mathcal{P}}_{g,n}$$

$$\pi \circ \text{gr}_\infty: \mathcal{M}_{g,n,L} \rightarrow \mathcal{M}_{g,n}$$

$$\tilde{\mathcal{V}}_{g,n}(L) \equiv \{\Sigma \in \mathcal{M}_{g,n,L} \mid \text{sys } \Sigma \geq L\}$$

$$\mathcal{V}_{g,n}(L) \equiv \text{gr}_\infty(\tilde{\mathcal{V}}_{g,n}(L))$$

$$\sinh\left(\frac{1}{2}w_i\right)\sinh\left(\frac{1}{2}\ell_i\right)=1$$

$$L \leq L_* \rightarrow w \geq L$$

$$L_c \leq L'_*, \text{ and } \sinh L_c \sinh L_o \leq 1$$

$$\mathbf{B}^{[0]}(t) = \sum_{n=0}^{\infty} t^n \mathbf{b}_{n+1}^{(n)}$$

$$\left(\mathbf{B}^{[0]}(t)\right)^2=0$$

$$\{\boldsymbol{\eta}_0,\mathbf{B}^{[0]}(t)\}=0$$

$$\mathbf{B}^{[m]}(t)=\sum_{n=0}^{\infty}\;t^n\mathbf{b}_{m+n+1}^{(n)}, m=0,1,\ldots$$

$$\mathbf{B}(s,t)=\sum_{m=0}^{\infty}s^m\mathbf{B}^{[m]}(t)=\sum_{m,n=0}^{\infty}s^mt^n\mathbf{b}_{m+n+1}^{(n)}$$

$$\boldsymbol{\mu}(s,t)=\sum_{m=0}^{\infty}s^m\boldsymbol{\mu}^{[m]}(t)=\sum_{m,n=0}^{\infty}s^mt^n\boldsymbol{\mu}_{m+n+2}^{(n+1)}$$

$$\{\mathbf{B}(s,t),\mathbf{B}(s,t)\}=0, \text{ and } \{\boldsymbol{\eta}_0,\mathbf{B}(s,t)\}=0$$

$$\frac{\partial}{\partial t}\mathbf{B}(s,t)=[\mathbf{B}(s,t),\boldsymbol{\mu}(s,t)],\frac{\partial}{\partial s}\mathbf{B}(s,t)=[\boldsymbol{\eta}_0,\boldsymbol{\mu}(s,t)].$$

$$\mathbf{B}(s,t)=\mathbf{Q}+t\mathbf{b}_2^{(1)}+s\mathbf{b}_2^{(0)}+\cdots,\boldsymbol{\mu}(s,t)=\boldsymbol{\mu}_2^{(1)}+\cdots$$

$$\mathbf{b}_2^{(1)}=\left[\mathbf{Q},\boldsymbol{\mu}_2^{(1)}\right],\mathbf{b}_2^{(0)}=\left[\boldsymbol{\eta}_0,\boldsymbol{\mu}_2^{(1)}\right].$$

$$\boldsymbol{\mu}_2^{(1)}=\frac{1}{3}\Big(\boldsymbol{\xi}_0\mathbf{b}_2^{(0)}-\mathbf{b}_2^{(0)}\boldsymbol{\xi}_0\Big).$$

$$\boldsymbol{\mu}_2^{(1)}(A\otimes B)=\frac{1}{3}\big(\boldsymbol{\xi}_0(A,B)-(\boldsymbol{\xi}_0 A,B)-(-1)^A(A,\boldsymbol{\xi}_0 B)\big).$$

$$\begin{aligned} \left[\boldsymbol{\eta}_0,\boldsymbol{\mu}_2^{(1)}\right](A\otimes B)&=\boldsymbol{\eta}_0\boldsymbol{\mu}_2^{(1)}(A\otimes B)=\frac{1}{3}\boldsymbol{\eta}_0\big(\boldsymbol{\xi}_0(A,B)-(\boldsymbol{\xi}_0 A,B)-(-1)^A(A,\boldsymbol{\xi}_0 B)\big)\\ &=\frac{1}{3}\big(\boldsymbol{\eta}_0\boldsymbol{\xi}_0(A,B)+(\boldsymbol{\eta}_0\boldsymbol{\xi}_0 A,B)+(A,\boldsymbol{\eta}_0\boldsymbol{\xi}_0 B)\big)=\frac{1}{3}\big((A,B)+(A,B)+(A,B)\big)\\ &=(A,B)=\mathbf{b}_2^{(0)}(A\otimes B) \end{aligned}$$

$$\begin{aligned} (A,B)^*\equiv\mathbf{b}_2^{(1)}(A\otimes B)&=\left[\mathbf{Q},\boldsymbol{\mu}_2^{(1)}\right](A\otimes B)\\ &=\mathbf{Q}\boldsymbol{\mu}_2^{(1)}(A\otimes B)-\boldsymbol{\mu}_2^{(1)}(QA\otimes B)-(-1)^A\boldsymbol{\mu}_2^{(1)}(A\otimes QB)\\ &=\frac{1}{3}Q\big(\boldsymbol{\xi}_0(A,B)-(\boldsymbol{\xi}_0 A,B)-(-1)^A(A,\boldsymbol{\xi}_0 B)\big)\\ &\quad-\frac{1}{3}\big(\boldsymbol{\xi}_0(QA,B)-(\boldsymbol{\xi}_0 QA,B)+(-1)^A(QA,\boldsymbol{\xi}_0 B)\big)\\ &\quad-\frac{1}{3}(-1)^A\big(\boldsymbol{\xi}_0(A,QB)-(\boldsymbol{\xi}_0 A,QB)-(-1)^A(A,\boldsymbol{\xi}_0 QB)\big) \end{aligned}$$



$$(A,B)^*\equiv \frac{1}{3}\left(\mathcal{X}_0(A,B)+(\mathcal{X}_0A,B)+(A,\mathcal{X}_0B)\right)$$

$$S(\Phi)=\frac{1}{2}\langle\Phi,Q\Phi\rangle'+\sum_{n=1}^{\infty}\frac{1}{n+2}\Big\langle\Phi,b_{n+1}^{(n)}(\Phi,\cdots,\Phi)\Big\rangle'$$

$$[A,B]^*\equiv \frac{1}{3}\left(\mathcal{X}_0[A,B]+[\mathcal{X}_0A,B]+[A,\mathcal{X}_0B]\right)$$

Métrica de Berkovits para espacios de Hilbert – Einstein deformados.

$$\{\eta_0,\xi_0\}=1$$

$$\{Q,\eta_0\}=0$$

$$S=\frac{1}{2g^2}\left|\left\langle(e^{-\Phi}\eta_0e^\Phi)(e^{-\Phi}Qe^\Phi)+\int_0^1dt\Phi\{e^{-t\Phi}Qe^{t\Phi},e^{-t\Phi}\eta_0e^{t\Phi}\}\right\rangle\right|$$

$$\left\langle\langle A_1\dots A_n\rangle\right\rangle=\left\langle h^{-1}\circ f_1^{(n)}\circ A_1(0)\dots h^{-1}\circ f_n^{(n)}\circ A_n(0)\right\rangle'_L, n\geqslant 2,$$

$$f_1^{(n)}(z)=\left(\frac{1+iz}{1-iz}\right)^{2/n}, f_k^{(n)}(z)=e^{2\pi i(k-1)/n}f_1^{(n)}(z), k=2,\dots n$$

$$\left\langle\xi c\partial c\partial^2ce^{-2\phi}\right\rangle'_L\neq 0$$

$$\left\langle\langle A_1\dots A_n\rangle\right\rangle=\langle A_1,A_2\star A_3\star\dots\star A_n\rangle'_L$$

$$\eta_0(e^{-\Phi}Qe^\Phi)=0$$

$$\delta e^\Phi=(Q\Lambda)e^\Phi+e^\Phi(\eta_0\Omega)$$

$$S=\frac{1}{2g^2}\left|\left\langle\frac{1}{2}(Q\Phi)(\eta_0\Phi)+\frac{1}{6}(Q\Phi)(\Phi(\eta_0\Phi)-(\eta_0\Phi)\Phi)\right\rangle\right|+\mathcal{O}(\Phi^3)$$

$$S_{\text{kin}}\sim\left\langle\langle(Q\Phi)(\eta_0\Phi)\rangle\right\rangle=\langle Q\Phi,\eta_0\Phi\rangle'_L,$$

$$Q\eta_0|\Phi\rangle=0$$

$$\delta\Phi=Q|\Lambda\rangle, \delta\Phi=\eta_0|\Omega\rangle$$

$$|\Phi\rangle=|\phi'\rangle+\xi_0|\phi\rangle, \text{ with } \eta_0|\phi'\rangle=\eta_0|\phi\rangle=0$$

$$\delta|\Phi\rangle=\eta_0|\Omega\rangle=\eta_0(|\omega\rangle+\xi_0|\omega'\rangle)=|\omega'\rangle$$

$$0=Q\eta_0|\Phi\rangle=Q\eta_0\xi_0|\phi\rangle=Q\{\eta_0,\xi_0\}|\phi\rangle=Q|\phi\rangle$$



Sector holomórfico.

$$\left\langle \xi c\partial c\partial^2c\bar{c}\partial\bar{c}\partial^2\bar{c}e^{-2\phi}\right\rangle _L\neq0$$

$$S=2\left[\frac{1}{2}\langle\eta_0V,QV\rangle_L+\frac{\kappa}{3!}\langle\eta_0V,[V,QV]_0\rangle_L\right]+\mathcal{O}(V^3)$$

$$\delta V=Q\Lambda+\eta_0\Omega$$

$$G(\tau V+d\tau V)=G(\tau V)+Q(d\tau V)+\sum_{n=1}^\infty\frac{\kappa^n}{n!}[G(\tau V)^n,d\tau V]_0+\mathcal{O}(d\tau^2)$$

$$\partial_\tau G(\tau V)=QV+\sum_{n=1}^\infty\frac{\kappa^n}{n!}[G(\tau V)^n,V]_0$$

$$\partial_\tau G=QV+\kappa[G,V]_0+\frac{\kappa^2}{2}[G,G,V]_0+\mathcal{O}(\kappa^3)$$

$$G=G^{(0)}+\kappa G^{(1)}+\kappa^2 G^{(2)}+\mathcal{O}(\kappa^3)$$

$$G(V)=QV+\frac{\kappa}{2}[V,QV]_0+\frac{\kappa^2}{3!}([V,QV,QV]_0+[V,[V,QV]_0]_0)+\mathcal{O}(\kappa^3)$$

$$S=2\int_0^1dt\langle\eta_0V,G(tV)\rangle_L$$

$$\left\langle \xi\bar{\xi}c\partial c\partial^2c\bar{c}\partial\bar{c}\partial^2\bar{c}e^{-2\phi}e^{-2\bar{\phi}}\right\rangle _L\neq0$$

$$S_2=\frac{1}{2}\langle\eta_0\Psi,Q\bar{\eta}_0\Psi\rangle_L$$

$$\delta\Psi=Q\Lambda+\eta_0\Omega+\bar{\eta}_0\bar{\Omega}$$

$$S_3=\frac{1}{3!}\langle\eta_0\Psi,[Q\bar{\eta}_0\Psi,\bar{\eta}_0\Psi]^*\rangle_L$$

$$[A,B]^*\equiv\frac{1}{3}\Big(\overline{\mathcal{X}}_0[A,B]_0+\big[\overline{\mathcal{X}}_0A,B\big]_0+\big[A,\overline{\mathcal{X}}_0B\big]_0\Big)$$

$$S=\int_0^1dt\langle\eta_0\Psi_t,\mathcal{G}^*(\Psi(t))\rangle_L$$

$$\Psi_t=\Psi+\frac{\kappa}{2}t[\bar{\eta}_0\Psi,\Psi]^*+\mathcal{O}(\kappa^2)$$

$$\mathcal{G}^*(\Psi)=Q\bar{\eta}_0\Psi+\frac{\kappa}{2}t[Q\bar{\eta}_0\Psi,\bar{\eta}_0\Psi]^*+\mathcal{O}(\kappa^2)$$



Sector de Ramond.

$$\beta(z) = \sum_{n \in \mathbb{Z}} \beta_n z^{-n-\frac{3}{2}}, \gamma(z) = \sum_{n \in \mathbb{Z}} \gamma_n z^{-n+\frac{1}{2}}$$

$$\beta_n | -1/2 \rangle = 0 \text{ for } n \geq 0, \gamma_n | -1/2 \rangle = 0 \text{ for } n \geq 1$$

$$X = -\delta(\beta_0)G_0^m + \delta'(\beta_0)b_0, Y = -c_0\delta'(\gamma_0)$$

$$\Psi = \phi - (\gamma_0 + 2c_0b_0\gamma_0 + c_0G_0^m)\psi$$

$$b_0\phi = 0, \beta_0\phi = 0, \eta_0\phi = 0, b_0\psi = 0, \beta_0\psi = 0, \eta_0\psi = 0$$

$$XY\Psi = \Psi, \eta_0\Psi = 0$$

$$S_R^{(0)} = -\frac{1}{2}\langle \Psi, YQ\Psi \rangle'$$

$$\delta|\Psi\rangle = Q|\lambda\rangle,$$

$$XY\lambda = \lambda, \eta_0\lambda = 0$$

$$S = -\frac{1}{2}\langle \Psi, YQ\Psi \rangle' - \int_0^1 dt \langle A_t(t), QA_\eta(t) + (F(t)\Psi)^2 \rangle'_L$$

$$A_\eta(t) = \eta e^{\Phi(t)}e^{-\Phi(t)}, A_t(t) = \partial_t e^{\Phi(t)}e^{-\Phi(t)}$$

$$F(t)\Psi = \Theta(\beta_0) \left\{ A_\eta(t), \Theta(\beta_0) \{ A_\eta(t), \dots \Theta(\beta_0) \{ A_\eta(t), \Psi \} \dots \} \right\},$$

Hiperpartículas – Condensación de campo.

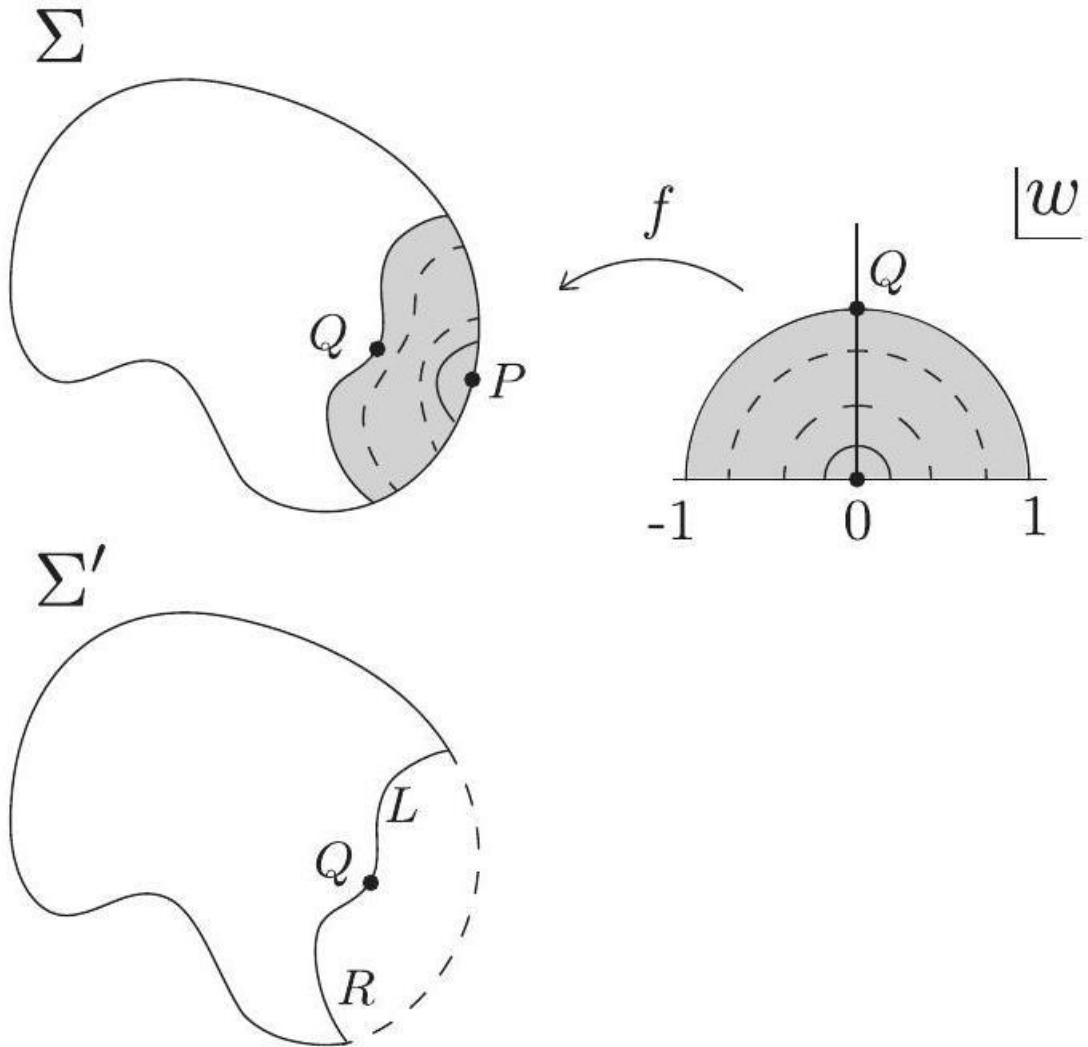
$$-S = Vg_o^{-2} \left[-\frac{1}{2}\phi^2 + \frac{1}{3}\left(\frac{3\sqrt{3}}{4}\right)^3 \phi^3 \right]$$

$$\mathcal{T} = \frac{1}{2\pi^2 g_o^2}$$

$$-S = -Vg_o^{-2} \frac{1}{6} \left(\frac{16}{27}\right)^3 = -V\mathcal{T}2\pi^2 \frac{1}{6} \left(\frac{16}{27}\right)^3 \simeq -0.68V\mathcal{T}$$

$$\langle \Sigma \mid \phi \rangle' = \langle \phi(w=0) \rangle'_\Sigma = \langle f \circ \phi(0) \rangle'_\Sigma$$





$$z = \frac{2}{\pi} \arctan w$$

$$|\Omega_\alpha\rangle \star |\Omega_\beta\rangle = |\Omega_{\alpha+\beta}\rangle,$$

$$\langle \phi, \Omega_\alpha \rangle' = \langle f \circ \phi(0) \rangle'_{C_{\alpha+1}}$$

$$\langle \phi, \Omega_\alpha \rangle' = \langle f_\alpha \circ f \circ \phi(0) \rangle'_{C_1}$$

$$f_{\alpha+\delta\alpha} = f_{\delta\alpha/(1+\alpha)} \circ f_\alpha$$

$$\langle \phi, \Omega_{\alpha+\delta\alpha} \rangle' = \langle f_{\delta\alpha/(1+\alpha)} \circ f_\alpha \circ f \circ \phi(0) \rangle'_{C_1}$$

$$\langle \phi, \Omega_{\alpha+\delta\alpha} \rangle' = \langle \phi, \Omega_\alpha \rangle' - \frac{\delta\alpha}{1+\alpha} \left\langle \oint_0 \frac{dz}{2\pi i} z T(z) f_\alpha \circ f \circ \phi(0) \right\rangle'_{C_1}$$

$$= \langle \phi, \Omega_\alpha \rangle' - \frac{\delta\alpha}{1+\alpha} \left\langle \oint_0 \frac{dz}{2\pi i} z T(z) f \circ \phi(0) \right\rangle'_{C_{\alpha+1}}$$

$$\left\langle \phi, \frac{d}{d\alpha} \Omega_\alpha \right\rangle' = -\frac{1}{\alpha+1} \left\langle \phi \circ \frac{dz}{2\pi i} z T(z) f \circ \phi(0) \right\rangle'_{C_{\alpha+1}}$$

$$\left\langle \phi, \frac{d}{d\alpha} \Omega_\alpha \right\rangle' = - \left\langle \int_R \frac{dz}{2\pi i} T(z) f \circ \phi(0) \right\rangle'_{C_{\alpha+1}}$$

$$\frac{d}{d\alpha} \Omega_\alpha = -\Omega_\alpha \mathcal{K} = -\mathcal{K} \Omega_\alpha$$

$$\mathcal{K}=\left(\int_L T\right)\mathcal{I}$$

$$\Omega_\alpha=e^{-\alpha\mathcal{K}}$$

$$B=\left(\int_L b\right)\mathcal{I}$$

$$\begin{aligned}c^2 &= 0, B^2 = 0, \{c, B\} = 1 \\ [\mathcal{K}, B] &= 0, [\mathcal{K}, c] = \partial c \\ Q\mathcal{K} &= 0, QB = \mathcal{K}, Qc = c\mathcal{K}c\end{aligned}$$

$$\Psi = Fc \frac{\mathcal{K}}{1-F^2} BcF, \text{ with } F = F(\mathcal{K})$$

$$(1-FBcF)^{-1}=1+\frac{F}{1-F^2}BcF$$

$$\Psi = (1-FBcF)Q(1-FBcF)^{-1}$$

$$A = \frac{1-F^2}{\mathcal{K}} B, Q_\Psi A = \mathcal{I}$$

$$F(\mathcal{K}) = e^{-\mathcal{K}/2} = \Omega_{1/2}$$

$$\Psi = \Omega_{1/2} c \frac{\mathcal{K}}{1-\Omega_1} Bc \Omega_{1/2} = \Omega_{1/2} c (1 + \Omega_1 + \Omega_2 + \dots) \mathcal{K} Bc \Omega_{1/2}$$

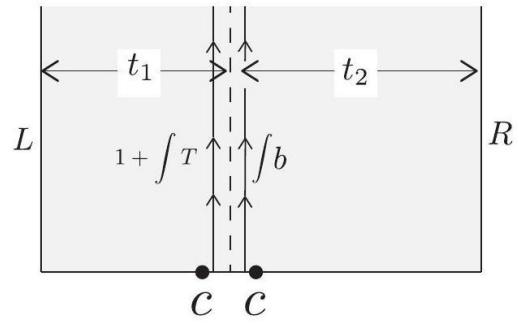
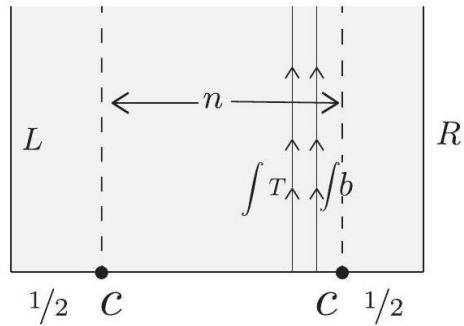
$$= \Omega_{1/2} c \mathcal{K} Bc \Omega_{1/2} + \Omega_{1/2} c \left(\sum_{n=1}^{\infty} \Omega_n \right) \mathcal{K} Bc \Omega_{1/2}$$

$$\Psi = \frac{1}{\sqrt{1+\mathcal{K}}} c (1 + \mathcal{K}) Bc \frac{1}{\sqrt{1+\mathcal{K}}}$$

$$\frac{1}{\sqrt{1+\mathcal{K}}} = \frac{1}{\sqrt{\pi}} \int_0^\infty \frac{dt}{\sqrt{t}} e^{-t} \Omega_t$$

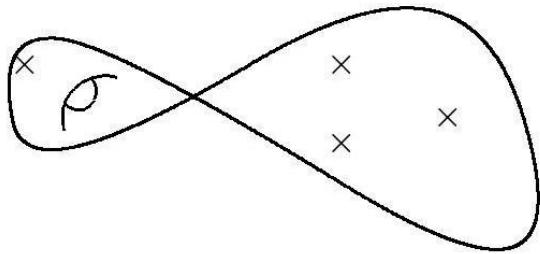
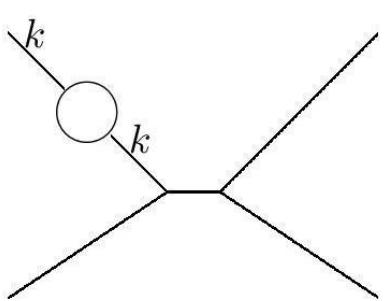
$$\Psi = \frac{1}{\pi} \int_0^\infty \int_0^\infty \frac{dt_1}{\sqrt{t_1}} \frac{dt_2}{\sqrt{t_2}} e^{-(t_1+t_2)} \Omega_{t_1} c (1 + \mathcal{K}) Bc \Omega_{t_2}$$



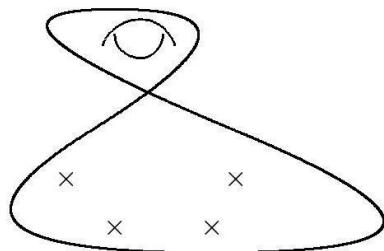
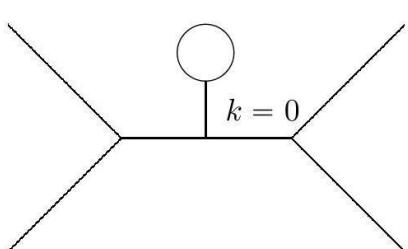


$$-4\pi i \langle E[\mathcal{O}_c] | \hat{\Phi} \rangle = \langle \mathcal{O}_c | c_0^- | B_{\hat{\Phi}} \rangle - \langle \mathcal{O}_c | c_0^- | B_0 \rangle.$$

Renormalización de masa.



$$\frac{1}{L_0 + \bar{L}_0} \delta_{L_0, \bar{L}_0} = \frac{1}{2\pi} \int_{|q| \leq 1} \frac{d^2 q}{|q|^2} q^{L_0} \bar{q}^{\bar{L}_0}$$



$$S = - \int d^D x \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{g_s}{3!} \phi^3 \right]$$

$$S = - \int d^D x \left[\frac{1}{2} \partial_\mu \phi \partial^\mu \phi + \frac{g_s}{3!} \phi^3 + c g_s \phi + \mathcal{O}(\phi^2) \right]$$

$$\phi = \pm \sqrt{-2c} + \mathcal{O}(g_s)$$

$$\begin{aligned} Q(|\Psi\rangle - \mathcal{G}|\widetilde{\Psi}\rangle) &= 0 \\ Q|\widetilde{\Psi}\rangle + \sum_{n=0}^{\infty} \frac{1}{n!} [\Psi^n]_{\text{1PI}} &= 0 \end{aligned}$$

$$|\widetilde{\Psi}\rangle=|\Psi\rangle,$$

$$Q|\Psi\rangle+\sum_{n=0}^\infty\frac{1}{n!}[\Psi^n]_{\text{1PI}}=0$$

$$|\Psi_{k+1}\rangle=-\frac{b_0^+}{L_0^+}\sum_{n=0}^{k+1}\frac{1}{n!}[\Psi_k^n]_{\text{1PI}}+\mathcal{O}(g_s^{k+2})$$

$$|\Psi_{k+1}\rangle=-\frac{b_0^+}{L_0^+}(1-\mathbf{P})\sum_{n=0}^{k+1}\frac{1}{n!}[\Psi_k^n]_{\text{1PI}}+|\psi_{k+1}\rangle+\mathcal{O}(g_s^{k+2}), \mathbf{P}|\psi_{k+1}\rangle=|\psi_{k+1}\rangle$$

$$Q|\psi_{k+1}\rangle=-\mathbf{P}\sum_{n=0}^{k+1}\frac{1}{n!}[\Psi_k^n]_{\text{1PI}}+\mathcal{O}(g_s^{k+2})$$

$$|\Psi\rangle=|\Psi_v\rangle+|\Phi\rangle, |\widetilde{\Psi}\rangle=|\Psi_v\rangle+|\widetilde{\Phi}\rangle$$

$$Q(|\Phi\rangle-\mathcal{G}|\widetilde{\Phi}\rangle)=0, Q|\widetilde{\Phi}\rangle+K|\Phi\rangle=0$$

$$K|A\rangle\equiv\sum_{n=0}^\infty\frac{1}{n!}[\Psi_v^n A]$$

$$(Q+\mathcal{G}K)|\Phi\rangle=0$$

$$QK+KQ+K\mathcal{G}K=0$$

$$2L_0^+=\left(k^2+\hat M^2\right)$$

$$|\Phi_0\rangle=|\phi_n\rangle, |\Phi_{\ell+1}\rangle=-\frac{b_0^+}{L_0^+}(1-P)\mathcal{G}K|\Phi_\ell\rangle+|\phi_n\rangle$$

$$P|\phi_n\rangle=|\phi_n\rangle, Q|\phi_n\rangle=-P\mathcal{G}K|\Phi_{n-1}\rangle+\mathcal{O}(g_s^{n+1})$$

$$(Q+\mathcal{G}K)|\Phi_n\rangle=\mathcal{O}(g_s^{n+1})$$

$$N\equiv\exp\left[\int_0^\infty\frac{dt}{t}F(t)\right]$$

$$F(t)=\frac{1}{2}\mathrm{Tr}\bigl\{(-1)^fe^{-2\pi tL_0}\bigr\},$$

$$N=\exp\left[\int_0^\infty\frac{dt}{2t}\Biggl(\sum_be^{-2\pi th_b}-\sum_f e^{-2\pi th_f}\Biggr)\right]$$

$$\int_0^\infty\frac{dt}{2t}\bigl(e^{-2\pi h_bt}-e^{-2\pi h_ft}\bigr)=\ln\sqrt{\frac{h_f}{h_b}}$$



$$N=\prod_{b,f}\sqrt{\frac{h_f}{h_b}}=\left(\mathrm{sdet}(L_0)\right)^{-1/2}$$

$$\mathrm{sdet}M=1,\text{ with }M_{rs}\equiv\langle\phi_r|c_0|\phi_s\rangle'.$$

$$\int \; \frac{d\phi_b}{\sqrt{2\pi}} e^{-\frac{1}{2} h_b \phi_b^2} = h_b^{-1/2}, \int \; dp_f dq_f e^{-h_f p_f q_f} = h_f$$

$$|\psi_o\rangle=\sum_r\;|\phi_r\rangle\psi_r$$

$$N=\left(\mathrm{sdet}(L_0)\right)^{-1/2}=\int\;\prod_r\;D\psi_r.\mathrm{exp}\left[\frac{1}{2}\langle\psi_o|c_0L_0|\psi_o\rangle'\right],$$

$$\xi^\mu=\frac{1}{\sqrt{2}\pi g_o}\delta y^\mu$$

$$|\psi_o\rangle=p|0\rangle+qc_1c_{-1}|0\rangle+\cdots$$

$$\frac{1}{2}\langle\psi_o|c_0L_0|\psi_o\rangle'=-hpq+\cdots$$

$$|\psi_o\rangle=i\phi c_0|0\rangle+\cdots$$

$$\delta\phi=h\theta$$

$$\int \;dpdq\rightarrow \int \;d\phi e^{S_{\phi,\xi^0}}/\int \;d\theta$$

$$S_{\phi,\xi^0}=\frac{1}{2}\langle\psi_o|Q|\psi_o\rangle'\Big|_{\phi,\xi^0}=-\left(\phi-\sqrt{\frac{h}{2}}\xi^0\right)^2$$

$$\delta\phi=h,\delta\xi^0=\sqrt{2h}$$

$$\int \;d\phi e^{-\phi^2}=\sqrt{\pi}$$

$$\theta=\alpha/g_o$$

$$\int \;d\theta=2\pi/g_o$$

$$\int \;dpdq\prod_\mu\frac{d\xi^\mu}{\sqrt{2\pi}}\rightarrow\big(2\pi\sqrt{\pi}g_o\big)^{-d}\frac{g_o}{2\pi}\sqrt{\pi}\int \;\prod_\mu\;dy^\mu$$

$$N=\big(2\pi\sqrt{\pi}g_o\big)^{-d}\frac{g_o}{2\pi}\sqrt{\pi}\prod_b\;{}'h_b^{-1/2}\prod_f\;{}'h_f^{1/2}(2\pi)^d\delta^{(d)}\biggl(\sum_i\;p_i\biggr)$$



$$\prod_b' h_b^{-1/2} \prod_f' h_f^{1/2} = \prod_b \prod_b'' h_b^{-1/2} \prod_f'' h_f^{1/2} \exp \left[- \int_0^\infty \frac{dt}{2t} \left(\sum_b''' e^{-2\pi t h_b} - \sum_f''' e^{-2\pi t h_f} \right) \right]$$

$$\mathcal{N}(e^{S_1}-1)$$

$$S=S_0+\mathcal{N}(e^{S_1}-1)+\mathcal{O}(\mathcal{N}^2)$$

$$\frac{1}{2}\{S_0,S_0\}+\Delta S_0=0$$

$$\frac{1}{2}\{S_0+S_1,S_0+S_1\}+\Delta(S_0+S_1)=0$$

$$\{S_0,S_1\}+\frac{1}{2}\{S_1,S_1\}+\Delta S_1=0$$

$$\frac{1}{2}\{S,S\}=\frac{1}{2}\{S_0,S_0\}+\mathcal{N}\{S_0,S_1\}e^{S_1}+\mathcal{O}(\mathcal{N}^2)$$

$$\Delta S=\Delta S_0+\mathcal{N}\Delta e^{S_1}=\Delta S_0+\mathcal{N}\left(\frac{1}{2}\{S_1,S_1\}+\Delta S_1\right)e^{S_1}$$

$$\frac{1}{2}\{S,S\}+\Delta S=\mathcal{O}(\mathcal{N}^2)$$

$$\frac{1}{2}\{S_0+2S_1+S_2,S_0+2S_1+S_2\}+\Delta(S_0+2S_1+S_2)=0$$

$$\{S_0,S_2\}+2\{S_1,S_2\}+\{S_1,S_1\}+\frac{1}{2}\{S_2,S_2\}=0$$

$$S=S_0+\mathcal{N}(e^{S_1}-1)+\frac{1}{2}\mathcal{N}^2(e^{S_2}-1)e^{2S_1}+\mathcal{O}(\mathcal{N}^3)$$

$$\frac{1}{2}\{S,S\}+\Delta S=\mathcal{O}(\mathcal{N}^3)$$

Simetrías.

$$T-T^\dagger=i\sum_n T^\dagger|n\rangle\langle n|T\Leftrightarrow\sum_n S^\dagger|n\rangle\langle n|S=1$$

$$\bar{\Delta}(A_1\otimes\cdots\otimes A_n)=\sum_{k=1}^{n-1}(A_1\otimes\cdots A_k)\otimes'(A_{k+1}\cdots A_n), n\geqslant 2, \bar{\Delta}(A_1)=0$$



Apéndice A.

Gravedad cuántica endógena, es decir, por acción de la partícula supermasiva, a razón de su masa extremadamente densa. Modelo Weyl – Polyakov.

$$d^2z \equiv 2dxdy, \delta^2(z, \bar{z}) \equiv \frac{1}{2}\delta(x)\delta(y)$$

$$M_P = G_N^{-1/2} = 1.22 \times 10^{19} \text{GeV}$$

$$M_P^{-1} = 1.6 \times 10^{-33} \text{ cm.}$$

$$G_N^2 E^2 \int dE'E' = \frac{E^2}{M_P^4} \int dE'E'$$

$$X'^\mu(\tau'(\tau)) = X^\mu(\tau)$$

$$S_{\text{pp}} = -m \int d\tau (-\dot{X}^\mu \dot{X}_\mu)^{1/2}$$

$$\delta S_{\text{pp}} = -m \int d\tau \dot{u}_\mu \delta X^\mu$$

$$u^\mu = \dot{X}^\mu (-\dot{X}^\nu \dot{X}_\nu)^{-1/2}$$

$$S'_{\text{pp}} = \frac{1}{2} \int d\tau (\eta^{-1} \dot{X}^\mu \dot{X}_\mu - \eta m^2)$$

$$\eta'(\tau')d\tau' = \eta(\tau)d\tau$$

$$\eta^2 = -\dot{X}^\mu \dot{X}_\mu / m^2$$

$$h_{ab} = \partial_a X^\mu \partial_b X_\mu$$

$$\begin{aligned} S_{\text{NG}} &= \int_M d\tau d\sigma \mathcal{L}_{\text{NG}} \\ \mathcal{L}_{\text{NG}} &= -\frac{1}{2\pi\alpha'} (-\det h_{ab})^{1/2} \end{aligned}$$

$$T = \frac{1}{2\pi\alpha'}$$

$$X'^\mu(\tau, \sigma) = \Lambda_v^\mu X^v(\tau, \sigma) + a^\mu$$

$$X'^\mu(\tau', \sigma') = X^\mu(\tau, \sigma)$$

$$S_P[X, \gamma] = -\frac{1}{4\pi\alpha'} \int_M d\tau d\sigma (-\gamma)^{1/2} \gamma^{ab} \partial_a X^\mu \partial_b X_\mu$$



$$\delta_\gamma S_{\text{P}}[X,\gamma] = -\frac{1}{4\pi\alpha'} \int_M d\tau d\sigma (-\gamma)^{1/2} \delta\gamma^{ab} \left(h_{ab} - \frac{1}{2} \gamma_{ab} \gamma^{cd} h_{cd} \right)$$

$$\delta\gamma=\gamma\gamma^{ab}\delta\gamma_{ab}=-\gamma\gamma_{ab}\delta\gamma^{ab}$$

$$h_{ab}=\frac{1}{2}\gamma_{ab}\gamma^{cd}h_{cd}$$

$$h_{ab}(-h)^{-1/2}=\gamma_{ab}(-\gamma)^{-1/2}$$

$$S_{\text{P}}[X,\gamma]\rightarrow-\frac{1}{2\pi\alpha'}\int~d\tau d\sigma (-h)^{1/2}=S_{\text{NG}}[X]$$

$$\begin{aligned} X'^{\mu}(\tau,\sigma) &= \Lambda^{\mu}_{\nu}X^{\nu}(\tau,\sigma)+a^{\mu}\\ \gamma'_{ab}(\tau,\sigma) &= \gamma_{ab}(\tau,\sigma) \end{aligned}$$

$$\begin{aligned} X'^{\mu}(\tau',\sigma') &= X^{\mu}(\tau,\sigma)\\ \frac{\partial\sigma'^c}{\partial\sigma^a}\frac{\partial\sigma'^d}{\partial\sigma^b}\gamma'_{cd}(\tau',\sigma') &= \gamma_{ab}(\tau,\sigma) \end{aligned}$$

$$\begin{aligned} X'^{\mu}(\tau,\sigma) &= X^{\mu}(\tau,\sigma)\\ \gamma'_{ab}(\tau,\sigma) &= \exp{(2\omega(\tau,\sigma))}\gamma_{ab}(\tau,\sigma) \end{aligned}$$

$$T^{ab}(\tau,\sigma)=-4\pi(-\gamma)^{-1/2}\frac{\delta}{\delta\gamma_{ab}}S_{\text{P}}=-\frac{1}{\alpha'}\Big(\partial^aX^\mu\partial^bX_\mu-\frac{1}{2}\gamma^{ab}\partial_cX^\mu\partial^cX_\mu\Big)$$

$$\gamma_{ab}\frac{\delta}{\delta\gamma_{ab}}S_{\text{P}}=0\Rightarrow T^a_a=0$$

$$T_{ab}=0$$

$$\partial_a [(-\gamma)^{1/2} \gamma^{ab} \partial_b X^\mu] = (-\gamma)^{1/2} \nabla^2 X^\mu = 0$$

$$-\infty<\tau<\infty, 0\leq\sigma\leq\ell.$$

$$\begin{aligned} \delta S_{\text{P}} &= \frac{1}{2\pi\alpha'} \int_{-\infty}^{\infty} d\tau \int_0^{\ell} d\sigma (-\gamma)^{1/2} \delta X^\mu \nabla^2 X_\mu \\ &\quad - \frac{1}{2\pi\alpha'} \int_{-\infty}^{\infty} d\tau (-\gamma)^{1/2} \delta X^\mu \partial^\sigma X_\mu \Big|_{\sigma=0}^{\sigma=\ell} \end{aligned}$$

$$\partial^\sigma X^\mu(\tau,0)=\partial^\sigma X^\mu(\tau,\ell)=0$$

$$n^a\partial_aX_\mu=0\text{ on }\partial M$$

$$\begin{aligned} X^\mu(\tau,\ell) &= X^\mu(\tau,0), \partial^\sigma X^\mu(\tau,\ell)=\partial^\sigma X^\mu(\tau,0)\\ \gamma_{ab}(\tau,\ell) &= \gamma_{ab}(\tau,0) \end{aligned}$$

$$\chi=\frac{1}{4\pi}\int_M d\tau d\sigma (-\gamma)^{1/2}R$$



$$(-\gamma')^{1/2}R'=(-\gamma)^{1/2}(R-2\nabla^2\omega).$$

$$S'_{\rm P}=S_{\rm P}-\lambda \chi=-\int_M d\tau d\sigma (-\gamma)^{1/2}\left(\frac{1}{4\pi\alpha'}\gamma^{ab}\partial_aX^\mu\partial_bX_\mu+\frac{\lambda}{4\pi}R\right)$$

$$x^\pm=2^{-1/2}(x^0\pm x^1), x^i, i=2,\ldots,D-1$$

$$\begin{gathered} a^\mu b_\mu = -a^+b^- - a^-b^+ + a^ib^i \\ a_- = -a^+, a_+ = -a^-, a_i = a^i \end{gathered}$$

$$X^+(\tau)=\tau$$

$$S'_{\rm pp}=\frac{1}{2}\int~~d\tau\bigl(-2\eta^{-1}\dot X^-+\eta^{-1}\dot X^i\dot X^i-\eta m^2\bigr)$$

$$p_- = -\eta^{-1}, p_i = \eta^{-1} \dot X^i$$

$$H=p_-\dot X^-+p_i\dot X^i-L=\frac{p^ip^i+m^2}{2p^+}$$

$$[p_i,X^j]=-i\delta^j_i,[p_-,X^-]=-i$$

$$\begin{gathered} X^+=\tau \\ \partial_\sigma\gamma_{\sigma\sigma}=0 \\ \det\!\gamma_{ab}=-1 \end{gathered}$$

$$f'd\sigma'=fd\sigma$$

$$\begin{bmatrix} \gamma^{\tau\tau} & \gamma^{\tau\sigma} \\ \gamma^{\sigma\tau} & \gamma^{\sigma\sigma} \end{bmatrix} = \begin{bmatrix} -\gamma_{\sigma\sigma}(\tau) & \gamma_{\tau\sigma}(\tau,\sigma) \\ \gamma_{\tau\sigma}(\tau,\sigma) & \gamma_{\sigma\sigma}^{-1}(\tau)(1-\gamma_{\tau\sigma}^2(\tau,\sigma)) \end{bmatrix}.$$

$$L=-\frac{1}{4\pi\alpha'}\int_0^\ell d\sigma\big[\gamma_{\sigma\sigma}\big(2\partial_\tau x^- -\partial_\tau X^i\partial_\tau X^i\big)-2\gamma_{\sigma\tau}\big(\partial_\sigma Y^- -\partial_\tau X^i\partial_\sigma X^i\big)+\gamma_{\sigma\sigma}^{-1}(1-\gamma_{\tau\sigma}^2)\partial_\sigma X^i\partial_\sigma X^i\big]$$

$$\begin{gathered} x^-(\tau)\,=\frac{1}{\ell}\int_0^\ell d\sigma X^-(\tau,\sigma) \\ Y^-(\tau,\sigma)\,=X^-(\tau,\sigma)-x^-(\tau) \end{gathered}$$

$$\gamma_{\tau\sigma}\partial_\tau X^\mu-\gamma_{\tau\tau}\partial_\sigma X^\mu=0\;\;\text{at}\;\sigma=0,\ell.$$

$$\gamma_{\tau\sigma}=0\;\;\text{at}\;\sigma=0,\ell$$

$$\partial_\sigma X^i=0\;\;\text{at}\;\sigma=0,\ell$$

$$L=-\frac{\ell}{2\pi\alpha'}\gamma_{\sigma\sigma}\partial_\tau x^-+\frac{1}{4\pi\alpha'}\int_0^\ell d\sigma\big(\gamma_{\sigma\sigma}\partial_\tau X^i\partial_\tau X^i-\gamma_{\sigma\sigma}^{-1}\partial_\sigma X^i\partial_\sigma X^i\big)$$



$$p_- = -p^+ = \frac{\partial L}{\partial (\partial_\tau x^-)} = -\frac{\ell}{2\pi\alpha'}\gamma_{\sigma\sigma}$$

$$\Pi^i=\frac{\delta L}{\delta(\partial_\tau X^i)}=\frac{1}{2\pi\alpha'}\gamma_{\sigma\sigma}\partial_\tau X^i=\frac{p^+}{\ell}\partial_\tau X^i$$

$$H=p_-\partial_\tau x^- -L+\int_0^\ell d\sigma \Pi_i\partial_\tau X^i=\frac{\ell}{4\pi\alpha' p^+}\int_0^\ell d\sigma\left(2\pi\alpha'\Pi^i\Pi^i+\frac{1}{2\pi\alpha'}\partial_\sigma X^i\partial_\sigma X^i\right)$$

$$\begin{gathered} \partial_\tau x^- = \frac{\partial H}{\partial p_-} = \frac{H}{p^+}, \partial_\tau p^+ = \frac{\partial H}{\partial x^-} = 0 \\ \partial_\tau X^i = \frac{\delta H}{\delta \Pi^i} = 2\pi\alpha' c \Pi^i, \partial_\tau \Pi^i = -\frac{\delta H}{\delta X^i} = \frac{c}{2\pi\alpha'} \partial_\sigma^2 X^i \end{gathered}$$

$$\partial_\tau^2 X^i = c^2 \partial_\sigma^2 X^i$$

$$X^i(\tau,\sigma)=x^i+\frac{p^i}{p^+}\tau+i(2\alpha')^{1/2}\sum_{\substack{n=-\infty\\n\neq 0}}^{\infty}\frac{1}{n}\alpha_n^i\exp\left(-\frac{\pi inc\tau}{\ell}\right)\cos\frac{\pi n\sigma}{\ell}$$

$$\begin{gathered} x^i(\tau) = \frac{1}{\ell}\int_0^\ell d\sigma X^i(\tau,\sigma) \\ p^i(\tau) = \int_0^\ell d\sigma \Pi^i(\tau,\sigma) = \frac{p^+}{\ell}\int_0^\ell d\sigma \partial_\tau X^i(\tau,\sigma) \end{gathered}$$

$$\begin{gathered} [x^-,p^+] = i\eta^{-+} = -i \\ [X^i(\sigma),\Pi^j(\sigma')] = i\delta^{ij}\delta(\sigma-\sigma') \end{gathered}$$

$$\begin{gathered} [x^i,p^j] = i\delta^{ij} \\ [\alpha_m^i,\alpha_n^j] = m\delta^{ij}\delta_{m,-n} \end{gathered}$$

$$\alpha_m^i \sim m^{1/2}a, \alpha_{-m}^i \sim m^{1/2}a^\dagger, m > 0$$

$$\begin{gathered} p^+|0;k\rangle = k^+|0;k\rangle, p^i|0;k\rangle = k^i|0;k\rangle \\ \alpha_m^i|0;k\rangle = 0, m > 0 \end{gathered}$$

$$|N;k\rangle = \left[\prod_{i=2}^{D-1} \prod_{n=1}^{\infty} \frac{(\alpha_{-n}^i)^{N_{in}}}{(n^{N_{in}} N_{in}!)^{1/2}} \right] |0;k\rangle.$$

$$\mathcal{H}=|\text{ vacuum }\rangle\oplus\mathcal{H}_1\oplus\mathcal{H}_2\oplus...$$

$$H=\frac{p^ip^i}{2p^+}+\frac{1}{2p^+\alpha'}\left(\sum_{n=1}^{\infty}\alpha_{-n}^i\alpha_n^i+A\right)$$

$$A=\frac{D-2}{2}\sum_{n=1}^{\infty}~n$$



$$\sum_{n=1}^\infty n \rightarrow -\frac{1}{12}$$

$$\exp\left(-\epsilon\gamma_{\sigma\sigma}^{-1/2}|k_\sigma|\right)$$

$$A\rightarrow \frac{D-2}{2}\sum_{n=1}^{\infty}~n{\rm exp}\left[-\epsilon n(\pi/2p^{+}\alpha'\ell)^{1/2}\right]=\frac{D-2}{2}\biggl(\frac{2\ell p^{+}\alpha'}{\epsilon^2\pi}-\frac{1}{12}+O(\epsilon)\biggr)$$

$$A=\frac{2-D}{24}$$

$$m^2 = 2 p^+ H - p^i p^i = \frac{1}{\alpha'} \Bigl(N + \frac{2-D}{24} \Bigr),$$

$$N=\sum_{i=2}^{D-1}\sum_{n=1}^{\infty}~nN_{in}$$

$$|0;k\rangle,m^2=\frac{2-D}{24\alpha'}$$

$$\alpha_{-1}^i|0;k\rangle,m^2=\frac{26-D}{24\alpha'}$$

$$A=-1,D=26$$

$$S^{ij}=-i\sum_{n=1}^{\infty}\frac{1}{n}(\alpha_{-n}^i\alpha_n^j-\alpha_{-n}^j\alpha_n^i)$$

$${\bf v}=(v^1,0,\dots,0)+(0,v^2,\dots,v^{D-1})$$

$$S^{23}\leq 1+\alpha'm^2$$

$$\sigma'=\sigma+s(\tau){\rm mod}\ell$$

$$\gamma_{\tau\sigma}(\tau,0)=0$$

$$\sigma'=\sigma+s{\rm mod}\ell.$$

$$X^i(\tau,\sigma)=x^i+\frac{p^i}{p^+}\tau+i\left(\frac{\alpha'}{2}\right)^{1/2}\times\sum_{\substack{n=-\infty\\n\neq 0}}^{\infty}\left\{\frac{\alpha_n^i}{n}\text{exp}\left[-\frac{2\pi i n(\sigma+c\tau)}{\ell}\right]+\frac{\tilde{\alpha}_n^i}{n}\text{exp}\left[\frac{2\pi i n(\sigma-c\tau)}{\ell}\right]\right\}$$

$$\alpha_n^i,\tilde{\alpha}_n^i,x^i,p^i,x^-,p^+$$



$$\begin{gathered} [x^-, p^+] = -i \\ \left[x^i, p^j\right] = i\delta^{ij} \\ \left[\alpha_m^i, \alpha_n^j\right] = m\delta^{ij}\delta_{m,-n} \\ \left[\tilde{\alpha}_m^i, \tilde{\alpha}_n^j\right] = m\delta^{ij}\delta_{m,-n} \end{gathered}$$

$$|N,\tilde N;k\rangle=\left[\prod_{i=2}^{D-1}\prod_{n=1}^\infty\frac{(\alpha_{-n}^i)^{N_{in}}(\tilde\alpha_{-n}^i)^{\tilde N_{in}}}{\left(n^{N_{in}}N_{in}!\,n^{\tilde N_{in}}\tilde N_{in}!\right)^{1/2}}\right]|0,0;k\rangle.$$

$$m^2=2p^+H-p^ip^i=\frac{2}{\alpha'}\Biggl[\sum_{n=1}^\infty\left(\alpha_{-n}^i\alpha_n^i+\tilde\alpha_{-n}^i\tilde\alpha_n^i\right)+A+\tilde A\Biggr]=\frac{2}{\alpha'}(N+\tilde N+A+\tilde A)$$

$$A=\tilde{A}=\frac{2-D}{24}$$

$$P=-\int_0^\ell d\sigma \Pi^i\partial_\sigma X^i=-\frac{2\pi}{\ell}\left[\sum_{n=1}^\infty\left(\alpha_{-n}^i\alpha_n^i-\tilde\alpha_{-n}^i\tilde\alpha_n^i\right)+A-\tilde A\right]=-\frac{2\pi}{\ell}(N-\tilde N)$$

$$N=\tilde{N}$$

$$|0,0;k\rangle,m^2=\frac{2-D}{6\alpha'}$$

$$\alpha_{-1}^i\tilde\alpha_{-1}^j|0,0;k\rangle,m^2=\frac{26-D}{6\alpha'}$$

$$A=\tilde{A}=-1,D=26$$

$$e^{ij}=\frac{1}{2}\Big(e^{ij}+e^{ji}-\frac{2}{D-2}\delta^{ij}e^{kk}\Big)+\frac{1}{2}\big(e^{ij}-e^{ji}\big)+\frac{1}{D-2}\delta^{ij}e^{kk},$$

$$\left[\alpha_m^\mu,\alpha_n^\nu\right]=m\eta^{\mu\nu}\delta_{m,-n}$$

$$\sigma'=\ell-\sigma, \tau'=\tau$$

$$\Omega\alpha_n^i\Omega^{-1}=(-1)^n\alpha_n^i$$

$$\begin{gathered}\Omega\alpha_n^i\Omega^{-1}=\tilde\alpha_n^i\\\Omega\tilde\alpha_n^i\Omega^{-1}=\alpha_n^i\end{gathered}$$

$$\begin{gathered}\Omega|N;k\rangle=(-1)^N|N;k\rangle\\\Omega|N,\tilde N;k\rangle=|\tilde N,N;k\rangle\end{gathered}$$

$$\chi=\frac{1}{4\pi}\int_Md\tau d\sigma (-\gamma)^{1/2}R+\frac{1}{2\pi}\int_{\partial M}ds k$$



$$k=\pm t^an_b\nabla_at^b$$

$$\sum_{n=1}^\infty\,(n-\theta)$$

$$X^{25}(\tau,0)=0,X^{25}(\tau,\ell)=y$$

$$X^{25}(\tau,0)=0,\partial^\sigma X^{25}(\tau,\ell)=0$$

$$X^{25}(\tau,\sigma+\ell)=X^{25}(\tau,\sigma)+2\pi R$$

$$X^{25}(\tau,\sigma+\ell)=-X^{25}(\tau,\sigma)$$

$$S = \frac{1}{4\pi\alpha'} \int ~d^2\sigma (\partial_1 X^\mu \partial_1 X_\mu + \partial_2 X^\mu \partial_2 X_\mu)$$

$$z=\sigma^1+i\sigma^2, \bar z=\sigma^1-i\sigma^2$$

$$\partial_z=\frac{1}{2}(\partial_1-i\partial_2), \partial_{\bar{z}}=\frac{1}{2}(\partial_1+i\partial_2)$$

$$\partial_z z=1, \partial_z \bar{z}=0, \partial_{\bar{z}} z=0, \partial_{\bar{z}} \bar{z}=1$$

$$v^z=v^1+iv^2,v^{\bar{z}}=v^1-iv^2,v_z=\frac{1}{2}(v^1-iv^2),v_{\bar{z}}=\frac{1}{2}(v^1+iv^2)$$

$$g_{z\bar{z}}=g_{\bar{z}z}=\frac{1}{2}, g_{zz}=g_{\bar{z}\bar{z}}=0, g^{z\bar{z}}=g^{\bar{z}z}=2, g^{zz}=g^{\bar{z}\bar{z}}=0$$

$$d^2z=2d\sigma^1d\sigma^2$$

$$\int ~d^2z \delta^2(z,\bar{z})=1$$

$$\int_R d^2z (\partial_z v^z + \partial_{\bar{z}} v^{\bar{z}}) = i \oint_{\partial R} (v^z d\bar{z} - v^{\bar{z}} dz)$$

$$S=\frac{1}{2\pi\alpha'}\int ~d^2z \partial X^\mu \bar{\partial} X_\mu$$

$$\partial\bar{\partial} X^\mu(z,\bar{z})=0$$

$$\partial\bigl(\bar{\partial} X^\mu\bigr)=\bar{\partial}(\partial X^\mu)=0$$

$$\langle {\cal F}[X]\rangle = \int ~[dX] \exp{(-S)} {\cal F}[X]$$

$$0=\int ~[dX]\frac{\delta}{\delta X_\mu(z,\bar{z})}\exp{(-S)}=-\int ~[dX]\exp{(-S)}\frac{\delta S}{\delta X_\mu(z,\bar{z})}=-\left\langle\frac{\delta S}{\delta X_\mu(z,\bar{z})}\right\rangle=\frac{1}{\pi\alpha'}\big\langle\partial\bar{\partial} X^\mu(z,\bar{z})\big\rangle$$



$$\left\langle \partial\bar{\partial}X^{\mu}(z,\bar{z})\dots\right\rangle =0$$

$$\partial\bar{\partial}\hat{X}^{\mu}(z,\bar{z})=0$$

$$\begin{aligned} 0 &= \int [dX] \frac{\delta}{\delta X_{\mu}(z, \bar{z})} [\exp(-S) X^{\nu}(z', \bar{z}')] \\ &= \int [dX] \exp(-S) \left[\eta^{\mu\nu} \delta^2(z - z', \bar{z} - \bar{z}') + \frac{1}{\pi\alpha'} \partial_z \partial_{\bar{z}} X^{\mu}(z, \bar{z}) X^{\nu}(z', \bar{z}') \right] \\ &= \eta^{\mu\nu} \langle \delta^2(z - z', \bar{z} - \bar{z}') \rangle + \frac{1}{\pi\alpha'} \partial_z \partial_{\bar{z}} \langle X^{\mu}(z, \bar{z}) X^{\nu}(z', \bar{z}') \rangle \\ \frac{1}{\pi\alpha'} \partial_z \partial_{\bar{z}} \langle X^{\mu}(z, \bar{z}) X^{\nu}(z', \bar{z}') \dots \rangle &= -\eta^{\mu\nu} \langle \delta^2(z - z', \bar{z} - \bar{z}') \dots \rangle \\ \frac{1}{\pi\alpha'} \partial_z \partial_{\bar{z}} X^{\mu}(z, \bar{z}) X^{\nu}(z', \bar{z}') &= -\eta^{\mu\nu} \delta^2(z - z', \bar{z} - \bar{z}') \\ :X^{\mu}(z, \bar{z}): &= X^{\mu}(z, \bar{z}) \\ :X^{\mu}(z_1, \bar{z}_1) X^{\nu}(z_2, \bar{z}_2): &:= X^{\mu}(z_1, \bar{z}_1) X^{\nu}(z_2, \bar{z}_2) + \frac{\alpha'}{2} \eta^{\mu\nu} \ln |z_{12}|^2 \end{aligned}$$

$$z_{ij}=z_i-z_j$$

$$\partial_1 \bar{\partial}_1 :X^{\mu}(z_1, \bar{z}_1) X^{\nu}(z_2, \bar{z}_2): = 0$$

$$\partial\bar{\partial} \ln |z|^2 = 2\pi \delta^2(z, \bar{z}).$$

$$\langle \mathcal{A}_{i_1}(z_1, \bar{z}_1) \mathcal{A}_{i_2}(z_2, \bar{z}_2) \dots \mathcal{A}_{i_n}(z_n, \bar{z}_n) \rangle,$$

$$\mathcal{A}_i(\sigma_1) \mathcal{A}_j(\sigma_2) = \sum_k c_{ij}^k (\sigma_1 - \sigma_2) \mathcal{A}_k(\sigma_2).$$

$$\langle \mathcal{A}_i(\sigma_1) \mathcal{A}_j(\sigma_2) \dots \rangle = \sum_k c_{ij}^k (\sigma_1 - \sigma_2) \langle \mathcal{A}_k(\sigma_2) \dots \rangle$$

$$\begin{aligned} X^{\mu}(z_1, \bar{z}_1) X^{\nu}(z_2, \bar{z}_2) &= -\frac{\alpha'}{2} \eta^{\mu\nu} \ln |z_{12}|^2 + :X^{\nu} X^{\mu}(z_2, \bar{z}_2) \\ &\quad + \sum_{k=1}^{\infty} \frac{1}{k!} [(z_{12})^k :X^{\nu} \partial^k X^{\mu}(z_2, \bar{z}_2): + (\bar{z}_{12})^k :X^{\nu} \bar{\partial}^k X^{\mu}(z_2, \bar{z}_2):] \\ :X^{\mu_1}(z_1, \bar{z}_1) \dots X^{\mu_n}(z_n, \bar{z}_n): &:= X^{\mu_1}(z_1, \bar{z}_1) \dots X^{\mu_n}(z_n, \bar{z}_n) + \sum \text{subtractions} \end{aligned}$$



$$:X^{\mu_1}(z_1,\bar{z}_1)X^{\mu_2}(z_2,\bar{z}_2)X^{\mu_3}(z_3,\bar{z}_3):$$

$$= X^{\mu_1}(z_1,\bar{z}_1)X^{\mu_2}(z_2,\bar{z}_2)X^{\mu_3}(z_3,\bar{z}_3)$$

$$+ \left(\frac{\alpha'}{2} \eta^{\mu_1\mu_2} \ln |z_{12}|^2 X^{\mu_3}(z_3,\bar{z}_3) + 2 \text{ permutations } \right)$$

$$:\mathcal{F}: = \exp \left(\frac{\alpha'}{4} \int \; d^2 z_1 d^2 z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X^\mu(z_1,\bar{z}_1)} \frac{\delta}{\delta X_\mu(z_2,\bar{z}_2)} \right) \mathcal{F}$$

$$\mathcal{F} = \exp \left(-\frac{\alpha'}{4} \int \; d^2 z_1 d^2 z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X^\mu(z_1,\bar{z}_1)} \frac{\delta}{\delta X_\mu(z_2,\bar{z}_2)} \right) :\mathcal{F}: + \sum \text{ contractions}$$

$$:\mathcal{F}::\mathcal{G}: = :\mathcal{F}\mathcal{G}: + \sum \text{ cross-contractions}$$

$$:\mathcal{F}::\mathcal{G}: = \exp \left(-\frac{\alpha'}{2} \int \; d^2 z_1 d^2 z_2 \ln |z_{12}|^2 \frac{\delta}{\delta X_F^\mu(z_1,\bar{z}_1)} \frac{\delta}{\delta X_{G\mu}(z_2,\bar{z}_2)} \right) :\mathcal{F}\mathcal{G}:,$$

$$:\partial X^\mu(z)\partial X_\mu(z)::\partial' X^\nu(z')\partial' X_\nu(z'):$$

$$=: \partial X^\mu(z) \partial X_\mu(z) \partial' X^\nu(z') \partial' X_\nu(z') : - 4 \cdot \frac{\alpha'}{2} (\partial \partial' \ln |z-z'|^2) :\partial X^\mu(z) \partial' X_\mu(z') : + 2$$

$$\cdot \eta^\mu_\mu \left(-\frac{\alpha'}{2} \partial \partial' \ln |z-z'|^2\right)^2$$

$$\sim \frac{D\alpha'^2}{2(z-z')^4}-\frac{2\alpha'}{(z-z')^2}:\partial' X^\mu(z') \partial' X_\mu(z'):-\frac{2\alpha'}{z-z'}:\partial'^2 X^\mu(z') \partial' X_\mu(z'):$$

$$\mathcal{F}=e^{ik_1\cdot X(z,\bar{z})}, \mathcal{G}=e^{ik_2\cdot X(0,0)}.$$

$$:e^{ik_1\cdot X(z,\bar{z})}::e^{ik_2\cdot X(0,0)}:=\exp\left(\frac{\alpha'}{2}k_1\cdot k_2\ln|z|^2\right):e^{ik_1\cdot X(z,\bar{z})}e^{ik_2\cdot X(0,0)}=|z|^{\alpha'k_1\cdot k_2}:e^{ik_1\cdot X(z,\bar{z})}e^{ik_2\cdot X(0,0)}:$$

$$:e^{ik_1\cdot X(z,\bar{z})}::e^{ik_2\cdot X(0,0)}:=|z|^{\alpha'k_1\cdot k_2}:e^{i(k_1+k_2)\cdot X(0,0)}[1+O(z,\bar{z})]:.$$

$$c^k_{ij}(\sigma_1-\sigma_2)_{\rm sym}=\pm c^k_{ji}(\sigma_2-\sigma_1)_{\rm sym}$$

$$\phi'_\alpha(\sigma)=\phi_\alpha(\sigma)+\delta\phi_\alpha(\sigma),$$

$$[d\phi']\mathrm{exp}\,(-S[\phi'])=[d\phi]\mathrm{exp}\,(-S[\phi]).$$

$$\phi'_\alpha(\sigma)=\phi_\alpha(\sigma)+\rho(\sigma)\delta\phi_\alpha(\sigma)$$

$$[d\phi']\mathrm{exp}\,(-S[\phi'])=[d\phi]\mathrm{exp}\,(-S[\phi])\left[1+\frac{i\epsilon}{2\pi}\int\;d^d\sigma g^{1/2}j^a(\sigma)\partial_a\rho(\sigma)+O(\epsilon^2)\right]$$

$$\cdot$$



$$0=\int~[d\phi']\mathrm{exp}~(-S[\phi'])\ldots-\int~[d\phi]\mathrm{exp}~(-S[\phi])\ldots=\frac{\epsilon}{2\pi i}\int~d^d\sigma g^{1/2}\rho(\sigma)\langle\nabla_aj^a(\sigma)\ldots\rangle$$

$$\nabla_a j^a = 0$$

$$\delta \mathcal{A}(\sigma_0) + \frac{\epsilon}{2\pi i} \int_R d^d\sigma g^{1/2} \nabla_a j^a(\sigma) \mathcal{A}(\sigma_0) = 0$$

$$\nabla_a j^a(\sigma) \mathcal{A}(\sigma_0) = g^{-1/2} \delta^d(\sigma - \sigma_0) \frac{2\pi}{i\epsilon} \delta \mathcal{A}(\sigma_0) + \text{ total } \sigma\text{-derivative}$$

$$\int_{\partial R} dA n_a j^a \mathcal{A}(\sigma_0) = \frac{2\pi}{i\epsilon} \delta \mathcal{A}(\sigma_0)$$

$$\oint_{\partial R} (j dz - \tilde{j} d\bar{z}) \mathcal{A}(z_0,\bar{z}_0) = \frac{2\pi}{\epsilon} \delta \mathcal{A}(z_0,\bar{z}_0)$$

$$\mathrm{Res}_{z\rightarrow z_0} j(z) \mathcal{A}(z_0,\bar{z}_0) + \overline{\mathrm{Res}}_{\bar{z}\rightarrow \bar{z}_0} \tilde{j}(\bar{z}) \mathcal{A}(z_0,\bar{z}_0) = \frac{1}{i\epsilon} \delta \mathcal{A}(z_0,\bar{z}_0)$$

$$\delta S = \frac{\epsilon a_\mu}{2\pi\alpha'} \int ~d^2\sigma \partial^a X^\mu \partial_a \rho$$

$$j_a^\mu = \frac{i}{\alpha'} \partial_a X^\mu$$

$$\begin{gathered} j^\mu(z) : e^{ik\cdot X(0,0)} : \sim \frac{k^\mu}{2z} : e^{ik\cdot X(0,0)} : \\ \tilde{j}^\mu(\bar{z}) : e^{ik\cdot X(0,0)} : \sim \frac{k^\mu}{2\bar{z}} : e^{ik\cdot X(0,0)} : \end{gathered}$$

$$\begin{gathered} j_a = i\nu^b T_{ab} \\ T_{ab} = -\frac{1}{\alpha'} :\Big(\partial_a X^\mu \partial_b X_\mu - \frac{1}{2} \delta_{ab} \partial_c X^\mu \partial^c X_\mu\Big): . \end{gathered}$$

$$T_{z\bar{z}}=0$$

$$\bar{\partial}T_{zz}=\partial T_{\bar{z}\bar{z}}=0$$

$$T(z)\equiv T_{zz}(z), \tilde{T}(\bar{z})\equiv T_{\bar{z}\bar{z}}(\bar{z})$$

$$T(z)= -\frac{1}{\alpha'} :\partial X^\mu \partial X_\mu:, \tilde{T}(\bar{z})= -\frac{1}{\alpha'} :\bar{\partial} X^\mu \bar{\partial} X_\mu:,$$

$$j(z)=i\nu(z)T(z), \tilde{j}(\bar{z})=i\nu(z)^*\tilde{T}(\bar{z})$$

$$T(z)X^\mu(0)\sim \frac{1}{z}\partial X^\mu(0), \tilde{T}(\bar{z})X^\mu(0)\sim \frac{1}{\bar{z}}\bar{\partial} X^\mu(0)$$

$$\delta X^\mu = -\epsilon\nu(z)\partial X^\mu - \epsilon\nu(z)^*\bar{\partial} X^\mu$$



$$X'^\mu(z',\bar{z}')=X^\mu(z,\bar{z}), z'=f(z)$$

$$z'=\zeta z$$

$$ds'^2 = dz'd\bar{z}' = \frac{\partial z'}{\partial z}\frac{\partial \bar{z}'}{\partial \bar{z}}dzd\bar{z}$$

$$T(z)\mathcal{A}(0,0) \sim \sum_{n=0}^{\infty} \frac{1}{z^{n+1}} \mathcal{A}^{(n)}(0,0)$$

$$\delta \mathcal{A}(z,\bar{z})=-\epsilon \sum_{n=0}^{\infty} \frac{1}{n!} \big[\partial^n v(z) \mathcal{A}^{(n)}(z,\bar{z}) + \bar{\partial}^n v(z)^* \tilde{\mathcal{A}}^{(n)}(z,\bar{z})\big]$$

$$\mathcal{A}'(z',\bar{z}')=\zeta^{-h\bar{\zeta}-\tilde{h}}\mathcal{A}(z,\bar{z})$$

$$T(z)\mathcal{A}(0,0)=\cdots+\frac{h}{z^2}\mathcal{A}(0,0)+\frac{1}{z}\partial\mathcal{A}(0,0)+\cdots,$$

$$\mathcal{O}'(z',\bar{z}')=(\partial_z z')^{-h}(\partial_{\bar{z}}\bar{z}')^{-\tilde{h}}\mathcal{O}(z,\bar{z})$$

$$T(z)\mathcal{O}(0,0)=\frac{h}{z^2}\mathcal{O}(0,0)+\frac{1}{z}\partial\mathcal{O}(0,0)+\cdots$$

$$\left.\begin{array}{lll} X^\mu & (0,0), & \partial X^\mu \\ \bar{\partial} X^\mu & (0,1), & \partial^2 X^\mu \\ :e^{ik\cdot X}: & & \left(\frac{\alpha' k^2}{4},\frac{\alpha' k^2}{4}\right) \end{array}\right\}$$

$$:\!\left(\prod_i~\partial^{m_i}X^{\mu_i}\right)\!\left(\prod_j~\bar{\partial}^{n_j}X^{v_j}\right)e^{ik\cdot X}:\,,$$

$$\left(\frac{\alpha' k^2}{4}+\sum_im_i,\frac{\alpha' k^2}{4}+\sum_jn_j\right)$$

$$\mathcal{A}_i(z_1,\bar{z}_1)\mathcal{A}_j(z_2,\bar{z}_2)=\sum_kz_{12}^{h_k-h_i-h_j}\bar{z}_{12}^{\tilde{h}_k-\tilde{h}_i-\tilde{h}_j}c_{ij}^k\mathcal{A}_k(z_2,\bar{z}_2)$$

$$\delta e^{ik\cdot X}=-\epsilon v(z)\partial e^{ik\cdot X}-\epsilon v(z)^*\bar{\partial} e^{ik\cdot X}$$

$$T(z)T(0)=\frac{\eta_\mu^\mu}{2z^4}-\frac{2}{\alpha' z^2}:\partial X^\mu(z)\partial X_\mu(0):+ :T(z)T(0):\sim \frac{D}{2z^4}+\frac{2}{z^2}T(0)+\frac{1}{z}\partial T(0)$$

$$\epsilon^{-1}\delta T(z)=-\frac{D}{12}\partial_z^3v(z)-2\partial_zv(z)T(z)-v(z)\partial_zT(z).$$



$$\epsilon^{-1}\delta T(z)=-\frac{c}{12}\partial_z^3v(z)-2\partial_zv(z)T(z)-v(z)\partial_zT(z)$$

$$T(z)T(0)\sim \frac{c}{2z^4}+\frac{2}{z^2}\,T(0)+\frac{1}{z}\partial T(0)$$

$$(\partial_z z')^2 T'(z') = T(z) - \frac{c}{12} \{z',z\}$$

$$\{f,z\}=\frac{2\partial_z^3f\partial_zf-3\partial_z^2f\partial_z^2f}{2\partial_zf\partial_zf}$$

$$T(z)=-\frac{1}{\alpha'}\colon\!\!\partial X^\mu\partial X_\mu\!\!:+V_\mu\partial^2X^\mu,\qquad\qquad\qquad\qquad\\ \tilde{T}(\bar{z})=-\frac{1}{\alpha'}\colon\!\!\bar{\partial} X^\mu\bar{\partial} X_\mu\!\!:+V_\mu\bar{\partial}^2X^\mu,$$

$$c=\tilde c=D+6\alpha' V_\mu V^\mu$$

$$\delta X^\mu=-\epsilon v\partial X^\mu-\epsilon v^*\bar\partial X^\mu-\frac{\epsilon}{2}\alpha' V^\mu[\partial v+(\partial v)^*]$$

$$S=\frac{1}{2\pi}\int\,\,d^2zb\bar\partial c$$

$$h_b=\lambda,h_c=1-\lambda$$

$$\bar\partial c(z)=\bar\partial b(z)=0\\ \bar\partial b(z)c(0)=2\pi\delta^2(z,\bar z)$$

$$\colon\!\! b(z_1)c(z_2)\!\!: = b(z_1)c(z_2) - \frac{1}{z_{12}}$$

$$\bar\partial\frac{1}{z}=\partial\frac{1}{\bar z}=2\pi\delta^2(z,\bar z)$$

$$b(z_1)c(z_2)\sim \frac{1}{z_{12}}, c(z_1)b(z_2)\sim \frac{1}{z_{12}}$$

$$b(z_1)b(z_2)=O(z_{12}), c(z_1)c(z_2)=O(z_{12})$$

$$T(z)=:(\partial b)c:-\lambda\partial(:bc:)\\ \tilde{T}(\bar{z})=0$$

$$c=-3(2\lambda-1)^2+1, \tilde{c}=0$$

$$S=\frac{1}{2\pi}\int\,\,d^2z\tilde b\partial\tilde c$$

$$j=-:bc:$$

$$T(z)j(0)\sim \frac{1-2\lambda}{z^3}+\frac{1}{z^2}j(0)+\frac{1}{z}\partial j(0).$$



$$\epsilon^{-1}\delta j=-v\partial j-j\partial v+\frac{2\lambda-1}{2}\partial^2v,$$

$$(\partial_z z')j_{z'}(z')=j_z(z)+\frac{2\lambda-1}{2}\frac{\partial_z^2z'}{\partial_z z'}.$$

$$\begin{array}{l} \psi = 2^{-1/2} (\psi_1 + i \psi_2), \bar \psi = 2^{-1/2} (\psi_1 - i \psi_2) \\ S = \dfrac{1}{4\pi} \int \; d^2 z (\psi_1 \bar \partial \psi_1 + \psi_2 \bar \partial \psi_2) \\ T = - \dfrac{1}{2} \psi_1 \partial \psi_1 - \dfrac{1}{2} \psi_2 \partial \psi_2 \end{array}$$

$$\beta\gamma CFT$$

$$h_\beta=\lambda, h_\gamma=1-\lambda.$$

$$S=\frac{1}{2\pi}\int\;\;d^2z\beta\bar\partial\gamma$$

$$\bar\partial\gamma(z)=\bar\partial\beta(z)=0.$$

$$\beta(z_1)\gamma(z_2)\sim -\frac{1}{z_{12}}, \gamma(z_1)\beta(z_2)\sim \frac{1}{z_{12}}.$$

$$\begin{array}{c} T=: (\partial \beta) \gamma : - \lambda \partial (:\beta \gamma:) \\ \tilde{T}=0 \end{array}$$

$$c=3(2\lambda-1)^2-1,\tilde c=0$$

$$\sigma^1\sim\sigma^1+2\pi$$

$$-\infty<\sigma^2<\infty$$

$$w=\sigma^1+i\sigma^2$$

$$z=\exp{(-iw)}=\exp{(-i\sigma^1+\sigma^2)}$$

$$T_{zz}(z)=\sum_{m=-\infty}^\infty \frac{L_m}{z^{m+2}}, \tilde{T}_{\bar{z}\bar{z}}(\bar{z})=\sum_{m=-\infty}^\infty \frac{\tilde{L}_m}{\bar{z}^{m+2}}.$$

$$L_m=\oint\limits_c\frac{dz}{2\pi iz}z^{m+2}T_{zz}(z)$$

$$\begin{array}{l} T_{ww}(w)=-\sum\limits_{m=-\infty}^\infty \exp{(im\sigma^1-m\sigma^2)}T_m \\ T_{\bar{w}\bar{w}}(\bar{w})=-\sum\limits_{m=-\infty}^\infty \exp{(-im\sigma^1-m\sigma^2)}\tilde{T}_m \end{array}$$

$$T_m=L_m-\delta_{m,0}\frac{c}{24}, \tilde{T}_m=\tilde{L}_m-\delta_{m,0}\frac{\tilde{c}}{24}$$



$$T_{ww}=(\partial_wz)^2T_{zz}+\frac{c}{24}$$

$$H=\int_0^{2\pi}\frac{d\sigma^1}{2\pi}T_{22}=L_0+\tilde L_0-\frac{c+\tilde c}{24}$$

$$Q_i\{C\} = \oint_c \frac{dz}{2\pi i} j_i$$

$$Q_1\{C_1\}Q_2\{C_2\}-Q_1\{C_3\}Q_2\{C_2\}$$

$$\hat Q_1\hat Q_2 - \hat Q_2\hat Q_1 \equiv [\hat Q_1,\hat Q_2]$$

$$[Q_1,Q_2]\{C_2\}=\oint_{c_2}\frac{dz_2}{2\pi i}\text{Res}_{z_1\rightarrow z_2}j_1(z_1)j_2(z_2)$$

$$[Q,\mathcal{A}(z_2,\bar{z}_2)]=\text{Res}_{z_1\rightarrow z_2}j(z_1)\mathcal{A}(z_2,\bar{z}_2)=\frac{1}{i\epsilon}\delta\mathcal{A}(z_2,\bar{z}_2)$$

$$\tilde Q\{C\}=-\oint_c \frac{d\bar z}{2\pi i} \tilde j$$

$$\left[\tilde Q,\mathcal{A}(z_2,\bar{z}_2)\right]=\overline{\text{Res}}_{\bar z_1\rightarrow \bar z_2}\tilde j(\bar z_1)\mathcal{A}(z_2,\bar{z}_2)=\frac{1}{i\epsilon}\delta\mathcal{A}(z_2,\bar{z}_2)$$

$$\begin{aligned}&\text{Res}_{z_1\rightarrow z_2}z_1^{m+1}T(z_1)z_2^{n+1}T(z_2)=\text{Res}_{z_1\rightarrow z_2}z_1^{m+1}z_2^{n+1}\left(\frac{c}{2z_{12}^4}+\frac{2}{z_{12}^2}T(z_2)+\frac{1}{z_{12}}\partial T(z_2)\right)\\&=\frac{c}{12}(\partial^3z_2^{m+1})z_2^{n+1}+2(\partial z_2^{m+1})z_2^{n+1}T(z_2)+z_2^{m+n+2}\partial T(z_2)=\frac{c}{12}(m^3-m)z_2^{m+n-1}+(m-n)z_2^{m+n+1}T(z_2)\end{aligned}$$

$$[L_m,L_n]=(m-n)L_{m+n}+\frac{c}{12}(m^3-m)\delta_{m,-n}$$

$$[L_0,L_n]=-nL_n$$

$$L_0L_n|\psi\rangle=L_n(L_0-n)|\psi\rangle=(h-n)L_n|\psi\rangle$$

$$[L_0,L_1]=-L_1,[L_0,L_{-1}]=L_{-1},[L_1,L_{-1}]=2L_0$$

$$\mathcal{O}(z)=\sum_{m=-\infty}^\infty \frac{\mathcal{O}_m}{z^{m+h}}$$

$$[L_m,\mathcal{O}_n]=[(h-1)m-n]\mathcal{O}_{m+n}$$

$$0\leq \mathrm{Re} w\leq \pi \iff \mathrm{Im} z\geq 0$$

$$T_{ab}n^at^b=0$$

$$T_{ww}=T_{\bar w\bar w}, \mathrm{Re} w=0,\pi \iff T_{zz}=T_{\bar z\bar z}, \mathrm{Im} z=0$$

$$T_{zz}(z)\equiv T_{\bar z\bar z}(\bar z'), \mathrm{Im} z<0$$



$$L_m = \frac{1}{2\pi i} \int_C (dzz^{m+1}T_{zz} - d\bar{z}\bar{z}^{m+1}T_{\bar{z}\bar{z}}) = \frac{1}{2\pi i} \oint dz z^{m+1} T_{zz}(z)$$

$$[L_m, L_n] = (m-n)L_{m+n} + \frac{c}{12}(m^3-m)\delta_{m,-n}$$

$$\partial X^\mu(z) = -i\left(\frac{\alpha'}{2}\right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\alpha_m^\mu}{z^{m+1}}, \bar{\partial} X^\mu(\bar{z}) = -i\left(\frac{\alpha'}{2}\right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\tilde{\alpha}_m^\mu}{\bar{z}^{m+1}}.$$

$$\begin{aligned}\alpha_m^\mu &= \left(\frac{2}{\alpha'}\right)^{1/2} \oint \frac{dz}{2\pi} z^m \partial X^\mu(z) \\ \tilde{\alpha}_m^\mu &= -\left(\frac{2}{\alpha'}\right)^{1/2} \oint \frac{d\bar{z}}{2\pi} \bar{z}^m \bar{\partial} X^\mu(\bar{z})\end{aligned}$$

$$p^\mu = \frac{1}{2\pi i} \oint dz j^\mu - d\bar{z} \tilde{j}^\mu = \left(\frac{2}{\alpha'}\right)^{1/2} \alpha_0^\mu = \left(\frac{2}{\alpha'}\right)^{1/2} \tilde{\alpha}_0^\mu$$

$$X^\mu(z, \bar{z}) = x^\mu - i \frac{\alpha'}{2} p^\mu \ln |z|^2 + i \left(\frac{\alpha'}{2}\right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{1}{m} \left(\frac{\alpha_m^\mu}{z^m} + \frac{\tilde{\alpha}_m^\mu}{\bar{z}^m} \right)$$

$$\begin{aligned} [\alpha_m^\mu, \alpha_n^\nu] &= [\tilde{\alpha}_m^\mu, \tilde{\alpha}_n^\nu] = m \delta_{m,-n} \eta^{\mu\nu} \\ [x^\mu, p^\nu] &= i \eta^{\mu\nu}\end{aligned}$$

$$L_m \sim \frac{1}{2} \sum_{n=-\infty}^{\infty} \alpha_{m-n}^\mu \alpha_{\mu n}^\circ$$

$$L_0 = \frac{\alpha' p^2}{4} + \sum_{n=1}^{\infty} (\alpha_{-n}^\mu \alpha_{\mu n}^\circ) + a^X$$

$$2L_0|0;0\rangle = (L_1L_{-1} - L_{-1}L_1)|0;0\rangle = 0$$

$$a^X = 0$$

$$L_m = \frac{1}{2} \sum_{n=-\infty}^{\infty} \alpha_{m-n}^\mu \alpha_{\mu n}^\circ$$

$$X^\mu(z, \bar{z}) X^\nu(z', \bar{z}') = \overset{\circ}{X}{}^\mu(z, \bar{z}) X^\nu(z', \bar{z}') \overset{\circ}{2} + \frac{\alpha'}{2} \eta^{\mu\nu} \left[-\ln |z|^2 + \sum_{m=1}^{\infty} \frac{1}{m} \left(\frac{z'^m}{z^m} + \frac{\bar{z}'^m}{\bar{z}^m} \right) \right]$$

$$= \circ X^\mu(z, \bar{z}) X^\nu(z', \bar{z}') \circ - \frac{\alpha'}{2} \eta^{\mu\nu} \text{I} = \text{n} |z - z'|^2$$

$$\circ \circ X^\mu(z, \bar{z}) X^\nu(z', \bar{z}') \circ =: X^\mu(z, \bar{z}) X^\nu(z', \bar{z}'): \quad$$

$$[X^\mu(z, \bar{z}) X^\nu(z', \bar{z}')]_1 = [X^\mu(z, \bar{z}) X^\nu(z', \bar{z}')]_2 + \eta^{\mu\nu} \Delta(z, \bar{z}, z', \bar{z}'),$$



$$[\mathcal{F}]_1=\exp\left(\frac{1}{2}\int\;d^2zd^2z'\Delta(z,\bar{z},z',\bar{z}')\frac{\delta}{\delta X^\mu(z,\bar{z})}\frac{\delta}{\delta X_\mu(z',\bar{z}')} \right)[\mathcal{F}]_2$$

$$L_m = \frac{1}{2}\sum_{n=-\infty}^{\infty} \circ\alpha_{m-n}^{\mu}\alpha_{\mu n}^{\circ} + i\left(\frac{\alpha'}{2}\right)^{1/2}(m+1)V^{\mu}\alpha_{\mu m}$$

$$b(z)=\sum_{m=-\infty}^{\infty}\frac{b_m}{z^{m+\lambda}}, c(z)=\sum_{m=-\infty}^{\infty}\frac{c_m}{z^{m+1-\lambda}}$$

$$\{b_m,c_n\}=\delta_{m,-n}$$

$$\begin{array}{l} b_0|\downarrow\rangle=0,b_0|\uparrow\rangle=|\downarrow\rangle,\\ c_0|\downarrow\rangle=|\uparrow\rangle,c_0|\uparrow\rangle=0,\\ b_n|\downarrow\rangle=b_n|\uparrow\rangle=c_n|\downarrow\rangle=c_n|\uparrow\rangle=0,n>0.\end{array}$$

$$L_m=\sum_{n=-\infty}^{\infty}(m\lambda-n)\circ b_nc_{m-n}^{\circ}+\delta_{m,0}a^g$$

$$2L_0|\downarrow\rangle=(L_1L_{-1}-L_{-1}L_1)|\downarrow\rangle=(\lambda b_0c_1)[(1-\lambda)b_{-1}c_0]|\downarrow\rangle=\lambda(1-\lambda)|\downarrow\rangle$$

$$L_m=\sum_{n=-\infty}^{\infty}(m\lambda-n)\circ b_nc_{m-n}^{\circ}+\frac{\lambda(1-\lambda)}{2}\delta_{m,0}$$

$$N^g=-\frac{1}{2\pi i}\int_0^{2\pi}dwj_w=\sum_{n=1}^{\infty}\left(c_{-n}b_n-b_{-n}c_n\right)+c_0b_0-\frac{1}{2}$$

$$[N^g,b_m]=-b_m,[N^g,c_m]=c_m$$

$$N^g|\downarrow\rangle=-\frac{1}{2}|\downarrow\rangle,N^g|\uparrow\rangle=\frac{1}{2}|\uparrow\rangle.$$

$$\alpha_0^\mu=(2\alpha')^{1/2}p^\mu$$

$$X^\mu(z,\bar{z})=x^\mu-i\alpha' p^\mu\text{ln}\;|z|^2+i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{\substack{m=-\infty \\ m\neq 0}}^{\infty}\frac{\alpha_m^\mu}{m}(z^{-m}+\bar{z}^{-m})$$

$$L_0=\alpha'p^2+\sum_{n=1}^{\infty}\alpha_{-n}^{\mu}\alpha_{\mu n}$$

$$\left[\alpha_m^\mu,\alpha_n^\nu\right]=m\delta_{m,-n}\eta^{\mu\nu},\left[x^\mu,p^\nu\right]=i\eta^{\mu\nu}$$

$$c(z) = \tilde{c}(\bar{z}), b(z) = \tilde{b}(\bar{z}), \mathrm{Im} z = 0$$

$$c(z)\equiv \tilde{c}(\bar{z}'), b(z)\equiv \tilde{b}(\bar{z}'), \mathrm{Im}(z)\leq 0, z'=\bar{z}$$



$$0\leq \mathrm{Re} w\leq 2\pi, w\sim w+2\pi, \mathrm{Im} w\leq 0$$

$$|1\rangle=|0;0\rangle$$

$$\alpha_{-m}^\mu=\left(\frac{2}{\alpha'}\right)^{1/2}\oint\frac{dz}{2\pi}z^{-m}\partial X^\mu(z)\rightarrow\left(\frac{2}{\alpha'}\right)^{1/2}\frac{i}{(m-1)!}\partial^mX^\mu(0)$$

$$\alpha_{-m}^\mu|1\rangle\cong\left(\frac{2}{\alpha'}\right)^{1/2}\frac{i}{(m-1)!}\partial^mX^\mu(0), m\geq 1$$

$$\tilde{\alpha}_{-m}^\mu|1\rangle\cong\left(\frac{2}{\alpha'}\right)^{1/2}\frac{i}{(m-1)!}\bar{\partial}^mX^\mu(0), m\geq 1.$$

$$\alpha_{-m}^\mu:\mathcal{A}(0,0):=\alpha_{-m}^\mu\mathcal{A}(0,0):$$

$$\begin{aligned}\alpha_{-m}^\mu&\rightarrow i\left(\frac{2}{\alpha'}\right)^{1/2}\frac{1}{(m-1)!}\partial^mX^\mu(0),\quad m\geq 1\\\tilde{\alpha}_{-m}^\mu&\rightarrow i\left(\frac{2}{\alpha'}\right)^{1/2}\frac{1}{(m-1)!}\bar{\partial}^mX^\mu(0),\quad m\geq 1\end{aligned}$$

$$x_0^\mu\rightarrow X^\mu(0,0)$$

$$|0;k\rangle\cong:e^{ik\cdot X(0,0)}:$$

$$b_m|1\rangle=0, m\geq -1, c_m|1\rangle=0, m\geq 2$$

$$|1\rangle=b_{-1}|\downarrow\rangle$$

$$\begin{aligned}b_{-m}&\rightarrow\frac{1}{(m-2)!}\partial^{m-2}b(0), m\geq 2\\c_{-m}&\rightarrow\frac{1}{(m+1)!}\partial^{m+1}c(0), m\geq -1\end{aligned}$$

$$(\partial_z w)j_w(w)=j_z(z)+q_0\frac{\partial_z^2w}{\partial_z w}=j_z(z)-\frac{q_0}{z}$$

$$Q^{\mathrm{g}}\equiv\frac{1}{2\pi i}\oint dz j_z=N^{\mathrm{g}}+q_0$$

$$0\leq \mathrm{Re} w\leq \pi, \mathrm{Im} w\leq 0$$

$$\Psi_{\mathcal{A}}[\phi_{\mathrm{b}}]=\int\,\,\,[d\phi_{\mathrm{i}}]_{\phi_{\mathrm{b}}}\exp{(-S[\phi_{\mathrm{i}}])}\mathcal{A}(0)$$

$$\int\,\,\,[d\phi'_{\mathrm{b}}][d\phi_{\mathrm{i}}]_{\phi_{\mathrm{b}},\phi'_{\mathrm{b}}}\exp{(-S[\phi_{\mathrm{i}}])}r^{-L_0-\tilde{L}_0}\Psi[\phi'_{\mathrm{b}}]$$

$$X_{\mathrm{b}}(\theta)=\sum_{n=-\infty}^{\infty}X_ne^{in\theta}, X_n^*=X_{-n}$$



$$\Psi_1[X_{\text{b}}] = \int [dX_{\text{i}}]_{X_{\text{b}}} \exp \left(-\frac{1}{2\pi\alpha'} \int d^2z \partial X \bar{\partial} X \right)$$

$$\begin{aligned} X_{\text{i}} &= X_{\text{cl}} + X'_{\text{i}} \\ X_{\text{cl}}(z, \bar{z}) &= X_0 + \sum_{n=1}^{\infty} (z^n X_n + \bar{z}^n X_{-n}) \end{aligned}$$

$$\Psi_1[X_{\text{b}}] = \exp(-S_{\text{cl}}) \int [dX'_{\text{i}}]_{X_{\text{b}}=0} \exp \left(-\frac{1}{2\pi\alpha'} \int d^2z \partial X' \bar{\partial} X' \right)$$

$$S_{\text{cl}} = \frac{1}{\alpha'} \sum_{m=1}^{\infty} m X_m X_{-m} = \frac{1}{\alpha'} \sum_{m=1}^{\infty} m X_m X_{-m}$$

$$\Psi_1[X_{\text{b}}] \propto \exp \left(-\frac{1}{\alpha'} \sum_{m=1}^{\infty} m X_m X_{-m} \right)$$

$$\begin{aligned} \alpha_n &= -\frac{in}{(2\alpha')^{1/2}} X_{-n} - i \left(\frac{\alpha'}{2}\right)^{1/2} \frac{\partial}{\partial X_n} \\ \tilde{\alpha}_n &= -\frac{in}{(2\alpha')^{1/2}} X_n - i \left(\frac{\alpha'}{2}\right)^{1/2} \frac{\partial}{\partial X_{-n}} \end{aligned}$$

$$\alpha_n \Psi_1[X_{\text{b}}] = \tilde{\alpha}_n \Psi_1[X_{\text{b}}] = 0, n \geq 0,$$

$$|1\rangle \propto |0;0\rangle.$$

$$|\partial^k X\rangle = k! X_k \Psi_1 = -i \left(\frac{\alpha'}{2}\right)^{1/2} (k-1)! \alpha_{-k} |0;0\rangle,$$

$$\mathcal{A}_i(z, \bar{z}) \mathcal{A}_j(0,0),$$

$$\mathcal{A}_{ij,z,\bar{z}} = \sum_k z^{h_k - h_i - h_j} \bar{z}^{\tilde{h}_k - \tilde{h}_i - \tilde{h}_j} c_{ij}^k \mathcal{A}_k,$$

$$\sum_l \begin{array}{c} i \\ \diagdown \quad \diagup \\ k \quad l \quad j \\ \diagup \quad \diagdown \\ m \end{array} = \sum_l \begin{array}{c} i \\ \diagdown \quad \diagup \\ k \quad l \quad j \\ \diagup \quad \diagdown \\ m \end{array}$$



$$\mathcal{A}_i(0,0)\mathcal{A}_j(1,1)\mathcal{A}_k(z,\bar{z}),$$

$$L_m|\mathcal{A}\rangle\cong \oint\frac{dz}{2\pi i}z^{m+1}T(z)\mathcal{A}(0,0)\cong L_m\cdot\mathcal{A}(0,0)$$

$$T(z)\mathcal{A}(0,0)=\sum_{m=-\infty}^{\infty}z^{-m-2}L_m\cdot\mathcal{A}(0,0).$$

$$\delta\mathcal{A}(z,\bar z)=-\epsilon\sum_{n=0}^\infty\frac{1}{n!}\big[\partial^n\nu(z)L_{n-1}+(\partial^n\nu(z))^*\tilde L_{n-1}\big]\cdot\mathcal{A}(z,\bar z)$$

$$\begin{gathered}L_{-1}\cdot\mathcal{A}=\partial\mathcal{A},\tilde L_{-1}\cdot\mathcal{A}=\bar\partial\mathcal{A},\\ L_0\cdot\mathcal{A}=h\mathcal{A},\tilde L_0\cdot\mathcal{A}=\tilde h\mathcal{A}.\end{gathered}$$

$$\begin{gathered}L_0|\mathcal{O}\rangle=h|\mathcal{O}\rangle,\tilde L_0|\mathcal{O}\rangle=\tilde h|\mathcal{O}\rangle\\ L_m|\mathcal{O}\rangle=\tilde L_m|\mathcal{O}\rangle=0,m>0\end{gathered}$$

$$L_m|1\rangle = \tilde L_m|1\rangle = 0, m\geq -1$$

$$SL(2,{\bf C}).$$

$$L_m^\dagger=L_{-m},\tilde L_m^\dagger=\tilde L_{-m}$$

$$\langle 0;k\mid 0;k'\rangle=2\pi\delta(k-k')$$

$$\alpha_m^\dagger=\alpha_{-m},\tilde\alpha_m^\dagger=\tilde\alpha_{-m}$$

$$2h_{\mathcal O}\langle {\mathcal O}\mid {\mathcal O}\rangle=2\langle {\mathcal O}|L_0|{\mathcal O}\rangle=\langle {\mathcal O}|[L_1,L_{-1}]|{\mathcal O}\rangle=\|L_{-1}|{\mathcal O}\rangle\|^2\geq 0$$

$$L_{-1}\cdot\mathcal{O}=\tilde L_{-1}\cdot\mathcal{O}=0$$

$$\partial\mathcal{A}=0\Leftrightarrow h=0,\bar\partial\mathcal{A}=0\Leftrightarrow\tilde h=0.$$

$$E=-\frac{c+\tilde c}{24}$$

$$E=-\frac{\pi(c+\tilde c)}{12\ell}$$

$$\sum_{n=1}^\infty(n-\theta)=\frac{1}{24}-\frac{1}{8}(2\theta-1)^2$$

$$\partial\bar\partial\mathrm{ln}~|z|^2=\partial\frac{1}{\bar z}=\bar\partial\frac{1}{z}=2\pi\delta^2(z,\bar z)$$

$$\left\langle \prod_{i=1}^n: e^{ik_i\cdot X(z_i,\bar z_i)}:\right\rangle=iC^X(2\pi)^D\delta^D\!\left(\sum_{i=1}^nk_i\right)\!\prod_{\substack{i,j=1\\i< j}}^n|z_{ij}|^{\alpha' k_i\cdot k_j}$$



$$S[\phi]=\int \; d^d\sigma {\mathcal L}(\phi(\sigma),\partial_a\phi(\sigma))$$

$$\delta \mathcal{L} = \epsilon \partial_a \mathcal{K}^a$$

$$j^a=2\pi i\left(\frac{\partial \mathcal{L}}{\partial \phi_{\alpha,a}}\epsilon^{-1}\delta\phi_\alpha-\mathcal{K}^a\right).$$

$$\begin{array}{l} \stackrel{\star}{\ast} e^{ik_1\cdot X(y_1)_\star\star\star}e^{ik_2\cdot X(y_2)_\star},y_1\rightarrow y_2(y)\\ :e^{ik\cdot X(z,\bar z)}:, {\rm Im}(z)\rightarrow 0\end{array}$$

$$L_m(L_{-m}|0;0\rangle)-L_{-m}(L_m|0;0\rangle)$$

$$:b(z)c(z'):-\stackrel{\circ}{:b(z)c(z')}\stackrel{\circ}{:}=\frac{(z/z')^{1-\lambda}-1}{z-z'}$$

$$\exp{(i S_{\mathrm{cl}}/\hbar)}$$

$$X^\mu_{\text{upper}}\left(\sigma\right)=X^\mu_{\text{lower}}\left(\pi-\sigma\right)$$

$$\int \; [dXdg] \mathrm{exp}\left(-S \right)$$

$$S=S_X+\lambda\chi$$

$$\begin{gathered} S_X=\frac{1}{4\pi\alpha'}\int_Md^2\sigma g^{1/2}g^{ab}\partial_aX^\mu\partial_bX_\mu \\ \chi=\frac{1}{4\pi}\int_Md^2\sigma g^{1/2}R+\frac{1}{2\pi}\int_{\partial M}ds k \end{gathered}$$

$$\int \; [d\eta dX] \mathrm{exp}\left[\frac{i}{2}\int \; d\tau (\eta^{-1}\dot{X}^\mu\dot{X}_\mu-\eta m^2)\right]$$

$$\eta(\tau)\rightarrow e^{-i\theta}\eta(\tau), X^0(\tau)\rightarrow e^{-i\theta}X^0(\tau)$$

$$\int \; [dedX] \mathrm{exp}\left[-\frac{1}{2}\int \; d\tau \left(e^{-1}\sum_{\mu=1}^D\dot{X}^\mu\dot{X}^\mu+em^2\right)\right]$$

$$g_0^2 \sim g_{\mathrm c} \sim e^{\lambda}$$

$$\chi=\tilde{\chi}+\frac{1}{4}n_{\mathrm c},$$

$$\exp{(-\lambda\tilde{\chi})}=\exp{(-\lambda\chi+\lambda n_{\mathrm c}/4)}$$

$$\int \; \frac{[dXdg]}{V_{\mathrm{diff}\times\mathrm{Weyl}}} \mathrm{exp}\left(-S \right) \equiv Z$$

$$g_{ab}(\sigma)\rightarrow \hat{g}_{ab}(\sigma)$$

$$\hat{g}_{ab}(\sigma)=\delta_{ab}$$



$$\hat{g}_{ab}(\sigma) = \exp [2\omega(\sigma)]\delta_{ab}$$

$$g'^{1/2}R'=g^{1/2}(R-2\nabla^2\omega).$$

$$R_{abcd}=\frac{1}{2}(g_{ac}g_{bd}-g_{ad}g_{bc})R$$

$$z'\equiv \sigma'^1+i\sigma'^2=f(z)$$

$$ds'^2=\exp{(2\omega)}|\partial_zf|^{-2}dz'd\bar{z}'$$

$$\omega=\ln|\partial_z f|$$

$$\zeta\colon g\rightarrow g^\zeta, g^\zeta_{ab}(\sigma')=\exp{[2\omega(\sigma)]}\frac{\partial\sigma^c}{\partial\sigma'^a}\frac{\partial\sigma^d}{\partial\sigma'^b}g_{cd}(\sigma)$$

$$1=\Delta_{\rm FP}(g)\int\,\,[d\zeta]\delta(g-\hat{g}^\zeta)$$

$$Z[\hat{g}]=\int\,\frac{[d\zeta dXdg]}{V_{{\rm diff}\times {\rm Weyl}}}\Delta_{\rm FP}(g)\delta(g-\hat{g}^\zeta){\rm exp}\,(-S[X,g]).$$

$$Z[\hat{g}]=\int\,\frac{[d\zeta dX^\zeta]}{V_{{\rm diff}\times {\rm Weyl}}}\Delta_{\rm FP}(\hat{g}^\zeta){\rm exp}\,(-S[X^\zeta,\hat{g}^\zeta])$$

$$Z[\hat{g}]=\int\,\frac{[d\zeta dX]}{V_{{\rm diff}\times {\rm Weyl}}}\Delta_{\rm FP}(\hat{g}){\rm exp}\,(-S[X,\hat{g}])$$

$$\begin{aligned}\Delta_{\rm FP}(g^\zeta)^{-1}&=\int\,\,[d\zeta']\delta(g^\zeta-\hat{g}^{\zeta'})=\int\,\,[d\zeta']\delta(g-\hat{g}^{\zeta^{-1}\cdot\zeta'})\\&=\int\,\,[d\zeta'']\delta(g-\hat{g}^{\zeta''})=\Delta_{\rm FP}(g)^{-1},\end{aligned}$$

$$Z[\hat{g}]=\int\,\,[dX]\Delta_{\rm FP}(\hat{g}){\rm exp}\,(-S[X,\hat{g}])$$

$$\delta g_{ab}=2\delta\omega g_{ab}-\nabla_a\delta\sigma_b-\nabla_b\delta\sigma_a=(2\delta\omega-\nabla_c\delta\sigma^c)g_{ab}-2(P_1\delta\sigma)_{ab}$$

$$(P_1\delta\sigma)_{ab}=\frac{1}{2}(\nabla_a\delta\sigma_b+\nabla_b\delta\sigma_a-g_{ab}\nabla_c\delta\sigma^c)$$

$$\begin{aligned}\Delta_{\rm FP}(\hat{g})^{-1}&=\int\,\,[d\delta\omega d\delta\sigma]\delta\big[-(2\delta\omega-\hat{\nabla}\cdot\delta\sigma)\hat{g}+2\hat{P}_1\delta\sigma\big]\\&=\int\,\,[d\delta\omega d\beta d\delta\sigma]{\rm exp}\,\left\{2\pi i\int\,\,d^2\sigma\hat{g}^{1/2}\beta^{ab}\big[-(2\delta\omega-\hat{\nabla}\cdot\delta\sigma)\hat{g}+2\hat{P}_1\delta\sigma\big]ab\right\}\\&=\int\,\,[d\beta'd\delta\sigma]{\rm exp}\,\left\{4\pi i\int\,\,d^2\sigma\hat{g}^{1/2}\beta'^{ab}\big(\hat{P}_1\delta\sigma\big)_{ab}\right\}\end{aligned}$$



$$\begin{array}{l} \delta\sigma^a\rightarrow c^a \\ \beta'_{ab}\rightarrow b_{ab} \end{array}$$

$$\Delta_{\rm FP}(\hat g)=\int\,\,\,[dbdc]{\rm exp}\left(-S_{\rm g}\right)$$

$$S_{\rm g}=\frac{1}{2\pi}\int\,\,d^2\sigma\hat g^{1/2}b_{ab}\hat\nabla^a c^b=\frac{1}{2\pi}\int\,\,d^2\sigma\hat g^{1/2}b_{ab}\big(\hat P_1c\big)^{ab}$$

$$Z[\hat g]=\int\,\,\,[dXd bdc]{\rm exp}\left(-S_X-S_{\rm g}\right)$$

$$Z[\hat g]=\left({\rm det}\hat\nabla^2\right)^{-D/2}{\rm det}\hat P_1,$$

$$S_{\rm g}=\frac{1}{2\pi}\int\,\,d^2z(b_{zz}\nabla_{\bar z}c^z+b_{\bar z\bar z}\nabla_zc^{\bar z})=\frac{1}{2\pi}\int\,\,d^2z(b_{zz}\partial_{\bar z}c^z+b_{\bar z\bar z}\partial_zc^{\bar z})$$

$$n_a\delta\sigma^a=0$$

$$n_ac^a=0$$

$$\int_{\partial M}ds n^a b_{ab} \delta c^b = 0$$

$$n_at_bb^{ab}=0$$

$$Z\big[g^\zeta\big] = Z[g]$$

$$\langle...\rangle_g\equiv\int\,\,\,[dXd bdc]{\rm exp}\left(-S[X,b,c,g]\right)...$$

$$\langle...\rangle_{g^\zeta}=\langle...\rangle_g$$

$$\delta\langle...\rangle_g=-\frac{1}{4\pi}\int\,\,d^2\sigma g(\sigma)^{1/2}\delta g_{ab}(\sigma)\big\langle T^{ab}(\sigma)\,...\,\big\rangle_g$$

$$T^{ab}(\sigma)\stackrel{\text{classical}}{=}\frac{4\pi}{g(\sigma)^{1/2}}\frac{\delta}{\delta g_{ab}(\sigma)}S$$

$$\delta_W\langle...\rangle_g=-\frac{1}{2\pi}\int\,\,d^2\sigma g(\sigma)^{1/2}\delta\omega(\sigma)\langle T^a_a(\sigma)\,...\,\rangle_g$$

$$T^a_a\stackrel{?}{=}0$$

$$T^a_a=a_1R$$

$$T_{z\bar z}=\frac{a_1}{2}g_{z\bar z}R$$

$$\nabla^{\bar z}T_{\bar z z}=\frac{a_1}{2}\nabla^{\bar z}(g_{z\bar z}R)=\frac{a_1}{2}\partial_zR,$$



$$\nabla^zT_{zz}=-\nabla^{\bar{z}}T_{\bar{z}z}=-\frac{a_1}{2}\partial_zR$$

$$a_1 \partial_z \nabla^2 \delta \omega \approx 4 a_1 \partial_z^2 \partial_{\bar{z}} \delta \omega,$$

$$\epsilon^{-1}\delta T_{zz}(z) = -\frac{c}{12}\partial_z^3v^z(z)-2\partial_zv^z(z)T_{zz}(z)-v^z(z)\partial_zT_{zz}(z)$$

$$\delta_W T_{zz} = -\frac{c}{6}\partial_z^2\delta\omega$$

$$c=-12a_1,T_a^a=-\frac{c}{12}R$$

$$\begin{array}{l} R\,=\,-2\mathrm{exp}\,(-2\omega)\partial_a\partial_a\omega\\ \nabla^2=\mathrm{exp}\,(-2\omega)\partial_a\partial_a\end{array}$$

$$\delta_W Z[\exp{(2\omega)}\delta]=\frac{a_1}{\pi}Z[\exp{(2\omega)}\delta]\int\,\,d^2\sigma\delta\omega\partial_a\partial_a\omega$$

$$Z[\exp{(2\omega)}\delta]=Z[\delta]\mathrm{exp}\left(-\frac{a_1}{2\pi}\int\,\,d^2\sigma\partial_a\omega\partial_a\omega\right)$$

$$Z[g]=Z[\delta]\mathrm{exp}\,\left[\frac{a_1}{8\pi}\int\,\,d^2\sigma\int\,\,d^2\sigma'g^{1/2}R(\sigma)G(\sigma,\sigma')g^{1/2}R(\sigma')\right]$$

$$g(\sigma)^{1/2}\nabla^2 G(\sigma,\sigma')=\delta^2(\sigma-\sigma')$$

$$\ln\frac{Z[\delta+h]}{Z[\delta]}\approx\frac{a_1}{8\pi^2}\int\,\,d^2z\int\,\,d^2z'(\partial_z^2\ln|z-z'|^2)h_{\bar{z}\bar{z}}(z,\bar{z})\partial_z^2h_{\bar{z}\bar{z}}(z',\bar{z}')$$

$$= -\frac{3a_1}{4\pi^2}\int\,\,d^2z\int\,\,d^2z'\frac{h_{\bar{z}\bar{z}}(z,\bar{z})h_{\bar{z}\bar{z}}(z',\bar{z}')}{(z-z')^4}$$

$$\ln\frac{Z[\delta+h]}{Z[\delta]}\approx\frac{1}{8\pi^2}\int\,\,d^2z\int\,\,d^2z'h_{\bar{z}\bar{z}}(z,\bar{z})h_{\bar{z}\bar{z}}(z',\bar{z}')\langle T_{zz}(z)T_{zz}(z')\rangle_\delta$$

$$c=c^X+c^{\rm g}=D-26$$

$$T_a^a=a_1R+a_2$$

$$S_{\rm ct}=b\int\,\,d^2\sigma g^{1/2}$$

$$T_a^a=a_1R+(a_2+4\pi b)$$

$$\delta_W\ln\langle...\rangle_g=-\frac{1}{2\pi}\int_M d^2\sigma g^{1/2}(a_1R+a_2)\delta\omega-\frac{1}{2\pi}\int_{\partial M}ds(a_3+a_4k+a_5n^a\partial_a)\delta\omega$$

$$S_{\rm ct}=\int_M d^2\sigma g^{1/2}b_1+\int_{\partial M}ds(b_2+b_3k)$$



$$\delta_W S_{\rm ct} = 2 \int_M d^2\sigma g^{1/2} b_1 \delta \omega + \int_{\partial M} ds (b_2 + b_3 n^a \partial_a) \delta \omega$$

$$\begin{aligned}\delta_{W_1}\big(\delta_{W_2}\ln\langle\dots\rangle_g\big)&=\frac{a_1}{\pi}\int_Md^2\sigma g^{1/2}\delta\omega_2\nabla^2\delta\omega_1-\frac{a_4}{2\pi}\int_{\partial M}ds\delta\omega_2n^a\partial_a\delta\omega_1\\&=-\frac{a_1}{\pi}\int_Md^2\sigma g^{1/2}\partial_a\delta\omega_2\partial^a\delta\omega_1+\frac{2a_1-a_4}{2\pi}\int_{\partial M}dsn^a\delta\omega_2\partial_a\delta\omega_1\end{aligned}$$

$$T^a_a(\sigma)=-\frac{\mathcal{C}(\sigma)}{12}R(\sigma)$$

$$-2\pi t\leq {\rm Im} w\leq 0, w\cong w+2\pi$$

$$z=\exp{(-iw)}, \exp{(-2\pi t)}\leq |z|\leq 1$$

$$-2\pi t\leq {\rm Im} w\leq 0, 0\leq {\rm Re} w\leq \pi$$

$$\exp{(-2\pi t)}\leq |z|\leq 1, {\rm Im} z\geq 0$$

$$S_{j_1\dots j_n}(k_1,\dots,k_n)=\sum_{\text{compact topologies}}\int\frac{[dXdg]}{V_{\text{diff}\times\text{Weyl}}}\exp{(-S_X-\lambda\chi)}\prod_{i=1}^n\int~d^2\sigma_i g(\sigma_i)^{1/2}\mathcal{V}_{j_i}(k_i,\sigma_i)$$

$$V_0=2g_{\mathrm c}\int~d^2\sigma g^{1/2}e^{ik\cdot X}\rightarrow g_{\mathrm c}\int~d^2z\colon e^{ik\cdot X}\colon$$

$$m^2=-k^2=-\frac{4}{\alpha'}$$

$$\frac{2g_c}{\alpha'}\int~d^2z\colon\partial X^\mu\bar\partial X^\nu e^{ik\cdot X}\colon.$$

$$h=\tilde h=1+\frac{\alpha' k^2}{4}$$

$$[\mathcal{F}]_{\rm r}=\exp\left(\frac{1}{2}\int~d^2\sigma d^2\sigma'\Delta(\sigma,\sigma')\frac{\delta}{\delta X^\mu(\sigma)}\frac{\delta}{\delta X_\mu(\sigma')}\right)\mathcal{F}$$

$$\Delta(\sigma,\sigma')=\frac{\alpha'}{2}\textrm{ln}~d^2(\sigma,\sigma')$$

$$\delta_W [\mathcal{F}]_{\rm r}=[\delta_W \mathcal{F}]_{\rm r}+\frac{1}{2}\int~d^2\sigma d^2\sigma'\delta_W\Delta(\sigma,\sigma')\frac{\delta}{\delta X^\mu(\sigma)}\frac{\delta}{\delta X_\mu(\sigma')}[\mathcal{F}]_{\rm r}$$

$$\delta_W V_0=2g_{\mathrm c}\int~d^2\sigma g^{1/2}\biggl(2\delta\omega(\sigma)-\frac{k^2}{2}\delta_W\Delta(\sigma,\sigma)\biggr)\bigl[e^{ik\cdot X(\sigma)}\bigr]_{\rm r}$$

$$d^2(\sigma,\sigma')\approx (\sigma-\sigma')^2\mathrm{exp}\,(2\omega(\sigma))$$



$$\Delta(\sigma, \sigma') \approx \alpha' \omega(\sigma) + \frac{\alpha'}{2} \ln (\sigma - \sigma')^2$$

$$\delta_W \Delta(\sigma, \sigma) = \alpha' \delta \omega(\sigma)$$

$$\delta^D(X(\sigma) - x_0) = \int \frac{d^D k}{(2\pi)^D} \exp [ik \cdot (X(\sigma) - x_0)]$$

$$\int_M d^2\sigma g(\sigma)^{1/2} \int_M d^2\sigma' g(\sigma')^{1/2} \delta^D(X(\sigma) - X(\sigma'))$$

$$V_1 = \frac{g_c}{\alpha'} \int d^2\sigma g^{1/2} \left\{ (g^{ab} S_{\mu\nu} + i\epsilon^{ab} A_{\mu\nu}) [\partial_a X^\mu \partial_b X^\nu e^{ik \cdot X}]_r + \alpha' \phi R [e^{ik \cdot X}]_r \right\}$$

$$\begin{aligned}\partial_a \delta_W \Delta(\sigma, \sigma')|_{\sigma'=\sigma} &= \frac{1}{2} \alpha' \partial_a \delta \omega(\sigma) \\ \partial_a \partial'_b \delta_W \Delta(\sigma, \sigma')|_{\sigma'=\sigma} &= \frac{1+\gamma}{2} \alpha' \nabla_a \partial_b \delta \omega(\sigma) \\ \nabla_a \partial_b \delta_W \Delta(\sigma, \sigma')|_{\sigma'=\sigma} &= -\frac{\gamma}{2} \alpha' \nabla_a \partial_b \delta \omega(\sigma)\end{aligned}$$

$$\delta_W V_1 = \frac{g_c}{2} \int d^2\sigma g^{1/2} \delta \omega \left\{ (g^{ab} S_{\mu\nu} + i\epsilon^{ab} A_{\mu\nu}) [\partial_a X^\mu \partial_b X^\nu e^{ik \cdot X}]_r + \alpha' F R [e^{ik \cdot X}]_r \right\}$$

$$\begin{aligned}S_{\mu\nu} &= -k^2 S_{\mu\nu} + k_\nu k^\omega S_{\mu\omega} + k_\mu k^\omega S_{\nu\omega} - (1+\gamma) k_\mu k_\nu S_\omega^\omega + 4k_\mu k_\nu \phi \\ A_{\mu\nu} &= -k^2 A_{\mu\nu} + k_\nu k^\omega A_{\mu\omega} - k_\mu k^\omega A_{\nu\omega} \\ F &= (\gamma-1)k^2 \phi + \frac{1}{2}\gamma k^\mu k^\nu S_{\mu\nu} - \frac{1}{4}\gamma(1+\gamma)k^2 S_\nu^\nu\end{aligned}$$

$$[\nabla^2 X^\mu e^{ik \cdot X}]_r = i \frac{\alpha' \gamma}{4} k^\mu R [e^{ik \cdot X}]_r$$

$$S_{\mu\nu} = A_{\mu\nu} = F = 0$$

$$\begin{aligned}s_{\mu\nu} &\rightarrow s_{\mu\nu} + \xi_\mu k_\nu + k_\mu \xi_\nu \\ a_{\mu\nu} &\rightarrow a_{\mu\nu} + \zeta_\mu k_\nu - k_\mu \zeta_\nu \\ \phi &\rightarrow \phi + \frac{\gamma}{2} k \cdot \xi\end{aligned}$$

$$n^\mu S_{\mu\nu} = n^\mu A_{\mu\nu} = 0$$

$$\begin{aligned}k^2 &= 0 \\ k^\nu S_{\mu\nu} &= k^\mu A_{\mu\nu} = 0 \\ \phi &= \frac{1+\gamma}{4} s_\mu^\mu\end{aligned}$$

$$k^\mu = (1, 1, 0, 0, \dots, 0), n^\mu = \frac{1}{2} (-1, 1, 0, 0, \dots, 0)$$



$$\begin{aligned}\left[e^{ik\cdot X}\right]_{\text{DR}} &= \left[e^{ik\cdot X}\right]_{\text{r}} \\ \left[\partial_a X^\mu e^{ik\cdot X}\right]_{\text{DR}} &= \left[\partial_a X^\mu e^{ik\cdot X}\right]_{\text{r}} \\ \left[\partial_a X^\mu \partial_b X^\nu e^{ik\cdot X}\right]_{\text{DR}} &= \left[\partial_a X^\mu \partial_b X^\nu e^{ik\cdot X}\right]_{\text{r}} - \frac{\alpha'}{12} g_{ab} \eta^{\mu\nu} R \left[e^{ik\cdot X}\right]_{\text{r}} \\ \left[\nabla_a \partial_b X^\mu e^{ik\cdot X}\right]_{\text{DR}} &= \left[\nabla_a \partial_b X^\mu e^{ik\cdot X}\right]_{\text{r}} + i \frac{\alpha'}{12} g_{ab} k^\mu R \left[e^{ik\cdot X}\right]_{\text{r}}\end{aligned}$$

$$g_0 \int_{\partial M} ds \left[e^{ik\cdot X} \right]_{\text{r}}$$

$$-i\frac{g_o}{(2\alpha')^{1/2}}e_\mu\int_{\partial M}ds\left[\dot{X}^\mu e^{ik\cdot X}\right]_{\text{r}}$$

$$S_{\text{pp}} = \frac{1}{2} \int d\tau (\eta^{-1} G_{\mu\nu}(X) \dot{X}^\mu \dot{X}^\nu - \eta m^2)$$

$$S_\sigma = \frac{1}{4\pi\alpha'} \int_M d^2\sigma g^{1/2} g^{ab} G_{\mu\nu}(X) \partial_a X^\mu \partial_b X^\nu$$

$$G_{\mu\nu}(X) = \eta_{\mu\nu} + \chi_{\mu\nu}(X)$$

$$\exp(-S_\sigma) = \exp(-S_{\text{P}}) \left(1 - \frac{1}{4\pi\alpha'} \int_M d^2\sigma g^{1/2} g^{ab} \chi_{\mu\nu}(X) \partial_a X^\mu \partial_b X^\nu + \dots \right)$$

$$\chi_{\mu\nu}(X) = -4\pi g_c e^{ik\cdot X} s_{\mu\nu}.$$

$$S_\sigma = \frac{1}{4\pi\alpha'} \int_M d^2\sigma g^{1/2} \left[(g^{ab} G_{\mu\nu}(X) + i\epsilon^{ab} B_{\mu\nu}(X)) \partial_a X^\mu \partial_b X^\nu + \alpha' R \Phi(X) \right]$$

$$\delta B_{\mu\nu}(X) = \partial_\mu \zeta_\nu(X) - \partial_\nu \zeta_\mu(X)$$

$$H_{\omega\mu\nu}=\partial_\omega B_{\mu\nu}+\partial_\mu B_{\nu\omega}+\partial_\nu B_{\omega\mu}$$

$$G_{\mu\nu}(X) \partial_a X^\mu \partial_b X^\nu = \left[G_{\mu\nu}(x_0) + G_{\mu\nu,\omega}(x_0) Y^\omega + \frac{1}{2} G_{\mu\nu,\omega\rho}(x_0) Y^\omega Y^\rho + \dots \right] \partial_a Y^\mu \partial_b Y^\nu$$

$$\begin{aligned}G_{\mu\nu}(X) &= \eta_{\mu\nu} - 4\pi g_c s_{\mu\nu} e^{ik\cdot X} \\ B_{\mu\nu}(X) &= -4\pi g_c a_{\mu\nu} e^{ik\cdot X} \\ \Phi(X) &= -4\pi g_c \phi e^{ik\cdot X}\end{aligned}$$

$$T_a^a = -\frac{1}{2\alpha'} \beta_{\mu\nu}^G g^{ab} \partial_a X^\mu \partial_b X^\nu - \frac{i}{2\alpha'} \beta_{\mu\nu}^B \epsilon^{ab} \partial_a X^\mu \partial_b X^\nu - \frac{1}{2} \beta^\Phi R$$



$$\begin{aligned}\beta_{\mu\nu}^G \approx & -\frac{\alpha'}{2}\big(\partial^2\chi_{\mu\nu}-\partial_\nu\partial^\omega\chi_{\mu\omega}-\partial_\mu\partial^\omega\chi_{\omega\nu}+\partial_\mu\partial_\nu\chi_\omega^\omega\big)+2\alpha'\partial_\mu\partial_\nu\Phi \\ \beta_{\mu\nu}^B \approx & -\frac{\alpha'}{2}\partial^\omega H_{\omega\mu\nu} \\ \beta^\Phi \approx & \frac{D-26}{6}-\frac{\alpha'}{2}\partial^2\Phi\end{aligned}$$

$$\beta_{\mu\nu}^G=\alpha' {\pmb R}_{\mu\nu}+2\alpha'\nabla_\mu\nabla_\nu\Phi-\frac{\alpha'}{4}H_{\mu\lambda\omega}H_\nu{}^{\lambda\omega}+O(\alpha'^2)$$

$$\beta_{\mu\nu}^B=-\frac{\alpha'}{2}\nabla^\omega H_{\omega\mu\nu}+\alpha'\nabla^\omega\Phi H_{\omega\mu\nu}+O(\alpha'^2)$$

$$\beta^\Phi=\frac{D-26}{6}-\frac{\alpha'}{2}\nabla^2\Phi+\alpha'\nabla_\omega\Phi\nabla^\omega\Phi-\frac{\alpha'}{24}H_{\mu\nu\lambda}H^{\mu\nu\lambda}+O(\alpha'^2)$$

$$\beta_{\mu\nu}^G=\beta_{\mu\nu}^B=\beta^\Phi=0$$

$$G_{\mu\nu}(X)=\eta_{\mu\nu}, B_{\mu\nu}(X)=0, \Phi(X)=\Phi_0$$

$$\lambda = \Phi_0$$

$$G_{\mu\nu}(X)=\eta_{\mu\nu}, B_{\mu\nu}(X)=0, \Phi(X)=V_\mu X^\mu$$

$$V_\mu V^\mu = \frac{26-D}{6\alpha'}.$$

$${\cal S}=\frac{1}{2\kappa_0^2}\int\,\,d^Dx(-G)^{1/2}e^{-2\Phi}[-\frac{2(D-26)}{3\alpha'}+{\pmb R}-\frac{1}{12}H_{\mu\nu\lambda}H^{\mu\nu\lambda}\!+\!4\partial_\mu\Phi\partial^\mu\Phi+O(\alpha')]$$

$$\begin{aligned}\delta {\cal S} = & -\frac{1}{2\kappa_0^2\alpha'}\int\,\,d^Dx(-G)^{1/2}e^{-2\Phi}\big[\delta G_{\mu\nu}\beta^{G\mu\nu} \\ & +\delta B_{\mu\nu}\beta^{B\mu\nu}+\Big(2\delta\Phi-\frac{1}{2}G^{\mu\nu}\delta G_{\mu\nu}\Big)\big(\beta_\omega^{G\omega}-4\beta^\Phi\big)\Big]\end{aligned}$$

$$\tilde{G}_{\mu\nu}(x) = \exp{(2\omega(x))}G_{\mu\nu}(x)$$

$$\tilde{{\pmb R}}=\exp{(-2\omega)}\big[{\pmb R}-2(D-1)\nabla^2\omega-(D-2)(D-1)\partial_\mu\omega\partial^\mu\omega\big]$$

$$\tilde{\Phi}=\Phi-\Phi_0$$

$${\cal S}=\frac{1}{2\kappa^2}\int\,\,d^DX(-\tilde{G})^{1/2}\left[-\frac{2(D-26)}{3\alpha'}e^{4\tilde{\Phi}/(D-2)}+\tilde{{\pmb R}}-\frac{1}{12}e^{-8\tilde{\Phi}/(D-2)}H_{\mu\nu\lambda}\tilde{H}^{\mu\nu\lambda}-\frac{4}{D-2}\partial_\mu\tilde{\Phi}\tilde{\partial}^\mu\tilde{\Phi}+O(\alpha')\right]$$

$$\kappa=(8\pi G_{\rm N})^{1/2}=\frac{(8\pi)^{1/2}}{M_{\rm P}}=(2.43\times 10^{18}{\rm GeV})^{-1}$$

$$\begin{aligned}{\cal S}\,=&-\frac{1}{4}\!\int\,\,d^Dx{\rm Tr}\!\left(F_{\mu\nu}F^{\mu\nu}\right)\\ F_{\mu\nu}\,=&\partial_\mu A_\nu-\partial_\nu A_\mu-ig\big[A_\mu,A_\nu\big]\end{aligned}$$



$$gA_\mu = A'_\mu, gF_{\mu\nu} = F'_{\mu\nu}$$

$$g_{MN}=\begin{bmatrix}\eta_{\mu\nu}&0\\0&g_{mn}(x^p)\end{bmatrix}$$

$$-\frac{1}{4\pi \alpha'}\int_M d^2\sigma g^{1/2}g^{ab}\partial_aX^0\partial_bX^0$$

$$(t,t')=\int\,\,d^2\sigma g^{1/2}t\cdot t'$$

$$\nabla_a j^a=aR$$

$$\delta \langle f \mid i \rangle = -\frac{1}{4\pi} \int \,\, d^2\sigma g(\sigma)^{1/2} \delta g_{ab}(\sigma) \langle f | T^{ab}(\sigma) | i \rangle.$$

$$\langle \psi | T^{ab}(\sigma) | \psi' \rangle = 0$$

$$T_{ab}=T^X_{ab}+T^{\mathrm{g}}_{ab}$$

$$T_{ab}=T^{\mathrm{m}}_{ab}+T^{\mathrm{g}}_{ab}$$

$$\left(L_n^{\mathrm{m}} + A\delta_{n,0}\right)|\psi\rangle=0\,\text{ for } n\geq 0$$

$$\langle \psi | L_n^{\mathrm{m}} | \psi' \rangle = \langle L_{-n}^{\mathrm{m}} \psi \mid \psi' \rangle = 0$$

$$L_n^{\mathrm{m}\dagger}=L_{-n}^{\mathrm{m}}$$

$$|\chi\rangle=\sum_{n=1}^\infty L_{-n}^{\mathrm{m}} |\chi_n\rangle$$

$$|\psi\rangle\cong|\psi\rangle+|\chi\rangle.$$

$$\mathcal{H}_{\mathrm{OCQ}}=\frac{\mathcal{H}_{\mathrm{phys}}}{\mathcal{H}_{\mathrm{null}}}.$$

$$\begin{array}{l} L_0^{\mathrm{m}} = \alpha' p^2 + \alpha_{-1} \cdot \alpha_1 + \cdots \\ L_{\pm 1}^{\mathrm{m}} = (2\alpha')^{1/2} p \cdot \alpha_{\pm 1} + \cdots \end{array}$$

$$|e;k\rangle=e\cdot\alpha_{-1}|0;k\rangle.$$

$$\langle e;k\mid e;k'\rangle=\langle 0;k|e^*\cdot\alpha_1e\cdot\alpha_{-1}|0;k'\rangle=\langle 0;k|(e^*\cdot e+e^*\cdot\alpha_{-1}e\cdot\alpha_1)|0;k'\rangle$$

$$=e^{\mu *}e_{\mu }(2\pi)^D\delta ^D(k-k')$$

$$\alpha_n^{\mu\dagger}=\alpha_{-n}^\mu,$$

$$\langle 0;k\mid 0;k'\rangle=(2\pi)^D\delta^D(k-k'),$$

$$m^2=\frac{1+A}{\alpha'}$$



$$L_1^{\text{m}}|e;k\rangle \propto p\cdot\alpha_1 e\cdot\alpha_{-1}|0;k\rangle = e\cdot k|0;k\rangle = 0$$

$$L_{-1}^{\text{m}}|0;k\rangle = (2\alpha')^{1/2}k\cdot\alpha_{-1}|0;k\rangle.$$

$$k\cdot e=0, e_\mu\cong e_\mu+\gamma k_\mu$$

$$\mathcal{H}_{\text{OCQ}}=\mathcal{H}_{\text{light-cone}},$$

$$\begin{aligned} |0;k\rangle, m^2 &= -\frac{4}{\alpha'} \\ e_{\mu\nu}\alpha_{-1}^\mu\tilde{\alpha}_{-1}^\nu|0;k\rangle, m^2 &= 0, k^\mu e_{\mu\nu} = k^\nu e_{\mu\nu} = 0 \\ e_{\mu\nu} &\cong e_{\mu\nu} + a_\mu k_\nu + k_\mu b_\nu, a\cdot k = b\cdot k = 0 \end{aligned}$$

$$(L_0^{\text{m}}+L_0^{\text{g}})|\psi,\downarrow\rangle=0$$

$$(H^{\text{m}}+H^{\text{g}})|\psi,\downarrow\rangle=0$$

Cuantización BRST:

$$[\delta_\alpha, \delta_\beta] = f^\gamma{}_{\alpha\beta}\delta_\gamma$$

$$F^A(\phi) = 0$$

$$\int \frac{[d\phi_i]}{V_{\text{gauge}}} \exp(-S_1) \rightarrow \int [d\phi_i dB_A db_A dc^\alpha] \exp(-S_1 - S_2 - S_3)$$

$$S_2 = -iB_A F^A(\phi)$$

$$S_3 = b_A c^\alpha \delta_\alpha F^A(\phi)$$

$$\begin{aligned} \delta_B \phi_i &= -i\epsilon c^\alpha \delta_\alpha \phi_i \\ \delta_B B_A &= 0 \\ \delta_B b_A &= \epsilon B_A \\ \delta_B c^\alpha &= \frac{i}{2}\epsilon f_{\beta\gamma}^\alpha c^\beta c^\gamma \end{aligned}$$

$$\delta_B(b_A F^A) = i\epsilon(S_2 + S_3).$$

$$\epsilon \delta \langle f \mid i \rangle = i \langle f \mid \delta_B(b_A \delta F^A) \mid i \rangle = -\epsilon \langle f \mid \{Q_B, b_A \delta F^A\} \mid i \rangle$$

$$\langle \psi \mid \{Q_B, b_A \delta F^A\} \mid \psi' \rangle = 0$$

$$Q_B |\psi\rangle = Q_B |\psi'\rangle = 0$$

$$\epsilon^{-1} \delta_B(b_A B_B M^{AB}) = -B_A B_B M^{AB}$$

$$0 = [Q_B, \{Q_B, b_A \delta F^A\}] = Q_B^2 b_A \delta F^A - Q_B b_A \delta F^A Q_B + Q_B b_A \delta F^A Q_B - b_A \delta F^A Q_B^2 = [Q_B^2, b_A \delta F^A]$$

$$Q_B^2 = 0$$



$$\delta_{\text{B}}(\delta'_{\text{B}}c^{\alpha})=-\frac{1}{2}\epsilon\epsilon'f^{\alpha}_{\beta\gamma}f^{\gamma}_{\delta\epsilon}c^{\beta}c^{\delta}c^{\epsilon}=0$$

$$Q_{\text{B}}|\chi\rangle$$

$$\langle \psi | (Q_{\text{B}} |\chi\rangle) = (\langle \psi | Q_{\text{B}}) |\chi\rangle = 0$$

$$|\psi'\rangle=|\psi\rangle+Q_{\text{B}}|\chi\rangle$$

$$\mathcal{H}_{\text{BRST}}=\frac{\mathcal{H}_{\text{closed}}}{\mathcal{H}_{\text{exact}}}$$

$$\delta_{\tau_1}X^\mu(\tau)=-\delta(\tau-\tau_1)\partial_\tau X^\mu(\tau), \delta_{\tau_1}e(\tau)=-\partial_\tau[\delta(\tau-\tau_1)e(\tau)]$$

$$[\delta_{\tau_1},\delta_{\tau_2}]X^\mu(\tau)=-[\delta(\tau-\tau_1)\partial_\tau\delta(\tau-\tau_2)-\delta(\tau-\tau_2)\partial_\tau\delta(\tau-\tau_1)]\partial_\tau X^\mu(\tau)$$

$$\equiv \int~d\tau_3 f^{\tau_3}{}_{\tau_1\tau_2}\delta_{\tau_3}X^\mu(\tau)$$

$$f^{\tau_3}{}_{\tau_1\tau_2}=\delta(\tau_3-\tau_1)\partial_{\tau_3}\delta(\tau_3-\tau_2)-\delta(\tau_3-\tau_2)\partial_{\tau_3}\delta(\tau_3-\tau_1)$$

$$\begin{aligned}\delta_{\text{B}}X^\mu &= i\epsilon c\dot{X}^\mu \\ \delta_{\text{B}}e &= i\epsilon(c\dot{e}) \\ \delta_{\text{B}}B &= 0 \\ \delta_{\text{B}}b &= \epsilon B \\ \delta_{\text{B}}c &= i\epsilon c\dot{c}\end{aligned}$$

$$S=\int~d\tau\left(\frac{1}{2}e^{-1}\dot{X}^\mu\dot{X}_\mu+\frac{1}{2}em^2+iB(e-1)-e\dot{b}c\right)$$

$$S=\int~d\tau\left(\frac{1}{2}\dot{X}^\mu\dot{X}_\mu+\frac{1}{2}m^2-\dot{b}c\right)$$

$$\begin{aligned}\delta_{\text{B}}X^\mu &= i\epsilon c\dot{X}^\mu \\ \delta_{\text{B}}b &= i\epsilon\left(-\frac{1}{2}\dot{X}^\mu\dot{X}_\mu+\frac{1}{2}m^2-\dot{b}c\right) \\ \delta_{\text{B}}c &= i\epsilon c\dot{c}\end{aligned}$$

$$[p^\mu,X^\nu]=-i\eta^{\mu\nu},\{b,c\}=1$$

$$Q_{\text{B}}=cH$$

$$\begin{aligned}p^\mu|k,\downarrow\rangle &= k^\mu|k,\downarrow\rangle, p^\mu|k,\uparrow\rangle = k^\mu|k,\uparrow\rangle \\ b|k,\downarrow\rangle &= 0, b|k,\uparrow\rangle = |k,\downarrow\rangle \\ c|k,\downarrow\rangle &= |k,\uparrow\rangle, c|k,\uparrow\rangle = 0\end{aligned}$$

$$Q_{\text{B}}|k,\downarrow\rangle=\frac{1}{2}(k^2+m^2)|k,\uparrow\rangle, Q_{\text{B}}|k,\uparrow\rangle=0$$

$$\begin{aligned}|k,\downarrow\rangle, \quad k^2+m^2 &= 0 \\ |k,\uparrow\rangle, \quad \text{all } k^\mu, \quad &\end{aligned}$$

$$|k,\uparrow\rangle, k^2+m^2 \neq 0$$



$$|k,\downarrow\rangle,k^2+m^2=0;\,|k,\uparrow\rangle,k^2+m^2=0$$

$$b|\psi\rangle = 0$$

$$\begin{array}{l} \delta_{\rm B} X^\mu \,=\, i\epsilon(c\partial + \tilde c\bar\partial)X^\mu \\ \delta_{\rm B} b \,=\, i\epsilon(T^X+T^{\rm g}), \delta_{\rm B}\tilde b = i\epsilon(\tilde T^X+\tilde T^{\rm g}) \\ \delta_{\rm B} c \,=\, i\epsilon c\partial c, \delta_{\rm B}\tilde c = i\epsilon \tilde c\bar\partial \tilde c \end{array}$$

$$\frac{i}{4\pi}\int\;d^2\sigma g^{1/2}B^{ab}(\delta_{ab}-g_{ab})$$

$$j_{\rm B}=cT^{\rm m}+\frac{1}{2}\colon cT^{\rm g}\colon+\frac{3}{2}\partial^2c=cT^{\rm m}\!+\!\colon bc\partial c\colon+\frac{3}{2}\partial^2c$$

$$\begin{array}{l} j_{\rm B}(z)b(0)\,\sim\frac{3}{z^3}\!+\!\frac{1}{z^2}j^{\rm g}(0)+\frac{1}{z}T^{{\rm m}+{\rm g}}(0),\\[1mm] j_{\rm B}(z)c(0)\,\sim\frac{1}{z}c\partial c(0),\\[1mm] j_{\rm B}(z){\cal O}^{\rm m}(0,0)\,\sim\frac{h}{z^2}c\theta^{\rm m}(0,0)+\frac{1}{z}[h(\partial c){\cal O}^{\rm m}(0,0)+c\partial\theta^{\rm m}(0,0)]. \end{array}$$

$$Q_{\rm B}=\frac{1}{2\pi i}\oint\,(dzj_{\rm B}-d\bar z\tilde j_{\rm B})$$

$$\{Q_{\rm B},b_m\}=L^{\rm m}_m+L^{\rm g}_m$$

$$Q_{\rm B}=\sum_{n=-\infty}^\infty\left(c_nL^{\rm m}_{-n}+\tilde c_n\tilde L^{\rm m}_{-n}\right)+\sum_{m,n=-\infty}^\infty\frac{(m-n)}{2}\circ^{(c_mc_nb_{-m-n}+\tilde c_m\tilde c_n\tilde b_{-m-n})\circ}+a^{\rm B}(c_0+\tilde c_0)$$

$$\{Q_{\rm B},b_0\}=L^{\rm m}_0+L^{\rm g}_0$$

$$\{Q_{\rm B},Q_{\rm B}\}=0 \text{ only if } c^{\rm m}=26$$

$$j_{\rm B}(z)j_{\rm B}(0)\sim-\frac{c^{\rm m}-18}{2z^3}c\partial c(0)-\frac{c^{\rm m}-18}{4z^2}c\partial^2c(0)-\frac{c^{\rm m}-26}{12z}c\partial^3c(0)$$

$$T(z)j_{\rm B}(0)\sim\frac{c^{\rm m}-26}{2z^4}c(0)+\frac{1}{z^2}j_{\rm B}(0)+\frac{1}{z}\partial j_{\rm B}(0),$$

$$\left[G_I,G_J\right]=ig^K{}_{IJ}G_K$$

$$\{c^I,b_J\}=\delta^I_J,\{c^I,c^J\}=\{b_I,b_J\}=0$$

$$Q_{\rm B}=c^IG^{\rm m}_I-\frac{i}{2}g^K_{IJ}c^Ic^Jb_K=c^I\left(G^{\rm m}_I+\frac{1}{2}G^{\rm g}_I\right)$$

$$G^{\rm g}_I=-ig^K{}_{IJ}c^Jb_K$$



$$Q_{\rm B}^2=\frac{1}{2}\{Q_{\rm B},Q_{\rm B}\}=-\frac{1}{2}\,g^K_{IJ}g^M_{KL}c^Ic^Jc^Lb_M=0$$

$$\begin{array}{ll} \left(\alpha_m^\mu\right)^\dagger = \alpha_{-m}^\mu, & \left(\tilde{\alpha}_m^\mu\right)^\dagger = \tilde{\alpha}_{-m}^\mu \\ \left(b_m\right)^\dagger = b_{-m}, & \left(\tilde{b}_m\right)^\dagger = \tilde{b}_{-m} \\ \left(c_m\right)^\dagger = c_{-m}, & \left(\tilde{c}_m\right)^\dagger = \tilde{c}_{-m}. \end{array}$$

$$\begin{aligned}\langle 0;k|c_0|0;k'\rangle &= (2\pi)^{26}\delta^{26}(k-k')\\ \langle 0;k|\tilde{c}_0c_0|0;k'\rangle &= i(2\pi)^{26}\delta^{26}(k-k')\end{aligned}$$

$$b_0|\psi\rangle=0$$

$$L_0|\psi\rangle=\{Q_{\rm B},b_0\}|\psi\rangle=0$$

$$L_0=\alpha'(p^\mu p_\mu+m^2)$$

$$\alpha'm^2=\sum_{n=1}^{\infty}n\Biggl(N_{bn}+N_{cn}+\sum_{\mu=0}^{25}N_{\mu n}\Biggr)-1.$$

$$|0;\mathbf{k}\rangle,-k^2=-\frac{1}{\alpha'}.$$

$$Q_{\rm B}|0;\mathbf{k}\rangle=0$$

$$|\psi_1\rangle=(e\cdot\alpha_{-1}+\beta b_{-1}+\gamma c_{-1})|0;\mathbf{k}\rangle,-k^2=0$$

$$\begin{aligned}\langle\psi_1\|\psi_1\rangle &= \langle 0;\mathbf{k}\|(e^*\cdot\alpha_1+\beta^*b_1+\gamma^*c_1)(e\cdot\alpha_{-1}+\beta b_{-1}+\gamma c_{-1})\mid 0;\mathbf{k}'\rangle\\ &= (e^*\cdot e+\beta^*\gamma+\gamma^*\beta)\langle 0;\mathbf{k}\|0;\mathbf{k}'\rangle\end{aligned}$$

$$0=Q_{\rm B}|\psi_1\rangle=(2\alpha')^{1/2}(c_{-1}k\cdot\alpha_1+c_1k\cdot\alpha_{-1})|\psi_1\rangle=(2\alpha')^{1/2}(k\cdot e c_{-1}+\beta k\cdot\alpha_{-1})|0;\mathbf{k}\rangle$$

$$Q_{\rm B}|\chi\rangle=(2\alpha')^{1/2}(k\cdot e' c_{-1}+\beta' k\cdot\alpha_{-1})|0;\mathbf{k}\rangle.$$

$$b_0|\psi\rangle=\tilde{b}_0|\psi\rangle=0$$

$$L_0|\psi\rangle=\tilde{L}_0|\psi\rangle=0$$

$$L_0=\frac{\alpha'}{4}(p^2+m^2),\tilde{L}_0=\frac{\alpha'}{4}(p^2+\tilde{m}^2)$$

$$\begin{aligned}\frac{\alpha'}{4}m^2 &= \sum_{n=1}^{\infty}n\Biggl(N_{bn}+N_{cn}+\sum_{\mu=0}^{25}N_{\mu n}\Biggr)-1,\\ \frac{\alpha'}{4}\tilde{m}^2 &= \sum_{n=1}^{\infty}n\Biggl(\tilde{N}_{bn}+\tilde{N}_{cn}+\sum_{\mu=0}^{25}\tilde{N}_{\mu n}\Biggr)-1.\end{aligned}$$

$$L_m=L_m^X+L_m^K+L_m^g$$



$$|N,I;k\rangle, |N,\tilde{N},I;k\rangle,$$

$$\begin{aligned} -\sum_{\mu=0}^{d-1} k_\mu k^\mu &= m^2 \\ \alpha' m^2 &= \sum_{n=1}^{\infty} n \left(N_{bn} + N_{cn} + \sum_{\mu=0}^{d-1} N_{\mu n} \right) + L_0^K - 1 \\ -\sum_{\mu=0}^{d-1} k_\mu k^\mu &= m^2 = \tilde{m}^2 \\ \frac{\alpha'}{4} m^2 &= \sum_{n=1}^{\infty} n \left(N_{bn} + N_{cn} + \sum_{\mu=0}^{d-1} N_{\mu n} \right) + L_0^K - 1 \\ \frac{\alpha'}{4} \tilde{m}^2 &= \sum_{n=1}^{\infty} n \left(\tilde{N}_{bn} + \tilde{N}_{cn} + \sum_{\mu=0}^{d-1} \tilde{N}_{\mu n} \right) + \tilde{L}_0^K - 1 \end{aligned}$$

$$\langle 0,I;\mathbf{k}\|0,I';\mathbf{k}'\rangle=\langle 0,0,I;\mathbf{k}\|0,0,I';\mathbf{k}'\rangle=2k^0(2\pi)^{d-1}\delta^{d-1}(\mathbf{k}-\mathbf{k}')\delta_{I,I'}$$

$$\alpha_m^\pm = 2^{-1/2} (\alpha_m^0 \pm \alpha_m^1)$$

$$[\alpha_m^+, \alpha_n^-] = -m\delta_{m,-n}, [\alpha_m^+, \alpha_n^+] = [\alpha_m^-, \alpha_n^-] = 0$$

$$N^{\text{lc}} = \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{1}{m} : \alpha_{-m}^+ \alpha_m^- :$$

$$Q_{\text{B}} = Q_1 + Q_0 + Q_{-1}$$

$$[N^{\text{lc}}, Q_j] = j Q_j$$

$$[N^{\text{g}}, Q_j] = Q_j$$

$$(Q_1^2) + (\{Q_1, Q_0\}) + (\{Q_1, Q_{-1}\} + Q_0^2) + (\{Q_0, Q_{-1}\}) + (Q_{-1}^2) = 0$$

$$Q_1 = -(2\alpha')^{1/2} k^+ \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \alpha_{-m}^- c_m$$

$$R = \frac{1}{(2\alpha')^{1/2} k^+} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \alpha_{-m}^+ b_m$$

$$S \equiv \{Q_1, R\} = \sum_{m=1}^{\infty} (mb_{-m}c_m + mc_{-m}b_m - \alpha_{-m}^+\alpha_m^- - \alpha_{-m}^-\alpha_m^+) = \sum_{m=1}^{\infty} m(N_{bm} + N_{cm} + N_m^+ + N_m^-)$$



$$|\psi\rangle = \frac{1}{s}\{Q_1,R\}|\psi\rangle = \frac{1}{s}Q_1R|\psi\rangle,$$

$$0=Q_1S|\psi\rangle=SQ_1|\psi\rangle$$

$$S+U\equiv\{Q_{\rm B},R\}$$

$$|\psi\rangle=(1-S^{-1}U+S^{-1}US^{-1}U-\cdots)|\psi_0\rangle$$

$$\langle\psi||\psi'\rangle=\langle\psi_0||\psi'_0\rangle$$

Cuantización BRST-OCQ.

$$\mathcal{H}_{\text{OCQ}} = \mathcal{H}_{\text{BRST}} = \mathcal{H}_{\text{light-cone}}$$

$$|\psi,\downarrow\rangle$$

$$Q_{\rm B}|\psi,\downarrow\rangle=\sum_{n=0}^{\infty}c_{-n}(L_n^{\rm m}-\delta_{n,0})|\psi,\downarrow\rangle=0$$

$$|\psi,\downarrow\rangle-|\psi',\downarrow\rangle$$

$$|\psi,\downarrow\rangle-|\psi',\downarrow\rangle=Q_{\rm B}|\chi\rangle,$$

$$|\chi\rangle=\sum_{n=1}^{\infty}b_{-n}|\chi_n,\downarrow\rangle+\cdots;$$

$$|\psi,\downarrow\rangle-|\psi',\downarrow\rangle=\sum_{m,n=1}^{\infty}c_mL_{-m}^{\rm m}b_{-n}|\chi_n,\downarrow\rangle=\sum_{n=1}^{\infty}L_{-n}^{\rm m}|\chi_n,\downarrow\rangle$$

$$\{[Q_{\rm B},L_m],b_n\}-\{[L_m,b_n],Q_{\rm B}\}-[\{b_n,Q_{\rm B}\},L_m]=0.$$

$$G_I G_J - (-1)^{F_I F_J} G_J G_I = i g^K{}_{IJ} G_K$$

$$(b_{-1})^{N_b}(c_{-1})^{N_c}(\alpha_{-1}^+)^{N^+}(\alpha_{-1}^-)^{N^-}|0\rangle.$$

Simetrías de gauge.

$$\int [dedX]\exp\left[-\frac{1}{2}\int d\tau(e^{-1}\dot{X}^\mu\dot{X}_\mu+em^2)\right]$$

$$\frac{\partial\tau'}{\partial\tau}=e(\tau).$$

$$\tau'(\tau)=\int_0^\tau d\tau'' e(\tau'')$$

$$\tau'(1)=\int_0^1d\tau e(\tau)=l$$



$$e'=l, 0\leq \tau \leq 1$$

$$e'=1, 0\leq \tau \leq l$$

$$\tau \rightarrow \tau + v {\rm mod} 1.$$

$$0\leq \sigma^1\leq 2\pi, 0\leq \sigma^2\leq 2\pi$$

$$(\sigma^1,\sigma^2)\cong (\sigma^1,\sigma^2)+2\pi(m,n)$$

$$ds^2=|d\sigma^1+\tau d\sigma^2|^2$$

$$\tilde{\sigma}^a\cong \tilde{\sigma}^a+2\pi(mu^a+nv^a)$$

$$w\cong w+2\pi(m+n\tau)$$

$$T\colon \tau'=\tau+1, S\colon \tau'=-1/\tau$$

$$\tau' = \frac{a\tau+b}{c\tau+d}$$

$$\begin{bmatrix}\sigma^1\\\sigma^2\end{bmatrix}=\begin{bmatrix}d&b\\c&a\end{bmatrix}\begin{bmatrix}\sigma'^1\\\sigma'^2\end{bmatrix}$$

$$-\frac{1}{2}\leq \mathrm{Re}\tau\leq \frac{1}{2}, |\tau|\geq 1$$

$$\sigma^a\rightarrow \sigma^a+v^a$$

Superficies de Moduli y Riemann.

$$\mathcal{M}_r=\frac{\mathcal{G}_r}{(\text{diff}\times\text{Weyl})_r}$$

$$\mathcal{M}_{r,n}=\frac{\mathcal{G}_r\times\mathcal{M}^n}{(\text{ diff }\times\text{ Weyl })_r}$$

$$\frac{\text{diff}_r}{\text{diff}_{r0}}$$

$$\delta g_{ab}=-2(P_1\delta\sigma)_{ab}+(2\delta\omega-\nabla\cdot\delta\sigma)g_{ab}$$

$$\begin{aligned}0 &= \int \, d^2\sigma g^{1/2} \delta' g_{ab} [-2(P_1\delta\sigma)^{ab}+(2\delta\omega-\nabla\cdot\delta\sigma)g^{ab}] \\&= \int \, d^2\sigma g^{1/2} [-2(P_1^T\delta' g)_a\delta\sigma^a+\delta' g_{ab}g^{ab}(2\delta\omega-\nabla\cdot\delta\sigma)] \\&\quad g^{ab}\delta' g_{ab}=0 \\&\quad (P_1^T\delta' g)_a=0 \\&\quad (P_1\delta\sigma)_{ab}=0\end{aligned}$$



$$\begin{array}{l} \partial_{\bar z}\delta'g_{zz}=\partial_z\delta'g_{\bar z\bar z}=0\\ \partial_{\bar z}\delta z=\partial_z\delta\bar z=0\end{array}$$

$$\mu-\kappa=-3\chi$$

$$\int~d^2\sigma g^{1/2}(P_1\delta\sigma)_{ab}(P_1\delta\sigma)^{ab}=\int~d^2\sigma g^{1/2}\delta\sigma_a(P_1^TP_1\delta\sigma)^a=\int~d^2\sigma g^{1/2}\left(\frac{1}{2}\nabla_a\delta\sigma_b\nabla^a\delta\sigma^b-\frac{R}{4}\delta\sigma_a\delta\sigma^a\right)$$

$$\begin{array}{ll} \chi>0: & \kappa=3\chi,\mu=0 \\ \chi<0: & \kappa=0,\mu=-3\chi \end{array}$$

$$\sigma_m^a=f_{mn}^a(\sigma_n)$$

$$z_m^a=f_{mn}^a(z_n)$$

$$ds^2 \propto dz_m d\bar{z}_m$$

$$w\cong w+2\pi\cong w+2\pi\tau$$

$$S_{j_1\dots j_n}(k_1,\dots,k_n)=\sum_{\substack{\text{compact}\\ \text{topologies}}}\int\frac{[d\phi dg]}{V_{\text{diff}\times\text{Weyl}}}\exp\left(-S_\text{m}-\lambda\chi\right)\prod_{i=1}^n\int~d^2\sigma_i g(\sigma_i)^{1/2}\mathcal{V}_{j_i}(k_i,\sigma_i)$$

$$[dg]d^{2n}\sigma\rightarrow [d\zeta]d^\mu t d^{2n-\kappa}\sigma$$

$$1=\Delta_{\rm FP}(g,\sigma)\int_F d^\mu t\int_{\text{diff}\times\text{Weyl}}[d\zeta]\delta\big(g-\hat{g}(t)^\zeta\big)\prod_{(a,i)\in f}\delta\left(\sigma_i^a-\hat{\sigma}_i^{\zeta a}\right)$$

$$\begin{aligned} S_{j_1\dots j_n}(k_1,\dots,k_n)=&\sum_{\substack{\text{compact}\\ \text{topologies}}}\int_F d^\mu t\Delta_{\rm FP}(\hat{g}(t),\hat{\sigma})\int~[d\phi]\int\prod_{(a,i)\notin f}d\sigma_i^a\\ &\times\exp\left(-S_\text{m}[\phi,\hat{g}(t)]-\lambda\chi\right)\prod_{i=1}^n\left[\hat{g}(\sigma_i)^{1/2}\mathcal{V}_{j_i}(k_i;\sigma_i)\right] \end{aligned}$$

$$\delta g_{ab}=\sum_{k=1}^\mu\delta t^k\partial_{t^k}\hat{g}_{ab}-2\big(\hat{P}_1\delta\sigma\big)_{ab}+(2\delta\omega-\hat{\nabla}\cdot\delta\sigma)\hat{g}_{ab}$$

$$\begin{aligned} \Delta_{\rm FP}(\hat{g},\hat{\sigma})^{-1}&=n_R\int~d^\mu\delta t[d\delta\omega d\delta\sigma]\delta(\delta g_{ab})\prod_{(a,i)\in f}\delta\big(\delta\sigma^a(\hat{\sigma}_i)\big)\\ &=n_R\int~d^\mu\delta t d^\kappa x[d\beta'd\delta\sigma]\times\exp\left[2\pi i\big(\beta',2\hat{P}_1\delta\sigma-\delta t^k\partial_k\hat{g}\big)+2\pi i\sum_{(a,i)\in f}x_{ai}\delta\sigma^a(\hat{\sigma}_i)\right] \end{aligned}$$

$$\begin{array}{l} \delta\sigma^a\rightarrow c^a\\ \beta'_{ab}\rightarrow b_{ab}\\ x_{ai}\rightarrow\eta_{ai}\\ \delta t^k\rightarrow\xi^k \end{array}$$



$$\Delta_{\text{FP}}(\hat{g}, \hat{\sigma}) = \frac{1}{n_R} \int [dbdc] d^\mu \xi d^\kappa \eta \times \exp \left[-\frac{1}{4\pi} (b, 2\hat{P}_1 c - \xi^\kappa \partial_\kappa \hat{g}) + \sum_{(a,i) \in f} \eta_{ai} c^a(\hat{\sigma}_i) \right]$$

$$= \frac{1}{n_R} \int [dbdc] \exp(-S_g) \prod_{k=1}^{\mu} \frac{1}{4\pi} (b, \partial_k \hat{g}) \prod_{(a,i) \in f} c^a(\hat{\sigma}_i)$$

$$S_{j_1\dots j_n}(k_1,\dots,k_n)$$

$$= \sum_{\substack{\text{compact} \\ \text{topologies}}} \int_F \frac{d^\mu t}{n_R} \int [d\phi dbdc] \exp(-S_m - S_g - \lambda \chi)$$

$$\times \prod_{(a,i) \notin f} \int d\sigma_i^a \prod_{k=1}^{\mu} \frac{1}{4\pi} (b, \partial_k \hat{g}) \prod_{(a,i) \in f} c^a(\hat{\sigma}_i) \prod_{i=1}^n \hat{g}(\sigma_i)^{1/2} \mathcal{V}_{j_i}(k_i, \sigma_i)$$

$$c^a(\sigma) = \sum_J c_J C_J^a(\sigma), b_{ab}(\sigma) = \sum_K b_K B_{Kab}(\sigma).$$

$$S_g = \frac{1}{2\pi} (b, P_1 c) = \frac{1}{2\pi} (P_1^T b, c)$$

$$P_1^T P_1 C_J^a = v_J'^2 C_J^a, P_1 P_1^T B_{Kab} = v_K^2 B_{Kab}$$

$$(C_J, C_{J'}) = \int d^2 \sigma g^{1/2} C_J^a C_{J'}{}^a = \delta_{JJ'} \\ (B_K, B_{K'}) = \int d^2 \sigma g^{1/2} B_{Kab} B_{K'}^{ab} = \delta_{KK'}$$

$$(P_1 P_1^T) P_1 C_J = P_1 (P_1^T P_1) C_J = v_J'^2 P_1 C_J$$

$$B_{Jab} = \frac{1}{v_J} (P_1 C_J)_{ab}, v_J = v_J' \neq 0$$

$$\int \prod_{k=1}^{\mu} db_{0k} \prod_{j=1}^{\kappa} dc_{0j} \prod_j db_j dc_j \exp \left(-\frac{v_J b_J c_J}{2\pi} \right) \prod_{k'=1}^{\mu} \frac{1}{4\pi} (b, \partial_{k'} \hat{g}) \prod_{(a,i) \in f} c^a(\sigma_i)$$

$$\Delta_{\text{FP}} = \int \prod_{k=1}^{\mu} db_{0k} \prod_{k'=1}^{\mu} \left[\sum_{k''=1}^{\mu} \frac{b_{0k''}}{4\pi} (B_{0k''}, \partial_{k'} \hat{g}) \right] \times \int \prod_{j=1}^{\kappa} dc_{0j} \prod_{(a,i) \in f} \left[\sum_{j'=1}^{\kappa} c_{0j'} C_{0j'}^a(\sigma_i) \right] \\ \times \int \prod_J db_j dc_j \exp \left(-\frac{v_J b_J c_J}{2\pi} \right)$$

$$\Delta_{\text{FP}} = \det \frac{(B_{0k}, \partial_{k'} \hat{g})}{4\pi} \det C_{0j}^a(\sigma_i) \det' \left(\frac{P_1^T P_1}{4\pi^2} \right)^{1/2}$$



Teorema de Riemann-Roch.

$$\nabla_a j^a = \frac{1 - 2\lambda}{4} R$$

$$\frac{\delta([d\phi]\exp(-S))}{[d\phi]\exp(-S)} = \frac{i\epsilon}{2\pi} \int d^2\sigma g^{1/2} \nabla_a j^a \rightarrow -i\epsilon \frac{2\lambda - 1}{2} \chi$$

$$\frac{1}{2}(\kappa - \mu) = \frac{1}{2}(\dim \ker P_1 - \dim \ker P_1^T).$$

$$\dim \ker P_n - \dim \ker P_n^T = (2n + 1)\chi,$$

$$d^\mu t' = \left| \det \frac{\partial t'}{\partial t} \right| d^\mu t \prod_{k=1}^{\mu} \frac{1}{4\pi} (b, \partial'_k \hat{g}) = \det \left(\frac{\partial t}{\partial t'} \right) \prod_{k=1}^{\mu} \frac{1}{4\pi} (b, \partial_k \hat{g})$$

$$\delta \hat{g}'_{ab}(t; \sigma) = 2\delta\omega(t; \sigma) \hat{g}_{ab}(t; \sigma),$$

$$(b, \partial_k \hat{g}') = \int d^2\sigma \hat{g}'^{1/2} b_{ab} \hat{g}'^{ac} \hat{g}'^{bd} \partial_k \hat{g}'_{cd} = \int d^2\sigma \hat{g}^{1/2} b_{ab} (\hat{g}^{ac} \hat{g}^{bd} \partial_k \hat{g}_{cd} + 2\hat{g}^{ab} \partial_k \omega)$$

$$= (b, \partial_k \hat{g})$$

$$\delta(b, \partial_k \hat{g}) = -2(b, P_1 \partial_k \xi) = -2(P_1^T b, \partial_k \xi) = 0$$

Invariancia BRST.

$$\delta_B \mathcal{V}_m = i\epsilon \partial_a (c^a \mathcal{V}_m)$$

$$\delta_B (b, \partial_k \hat{g}) = i\epsilon (T, \partial_k \hat{g})$$

$$\mu_{ka}^b = \frac{1}{2} \hat{g}^{bc} \partial_k \hat{g}_{ac}$$

$$\frac{1}{2\pi} (b, \mu_k) = \frac{1}{2\pi} \int d^2z (b_{zz} \mu_{k\bar{z}}^z + b_{\bar{z}\bar{z}} \mu_{kz}^{\bar{z}})$$

$$z'_m = z_m + \delta t^k v_{km}^{z_m}(z_m, \bar{z}_m)$$

$$dz'_m d\bar{z}'_m \propto dz_m d\bar{z}_m + \delta t^k \left(\mu_{kz_m}^{\bar{z}_m} dz_m dz_m + \mu_{k\bar{z}_m}^{z_m} d\bar{z}_m d\bar{z}_m \right)$$

$$\mu_{kz_m}^{\bar{z}_m} = \partial_{z_m} v_{km}^{\bar{z}_m}, \mu_{k\bar{z}_m}^{z_m} = \partial_{\bar{z}_m} v_{km}^{z_m}$$

$$\frac{1}{2\pi} (b, \mu_k) = \frac{1}{2\pi i} \sum_m \oint c_m (dz_m v_{km}^{z_m} b_{z_m z_m} - d\bar{z}_m v_{km}^{\bar{z}_m} b_{\bar{z}_m \bar{z}_m})$$

$$\frac{dz_m}{dt^k} = v_{km}^{z_m}$$



$$\left.\frac{\partial z_m}{\partial t^k}\right|_{z_n}=v^{z_m}_{km}-\left.\frac{\partial z_m}{\partial z_n}\right|_t v^{z_n}_{kn}=v^{z_m}_{km}-v^{z_m}_{kn}$$

$$\frac{1}{2\pi}(b,\mu_k)=\frac{1}{2\pi i}\sum_{(mn)}\int_{C_{mn}}\left(dz_m\left.\frac{\partial z_m}{\partial t^k}\right|_{z_n}b_{z_mz_m}-d\bar{z}_m\left.\frac{\partial \bar{z}_m}{\partial t^k}\right|_{z_n}b_{\bar{z}_m\bar{z}_m}\right)$$

$$z=z'+z_v$$

$$\left.\frac{\partial z'}{\partial z_v}\right|_z=-1.$$

$$\int_C \frac{dz'}{2\pi i} b_{z' z'} \int_C \frac{d\bar{z}'_m}{-2\pi i} b_{\bar{z}' \bar{z}'} = b_{-1} \tilde{b}_{-1}$$

$$S(1;...;n) = \sum_{\substack{\text{compact} \\ \text{topologies}}} e^{-\lambda \chi} \int_F \frac{d^mt}{n_R} \biggl\langle \prod_{k=1}^m \; B_k \prod_{i=1}^n \; \hat{\mathcal{V}}_i \biggr\rangle$$

$$m=\mu+2n_\mathrm{c}+n_\mathrm{o}-\kappa=-3\chi+2n_\mathrm{c}+n_\mathrm{o}$$

$$b_{-1}\tilde{b}_{-1}\cdot \tilde{c}c\mathcal{V}_{\mathbf{m}}=\mathcal{V}_{\mathbf{m}}$$

$$u=1/z.$$

$$ds^2=\exp{(2\omega(z,\bar{z}))}dzd\bar{z}$$

$$ds^2=\frac{4r^2dzd\bar{z}}{(1+z\bar{z})^2}=\frac{4r^2dud\bar{u}}{(1+u\bar{u})^2}$$

$$\begin{gathered}\delta u\,=\frac{\partial u}{\partial z}\delta z=-z^{-2}\delta z\\\delta g_{uu}\,=\left(\frac{\partial u}{\partial z}\right)^{-2}\delta g_{zz}=z^4\delta g_{zz}\end{gathered}$$

$$\begin{gathered}\delta z=a_0+a_1z+a_2z^2\\\delta\bar{z}=a_0^*+a_1^*\bar{z}+a_2^*\bar{z}^2\end{gathered}$$

$$z'=\frac{\alpha z+\beta}{\gamma z+\delta}$$

$$z'=1/\bar{z}$$

$$z'=\bar{z}$$

$$z'=-1/\bar{z}$$

$$Z[J]=\Bigl\langle\exp\Bigl(i\int\;d^2\sigma J(\sigma)\cdot X(\sigma)\Bigr)\Bigr\rangle$$



$$\begin{aligned} X^\mu(\sigma) &= \sum_I \, x_I^\mu \mathrm{x}_I(\sigma) \\ \nabla^2 \mathrm{x}_I &= -\omega_I^2 \mathrm{x}_I \\ \int_M d^2\sigma g^{1/2} \mathrm{x}_I \mathrm{x}_{I'} &= \delta_{II'} \end{aligned}$$

$$Z[J]=\prod_{I,\mu}\,\int\,\,dx_I^\mu\exp\left(-\frac{\omega_I^2x_I^\mu x_{I\mu}}{4\pi\alpha'}+ix_I^\mu J_{I\mu}\right)$$

$$J_I^\mu = \int \,\,d^2\sigma J^\mu(\sigma)\mathrm{X}_I(\sigma)$$

$$\mathrm{X}_0=\left(\int \,\,d^2\sigma g^{1/2}\right)^{-1/2}$$

$$\begin{aligned} Z[J] &= i(2\pi)^d \delta^d(J_0) \prod_{I\neq 0} \left(\frac{4\pi^2\alpha'}{\omega_I^2}\right)^{\frac{d}{2}} \exp\left(-\frac{\pi\alpha' J_I\cdot J_I}{\omega_I^2}\right) \\ &= i(2\pi)^d \delta^d(J_0) \left(\det'\frac{-\nabla^2}{4\pi^2\alpha'}\right)^{-d/2} \times \exp\left(-\frac{1}{2}\int \,\,d^2\sigma d^2\sigma' J(\sigma)\cdot J(\sigma') G'(\sigma,\sigma')\right) \end{aligned}$$

$$G'(\sigma_1,\sigma_2)=\sum_{I\neq 0}\frac{2\pi\alpha'}{\omega_I^2}\mathrm{X}_I(\sigma_1)\mathrm{X}_I(\sigma_2)$$

$$-\frac{1}{2\pi\alpha'}\nabla^2G'(\sigma_1,\sigma_2)=\sum_{I\neq 0}\,\mathrm{X}_I(\sigma_1)\mathrm{X}_I(\sigma_2)=g^{-1/2}\delta^2(\sigma_1-\sigma_2)-\mathrm{X}_0^2$$

$$G'(\sigma_1,\sigma_2)=-\frac{\alpha'}{2}\ln|z_{12}|^2+f(z_1,\bar z_1)+f(z_2,\bar z_2)$$

$$f(z,\bar z)=\frac{\alpha'X_0^2}{4}\int \,\,d^2z'\mathrm{exp}\,[2\omega(z',\bar z')]\mathrm{ln}\,|z-z'|^2+k$$

$$A_{S_2}^n(k,\sigma)=\Big\langle\big[e^{ik_1\cdot X(\sigma_1)}\big]_{\bf r}\big[e^{ik_2\cdot X(\sigma_2)}\big]_{\bf r}\dots\big[e^{ik_n\cdot X(\sigma_n)}\big]_{\bf r}\Big\rangle_{S_2}$$

$$J(\sigma)=\sum_{i=1}^nk_i\delta^2(\sigma-\sigma_i)$$

$$A_{S_2}^n(k,\sigma)=iC_{S_2}^X(2\pi)^d\delta^d\left(\sum_ik_i\right)\times\exp\left(-\sum_{\substack{i,j=1\\i< j}}^nk_i\cdot k_jG'\big(\sigma_i,\sigma_j\big)-\frac{1}{2}\sum_{i=1}^nk_i^2G'_{\bf r}(\sigma_i,\sigma_i)\right)$$

$$C_{S_2}^X=\mathrm{X}_0^{-d}\left(\det'\frac{-\nabla^2}{4\pi^2\alpha'}\right)^{-d/2}_{S_2}$$



$$G'_{\text{r}}(\sigma,\sigma')=G'(\sigma,\sigma')+\frac{\alpha'}{2}\ln~d^2(\sigma,\sigma')$$

$$G'_{\text{r}}(\sigma,\sigma)=2f(z,\bar z)+\alpha'\omega(z,\bar z)$$

$$A^n_{S_2}(k,\sigma) = iC^X_{S_2}(2\pi)^d \delta^d\left(\sum_i~k_i\right)\exp\left(-\frac{\alpha'}{2}\sum_i~k_i^2\omega(\sigma_i)\right)\prod_{\substack{i,j=1\\i < j}}^n|z_{ij}|^{\alpha' k_i \cdot k_j}$$

$$\left\langle \prod_{i=1}^n~\left[e^{ik_i\cdot X(z_i,\bar z_i)}\right]_{\text{r}}\prod_{j=1}^p~\partial X^{\mu_j}(z'_j)\prod_{k=1}^q~\bar\partial X^{\nu_k}(\bar z''_k)\right\rangle_{S_2}$$

$$iC^X_{S_2}(2\pi)^d \delta^d\left(\sum_i~k_i\right)\prod_{\substack{i,j=1\\i < j}}^n|z_{ij}|^{\alpha' k_i \cdot k_j}\times\left\langle \prod_{j=1}^p~\left[v^{\mu_j}(z'_j)+q^{\mu_j}(z'_j)\right]\prod_{k=1}^q~[\tilde v^{\nu_k}(\bar z''_k)+\tilde q^{\nu_k}(\bar z''_k)]\right\rangle$$

$$v^\mu(z)=-i\frac{\alpha'}{2}\sum_{i=1}^n\frac{k_i^\mu}{z-z_i}, \tilde v^\mu(\bar z)=-i\frac{\alpha'}{2}\sum_{i=1}^n\frac{k_i^\mu}{\bar z-\bar z_i}$$

$$\langle \partial X^\mu(z_1) \partial X^\nu(z_2) \rangle_{S_2}$$

$$-\frac{\alpha'\eta^{\mu\nu}}{2z_{12}^2}\langle 1\rangle_{S_2}+g(z_1,z_2)$$

$$\partial_u X^\mu = -z^2 \partial_z X^\mu$$

$$A^n_{S_2}(k,\sigma)=\left\langle \prod_{i=1}^n: e^{ik_i\cdot X(z_i,\bar z_i)}:\right\rangle_{S_2}$$

$$\left\langle \partial X^\mu(z)\prod_{i=1}^n: e^{ik_i\cdot X(z_i,\bar z_i)}:\right\rangle_{S_2}=-\frac{i\alpha'}{2}A^n_{S_2}(k,\sigma)\sum_{i=1}^n\frac{k_i^\mu}{z-z_i}+\text{ }\lambda_{\text{terms holomorphic in }z}$$

$$A^n_{S_2}(k,\sigma)\sum_{i=1}^n~k_i^\mu=0$$

$$p^\mu=\frac{1}{2\pi i}\oint_c\left(dzj_z^\mu-d\bar zj_{\bar z}^\mu\right)$$

$$-\frac{i\alpha'}{2}A^n_{S_2}(k,\sigma)\left(\frac{k_1^\mu}{z-z_1}+\sum_{i=2}^n\frac{k_i^\mu}{z_1-z_i}+O(z-z_1)\right).$$

$$ik_1\cdot\partial X(z):e^{ik_1\cdot X(z_1,\bar z_1)}:=\frac{\alpha' k_1^2}{2(z-z_1)}:e^{ik_1\cdot X(z_1,\bar z_1)}:+\partial_{z_1}:e^{ik_1\cdot X(z_1,\bar z_1)}:+O(z-z_1)$$



$$\partial_{z_1} A_{S_2}^n(k,\sigma) = \frac{\alpha'}{2} A_{S_2}^n(k,\sigma) \sum_{i=2}^n \frac{k_1 \cdot k_i}{z_{1i}}.$$

$$A_{S_2}^n(k,\sigma) \propto \delta^d \left(\sum_i k_i \right) \prod_{\substack{i,j=1 \\ i < j}}^n |z_{ij}|^{\alpha' k_i \cdot k_j},$$

$$G'(\sigma_1, \sigma_2) = -\frac{\alpha'}{2} \ln |z_1 - z_2|^2 - \frac{\alpha'}{2} \ln |z_1 - \bar{z}_2|^2,$$

$$\left\langle \prod_{i=1}^n :e^{ik_i \cdot X(z_i, \bar{z}_i)}:\right\rangle_{D_2}$$

$$= i C_{D_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \prod_{i=1}^n |z_i - \bar{z}_i|^{\alpha' k_i^2 / 2} \times \prod_{\substack{i,j=1 \\ i < j}}^n |z_i - z_j|^{\alpha' k_i \cdot k_j} |z_i - \bar{z}_j|^{\alpha' k_i \cdot k_j}$$

$${}_\star^* X^\mu(y_1) X^\nu(y_2) {}_\star^* = X^\mu(y_1) X^\nu(y_2) + 2\alpha' \eta^{\mu\nu} \ln |y_1 - y_2|,$$

$$\left\langle \prod_{i=1}^n {}_\star^* e^{ik_i \cdot X(y_i)} \right\rangle_{D_2} = i C_{D_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \prod_{\substack{i,j=1 \\ i < j}}^n |y_i - y_j|^{2\alpha' k_i \cdot k_j}$$

$$\left\langle \prod_{i=1}^n {}_\star^* e^{ik_i \cdot X(y_i)} \prod_{j=1}^p \partial_y X^{\mu_j}(y'_j) \right\rangle_{D_2}$$

$$= i C_{D_2}^X (2\pi)^d \delta^d \left(\sum_i k_i \right) \times \prod_{i,j=1}^n |y_{ij}|^{2\alpha' k_i \cdot k_j} \left\langle \prod_{j=1}^p [v^{\mu_j}(y'_j) + q^{\mu_j}(y'_j)] \right\rangle_{D_2}$$

$$v^\mu(y) = -2i\alpha' \sum_{i=1}^n \frac{k_i^\mu}{y - y_i}$$

$$G'(\sigma_1, \sigma_2) = -\frac{\alpha'}{2} \ln |z_1 - z_2|^2 - \frac{\alpha'}{2} \ln |1 + z_1 \bar{z}_2|^2.$$



$$\left\langle \prod_{i=1}^n :e^{ik_i\cdot X(z_i,\bar{z}_i)}:\right\rangle_{RP_2}$$

$$= i C_{RP_2}^X (2\pi)^d \delta^d \biggl(\sum_i ~ k_i \biggr) \prod_{i=1}^n ~ |1 + z_i \bar{z}_i|^{\alpha' k_i^2/2}$$

$$\times \prod_{\stackrel{i,j=1}{i < j}}^n |z_i - z_j|^{\alpha' k_i \cdot k_j} |1 + z_i \bar{z}_j|^{\alpha' k_i \cdot k_j}$$

$${\bf M\'etrica\;CFT}.$$

$$\langle c(z_1)c(z_2)c(z_3)\tilde{c}(\bar{z}_4)\tilde{c}(\bar{z}_5)\tilde{c}(\bar{z}_6)\rangle_{S_2}$$

$$\det\! C_{0j}^\alpha(\sigma_i)$$

$$\begin{array}{l} {\mathsf C}^z = 1,z,z^2,{\mathsf C}^{\bar{z}} = 0 \\ {\mathsf C}^z = 0,{\mathsf C}^{\bar{z}} = 1,\bar{z},\bar{z}^2 \end{array}$$

$$C_{S_2}^g \det \begin{vmatrix} 1 & 1 & 1 \\ z_1 & z_2 & z_3 \\ z_1^2 & z_2^2 & z_3^2 \end{vmatrix} \det \begin{vmatrix} 1 & 1 & 1 \\ \bar{z}_4 & \bar{z}_5 & \bar{z}_6 \\ \bar{z}_4^2 & \bar{z}_5^2 & \bar{z}_6^2 \end{vmatrix} = C_{S_2}^g z_{12} z_{13} z_{23} \bar{z}_{45} \bar{z}_{46} \bar{z}_{56}$$

$$\left\langle \prod_{i=1}^{p+3}~c(z_i) \prod_{j=1}^p~b(z'_j) \cdot (\text{anti}) \right\rangle_{S_2} = C_{S_2}^g \frac{z_{p+1,p+2} z_{p+1,p+3} z_{p+2,p+3}}{(z_1-z'_1) \dots (z_p-z'_p)} \cdot (\beth_{\text{anti}}) \pm \boxtimes_{\text{permutations}}$$

$$\oint ~c \frac{dz}{2\pi i} j_z = - \oint ~c \frac{du}{2\pi i} j_u + 3 \rightarrow 3$$

$$z_{12} z_{13} z_{23} \bar{z}_{45} \bar{z}_{46} \bar{z}_{56} F(z_1,z_2,z_3) \tilde{F}(\bar{z}_4,\bar{z}_5,\bar{z}_6)$$

$$C_{S_2}^g \prod_{\stackrel{i,i'=1}{i < i'}}^{p+3} z_{ii'} \prod_{\stackrel{j,j'=1}{j < j'}}^p z'_{jj'} \prod_{i=1}^{p+3} \prod_{j=1}^p \left(z_i - z'_j\right)^{-1} \cdot (\text{anti}),$$

$$\tilde{b}(\bar{z})=b(z'),\tilde{c}(\bar{z})=c(z'),z'=\bar{z},\mathrm{Im}z>0$$

$$\langle c(z_1)c(z_2)c(z_3)\rangle_{D_2}=C_{D_2}^g z_{12} z_{13} z_{23}$$

$$\langle c(z_1)c(z_2)\tilde{c}(\bar{z}_3)\rangle_{D_2}=\langle c(z_1)c(z_2)c(z'_3)\rangle_{D_2}=C_{D_2}^g z_{12}(z_1-\bar{z}_3)(z_2-\bar{z}_3)$$

$$\tilde{b}(\bar{z})=\left(\frac{\partial z'}{\partial \bar{z}}\right)^2 b(z')=z'^4 b(z'), \tilde{c}(\bar{z})=\left(\frac{\partial z'}{\partial \bar{z}}\right)^{-1} c(z')=z'^{-2} c(z')$$

$$\langle c(z_1)c(z_2)c(z_3)\rangle_{RP_2}=C_{RP_2}^g z_{12} z_{13} z_{23}$$



$$\langle c(z_1)c(z_2)\tilde{c}(\bar{z}_3) \rangle_{RP_2} = z_3'^{-2}\langle c(z_1)c(z_2)c(z_3') \rangle_{RP_2} = C_{RP_2}^g z_{12}(1+z_1\bar{z}_3)(1+z_2\bar{z}_3)$$

$$S_{D_2}(k_1; k_2; k_3) = g_0^3 e^{-\lambda} \left\langle {}^{\star}_{\star}c^1 e^{ik_1 \cdot X}(y_1) {}^{\star}_{\star\star}c^1 e^{ik_2 \cdot X}(y_2) {}^{\star}_{\star\star}c^1 e^{ik_3 \cdot X}(y_3) \right\rangle_{D_2} + (k_2 \leftrightarrow k_3)$$

$$S_{D_2}(k_1; k_2; k_3) = ig_0^3 C_{D_2} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) \times |y_{12}|^{1+2\alpha' k_1 \cdot k_2} |y_{13}|^{1+2\alpha' k_1 \cdot k_3} |y_{23}|^{1+2\alpha' k_2 \cdot k_3} \\ + (k_2 \leftrightarrow k_3)$$

$$2\alpha' k_1 \cdot k_2 = \alpha'(k_3^2 - k_1^2 - k_2^2) = -1$$

$$S_{D_2}(k_1; k_2; k_3) = 2ig_0^3 C_{D_2} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right)$$

$$S_{D_2}(k_1; k_2; k_3; k_4) = g_0^4 e^{-\lambda} \int_{-\infty}^{\infty} dy_4 \left\langle \prod_{i=1}^3 {}^{\star}_{\star}c^1(y_i) e^{ik_i \cdot X(y_i)} {}^{\star\star}e^{ik_4 \cdot X(y_4)} \right\rangle + (k_2 \leftrightarrow k_3) \\ = ig_0^4 C_{D_2} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) |y_{12}y_{13}y_{23}| \int_{-\infty}^{\infty} dy_4 \prod_{i < j} |y_{ij}|^{2\alpha' k_i \cdot k_j} + (k_2 \leftrightarrow k_3)$$

$$s = -(k_1 + k_2)^2, t = -(k_1 + k_3)^2, u = -(k_1 + k_4)^2.$$

$$s+t+u = \sum_i m_i^2 = -\frac{4}{\alpha'}$$

$$S_{D_2}(k_1; k_2; k_3; k_4) = ig_0^4 C_{D_2} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) \times \left[\int_{-\infty}^{\infty} dy_4 |y_4|^{-\alpha' u - 2} |1 - y_4|^{-\alpha' t - 2} + (t \rightarrow s) \right]$$

$$S_{D_2}(k_1; k_2; k_3; k_4) = 2ig_0^4 C_{D_2} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) [I(s, t) + I(t, u) + I(u, s)]$$

$$I(s, t) = \int_0^1 dy y^{-\alpha' s - 2} (1 - y)^{-\alpha' t - 2}$$

$$I(s, t) = \int_0^r dy y^{-\alpha' s - 2} + \underset{\text{terms analytic at } \alpha' s}{\varsigma} = -1 = -\frac{r^{-\alpha' s - 1}}{\alpha' s + 1} + \underset{\text{terms analytic at } \alpha' s}{\varsigma} = -1$$

$$= -\frac{1}{\alpha' s + 1} + \underset{\text{terms analytic at } \alpha' s}{\varsigma} = -1$$

$$\frac{1}{\alpha' s + 1} \equiv \frac{1}{\alpha' s + 1 + i\epsilon} \equiv P \frac{1}{\alpha' s + 1} - i\pi\delta(\alpha' s + 1)$$



$$S_{D_2}(k_1; k_2; k_3; k_4) = i \int \frac{d^{26}k}{(2\pi)^{26}} \frac{S_{D_2}(k_1; k_2; k) S_{D_2}(-k; k_3; k_4)}{-k^2 + \alpha'^{-1} + i\epsilon} + \text{Q terms analytic at } k^2 = 1/\alpha'$$

$$C_{D_2} = e^{-\lambda} C_{D_2}^X C_{D_2}^{\text{g}} = \frac{1}{\alpha' g_0^2}$$

$$S_{D_2}(k_1; k_2; k_3) = \frac{2ig_o}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right)$$

$$I(s,t) = \int_0^r dy [y^{-\alpha's-2} + (\alpha't+2)y^{-\alpha's-1} + \dots]$$

$$I(s,t) = \frac{u-t}{2s} + \text{ terms analytic at } \alpha's = 0$$

$$\alpha's = -1, 0, 1, 2, \dots$$

$$B(a,b) = \int_0^1 dy y^{a-1} (1-y)^{b-1}$$

$$I(s,t) = B(-\alpha_0(s), -\alpha_0(t)), \alpha_0(x) = 1 + \alpha'x$$

$$w^{a+b-1}B(a,b) = \int_0^w dv v^{a-1} (w-v)^{b-1}$$

$$\Gamma(a+b)B(a,b) = \int_0^\infty dv v^{a-1} e^{-v} \int_0^\infty d(w-v) (w-v)^{b-1} e^{-(w-v)} = \Gamma(a)\Gamma(b)$$

$$S_{D_2}(k_1; k_2; k_3; k_4)$$

$$= \frac{2ig_0^2}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right)$$

$$\times [B(-\alpha_0(s), -\alpha_0(t)) + B(-\alpha_0(s), -\alpha_0(u)) + B(-\alpha_0(t), -\alpha_0(u))]$$

$$B(-\alpha_0(x), -\alpha_0(y)) = \frac{\Gamma(-\alpha'x-1)\Gamma(-\alpha'y-1)}{\Gamma(-\alpha'x-\alpha'y-2)}$$

$$s \rightarrow \infty, t \text{ fixed}$$

$$s \rightarrow \infty, t/s \text{ fixed}$$

$$s = E^2, t = (4m^2 - E^2)\sin^2 \frac{\theta}{2}, u = (4m^2 - E^2)\cos^2 \frac{\theta}{2}$$

$$S_{D_2}(k_1; k_2; k_3; k_4) \propto s^{\alpha_0(t)} \Gamma(-\alpha_0(t))$$

$$S_{D_2}(k_1; k_2; k_3; k_4) \approx \exp [-\alpha'(s \ln s \alpha' + t \ln t \alpha' + u \ln u \alpha')] = \exp [-\alpha' s f(\theta)]$$



$$f(\theta) \approx -\sin^2 \frac{\theta}{2} \ln \sin^2 \frac{\theta}{2} - \cos^2 \frac{\theta}{2} \ln \cos^2 \frac{\theta}{2}$$

Factores Chan-Paton e interacciones de gauge.

$$|N; k; ij\rangle,$$

$$\text{Tr}(\lambda^a \lambda^b) = \delta^{ab},$$

$$|N; k; a\rangle = \sum_{i,j=1}^n |N; k; ij\rangle \lambda_{ij}^a$$

$$\text{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3} \lambda^{a_4})$$

$$S_{D_2}(k_1, a_1; k_2, a_2; k_3, a_3) = \frac{i g_o}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right) \text{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3} + \lambda^{a_1} \lambda^{a_3} \lambda^{a_2})$$

$$S_{D_2}(k_1, a_1; k_2, a_2; k_3, a_3; k_4, a_4)$$

$$= \frac{i g_o^2}{\alpha'} (2\pi)^{26} \delta^{26} \left(\sum_i k_i \right)$$

$$\begin{aligned} & \times [\text{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_4} \lambda^{a_3} + \lambda^{a_1} \lambda^{a_3} \lambda^{a_4} \lambda^{a_2}) B(-\alpha_o(s), -\alpha_o(t)) \\ & + \text{Tr}(\lambda^{a_1} \lambda^{a_3} \lambda^{a_2} \lambda^{a_4} \\ & + \lambda^{a_1} \lambda^{a_4} \lambda^{a_2} \lambda^{a_3}) B(-\alpha_o(t), -\alpha_o(u)) + \text{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3} \lambda^{a_4} \\ & + \lambda^{a_1} \lambda^{a_4} \lambda^{a_3} \lambda^{a_2}) B(-\alpha_o(s), -\alpha_o(u))]. \end{aligned}$$

$$\frac{1}{4} \text{Tr}(\{\lambda^{a_1}, \lambda^{a_2}\} \{\lambda^{a_3}, \lambda^{a_4}\})$$

$$\frac{1}{4} \sum_a \text{Tr}(\{\lambda^{a_1}, \lambda^{a_2}\} \lambda^a) \text{Tr}(\{\lambda^{a_3}, \lambda^{a_4}\} \lambda^a)$$

$$\text{Tr}(A \lambda^a) \text{Tr}(B \lambda^a) = \text{Tr}(AB)$$

$$S_{D_2}(k_1, a_1, e_1; k_2, a_2; k_3, a_3)$$

$$= -i g'_o g_o^2 e^{-\lambda} e_{1\mu}$$

$$\times \langle {}^* c^1 \dot{X}^\mu e^{ik_1 \cdot X} (y_1) {}^*_\star {}^* c^1 e^{ik_2 \cdot X} (y_2) {}^*_\star {}^* c^1 e^{ik_3 \cdot X} (y_3) {}^*_\star \rangle_{D_2} \text{Tr}(\lambda^{a_1} \lambda^{a_2} \lambda^{a_3}) + (k_2, a_2)$$

$$\leftrightarrow (k_3, a_3)$$



$$\langle \stackrel{\star}{\ast}X^\mu e^{ik_1\cdot X}(y_1)^\star_{\star\star}e^{ik_2\cdot X}(y_2)^\star_{\star\star}e^{ik_3\cdot X}(y_3)^\star_{\star}\rangle_{D_2}$$

$$= -2i\alpha'\left(\frac{k_2^{\mu}}{y_{12}}+\frac{k_3^{\mu}}{y_{13}}\right)$$

$$\times\, iC_{D_2}^X(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr)|y_{12}|^{2\alpha'k_1\cdot k_2}|y_{13}|^{2\alpha'k_1\cdot k_3}|y_{23}|^{2\alpha'k_2\cdot k_3}$$

$$S_{D_2}(k_1,a_1,e_1;k_2,a_2,e_2;k_3,a_3)= -ig'_oe_1\cdot k_{23}(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr)\text{Tr}(\lambda^{a_1}[\lambda^{a_2},\lambda^{a_3}])$$

$$\text{Tr}([\lambda^{a_1},\lambda^{a_2}][\lambda^{a_3},\lambda^{a_4}])$$

$$g_o'=(2\alpha')^{-1/2}g_o$$

$$S_{D_2}(k_1,a_1,e_1;k_2,a_2,e_2;k_3,a_3,e_3)$$

$$=ig'_o(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr)(e_1\cdot k_{23}e_2\cdot e_3+e_2\cdot k_{31}e_3\cdot e_1+e_3\cdot k_{12}e_1\cdot e_2+\frac{\alpha'}{2}e_1\\ \cdot k_{23}e_2\cdot k_{31}e_3\cdot k_{12}\biggr)\text{Tr}(\lambda^{a_1}[\lambda^{a_2},\lambda^{a_3}])$$

$${\cal S}=\frac{1}{g_o'^2}\int\,\,d^{26}x\left[-\frac{1}{2}\text{Tr}\big(D_\mu\varphi D^\mu\varphi\big)+\frac{1}{2\alpha'}\text{Tr}(\varphi^2)+\frac{2^{1/2}}{3\alpha'^{1/2}}\text{Tr}(\varphi^3)-\frac{1}{4}\text{Tr}\big(F_{\mu\nu}F^{\mu\nu}\big)\right]$$

$$-\frac{2i\alpha'}{3g_o'^2}\text{Tr}\big(F_\mu^{\,\,\,v}F_v^{\,\,\,\omega}F_\omega^{\,\,\,\mu}\big)$$

$$\lambda^a \rightarrow U \lambda^a U^\dagger,$$

$$\Omega\alpha_n^\mu\Omega^{-1}=(-1)^n\alpha_n^\mu$$

$$\Omega\alpha_n^\mu\Omega^{-1}=\tilde\alpha_n^\mu$$

$$\Omega|N;k\rangle=\omega_N|N;k\rangle, \omega_N=(-1)^{1+\alpha'm^2}$$

$$\Omega|N;k;ij\rangle=\omega_N|N;k;ji\rangle.$$

$$\Omega|N;k;a\rangle=\omega_N s^a|N;k;a\rangle$$

$$\begin{array}{c} \alpha'm^2 \text{ even: } \lambda^a \text{ antisymmetric} \\ \alpha'm^2 \text{ odd: } \lambda^a \text{ symmetric} \end{array}$$

$$\Omega_\gamma|N;k;ij\rangle=\omega_N\gamma_{jj'}|N;k;j'i'\rangle\gamma_{i'i}^{-1}$$

$$\Omega_\gamma^2|N;k;ij\rangle=[(\gamma^T)^{-1}\gamma]_{ii'}|N;k;i'j'\rangle(\gamma^{-1}\gamma^T)_{j'j}.$$



$$\gamma^T=\pm\gamma$$

$$|N;k;ij\rangle'=U_{ii'}^{-1}|N;k;i'j'\rangle U_{j'j},$$

$$\gamma' = U^T \gamma U.$$

$$\gamma=M\equiv i\begin{bmatrix}0&I\\-I&0\end{bmatrix}$$

$$\begin{array}{ll}\alpha'm^2\text{ even: }M(\lambda^a)^TM=-\lambda^a\\\alpha'm^2\text{ odd: }M(\lambda^a)^TM=+\lambda^a\end{array}$$

$$\lambda\rightarrow (\gamma^T)^{-1}\gamma\lambda\gamma^{-1}\gamma^T=\lambda,$$

$$S_{S_2}(k_1;k_2;k_3)=g_{\rm c}^3e^{-2\lambda}\left\langle\prod_{i=1}^3:\tilde{c}ce^{ik_i\cdot X}(z_i,\bar{z}_i):\right\rangle_{S_2}$$

$$S_{S_2}(k_1;k_2;k_3)=ig_{\rm c}^3C_{S_2}(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr)$$

$$S_{S_2}(k_1;k_2;k_3;k_4)=g_{\rm c}^4e^{-2\lambda}\int_{\bf C}d^2z_4\left\langle\prod_{i=1}^3:\tilde{c}ce^{ik_i\cdot X}(z_i,\bar{z}_i):\right.:e^{ik_4\cdot X}(z_4,\bar{z}_4):\biggr\rangle_{S_2}$$

$$S_{S_2}(k_1;k_2;k_3;k_4)=ig_{\rm c}^4C_{S_2}(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr)J(s,t,u)$$

$$J(s,t,u)=\int_{\bf C}d^2z_4|z_4|^{-\alpha'u/2-4}|1-z_4|^{-\alpha't/2-4}$$

$$\alpha' s, \alpha' t, \alpha' u = -4, 0, 4, 8, \ldots,$$

$$ig_{\rm c}^4C_{S_2}\int_{|z_4|>1/\epsilon}d^2z_4|z_4|^{\alpha's/2}\sim-\frac{8\pi ig_{\rm c}^4C_{S_2}}{\alpha's+4}$$

$$C_{S_2}=\frac{8\pi}{\alpha' g_{\rm c}^2}$$

$$S_{S_2}(k_1;k_2;k_3)=\frac{8\pi ig_{\rm c}}{\alpha'}(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr).$$

$$S_{S_2}(k_1;k_2;k_3;k_4)=\frac{8\pi ig_{\rm c}^2}{\alpha'}(2\pi)^{26}\delta^{26}\biggl(\sum_ik_i\biggr)C(-\alpha_{\rm c}(t),-\alpha_{\rm c}(u)),$$

$$C(a,b)=\int_{\bf C}d^2z|z|^{2a-2}|1-z|^{2b-2}=2\pi\frac{\Gamma(a)\Gamma(b)\Gamma(c)}{\Gamma(a+b)\Gamma(a+c)\Gamma(b+c)},a+b+c=1$$



$$S_{S_2}(k_1;k_2;k_3;k_4) \propto s^{2\alpha_c(t)}\frac{\Gamma(-\alpha_c(t))}{\Gamma(1+\alpha_c(t))},$$

$$S_{S_2}(k_1;k_2;k_3;k_4)\propto \exp\left[-\frac{\alpha'}{2}(\operatorname{sln}\,s\alpha'+t\ln\,t\alpha'+u\ln\,u\alpha')\right]$$

$$S_{S_2}(k_1,e_1;k_2;k_3)=g^2_{\rm c}g'_{\rm c}e^{-2\lambda}e_{1\mu\nu}\langle :\tilde cc\partial X^\mu\bar\partial X^\nu e^{ik_1\cdot X}(z_1,\bar z_1)\langle :\tilde cc e^{ik_2\cdot X}(z_2,\bar z_2)\rangle :\tilde cc e^{ik_3\cdot X}(z_3,\bar z_3)\rangle\rangle_{S_2}$$

$$= -\frac{\pi i \alpha'}{2} g'_{\rm c} e_{1\mu\nu} k^{\mu}_{23} k^{\nu}_{23} (2\pi)^{26} \delta^{26} \Biggl(\sum_i \; k_i \Biggr)$$

$$g'_{\rm c} = \frac{2}{\alpha'} g_{\rm c}$$

$$S_T=-\frac{1}{2}\int\;d^{26}x (-G)^{1/2}e^{-2\tilde{\Phi}}\left(G^{\mu\nu}\partial_\mu T\partial_\nu T-\frac{4}{\alpha'}T^2\right)$$

$$\tilde G_{\mu\nu}=\eta_{\mu\nu}-2\kappa e_{\mu\nu}e^{ik\cdot x}$$

$$\kappa=\pi\alpha'g'_{\rm c}=2\pi g_{\rm c}$$

$$S_{S_2}(k_1,e_1;k_2,e_2;k_3,e_3)=\frac{i\kappa}{2}(2\pi)^{26}\delta^{26}\left(\sum_i \; k_i \right)e_{1\mu\nu}e_{2\alpha\beta}e_{3\gamma\delta}T^{\mu\alpha\gamma}T^{\nu\beta\delta}$$

$$T^{\mu\alpha\gamma}=k^{\mu}_{23}\eta^{\alpha\gamma}+k^{\alpha}_{31}\eta^{\gamma\mu}+k^{\gamma}_{12}\eta^{\mu\alpha}+\frac{\alpha'}{8}k^{\mu}_{23}k^{\alpha}_{31}k^{\gamma}_{12}$$

$$J(s,t,u,\alpha')=-2\sin\,\pi\alpha_c(t)I(s,t,4\alpha')I(t,u,4\alpha')$$

$$\int_{\mathsf C} d^2zz^{a-1+m_1}\bar z^{a-1+n_1}(1-z)^{b-1+m_2}(1-\bar z)^{b-1+n_2}$$

$$=2\sin\,[\pi(b+n_2)]B(a+m_1,b+m_2)B(b+n_2,1-a-b-n_1-n_2)m_1-n_1$$

$$\in {\bf Z}, m_2-n_2 \in {\bf Z}$$

$$A_{\rm c}(s,t,u,\alpha',g_{\rm c})=\frac{\pi i g_{\rm c}^2\alpha'}{g_0^4}\sin\,[\pi\alpha_{\rm c}(t)]A_{\rm o}\left(s,t,\frac{1}{4}\alpha',g_{\rm o}\right)A_{\rm o}\left(t,u,\frac{1}{4}\alpha',g_{\rm o}\right)^*$$

$$:e^{ik_1\cdot X(z_1,\bar z_1)}\langle :e^{ik_4\cdot X(z_4,\bar z_4)}\cdot$$

$$=|z_{14}|^{\alpha'k_1\cdot k_4}\langle\big(1+iz_{14}k_1\cdot\partial X+i\bar z_{14}k_1\cdot\bar\partial X-z_{14}\bar z_{14}k_1\cdot\partial Xk_1\cdot\bar\partial X\\+\cdots\big)e^{i(k_1+k_4)\cdot X}(z_4,\bar z_4)\rangle$$

$$\alpha'k_1\cdot k_4=\frac{\alpha'}{2}(k_1+k_4)^2-4>-2$$

$$-\frac{1}{4g_0'^2}\int\;d^{26}x{\rm Tr}\big(F_{\mu\nu}F^{\mu\nu}\big)$$



$$-\frac{1}{4g_o'^2}\int\;d^{26}x(-G)^{1/2}e^{-\tilde{\Phi}}\text{Tr}\big(F_{\mu\nu}F^{\mu\nu}\big)$$

$$-\Lambda \int\;d^{26}x(-G)^{1/2}e^{-\tilde{\Phi}}$$

Invariancia Möbius.

$$z'=\frac{\alpha z+\beta}{\gamma z+\delta}$$

$$\langle \mathcal{A}_i(z_1,\bar{z}_1) \dots \mathcal{A}_k(z_n,\bar{z}_n) \rangle_{S_2} = \langle \mathcal{A}'_i(z_1,\bar{z}_1) \dots \mathcal{A}'_k(z_n,\bar{z}_n) \rangle_{S_2}.$$

$$\langle \mathcal{A}_i(0,0) \rangle_{S_2} = \langle \mathcal{A}'_i(0,0) \rangle_{S_2} = \gamma^{-h_i} \bar{\gamma}^{-\tilde{h}_i} \langle \mathcal{A}_i(0,0) \rangle_{S_2}.$$

$$\langle \mathcal{A}_i(z_1,\bar{z}_1)\mathcal{A}_j(z_2,\bar{z}_2) \rangle_{S_2} = z_{12}^{-h_i-h_j} \bar{z}_{12}^{-\tilde{h}_i-\tilde{h}_j} \langle \mathcal{A}_i(1,1)\mathcal{A}_j(0,0) \rangle_{S_2},$$

$$\left\langle \mathcal{O}_p(z_1,\bar{z}_1)\mathcal{O}_q(z_2,\bar{z}_2) \right\rangle_{S_2} = 0 \text{ unless } h_p = h_q, \tilde{h}_p = \tilde{h}_q.$$

$$\left\langle \prod_{i=1}^3 \mathcal{O}_{p_i}(z_i,\bar{z}_i) \right\rangle_{S_2} = C_{p_1 p_2 p_3} \prod_{\substack{i,j=1 \\ i < j}}^3 z_{ij}^{h-2(h_i+h_j)} z_{ij}^{\tilde{h}-2(\tilde{h}_i+\tilde{h}_j)}$$

$$\left\langle \prod_{i=1}^4 \mathcal{O}_{p_i}(z_i,\bar{z}_i) \right\rangle_{S_2} = C_{p_1 p_2 p_3 p_4} (z_c,\bar{z}_c) (z_{12} z_{34})^h (\bar{z}_{12} \bar{z}_{34})^{\tilde{h}} \times \prod_{\substack{i,j=1 \\ i < j}}^4 z_{ij}^{-h_i-h_j} \bar{z}_{ij}^{-\tilde{h}_i-\tilde{h}_j}$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_j(0,0) \rangle_{S_2}.$$

$$\int\;[d\phi_{\text{b}}]\Psi_{\mathcal{A}_i}[\phi_{\text{b}}^{\Omega}]\Psi_{\mathcal{A}_j}[\phi_{\text{b}}]$$

$$\left\langle \langle i \mid j \rangle = \langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_j(0,0) \rangle_{S_2} .\right.$$

$$\mathcal{O}_i(z,\bar{z})\mathcal{O}_j(0,0) = \frac{\langle \langle i \mid j \rangle}{z^{h_i+h_j} \tilde{Z}_i + \tilde{h}_j \langle 1 \rangle_{S_2}} + \cdots$$

$$\langle \langle i \mid j \rangle = \pm \langle \langle j \mid i \rangle,$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_k(1,1)\mathcal{A}_l(0,0) \rangle_{S_2} = \langle \langle i | \hat{\mathcal{A}}_k(1,1) \mid j \rangle,$$

$$\sum_l\; c_{kj}^l \langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_l(0,0) \rangle_{S_2} = c_{ikj}$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_k(z_1,\bar{z}_1)\mathcal{A}_l(0,0) \rangle_{S_2} = z_1^{h_i-h_k-h_j} \bar{z}_1^{\tilde{h}_i-\tilde{h}_k-\tilde{h}_j} c_{ikj}$$

$$\langle \mathcal{A}'_i(\infty,\infty)\mathcal{A}_k(z_1,\bar{z}_1)\mathcal{A}_l(z_2,\bar{z}_2)\mathcal{A}_j(0,0) \rangle_{S_2} = \langle \langle i | \text{T}[\hat{\mathcal{A}}_k(z_1,\bar{z}_1)\hat{\mathcal{A}}_l(z_2,\bar{z}_2)] \mid j \rangle$$



$$1=|m\rangle\mathcal{G}^{mn}\langle\langle n|.$$

$$\sum_m z_1^{h_i-h_k-h_m}\bar{z}_1^{\tilde{h}_i-\tilde{h}_k-\tilde{h}_m}z_2^{h_m-h_l-h_j}\bar{z}_2^{\tilde{h}_m-\tilde{h}_l-\tilde{h}_j}c_{ikm}c^m{}_{lj}$$

Cálculos de operador.

$$\left\langle :e^{ik_4\cdot X(\infty,\infty)}:':e^{ik_1\cdot X(z_1,\bar{z}_1)}: :e^{ik_2\cdot X(z_2,\bar{z}_2)}: :e^{ik_3\cdot X(0,0)}:\right\rangle_{S_2}=\left\langle \langle 0;k_4|\mathrm{T}\left[\stackrel{\circ}{e}^{ik_1\cdot X_1\circ}\stackrel{\circ}{e}^{ik_2\cdot X_2\circ}\right]|0;k_3\rangle\right.$$

$$e^{ik\cdot X_\circ}=e^{ik\cdot X_C}e^{ik\cdot X_A},$$

$$\begin{aligned}X_C^\mu(z,\bar{z})&=x^\mu-i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{m=1}^{\infty}\frac{1}{m}\left(\alpha_{-m}^\mu z^m+\tilde{\alpha}_{-m}^\mu\bar{z}^m\right)\\X_A^\mu(z,\bar{z})&=-i\frac{\alpha'}{2}p^\mu\ln|z|^2+i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{m=1}^{\infty}\frac{1}{m}\left(\frac{\alpha_m^\mu}{z^m}+\frac{\tilde{\alpha}_m^\mu}{\bar{z}^m}\right).\end{aligned}$$

$$\langle 0;k_4|e^{ik_1\cdot X_{1C}}e^{ik_1\cdot X_{1A}}e^{ik_2\cdot X_{2C}}e^{ik_2\cdot X_{2A}}|0;k_3\rangle.$$

$$e^{ik_1\cdot X_{1A}}e^{ik_2\cdot X_{2C}}=e^{ik_2\cdot X_{2C}}e^{ik_1\cdot X_{1A}}e^{-[k_1\cdot X_{1A},k_2\cdot X_{2C}]}=e^{ik_2\cdot X_{2C}}e^{ik_1\cdot X_{1A}}|z_{12}|^{\alpha' k_1\cdot k_2}$$

$$\begin{aligned}&|z_{12}|^{\alpha'}k_1\cdot k_2\langle 0;k_4|e^{ik_1\cdot X_{1C}+ik_2\cdot X_{2C}}e^{ik_1\cdot X_{1A}+ik_2\cdot X_{2A}}|0;k_3\rangle\\&=|z_{12}|^{\alpha' k_1\cdot k_2}\left\langle \langle 0;k_4|e^{i(k_1+k_2)\cdot x}e^{\alpha'(k_1\ln|z_1|+k_2\ln|z_2|)\cdot p}|0;k_3\rangle\right.\\&=|z_{12}|^{\alpha' k_1\cdot k_2}|z_1|^{\alpha' k_1\cdot k_3}|z_2|^{\alpha' k_2\cdot k_3}\langle \langle 0;k_1+k_2+k_4|0;k_3\rangle\\&=iC_{S_2}^X(2\pi)^d\delta^d\left(\sum_i k_i\right)|z_{12}|^{\alpha' k_1\cdot k_2}|z_1|^{\alpha' k_1\cdot k_3}|z_2|^{\alpha' k_2\cdot k_3}\end{aligned}$$

$$\langle 0;k|0;l\rangle=iC_{S_2}^X(2\pi)^d\delta^d(k+l)$$

$$\langle 0;k|0;l\rangle=(2\pi)^d\delta^d(l-k)$$

$$\langle \langle 0;k|=iC_{S_2}^X\langle 0;-k|.$$

$$\overline{\mathcal{A}(p)}=\mathcal{A}'(p')^\dagger.$$

$$\mathcal{O}(z)=i^h\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_n}{z^{n+h}}$$

$$\mathcal{O}(z)^\dagger=i^{-h}\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_n^\dagger}{\bar{z}^{n+h}}.$$

$$\overline{\mathcal{O}(z)}=i^{-h}(-z^{-2})^h\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_n^\dagger}{z^{-n-h}}=i^h\sum_{n=-\infty}^{\infty}\frac{\mathcal{O}_{-n}^\dagger}{z^{n+h}}.$$



$$\left\langle \overline{\mathcal{A}_i} \right| = K \langle \mathcal{A}_i |$$

$$\langle \bar{i} \mid j \rangle \equiv \left\langle \langle 1 | \overline{\mathcal{A}'_i(\infty,\infty)} \mid j \right\rangle = K \langle 1 | \overline{\mathcal{A}'_i(\infty,\infty)} | j \rangle = K \langle 1 | \mathcal{A}_i(0,0)^\dagger | j \rangle = K \langle j | \mathcal{A}_i(0,0) | 1 \rangle^* = K \langle j \mid i \rangle^* = K \langle i \mid j \rangle$$

$$k_1^ik_1^jP_{ij,kl}^Jk_3^kk_3^l$$

$$|z|^{-2a}=\frac{1}{\Gamma(a)}\int_0^{\infty}dt t^{a-1}\exp{(-tz\bar{z})}$$

$$w\cong w+2\pi\cong w+2\pi\tau$$

$$(\sigma^1,\sigma^2)\cong (\sigma^1+2\pi,\sigma^2)\cong (\sigma^1+2\pi\tau_1,\sigma^2+2\pi\tau_2)$$

$$z\cong z\exp{(-2\pi i\tau)}$$

$$1\leq |z|\leq \exp{(2\pi\tau_2)}$$

$$0\leq \text{Re} w\leq \pi, w\cong w+2\pi it$$

$$w'=-\bar{w}$$

$$(\sigma^1,0)\cong (\sigma^1,2\pi t)$$

$$w\cong w+2\pi\cong -\bar{w}+2\pi it$$

$$(\sigma^1,\sigma^2)\cong (\sigma^1+2\pi,\sigma^2)\cong (-\sigma^1,\sigma^2+2\pi t)$$

$$w'=-\bar{w}+2\pi it$$

$$0\leq \text{Re} w\leq \pi, w\cong -\bar{w}+\pi+2\pi it$$

$$w'=-\bar{w} \text{ and } w'=w+\pi(2it+1)$$

Correlaciones escalares.

$$\frac{2}{\alpha'}\bar{\partial}\partial G'(w,\bar{w};w',\bar{w}')=-2\pi\delta^2(w-w')+\frac{1}{4\pi\tau_2}.$$

$$G'(w,\bar{w};w',\bar{w}')\sim-\frac{\alpha'}{2}\ln\left|\vartheta_1\left(\frac{w-w'}{2\pi}\middle|\tau\right)\right|^2.$$

$$G'(w,\bar{w};w',\bar{w}')=-\frac{\alpha'}{2}\ln\left|\vartheta_1\left(\frac{w-w'}{2\pi}\middle|\tau\right)\right|^2+\alpha'\frac{[\text{Im}(w-w')]^2}{4\pi\tau_2}+k(\tau,\bar{\tau}).$$



$$\left\langle \prod_{i=1}^n :e^{ik_i\cdot X(z_i,\bar{z}_i)}:\right\rangle_{T^2}$$

$$= i C_{T^2}^X(\tau) (2\pi)^d \delta^d \left(\sum_i ~ k_i \right)$$

$$\times \prod_{i < j} \left| \frac{2\pi}{\partial_\nu \vartheta_1(0 \mid \tau)} \vartheta_1\left(\frac{w_{ij}}{2\pi} \Big| \tau\right) \exp\left[-\frac{\left({\rm Im} w_{ij}\right)^2}{4\pi \tau_2}\right] \right|^{\alpha' k_i \cdot k_j}$$

$$Z(\tau)=\text{Tr}[\exp{(2\pi i\tau_1 P - 2\pi \tau_2 H)}]=(q\bar q)^{-d/24}\text{Tr}(q^{L_0}\bar q^{\tilde L_0})$$

$$V_d(q\bar q)^{-d/24}\int~\frac{d^dk}{(2\pi)^d}\exp{(-\pi\tau_2\alpha'k^2)}\prod_{\mu,n}~\sum_{N_{\mu n},\tilde N_{\mu n}=0}^\infty~q^{nN_{\mu n}}\bar q^{n\tilde N_{\mu n}}$$

$$\sum_{N=0}^\infty q^{nN}=(1-q^n)^{-1}$$

$$Z(\tau)=i V_d Z_X(\tau)^d$$

$$Z_X(\tau)=(4\pi^2\alpha'\tau_2)^{-1/2}|\eta(\tau)|^{-2}$$

$$\eta(\tau)=q^{1/24}\prod_{n=1}^\infty ~(1-q^n)$$

$$ds^2=dwd\bar{w}+\epsilon^*dw^2+\epsilon d\bar{w}^2=(1+\epsilon^*+\epsilon)d[w+\epsilon(\bar{w}-w)]d[\bar{w}+\epsilon^*(w-\bar{w})]+O(\epsilon^2)$$

$$w'\cong w'+2\pi\cong w'+2\pi(\tau-2i\tau_2\epsilon)$$

$$\delta\tau=-2i\tau_2\epsilon$$

$$\delta Z(\tau)\, = -\frac{1}{2\pi}\int\,\, d^2w [\delta g_{\bar{w}\bar{w}}\langle T_{ww}(w)\rangle + \delta g_{ww}\langle T_{\bar{w}\bar{w}}(\bar{w})\rangle] = -2\pi i [\delta \tau \langle T_{ww}(0)\rangle - \delta \bar{\tau} \langle T_{\bar{w}\bar{w}}(0)\rangle]$$

$$\partial_w X^\mu(w) \partial_w X_\mu(0) = -\frac{\alpha' d}{2w^2} - \alpha' T_{ww}(0) + O(w)$$

$$Z(\tau)^{-1}\bigl\langle \partial_w X^\mu(w) \partial_w X_\mu(0) \bigr\rangle = d \partial_w \partial_{w'} G'(w,\bar{w};w',\bar{w}')|_{w'=0} = \frac{\alpha' d}{2} \frac{\vartheta_1 \partial_w^2 \vartheta_1 - \partial_w \vartheta_1 \partial_w \vartheta_1}{\vartheta_1^2} + \frac{\alpha' d}{8\pi \tau_2}$$

$$\frac{\alpha' d}{6} \frac{\partial_w^3 \vartheta_1}{\partial_w \vartheta_1} + \frac{\alpha' d}{8\pi \tau_2}$$



$$\langle T_{ww}(0)\rangle = \left(-\frac{d}{6}\frac{\partial_w^3\vartheta_1(0\mid\tau)}{\partial_w\vartheta_1(0\mid\tau)}-\frac{d}{8\pi\tau_2}\right)Z(\tau)$$

$$\partial_\tau \ln Z(\tau)=\frac{\pi id}{3}\frac{\partial_w^3\vartheta_1(0\mid\tau)}{\partial_w\vartheta_1(0\mid\tau)}+\frac{id}{4\tau_2}$$

$$\partial_w^2\vartheta_1\left(\frac{w}{2\pi}\Big|\;\tau\right)=\frac{i}{\pi}\,\partial_\tau\vartheta_1\left(\frac{w}{2\pi}\Big|\;\tau\right)$$

$$\partial_\tau \ln Z(\tau)=-\frac{d}{3}\partial_\tau \ln \partial_w\vartheta_1(0\mid\tau)+\frac{id}{4\tau_2}$$

$$Z(\tau)=|\partial_w\vartheta_1(0\mid\tau)|^{-2d/3}\tau_2^{-d/2}$$

Propagador CFT.

$$\text{Tr}[\exp(2\pi i \tau_1 P - 2\pi \tau_2 H)] = (q\bar{q})^{13/12} \text{Tr}(q^{L_0}\bar{q}^{\tilde{L}_0}) = 4(q\bar{q})^{1/12} \prod_{n=1}^{\infty} |1+q^n|^4$$

$$Z(\tau) = \text{Tr}[(-1)^F \exp(2\pi i \tau_1 P - 2\pi \tau_2 H)] = 0$$

$$\big\langle c(w_1)b(w_2)\tilde{c}(\bar{w}_3)\tilde{b}(\bar{w}_4)\big\rangle.$$

$$\text{Tr}\big[(-1)^Fc_0b_0\tilde{c}_0\tilde{b}_0\exp(2\pi i \tau_1 P - 2\pi \tau_2 H)\big]$$

$$(q\bar{q})^{1/12} \prod_{n=1}^{\infty} |1-q^n|^4 = |\eta(\tau)|^4$$

$$Z(\tau) = \sum_i q^{h_i - c/24} \bar{q}^{\tilde{h}_i - \tilde{c}/24} (-1)^{F_i}$$

$$h_i-\tilde h_i\in {\bf Z}$$

$$Z(i\ell) \stackrel{\ell \rightarrow 0}{\approx} \exp\left[\frac{\pi(c + \tilde{c})}{12\ell}\right]$$

$$\vartheta(v,\tau)=\sum_{n=-\infty}^\infty \exp{(\pi i n^2\tau+2\pi i nv)}$$

$$\begin{aligned}\vartheta(v+1,\tau)&=\vartheta(v,\tau)\\\vartheta(v+\tau,\tau)&=\exp{(-\pi i \tau-2\pi i v)}\vartheta(v,\tau)\end{aligned}$$

$$\begin{aligned}\vartheta(v,\tau+1)&=\vartheta(v+1/2,\tau)\\\vartheta(v/\tau,-1/\tau)&=(-i\tau)^{1/2}\exp{(\pi i v^2/\tau)}\vartheta(v,\tau)\end{aligned}$$

$$\vartheta(v,\tau)=\prod_{m=1}^\infty (1-q^m)(1+zq^{m-1/2})(1+z^{-1}q^{m-1/2})$$



$$q = \exp(2\pi i \tau), z = \exp(2\pi i v)$$

$$\vartheta \begin{bmatrix} a \\ b \end{bmatrix}(\nu, \tau) = \exp [\pi i a^2 \tau + 2\pi i a(\nu + b)] \vartheta(\nu + a\tau + b, \tau) = \sum_{n=-\infty}^{\infty} \exp [\pi i(n+a)^2 \tau + 2\pi i(n+a)(\nu + b)]$$

$$\begin{aligned}\vartheta_{00}(\nu, \tau) &= \vartheta_3(\nu | \tau) = \vartheta \begin{bmatrix} 0 \\ 0 \end{bmatrix}(\nu, \tau) = \sum_{n=-\infty}^{\infty} q^{n^2/2} z^n \\ \vartheta_{01}(\nu, \tau) &= \vartheta_4(\nu | \tau) = \vartheta \begin{bmatrix} 0 \\ 1/2 \end{bmatrix}(\nu, \tau) = \sum_{n=-\infty}^{\infty} (-1)^n q^{n^2/2} z^n\end{aligned}$$

$$\begin{aligned}\vartheta_{10}(\nu, \tau) &= \vartheta_2(\nu | \tau) = \vartheta \begin{bmatrix} 1/2 \\ 0 \end{bmatrix}(\nu, \tau) = \sum_{n=-\infty}^{\infty} q^{(n-1/2)^2/2} z^{n-1/2} \\ \vartheta_{11}(\nu, \tau) &= -\vartheta_1(\nu | \tau) = \vartheta \begin{bmatrix} 1/2 \\ 1/2 \end{bmatrix}(\nu, \tau) \\ &= -i \sum_{n=-\infty}^{\infty} (-1)^n q^{(n-1/2)^2/2} z^{n-1/2}\end{aligned}$$

$$\begin{aligned}\vartheta_{00}(\nu, \tau) &= \prod_{m=1}^{\infty} (1 - q^m)(1 + zq^{m-1/2})(1 + z^{-1}q^{m-1/2}) \\ \vartheta_{01}(\nu, \tau) &= \prod_{m=1}^{\infty} (1 - q^m)(1 - zq^{m-1/2})(1 - z^{-1}q^{m-1/2}) \\ \vartheta_{10}(\nu, \tau) &= 2 \exp(\pi i \tau / 4) \cos \pi \nu \prod_{m=1}^{\infty} (1 - q^m)(1 + zq^m)(1 + z^{-1}q^m) \\ \vartheta_{11}(\nu, \tau) &= -2 \exp(\pi i \tau / 4) \sin \pi \nu \prod_{m=1}^{\infty} (1 - q^m)(1 - zq^m)(1 - z^{-1}q^m)\end{aligned}$$

$$\begin{aligned}\vartheta_{00}(\nu, \tau + 1) &= \vartheta_{01}(\nu, \tau) \\ \vartheta_{01}(\nu, \tau + 1) &= \vartheta_{00}(\nu, \tau) \\ \vartheta_{10}(\nu, \tau + 1) &= \exp(\pi i / 4) \vartheta_{10}(\nu, \tau) \\ \vartheta_{11}(\nu, \tau + 1) &= \exp(\pi i / 4) \vartheta_{11}(\nu, \tau)\end{aligned}$$

$$\begin{aligned}\vartheta_{00}(\nu/\tau, -1/\tau) &= (-i\tau)^{1/2} \exp(\pi i \nu^2 / \tau) \vartheta_{00}(\nu, \tau) \\ \vartheta_{01}(\nu/\tau, -1/\tau) &= (-i\tau)^{1/2} \exp(\pi i \nu^2 / \tau) \vartheta_{10}(\nu, \tau) \\ \vartheta_{10}(\nu/\tau, -1/\tau) &= (-i\tau)^{1/2} \exp(\pi i \nu^2 / \tau) \vartheta_{01}(\nu, \tau) \\ \vartheta_{11}(\nu/\tau, -1/\tau) &= -i(-i\tau)^{1/2} \exp(\pi i \nu^2 / \tau) \vartheta_{11}(\nu, \tau)\end{aligned}$$

$$\vartheta_{00}^4(0, \tau) - \vartheta_{01}^4(0, \tau) - \vartheta_{10}^4(0, \tau) = 0$$

$$\vartheta_{11}(0, \tau) = 0$$

$$\eta(\tau) = q^{1/24} \prod_{m=1}^{\infty} (1 - q^m) = \left[\frac{\partial_v \vartheta_{11}(0, \tau)}{-2\pi} \right]^{1/3}$$



$$\begin{aligned}\eta(\tau+1) &= \exp(i\pi/12)\eta(\tau) \\ \eta(-1/\tau) &= (-i\tau)^{1/2}\eta(\tau)\end{aligned}$$

$$S_{T^2}(1;2;\dots;n)=\frac{1}{2}\int_{F_0}d\tau d\bar{\tau}\left\langle B\tilde{B}\tilde{c}c\mathcal{V}_1(w_1,\bar{w}_1)\prod_{i=2}^n\int\limits_{T^2}dw_id\bar{w}_i\mathcal{V}_i(w_i,\bar{w}_i)\right\rangle$$

$$B=\frac{1}{4\pi}(b,\partial_\tau g)=\frac{1}{2\pi}\int\;d^2wb_{ww}(w)\partial_\tau g_{\bar{w}\bar{w}}=\frac{i}{4\pi\tau_2}\int\;d^2wb_{ww}(w)\rightarrow 2\pi i b_{ww}$$

$$\int\frac{dwd\bar{w}}{2(2\pi)^2\tau_2}$$

$$S_{T^2}(1;2;\dots;n)=\int_{F_0}\frac{d\tau d\bar{\tau}}{4\tau_2}\left\langle b(0)\tilde{b}(0)\tilde{c}(0)c(0)\prod_{i=1}^n\int\limits_{T^2}dw_id\bar{w}_i\mathcal{V}_i(w_i,\bar{w}_i)\right\rangle$$

$$Z_{T^2}=\int_{F_0}\frac{d\tau d\bar{\tau}}{4\tau_2}\langle b(0)\tilde{b}(0)\tilde{c}(0)c(0)\rangle_{T^2}$$

$$Z_{T^2}=iV_{26}\int_{F_0}\frac{d\tau d\bar{\tau}}{4\tau_2}(4\pi^2\alpha'\tau_2)^{-13}|\eta(\tau)|^{-48}$$

$$\frac{d\tau d\bar{\tau}}{\tau_2^2}$$

$$\begin{aligned}Z_{T^2}&=V_d\int_{F_0}\frac{d\tau d\bar{\tau}}{4\tau_2}\int\frac{d^dk}{(2\pi)^d}\exp{(-\pi\tau_2\alpha'k^2)}\sum_{i\in\mathcal{H}^\perp}q^{h_i-1}\bar{q}^{\tilde{h}_i-1}\\&=iV_d\int_{F_0}\frac{d\tau d\bar{\tau}}{4\tau_2}(4\pi^2\alpha'\tau_2)^{-d/2}\sum_{i\in\mathcal{H}^\perp}q^{h_i-1}\bar{q}^{\tilde{h}_i-1}\end{aligned}$$

$$Z_{S_1}(m^2)=V_d\int\frac{d^dk}{(2\pi)^d}\int_0^\infty\frac{dl}{2l}\exp{[-(k^2+m^2)l/2]}=iV_d\int_0^\infty\frac{dl}{2l}(2\pi l)^{-d/2}\exp{(-m^2l/2)}$$

$$m^2=\frac{2}{\alpha'}(h+\tilde{h}-2)$$

$$\delta_{h,\tilde{h}}=\int_{-\pi}^\pi\frac{d\theta}{2\pi}\exp{[i(h-\tilde{h})\theta]}$$

$$\begin{aligned}\sum_{i\in\mathcal{H}^\perp}Z_{S_1}(m_i^2)&=iV_d\int_0^\infty\frac{dl}{2l}\int_{-\pi}^\pi\frac{d\theta}{2\pi}(2\pi l)^{-\frac{d}{2}}\times\sum_{i\in\mathcal{H}^\perp}\exp\left[-\frac{(h_i+\tilde{h}_i-2)l}{\alpha'}+i(h_i-\tilde{h}_i)\theta\right]\\&=iV_d\int_R\frac{d\tau d\bar{\tau}}{4\tau_2}(4\pi^2\alpha'\tau_2)^{-d/2}\sum_{i\in\mathcal{H}^\perp}q^{h_i-1}\bar{q}^{\tilde{h}_i-1}\end{aligned}$$



$$R\colon \tau_2>0, |\tau_1|<\frac{1}{2}$$

$$F_0\colon |\tau|>1, |\tau_1|<\frac{1}{2}, \tau_2>0$$

$$iV_{26}\int^{\infty}\frac{d\tau_2}{2\tau_2}\frac{d\tau_2}{2\tau_2}(4\pi^2\alpha'\tau_2)^{-13}[\exp{(4\pi\tau_2)}+24^2+\cdots]$$

$$iV_d\int^{\infty}\frac{d\tau_2}{2\tau_2}\frac{d\tau_2}{2\tau_2}(4\pi^2\alpha'\tau_2)^{-d/2}\sum_i~\exp{(-\pi\alpha'm_i^2\tau_2)}$$

$$\mathbf{Z}_{\mathrm{vac}}(m^2)=\exp\left[Z_{S_1}(m^2)\right]$$

$$\mathbf{Z}_{\mathrm{vac}}(m^2)=\langle 0|\exp{(-iHT)}|0\rangle=\exp{(-i\rho_0 V_d)}$$

$$\rho_0=\frac{i}{V_d}Z_{S_1}(m^2).$$

$$\int_0^\infty \frac{dl}{2l} \exp{[-(k^2+m^2)l/2]}\rightarrow -\frac{1}{2}\ln{(k^2+m^2)}$$

$$i\int_0^\infty \frac{dl}{2l} \int_{-\infty}^\infty \frac{dk^0}{2\pi} \exp{[-(k^2+m^2)l/2]}\rightarrow \frac{\omega_\mathbf{k}}{2}$$

$$\rho_0=\int\,\frac{d^{d-1}\mathbf{k}}{(2\pi)^{d-1}}\frac{\omega_\mathbf{k}}{2}$$

$$\ln\,\mathbf{Z}_{\mathrm{vac}}(m^2)=-\frac{1}{2}\mathrm{Tr}\mathrm{ln}\,(-\partial^2+m^2)=-\frac{V_d}{2}\int\,\frac{d^dk}{(2\pi)^d}\ln{(k^2+m^2)}$$

$$\rho_0=\frac{i}{V_d}\sum_i~(-1)^{\mathbf{F}_i}Z_{S_1}(m_i^2).$$

$$|\rho_0|\lesssim 10^{-44}\mathrm{GeV}^4$$

$$m_{\mathrm{ew}}^4\approx 10^8\mathrm{GeV}^4$$

$$Z_{C_2}=\int_0^\infty \frac{dt}{2t}\mathrm{Tr}'_0[\exp{(-2\pi tL_0)}]=iV_d\int_0^\infty \frac{dt}{2t}(8\pi^2\alpha't)^{-d/2}\sum_{i\in\mathcal{H}_0^1}\exp{[-2\pi t(h_i-1)]}$$

$$\rightarrow iV_{26}n^2\int_0^\infty \frac{dt}{2t}(8\pi^2\alpha't)^{-13}\eta(it)^{-24}$$

$$\eta(it)=t^{-1/2}\eta(i/t)$$

$$Z_{C_2}=i\frac{V_{26}n^2}{2\pi(8\pi^2\alpha')^{13}}\int_0^\infty ds\eta(is/\pi)^{-24}$$



$$\eta{(is/\pi)^{-24}}=\exp{(2s)}\prod_{n=1}^{\infty}\left[1-\exp{(-2ns)}\right]^{-24}=\exp{(2s)}+24+O(\exp{(-2s)})$$

$$\int_0^\infty ds \mathrm{exp}\left(\beta s \right) \equiv -\frac{1}{\beta}$$

$$\left.\frac{1}{k^2}\right|_{k^\mu=0}$$

$$-\frac{1}{g^2}\!\int~d^dx\left(\frac{1}{2}\partial_\mu\phi\partial^\mu\phi+g\Lambda\phi\right)$$

$$\partial^2\phi=g\Lambda$$

$$-\Lambda \int ~d^{26}x (-G)^{1/2} e^{-\tilde{\Phi}}$$

$$\exp{(-\alpha' k^2 s/2)}$$

$$\langle B | c_0 b_0 \mathrm{exp}\left[-s\big(L_0+\tilde{L}_0\big)\right] | B \rangle.$$

$$(\alpha_n^{\mu}+\tilde{\alpha}_{-n}^{\mu})|B\rangle=(c_n+\tilde{c}_{-n})|B\rangle=\big(b_n-\tilde{b}_{-n}\big)|B\rangle=0,\text{ all }n$$

$$|B\rangle\propto(c_0+\tilde{c}_0)\mathrm{exp}\left[-\sum_{n=1}^{\infty}\left(n^{-1}\alpha_{-n}\cdot\tilde{\alpha}_{-n}+b_{-n}\tilde{c}_{-n}+\tilde{b}_{-n}c_{-n}\right)\right]|0;0\rangle.$$

$$Z_{K_2}=\int_0^\infty\frac{dt}{4t}\mathrm{Tr}'_{\mathrm{c}}\{\Omega\mathrm{exp}\left[-2\pi t\big(L_0+\tilde{L}_0\big)\right]\}=iV_d\int_0^\infty\frac{dt}{4t}(4\pi^2\alpha't)^{-d/2}\sum_{i\in\mathcal{H}_c^\perp}\Omega_i\mathrm{exp}\left[-2\pi t\big(h_i+\tilde{h}_i-2\big)\right]$$

$$Z_{K_2}\rightarrow iV_{26}\int_0^\infty\frac{dt}{4t}(4\pi^2\alpha't)^{-13}\eta(2it)^{-24}$$

$$\begin{array}{l} 0\leq\sigma^1\leq2\pi, 0\leq\sigma^2\leq2\pi t\\ 0\leq\sigma^1\leq\pi, 0\leq\sigma^2\leq4\pi t \end{array}$$

$$w\cong w+2\pi\cong -\bar{w}+2\pi it$$

$$w\cong w+4\pi it, w+\pi\cong-(\bar{w}+\pi)+2\pi it$$

$$Z_{K_2}=i\frac{2^{26}V_{26}}{4\pi(8\pi^2\alpha')^{13}}\int_0^\infty ds\eta{(is/\pi)^{-24}}$$

$$Z_{M_2}=iV_d\int_0^\infty\frac{dt}{4t}(8\pi^2\alpha't)^{-d/2}\sum_{i\in\mathcal{H}_0^\perp}\Omega_i\mathrm{exp}\left[-2\pi t(h_i-1)\right]$$

$$\exp{(2\pi t)}\prod_{n=1}^{\infty}\left[1-(-1)^n\mathrm{exp}{(-2\pi nt)}\right]^{-24}=\vartheta_{00}(0,2it)^{-12}\eta(2it)^{-12}$$



$$Z_{M_2}=\pm inV_{26}\int_0^{\infty}\frac{dt}{4t}(8\pi^2\alpha't)^{-13}\vartheta_{00}(0,2it)^{-12}\eta(2it)^{-12}$$

$$Z_{M_2}=\pm 2in\frac{2^{13}V_{26}}{4\pi(8\pi^2\alpha')^{13}}\int_0^{\infty}ds\vartheta_{00}(0,2is/\pi)^{-12}\eta(2is/\pi)^{-12}$$

$$i\frac{24V_{26}}{4\pi(8\pi^2\alpha')^{13}}(2^{13}\mp n)^2\int_0^{\infty}ds$$

$$\mathrm{Tr}(\lambda^a \lambda^b) \mathrm{Tr}(\lambda^c \lambda^d) = \mathrm{Tr}(\lambda^a \lambda^b \lambda^e \lambda^c \lambda^d \lambda^e).$$

$$x^4\cong x^4+2\pi R$$

$$ds^2=G^D_{MN}dx^M dx^N=G_{\mu\nu}dx^\mu dx^\nu +G_{dd}\bigl(dx^d+A_\mu dx^\mu\bigr)^2$$

$$x'^d=x^d+\lambda(x^\mu).$$

$$A'_\mu=A_\mu-\partial_\mu\lambda$$

$$\phi(x^M)=\sum_{n=-\infty}^{\infty}\phi_n(x^{\mu})\text{exp}\left(inx^d/R\right)$$

$$\partial_\mu\partial^\mu\phi_n(x^\mu)=\frac{n^2}{R^2}\phi_n(x^\mu)$$

$$-p^\mu p_\mu=\frac{n^2}{R^2}$$

$${\boldsymbol R}={\boldsymbol R}_d-2e^{-\sigma}\nabla^2e^{\sigma}-\frac{1}{4}e^{2\sigma}F_{\mu\nu}F^{\mu\nu}$$

$$\begin{aligned} S_1 &= \frac{1}{2\kappa_0^2} \int ~ d^Dx (-G_D)^{1/2} e^{-2\Phi} \big({\boldsymbol R} + 4\nabla_\mu \Phi \nabla^\mu \Phi \big) \\ &= \frac{\pi R}{\kappa_0^2} \int ~ d^d x (-G_d)^{1/2} e^{-2\Phi+\sigma} \times \Big({\boldsymbol R}_d - 4\partial_\mu \Phi \partial^\mu \sigma + 4\partial_\mu \Phi \partial^\mu \Phi - \frac{1}{4} e^{2\sigma} F_{\mu\nu} F^{\mu\nu} \Big) \\ &= \frac{\pi R}{\kappa_0^2} \int ~ d^d x (-G_d)^{1/2} e^{-2\Phi_d} \times \Big({\boldsymbol R}_d - \partial_\mu \sigma \partial^\mu \sigma + 4\partial_\mu \Phi_d \partial^\mu \Phi_d - \frac{1}{4} e^{2\sigma} F_{\mu\nu} F^{\mu\nu} \Big) \end{aligned}$$

$$\partial_\mu + i p_d A_\mu = \partial_\mu + i n \tilde{A}_\mu,$$

$$g_d^2=\frac{\kappa_0^2 e^{2\Phi_d}}{\pi R^3 e^{2\sigma}}=\frac{2\kappa_d^2}{\rho^2}$$

$$\frac{1}{\kappa_d^2}=\frac{2\pi\rho}{\kappa^2}$$



$$\begin{aligned} S_2 &= -\frac{1}{24\kappa_0^2} \int d^Dx (-G_D)^{1/2} e^{-2\Phi} H_{MNL} H^{MNL} \\ &= -\frac{\pi R}{12\kappa_0^2} \int d^d x (-G_d)^{1/2} e^{-2\Phi_d} (\tilde{H}_{\mu\nu\lambda} \tilde{H}^{\mu\nu\lambda} + 3e^{-2\sigma} H_{d\mu\nu} H_d^{\mu\nu}) \end{aligned}$$

$$\tilde{H}_{\mu\nu\lambda} = (\partial_\mu B_{\nu\lambda} - A_\mu H_{d\nu\lambda}) + \text{Cyclic permutations}$$

$$B'_{v\lambda} = B_{v\lambda} - \lambda H_{dv\lambda}$$

$$X \cong X+2\pi R$$

$$k=\frac{n}{R}, n\in \mathbf{Z}$$

$$X(\sigma+2\pi)=X(\sigma)+2\pi R w, w\in\mathbf{Z}$$

$$\partial X(z) = -i\left(\frac{\alpha'}{2}\right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\alpha_m}{z^{m+1}}, \bar{\partial} X(\bar{z}) = -i\left(\frac{\alpha'}{2}\right)^{1/2} \sum_{m=-\infty}^{\infty} \frac{\tilde{\alpha}_m}{\bar{z}^{m+1}}.$$

$$2\pi R w = \oint (dz \partial X + d\bar{z} \bar{\partial} X) = 2\pi (\alpha'/2)^{1/2} (\alpha_0 - \tilde{\alpha}_0)$$

$$p = \frac{1}{2\pi\alpha'} \oint (dz \partial X - d\bar{z} \bar{\partial} X) = (2\alpha')^{-1/2} (\alpha_0 + \tilde{\alpha}_0)$$

$$\begin{aligned} p_L &\equiv (2/\alpha')^{1/2} \alpha_0 = \frac{n}{R} + \frac{wR}{\alpha'} \\ p_R &\equiv (2/\alpha')^{1/2} \tilde{\alpha}_0 = \frac{n}{R} - \frac{wR}{\alpha'} \end{aligned}$$

$$\begin{aligned} L_0 &= \frac{\alpha' p_L^2}{4} + \sum_{n=1}^{\infty} \alpha_{-n} \alpha_n, \\ \tilde{L}_0 &= \frac{\alpha' p_R^2}{4} + \sum_{n=1}^{\infty} \tilde{\alpha}_{-n} \tilde{\alpha}_n. \end{aligned}$$

Función de Partición.

$$\begin{aligned} (q\bar{q})^{-1/24} \text{Tr}(q^{L_0} \bar{q}^{\tilde{L}_0}) &= |\eta(\tau)|^{-2} \sum_{n,w=-\infty}^{\infty} q^{\alpha' p_L^2/4} \bar{q}^{\alpha' p_R^2/4} \\ &= |\eta(\tau)|^{-2} \sum_{n,w=-\infty}^{\infty} \exp \left[-\pi \tau_2 \left(\frac{\alpha' n^2}{R^2} + \frac{w^2 R^2}{\alpha'} \right) + 2\pi i \tau_1 n w \right] \\ &\sum_{n=-\infty}^{\infty} \exp(-\pi a n^2 + 2\pi i b n) = a^{-1/2} \sum_{m=-\infty}^{\infty} \exp \left[-\frac{\pi(m-b)^2}{a} \right]. \end{aligned}$$



$$2\pi R Z_X(\tau) \sum_{m,w=-\infty}^{\infty} \exp\left(-\frac{\pi R^2 |m-w\tau|^2}{\alpha' \tau_2}\right)$$

$$\begin{aligned} X(\sigma^1 + 2\pi, \sigma^2) &= X(\sigma^1, \sigma^2) + 2\pi wR \\ X(\sigma^1 + 2\pi\tau_1, \sigma^2 + 2\pi\tau_2) &= X(\sigma^1, \sigma^2) + 2\pi mR \end{aligned}$$

$$X_{\text{cl}} = \sigma^1 wR + \sigma^2 (m - w\tau_1)R/\tau_2$$

Operadores de vórtice.

$$[x_L, p_L] = [x_R, p_R] = i$$

$$X(z, \bar{z}) = X_L(z) + X_R(\bar{z})$$

$$\begin{aligned} X_L(z) &= x_L - i \frac{\alpha'}{2} p_L \ln z + i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\alpha_m}{mz^m}, \\ X_R(\bar{z}) &= x_R - i \frac{\alpha'}{2} p_R \ln \bar{z} + i \left(\frac{\alpha'}{2} \right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\tilde{\alpha}_m}{m\bar{z}^m}. \end{aligned}$$

$$\begin{aligned} X_L(z_1)X_L(z_2) &\sim -\frac{\alpha'}{2} \ln z_{12}, X_R(\bar{z}_1)X_R(\bar{z}_2) \sim -\frac{\alpha'}{2} \ln \bar{z}_{12} \\ X_L(z_1)X_R(\bar{z}_2) &\sim 0 \end{aligned}$$

$$\mathcal{V}_{k_L k_R}(z, \bar{z}) =: e^{ik_L X_L(z) + ik_R X_R(\bar{z})} :$$

$$\mathcal{V}_{k_L k_R}(z_1, \bar{z}_1) \mathcal{V}_{k'_L k'_R}(z_2, \bar{z}_2) \sim z_{12}^{\alpha' k'_L / 2} \bar{z}_{12}^{\alpha' k_R k'_R / 2} \mathcal{V}_{(k+k')_L (k+k')_R}(z_2, \bar{z}_2).$$

$$\exp [\pi i \alpha' (k_L k'_L - k_R k'_R)] = \exp [2\pi i (nw' + wn')] = 1.$$

$$[X_L(z_1), X_L(z_2)] = \frac{\pi i \alpha'}{2} \text{sign}(\sigma_1^1 - \sigma_2^1).$$

$$\mathcal{V}_{k_L k_R}(z, \bar{z}) = \exp [\pi i (k_L - k_R)(p_L + p_R)\alpha'/4] \circ e^{ik_L X_L(z) + ik_R X_R(\bar{z})} \circ$$

$$\begin{aligned} \exp \{ \pi i [(k_L - k_R)(k'_L + k'_R) - (k'_L - k'_R)(k_L + k_R)]\alpha'/4 \} \\ = \exp [\pi i (nw' - wn')] \end{aligned}$$

$$\prod_{\substack{i,j=1 \\ i < j}}^n |z_{ij}|^{\alpha' k_i k_j} \rightarrow \prod_{\substack{i,j=1 \\ i < j}}^n z_{ij}^{\alpha' k_{Li} k_{Lj} / 2} \bar{z}_{ij}^{\alpha' k_{Ri} k_{Rj} / 2}$$

$$2\pi R \delta_{\Sigma_i n_i, 0} \delta_{\Sigma_t w_t, 0}$$



Operadores DDF.

$$X^+(z)X^-(0) \sim \frac{\alpha'}{2} \ln z, X^+(z)X^+(0) \sim X^-(z)X^-(0) \sim 0.$$

$$V^i(nk_0, z) = \partial X^i(z) e^{ink_0 X^+(z)} (2/\alpha')^{1/2}$$

$$V^i(nk_0, z)V^j(mk_0, 0) \sim -\frac{\delta^{ij}}{z^2} e^{i(n+m)k_0 X^+(0)} - \frac{ink_0 \delta^{ij}}{z} \partial X^+(0) e^{i(n+m)k_0 X^+(0)}$$

$$A_n^i = \oint \frac{dz}{2\pi} V^i(nk_0, z)$$

$$[A_m^i, A_n^j] = m\delta^{ij}\delta_{m,-n} \frac{\alpha' k_0 p^+}{2}$$

$$m^2 = -k^\mu k_\mu = (k_L^{25})^2 + \frac{4}{\alpha'}(N-1) = (k_R^{25})^2 + \frac{4}{\alpha'}(\tilde{N}-1)$$

$$\begin{aligned} m^2 &= \frac{n^2}{R^2} + \frac{w^2 R^2}{\alpha'^2} + \frac{2}{\alpha'}(N + \tilde{N} - 2) \\ 0 &= nw + N - \tilde{N} \end{aligned}$$

$$\alpha_{-1}^\mu \tilde{\alpha}_{-1}^\nu |0;k\rangle, (\alpha_{-1}^\mu \tilde{\alpha}_{-1}^{25} + \alpha_{-1}^{25} \tilde{\alpha}_{-1}^\mu) |0;k\rangle (\alpha_{-1}^\mu \tilde{\alpha}_{-1}^{25} - \alpha_{-1}^{25} \tilde{\alpha}_{-1}^\mu) |0;k\rangle, \alpha_{-1}^{25} \tilde{\alpha}_{-1}^{25} |0;k\rangle$$

$$\partial X^\mu \bar{\partial} X^{25} - \partial X^{25} \bar{\partial} X^\mu = \bar{\partial}(X^{25} \partial X^\mu) - \partial(X^{25} \bar{\partial} X^\mu)$$

$$\frac{2^{1/2} g_{c,25}}{\alpha'} : (\partial X^\mu \bar{\partial} X^{25} \pm \partial X^{25} \bar{\partial} X^\mu) e^{ik \cdot X} :$$

$$g_{c,25} : e^{ik_L \cdot X_L(z) + ik_R \cdot X_R(\bar{z})} :$$

$$\begin{aligned} &-2^{-1/2} \pi i g_{c,25} (2\pi)^{25} \delta^{25} \left(\sum_i k_i \right) k_{23}^\mu (k_{L23}^{25} \pm k_{R23}^{25}) \\ &\rightarrow -2^{3/2} \pi i g_{c,25} (2\pi)^{25} \delta^{25} \left(\sum_i k_i \right) k_2^\mu (k_{L2}^{25} \pm k_{R2}^{25}) \end{aligned}$$

Simetrías de gauge.

$$(n+w)^2 + 4N = (n-w)^2 + 4\tilde{N} = 4$$

$$n=w=\pm 1, N=0, \tilde{N}=1, n=-w=\pm 1, N=1, \tilde{N}=0$$

$$n=\pm 2, w=N=\tilde{N}=0, w=\pm 2, n=N=\tilde{N}=0$$

$$: \bar{\partial} X^\mu e^{ik \cdot X} \exp [\pm 2i\alpha'^{-1/2} X_L^{25}] : , : \partial X^\mu e^{ik \cdot X} \exp [\pm 2i\alpha'^{-1/2} X_R^{25}] : .$$



$$\begin{aligned} j^1(z) &=: \cos [2\alpha'^{-1/2}X_L^{25}(z)]: \\ j^2(z) &=: \sin [2\alpha'^{-1/2}X_L^{25}(z)]: \\ j^3(z) &= i\partial X_L^{25}(z)/\alpha'^{1/2} \end{aligned}$$

$$j^i(z)j^j(0) \sim \frac{\delta^{ij}}{2z^2} + i\frac{\epsilon^{ijk}}{z}j^k(0)$$

$$\begin{aligned} j^i(z) &= \sum_{m=-\infty}^{\infty} \frac{j_m^i}{z^{m+1}} \\ [j_m^i, j_n^j] &= \frac{m}{2}\delta_{m,-n}\delta^{ij} + i\epsilon^{ijk}j_{m+n}^k \end{aligned}$$

Escalares y acoplamientos.

$$g_{25}^2 = 2\kappa_{25}^2/\alpha'$$

$$g_4^2 = 2\kappa_4^2/\alpha'$$

$$g_{G,4}^2(E) = \kappa_4^2 E^2$$

$$g_5^2 = 2\pi\rho_5 g_4^2$$

$$\hat{g}_5^2 = g_5^2 E = 2\pi\rho_5 E g_4^2$$

Mecanismo de Higgs.

$$m = \frac{|R^2 - \alpha'|}{R\alpha'} \approx \frac{2}{\alpha'} |R - \alpha'^{1/2}|,$$

$$:j^i(z)\tilde{j}^j(\bar{z})e^{ik\cdot X(z,\bar{z})}:.$$

$$U(M) \propto \epsilon^{ijk}\epsilon^{i'j'k'}M_{ii'}M_{jj'}M_{kk'} = \det M$$

$$U(M) = \frac{\partial U(M)}{\partial M_{ij}} = 0$$

$$M_{11}M_{22}M_{33}=M_{11}M_{22}=M_{11}M_{33}=M_{22}M_{33}=0$$

$$m^2 = \frac{n^2}{R^2} + \frac{w^2 R^2}{\alpha'^2} + \frac{2}{\alpha'} (N + \tilde{N} - 2),$$

$$R \rightarrow R' = \frac{\alpha'}{R}, n \leftrightarrow w.$$

$$p_L^{25} \rightarrow p_L^{25}, p_R^{25} \rightarrow -p_R^{25}$$

$$X'^{25}(z,\bar{z}) = X_L^{25}(z) - X_R^{25}(\bar{z})$$

$$R_{\text{self-dual}} = R_{SU(2) \times SU(2)} = \alpha'^{1/2}$$



$$\rho' = \frac{\alpha'}{\rho}, \kappa' = \frac{\alpha'^{1/2}}{\rho} \kappa$$

$$e^{\Phi'} = \frac{\alpha'^{1/2}}{\rho} e^\Phi$$

Dimensiones múltiples.

$$X^m \cong X^m + 2\pi R, 26 - k \leq m \leq 25$$

$$S = \frac{(2\pi R)^k}{2\kappa_0^2} \int d^d x (-G_d)^{1/2} e^{-2\Phi_d} [\mathbf{R}_d + 4\partial_\mu \Phi_d \partial^\mu \Phi_d - \frac{1}{4} G^{mn} G^{pq} (\partial_\mu G_{mp} \partial^\mu G_{nq} + \partial_\mu B_{mp} \partial^\mu B_{nq}) - \frac{1}{4} G_{mn} F_{\mu\nu}^m F^{n\mu\nu} - \frac{1}{4} G^{mn} H_{m\mu\nu} H_n^{\mu\nu} - \frac{1}{12} H_{\mu\nu\lambda} H^{\mu\nu\lambda}]$$

$$B_{mn} \partial_a (g^{1/2} \epsilon^{ab} X^m \partial_b X^n)$$

$$X^m(\sigma^1, \sigma^2) = x^m(\sigma^2) + w^m R \sigma^1$$

$$L = \frac{1}{2\alpha'} G_{mn} (\dot{x}^m \dot{x}^n + w^m w^n R^2) - \frac{i}{\alpha'} B_{mn} \dot{x}^m w^n R$$

$$p_m = - \frac{\partial L}{\partial v^m} = \frac{1}{\alpha'} (G_{mn} v^n + B_{mn} w^n R),$$

$$v_m = \alpha' \frac{n_m}{R} - B_{mn} w^n R$$

$$\frac{1}{2\alpha'} G_{mn} (v^m v^n + w^m w^n R^2)$$

$$m^2 = \frac{1}{2\alpha'^2} G_{mn} (v_L^m v_L^n + v_R^m v_R^n) + \frac{2}{\alpha'} (N + \tilde{N} - 2) \\ v_{L,R}^m = v^m \pm w^m R$$

$$0 = G_{mn} (v_L^m v_L^n - v_R^m v_R^n) + 4\alpha' (N - \tilde{N}) \\ = 4\alpha' (n_m w^m + N - \tilde{N}) \quad (8.4.10)$$

$$X^m = (w_1^m \sigma^1 + w_2^m \sigma^2) R$$

$$2\pi i b_{mn} w_1^m w_2^n$$

$$G_{mn} = e_m^r e_n^r$$

$$k_{rL} = e_r^m \frac{v_{mL}}{\alpha'}, k_{rR} = e_r^m \frac{v_{mR}}{\alpha'}$$



$$\begin{aligned}m^2 &= \frac{1}{2}(k_{rL}k_{rL} + k_{rR}k_{rR}) + \frac{2}{\alpha'}(N + \tilde{N} - 2) \\0 &= \alpha'(k_{rL}k_{rL} - k_{rR}k_{rR}) + 4(N - \tilde{N})\end{aligned}$$

Compactación de Narain.

$$:e^{ik_L\cdot X_L(z)+ik_R\cdot X_R(\bar{z})}::e^{ik'_L\cdot X_L(0)+ik'_R\cdot X_R(0)}:\sim z^{l_L\cdot l'_L}\bar{z}_Rl_R\cdot l'_R:e^{i(k_L+k'_L)\cdot X_L(0)+i(k_R+k'_R)\cdot X_R(0)}:$$

$$l_L\cdot l'_L-l_R\cdot l'_R\equiv l\circ l'\in{\bf Z}$$

$$\Gamma\subset\Gamma^*$$

$$l\circ l\in 2{\bf Z}\,\,\,{\rm for\,\,all}\,\,l\in\Gamma.$$

$$2l\circ l'=(l+l')\circ(l+l')-l\circ l-l'\circ l'\in 2{\bf Z}$$

$$Z_\Gamma(\tau)=|\eta(\tau)|^{-2k}\sum_{l\in\Gamma}\exp{(\pi i\tau l_L^2-\pi i\bar{\tau}l_R^2)}$$

$$\sum_{l'\in\Gamma}\delta(l-l')=V_\Gamma^{-1}\sum_{l''\in\Gamma^*}\exp{(2\pi i l''\circ l)}.$$

$$\begin{aligned}Z_\Gamma(\tau)&=V_\Gamma^{-1}|\eta(\tau)|^{-2k}\sum_{l''\in\Gamma^*}\int d^{2k}l\exp{(2\pi i l''\circ l+\pi i\tau l_L^2-\pi i\bar{\tau}l_R^2)}\\&=V_\Gamma^{-1}(\tau\bar{\tau})^{-k/2}|\eta(\tau)|^{-2k}\sum_{l''\in\Gamma^*}\exp{(-\pi i l''_L{}^2/\tau+\pi i l''_R{}^2/\bar{\tau})}=V_\Gamma^{-1}Z_{\Gamma^*}(-1/\tau)\end{aligned}$$

$$\Gamma=\Gamma^*$$

$$\Gamma' = \Lambda \Gamma$$

$$l_{L,R}=\frac{n}{r}\pm\frac{mr}{2}$$

$$l'_L=l_L\cosh\lambda+l_R\sinh\lambda,l'_R=l_L\sinh\lambda+l_R\cosh\lambda$$

$$\frac{O(k,k,\mathbf{R})}{O(k,\mathbf{R})\times O(k,\mathbf{R})}$$

$$\frac{2k(2k-1)}{2}-k(k-1)=k^2$$

$$\begin{gathered}\Lambda\Gamma_0\cong\Lambda'\Lambda\Lambda''\Gamma_0\\\Lambda'\in O(k,\mathbf{R})\times O(k,\mathbf{R}),\Lambda''\in O(k,k,\mathbf{Z})\end{gathered}$$

$$\frac{O(k,k,\mathbf{R})}{O(k,\mathbf{R})\times O(k,\mathbf{R})\times O(k,k,\mathbf{Z})}$$

$$x'^m=L^m{}_n x^n$$

$$b_{mn}\rightarrow b_{mn}+N_{mn}$$



$$-\frac{1}{2}\,g_{ij}(\phi)\partial_\mu\phi^i\partial^\mu\phi^j$$

$$\frac{16}{d-2}\partial_\mu \Phi \partial^\mu \Phi + G^{mn}G^{pq}\big(\partial_\mu G_{mp}\partial^\mu G_{nq} + \partial_\mu B_{mp}\partial^\mu B_{nq} \big)$$

$$\rho=\frac{R^2}{\alpha'}\big(B_{24,25}+i{\rm det}^{1/2}G_{mn}\big)=b_{24,25}+i\frac{V}{4\pi^2\alpha'}$$

$$ds^2 = \frac{\alpha'\rho_2}{R^2\tau_2}\big|dX^{24}+\tau dX^{25}\big|^2$$

$$\mathrm{PSL}(2,\mathbf{Z})\times \mathrm{PSL}(2,\mathbf{Z})\times \mathbf{Z}_2^2$$

$$\frac{\partial_\mu \tau \partial^\mu \bar{\tau}}{\tau_2^2} + \frac{\partial_\mu \rho \partial^\mu \bar{\rho}}{\rho_2^2}$$

$$\frac{PSL(2,\mathbf{R})\times PSL(2,\mathbf{R})}{U(1)\times U(1)\times PSL(2,\mathbf{Z})\times PSL(2,\mathbf{Z})\times \mathbf{Z}_2^2}.$$

$${\bf Orbifolds}.$$

$$X^{25}\cong -X^{25}$$

$$X^m\rightarrow-X^m, 26-k\leq m\leq 25.$$

$$\begin{gathered} t^m\!:\! X^{25}\cong X^{25}+2\pi Rm,\\ t^mr\!:\! X^{25}\cong 2\pi Rm-X^{25},\end{gathered}$$

$$X^{25}(\sigma^1+2\pi)=-X^{25}(\sigma^1)$$

$$\left|N,\tilde{N};k^{\mu},n,w\right\rangle\rightarrow(-1)^{\sum_{m=1}^{\infty}\left(N_m^{25}+\tilde{N}_m^{25}\right)}\left|N,\tilde{N};k^{\mu},-n,-w\right\rangle,$$

$$X^{25}(z,\bar z)=i\left(\frac{\alpha'}{2}\right)^{1/2}\sum_{m=-\infty}^{\infty}\frac{1}{m+1/2}\left(\frac{\alpha_{m+1/2}^{25}}{z^{m+1/2}}+\frac{\tilde\alpha_{m+1/2}^{25}}{\bar z^{m+1/2}}\right).$$

$$X^{25}(\sigma^1+2\pi)=2\pi R-X^{25}(\sigma^1).$$

$$m^2=\frac{4}{\alpha'}\Big(N-\frac{15}{16}\Big), N=\tilde{N}$$

$$(q\bar q)^{-1/24}\text{Tr}_U\left(\frac{1+r}{2}q^{L_0}\bar q^{\tilde L_0}\right)$$

$$\frac{1}{2} Z_{\text{tor}}(R,\tau) + \frac{1}{2} (q\bar q)^{-1/24} \prod_{m=1}^{\infty} |1+q^m|^{-2}$$

$$(q\bar q)^{1/48}\text{Tr}_T\left(\frac{1+r}{2}q^{L_0}\bar q^{\tilde L_0}\right)=(q\bar q)^{1/48}\left[\prod_{m=1}^{\infty}|1-q^{m-1/2}|^{-2}+\prod_{m=1}^{\infty}|1+q^{m-1/2}|^{-2}\right]$$



$$Z_{\text{orb}}\left(R,\tau\right)=\frac{1}{2}Z_{\text{tor}}\left(R,\tau\right)+\left|\frac{\eta(\tau)}{\vartheta_{10}(0,\tau)}\right|+\left|\frac{\eta(\tau)}{\vartheta_{01}(0,\tau)}\right|+\left|\frac{\eta(\tau)}{\vartheta_{00}(0,\tau)}\right|$$

$$\begin{gathered} X^{25}(\sigma^1+2\pi,\sigma^2)\,=(-1)^{a+1}X^{25}(\sigma^1,\sigma^2)\\ X^{25}(\sigma^1+2\pi\tau_1,\sigma^2+2\pi\tau_2)\,=(-1)^{b+1}X^{25}(\sigma^1,\sigma^2) \end{gathered}$$

$${\bf Twisting}.$$

$$\phi(\sigma^1+2\pi)=h\cdot\phi(\sigma^1)$$

$$P_H=\frac{1}{\text{order}(H)}\sum_{h_2\in H}\hat{h}_2$$

$$Z=\frac{1}{\text{order}(H)}\sum_{h_1,h_2\in H}Z_{h_1,h_2}$$

$$\phi'(\sigma^1)=h_2\cdot\phi(\sigma^1)$$

$$\phi'(\sigma^1+2\pi)=h'_1\cdot\phi'(\sigma^1)$$

$$h'_1=h_2h_1h_2^{-1}$$

$$c=1CFTs$$

$$r'\colon X^{25}\rightarrow X^{25}+\pi\alpha^{1/2}$$

$$Z_{\text{orb}}\left(\alpha'^{1/2},\tau\right)=Z_{\text{tor}}\left(2\alpha'^{1/2},\tau\right)$$

$$X^{25}\rightarrow X^{25}+\frac{2\pi\alpha'^{1/2}}{k}$$

$$j^3\tilde{j}^3,j^1\tilde{j}^1+j^2\tilde{j}^2,j^1\tilde{j}^2-j^2\tilde{j}^1$$

$$A_{25}(x^M)=-\frac{\theta}{2\pi R}=-i\Lambda^{-1}\frac{\partial\Lambda}{\partial x^{25}},\Lambda(x^{25})=\exp\left(-\frac{i\theta x^{25}}{2\pi R}\right)$$

$$W_q=\exp\left(iq\phi\;\;dx^{25}A_{25}\right)=\exp\left(-iq\theta\right)$$

$$S=\int\;d\tau\left(\frac{1}{2}\dot{X}^M\dot{X}_M+\frac{m^2}{2}-iqA_M\dot{X}^M\right)$$

$$p_{25}=-\frac{\partial L}{\partial v^{25}}=v^{25}-\frac{q\theta}{2\pi R}$$

$$v_{25}=\frac{2\pi l+q\theta}{2\pi R}$$

$$H=\frac{1}{2}\big(p_\mu p^\mu+v_{25}^2+m^2\big)$$



$$v_{25} = p_{25} = \frac{2\pi l + q\theta}{2\pi R}$$

$$A_{25} = -\frac{1}{2\pi R} \text{diag}(\theta_1, \theta_2, \dots, \theta_n)$$

$$v_{25} = \frac{2\pi l - \theta_j + \theta_i}{2\pi R}$$

$$m^2 = \frac{(2\pi l - \theta_j + \theta_i)^2}{4\pi^2 R^2} + \frac{1}{\alpha'}(N-1)$$

$$m^2 = \frac{(\theta_j - \theta_i)^2}{4\pi^2 R^2}$$

$$U(r_1) \times \dots \times U(r_s), \sum_{i=1}^s r_i = n$$

$$\text{Tr}([A_m, A_n]^2)$$

Dualidad T.

$$X'^{25}(z, \bar{z}) = X_L^{25}(z) - X_R^{25}(\bar{z}).$$

$$\partial_n X^{25} = -i\partial_t X'^{25}$$

$$X'^{25}(\pi) - X'^{25}(0) = \int_0^\pi d\sigma^1 \partial_1 X'^{25} = -i \int_0^\pi d\sigma^1 \partial_2 X^{25} = -2\pi\alpha' v^{25} = -\frac{2\pi\alpha' l}{R} = -2\pi l R'$$

$$\Delta X'^{25} = X'^{25}(\pi) - X'^{25}(0) = -(2\pi l - \theta_j + \theta_i)R'.$$

$$X'^{25} = \theta_i R' = -2\pi\alpha' A_{25,ii}.$$

$$\begin{aligned} X'^{25}(z, \bar{z}) &= \theta_i R' - \frac{iR'}{2\pi} (2\pi l - \theta_j + \theta_i) \ln(z/\bar{z}) + i \left(\frac{\alpha'}{2}\right)^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\alpha_m^{25}}{m} (z^{-m} - \bar{z}^{-m}) \\ &= \theta_i R' + \frac{\sigma^1}{\pi} \Delta X'^{25} - (2\alpha')^{1/2} \sum_{\substack{m=-\infty \\ m \neq 0}}^{\infty} \frac{\alpha_m^{25}}{m} \exp(-m\sigma^2) \sin m\sigma^1 \\ m^2 &= \left(\frac{\Delta X'^{25}}{2\pi\alpha'}\right)^2 + \frac{1}{\alpha'}(N-1) \end{aligned}$$

Membranas.

$$\begin{aligned} \alpha_{-1}^\mu |k; ii\rangle, \quad &\mathcal{V} = i\partial_t X^\mu \\ \alpha_{-1}^{25} |k; ii\rangle, \quad &\mathcal{V} = i\partial_t X^{25} = \partial_n X'^{25} \end{aligned}$$



$$\pmb{S}_p=-T_p\int~d^{p+1}\xi e^{-\Phi}[-{\rm det}(G_{ab}+B_{ab}+2\pi\alpha'F_{ab})]^{1/2}$$

$$G_{ab}(\xi)=\frac{\partial X^\mu}{\partial \xi^a}\frac{\partial X^\nu}{\partial \xi^b}G_{\mu\nu}(X(\xi)), B_{ab}(\xi)=\frac{\partial X^\mu}{\partial \xi^a}\frac{\partial X^\nu}{\partial \xi^b}B_{\mu\nu}(X(\xi))$$

$$X'^2 = -2\pi\alpha'X^1F_{12}$$

$$\int~dX^1[1+(\partial_1X'^2)^2]^{1/2}=\int~dX^1[1+(2\pi\alpha'F_{12})^2]^{1/2}$$

$$\frac{i}{4\pi\alpha'}\int_M d^2\sigma g^{1/2}\epsilon^{ab}\partial_aX^\mu\partial_bX^\nu B_{\mu\nu}+i\int_{\partial M} dX^\mu A_\mu$$

$$\delta A_\mu = \partial_\mu \lambda$$

$$\delta B_{\mu\nu}=\partial_\mu\zeta_\nu-\partial_\nu\zeta_\mu$$

$$\delta A_\mu = -\zeta_\mu/2\pi\alpha'$$

$$B_{\mu\nu}+2\pi\alpha'F_{\mu\nu}\equiv 2\pi\alpha'\mathcal{F}_{\mu\nu}$$

$$V\propto {\rm Tr}([X_m,X_n][X^m,X^n]).$$

$$\pmb{S}_p=-T_p\int~d^{p+1}\xi {\rm Tr}\{e^{-\Phi}[-{\rm det}(G_{ab}+B_{ab}+2\pi\alpha'F_{ab})]^{1/2}+O([X^m,X^n]^2)\}$$

$$T_pe^{-\Phi}\prod_{i=1}^p~(2\pi R_i)$$

$$2\pi\alpha'^{1/2}T_pe^{-\Phi'}\prod_{i=1}^{p-1}~(2\pi R_i)$$

$$T_{p-1}e^{-\Phi'}\prod_{i=1}^{p-1}~(2\pi R_i)$$

$$T_p=T_{p-1}/2\pi\alpha'^{1/2},$$

$$\mathcal{A}=iV_{p+1}\int_0^{\infty}\frac{dt}{t}(8\pi^2\alpha't)^{-(p+1)/2}\exp{(-ty^2/2\pi\alpha')}\eta(it)^{-24}$$

$$=\frac{iV_{p+1}}{(8\pi^2\alpha')^{(p+1)/2}}\int_0^{\infty}dt t^{(21-p)/2}\exp{(-ty^2/2\pi\alpha')}\times[\exp{(2\pi/t)}+24+\cdots]$$

$$\mathcal{A}=iV_{p+1}\frac{24}{2^{12}}(4\pi^2\alpha')^{11-p}\pi^{(p-23)/2}\Gamma\left(\frac{23-p}{2}\right)|y|^{p-23}=iV_{p+1}\frac{24\pi}{2^{10}}(4\pi^2\alpha')^{11-p}G_{25-p}(y)$$

$$\pmb{S}=\frac{1}{2\kappa^2}\int~d^{26}X(-\tilde{G})^{1/2}\left(\tilde{\pmb{R}}-\frac{1}{6}\nabla_\mu\tilde{\Phi}\tilde{\nabla}^\mu\tilde{\Phi}\right)$$



$$\mathcal{S}_p=-\tau_p\int\;d^{p+1}\xi\text{exp}\left(\frac{p-11}{12}\tilde{\Phi}\right)(-\text{det}\tilde{G}_{ab})^{1/2}$$

$$F_\nu \equiv \partial^{\hat{\mu}} h_{\mu\nu} - \frac{1}{2} \partial_\nu h^\hat{\mu}_\mu = 0$$

$$\mathcal{S}=-\frac{1}{8\kappa^2}\int\;d^{26}X\Big(\partial_\mu h_{\nu\lambda}\partial^{\hat{\mu}} h^{\hat{\nu}\hat{\lambda}}-\frac{1}{2}\partial_\mu h^{\hat{\nu}}_\nu\partial^{\hat{\mu}} h^{\hat{\lambda}}_\lambda+\frac{2}{3}\partial_\mu\tilde{\Phi}\partial^{\hat{\mu}}\tilde{\Phi}\Big)$$

$$\begin{gathered}\langle \tilde{\Phi}\tilde{\Phi}\rangle\, = -\frac{(D-2)i\kappa^2}{4k^2},\\ \langle h_{\mu\nu}h_{\sigma\rho}\rangle\, = -\frac{2i\kappa^2}{k^2}\Big(\eta_{\mu\sigma}\eta_{\nu\rho}+\eta_{\mu\rho}\eta_{\nu\sigma}-\frac{2}{D-2}\eta_{\mu\nu}\eta_{\sigma\rho}\Big).\end{gathered}$$

$$\mathcal{S}_p=-\tau_p\int\;d^{p+1}\xi\Big(\frac{p-11}{12}\tilde{\Phi}-\frac{1}{2}h_{aa}\Big)$$

$$\mathcal{A}=\frac{i\kappa^2\tau_p^2}{k_\perp^2}V_{p+1}\bigg\{6\left[\frac{p-11}{12}\right]^2+\frac{1}{2}\bigg[2(p+1)-\frac{1}{12}(p+1)^2\bigg]\bigg\}\!=\!\frac{6i\kappa^2\tau_p^2}{k_\perp^2}V_{p+1}$$

$$\tau_p^2=\frac{\pi}{256\kappa^2}(4\pi^2\alpha')^{11-p}.$$

$$\frac{\tau_{25}}{4}(2\pi\alpha')^2\text{Tr}\big(F_{\mu\nu}F^{\mu\nu}\big).$$

$$\frac{g_o^2}{g_c} \! = \! \frac{4 \pi \alpha' g_o'^2}{\kappa} = 2^{18} \pi^{25/2} \alpha'^6$$

$$\Omega\colon X_L^M(z)\leftrightarrow X_R^M(z)$$

$$\begin{array}{ll}\Omega: & X'^m(z,\bar z) \leftrightarrow -X'^m(\bar z,z) \\ & X^\mu(z,\bar z) \leftrightarrow X^\mu(\bar z,z)\end{array}$$

$$\begin{array}{l} G_{\mu\nu}(x')\,=G_{\mu\nu}(x),B_{\mu\nu}(x')=-B_{\mu\nu}(x)\\ G_{\mu n}(x')\,=-G_{\mu n}(x),B_{\mu n}(x')=B_{\mu n}(x)\\ G_{mn}(x')\,=G_{mn}(x),B_{mn}(x')=-B_{mn}(x)\end{array}$$

$$W=\text{diag}\big(e^{i\theta_1},e^{-i\theta_1},e^{i\theta_2},e^{-i\theta_2},\cdots,e^{i\theta_{n/2}},e^{-i\theta_{n/2}}\big)$$

$$2^{12-k}T_p\int\;d^{p+1}\xi e^{-\Phi}(-\text{det}G_{ab})^{1/2}$$

$$g_{c,d}\colon e^{ik\cdot X}\colon$$

$$\mathcal{V}_j(k;z,\bar{z})\overline{\mathcal{V}_{j'}(k;0,0)}=\frac{g_{c,d}^2}{z^2\bar{z}^2}+\cdots$$

$$e^{-2\lambda}\big\langle \tilde{c}(\bar{z}_1)c(z_1)\tilde{c}(\bar{z}_2)c(z_2)\tilde{c}(\bar{z}_3)c(z_3)\colon e^{ik\cdot X}(z,\bar{z})\colon\big\rangle=\frac{8\pi i}{\alpha'g_{c,d}^2}|z_{12}z_{13}z_{23}|^2(2\pi)^d\delta^d(k)$$



$$S_{S_2}(1;2;3)=e^{-2\lambda}\bigl\langle \hat{\mathcal{V}}'_1(\infty,\infty)\hat{\mathcal{V}}_2(1,1)\hat{\mathcal{V}}_3(0,0)\bigr\rangle_{S_2}$$

$$\hat{\mathcal{V}}_j = \tilde{c} c \mathcal{V}_j$$

$$S_{S_2}(1;2;3)=e^{-2\lambda}\left\langle \langle \hat{\mathcal{V}}_1|\hat{\mathcal{V}}_2(1,1)\mid \hat{\mathcal{V}}_3\right\rangle=e^{-2\lambda}c_{123}.$$

$$\big| \hat{\mathcal{V}}_3 \big\rangle = Q_\mathrm{B} |\chi\rangle.$$

$$e^{2\lambda} S_{S_2}(1;2;3)=\Big\langle \langle \hat{\mathcal{V}}_1| \big[\hat{\mathcal{V}}_2(1,1),Q_\mathrm{B} \big] \mid \chi \Big\rangle + \Big\langle \langle \hat{\mathcal{V}}_1| Q_\mathrm{B} \hat{\mathcal{V}}_2(1,1) \mid \chi \Big\rangle=0$$

$$\hat{\mathcal{V}}_3=Q_\mathrm{B}\cdot\mathcal{X}.$$

$$\sum_n\,S_{mn}S^*_{pn}\stackrel{?}{=}\delta_{mp}$$

$$S_{mn}=\delta_{mn}+iT_{mn}$$

$$T_{mp}-T_{pm}^*=i\sum_n\;T_{mn}T_{pn}^*$$

$$T_{S_2}(1;2;3)=T_{S_2}(-1;-2;-3)^*$$

$$\hat{\mathcal{V}}_{-i}=-\overline{\hat{\mathcal{V}}}_i.$$

$$\bigl\langle \hat{\mathcal{V}}'_1(\infty,\infty)\hat{\mathcal{V}}_2(1,1)\hat{\mathcal{V}}_3(0,0)\bigr\rangle_{S_2}=-\bigl\langle \hat{\mathcal{V}}'_{-1}(\infty,\infty)\hat{\mathcal{V}}_{-2}(1,1)\hat{\mathcal{V}}_{-3}(0,0)\bigr\rangle^*_{S_2},$$

$$S_{S_2}(1;2;3;4)=e^{-2\lambda}\int_{\mathbf{C}}d^2z\bigl\langle \hat{\mathcal{V}}'_1(\infty,\infty)\hat{\mathcal{V}}_2(1,1)\hat{\mathcal{V}}_3(0,0)\mathcal{V}_4(z,\bar{z})\bigr\rangle_{S_2}.$$

$$\Big\langle \langle \hat{\mathcal{V}}_1| \mathrm{T}\big[\hat{\mathcal{V}}_2(1,1)\mathcal{V}_4(z,\bar{z})\big] \mid \hat{\mathcal{V}}_3\Big\rangle.$$

$$\mathcal{V}_4(z,\bar{z})=(z\bar{z})^{-1}\{b_0,[\tilde{b}_0,\hat{\mathcal{V}}_4(z,\bar{z})]\}.$$

$$\theta(1-|z|)\left\langle \langle \hat{\mathcal{V}}_1| \hat{\mathcal{V}}_2(1,1)b_0\tilde{b}_0z^{L_0-1}\bar{z}^{\tilde{L}_0-1}\hat{\mathcal{V}}_4(1,1) \mid \hat{\mathcal{V}}_3\right\rangle+\theta(|z|$$

$$-1)\left\langle \langle \hat{\mathcal{V}}_1| \hat{\mathcal{V}}_4(1,1)b_0\tilde{b}_0z^{-L_0-1}\bar{z}^{-\tilde{L}_0-1}\hat{\mathcal{V}}_2(1,1) \mid \hat{\mathcal{V}}_3\right\rangle$$

$$\begin{aligned} &\theta(1-|z|)\sum_{i,i'}z^{\alpha'(k_i^2+m_i^2)/4-1}\bar{z}^{\alpha'(k_i^2+\tilde{m}_i^2)/4-1}\times\Big\langle \langle \hat{\mathcal{V}}_1| \hat{\mathcal{V}}_2(1,1)b_0\tilde{b}_0 \mid i\Big\rangle\mathcal{G}^{ii'}\left\langle \langle i'|\hat{\mathcal{V}}_4(1,1) \mid \hat{\mathcal{V}}_3\right\rangle+\theta(|z| \\ &-1)\sum_{i,i'}z^{-\alpha'(k_i^2+m_i^2)/4-1}\bar{z}^{-\alpha'(k_i^2+\tilde{m}_i^2)/4-1} \\ &\times\Big\langle \langle \hat{\mathcal{V}}_1| \hat{\mathcal{V}}_4(1,1)b_0\tilde{b}_0 \mid i\Big\rangle\mathcal{G}^{ii'}\left\langle \langle i'|\hat{\mathcal{V}}_2(1,1) \mid \hat{\mathcal{V}}_3\right\rangle \end{aligned}$$



$$k_i^2=(k_1+k_2)^2>\frac{4}{\alpha'},$$

$$\int_{|z|<1} d^2 z z^{\alpha'(k^2+m^2)/4-1} \bar{z}^{\alpha'(\bar{k}^2+\bar{m}^2)/4-1} = \frac{8\pi}{\alpha'} \frac{\delta_{m^2,\tilde{m}^2}}{k^2+m^2-i\epsilon}$$

$$\left[b_0\tilde{b}_0z^{\pm L_0-1}\bar{z}^{\pm \tilde{L}_0-1},Q_\mathrm{B}\right]=\pm b_0\partial_{\bar{z}}\left(z^{\pm L_0-1}\bar{z}^{\pm \tilde{L}_0}\right)\mp \tilde{b}_0\partial_z\left(z^{\pm L_0}\bar{z}^{\pm \tilde{L}_0-1}\right)$$

$$T_{S_2}(1;2;3;4)-T_{S_2}(-1;-2;-3;-4)^*$$

$$\begin{aligned}&= i \sum_{j_5} \int \frac{d^{d-1}\mathbf{k}_5}{2E_5(2\pi)^{d-1}} T_{S_2}(1;2;5) T_{S_2}(-5;3;4) + \text{2 permutations} \\&= -ie^{-4\lambda} \sum_{j_5} \int \frac{d^{d-1}\mathbf{k}_5}{2E_5(2\pi)^{d-1}} \left\langle \hat{\mathcal{V}}_1 \middle| \hat{\mathcal{V}}_2(1,1) \mid \hat{\mathcal{V}}_5 \right\rangle \left\langle \hat{\mathcal{V}}_{-5} \middle| \hat{\mathcal{V}}_3(1,1) \mid \hat{\mathcal{V}}_4 \right\rangle \\&\quad + \text{2 permutations}\end{aligned}$$

$$\frac{1}{k^2+m^2-i\epsilon}-\frac{1}{k^2+m^2+i\epsilon}=2\pi i\delta(k^2+m^2)$$

$$T_{S_2}(1;2;3;4)-T_{S_2}(-1;-2;-3;-4)^*$$

$$\begin{aligned}&= \frac{16\pi^2 i}{\alpha' e^{2\lambda}} \sum_{i,i'} \delta_{m_i^2, \tilde{m}_i^2} \delta(k_i^2 + m_i^2) \langle \hat{\mathcal{V}}_1 \mid \hat{\mathcal{V}}_2(1,1) b_0 \tilde{b}_0 \mid i \rangle \mathcal{G}^{ii'} \langle i' \mid \hat{\mathcal{V}}_3(1,1) \mid \hat{\mathcal{V}}_4 \rangle \\&\quad + \text{2 permutations}\end{aligned}$$

$$i\leftrightarrow(j,k)$$

$$\begin{aligned}&\sum_{i,i'} \delta_{m_i^2, \tilde{m}_i^2} \delta(k_i^2 + m_i^2) b_0 \tilde{b}_0 \mid i \rangle \mathcal{G}^{ii'} \langle i' \mid = \delta_{L_0, \tilde{L}_0} \delta(4L_0/\alpha') b_0 \tilde{b}_0 \\&= -\frac{i\alpha'}{16\pi^2 e^{2\lambda}} \int \frac{d^{d-1}\mathbf{k}}{2E(2\pi)^{d-1}} \sum_{\substack{j \in \hat{\mathcal{H}} \\ k^0 = \pm \omega_k}} \mid \hat{\mathcal{V}}_j(k) \rangle \langle \hat{\mathcal{V}}_j(k) \mid \\&\left\langle \overline{\hat{\mathcal{V}}_j(k)} \right| \tilde{c}_0 c_0 \mid \hat{\mathcal{V}}_{j'}(k') \rangle = \frac{8\pi i e^{2\lambda}}{\alpha'} \delta_{jj'} (2\pi)^d \delta^d(k-k')\end{aligned}$$

$$\begin{aligned}\mathcal{P}|b\rangle_{\mathrm{P}}&=|b\rangle_{\mathrm{P}},\mathcal{P}|a\rangle_{\mathrm{U}}=0,\mathcal{P}|a\rangle_{\mathrm{N}}=0;\\\mathcal{U}|b\rangle_{\mathrm{P}}&=0,\mathcal{U}|a\rangle_{\mathrm{U}}=0,\mathcal{U}|a\rangle_{\mathrm{N}}=|a\rangle_{\mathrm{U}}.\end{aligned}$$

$$1=\mathcal{P}+Q_{\mathrm{B}}\mathcal{U}+\mathcal{U}Q_{\mathrm{B}}.$$

$$-i\,\frac{\eta^{\mu\nu}}{k^2}$$

$$\eta^{\mu\nu}=(\eta^{\mu\nu}-k^\mu n^\nu-n^\mu k^\nu)+k^\mu n^\nu+n^\mu k^\nu.$$



$$\left.\frac{\partial z_1}{\partial z}\right|_{z_2}=\frac{z_1}{z}$$

$$\gamma \colon \frac{z'-U}{z'-V}=K\frac{z-U}{z-V},$$

$$|K|^{-1/2}<\left|\frac{z-U}{z-V}\right|<|K|^{1/2}.$$

$$ds^2 = \frac{dzd\bar{z}}{(\mathrm{Im} z)^2}$$

$$(A_1B_1^{-1}A_1^{-1}B_1)(A_2B_2^{-1}A_2^{-1}B_2)\ldots(A_gB_g^{-1}A_g^{-1}B_g)$$

$$\partial_{\bar{z}}\omega_z=0,\omega_{\bar{z}}=0$$

$$\dim\ker P_0^T-\dim\ker P_0=-\chi=2g-2.$$

$$\oint_{-A_i} dz \omega_{zj} = \delta_{ij}$$

$$\tau_{ij}=\oint_{-B_i} dz \omega_{zj}$$

$$y^2=\prod_{i=1}^{2k}\;(z-Z_i)$$

$$z_1z_2=q$$

$$(1-\epsilon)^{-1}|q|^{1/2}>|z_1|>(1-\epsilon)|q|^{1/2}$$

$$M=M_1\infty M_2(z_1,z_2,q)$$

$$M=M_08(z_1,z_2,q).$$

$$\ln\;|q|<\mathrm{Im} w<0, w\cong w+2\pi,$$

$$z^{(0)}=\frac{az}{z-2}, z^{(1)}=\frac{a(z-1)}{z+1}, z^{(\infty)}=\frac{a}{2z-1}$$

$$z_4=(q-a^2)^2/4a^2q$$

$$r=3g+n-2$$

$$\langle \cdots 1 \cdots 2 \rangle_M = \sum_{ij} \, \left\langle \cdots 1 \mathcal{A}_i^{(z_1)} \right\rangle_{M_1} G^{ij} \left\langle \mathcal{A}_j^{(z_2)} \cdots 2 \right\rangle_{M_2} \, .$$



$$\left\langle \cdots \mathcal{A}_k^{(z_1)} \right\rangle_{M_1} = \sum_{ij} \; \left\langle \cdots \mathcal{A}_i^{(z_1)} \right\rangle_{M_1} \mathsf{G}^{ij} \left\langle \mathcal{A}_j^{(u)} \mathcal{A}_k^{(z)} \right\rangle_{S_2}.$$

$$\mathsf{G}^{ij}\mathcal{G}_{jk}=\delta_k^i$$

$$\mathsf{G}^{ij}=\mathcal{G}^{ij}$$

$$\mathcal{A}_j^{(z'_2)}=q^{h_j}\bar q^{\tilde h_j}\mathcal{A}_j^{(z_2)}$$

$$\langle \cdots 1 \cdots 2 \rangle_M = \sum_{ij} \; q^{h_j}\bar q^{\tilde h_j} \left\langle \cdots 1 \mathcal{A}_i^{(z_1)} \right\rangle_{M_1} \mathcal{G}^{ij} \left\langle \mathcal{A}_j^{(z_2)} \cdots 2 \right\rangle_{M_2}$$

$$\langle \ldots \rangle_M = \sum_{ij} \; q^{h_j}\bar q^{\tilde h_j} \mathcal{G}^{ij} \left\langle \ldots \mathcal{A}_i^{(z_1)} \mathcal{A}_j^{(z_2)} \right\rangle_{M_0}$$

$$z_1'=z_1+\sum_{n=-\infty}^\infty \epsilon_n z_1^{n+1}\\ z_2'=z_2-\sum_{n=-\infty}^\infty \epsilon_{-n} q^{-n} z_2^{n+1}$$

$$\mathcal{A}_i^{(z'_1)}\mathcal{G}^{ij}\mathcal{A}_j^{(z'_2)} - \mathcal{A}_i^{(z_1)}\mathcal{G}^{ij}\mathcal{A}_j^{(z_2)} = -\sum_{n=-\infty}^\infty \epsilon_n \left[L_n \cdot \mathcal{A}_i^{(z_1)}\mathcal{G}^{ij}\mathcal{A}_j^{(z_2)} - \mathcal{A}_i^{(z_1)}\mathcal{G}^{ij}L_{-n} \cdot \mathcal{A}_j^{(z_2)} \right]$$

$$2\pi(m,n).$$

$$(\sigma^1,\sigma^2)=(\sigma'^1,\sigma'^2)\begin{bmatrix}m&n\\q&p\end{bmatrix}$$

$$(\sigma'^1,\sigma'^2)=2\pi(1,0)$$

$$\oint\;\frac{dz}{2\pi i}\,b_{z_1z_1}\frac{\partial z_1}{\partial q}\Big|_{z_2}=\frac{b_0}{q}$$

$$\int_{|q|<1}\frac{d^2q}{q\bar q}q^{\alpha'(k_i^2+m_i^2)/4}\bar q^{\alpha'(\tilde k_i^2+\tilde m_i^2)/4}\mathcal G^{ij}b_0\tilde b_0=\frac{8\pi\delta_{m_i^2,\tilde m_i^2}\mathcal G^{ij}b_0\tilde b_0}{\alpha'\big(k_i^2+m_i^2-i\epsilon)}.$$

$$g(i_1,\dots,i_n) = \sum_g \; \int_{\mathcal{V}_{g,n}} d^m t \left\langle \prod_{k=1}^m \; B_k \prod_{l=1}^n \; \hat{\mathcal{V}}_{i_l} \right\rangle_g$$

$$\partial X^\mu \bar{\partial} X_\mu \frac{dq d\bar{q}}{q \bar{q}}$$

$$-4\pi \partial X^\mu \bar{\partial} X_\mu \ln{(ae^{-\omega})}$$

$$\delta S = - \sum_\chi \; \Lambda_\chi \int \; d^dx (-G)^{1/2} e^{-x \Phi}$$



$$\int \; [d\Psi] {\exp\left(iS[\Psi]\right)}$$

$$\delta \Psi = Q_{\mathrm{B}} \Lambda$$

$$S_0=\frac{1}{2}\langle\langle\Psi|Q_{\mathrm{B}}\mid\Psi\rangle.$$

$$Q_{\mathrm{B}}\Psi=0$$

$$\Psi[X,c,\tilde c]=\sum_i\Phi_i(x)\Psi_i[X',c,\tilde c]$$

$$\int \; [d\Psi] \rightarrow \prod_i \int \; [d\Phi_i]$$

$$\begin{aligned} \Psi=[\varphi(x)+A_\mu(x)\alpha^\mu_{-1}+B(x)b_{-1}+C(x)c_{-1}\\ +\varphi'(x)c_0+A'_\mu(x)\alpha^\mu_{-1}c_0+B'(x)b_{-1}c_0+C'(x)c_{-1}c_0+\cdots]\Psi_0\\ \delta \Psi=Q_{\mathrm{B}}\Lambda+g\Psi*\Lambda-g\Lambda*\Psi. \end{aligned}$$

$$Q_{\mathrm{B}}(\Psi_1 * \Psi_2) = (Q_{\mathrm{B}}\Psi_1) * \Psi_2 + \Psi_1 * (Q_{\mathrm{B}}\Psi_2)$$

$$\int \;\Psi_1 * \Psi_2 = \int \;\Psi_2 * \Psi_1$$

$$S=\frac{1}{2}\int \;\Psi * Q_{\mathrm{B}}\Psi +\frac{2g}{3}\int \;\Psi * \Psi * \Psi$$

$$S=\frac{1}{2}\langle \mathcal{V}_{\Psi}Q_{\mathrm{B}}\cdot \mathcal{V}_{\Psi}\rangle_{D_2}+\frac{2g}{3}\langle \mathcal{V}_{\Psi}\mathcal{V}_{\Psi}\mathcal{V}_{\Psi}\rangle_{D_2}$$

$$\int_0^{\infty} dt \text{exp}\left(-tL_0\right)=L_0^{-1}$$

$$S_0'=\frac{1}{2}\langle\langle\Psi|(c_0-\tilde c_0)Q_{\mathrm{B}}\mid\Psi\rangle$$

$$\int_{-\infty}^{\infty}dy\text{exp}\left[-\frac{y^2}{2!}-\lambda\frac{y^3}{3!}\right]=\lambda^{-1}\int_{-\infty}^{\infty}dz\text{exp}\left[-\frac{1}{\lambda^2}\biggl(\frac{z^2}{2!}+\frac{z^3}{3!}\biggr)\right]$$

$$\sum_{n=0}^{\infty}\frac{(-\lambda)^n}{6^nn!}\int_{-\infty}^{\infty}dy y^{3n}\text{exp}\left(-y^2/2\right)=(2\pi)^{1/2}\sum_{k=0}^{\infty}\lambda^{2k}C_{2k}$$

$$C_{2k}=\frac{2^k\Gamma(3k+1/2)}{\pi^{1/2}3^{2k}(2k)!}$$

$$C_{2k}\approx k^k\approx k!,$$



$$\sum_{k=0}^{\infty}\,\lambda^{2k}k!\,f_{2k}$$

$$\exp\left[-O(1/\lambda^2)\right]$$

$$\int_0^\infty dt \text{exp}\left(-t\right) \sum_{k=0}^\infty \left(t\lambda^2\right)^k f_{2k}$$

$$\sum_{k=0}^{\infty}g_0^{2k}O(k!)=\sum_{k=0}^{\infty}g_c^kO(k!)$$

$$\exp\left[-O(1/g_{\mathrm c})\right]$$

$$\frac{8\pi i g_c^2}{\alpha'}(2\pi)^{26}\delta^{26}\left(\sum_i~k_i\right)\int_{\mathbf C}d^2z_4|z_4|^{-\alpha'u/2-4}|1-z_4|^{-\alpha't/2-4}$$

$$\frac{\partial}{\partial z_4}(u\ln |z_4|^2+t\ln |1-z_4|^2)=\frac{u}{z_4}+\frac{t}{z_4-1}=0$$

$$S\approx \exp\left[-\alpha' (\text{sln } s + t \text{ln } t + u \text{ln } u)/2\right]$$

$$X_{\rm cl}^\mu(\sigma)=i\sum_i~k_i^\mu G'(\sigma,\sigma_i),$$

$$\exp\left[-\sum_{i < j}~k_i\cdot k_jG'\big(\sigma_i,\sigma_j\big)\right]$$

$$z_2^N=\frac{(z_1-a_1)(z_1-a_2)}{(z_1-a_3)(z_1-a_4)}$$

$$S_g\approx \exp\left[-\alpha' (\text{sln } s + t \text{ln } t + u \text{ln } u)/2(g+1)\right]$$

$$\exp\left[\alpha' t \text{ln } s/2(g+1)\right]$$

$$\langle 0|[X^1(\sigma)-x^1]^2|0\rangle = \sum_{n=1}^{\infty}\,\frac{\alpha'}{2n^2}\langle 0|(\alpha_n\alpha_{-n}+\tilde{\alpha}_n\tilde{\alpha}_{-n})|0\rangle = \alpha'\sum_{n=1}^{\infty}\,\frac{1}{n}$$

$$s^{2+\alpha' t/2}\frac{\Gamma(-\alpha' t/4-1)}{\Gamma(\alpha' t/4+2)}\sim \frac{s^2}{t}\exp\left[-\frac{q^2\alpha'}{2}\text{ln}\left(\alpha' s\right)\right]$$

$$\int_0^\infty dm n(m) \text{exp}\left(-\alpha' \pi m^2 \ell\right) \approx \text{exp}\left(4\pi/\ell\right)$$

$$\int_0^\infty dm \text{exp}\left(4\pi m \alpha'^{1/2}\right) \text{exp}\left(-m/T\right)$$



$$T_{\rm H}=\frac{1}{4\pi \alpha'^{1/2}}$$

$$F(T,m^2)=T\int\,\frac{d^{d-1}\textbf{k}}{(2\pi)^{d-1}}\ln\left[1-\exp\left(-\omega_k/T\right)\right]=-\int_0^\infty\frac{dt}{t}(2\pi t)^{-d/2}\sum_{r=1}^\infty\,\exp\left(-\frac{m^2t}{2}-\frac{r^2}{2T^2t}\right)$$

$$F(T)=-\int_R\frac{d\tau d\bar{\tau}}{2\tau_2}(4\pi^2\alpha'\tau_2)^{-13}|\eta(\tau)|^{-48}\sum_{r=1}^\infty\,\exp\left(-\frac{r^2}{4\pi T^2\alpha'\tau_2}\right)$$

$$R\colon \tau_1\leq \frac{1}{2}, |\tau_2|>0$$

$$\eta(i\tau_2)=\eta(i/\tau_2)\tau_2^{-1/2}\approx \exp{(-\pi/12\tau_2)}\text{ as }\tau_2\rightarrow 0$$

$$\tilde F(T)=F(T)+\rho_0$$

$$\frac{1}{T}\tilde F(T)=\frac{T}{4T_{\rm H}^2}\tilde F\big(4T_{\rm H}^2/T\big)$$

$$\tilde F(T)\rightarrow \frac{T^2}{4T_{\rm H}^2}\rho_0$$

$$\begin{array}{l} G_{\mu\nu}(x) \, = \eta_{\mu\nu}, B_{\mu\nu}(x) = 0, \Phi(x) = V_\mu x^\mu \\ V_\mu \, = \delta_\mu{}^1 \Big(\frac{26-D}{6\alpha'} \Big)^{1/2} \end{array}$$

$$-\partial^\mu\partial_\mu T(x)+2V^\mu\partial_\mu T(x)-\frac{4}{\alpha'}T(x)=0$$

$$T(x)=\exp{(q\cdot x)}, (q-V)^2=\frac{2-D}{6\alpha'}$$

$$q_1=\alpha_{\pm}=\left(\frac{26-D}{6\alpha'}\right)^{1/2}\pm\left(\frac{2-D}{6\alpha'}\right)^{1/2}$$

$$S_{\sigma}=\frac{1}{4\pi\alpha'}\int_Md^2\sigma g^{1/2}\big[g^{ab}\eta_{\mu\nu}\partial_aX^{\mu}\partial_bX^{\nu}+\alpha'RV_1X^1+T_0\mathrm{exp}\left(\alpha_-X^1\right)\big]$$

$$S=\frac{1}{4\pi\alpha'}\int_Md^2\sigma g^{1/2}\big(g^{ab}\partial_aX^{\mu}\partial_bX_{\mu}+\mu\big)$$

$$g_{ab}(\sigma)=\exp{(2\phi(\sigma))}\hat{g}_{ab}(\sigma)$$

$$g^{1/2}R=\hat{g}^{1/2}\big(\hat{R}-2\hat{\nabla}^2\phi\big)$$



$$S = \frac{1}{4\pi\alpha'} \int_M d^2\sigma \hat{g}^{1/2} \left[\hat{g}^{ab} \partial_a X^\mu \partial_b X_\mu + \mu \exp 2\phi + \frac{13\alpha'}{3} (\hat{g}^{ab} \partial_a \phi \partial_b \phi + \hat{R}\phi) \right]$$

$$\begin{aligned}\hat{g}_{ab}(\sigma) &\rightarrow \exp(2\omega(\sigma))\hat{g}_{ab}(\sigma) \\ \phi(\sigma) &\rightarrow \phi(\sigma) - \omega(\sigma)\end{aligned}$$

Campos bosónicos.

$$[\hat{q}, \hat{p}] = i\hbar$$

$$\langle q_f | \exp(-i\hat{H}T/\hbar) | q_i \rangle$$

$$\hat{A}(t) = \exp(i\hat{H}t/\hbar) \hat{A} \exp(-i\hat{H}t/\hbar)$$

$$\hat{q}(t)|q,t\rangle = |q,t\rangle q$$

$$|q,t\rangle = \exp(i\hat{H}t/\hbar)|q\rangle$$

$$\langle q_f, T \mid q_i, 0 \rangle = \int dq \langle q_f, T \mid q, t \rangle \langle q, t \mid q_i, 0 \rangle$$

$$t_m = m\epsilon, \epsilon = T/N$$

$$\langle q_f, T \mid q_i, 0 \rangle = \int dq_{N-1} \dots dq_1 \prod_{m=0}^{N-1} \langle q_{m+1}, t_{m+1} \mid q_m, t_m \rangle$$

$$\langle q_{m+1}, t_{m+1} \mid q_m, t_m \rangle = \langle q_{m+1} | \exp(-i\hat{H}\epsilon/\hbar) | q_m \rangle = \int dp_m \langle q_{m+1} \mid p_m \rangle \langle p_m | \exp(-i\hat{H}\epsilon/\hbar) | q_m \rangle$$

$$\langle p_m | \hat{H}(\hat{p}, \hat{q}) | q_m \rangle = H(p_m, q_m) \langle p_m \mid q_m \rangle.$$

$$\begin{aligned}\int dp_m \exp \left[-\frac{iH(p_m, q_m)\epsilon}{\hbar} \right] \langle q_{m+1} \mid p_m \rangle \langle p_m \mid q_m \rangle \\ = \int \frac{dp_m}{2\pi\hbar} \exp \left\{ -\frac{i}{\hbar} [H(p_m, q_m)\epsilon - p_m(q_{m+1} - q_m)] + O(\epsilon^2) \right\}\end{aligned}$$

$$\begin{aligned}\langle q_f, T \mid q_i, 0 \rangle \approx \int \frac{dp_{N-1}}{2\pi\hbar} dq_{N-1} \dots \frac{dp_1}{2\pi\hbar} dq_1 \frac{dp_0}{2\pi\hbar} \times \exp \left\{ -\frac{i}{\hbar} \sum_{m=0}^{N-1} [H(p_m, q_m)\epsilon - p_m(q_{m+1} - q_m)] \right\} \\ \rightarrow \int [dp dq] \exp \left\{ \frac{i}{\hbar} \int_0^T dt [p\dot{q} - H(p, q)] \right\}\end{aligned}$$

$$0 = \frac{\partial}{\partial p} [p\dot{q} - H(p, q)] = \dot{q} - \frac{\partial}{\partial p} H(p, q).$$

$$\langle q_f, T \mid q_i, 0 \rangle = \int [dq] \exp \left[\frac{i}{\hbar} \int_0^T dt L(q, \dot{q}) \right]$$



$$\int [dq]_{q_i,0}^{q_f,T} \exp \left(i \int_0^T dt L \right) = \int dq \int [dq]_{q,t}^{q_f,T} \exp \left(i \int_t^T dt L \right) \int [dq]_{q_i,0}^{q,t} \exp \left(i \int_0^t dt L \right)$$

$$\int [dq]_{q_i,0}^{q_f,T} \exp(iS) q(t) = \int dq \langle q_f, T | q, t \rangle q \langle q, t | q_i, 0 \rangle = \int dq \langle q_f, T | \hat{q}(t) | q, t \rangle \langle q, t | q_i, 0 \rangle$$

$$= \langle q_f, T | \hat{q}(t) | q_i, 0 \rangle$$

$$\int [dq]_{q_i,0}^{q_f,T} \exp(iS) q(t) q(t') = \langle q_f, T | T[\hat{q}(t) \hat{q}(t')] | q_i, 0 \rangle.$$

$$T[\hat{A}(t) \hat{B}(t')] = \theta(t-t') \hat{A}(t) \hat{B}(t') + \theta(t'-t) \hat{B}(t') \hat{A}(t)$$

$$\int [dq]_{q_i,0}^{q_f,T} \exp(iS) \frac{\delta S}{\delta q(t)} \mathcal{F} = -i \int [dq]_{q_i,0}^{q_f,T} \left[\frac{\delta}{\delta q(t)} \exp(iS) \right] \mathcal{F} = i \int [dq]_{q_i,0}^{q_f,T} \exp(iS) \frac{\delta \mathcal{F}}{\delta q(t)}$$

$$\frac{\delta}{\delta q(t)} q(t') = \delta(t-t')$$

$$\frac{1}{\epsilon} \frac{\partial}{\partial q_m}, m = t/\epsilon$$

$$\left\langle q(t') \frac{\delta S}{\delta q(t)} \dots \right\rangle = i \langle \delta(t-t') \dots \rangle$$

$$q(t') \frac{\delta S}{\delta q(t)} = i \delta(t-t')$$

$$\begin{aligned} S &= \frac{1}{2} \int dt (\dot{q}^2 - \omega^2 q^2) \\ \frac{\delta S}{\delta q(t)} &= -\ddot{q}(t) - \omega^2 q(t) \end{aligned}$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) q(t) = 0$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) q(t) q(t') = -i \delta(t-t')$$

$$\begin{aligned} \left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) \hat{q}(t) &= 0 \\ \left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) T[\hat{q}(t) \hat{q}(t')] &= -i \delta(t-t') \end{aligned}$$

$$\left(\frac{\partial^2}{\partial t^2} + \omega^2 \right) T[\hat{q}(t) \hat{q}(t')] = T\{ [\ddot{q}(t) + \omega^2 \hat{q}(t)] \hat{q}(t') \} + \delta(t-t') [\dot{q}(t), \hat{q}(t)] = 0 - i \delta(t-t')$$

$$\langle q_f | \exp[-i\hat{H}(T-t_>)] \hat{q} \exp[-i\hat{H}(t_>-t_<)] \hat{q} \exp[-i\hat{H}t_<] | q_i \rangle,$$



$$\langle q_f, u | \text{T} \left[\prod_a \hat{q}(u_a) \right] | q_i, 0 \rangle_{\text{E}}$$

$$\langle q_f, t | \text{T} \left[\prod_a \hat{q}(t_a) \right] | q_i, 0 \rangle$$

$$\langle q_f, U | \text{T} \left[\prod_a \hat{q}(u_a) \right] | q_i, 0 \rangle_{\text{E}} = \int [dq]_{q_i, 0}^{q_f, U} \exp \left[\int_0^U du L(q, i\partial_u q) \right] \prod_a q(u_a).$$

$$S_{\text{E}} = \int du L_{\text{E}}(q, \partial_u q) = - \int du L(q, i\partial_u q)$$

$$\text{Tr} \exp(-\hat{H}U) = \sum_i \exp(-E_i U)$$

$$\int dq \langle q | \exp(-\hat{H}U) | q \rangle = \int dq \langle q, U | q, 0 \rangle_{\text{E}} = \int dq \int [dq]_{q, 0}^{q, U} \exp(-S_{\text{E}}) = \int [dq]_{\text{P}} \exp(-S_{\text{E}})$$

$$\int [dq]_{\text{A}} \exp(-S_{\text{E}}) = \int dq \int [dq]_{-q, 0}^{q, U} \exp(-S_{\text{E}}) = \int dq \langle q, U | -q, 0 \rangle_{\text{E}} = \int dq \langle q, U | \hat{R} | q, 0 \rangle_{\text{E}}$$

$$= \text{Tr}[\exp(-\hat{H}U)\hat{R}]$$

$$\frac{1}{2}(\text{ periodic } + \text{ antiperiodic }) = \sum_{R_i=\pm 1} \exp(-E_i U)$$

$$\hat{A}(u)^\dagger = \hat{A}(-u)$$

$$\overline{\hat{A}(u)} = \hat{A}(-u)^\dagger.$$

$$\frac{1}{2} \int d^d x \phi(x) \Delta \phi(x)$$

$$Z[J] = \left\langle \exp \left[i \int d^d x J(x) \phi(x) \right] \right\rangle = \int [d\phi] \exp \left[-\frac{1}{2} \int d^d x \phi(x) \Delta \phi(x) + i \int d^d x J(x) \phi(x) \right]$$

$$\begin{aligned} Z[J] &= Z[0] \exp \left[-\frac{1}{2} \int d^d x d^d y J(x) \Delta^{-1}(x, y) J(y) \right] \\ &= Z[0] \exp \left[\frac{1}{2} \int d^d x d^d y \Delta^{-1}(x, y) \frac{\delta^2}{\delta \phi(x) \delta \phi(y)} \right] \times \exp \left[i \int d^d x J(x) \phi(x) \right] \Big|_{\phi=0} \\ \mathcal{F}[\phi] &= \int [dJ] \exp \left[i \int d^d x J(x) \phi(x) \right] f[J] \\ \langle \mathcal{F} \rangle &= Z[0] \exp \left[\frac{1}{2} \int d^d x d^d y \Delta^{-1}(x, y) \frac{\delta^2}{\delta \phi(x) \delta \phi(y)} \right] \mathcal{F}[\phi] \Big|_{\phi=0}. \end{aligned}$$



$$\frac{\delta}{\delta \phi(x)} = \frac{\delta}{\delta \phi_1(x)} + \frac{\delta}{\delta \phi_2(x)} + \cdots$$

$$\frac{\delta^2}{\delta \phi_i(x) \delta \phi_i(y)} \text{ (no sum on } i)$$

$$\phi(x)=\sum_i~\phi_i\Phi_i(x)$$

$$\begin{aligned}\Delta\Phi_i(x)&=\lambda_i\Phi_i(x)\\\int d^dx\Phi_i(x)\Phi_j(x)&=\delta_{ij}\end{aligned}$$

$$[d\phi]=\prod_id\phi_i$$

$$\langle 1\rangle=\prod_i\left[\int~d\phi_i\exp\left(-\frac{1}{2}\phi_i^2\lambda_i\right)\right]=\prod_i\left(\frac{2\pi}{\lambda_i}\right)^{1/2}=\left(\det\frac{\Delta}{2\pi}\right)^{-1/2}$$

$$\int~[d\phi d\phi^*]\exp\left(-\int~d^dx\phi^*\Delta\phi\right)=(\det\Delta)^{-1}$$

$$\int~[d\phi^1d\phi^2]\exp\left(i\int~d^dx\phi^1\Delta\phi^2\right)=(\det\Delta)^{-1}$$

$$\int_{-\infty}^\infty dx \int_{-\infty}^\infty dy \exp{(i\lambda xy)}=2\pi \int_{-\infty}^\infty dx \delta(\lambda x)=\frac{2\pi}{\lambda}$$

$$\begin{aligned}\langle q_f,U\mid q_i,0\rangle_{\mathrm{E}}&=\int~[dq]^{q_f,U}_{q_i,0}\exp{(-S_{\mathrm{E}})}\\S_{\mathrm{E}}&=\frac{1}{2}\int_0^Udu[(\partial_uq)^2+\omega^2q^2]+S_{\mathrm{ct}}\end{aligned}$$

$$q(u)=q_{\mathrm{cl}}(u)+q'(u)$$

$$\begin{aligned}-\ddot{q}_{\mathrm{cl}}(u)+\omega^2q_{\mathrm{cl}}(u)&=0\\q_{\mathrm{cl}}(0)&=q_i,q_{\mathrm{cl}}(U)=q_f\end{aligned}$$

$$q'(0)=q'(U)=0$$

$$S_{\mathrm{E}}=S_{\mathrm{cl}}\big(q_i,q_f\big)+S'+S_{\mathrm{ct}}$$

$$S_{\mathrm{cl}}\big(q_i,q_f\big)=\frac{1}{2}\int_0^Udu\big[(\partial_uq_{\mathrm{cl}})^2+\omega^2q_{\mathrm{cl}}^2\big]=\omega\frac{\big(q_i^2+q_f^2\big)\cosh~\omega U-2q_iq_f}{2\sinh~\omega U}$$

$$S'=\frac{1}{2}\int_0^Udu[(\partial_uq')^2+\omega^2q'^2]$$



$$\langle q_f, U \mid q_i, 0 \rangle = \exp [-S_{\text{cl}}(q_i, q_f) - S_{\text{ct}}] \int [dq']_{0,0}^{0,U} \exp (-S')$$

$$f_j(u) = \left(\frac{2}{U}\right)^{1/2} \sin \frac{j\pi u}{U}, \lambda_j = \frac{j^2\pi^2}{U^2} + \omega^2$$

$$\det \frac{\Delta}{2\pi} = \prod_{j=1}^{\infty} \frac{j^2\pi^2 + \omega^2 U^2}{2\pi U^2}$$

$$\det \frac{\Delta}{2\pi} \stackrel{\text{regulate}}{\rightarrow} \prod_{j=1}^{\infty} \frac{j^2\pi^2 + \omega^2 U^2}{j^2\pi^2 + \Omega^2 U^2}$$

$$\frac{\Omega \sinh \omega U}{\omega \sinh \Omega U}$$

$$\langle q_f, U \mid q_i, 0 \rangle \rightarrow \left(\frac{\omega}{2\sinh \omega U}\right)^{1/2} \exp \left[-S_{\text{cl}}(q_i, q_f) + \frac{1}{2}(\Omega U - \ln \Omega) - S_{\text{ct}} \right]$$

$$\langle q_f, U \mid q_i, 0 \rangle \rightarrow \langle q_f \mid 0 \rangle \langle 0 \mid q_i \rangle \exp (-E_0 U) = \left(\frac{\omega}{\pi}\right)^{1/2} \exp \left[-\frac{\omega}{2}(q_f^2 + q_i^2 + U) \right]$$

$$\langle q_f, U \mid q_i, 0 \rangle \rightarrow \omega^{1/2} \exp \left[-\frac{\omega}{2}(q_f^2 + q_i^2 + U) + \frac{1}{2}(\Omega U - \ln \Omega) - S_{\text{ct}} \right]$$

$$S_{\text{ct}} = \frac{1}{2} \left[\int_0^U du \Omega \right] - \frac{1}{2} \ln \frac{\Omega}{\pi}$$

$$\langle q_f, U \mid q_i, 0 \rangle = \left(\frac{\omega}{2\pi \sinh \omega U}\right)^{1/2} \exp (-S_{\text{cl}}),$$

Campos Fermiónicos.

$$\begin{aligned} \hat{\psi} | \downarrow \rangle &= 0, \hat{\chi} | \downarrow \rangle = | \uparrow \rangle, \\ \hat{\psi} | \uparrow \rangle &= | \downarrow \rangle, \hat{\chi} | \uparrow \rangle = 0, \\ \{\hat{\psi}, \hat{\chi}\} &= 1, \hat{\psi}^2 = \hat{\chi}^2 = 0. \end{aligned}$$

$$\{\chi_m, \psi_n\} = \{\chi_m, \chi_n\} = \{\psi_m, \psi_n\} = 0.$$

$$\int d\theta = 0, \int d\theta \theta = - \int \theta d\theta = 1$$

$$\int d\theta \frac{d}{d\theta} (a + \theta b) = \int d\theta b = 0$$

$$\int d\theta_1 \dots d\theta_n f(\theta) = c$$

$$f(\theta) = \dots + \theta_n \theta_{n-1} \dots \theta_1 c.$$

$$\begin{aligned} |\psi\rangle &= |\downarrow\rangle + |\uparrow\rangle \psi, \\ \hat{\psi}|\psi\rangle &= |\psi\rangle \psi. \end{aligned}$$



$$\langle \psi\mid \psi'\rangle = \psi - \psi'.$$

$$(\psi-\psi')f(\psi)=(\psi-\psi')f(\psi')$$

$$\int \; d\psi (\psi - \psi') f(\psi) = f(\psi')$$

$$\langle \psi | \hat{\psi} = - \psi \langle \psi |,$$

$$\int \; |\psi\rangle d\psi \langle \psi| = 1$$

$$-\int \; d\chi {\rm exp}\left[\chi(\psi'-\psi)\right] = \psi - \psi' = \langle \psi \mid \psi' \rangle$$

$$\langle \psi \mid \downarrow \rangle = \psi, \langle \psi \mid \uparrow \rangle = -1$$

$$\langle \psi | \hat{\chi} | \psi' \rangle = \langle \psi \mid \uparrow \rangle = - \int \; d\chi \chi {\rm exp}\left[\chi(\psi'-\psi)\right]$$

$$\begin{aligned} -\langle \psi | f(\hat{\chi}, \hat{\psi}) | \psi' \rangle &= \int \; d\chi f(\chi, \psi') {\rm exp}\left[\chi(\psi'-\psi)\right] \\ \langle \psi_f, T \mid \psi_i, 0 \rangle &\approx - \int \; d\chi_0 d\psi_1 d\chi_1 \dots d\psi_{N-1} d\chi_{N-1} \\ &\times {\rm exp} \sum_{m=0}^{N-1} \left[-iH(\chi_m, \psi_m) \epsilon - \chi_m (\psi_{m+1} - \psi_m) \right] \\ &\rightarrow \int \; [d\chi d\psi] {\rm exp} \left\{ \int_0^T dt [-\chi \dot{\psi} - iH(\chi, \psi)] \right\} \end{aligned}$$

$$\begin{array}{ll} \text{(A)}: & \langle \uparrow \mid \uparrow \rangle = \langle \downarrow \mid \downarrow \rangle = 1, \langle \uparrow \mid \downarrow \rangle = \langle \downarrow \mid \uparrow \rangle = 0, \\ \text{(B)}: & \langle \uparrow \mid \downarrow \rangle = \langle \downarrow \mid \uparrow \rangle = 1, \langle \uparrow \mid \uparrow \rangle = \langle \downarrow \mid \downarrow \rangle = 0. \end{array}$$

$$\psi(t)\psi(t')=-\psi(t')\psi(t)$$

$${\rm T}\big[\hat{\psi}(t)\hat{\psi}(t')\big]\equiv \theta(t-t')\hat{\psi}(t)\hat{\psi}(t')-\theta(t'-t)\hat{\psi}(t')\hat{\psi}(t)$$

$$\hat{A}|i\rangle=|j\rangle A_{ji}$$

$$\int \; d\psi \langle \psi | \hat{A} | \psi \rangle = A_{\downarrow \downarrow} - A_{\uparrow \uparrow} = {\rm Tr}\big[(-1)^{\hat{F}} \hat{A}\big]$$

$$\int \; d\psi \langle \psi, U \mid \psi, 0 \rangle = {\rm Tr}\big[(-1)^{\hat{F}} {\rm exp}\left(-\hat{H}U\right)\big]$$

$$\langle 1 \rangle = \int \; [d\psi d\chi] {\rm exp} \left(\int \; d^d x \chi \Delta \psi \right)$$



$$\begin{aligned}\psi(x) &= \sum_i \psi_i \Psi_i(x), \chi(x) = \sum_i \chi_i \Upsilon_i(x) \\ \Delta \Psi_i(x) &= \lambda_i \Psi_i(x), \Delta^T \Upsilon_i(x) = \lambda_i \Upsilon_i(x) \\ \int d^d x \chi_i(x) \Psi_j(x) &= \delta_{ij}\end{aligned}$$

$$\langle 1 \rangle = \prod_i \left[\int d\psi_i d\chi_i \exp(\lambda_i \chi_i \psi_i) \right] = \prod_i \lambda_i = \det \Delta$$

$$\int [d\psi] \exp \left(\int d^d x \psi \Delta \psi \right) = (\det \Delta)^{1/2}$$

$$\int_{-\infty}^{\infty} dx \int_{-\infty}^{\infty} dy \exp(2\pi i \lambda xy) = \frac{1}{\lambda} = \left[\int d\psi \int d\chi \exp(\lambda \chi \psi) \right]^{-1}$$

$$\delta(d^m x d^n \theta) = d^m x d^n \theta \left(\sum_i \frac{\partial}{\partial x_i} \delta x_i - \sum_j \frac{\partial}{\partial \theta_j} \delta \theta_j \right)$$

$$\int [dq dq^*]_\theta \exp(-S_E)$$

$$S = \frac{1}{4\pi\alpha'} \int d^2\sigma (\partial_1 X \partial_1 X + \partial_2 X \partial_2 X + m^2 X^2)$$

$$0 \leq \sigma_1 \leq 2\pi, 0 \leq \sigma_2 \leq T$$

$$H = -\frac{1}{2} \frac{\partial^2}{\partial \phi^2}$$

$$\frac{1}{2\kappa^2} \int d^d x (-G)^{1/2} R$$

$$\chi = \sum_{p=0}^d (-1)^p B_p \frac{i}{2\pi\alpha'} \int_M B + i \int_{\partial M} \mathcal{A}$$

CONCLUSIONES.

En mérito a lo antes expuesto, concluyese que, las ecuaciones de campo de Einstein engranan, en tratándose de un espacio – tiempo cuántico, en tanto se esté ante un efecto cuántico – gravitacional ordinario, es decir, cuando el tejido del espacio – tiempo, en dimensión \mathbb{R}^4 , se deforma, sin que por esto, estemos ante pluridimensiones, como sí supone la supergravedad cuántica, lo que no ocupa este estudio. Sin embargo, esta deformación, supone la existencia de otros puntos de realidad en una misma



dimensión, en los que impera la simetría de gauge fija o corregida. Por tanto, la curvatura, es puramente local, por lo que, las transformaciones de gauge son locales, sin que existan contrapuntos pluridimensionales, sino pluriespaciales, es decir, interacciones de partículas focalizadas en puntos arbitrarios del espacio de Hilber – Einstein deformado. Adicionalmente y como ha quedado anotado, concebida así, la gravedad cuántica, ésta puede ser endógena o exógena, según la naturaleza fenomenológica de la partícula supermasiva en un campo cuántico específico y según las magnitudes escalares y simetrías tensoriales.

REFERENCIAS BIBLIOGRÁFICAS.

- Joseph Polchinski, STRING THEORY VOLUME I, Cambridge University Press 2001, 2005.
- Ashoke Sen y Barton Zwiebach, String Field Theory: A Review, arXiv:2405.19421v2 [hep-th] 19 Jun 2024.
- Harold Erbin, String Field Theory – A Modern Introduction, arXiv:2301.01686v1 [hep-th] 4 Jan 2023.
- Daniel Baumann y Liam McAllister, Inflation and String Theory, arXiv:1404.2601v1 [hep-th] 9 Apr 2014.
- Albuja Bustamante, M. I. (2024). Demostración del Espectro Hamiltoniano para un Campo de Yang–Mills no Abierto que Poseen una Brecha de Masa Finita con Respecto al Estado de Vacío. *Ciencia Latina Revista Científica Multidisciplinar*, 8(1).
- https://doi.org/10.37811/cl_rcm.v8i1.9738.
- Manuel Ignacio Albuja Bustamante, Teoría de Campos: Reforzamiento Teórico – Matemático al Modelo Estándar de Partículas, bajo la estructura cuacional de Yang – Mills, *Ciencia Latina Revista Científica Multidisciplinar*: Vol. 8 Núm. 2 (2024)
- Manuel Ignacio Albuja Bustamante, La brecha de masa y la curvatura de los campos cuánticos, *Ciencia Latina Revista Científica Multidisciplinar*: Vol. 8 Núm. 4 (2024)
- Manuel Ignacio Albuja Bustamante, Aportes Matemáticos en Econometría Teórica, *Ciencia Latina Revista Científica Multidisciplinar*: Vol. 8 Núm. 4 (2024)
- Manuel Ignacio Albuja Bustamante, Formalización Matemática y en Física de Partículas, en Relación a la Brecha de Masa y la Curvatura Geométrica de los Campos Cuánticos, *Ciencia Latina Revista Científica Multidisciplinar*: Vol. 8 Núm. 5 (2024)



Manuel Ignacio Albuja Bustamante, Campos Cuánticos Relativistas: Aproximaciones Teórico – Matemáticas Relativas a los Espacios Cuánticos Geométricamente Deformados o Perforados por Partículas y Antipartículas Supermasivas y Masivas E Hiperpartículas y Suprapartículas, Ciencia Latina Revista Científica Multidisciplinar: Vol. 9 Núm. 1 (2025)

Manuel Ignacio Albuja Bustamante, Teoría Cuántica de Campos Relativistas. Formalización Teórica, Ciencia Latina Revista Científica Multidisciplinar: Vol. 9 Núm. 3 (2025)

Albuja Bustamante, M. I. (2025). Las partículas supermasivas: naturaleza fenomenológica de la partícula oscura. Ciencia Latina Revista Científica Multidisciplinar, 9(3).
https://doi.org/10.37811/cl_rcm.v9i3.18172

